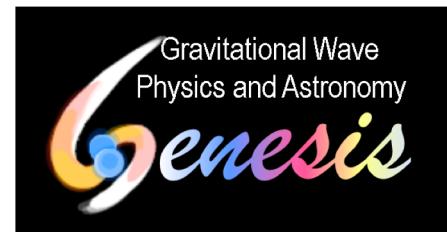
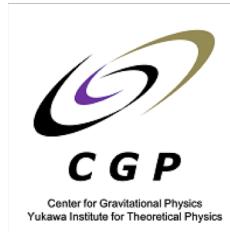
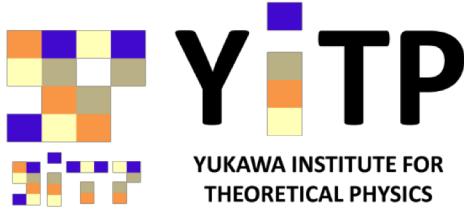


Constant-roll inflation in scalar-tensor theory

Hayato Motohashi

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2019.03.08 Accelerating Universe in the Dark, YITP



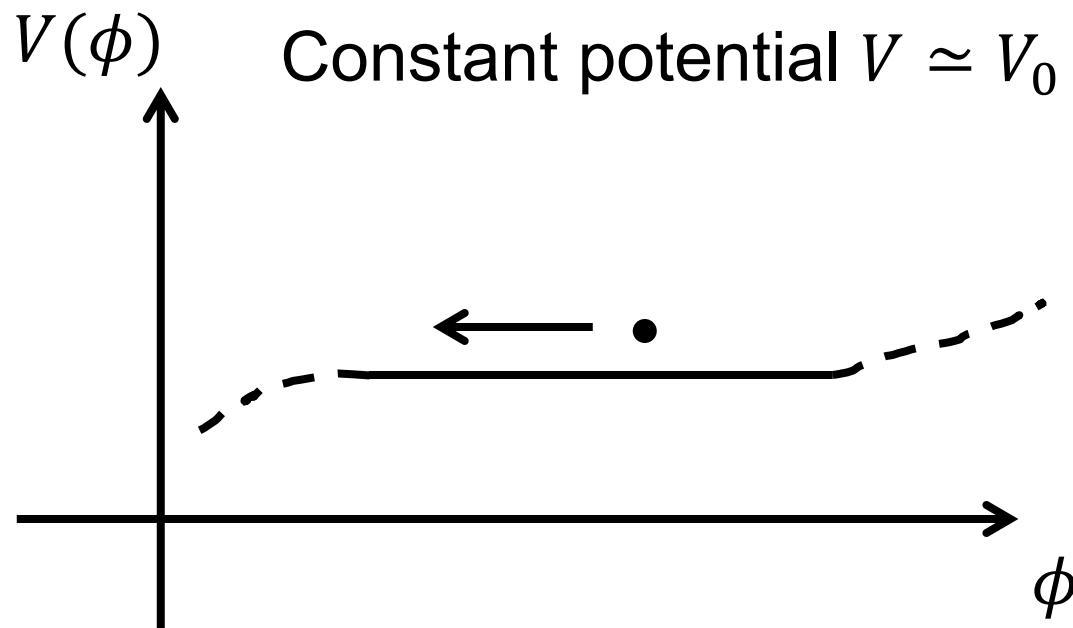
Canonical single field Inflation

	$\simeq 0$	Slow-roll
$\ddot{\phi}/(H\dot{\phi})$	$= -3$	Ultra slow-roll
	$= \beta$ (constant)	Constant-roll



Ultra slow-roll inflation

Kinney, gr-qc/0503017



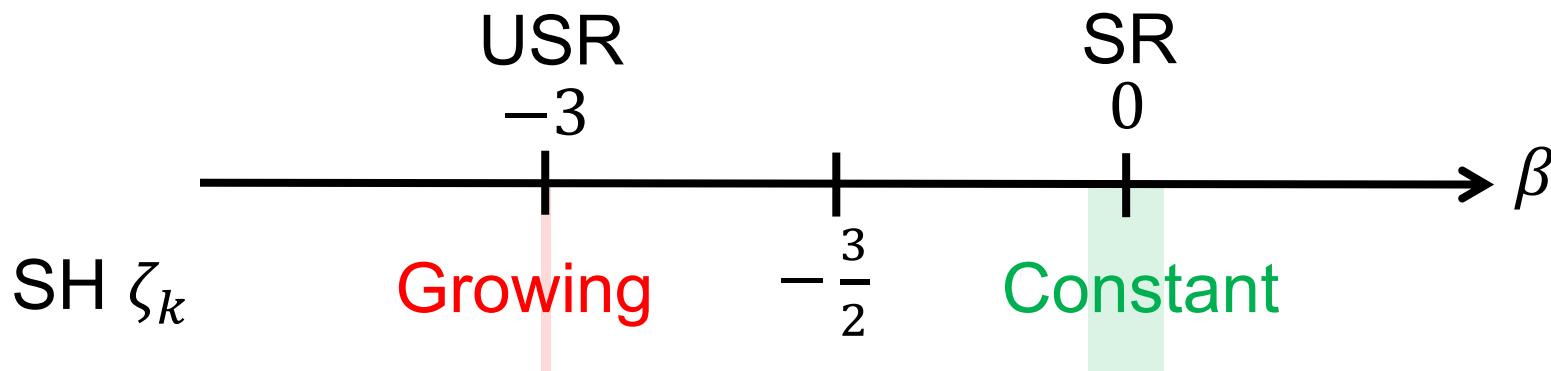
$$\begin{aligned}\ddot{\phi} = -3H\dot{\phi} &\Rightarrow \dot{\phi} \propto a^{-3} \\ &\Rightarrow \epsilon_H \propto a^{-6}\end{aligned}$$

Ultra slow-roll inflation

Superhorizon solution

$$\zeta_k = A_k + B_k \int \frac{dt}{a^3 \epsilon_H}$$

Constant mode ↗
Slow-roll $\epsilon_H \simeq \text{const.} \ll 1$ Decaying mode
Ultra slow-roll $\epsilon_H \propto a^{-6}$ Growing mode
While $\epsilon_H \ll 1$, $d \ln \epsilon_H / dN = -6$ violates slow roll.



Generalization

Martin, HM, Suyama, 1211.0083

Superhorizon solution

$$\zeta_k = A_k + B_k \int \frac{dt}{a^3 \epsilon_H}$$

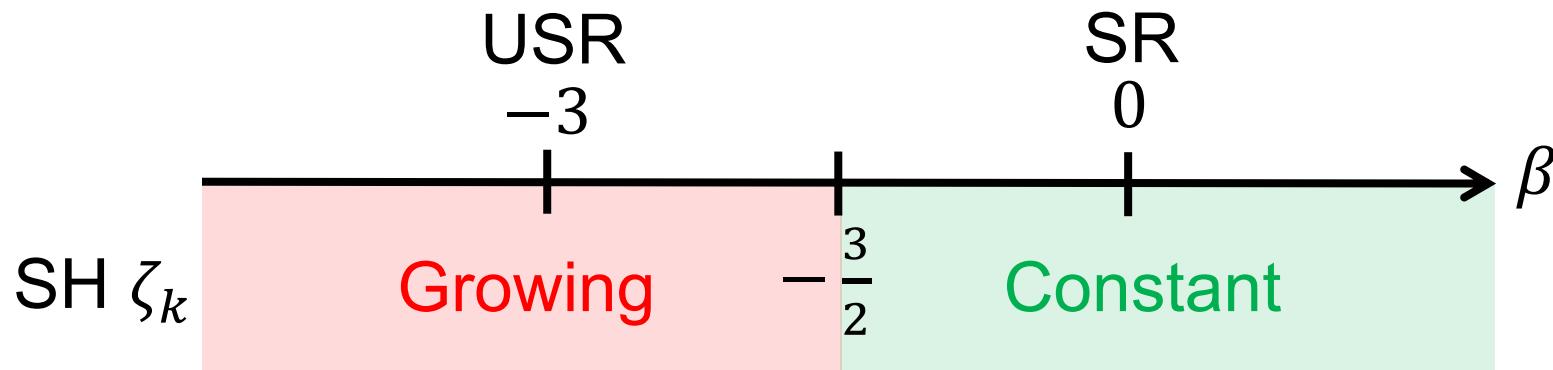
Constant mode

Decaying mode

Growing mode

Constant roll $\epsilon_H \propto a^{2\beta}$

$2\beta > -3$



Constant-roll potential

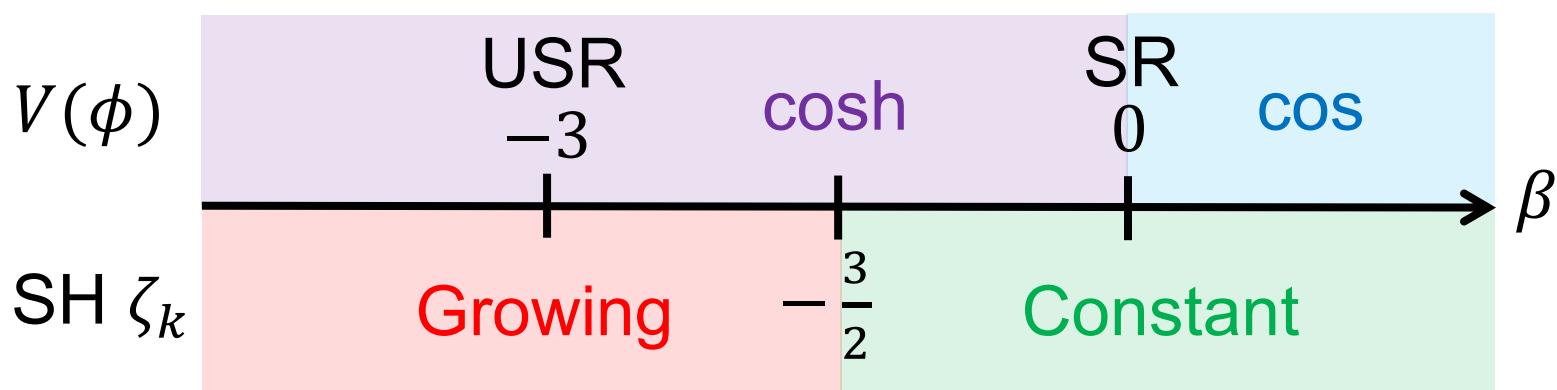
Allows constant-roll evolution as an exact solution

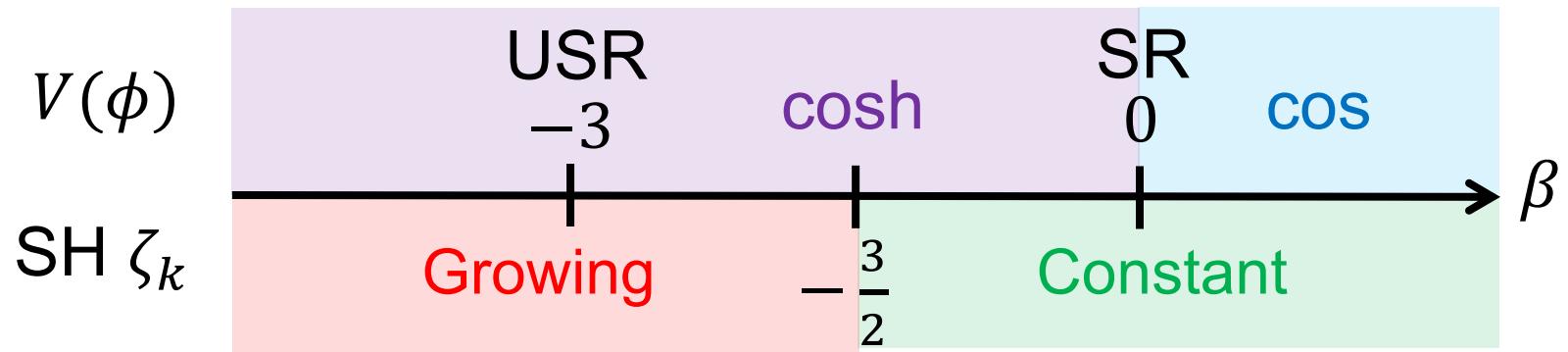
- a) $V \propto e^{\sqrt{-2\beta}\phi}$ with $\beta < 0$: Power-law inflation

X $r = 8(1 - n_s)$

- b) $V \propto \cosh(\sqrt{-2\beta}\phi) + \text{const.}$ with $\beta < 0$

- c) $V \propto \cos(\sqrt{2\beta}\phi) + \text{const.}$ with $\beta > 0$





✓ ζ_k frozen on SH

- Attractor

- n_s

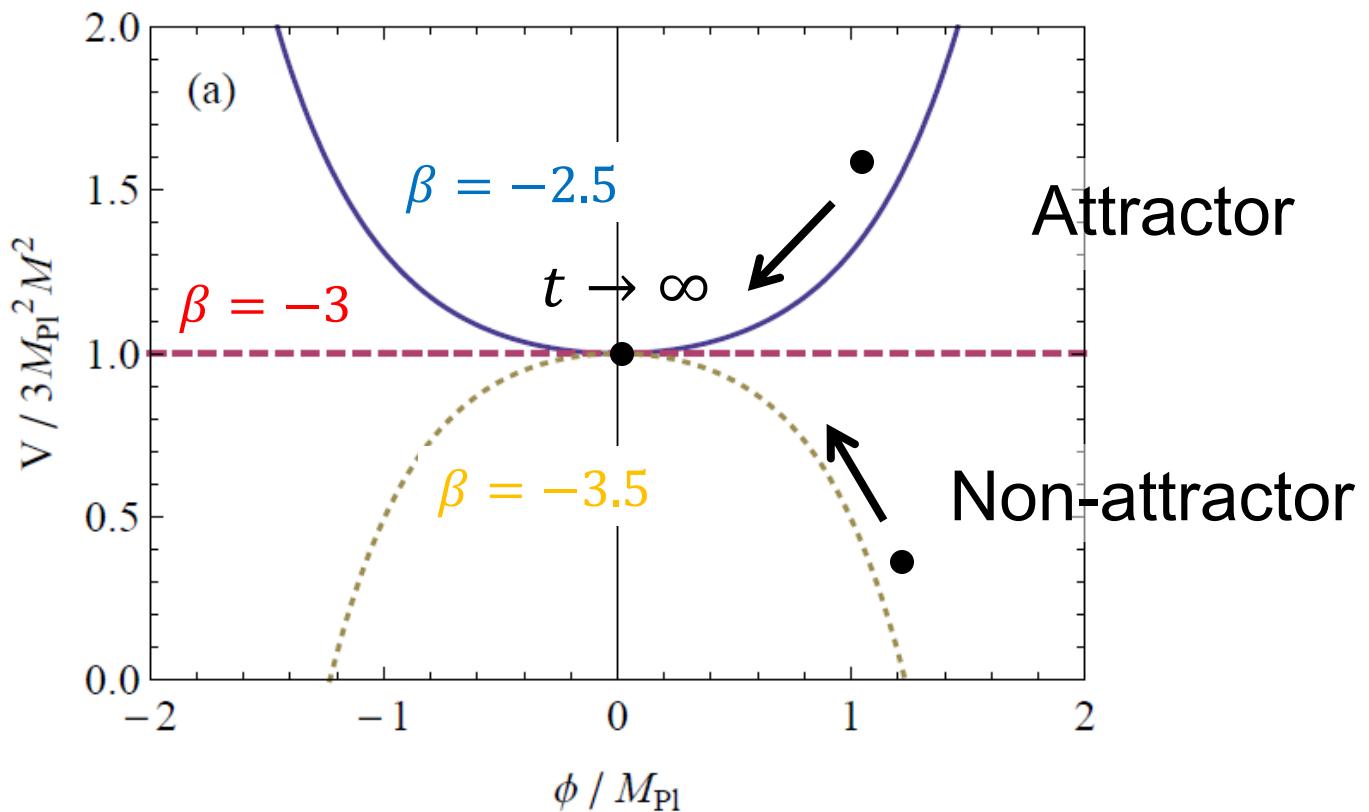
- r

- $N \sim 60$

cosh potential ($\beta < 0$)

$$\frac{V(\phi)}{M^2 M_{Pl}^2} = 3 \left[1 - \frac{3 + \beta}{6} \left\{ 1 - \cosh \left(\sqrt{-2\beta} \frac{\phi}{M_{Pl}} \right) \right\} \right]$$

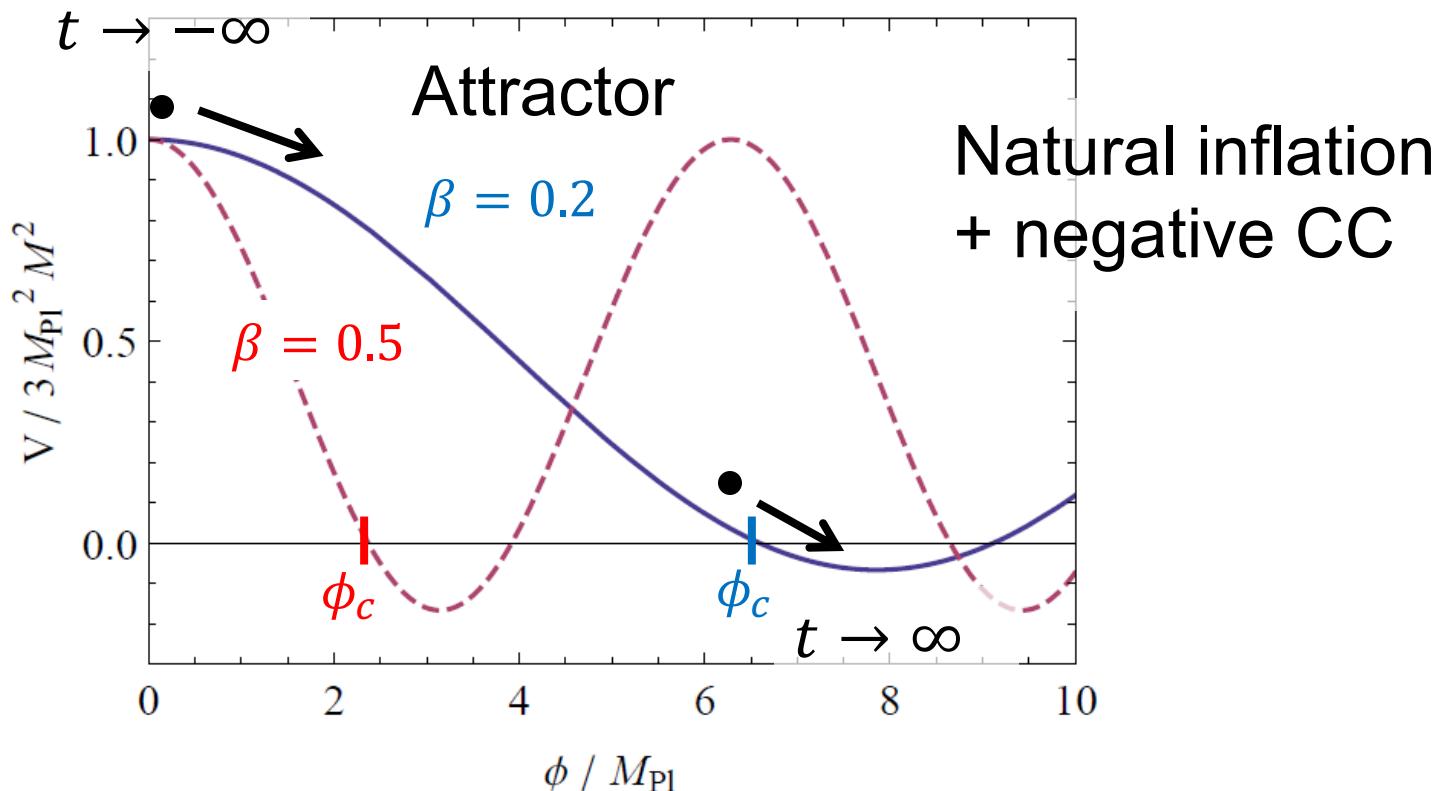
Assume a transition to reheating at $\phi > 0$.

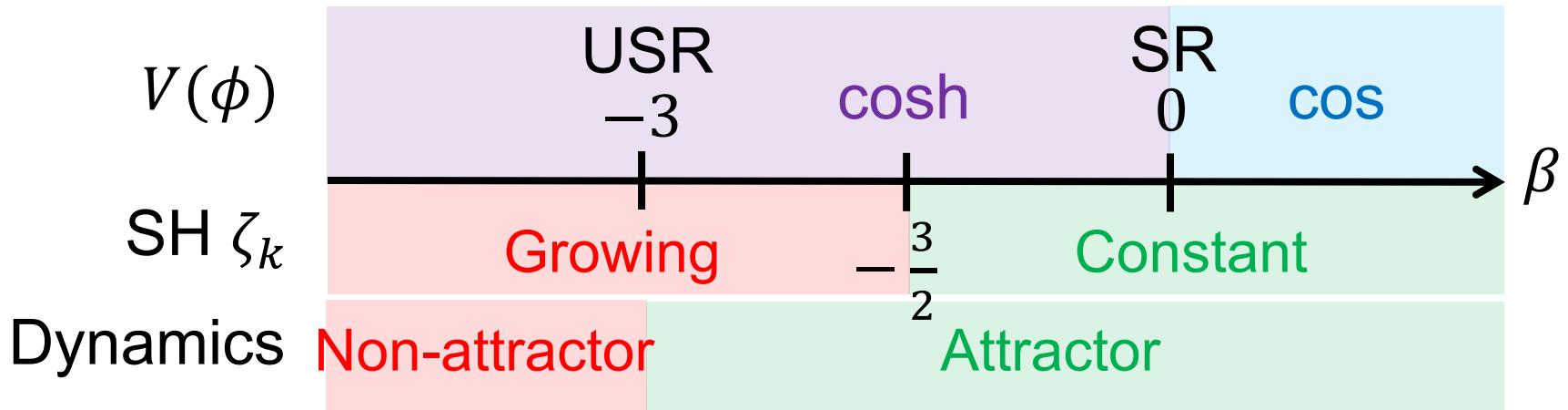


cos potential ($\beta > 0$)

$$\frac{V(\phi)}{M^2 M_{Pl}^2} = 3 \left[1 - \frac{3 + \beta}{6} \left\{ 1 - \cos \left(\sqrt{2\beta} \frac{\phi}{M_{Pl}} \right) \right\} \right]$$

Assume a transition to reheating at $\phi < \phi_c$.





✓ ζ_k frozen on SH

✓ Attractor

- n_s

- r

- $N \sim 60$

Curvature perturbation

Mukhanov-Sasaki equation

$$\nu_k'' + \left(k^2 - \frac{z''}{z} \right) \nu_k = 0$$

Using analytic solutions for cosh and cos potentials,

$$\frac{z''}{z} \rightarrow \frac{\nu^2 - 1/4}{\tau^2} \quad (t \rightarrow \pm\infty)$$

where $\nu \equiv |\beta + 3/2|$ without slow-roll approximation.

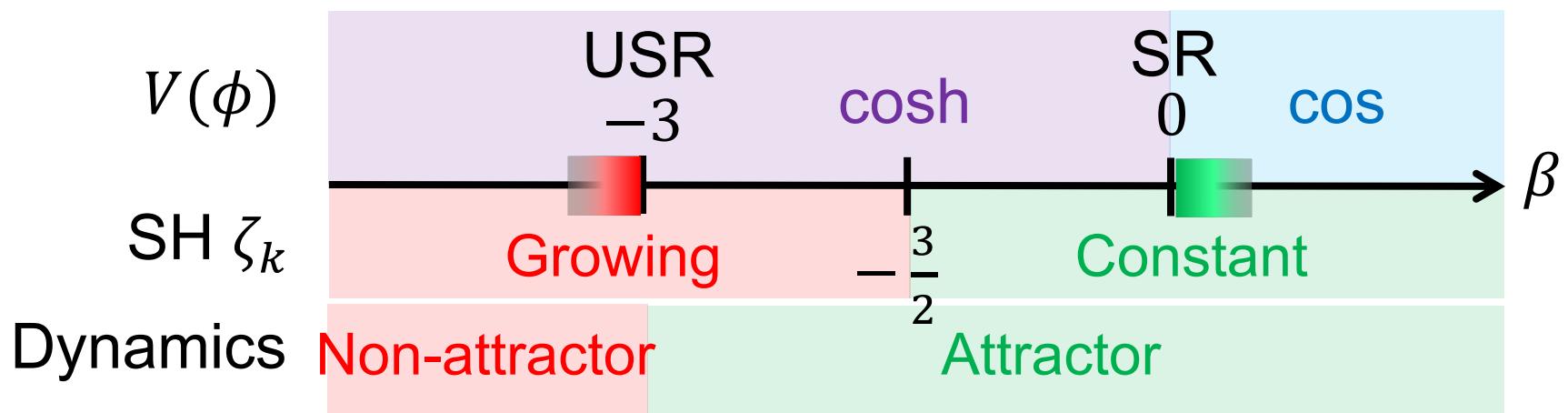
Curvature perturbation

Spectral index

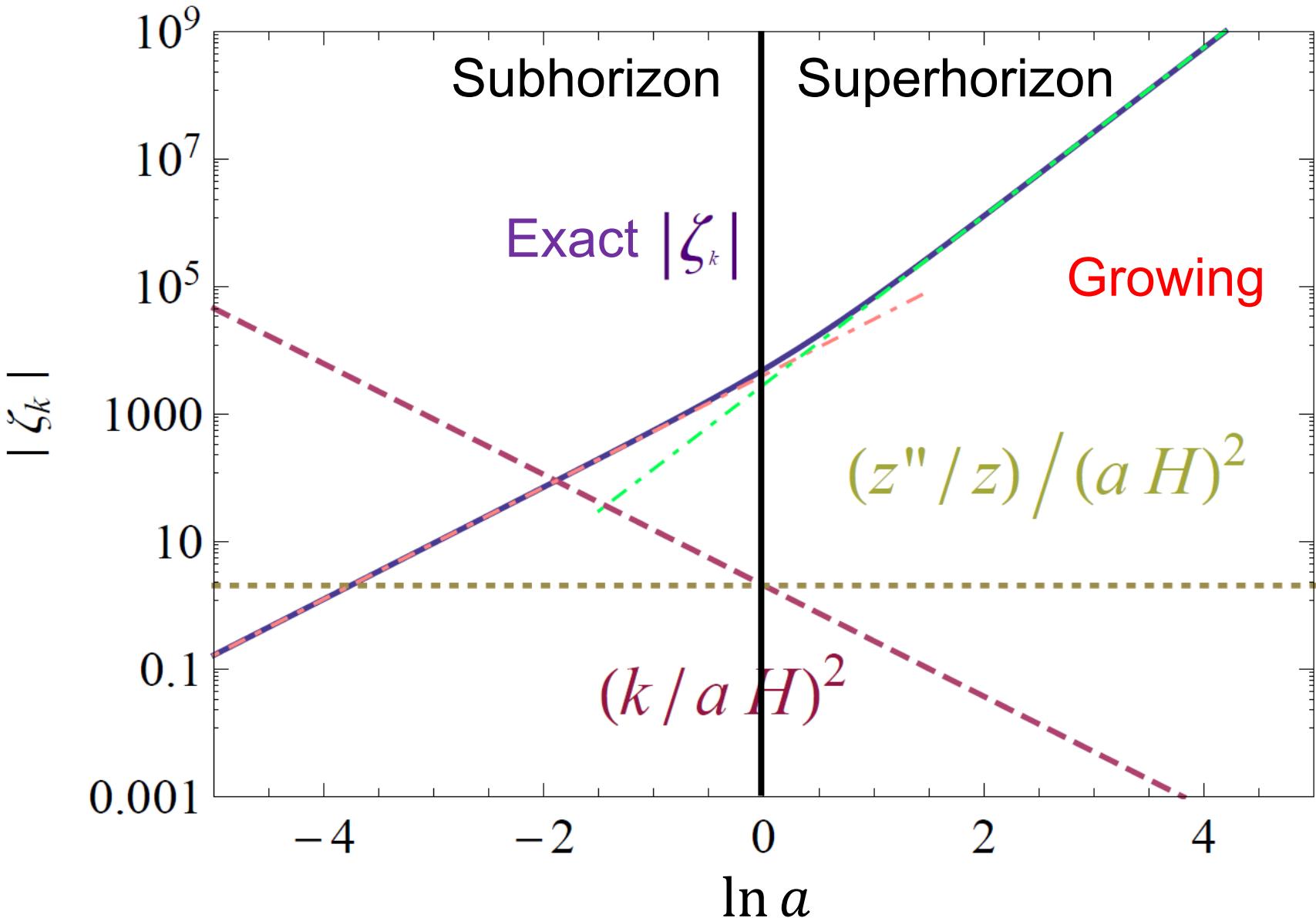
$$n_s - 1 = 3 - 2\nu = 3 - |2\beta + 3|$$

For slightly red-tilted spectrum e.g. $n_s = 0.96$,

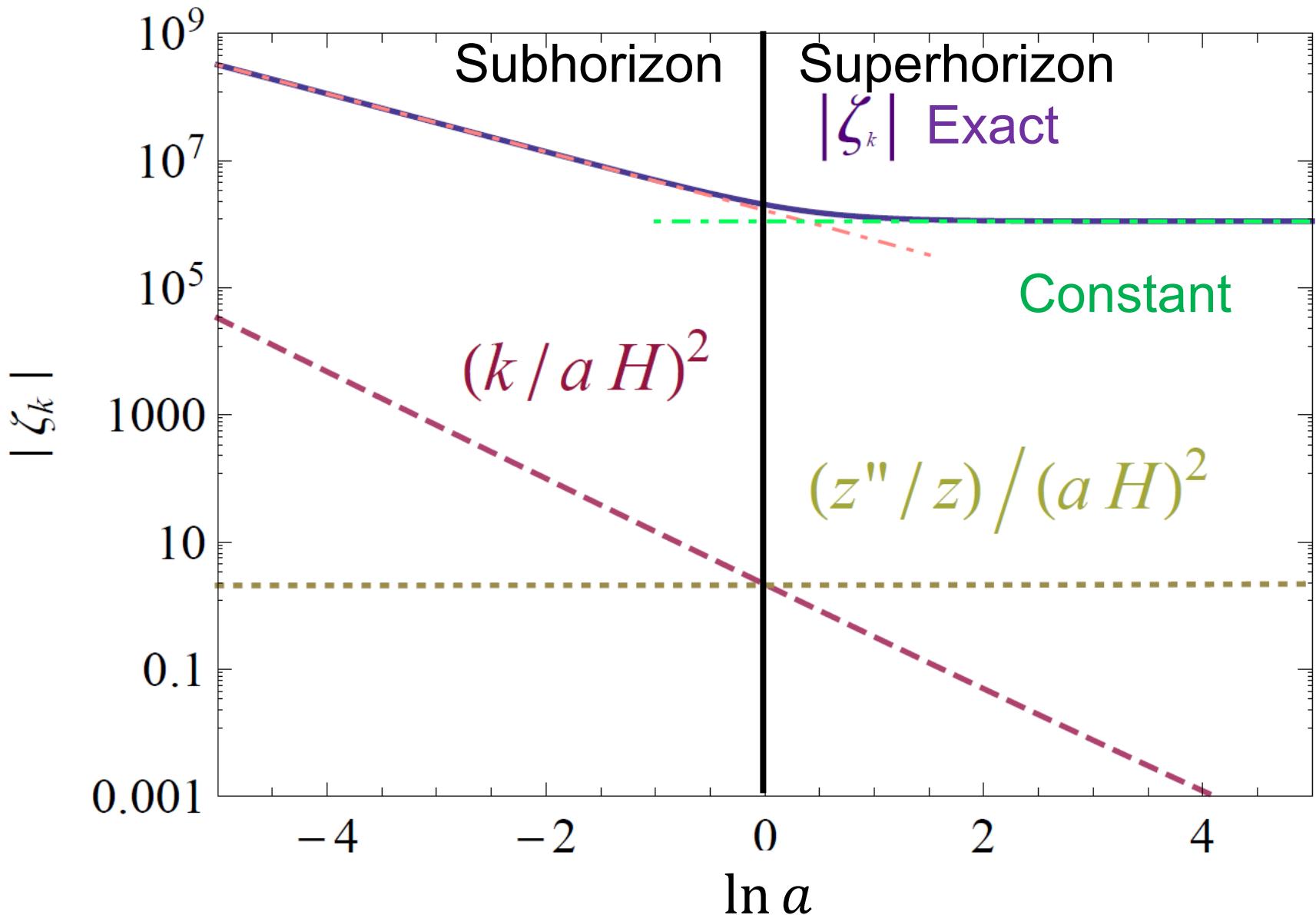
$$\beta = \textcolor{red}{-3.02} \text{ or } \textcolor{green}{0.02}$$

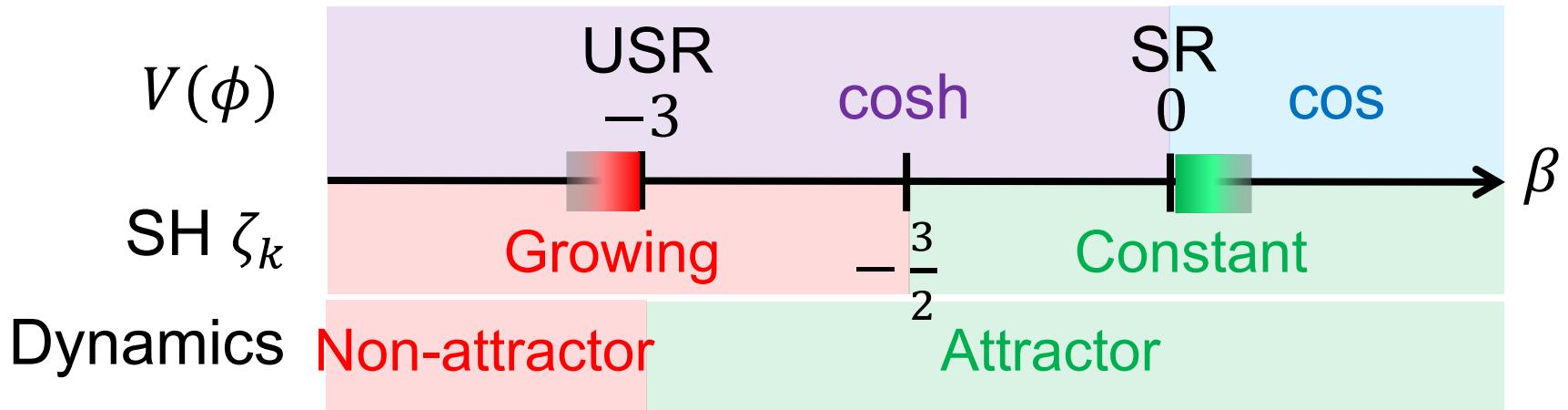


cosh potential $\beta = -3.02$



cos potential $\beta = 0.02$





✓ ζ_k frozen on SH

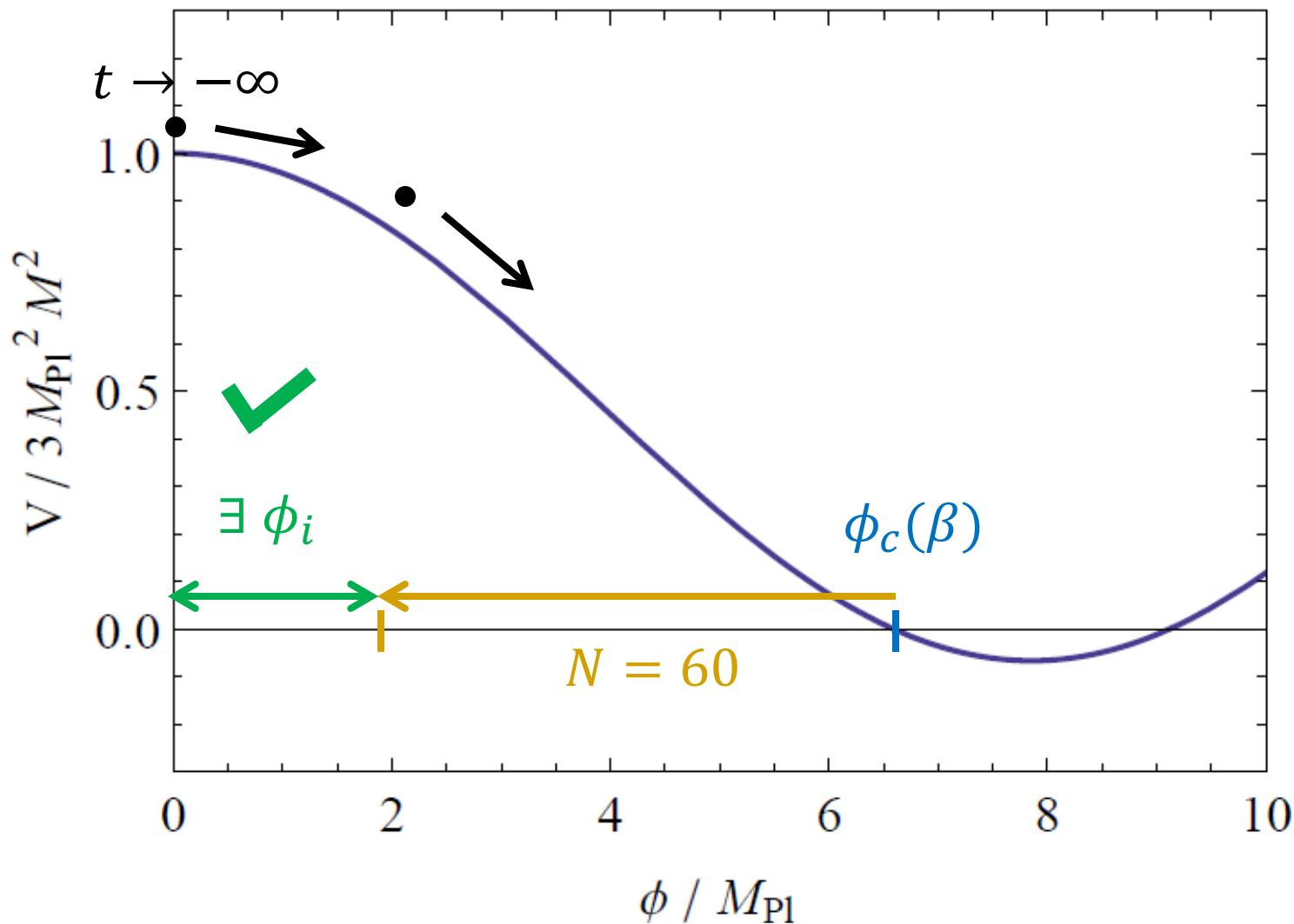
✓ Attractor

✓ n_s

- r

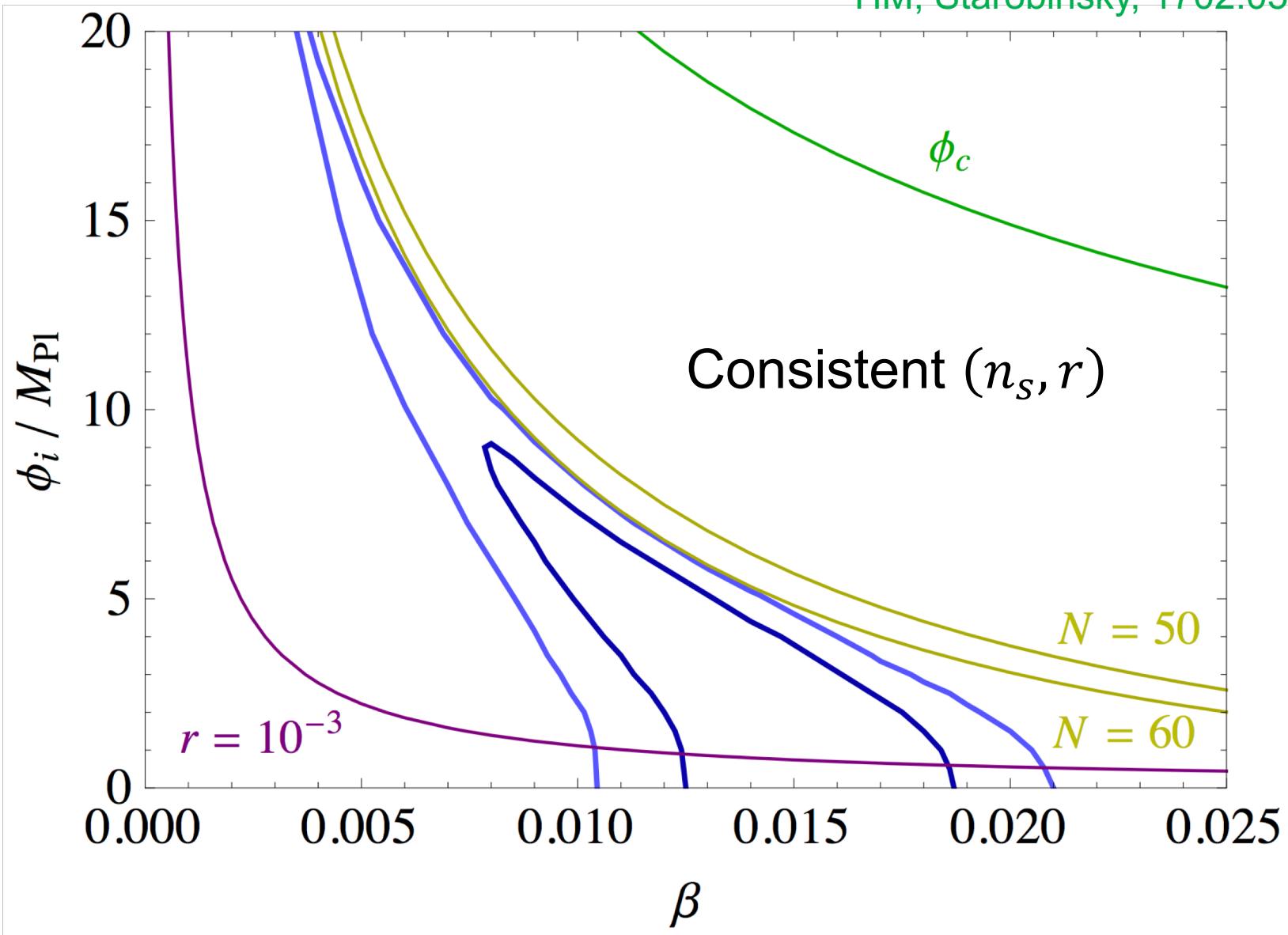
- $N \sim 60$

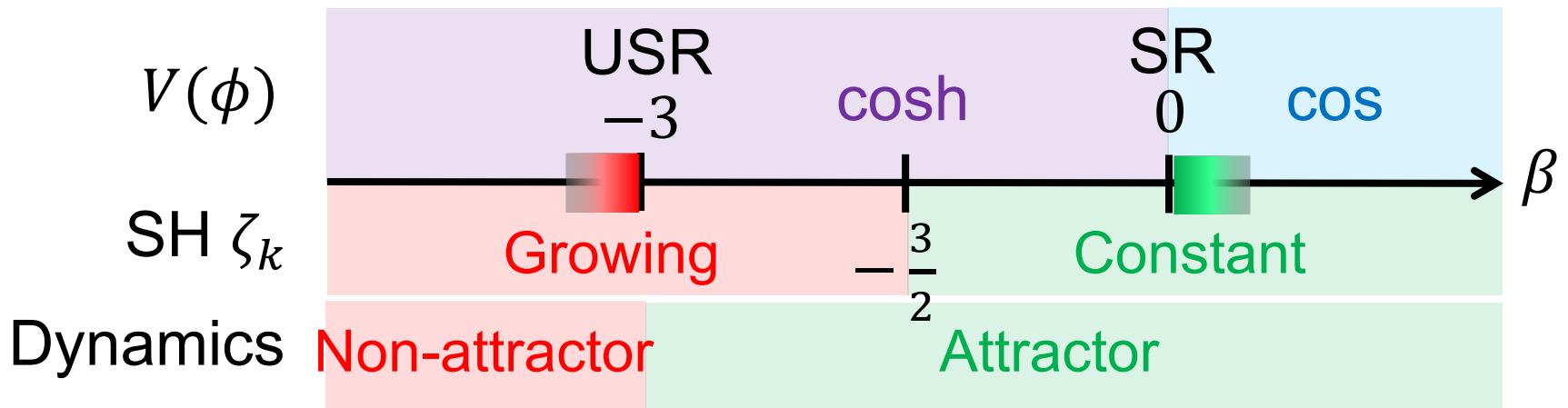
Sufficient number of efolds



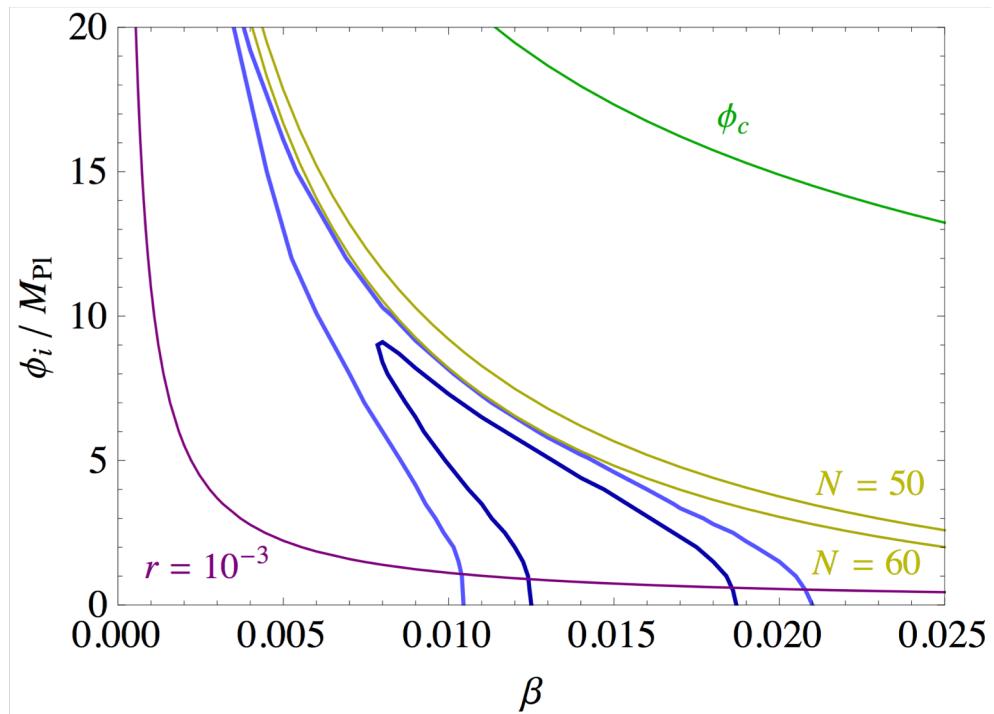
Observational constraint

HM, Starobinsky, 1702.05847





- ✓ ζ_k frozen on SH
- ✓ Attractor
- ✓ n_s
- ✓ r
- ✓ $N \sim 60$



$f(R)$ constant-roll inflation

HM, Starobinsky, 1704.08188

$$S = \int d^4x \sqrt{-g} \frac{f(R)}{2}$$

Constant-roll condition in JF (different from $\ddot{\phi} = \beta H \dot{\phi}$)

$$\ddot{F} = \beta H \dot{F} \quad (F \equiv df/dR)$$

Analytic solution for

✓ $f(R)$ (parametric form)

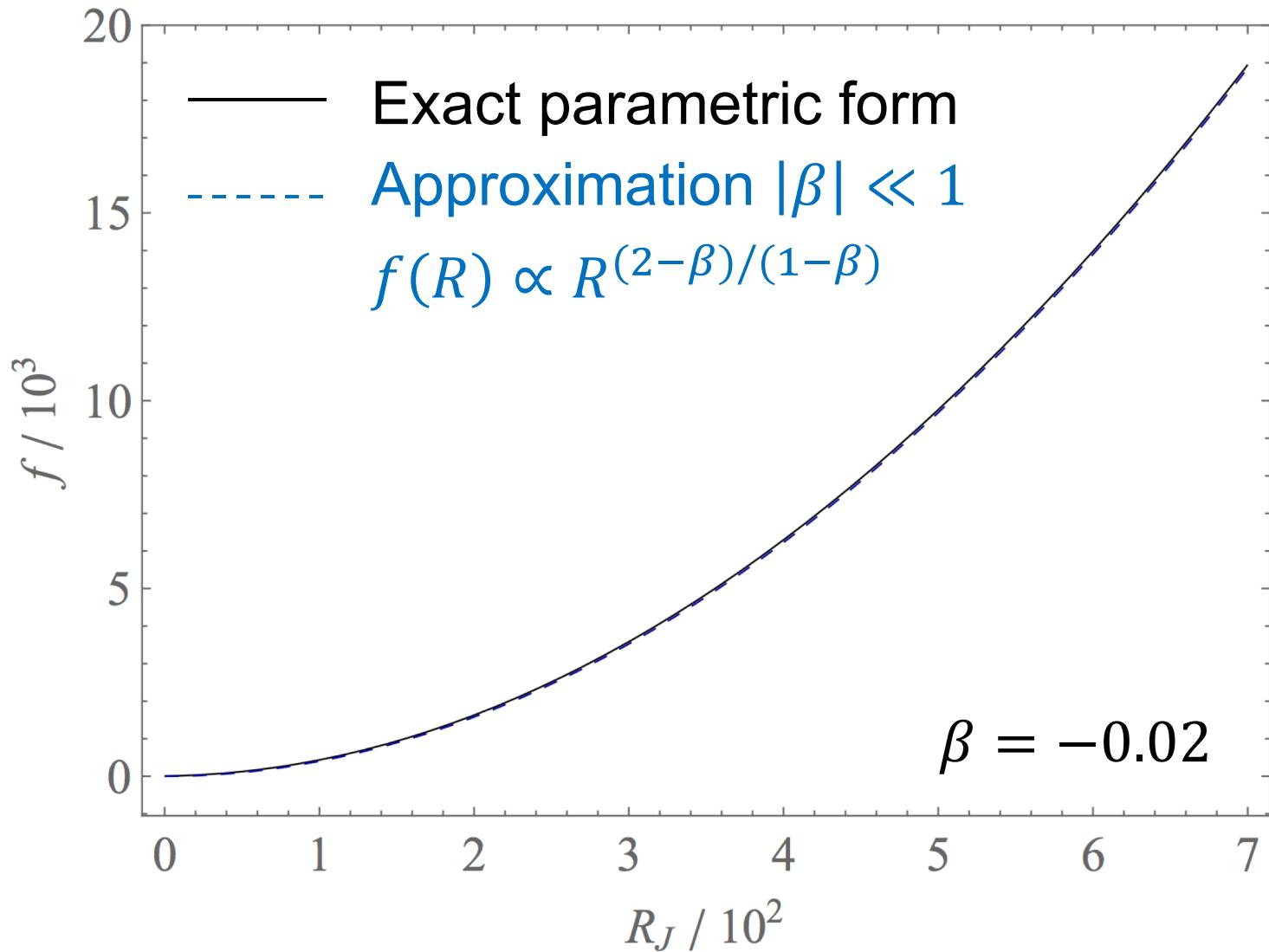
✓ $V(\phi)$

✓ $\phi(t_E)$

} Einstein frame $g_{\mu\nu}^E = F g_{\mu\nu}^J, F = e^{\sqrt{\frac{2}{3}}\phi}$

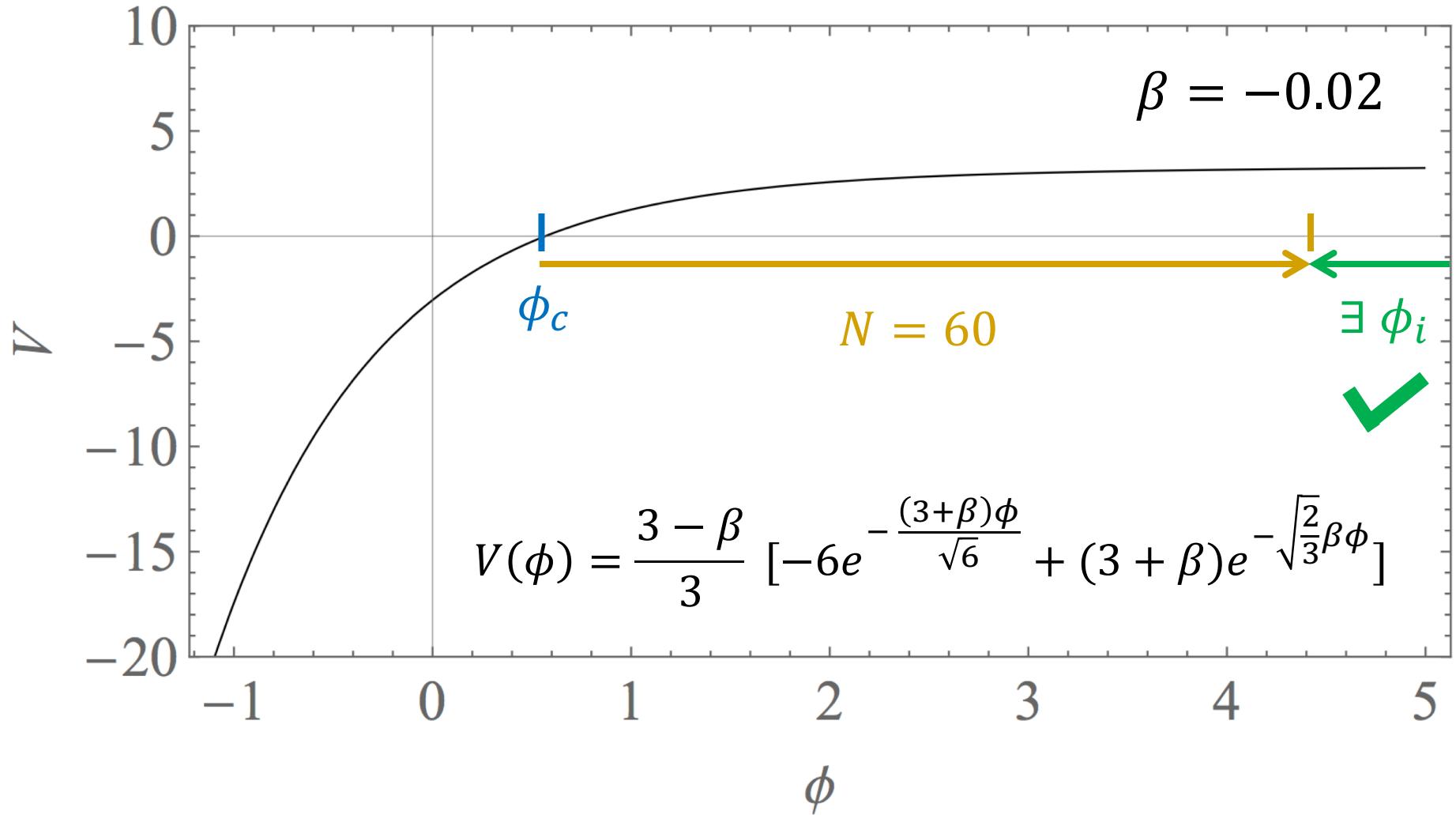
$f(R)$

HM, Starobinsky, 1704.08188



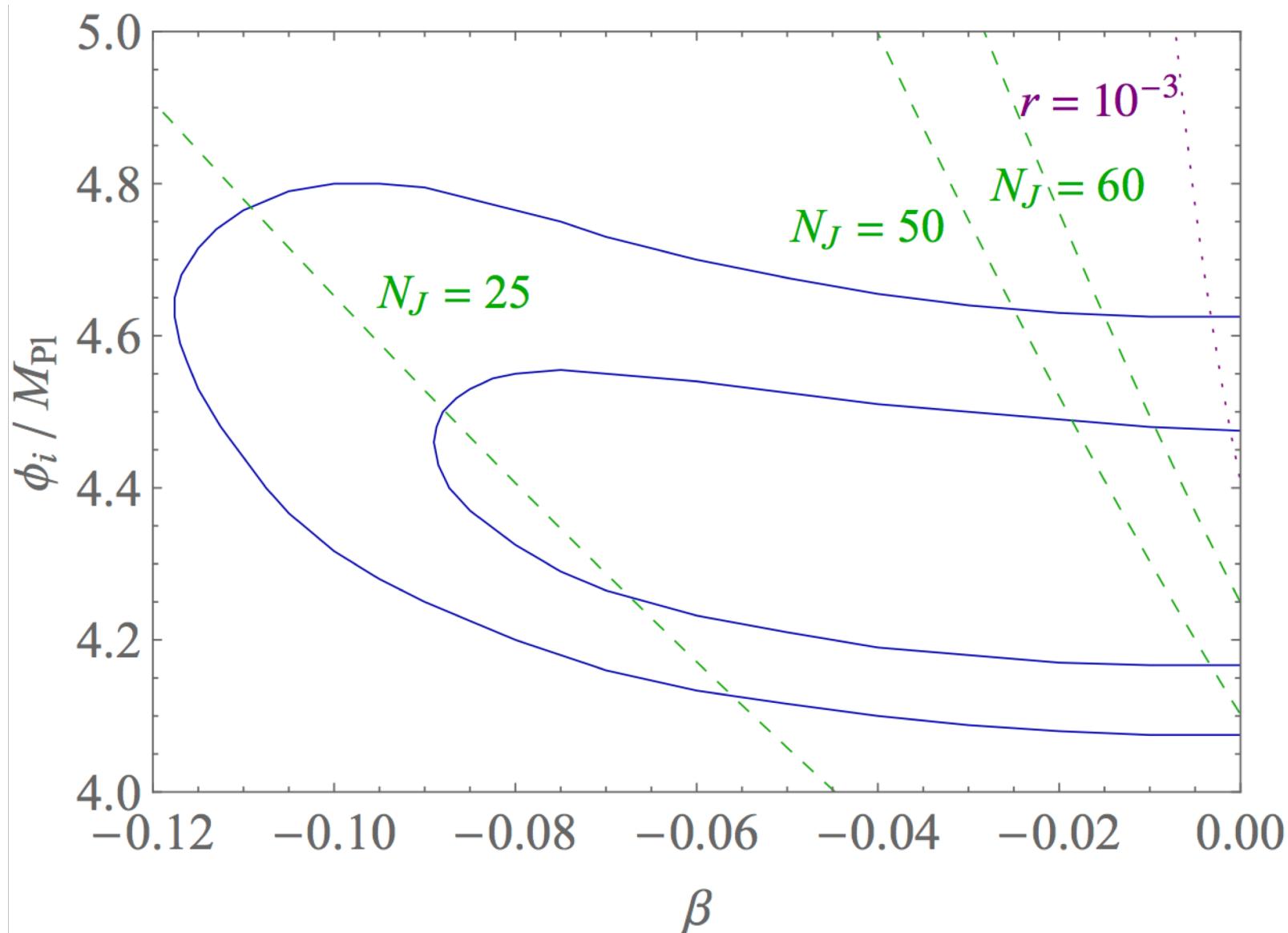
Einstein frame potential

HM, Starobinsky, 1704.08188



Observational constraint

HM, Starobinsky, 1704.08188



Summary

Generalizing SR and USR we explored constant roll

$$\ddot{\phi} = \beta H \dot{\phi} \quad \text{or} \quad \ddot{F} = \beta H \dot{F}$$

which allows analytic solutions of

$$V(\phi), \phi(t), H(t), a(t) \text{ and } f(R)$$

- ✓ Attractor
- ✓ ζ_k frozen on superhorizon scales
- ✓ $N \sim 60$
- ✓ n_s & r

Constant roll + obs. constraint
→ Slow roll is preferred.

$\beta \sim 0.02$	Canonical
-0.02	$f(R)$
± 0.04	ST