

# B/2 correlators in (0,2) hybrid models

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## A quest for (0,2) mirror symmetry

Just what is (0,2) mirror symmetry?

- (2,2) mirror symmetry: Let  $M$  be a CY manifold.

$$(M, T_M) \longleftrightarrow (\widetilde{M}, T_{\widetilde{M}})$$

A model (Kähler moduli)  $\longleftrightarrow$  B model, cx.str. moduli

- (0,2) mirror symmetry: Let  $M$  Calabi-Yau + holomorphic stable bundle  $\mathcal{E}$ .

$$(M, \mathcal{E}) \xleftarrow{?} (\widetilde{M}, \widetilde{\mathcal{E}})$$

A/2 model  $\xleftarrow{?}$  B/2 model

- (Some) open questions:

- Do B/2 model correlators receive quantum corrections?
- If yes, is there a subset of theories where they do not?
- If yes, is there a split in the moduli space, which is exchanged by a mirror map? What about the bundle moduli?

## (0,2) hybrid conformal field theories

Let  $Y$  be a Kähler manifold such that

$$Y = \text{tot} \left( X \xrightarrow{\pi} B \right).$$

A hybrid model can be thought as a fibration of a (0,2) Landau-Ginzburg (LG) model varying adiabatically over a compact and smooth manifold  $B$ . Let  $\mathcal{E} \rightarrow Y$  be a holomorphic vector bundle which couples to the left-moving fermions of the theory.

**Definition 1.** The quadruple  $(Y, \mathcal{E}, V, J)$  defines a nonsingular good hybrid model when it admits

1. Chiral symmetries.
2. Potential condition: the (0,2) superpotential  $J \in \Gamma(\mathcal{E}^*)$  is such that  $J^{-1}(0) = B$ .

This UV theory will RG flow to a family of IR SCFTs parametrized by

1. Kähler class of the Kähler metric of  $Y$ ;
2. Complex structure of  $Y$ ;
3. Complex structure of  $\mathcal{E}$  as a holomorphic bundle;
4. Superpotential  $J$ .

## B/2 heterotic topological ring

(0,2) SCFTs possess a (infinite) chiral ring defined by the cohomology of the supercharge  $\overline{Q}$ ,  $\mathcal{H}_{\overline{Q}}$ . Two projections are possible within  $\overline{Q}$ -cohomology

- $\mathcal{H}_{A/2}$  ring;
- $\mathcal{H}_{B/2}$  ring.

Unitarity bounds imply that the B/2 genus zero correlators

$$c_{\alpha\beta\gamma} = \langle \mathcal{O}_\alpha \mathcal{O}_\beta \mathcal{O}_\gamma \rangle,$$

are topological in the sense that they are independent from the worldsheet metric and from the insertion points of the operators.

**Result 1.** In (0,2) hybrid models, the  $\mathcal{H}_{\overline{Q}}$  chiral ring consists of elements

$$[\mathcal{O}] \in H_{\overline{Q}}^*(Y, \Lambda^* \mathcal{E}),$$

where  $\mathcal{O} \in \Omega^{0,*}(\Lambda^* \mathcal{E})$  is a anti-holomorphic horizontal form (in the sense of the fiber/base geometry of  $Y$ ) valued in the sheaf  $\Lambda^* \mathcal{E}$ .

## Localization and $S^2$ correlators

$S^2$  localization of the B/2 twisted theory yields the following:

**Result 2.** The  $S^2$  correlators are given by the following formula

$$\langle \mathcal{O}_\alpha \mathcal{O}_\beta \mathcal{O}_\gamma \rangle = \int_Y \Omega_Y \wedge \left( \omega_{\mathcal{E}^\perp} \left( \exp \left( -\frac{\mathbf{v}}{4} \|J\|^2 + \frac{\mathbf{v}}{4} \overline{\partial} J^A \partial_A \right) \mathcal{O}_\alpha \mathcal{O}_\beta \mathcal{O}_\gamma \right) \right),$$

where,  $\Omega_Y$  is the holomorphic top form on  $Y$ ,  $\partial_A$  is a basis for  $\mathcal{E}$ ,  $\omega_{\mathcal{E}}$  is a section of  $\wedge^R \mathcal{E}^\vee$ ,  $R = \text{rank } \mathcal{E}$ , and  $\mathbf{v}$  is the volume of the worldsheet  $\Sigma = \mathbb{P}^1$ . The operator  $\omega_{\mathcal{E}^\perp}$  acts by contraction with the section  $\omega_{\mathcal{E}}$ .

The formula above possesses some formal properties, which are expected for a B/2 correlator.

**Result 3.** Let  $\alpha = \mathcal{O}_1 \mathcal{O}_2 \mathcal{O}_3$ , it is possible to show the following

1. Independence from the choice of the representative of cohomology classes

$$\langle \overline{Q}\alpha \rangle = 0.$$

2. Independence from the worldsheet volume  $\mathbf{v}$

$$\langle \alpha \rangle_{\mathbf{v}+\delta\mathbf{v}} = \langle \alpha \rangle_{\mathbf{v}}.$$

3. Independence from Kähler class of the Kähler metric  $g$  of  $Y$

$$\langle \alpha \rangle_{g+\delta g} = \langle \alpha \rangle_g.$$

4. Holomorphicity

$$\langle \alpha \rangle_{\overline{J}+\delta\overline{J}} = \langle \alpha \rangle_{\overline{J}}.$$

(1)

## Transformation law

**Conjecture 1.** There exist  $\mathcal{B} \in \text{Hom}(\mathcal{E}^*, \mathcal{E}^*)$  such that  $T = \mathcal{B}J \in \Gamma(\mathcal{E}^*)$  does not depend on the complex structure parameters and  $T^{-1}(0) = B$ . Then

$$\langle \mathcal{O}_\alpha \mathcal{O}_\beta \mathcal{O}_\gamma \rangle = \sum_n \left( \frac{\mathbf{v}}{4} \right)^n \int_Y \Omega_Y \wedge \left( \mathcal{M}(\mathcal{B})_{n^\perp} \left( e^{-\frac{\mathbf{v}}{4} \|T\|^2} \left( \overline{\partial} T^A \partial_A \right)^n \mathcal{O}_\alpha \mathcal{O}_\beta \mathcal{O}_\gamma \right) \right),$$

where

$$\mathcal{M}(\mathcal{B})_n := \omega_{\mathcal{E}^\perp} \underbrace{\mathcal{B}^t \wedge \cdots \wedge \mathcal{B}^t}_{n\text{-times}}.$$

This form allows to compute the dependence of the correlators from the parameters of  $J$  explicitly.

## Conclusions

- We find perfect agreement with other GLSM phases.
- We obtain a non-trivial prediction for correlators involving twisted fields (from the LG point of view).
- Instanton corrections vanish in all the examples we have checked, suggesting that at least in a large class of theories B/2 correlators are classical. Can we find examples where this does not hold?