PBHs from Inflation and their Implications

Kyohei Mukaida

Kavli IPMU, Univ. of Tokyo \rightarrow DESY, Hamburg (Oct. 2017)

Based on 1611.06130, 1701.02544, 1711.xxxxx In collaboration with K.Inomata, M.Kawasaki, Y.Tada, T.T.Yanagida





Introduction

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Why Primordial Black Holes (PBHs)?

- Non-particle candidate of DM
- Candidate of gravitational wave events observed by LIGO.
- Constrain other DM models; WIMP by UCMH, axion by super-radiance,...

How do you produce them?

- Need Large density perturbations for Gravity > Pressure.
 - Collapse of localized configurations: bubble collision, cosmic string, Q-ball,...
 - Collapse of **primordial** density perturbations: **inflation**, curvaton,...

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Constraints independent of production mechanisms.

▶ Note: a **delta function** for PBH spectrum is assumed.



 Constraints from Neutron Star capture are evaded for a conservative value of DM inside the globular clusters. [See e.g. Kusenko+, 1310.8642; Carr+, 1607.06077]

Hawking radiation EGy: 0912.5297

Gravitational lensing

Femto: 1204.2056 HSC: 1701.02151 Kepler: PhysRevLett.111.181302 EROS/MACHO/OGLE: 0011506, 0607207, 1106.2925 Dynamical

WD: 1505.04444

UFD: 1605.03665, 1704.01668

Accretion

CMB: 1612.05644,1707.04206,...

Radio/Xray: 1612.00457, 1705.00791

Talk by Y. Inoue

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Assume a specific production mechanism (inflation). Are there any other ways to probe them? Can we construct concrete inflation models?

Outline of Talk

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Constraints on PBHs from Inflation

Double inflation: PBHs for LIGO or DM

Summary

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Constraints on PBHs from Inflation

Formation of PBHs

Need large $\delta \rho / \rho$ for Gravity > Pressure Talk by T. Harada







► PBH mass (M) \Leftrightarrow scale of perturbation (k) $M = \gamma \rho \frac{4\pi H^{-3}}{3} \simeq M_{\odot} \left(\frac{\gamma}{0.2}\right) \left(\frac{g_{*}}{3.36}\right)^{-\frac{1}{6}} \left(\frac{k/(2\pi)}{3 \times 10^{-9} \,\mathrm{Hz}}\right)^{-2}$ [Carr, '75]

▶ PBH abundance (β) \leftrightarrows amplitude of perturbation (P_{ζ})

$$\beta(M) = \int_{\delta_c} \mathrm{d}\delta \frac{\mathrm{e}^{-\frac{\delta^2}{2\sigma^2(M)}}}{\sqrt{2\pi\sigma^2(M)}} \sim \sigma(M) \,\mathrm{e}^{-\frac{\delta_c^2}{2\sigma^2(M)}} \qquad \sigma^2(M(k)) = \int \mathrm{d}\ln q \, W^2(q \, k^{-1}) \frac{16}{81} (q \, k^{-1})^4 \, \mathscr{P}_{\zeta}(q) \\ \propto \mathscr{P}_{\zeta}(k)$$

• Enhanced **non-Gaussianity** \rightarrow same amount of PBHs w/ smaller/larger P_{ζ}

Formation of PBHs

Typical probability we need

• 1% of DM @ O(10) solarmass: $\beta \sim 10^{-10}$.

 $\beta \ll I \rightarrow PBHs$

➡ independent of how they produced.

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$$\beta$$
 \rightarrow **plenty of** over-densities

➡ independent of how they produced.
➡ Use them!

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Energy injection from large small-scale perturbs.

How do they affect? → Depends on components and era.



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Constraints on the **Power spectrum** (P_{ζ})



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Energy injection from large small-scale perturbs.

How do they affect? → Depends on components and era.



Energy injection from large small-scale perturbs.

How are they dissipated among background? → Depends on Era.











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Large density perturbation as a source of GWs

Tensor perturbation obeys...

[Saito, Yokoyama,'09; Bugaev, Klimai,'10]

$$h_{ij}'' + 2\mathcal{H}h_{ij}' - \nabla^2 h_{ij} = -4\hat{\mathcal{T}}_{ij;kl}S_{kl}$$

Depends on the **density perturb.**, $\Psi \sim \zeta$

$$S_{ij} \equiv 4\Psi \partial_i \partial_j \Psi + 2\partial_i \Psi \partial_j \Psi - \frac{4}{3(1+w)} \partial_i \left(\frac{\Psi'}{\mathscr{H}} + \Psi\right) \partial_j \left(\frac{\Psi'}{\mathscr{H}} + \Psi\right)$$

Formation of PBH

Production of GW by second order effects $h_{i\,j}\,{\propto}\,\Psi^2\,{\sim}\,\zeta^2$

$$\Omega_{\rm GW}(k)h^2 \sim 10^{-9} \left(\frac{\mathscr{P}_{\zeta}(k)}{10^{-2}}\right)^2$$
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Large density perturbation as a source of GWs

Current and future observations of GWs



Large density perturbation as a source of GWs

• GW has a corresponding peak at the same k.



Constraints on the **Power spectrum** (P_{ζ})



Constraints on **PBHs from inflation** (*misleading)

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No enhanced **non-Gaussianity**. Nakama+1612.06264, 1710.06945

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Flatten your potential

Single-field inflation for the total e-folds of N = 50-60





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 Oscillation
 2nd
 Slow-roll is strongly violated!

 inflation
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Use other fields

Axion-like Inflation, Curvaton,...

Enhanced **non-Gaussianity** can be obtained.

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PBHs for LIGO or DM from **Double Inflation**



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- SUGRA: discrete R symmetry breaking model.

[Kawasaki+1606.07631, Inomata+1611.06130]

$$W = mX\Psi \qquad K = \frac{1}{2} \left(\Psi + \Psi^{\dagger} \right)^{2} + |X|^{2} + |\Phi|^{2} + \frac{\kappa}{4} |\Phi|^{4} \qquad \frac{\|\Psi\|X\|\Phi\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{2}\|v^{$$

Ψ : Ist-inflaton; X: stabilizer; Φ : R-breaking field, 2nd-inflaton





PBHs for LIGO or DM from **Double Inflation**

Total e-folds (N=50-60) = Ist-inflation + 2nd inflation



- PBHs for LIGO \rightarrow **SKA** and future CMB observation.
- PBHs for $DM \rightarrow eLISA$ and LISA.

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Summary

Inflation for PBHs needs LARGE $P_{\zeta}(k) \sim 10^{-2}$.

Many over-densities are generated per one PBH:

- CMB spectral distortion @ 104-1 Mpc-1; BBN @ 105-104 Mpc-1
- Induced GWs: PTA @ ~10⁶ Mpc⁻¹; eLISA @ 10¹¹-10¹³ Mpc⁻¹
- UCMHs...depends on models and DM profile.

PBHs for LIGO \rightarrow need a sharp peak@k~10⁻² Mpc⁻¹

Inomata, Kawasaki, **KM**, Tada, Yanagida; 1611.06130

PBHs for DM \rightarrow could be broad (wave eff. on HSC).

Double inflation (SUGRA) can explain both at once!

Inomata, Kawasaki, **KM**, Yanagida; in prep.

Enhance the **non-Gaussianity** → **multi-field**?

Nakama+1612.06264, 1710.06945