## PRIMORDIAL BLACK HOLES AS A PROBE OF DARK MATTER, DARK ENERGY & DARK DIMENSIONS



## Accelerating universe in the dark, 7/3/19

## PLAN OF TALK

- Introduction
- PBHs as probe of dark energy PBHs formation from inflation Accretion of quintessence X
- PBHs as probe of dark matter

PBHs as source of dark matterPBHs as source of LIGO eventsPBHs as generators of cosmic structure

PBHs as probe of dark dimensions
 PBHs in higher and lower dimensions X

## OVERWHELMING EVIDENCE FOR STELLAR BHS (M~10<sup>1-2</sup>M<sub>0</sub>)



X-ray binary Cygnus X1



LIGO detects gravity waves from coalescing BHs

## OVERWHELMING EVIDENCE FOR SMBH IN AGN (M~10<sup>6-9</sup>M<sub>0</sub>)



MW  $4x10^{6}M_{O}$ QSO  $10^{8}M_{O}$  $1.4x10^{10}M_{O}$ BH at z=6.3



BH mass proportional to stellar mass

## POSSIBLE EVIDENCE FOR IMBH (M~10<sup>3-5</sup>M<sub>0</sub>)



Ultraluminous X-ray sources NGC1313 may have 500M<sub>o</sub> BH



Globular clusters Omega Cen may have 4x10<sup>4</sup>M<sub>0</sub>BH

### **Properties of the binary black hole merger GW150914**

The LIGO Scientific Collaboration and The Virgo Collaboration



$$36^{+5}_{-4} M_{\odot} + 29^{+4}_{-4} M_{\odot} \rightarrow 62^{+4}_{-4} M_{\odot}$$

Largest is now 80 M<sub>O</sub> which is nearly in IMBH range

## **PRIMORDIAL BLACK HOLE FORMATION**

 $R_S = 2GM/c^2 = 3(M/M_O) \text{ km} \implies \rho_S = 10^{18}(M/M_O)^{-2} \text{ g/cm}^3$ 

**Small BHs can only form in early Universe** 

cf. cosmological density  $\rho \sim 1/(Gt^2) \sim 10^6 (t/s)^{-2} g/cm^3$ 

⇒ primordial BHs with horizon mass at formation

 $M_{PBH} \sim c^{3}t/G = \begin{cases} 10^{-5}g \text{ at } 10^{-43}s & (\text{minimum}) \\ 10^{15}g \text{ at } 10^{-23}s & (\text{evaporating now}) \\ 1M_{O} \text{ at } 10^{-5}s & (\text{QCD transition}) \\ 10^{5}M_{O} \text{ at } 1s & (\text{maximum}) \end{cases}$ 



Mon. Not. R. astr. Soc. (1971) 152, 75-78.

#### GRAVITATIONALLY COLLAPSED OBJECTS OF VERY LOW MASS

Stephen Hawking

(Communicated by M. J. Rees)

(Received 1970 November 9)

#### SUMMARY

It is suggested that there may be a large number of gravitationally collapsed objects of mass  $10^{-5}$  g upwards which were formed as a result of fluctuations in the early Universe. They could carry an electric charge of up to  $\pm 30$  electron units. Such objects would produce distinctive tracks in bubble chambers and could form atoms with orbiting electrons or protons. A mass of  $10^{17}$  g of such objects could have accumulated at the centre of a star like the Sun. If such a star later became a neutron star there would be a steady accretion of matter by a central collapsed object which could eventually swallow up the whole star in about ten million years.

# Cosmological effects of primordial black holes

#### GEORGE F. CHAPLINE

Nature **253**, 251–252 (24 January 1975) doi:10.1038/253251a0 Download Citation Received: 29 July 1974 Revised: 03 October 1974 Published online: 24 January 1975

#### Abstract

ALTHOUGH only black holes with masses  $\gtrsim$ ;  $1.5M_{\odot}$  are expected to result from stellar evolution<sup>1</sup> black holes with much smaller masses may be present throughout the Universe<sup>2</sup>. These small black holes are the result of density fluctuations in the very early Universe. Density fluctuations on very large mass scales were certainly present in the early universe as is evident from the irregular distribution of galaxies in the sky<sup>3</sup>. Evidence of density fluctuations on scales smaller than the size of galaxies is generally thought to have been destroyed during the era of radiation recombination<sup>4</sup>. But fluctuations in the metric of order unity may be fossilised in the form of black holes. Observation of black holes, particularly those with masses  $M < M_{\odot}$ , could thus provide information concerning conditions in the very early Universe.

#### First paper on PBHs as dark matter

#### letters to nature

Nature 248, 30 - 31 (01 March 1974); doi:10.1038/248030a0

#### **Black hole explosions?**

S. W. HAWKING

Department of Applied Mathematics and Theoretical Physics and Institute of Astronomy University of Cambridge

QUANTUM gravitational effects are usually ignored in calculations of the formation and evolution of black holes. The justification for this is that the radius of curvature of space-time outside the event horizon is very large compared to the Planck length  $(G\hbar/c^{-3})^{1/2} \approx 10^{-33}$  cm, the length scale on which quantum fluctuations of the metric are expected to be of order unity. This means that the energy density of particles created by the gravitational field is small compared to the space-time curvature. Even though quantum effects may be small locally, they may still, however, add up to produce a significant effect over the lifetime of the Universe  $\approx 10^{17}$  s which is very long compared to the Planck time  $\approx 10^{-43}$  s. The purpose of this letter is to show that this indeed may be the case: it seems that any black hole will create and emit particles such as neutrinos or photons at just the rate that one would expect if the black hole was a body with a temperature of  $(\varkappa/2\pi)$  ( $\hbar/2k$ )  $\approx 10^{-6}$  ( $M\odot/M$ )K where  $\varkappa$  is the surface gravity of the black hole<sup>1</sup>. As a black hole emits this thermal radiation one would expect it to lose mass. This in turn would increase the surface gravity and so increase the rate of emission. The black hole would therefore have a finite life of the order of  $10^{71}$  ( $M\odot/M$ )<sup>-3</sup> s. For a black hole of solar mass this is much longer than the age of the Universe. There might, however, be much smaller black holes which were formed by fluctuations in the early Universe<sup>2</sup>. Any such black hole of mass less than  $10^{15}$  g would have evaporated by now. Near the end of its life the rate of emission would be very high and about  $10^{30}$  erg would be released in the last 0.1 s. This is a fairly small explosion by astronomical standards but it is equivalent to about 1 million 1 Mton hydrogen bombs.





## PBHs are important even if they never formed!

## **PBH EVAPORATION**

**Black holes radiate thermally with temperature** 

$$\mathbf{T} = \frac{hc^3}{8\pi G k M} \sim \mathbf{10^{-7}} \left[\frac{M}{M_0}\right]^{-1} \mathbf{K}$$
  
=> evaporate completely in time  $\mathbf{t}_{evap} \sim \mathbf{10^{64}} \left[\frac{M}{M_0}\right]^3 \mathbf{y}$ 

 $M \sim 10^{15}g \Rightarrow$  final explosion phase today (10<sup>30</sup> ergs)

This can only be important for PBHs

γ-ray background at 100 MeV =>  $\Omega_{PBH}(10^{15}g) < 10^{-8}$ 

=> explosions undetectable in standard particle physics model T > T<sub>CMB</sub>=3K for M <  $10^{26}$ g => "quantum" black holes

## **BLACK HOLES**



## **FORMATION MECHANISMS**

Primordial inhomogeneities Inflation



Cosmic strings PBH constraints =>  $G \mu < 10^{-6}$ 



Bubble collisions Need fine-tuning of bubble formation rate Domain walls PBHs can be very large





 $\bigcirc$ 

## **PBH FORMATION => LARGE INHOMOGENEITIES**

To collapse against pressure, need (Carr 1975)

 $R > \sqrt{\alpha}$  ct when  $\delta \sim 1 \implies \delta_{\rm H} > \alpha$  (p= $\alpha \rho c^2$ )



 $\epsilon(M)$  decreases with M => exponential upper cut-off

## **Limit on fraction of Universe collapsing**

 $\beta(M)$  fraction of density in PBHs of mass M at formation

**General limit** 

$$\frac{\rho_{PBH}}{\rho_{CBR}} \approx \frac{\Omega_{PBH}}{10^{-4}} \left[ \frac{R}{R_0} \right] \Longrightarrow \beta \sim 10^{-6} \Omega_{PBH} \left[ \frac{t}{\text{sec}} \right]^{1/2} \sim 10^{-18} \Omega_{PBH} \left[ \frac{M}{10^{15} g} \right]^{1/2}$$

So both require and expect  $\beta(M)$  to be tiny => fine-tuning

UnevaporatedM>10^{15}g =>  $\Omega_{PBH} < 0.25$  (CDM)Evaporating nowM~10^{15}g =>  $\Omega_{PBH} < 10^{-8}$  (GRB)Evaporated in pastM<10^{15}g</td>

=> constraints from entropy, γ-background, BBNS

## PBHS AS PROBE OF PRIMORDIAL FLUCTUATIONS



PBHs are unique probe of  $\varepsilon$  on small scales.

Need blue spectrum or spectral feature to produce them.

## **CONSTRAINTS FOR EVAPORATING PBHS**

#### B. Carr, K. Kohri, Y. Sendouda & J. Yokoyama PRD 81(2010) 104019



## PBHS FROM NEAR-CRITICAL COLLAPSE

Critical phenomena => M = k  $M_H(\delta - \delta_c)^{\gamma}$  (Niemeyer & Jedamzik 1999) spectrum peaks at horizon mass with extended low mass tail

 $dN/dM \propto M^{1/\gamma-1} \exp[-(M/M_f)^{1/\gamma}]$  ( $\gamma = 0.35$ ) (Yokoyama 1998)

Later calculations and peak analysis =>

 $\delta_{\rm C} \sim 0.45$  and applies to  $\delta - \delta_{\rm C} \sim 10^{-10}$  (Musco & Miller 2013)



## **PBHS AND INFLATION**

**PBHs formed before reheat inflated away =>** 

 $M > M_{min} = M_{Pl} (T_{reheat} / T_{Pl})^{-2} > 1 \text{ gm}$ 

CMB quadrupole => T<sub>reheat</sub> < 10<sup>16</sup>GeV

**But inflation generates fluctuations** 

$$\frac{\delta\rho}{\rho} \sim \left[\frac{V^{3/2}}{M_{\rm Pl}^{3}V'}\right]_{H}$$



**Can these generate PBHs?** 

[HUGE NUMBER OF PAPERS ON THIS]

## QUANTUM DIFFUSION

★ Consider the possibility of a plateau in the inflaton potential:



### QUANTUM DIFFUSION: CURRENT HOT TOPIC

## Quantum diffusion during inflation and primordial black holes

#### arXiv:1705.04861

Chris Pattison,<sup>a</sup> Vincent Vennin,<sup>b,a</sup> Hooshyar Assadullahi,<sup>a,c</sup> and David Wands<sup>a</sup>

### Quantum diffusion beyond slow-roll: implications for primordial black-hole production

#### arXiv:1805.06731

Jose María Ezquiaga<sup>*a,b*</sup> and Juan García-Bellido<sup>*a,b*</sup>

# Single Field Double Inflation and Primordial Black Holes

arXiv:1705.06225 K. Kannike,<sup>*a*</sup> L. Marzola,<sup>*a,b*</sup> M. Raidal,<sup>*a*</sup> and H. Veermäe<sup>*a*</sup> Primordial black hole production in critical Higgs inflation arXiv:1705.04861 J Ezquiaga, J Garcia-Bellido, E Morales

#### Primordial black holes from an inflexion point arXiv:1706.042261

C Germani and T Prokopec

# Primordial black holes from inflation and quantum diffusion

arXiv:1804.07124

M. Biagetti,<sup>a</sup> G. Franciolini,<sup>b</sup> A. Kehagias<sup>c</sup> and A. Riotto<sup>b</sup>

### Primordial Black Holes With Multi-Spike Mass Spectra

Bernard Carr<sup>1, \*</sup> and Florian Kühnel<sup>2, 3, †</sup>

arXiv:1811.06532



cf. talk by Yuichiro Tada

## PRIMORDIAL BLACK HOLES AS DARK MATTER

## PRO

\* Black holes exist
\* No new physics needed
\* LIGO results

CON

\* Requires fine-tuning

# PBH can do it!



#### PRIMORDIAL BLACK HOLES AS DARK MATTER



 $\Omega_{vis}$ = 0.01,  $\Omega_{dm}$ = 0.25  $\Rightarrow$  need baryonic and non-baryonic DM MACHOs WIMPs

#### PBHs are non-baryonic with features of both WIMPs and MACHOs



Early microlensing searches suggested MACHOs with 0.5 M<sub>O</sub> => PBH formation at QCD transition? Pressure reduction => PBH mass function peak at 0.5 M<sub>O</sub>

Later found that at most 20% of DM can be in these objects



## PRIMORDIAL BLACK HOLES AS DARK MATTER

Bernard Carr,<sup>1,\*</sup> Florian Kühnel,<sup>2,†</sup> and Marit Sandstad<sup>3,‡</sup> PRD 94, 083504, arXiv:1607.06077



Three windows: (A) intermedate mass; (B) sublunar mass; (C) asteroid mass.

#### Also (D) Planck mass relics?

But some of these limits are now thought to be wrong

### WHICH MASS WINDOW IS MOST PLAUSIBLE?

PBH dark matter @10 M<sub>o</sub> from hybrid inflation

Clesse & Garcia-Bellido arXiv:1501.07565 PBH dark matter @10<sup>20</sup>g from double inflation Inomata et al arXiv:1701.02544

1

1035



cf. light versus heavy dark matter particle

# Microlensing constraints on primordial black holes with the Subaru/HSC Andromeda observation

Hiroko Niikura<sup>1,2</sup>, Masahiro Takada<sup>1</sup>, Naoki Yasuda<sup>1</sup>, Robert H. Lupton<sup>3</sup>, Takahiro Sumi<sup>4</sup>, Surhud More<sup>1,5</sup>, Toshiki Kurita<sup>1,2</sup>, Sunao Sugiyama<sup>1,2</sup>, Anupreeta More<sup>1</sup>, Masamune Oguri<sup>1,2,6</sup>, Masashi Chiba<sup>7</sup>

#### arXiv:1701.02151



## Earth-mass black holes? – Constraints on primordial black holes with 5-years OGLE microlensing events

Hiroko Niikura,<sup>1, 2, \*</sup> Masahiro Takada,<sup>2, †</sup> Shuichiro Yokoyama,<sup>3, 2</sup> Takahiro Sumi,<sup>4</sup> and Shogo Masaki<sup>5</sup>

#### arXiv:1901.07120



### Limits on stellar-mass compact objects as dark matter from gravitational lensing of type Ia supernovae

Miguel Zumalacárregui<sup>1, 2, 3</sup>, \* and Uroš Seljak<sup>1, 4</sup>, †

arXiv:1712.02240



But PBHs can evade these limits Garcia-Bellido et al, arXiv:1712.06574

# Formation and evolution of primordial black hole binaries in early universe

Raidal, Spethmann, Vaskonen, Veermae

arXiv:1812.01930 Kavanagh et al arXiv:1805.09034  $10^{0}$  $= \Omega_{PBH}/\Omega_{DM}$ 0.1LIGO **FBH** DM fraction  $f_{PBH}^{PBH} = 10^{-5}$ limit 10-2 LIGO EROS+MACHO Eridanus II ccretion - radio Accretion - X-ray  $10^{-3}$ CMB - PLANCK  $\sigma = 0.6$ CMB - FIRAS  $10^{-4}$  $10^{2}$  $10^{1}$  $10^{3}$  $10^{0}$ 100 1000 0.010.110 1  $M_{\rm PBH} [M_{\odot}]$  $m_c/M_{\odot}$ 

# Multi-wavelength astronomical searches for primordial black holes

Julien Manshanden<sup>a,b</sup> Daniele Gaggero<sup>a,c</sup> Gianfranco Bertone<sup>a</sup> Riley M. T. Connors<sup>d</sup> Massimo Ricotti<sup>e</sup>

arXiv:1812.07967



#### Sunyaev-Zel'dovich anisotropy due to Primordial black holes

Katsuya T. Abe,\* Hiroyuki Tashiro, and Toshiyuki Tanaka

arXiv:1901.06809



### **VOYAGER-1** $e^{\pm}$ further constrain Primordial Black Holes as Dark Matter

Mathieu Boudaud<sup>1</sup> and Marco Cirelli<sup>1</sup>

#### arXiv:1807.03075



- EGB limits (Fermi-LAT) Carr+(2010)
- red band: propagation uncertainty (magnetic halo size)
   4 < L < 20 kpc Reinert & Winkler(2018)</li>
- even better assuming a background for Voyager-1 data (SNRs e<sup>-</sup>)

$$\Phi_{e^-}(E) \propto E^{-1.3}$$

#### CONSTRAINTS ON PRIMORDIAL BLACK HOLES

Bernard Carr,<sup>1, 2, \*</sup> Kazunori Kohri,<sup>3, †</sup> Yuuiti Sendouda,<sup>4, ‡</sup> and Jun'ichi Yokoyama<sup>2, 5, §</sup>



## LENSING LIMITS



### DYNAMICAL LIMITS



### ACCRETION AND LIGO LIMITS



## LARGE-SCALE STRUCTURE LIMITS



## **EVAPORATION LIMITS**



## COSMOLOGICAL LIMITS



#### **EVOLUTION OF CONSTRAINTS ON COLLAPSE FRACTION**



1994





## These constraints are not just nails in a coffin!



All constraints have caveats and may change

Each constraint is a potential signature

PBHs are interesting even if f << 1

## **PBHS AND LIGO**



Do we need Population III or primordial BHs?

## **Prescience of Japanese!**

GRAVITATIONAL WAVES FROM COALESCING BLACK HOLE MACHO BINARIES Takashi Nakamura, Misao Sasaki, Takahiro Tanaka and Kip Thorne

THE ASTROPHYSICAL JOURNAL, 487:L139–L142, 1997



If MACHOs are black holes of mass ~0.5  $M_{\odot}$ , they must have been formed in the early universe when the temperature was ~1 GeV. We estimate that in this case in our Galaxy's halo out to ~ 50 kpc there exist ~5 × 10<sup>8</sup> black hole binaries the coalescence times of which are comparable to the age of the universe, so that the coalescence rate will be ~5 × 10<sup>-2</sup> events yr<sup>-1</sup> per galaxy. This suggests that we can expect a few events per year within 15 Mpc. The gravitational waves from such coalescing black hole MACHOs can be detected by the first generation of interferometers in the LIGO/VIRGO/TAMA/GEO network. Therefore, the existence of black hole MACHOs can be tested within the next 5 yr by gravitational waves.

## POSSIBLE INDIRECT CONFIRMATION OF THE EXISTENCE OF POP III MASSIVE STARS BY GRAVITATIONAL WAVES

Tomaya Kinagawa, Kohei Inayoshi, Kenta Hotokezaka, Daisuka Nakauchi and Tahashi Nakamura



We perform population synthesis simulations for Population III (Pop III) coalescing compact binary which merges within the age of the Universe. We found that the typical mass of Pop III binary black holes (BH–BHs) is  $\sim 30 \,\mathrm{M}_{\odot}$  so that the inspiral chirp signal of gravitational waves can be detected up to z = 0.28 by KAGRA, Adv. LIGO, Adv. Virgo and GEO

#### Did LIGO detect dark matter?

Simeon Bird,<sup>\*</sup> Ilias Cholis, Julian B. Muñoz, Yacine Ali-Haïmoud, Marc Kamionkowski, Ely D. Kovetz, Alvise Raccanelli, and Adam G. Riess<sup>1</sup> arXiv:1603.00464

Dark matter in 20-100 M<sub>O</sub> binaries may provide observed rate of 2-53 Gpc<sup>-1</sup>yr<sup>-1</sup>

#### Primordial Black Hole Scenario for the Gravitational-Wave Event GW150914

Misao Sasaki,<sup>1</sup> Teruaki Suyama,<sup>2</sup> Takahiro Tanaka,<sup>3,1</sup> and Shuichiro Yokoyama<sup>4</sup>

#### arXiv:1603.08338

Only need small f and comparable to limits from CMB distortion

LIGO gravitational wave detection, primordial black holes and the near-IR cosmic infrared background anisotropies

A. Kashlinsky<sup>1</sup>,

arXiv:1605.04023



![](_page_46_Picture_11.jpeg)

#### Spin Distribution of Primordial Black Holes

Takeshi Chiba<sup>1</sup> and Shuichiro Yokoyama<sup>2</sup>

#### arXiv:1704.06573

#### Abstract

We estimate the spin distribution of primordial black holes based on the recent study of the critical phenomena in the gravitational collapse of a rotating radiation fluid. We find that primordial black holes are mostly slowly rotating.

![](_page_47_Figure_5.jpeg)

#### Primordial black holes with an accurate QCD equation of state

Christian T. Byrnes,<sup>1,\*</sup> Mark Hindmarsh,<sup>1,2,†</sup> Sam Young,<sup>1,‡</sup> and Michael R. S. Hawkins<sup>3,§</sup>

arXiv:1801.06138

![](_page_48_Figure_3.jpeg)

![](_page_48_Figure_4.jpeg)

## PBHS, DARK MATTER AND BARYOGENESIS AT THE QUARK HADRON EPOCH: ADDRESSING THE FINE-TUNINGS

BC, Sebastien Clesse & Juan Garcia-Bellido (2019)

Stars have a mass in range (0.1 – 10)  $M_C$  where  $M_C \sim \alpha_G^{-3/2} m_P \sim 1 M_O$  and  $\alpha_G \sim Gm_P^2/hc \sim 10^{-38}$ 

PBHs forming at time t have mass and collapse fraction  $M \sim 10^5$ (t/s) M<sub>O</sub>,  $\beta$ (M) ~ 10<sup>-9</sup> f(M) (M/M<sub>O</sub>)<sup>1/2</sup>

So  $\beta$  must be fine-tuned and we must also explain why

$$\chi = \rho_{PBH} / \rho_B = f \rho_{DM} / \rho_B = 6 f$$
 is O(1).

QCD epoch => M ~ M<sub>C</sub>,  $\beta$ (M) ~  $\eta = n_B/n_\gamma$  ~10<sup>-9</sup>

dark matter and visible baryons have similar mass
>
PBHs may generate baryon asymmetry

 $\chi >> 1 => t_{eq} << t_{dec} =>$  not enough baryons to make galaxies  $\chi << 1 => t_{deq} >> t_{dec} =>$  fluctuations too small to make galaxies

Baryogenesis scenario

EW baryogenesis a QCD epoch

Baryon violation via sphaleron transitions and B+L chiral anomaly

CP violation via CKM matrix

Equilibrium violation via supercooling near QCD scale

PBH formation => large curvature perturbations => huge entropy production => out-of-equilibrium condition => baryogenesis

$$\eta_{\text{loc}} \sim \frac{\Gamma_{\text{sph}}(T_{\text{eff}})}{T_{\text{eff}} T_{\text{rh}}^3} \, \delta_{\text{CP}}(T_{\text{rh}}) \, \Delta \theta \sim O(1) \text{ for } T_{\text{eff}} \sim 60 \text{ GeV}, \, T_{\text{rh}} \sim 200 \text{ MeV}, \, \Delta \theta \sim 1$$

Diffusion of baryon asymmetry =>  $\eta \sim \beta \eta_{\text{loc}}$ 

## Curvature perturbation scenario

Natural peak in PBH mass function but need to fine-tune pert' amp'

Stochastic fluct'ns in spectator field during inflation (QCD axion)  $\Rightarrow$  different values in different patches

- $\Rightarrow$  frozen until PE dominates density long after inflation
- $\Rightarrow$  2<sup>nd</sup> inflation phase within region (few e-folds)
- $\Rightarrow$  non-linear perturbations => PBHs.

Axion field fluct' =>  $\Delta N \sim O(1) => O(1)$  curv' fluct' if f<sub>a</sub> ~ 0.2 M<sub>Pl</sub> => PQ breaking at GUT epoch => axions dominate at QCD epoch.

![](_page_51_Figure_7.jpeg)

1<sup>st</sup> peak at 1M<sub>O</sub> for DM plus 2<sup>nd</sup> peak at 10-20 M<sub>O</sub> for LIGO events

## PBHS AS GENERATORS OF COSMIC STRUCTURES B.J. Carr & J. Silk MNRAS 478 (2018) 3756; arXiv:1801.00672

What is maximum mass of PBH?

Could  $10^6 - 10^{10} M_0$  black holes in galactic nuclei be primordial?

BBNS => t < 1 s => M <  $10^{5}M_{\odot}$  .....but  $\beta$  <  $10^{-6}$  (t/s)<sup>1/2</sup>

Supermassive PBHs could also generate cosmic structures on larger scale through 'seed' or 'Poisson' effect

Upper limit on  $\mu$  distortion of CMB excludes  $10^4 < M/M_O < 10^{12}$  for Gaussian fluctuations but some models evades these limits. Otherwise need accretion factor of  $(M/10^4M_o)^{-1}$ 

## Limits on primordial black holes from $\mu$ distortions in cosmic microwave background

Tomohiro Nakama,<sup>1</sup> Bernard Carr,<sup>2,3</sup> and Joseph Silk<sup>1,4,5</sup>

#### PHYSICAL REVIEW D 97, 043525 (2018)

If primordial black holes (PBHs) form directly from inhomogeneities in the early Universe, then the number in the mass range  $10^5 - 10^{12} M_{\odot}$  is severely constrained by upper limits to the  $\mu$  distortion in the cosmic microwave background (CMB). This is because inhomogeneities on these scales will be dissipated by Silk damping in the redshift interval  $5 \times 10^4 \leq z \leq 2 \times 10^6$ . If the primordial fluctuations on a given mass scale have a Gaussian distribution and PBHs form on the high- $\sigma$  tail, as in the simplest scenarios, then the  $\mu$  constraints exclude PBHs in this mass range from playing any interesting cosmological role. Only if the fluctuations are highly non-Gaussian, or form through some mechanism unrelated to the primordial fluctuations, can this conclusion be obviated.

![](_page_53_Figure_4.jpeg)

## SEED AND POISSON FLUCTUATIONS

PBHs larger than  $10^{2}M_{O}$  cannot provide dark matter but can affect large-scale structure through seed effect on small scales or Poisson effect on large scales even if f small.

If region of mass M contains PBHs of mass m, initial fluctuation is

$$\delta_i \sim \begin{cases} m/M & \text{(seed)} \\ (fm/M)^{1/2} & \text{(Poisson)} \end{cases}$$

f = 1 => Poisson dominates; f <<1 => seed dominates for M < m/f. Fluctuation grows as  $z^{-1}$  from  $z_{eq} \sim 10^4$ , so mass binding at  $z_B$  is

$$M \approx \begin{cases} 4000 \, m z_B^{-1} \quad \text{(seed)} \\ 10^7 f m z_B^{-2} \quad \text{(Poisson)} \end{cases}$$

## SEED VERSUS POISSON

![](_page_55_Figure_1.jpeg)

 $f = 1 => m < 10^3 M_O => M < 10^{11} z_B^{-2} M_O < M_{gal}$  (Poisson)

=> can generate dwarf galaxies

f << 1 => M can be larger

=> PBHs can be seeds for galaxies

Can constrain PBH scenarios by requiring that various cosmic structures do not form too early

![](_page_56_Figure_1.jpeg)

#### First clouds bind earlier than in standard model

## SUPERMASSIVE PBHS AS SEEDS FOR GALAXIES

Seed effect =>  $M_B \sim 10^3 \text{ m} (z_B/10)$  $\Rightarrow$  naturally explain  $M_{BH}/M_{bulge}$  relation

Effect of mergers and accretion?

![](_page_57_Figure_3.jpeg)

Also predict mass function of galaxies (cf. Press-Schechter)

 $dN_g/dM \propto M^{-2} \exp(-M/M_*) \qquad M_* \sim 10^{12} M_{\odot}$ 

If M exceeds 10<sup>10</sup>M<sub>O</sub>, black hole accretes whole galaxy

### PRIMORDIAL BLACK HOLES IN A DIMENSIONALLY REDUCED UNIVERSE

## A. Tzikas, P. Nicolini, J. Mureika and B. Carr JCAP 1812(2018), 033; arXiv:1811.09518

#### Abstract

We investigate the spontaneous creation of primordial black holes in a lowerdimensional expanding early universe. We use the no-boundary proposal to construct instanton solutions for both the background and a black hole nucleated inside this background. The resulting creation rate could lead to a significant population of primordial black holes during the lower dimensional phase. We also consider the subsequent evaporation of these dimensionally reduced black holes and find that their temperature increases with mass, whereas it decreases with mass for 4-dimensional black holes. This means that they could leave stable sub-Planckian relics, which might in principle provide the dark matter. Universe may have been (1+1)-dimensional at early times

![](_page_59_Figure_1.jpeg)

=> PBH production through gravitational instanton effects

In 3+1 case  $\Gamma \sim \exp(-\pi/\Lambda) << 1$  now (Bousso & Hawking 1996) In 1+1 case:

![](_page_59_Figure_4.jpeg)

## PBH PRODUCTION RATE

$$M \ll \sqrt{|\Lambda|} \implies T \approx M/(2\pi), \quad \Gamma_{\rm L} \approx (\Lambda/M^2)^{1/2} \ll 1$$
  
$$\sqrt{|\Lambda|} < M < \sqrt{2|\Lambda|} \implies T \approx \sqrt{|\Lambda|}/(2\pi), \quad \Gamma_{\rm L} \approx 1$$
  
$$M = \sqrt{|\Lambda|}(1+\epsilon) \implies T \approx \sqrt{\epsilon|\Lambda|}/(\sqrt{2}\pi) \ll \sqrt{|\Lambda|}, \quad \Gamma_{\rm L} \approx \epsilon^{-1/2} \gg 1$$
  
$$M = \sqrt{|\Lambda|} \implies T = 0, \quad \Gamma_{\rm N} = (\mu_0/\sqrt{|\Lambda|}) \quad \text{(Nariai)}$$

![](_page_60_Figure_2.jpeg)

PBH EVAPORATION 
$$\beta = T^{-1} = \left(\frac{1}{2\pi}\sqrt{M^2 - |\Lambda|}\right)^{-1}$$
$$M \gg \sqrt{|\Lambda|} \implies R_{\rm S} \approx \frac{1}{2M}, \quad T \approx \frac{M}{2\pi} \implies L \sim R_{\rm S}^2 T^4 \sim \gamma M^2$$

=>  $M = \frac{M_{\rm i}}{1 + \gamma M_{\rm i}(t_o - t_{\rm i})} \approx \frac{M_{\rm i}}{1 + M_{\rm i}/M_*}$  with  $M_* \sim 1/(\gamma t_0) \sim \hbar/(c^2 t_0) \sim 10^{-65}$  g

## But CMB suppresses evaporation above

$$M_{\rm CMB} = 10^{-36} (T_{\rm CMB}/3K) \,{\rm g}$$

Non-evaporating Nariai PBHs have

$$\rho_N(\Lambda_i) \sim M^2 \frac{dn_N}{dM} \sim \Gamma(\Lambda_i) \sim \frac{\mu_0}{\sqrt{\Lambda_i}}$$

and could provide dark matter

![](_page_61_Figure_7.jpeg)

Transition to 3+1 may generate  $\Lambda$  with density comparable to DM

![](_page_62_Figure_0.jpeg)

POPULARITY

### CONCLUSIONS

PBHs have been invoked for four roles:

Cosmic rays

Dark matter

LIGO events

Cosmic structure

These are distinct roles but with extended mass function PBHs could fulfill all!

This talk is dedicated to the memory of Stephen Hawking. He wrote the first paper on primordial black holes in 1971. If they play any of the roles discussed here, this may have been his most prescient and important work

![](_page_63_Picture_8.jpeg)