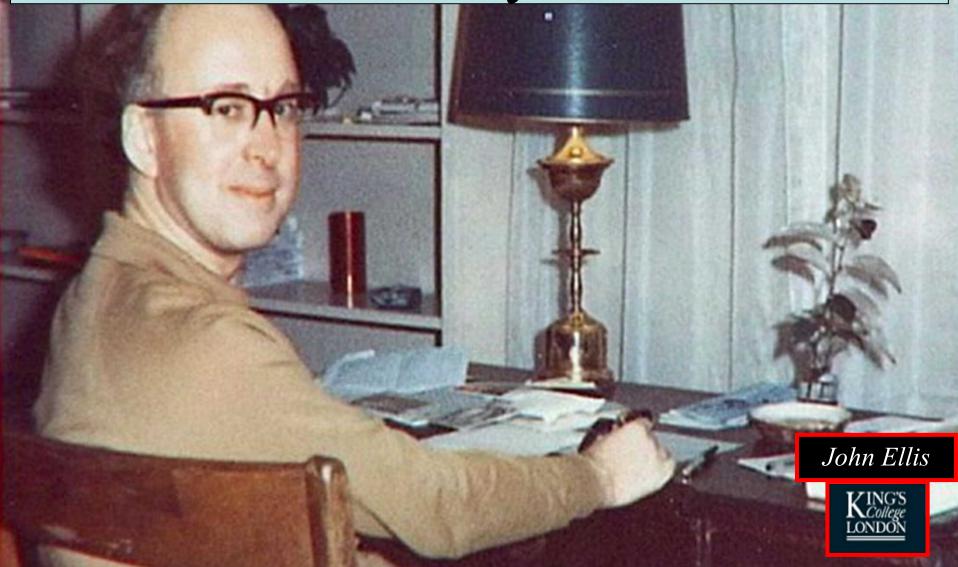
## Particle Physics Today, Tomorrow and Beyond

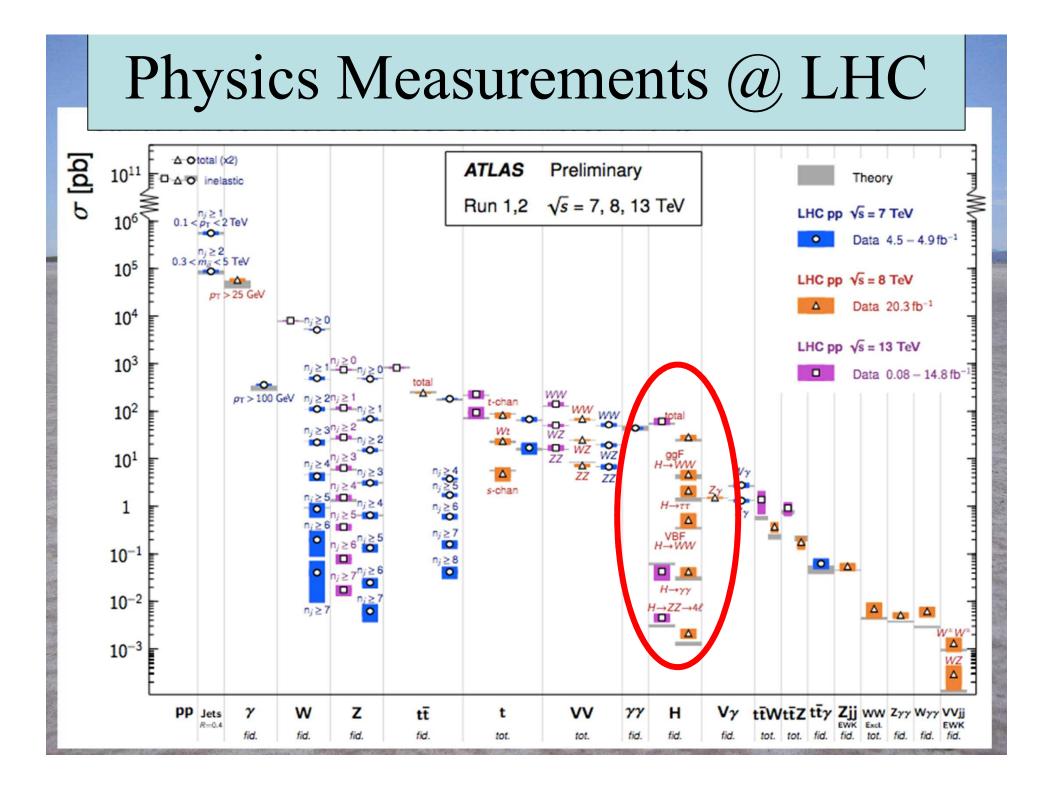


#### Summary of the Standard Model

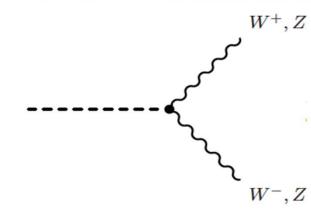
• Particles and SU(3) × SU(2) × U(1) quantum numbers:

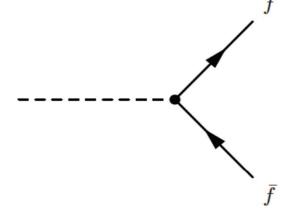
$L_L$ $E_R$	$ \begin{pmatrix} \nu_e \\ e^- \end{pmatrix}_L, \begin{pmatrix} \nu_\mu \\ \mu^- \end{pmatrix}_L, \begin{pmatrix} \nu_\tau \\ \tau^- \end{pmatrix}_L \\ e_R^-, \mu_R^-, \tau_R^- \end{pmatrix} $	( <b>1</b> , <b>2</b> ,-1) ( <b>1</b> , <b>1</b> ,-2)
$Q_L$ $U_R$ $D_R$	$ \begin{pmatrix} u \\ d \end{pmatrix}_{L}, \begin{pmatrix} c \\ s \end{pmatrix}_{L}, \begin{pmatrix} t \\ b \end{pmatrix}_{L} $ $ u_{R}, c_{R}, t_{R} $ $ d_{R}, s_{R}, b_{R} $	$({f 3,2,+1/3}) \ ({f 3,1,+4/3}) \ ({f 3,1,-2/3})$

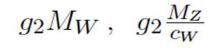
- Lagrangian:  $\mathcal{L} = -\frac{1}{4} F^a_{\mu\nu} F^{a\ \mu\nu}$  gauge interactions +  $i\bar{\psi} / D\psi + h.c.$  matter fermions +  $\psi_i y_{ij} \psi_j \phi + h.c.$  Yukawa interactions +  $|D_\mu \phi|^2 - V(\phi)$  Higgs potential
- Untested before 2012

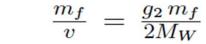


#### Higgs Boson Couplings









$$\Gamma(H \to f\bar{f}) = N_c \frac{G_F M_H}{4\pi\sqrt{2}} m_f^2, \quad N_C = 3 (1) \text{ for quarks (leptons)}$$
  
Weinberg 196  
$$\Gamma(H \to VV) = \frac{G_F M_H^3}{8\pi\sqrt{2}} F(r) \left(\frac{1}{2}\right)_Z, \quad r = \frac{M_V}{M_H}$$
  
Higgs 1966

## A Phenomenological Profile of the Higgs Boson

#### • First attempt at systematic survey

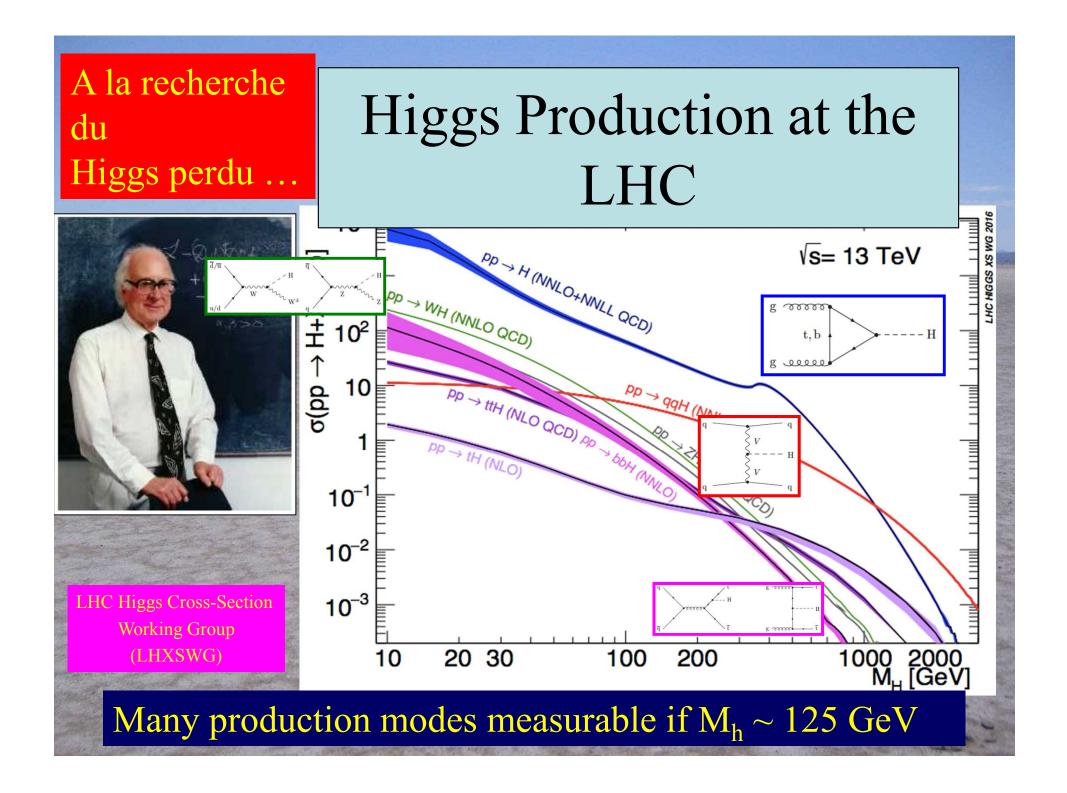
#### A PHENOMENOLOGICAL PROFILE OF THE HIGGS BOSON

John ELLIS, Mary K. GAILLARD \* and D.V. NANOPOULOS \*\* CERN, Geneva

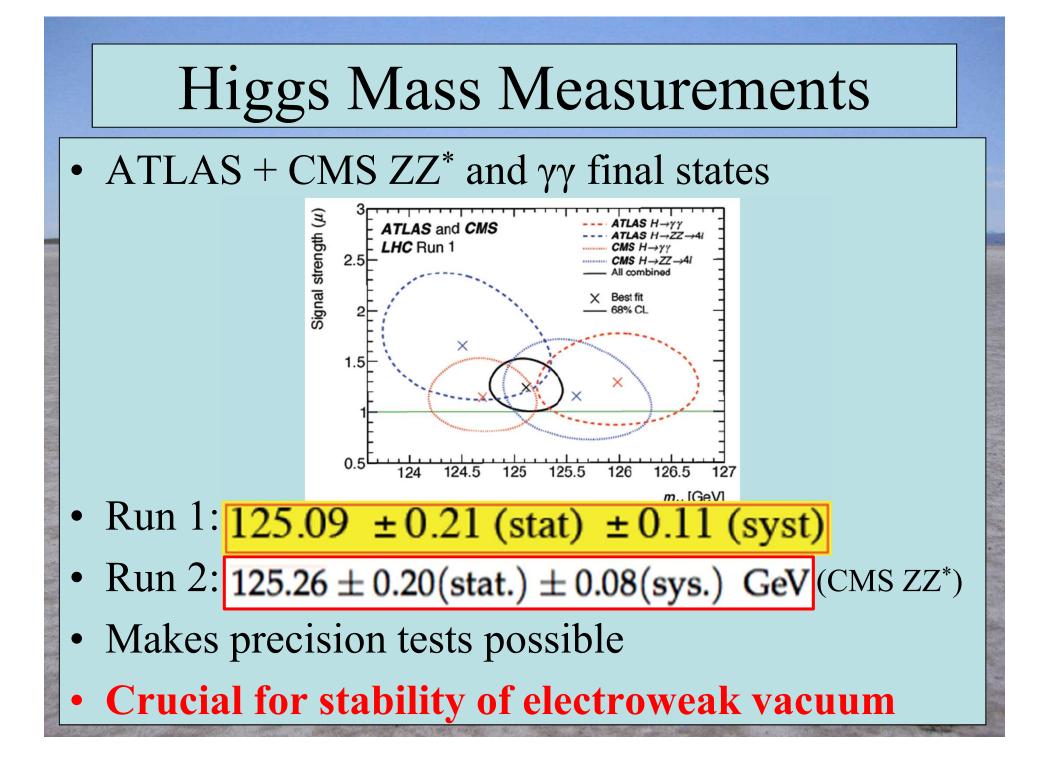
Received 7 November 1975

A discussion is given of the production, decay and observability of the scalar Higgs boson H expected in gauge theories of the weak and electromagnetic interactions such as the Weinberg-Salam model. After reviewing previous experimental limits on the mass of

We should perhaps finish with an apology and a caution. We apologize to experimentalists for having no idea what is the mass of the Higgs boson, unlike the case with charm [3,4] and for not being sure of its couplings to other particles, except that they are probably all very small. For these reasons we do not want to encourage big experimental searches for the Higgs boson, but we do feel that people performing experiments vulnerable to the Higgs boson should know how it may turn up.



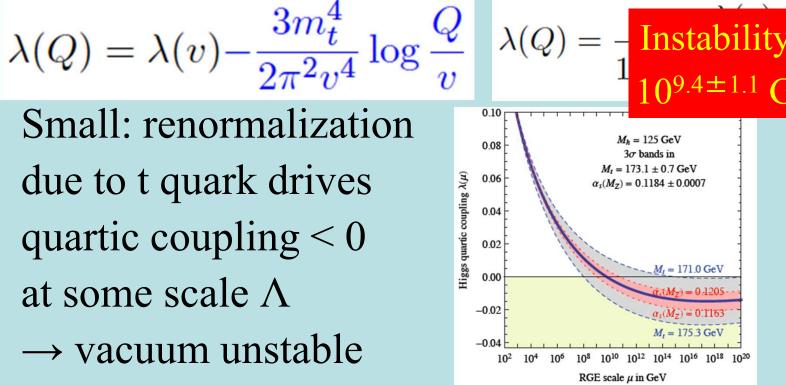
#### Large Hadron Collider at CERN



#### Theoretical Constraints on Higgs Mass

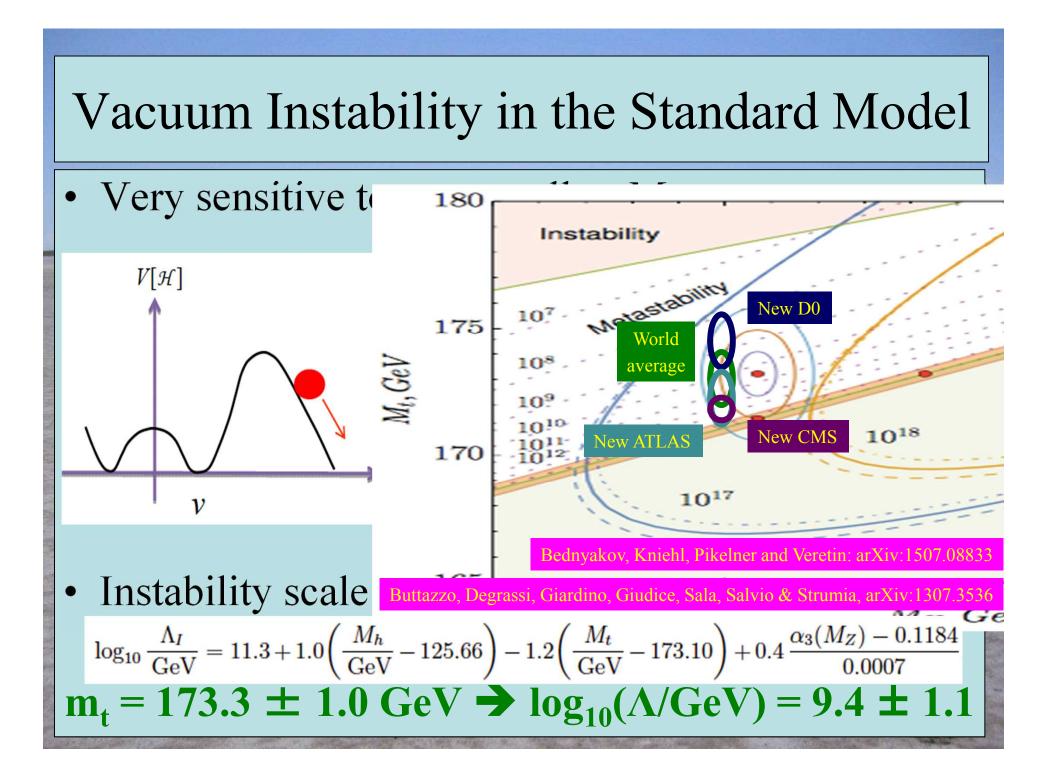
Large  $M_h \rightarrow$  large self-coupling  $\rightarrow$  blow up at

Small: renormalization due to t quark drives quartic coupling < 0at some scale  $\Lambda$  $\rightarrow$  vacuum unstable



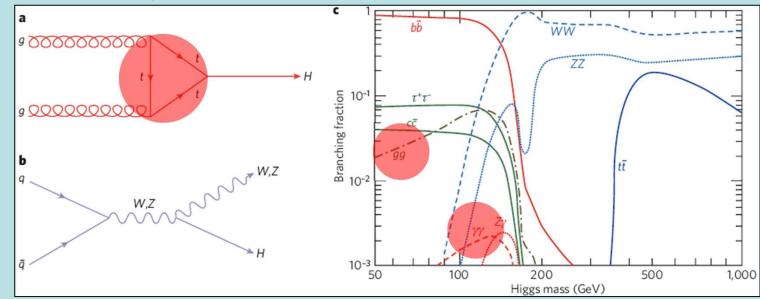
• Vacuum could be stabilized by **Supersymmetry** 

Degrassi, Di Vita, Elias-Miro, Giudice, Isodori & Strumia, arXiv:1205.6497



#### Higgs Decay Branching Ratios

#### • Couplings proportional to masses (?)

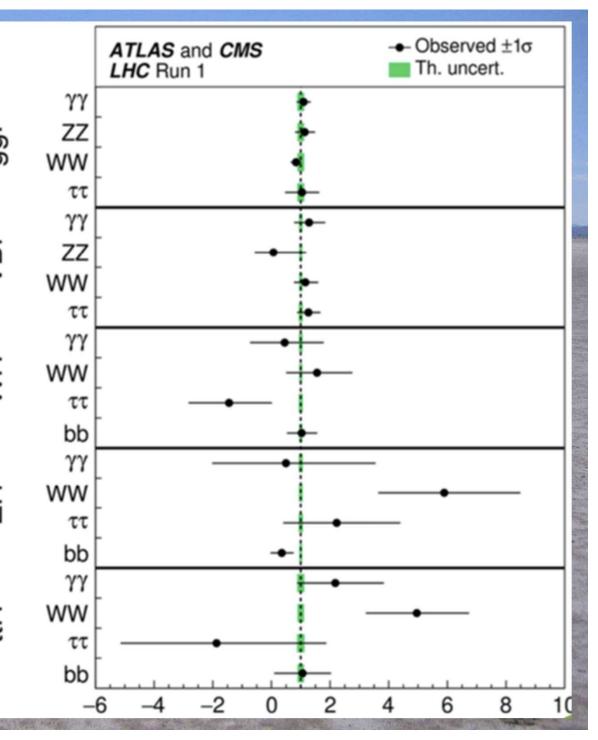


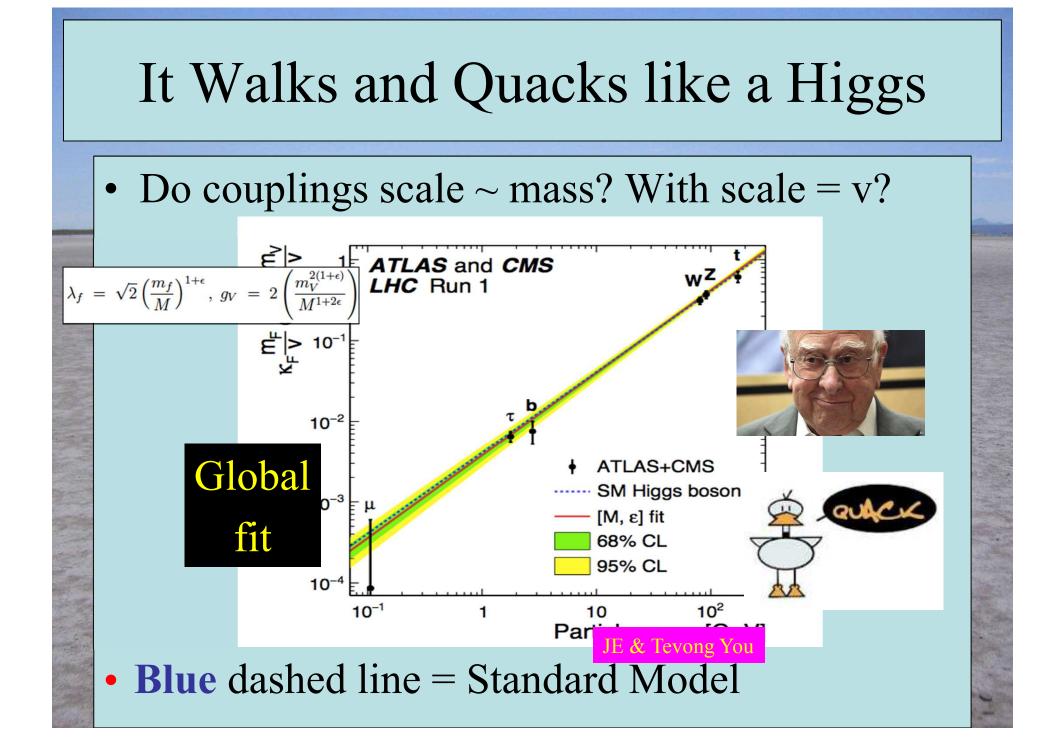
Important couplings through loops:  $-gluon + gluon \rightarrow Higgs \rightarrow \gamma\gamma$ 

Many decay modes measurable if  $M_h \sim 125 \text{ GeV}$ 

#### Production ggF Measurements in Run 1 VBF **Open questions:** - H→bb? MΗ • 2.6σ @ LHC • 2.8σ @ FNAL ZH - H**→ 7**µ?

- ttH production?
- tH production? ∎



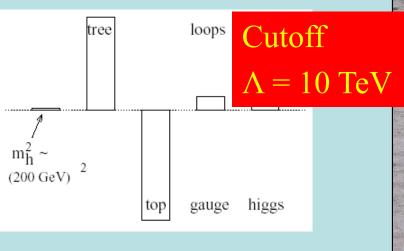




« Empty » space is unsta SUSY Dark matter SUSY **SUSY** Origin of matter Masses of neutrinos Hierarchy problem **SUSY** Inflation **SUSY** Quantum gravity **SUSY** Is Not The Standard Model

### Elementary Higgs or Composite?

- Higgs field:  $<0|H|0> \neq 0$
- Quantum loop problems



Cut-off Λ ~ 1 TeV with Supersymmetry?

- Fermion-antifermion condensate
- Just like QCD, BCS superconductivity

boson?

- New 'technicolour' force?
- Heavy scalar resonance?
   Problems with precision electroweak data
   Pseudo-Nambu-Goldstone

#### Phenomenological Framework

• Assume custodial symmetry:

 $SU(2) \times SU(2) \rightarrow SU(2)_V$   $(\rho \equiv M_W/M_Z \cos \theta_w \sim 1)$ 

• Parameterize gauge bosons by  $2 \times 2$  matrix  $\Sigma$ :

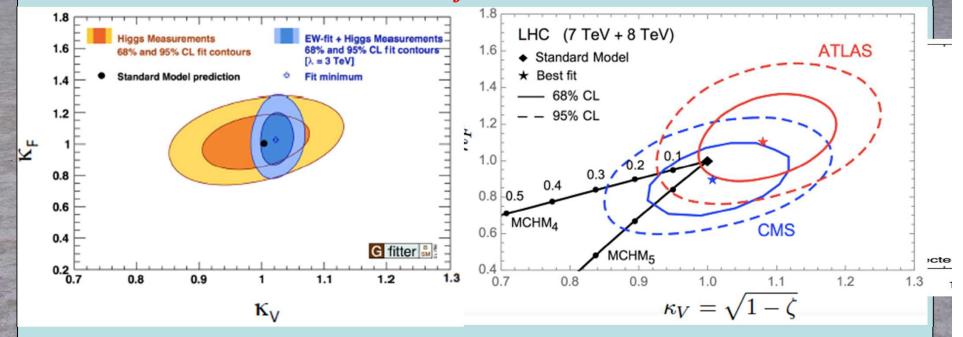
$$\begin{split} \mathcal{L} &= \frac{v^2}{4} \mathrm{Tr} D_{\mu} \Sigma^{\dagger} D^{\mu} \Sigma \left( 1 + 2 \mathbf{a} \frac{h}{v} + \mathbf{b} \frac{h^2}{v^2} + \dots \right) - m_i \bar{\psi}_L^i \Sigma \left( 1 + \mathbf{c} \frac{h}{v} + \dots \right) \psi_R^i + \mathrm{h.c.} \\ &+ \frac{1}{2} (\partial_{\mu} h)^2 + \frac{1}{2} m_h^2 h^2 + \mathbf{d}_3 \frac{1}{6} \left( \frac{3m_h^2}{v} \right) h^3 + \mathbf{d}_4 \frac{1}{24} \left( \frac{3m_h^2}{v^2} \right) h^4 + \dots \quad , \end{split}$$

$$\Sigma = \exp\left(i\frac{\sigma^a\pi^a}{v}\right) \quad \mathcal{L}_{\Delta} = -\left[\frac{\alpha_s}{8\pi}b_sG_{a\mu\nu}G_a^{\mu\nu} + \frac{\alpha_{em}}{8\pi}b_{em}F_{\mu\nu}F^{\mu\nu}\right]\left(\frac{h}{V}\right)$$

• Coefficients a = c = 1 in Standard Model

## Global Analysis of Higgs-like Models

- Rescale couplings: to bosons by  $\kappa_V$ , to fermions by  $\kappa_f$
- Standard Model:  $\kappa_V = \kappa_f = 1$

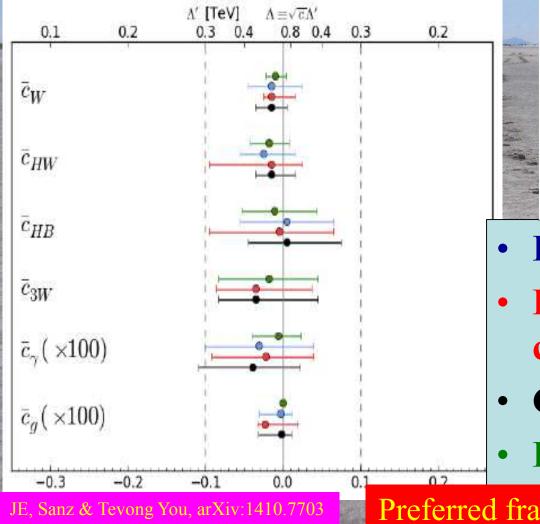


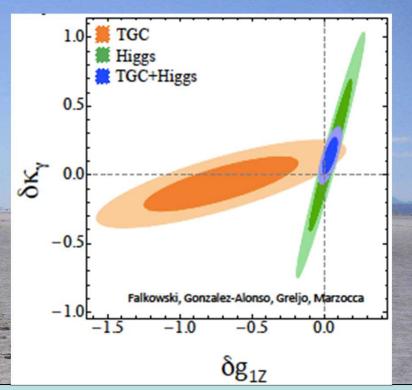
- Consistency between Higgs and EW measurements
- Must tune composite models to look like SM

#### Assuming H(125) is SM-like: Model-independent search for new physics Standard Model Effective Field Theory

- Higher-dimensional operators as relics of higherenergy physics, e.g., dimension 6:  $\mathcal{L}_{\text{eff}} = \sum \frac{f_n}{\Lambda^2} \mathcal{O}_n$
- Operators constrained by SU(2)  $\times$  U(1) symmetry:
  - $\mathcal{L} \supset \frac{\bar{c}_{H}}{2v^{2}} \partial^{\mu} [\Phi^{\dagger} \Phi] \partial_{\mu} [\Phi^{\dagger} \Phi] + \frac{g^{\prime 2} \bar{c}_{\gamma}}{m_{W}^{2}} \Phi^{\dagger} \Phi B_{\mu\nu} B^{\mu\nu} + \frac{g_{s}^{2} \bar{c}_{g}}{m_{W}^{2}} \Phi^{\dagger} \Phi G_{\mu\nu}^{a} G_{a}^{\mu\nu}$  $+ \frac{2ig \bar{c}_{HW}}{m_{W}^{2}} [D^{\mu} \Phi^{\dagger} T_{2k} D^{\nu} \Phi] W_{\mu\nu}^{k} + \frac{ig^{\prime} \bar{c}_{HB}}{m_{W}^{2}} [D^{\mu} \Phi^{\dagger} D^{\nu} \Phi] B_{\mu\nu}$ 
    - $+ \frac{ig \ \bar{c}_W}{m_W^2} \left[ \Phi^{\dagger} T_{2k} \overleftrightarrow{D}^{\mu} \Phi \right] D^{\nu} W_{\mu\nu}^k + \frac{ig' \ \bar{c}_B}{2m_W^2} \left[ \Phi^{\dagger} \overleftrightarrow{D}^{\mu} \Phi \right] \partial^{\nu} B_{\mu\nu}$  $+ \frac{\bar{c}_t}{v^2} y_t \Phi^{\dagger} \Phi \ \Phi^{\dagger} \cdot \bar{Q}_L t_R + \frac{\bar{c}_b}{v^2} y_b \Phi^{\dagger} \Phi \ \Phi \cdot \bar{Q}_L b_R + \frac{\bar{c}_\tau}{v^2} y_\tau \ \Phi^{\dagger} \Phi \ \Phi \cdot \bar{L}_L \tau_R$
- Constrain with precision EW, Higgs data, TGCs ...







- Higgs production
- LHC Triple-gauge couplings
- Global combination
- Individual operators

Preferred framework for analyzing Run 2

#### What lies beyond the Standard Model?

# Supersymmetry

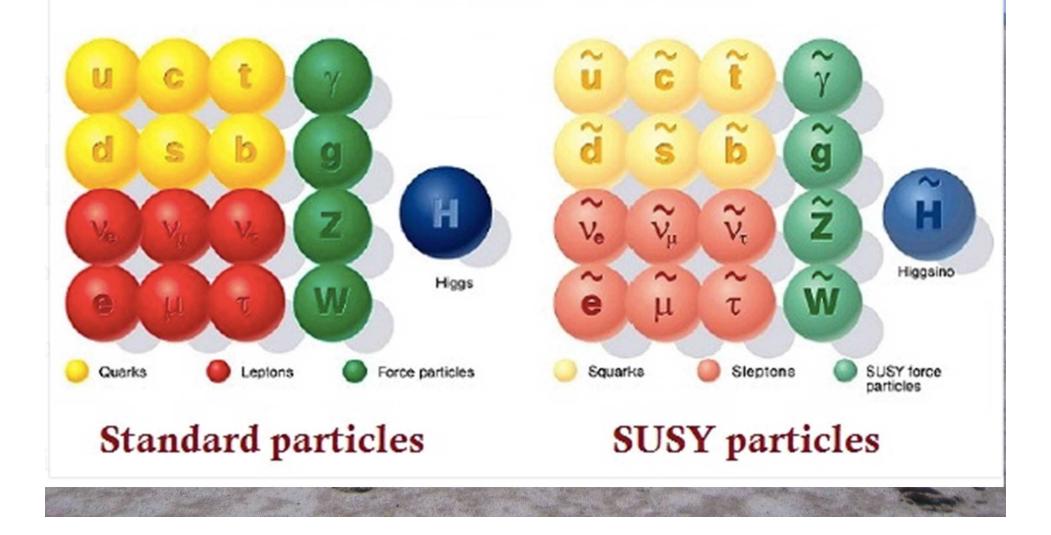
New motivations

From LHC Run 1

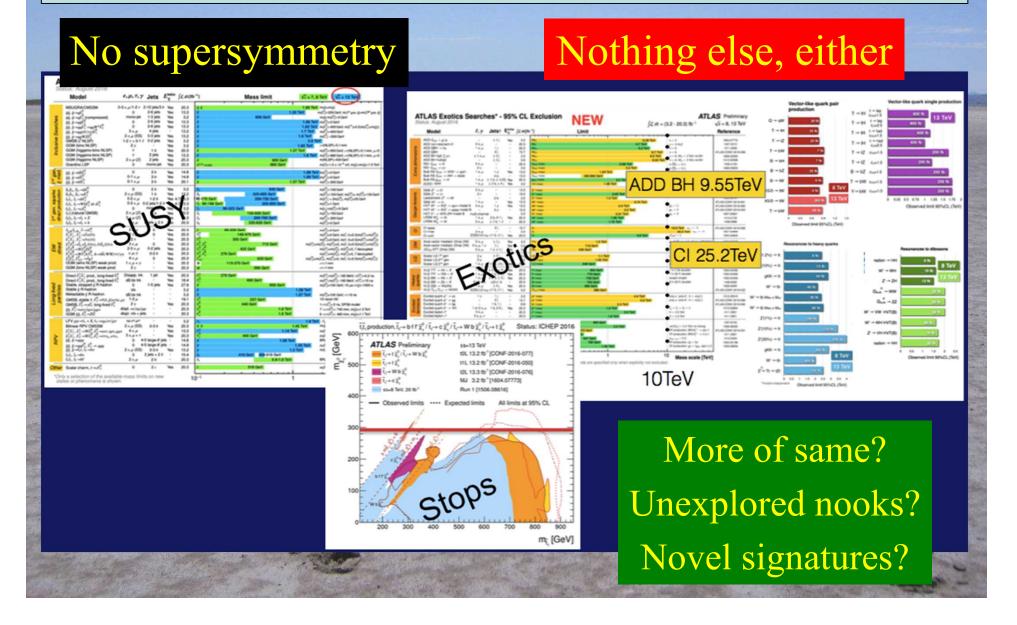
- Stabilize electroweak vacuum
- Successful prediction for Higgs mass
   Should be < 130 GeV in simple models</li>
  - Should be < 130 GeV in simple models</p>
- Successful predictions for couplings

   Should be within few % of SM values
- Naturalness, GUTs, string, ..., dark matter

## Minimal Supersymmetric Extension of the Standard Model



## Nothing (yet) at the LHC



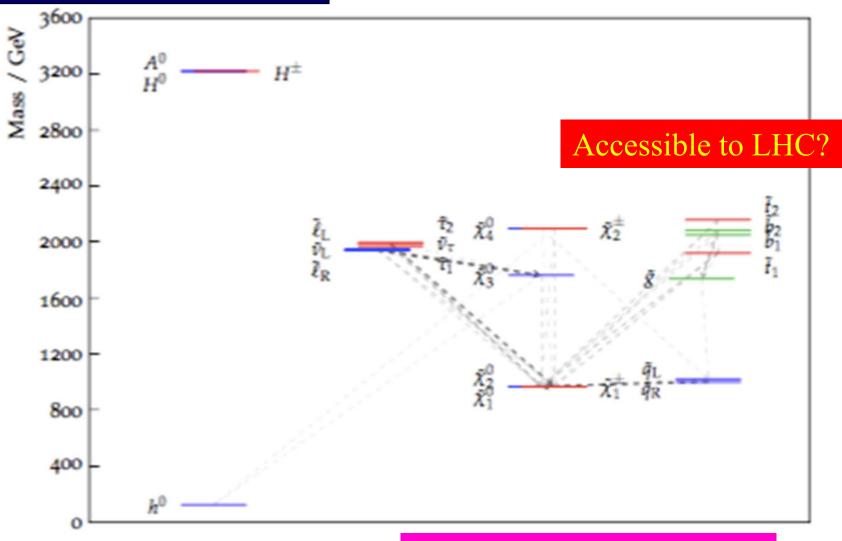
#### Inputs to Global Fits for New Physics



Electroweak	Observable	Source Th./Ex.	Constraint
observables	$M_W$ [GeV] $a_{\mu\nu}^{\rm EXP} - a_{\mu\nu}^{\rm SM}$	[00] / [57, 50] [59] / [60]	$\frac{30.270 \pm 0.012 \pm 0.010_{\text{MSSM}}}{(30.2 \pm 8.8 \pm 2.0_{\text{MSSM}}) \times 10^{-10}}$
Flavour	$R_{\mu\mu}$	[01-05]	2D likelihood, MFV
1 lavoui	$ au(B_s  o \mu^+ \mu^-)$	[63]	$2.04 \pm 0.44$ (stat.) $\pm 0.05$ (syst.) ps
observables:	$BR_{b \rightarrow s\gamma}^{EXP/SM}$	[65]/ [66]	$0.988 \pm 0.045_{\rm EXP} \pm 0.068_{\rm TH,SM} \pm 0.050_{\rm TH,SUSY}$
observables.	$BR_{B}^{EXP/SM}$	[00,07]	$0.992 \pm 0.58_{\rm EXP} \pm 0.096_{\rm SM}$
Internation	$BR_{B \to X_{s} \ell \ell}^{EXP/SM}$	[68]/[66]	$0.966 \pm 0.278_{\rm EXP} \pm 0.037_{\rm SM}$
Interpretation	$\Delta M_{B_{e}}$	[01, 00] / [00]	$0.000 \pm 0.001_{\rm EXP} \pm 0.078_{\rm SM}$
requires	$\frac{\Delta M_{B_g}^2}{\Delta M_{B_g}^{\text{EXP/SM}}}$	[34,69] / [66]	$1.007\pm 0.004_{\rm EXP}\pm 0.116_{\rm SM}$
	$BR_{K \rightarrow \mu\nu}^{EXP/SM}$	[34,70] / [71]	$1.0005\pm0.0017_{\rm EXP}\pm0.0093_{\rm TH}$
lattice inputs	$BR_{K \to \pi \nu \bar{\nu}}^{\text{EXP/SM}}$	[72]/ [73]	$2.01 \pm 1.30_{\rm EXP} \pm 0.18_{\rm SM}$
Doult Matter	$\sigma_p$	[3, 5, 0]	Combined likelihood in the $(m_{\tilde{\chi}_1^0}, \sigma_p^+)$ plane
Dark Matter	$\sigma_{n}^{SD}$	[4]	Likelihood in the $(m_{z^0}, \sigma_p^{SD})$ plane
LHC	$ ilde{g}  ightarrow q  ilde{q}  ilde{\chi}_1^{ m o}, bb  ilde{\chi}_1^{ m o}, tt  ilde{\chi}_1^{ m o}$	[16, 17]	Combined likelihood in the $(m_{\tilde{g}}, m_{\tilde{\chi}_1^0})$ plane
LIIC	${ ilde q}  o q { ilde \chi}_1^0$	[16]	Likelihood in the $(m_{\tilde{q}}, m_{\tilde{\chi}_1^0})$ plane
observables	$ ilde{b}  o b  ilde{\chi}_1^0$	[16]	Likelihood in the $(m_{\tilde{b}}, m_{\tilde{\chi}_1^0})$ , plane
ouser values	$ ilde{t}_1  o t  ilde{\chi}^0_1, c  ilde{\chi}^0_1, b  ilde{\chi}^\pm_1$	[16]	Likelihood in the $(m_{\tilde{t}_1}, m_{\tilde{\chi}_1^0})$ , plane
	$ ilde{\chi}_1^\pm  o  u \ell^\pm  ilde{\chi}_1^0,  u  au^\pm  ilde{\chi}_1^0, W^\pm  ilde{\chi}_1^0$	[18]	Likelihood in the $(m_{\tilde{\chi}^{\pm}_1}, m_{\tilde{\chi}^0_1})$ plane
	$\tilde{\chi}^0_2 \rightarrow \ell^+ \ell^- \tilde{\chi}^0_1, \tau^+ \tau^- \tilde{\chi}^0_1, Z \tilde{\chi}^0_1$	[18]	Likelihood in the $(m_{\tilde{\chi}_{1}^{0}}, m_{\tilde{\chi}_{1}^{0}})$ plane
	Heavy stable charged particles	[74]	Fast simulation based on [74, 75]
	$H/A \rightarrow \tau^+ \tau^-$	[28, 29, 76, 77]	Likelihood in the $(M_A, \tan\beta)$ plane

#### Best-Fit Sparticle Spectrum



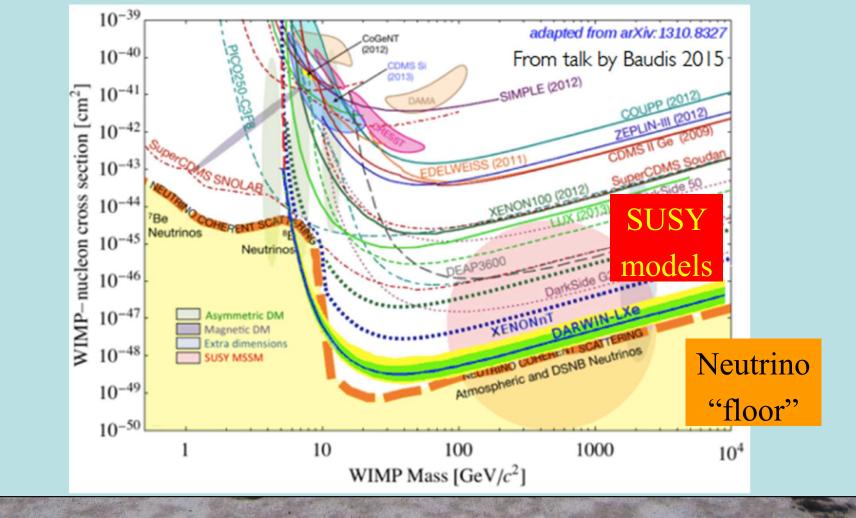


Bagnaschi, Sakurai, JE et al: to appear

mas Tencore

## Direct Dark Matter Searches

• Compilation of present and future sensitivities



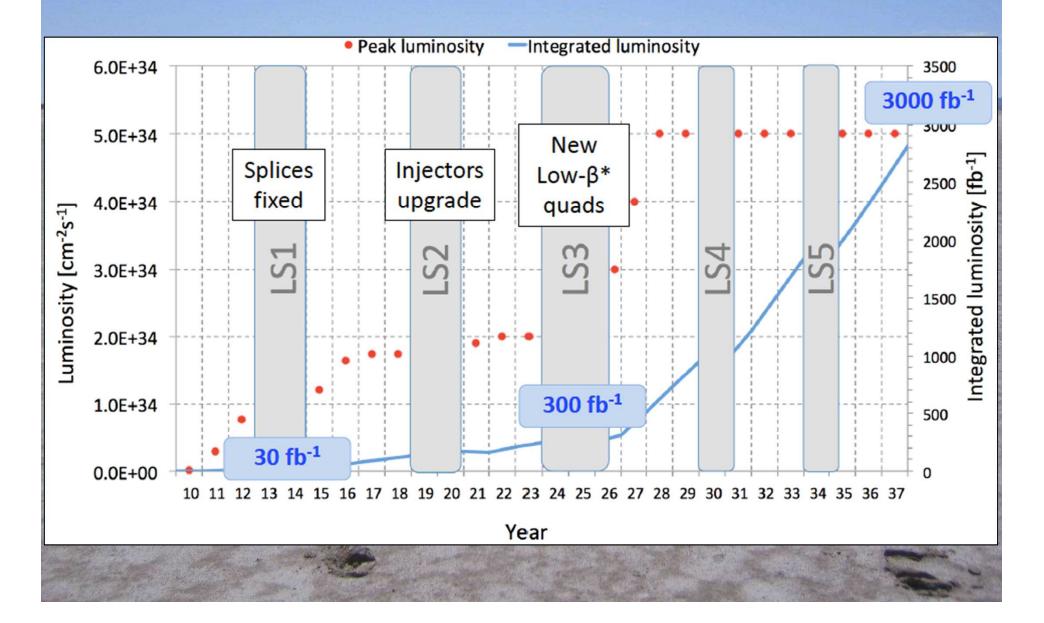
#### Direct Dark Matter Searches



#### Phenomenological MSSM pMSSM11 w/o $(g - 2)_{\mu}$ : best fit, $1\sigma$ , $2\sigma$ , $3\sigma$ pMSSM11 w/o $(g - 2)_{\mu}$ : best fit, $1\sigma$ , $2\sigma$ , $3\sigma$ $10^{-3}$ $10^{-38}$ master.... masteR:coe $10^{-3}$ PICO-60 CRESST-II 10-40 $10^{-4}$ $10^{-43}$ $10^{-42}$ $\sigma_p^{SD}[\mathrm{cm}^2]$ $\sigma_p^{SI}[cm^2]$ $10^{-4}$ $10^{-44}$ 10 PANDAX-II 10-4 10-40 $10^{-51}$ $10^{-48}$ $10^{-53}$ $10^{-5}$ $10^{-50}$ 100 10<sup>1</sup> $10^{2}$ $10^{3}$ 10<sup>1</sup> $10^{2}$ 100 $10^{3}$ $m_{\tilde{\chi}_1^0}$ [GeV] $m_{\tilde{\chi}_1^0}$ [GeV] $\tilde{\chi}_1^{\pm}$ coann. gluino coann. slep coann. stop coann. A/H funnel squark coann. sbot coann. stau coann. **Spin-independent scattering Spin-dependent scattering:** Strongest limit from cross-section close to PandaX upper limit? **PICO** experiment

Bagnaschi, Sakurai, ..., JE et al: in preparation

## The LHC in Future Years

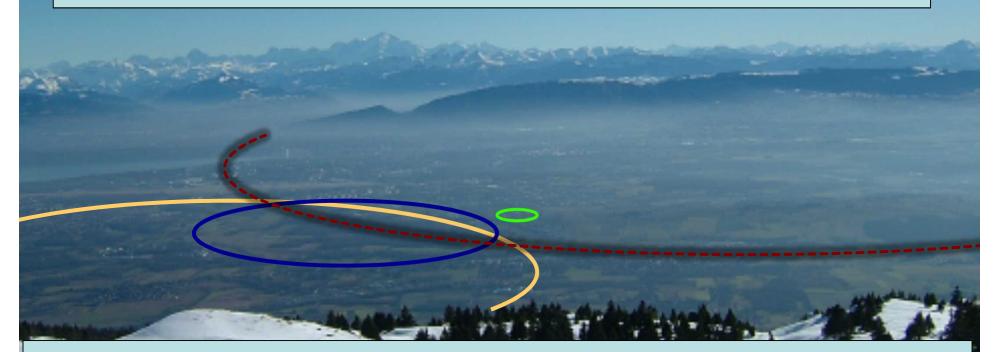




**CEPC-SPPC** 

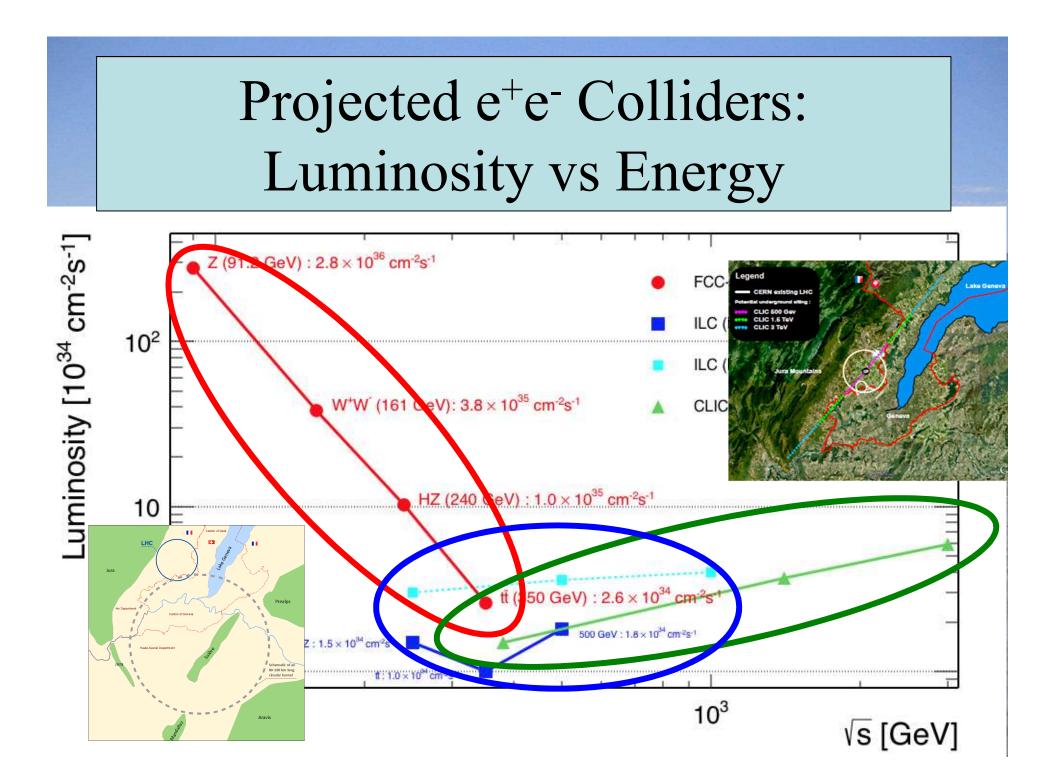
Preliminary Conceptual Design Report

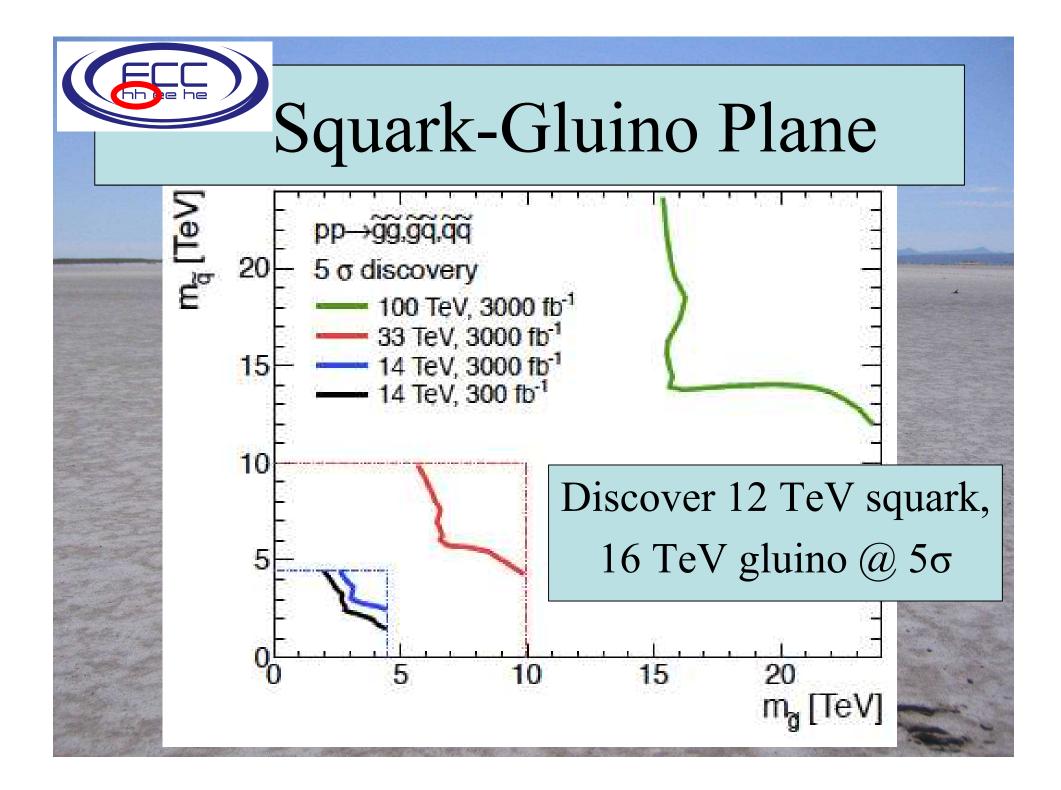
# Future Circular Colliders



#### The vision:

explore 10 TeV scale directly (100 TeV pp) + indirectly  $(e^+e^-)$ 





## Summary

- The discovery of the Higgs boson at the LHC is a big challenge for theoretical physics!
- The LHC may yet discover physics beyond the SM at ~ 13 TeV
- If it **does**, priority will be to study it
- If it does **not**, natural to focus on the Higgs
- In either case, a large circular collider offers the best prospects