Higgs Precision Measurement at the ILC







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well documented Higgs physics at the ILC

- Report by the Committee on the Scientific Case of the ILC operating at 250 GeV as a Higgs Factory
- Physics Case for the 250 GeV Stage of the ILC

• Physics Case for the International Linear Collider

- Fujii, et al, arXiv: 1506.05992
- Baer, et al, arXiv: 1306.6352

Asner, et al, arXiv: 1310.0763

- ILC TDR Volume 2: Physics
- ILC Higgs White Paper

Asai, et al, arXiv: 1710.08639

Fujii, et al, arXiv: 1710.07621

outline

• introduction: Higgs as a window for BSM

• Precision Higgs Physics at ILC250

(focus of this talk)

• Higgs Physics at ILC beyond 250 GeV

• ILC status

Higgs as a window for BSM



o What is the origin of EWSB?o Naturalness of the light scalar mass?

o Connection to Dark Matter, BAU, inflation?



can be revealed by looking in detail at Higgs properties

Higgs as a window for BSM

o fingerprint BSM by patterns of deviations

—> measure as many couplings as possible



deviation is typically small, 1-10% for mBSM~1TeV

—> need 1% or below measurement

fingerprint 4 types of 2HDM: gHbb .vs. gHττ



Kanemura, et al, arXiv: 1406.3294

synergy with LHC: direct and indirect discoveries



Cahill-Rowley, et al, arXiv:1308.0297

Figure 8: Histograms of the ratio $r_{bb} = \Gamma(h \to \bar{b}b)/\Gamma(h \to \bar{b}b)_{\rm SM}$ within a scan of the approximately 250,000 supersymmetry parameter sets after various stages of the LHC, assuming the LHC does not find direct evidence for supersymmetry. The purple histogram shows parameter points that would not be discovered at future upgrades of the LHC (14 TeV and $3 \, {\rm ab}^{-1}$ integrated luminosity). From [38].

ILC250: International Linear Collider at √s=250 GeV

o ILC was proposed as a 500 GeV e+e- machine in TDR (2013)

 in view of the new results by LHC Run 2 and the cost reduction issue, ILC is now proposed as a 250 GeV machine (ILC250 as a Higgs Factory) hopping for an early realization

∫Ldt=2 ab⁻¹, ~0.5M Higgs, P(e⁻,e⁺)=(80%,30%)

• Higgs couplings in effective theory formalism -> view of coupling measurement at 250 GeV is somewhat changed

How Higgs is produced at e+e-



• at $\sqrt{s} \sim 250$ GeV e+e- -> ZH attains its cross section maximum

what are the direct Higgs observables at ILC250

 $\mathbf{\overline{O}}$ $\mathbf{\overline{O}}$ ZH $\sigma_{ZH} \times Br(H \longrightarrow bb), \sigma_{\nu\nu H} \times Br(H \longrightarrow bb)$ $\sigma_{ZH} \times Br(H \rightarrow cc)$ $\mathbf{\sigma}_{ZH} \times Br(H \longrightarrow gg)$ $\sigma_{ZH} \times Br(H \longrightarrow WW^*)$ $\sigma_{ZH} \times Br(H \longrightarrow ZZ^*)$ $\boldsymbol{\boldsymbol{\sigma}} \sigma_{ZH} \times Br(H \rightarrow \tau \tau)$ $+ d\sigma/d\Omega$ σ $\sigma_{ZH} \times Br(H \longrightarrow \gamma \gamma)$ $\sigma \sigma_{ZH} \times Br(H \rightarrow \mu \mu)$ $\sigma_{ZH} \times Br(H \longrightarrow Invisible)$

what are the direct Higgs observables at ILC250



note the important complementarity with LHC

expected accuracies of Higgs observables at ILC250

(based on full detector simulations for ILD and SiD)

$-80\% \ e^-, \ +30\% \ e^+$	polarization:					
	$250 { m GeV}$		$350~{\rm GeV}$		$500 { m GeV}$	
	Zh	$ u \overline{ u} h$	Zh	$ u \overline{ u} h$	Zh	$ u\overline{ u}h$
σ [50–53]	2.0		1.8		4.2	
$h \rightarrow invis.$ [54, 55]	0.86		1.4		3.4	
$h \to b\overline{b}$ [56–59]	1.3	8.1	1.5	1.8	2.5	0.93
$h \to c\overline{c} \ [56, 57]$	8.3		11	19	18	8.8
$h \rightarrow gg \ [56, 57]$	7.0		8.4	7.7	15	5.8
$h \rightarrow WW \ [59-61]$	4.6		5.6 *	5.7 *	7.7	3.4
$h \to \tau \tau$ [63]	3.2		4.0 *	16 *	6.1	9.8
$h \rightarrow ZZ$ [2]	18		25 *	20 *	35 *	12 *
$h \to \gamma \gamma \ [64]$	34 *		39 *	45 *	47	27
$h \rightarrow \mu \mu \ [65, 66]$	72 *		87 *	160 *	120 *	100 *
a [27]	7.6		2.7 *		4.0	
b	2.7		0.69 *		0.70	
$\rho(a,b)$	-99.17		-95.6 *		-84.8	

(arXiv: 1708.08912; numbers are in %, for nominal ∫Ldt = 250 fb⁻¹)

expected accuracies of Higgs observables at ILC250

(based on full detector simulations for ILD and SiD)



(arXiv: 1708.08912; numbers are in %, for nominal ∫Ldt = 250 fb⁻¹)

- Precision Higgs Physics at ILC250
 - some key experimental observables
 - from observable to couplings: kappa & EFT
 - BSM models discrimination & next energy scale
 - Higgs CP, Higgs exotic decays

inclusive σ_{ZH} : the key of model independence





well defined initial states at e+erecoil mass technique -> tag Z only
Higgs is tagged without looking into H decay
absolute cross section of e+e⁻ -> ZH

for Z->II (leptonic recoil), Yan et al, arXiv:1604.07524; for Z->qq (hadronic recoil), Thomson, arXiv:1509.02853

what does model independence mean?



 \circ meas. of σ_{ZH} doesn't depend on how Higgs decays

O meas. of σ_{ZH} doesn't depend on underlying HZZ vertex

is it really possible? (Yes, but not trivial at all, see backup slides)

Higgs direct couplings to cc and gg at ILC250



e+e- -> ZH -> ff(jj): b-likeness .vs. c-likeness

Oclean environment at e+e-; excellent b- and c-tagging performance
 Obb/cc/gg modes can be separated simultaneously by template fitting

Ono, et. al, Euro. Phys. J. C73, 2343; F.Mueller, PhD thesis (DESY)

Higgs —> invisible at ILC250

$$e^+ + e^- \rightarrow ZH \rightarrow l^+l^-/q\bar{q} + \text{Missing}$$

Ishikawa, Kato, JT, et al



Orecoil technique: Higgs mass fully reconstructed even it decays invisibly

• right-handed beam polarization helps: much lower background • BR(H—>inv.) < 0.3% (CL95%)

From observables to couplings — Global Fit

$$\chi^{2} = \sum_{i=1}^{n} (\frac{Y_{i} - Y_{i}'}{\Delta Y_{i}})^{2}$$

Yi: measured values by experiments
Yi': predicted values by underlying theory
ΔYi: measurement uncertainty
n: number of independent observables

o kappa formalism

$$Y'_{i} = F_{i} \cdot \frac{g_{HA_{i}A_{i}}^{2} \cdot g_{HB_{i}B_{i}}^{2}}{\Gamma_{0}} \qquad (A_{i} = Z, W, t)$$
$$(B_{i} = b, c, \tau, \mu, g, \gamma, Z, W : decay)$$

$$g_{HXX} = \kappa_X \cdot g_{HXX}^{SM}$$

o effective field theory formalism

From observables to couplings - kappa formalism

1) recoil mass technique —> inclusive σ_{Zh}

2)
$$\sigma_{Zh} \longrightarrow \kappa_Z \longrightarrow \Gamma(h->ZZ^*)$$

- 3) WW-fusion $v_e v_e h \longrightarrow \kappa_W \longrightarrow \Gamma(h \rightarrow WW^*)$
- 4) total width $\Gamma_h = \Gamma(h \rightarrow ZZ^*)/BR(h \rightarrow ZZ^*)$
- 5) or $\Gamma_h = \Gamma(h \rightarrow WW^*)/BR(h \rightarrow WW^*)$
- 6) then all other couplings BR(h->AA) $^{*}\Gamma_{h} \rightarrow \kappa_{A}$

Q1: can we assume $\sigma(e+e-->Zh) \propto \Gamma(h->ZZ^*)$?



BSM territory -> can deviations be represented by single κ_Z ?

the answer is: No, if there is anomalous hZZ coupling

$$\delta \mathcal{L} = (1+\eta_Z) \frac{m_Z^2}{v} h Z_\mu Z^\mu + \zeta_Z \frac{h}{2v} Z_{\mu\nu} Z^{\mu\nu}$$





• $\sigma(e+e-->Zh) \propto \kappa^2 \propto \Gamma(h->ZZ^*)$ not any more: what would be a more model-independent formalism? Q2: if WW-fusion is so important, can we do precision Higgs physics at $\sqrt{s} = 250$ GeV?



WW-fusion is smaller by x10 than 500 GeV

a strategy: SM Effective Field Theory

$$\mathcal{L}_{\text{eff}} = \mathcal{L}_{\text{SM}} + \Delta \mathcal{L}$$

= $\mathcal{L}_{\text{SM}} + \sum_{i} \frac{c_i}{\Lambda^{d_i - 4}} O_i$

 O_i : dimension d_i operators, respect SU(3)xSU(2)xU(1) of L_{SM} c_i : Wilson coefficients Λ : EFT cutoff scale

 ΔL represent the most general effects of BSM physics

2, 84, 30, 993, 560, 15456, 11962, 261485,...: Higher dimension operators in the SM EFT

Henning, et al, arXiv:1512.03433

comprehensive review, arXiv:1610.0792

a strategy: SM Effective Field Theory

$$\mathcal{L}_{\text{eff}} = \mathcal{L}_{\text{SM}} + \Delta \mathcal{L}$$
$$= \mathcal{L}_{\text{SM}} + \sum_{i} \frac{c_i}{\Lambda^{d_i - 4}} O_i$$

the new particle searches at LHC suggest Λ >500 GeV

justify the analysis at dimension-6 operators

there are 84 of such operators for 1 fermion generation

assuming baryon number conservation, there are 59

luckily, there exists a smaller set relevant to physics at e+e-

SM Effective Field Theory: full formalism (23 pars.) ("Warsaw" basis)

$$\begin{split} \Delta \mathcal{L} &= \frac{c_H}{2v^2} \partial^{\mu} (\Phi^{\dagger} \Phi) \partial_{\mu} (\Phi^{\dagger} \Phi) + \frac{c_T}{2v^2} (\Phi^{\dagger} \overleftrightarrow{D}^{\mu} \Phi) (\Phi^{\dagger} \overleftrightarrow{D}_{\mu} \Phi) - \frac{c_6 \lambda}{v^2} (\Phi^{\dagger} \Phi)^3 \\ &+ \frac{g^2 c_{WW}}{m_W^2} \Phi^{\dagger} \Phi W^a_{\mu\nu} W^{a\mu\nu} + \frac{4gg' c_{WB}}{m_W^2} \Phi^{\dagger} t^a \Phi W^a_{\mu\nu} B^{\mu\nu} \\ &+ \frac{g'^2 c_{BB}}{m_W^2} \Phi^{\dagger} \Phi B_{\mu\nu} B^{\mu\nu} + \frac{g^3 c_{3W}}{m_W^2} \epsilon_{abc} W^a_{\mu\nu} W^{b\nu}{}_{\rho} W^{c\rho\mu} \\ &+ i \frac{c_{HL}}{v^2} (\Phi^{\dagger} \overleftrightarrow{D}^{\mu} \Phi) (\overline{L} \gamma_{\mu} L) + 4i \frac{c'_{HL}}{v^2} (\Phi^{\dagger} t^a \overleftrightarrow{D}^{\mu} \Phi) (\overline{L} \gamma_{\mu} t^a L) \\ &+ i \frac{c_{HE}}{v^2} (\Phi^{\dagger} \overleftrightarrow{D}^{\mu} \Phi) (\overline{e} \gamma_{\mu} e) \;. \end{split}$$

10 operators (h,W,Z, γ): CH, CT, C₆, CWW, CWB, CBB, C₃W, CHL, C'HL, CHE

- + 4 SM parameters: g, g', v, λ
- + 5 operators modifying h couplings to b, c, τ , μ , g
- + 2 parameters for h->invisible and exotic
- + 2 operators for contact interactions with quarks

strategy to determine all the 23 parameters simultaneously



(see papers the details of input observables and dependence on EFT operators)

recap 1: (synergy with LHC) two input observables from HL-LHC

$$\frac{\text{BR}(h \to \gamma \gamma)}{\text{BR}(h \to ZZ^*)} \qquad \frac{\text{BR}(h \to \gamma Z)}{\text{BR}(h \to ZZ^*)}$$

turn out to be very useful for constraining CWW, CWB, CBB

 $rac{g^2 c_{WW}}{m_W^2}$ $\Phi^{\dagger}\Phi W^{a}_{\mu
u}W$ $\frac{4gg'c_{WB}}{m_W^2}\Phi^{\dagger}t^a\Phi W^a_{\mu\nu}B^{\mu\nu}$ $\frac{g'^2 c_{BB}}{2} \Phi^{\dagger} \Phi B_{\mu\nu} B^{\mu\nu}$

recap 2: hWW is determined as precisely as hZZ @ $\sqrt{s} = 250$ GeV

 hWW/hZZ ratio can be determined to <0.1%: feature of a general SU(2) x U(1) gauge theory

$$\begin{split} \Gamma(h \to ZZ^*) &= (SM) \cdot (1 + 2\eta_Z - (0.50)\zeta_Z) \ ,\\ \Gamma(h \to WW^*) &= (SM) \cdot (1 + 2\eta_W - (0.78)\zeta_W) \\ \eta_W &= -\frac{1}{2}c_H \\ \eta_Z &= -\frac{1}{2}c_H - c_T \end{split} \qquad \text{custodial symmetry} \end{split}$$

SM-like hVV

$$C_{i} \sim O(10^{-4}-10^{-3})$$

 $\zeta_{Z} = c_{w}^{2}(8c_{WW}) + 2s_{w}^{2}(8c_{WB}) + (s_{w}^{4}/c_{w}^{2})(8c_{BB})$

anomalous hVV

recap 3: σ_{ZH} , σ_{ZHH} & beam polarizations

 $\frac{\partial^{\mu}}{\partial w^2} \partial^{\mu} (\Phi^{\dagger} \Phi) \partial_{\mu} (\Phi^{\dagger} \Phi)$

 $(\Phi^\dagger\Phi)^3$

- c_H has to be determined by inclusive σ_{Zh} measurement unique role of recoil mass analysis remains same
- c₆ has to be determined by double Higgs measurement

c6 is decoupled with single Higgs process (tree level), large deviation is allowed

 beam polarizations very powerful in EFT, separating hγZ and hZZ couplings: improved precisions and provide means to test EFT validity

typical precisions by EFT: combined EWPO+TGC+Higgs fit

ILC250: ∫Ldt = 2 ab⁻¹ @ 250 GeV

coupling $\Delta g/g$	kappa-fit	EFT-fit
hZZ	0.38%	0.68%
hWW	1.8%	0.67%
hbb	1.8%	1.1%
$\Gamma_{\rm h}$	3.9%	2.5%

(for hZZ and hWW couplings: 1/2 of partial width precision)

Higgs coupling precisions at ILC250



note the synergy: HL-LHC input is always included

benchmark BSM models

	Model	$b\overline{b}$	$c\overline{c}$	<i>gg</i>	WW	au au	ZZ	$\gamma\gamma$	$\mu\mu$
1	MSSM [34]	+4.8	-0.8	- 0.8	-0.2	+0.4	-0.5	+0.1	+0.3
2	Type II 2HD [36]	+10.1	-0.2	-0.2	0.0	+9.8	0.0	+0.1	+9.8
3	Type X 2HD [36]	-0.2	-0.2	-0.2	0.0	+7.8	0.0	0.0	+7.8
4	Type Y 2HD [36]	+10.1	-0.2	-0.2	0.0	-0.2	0.0	0.1	-0.2
5	Composite Higgs [38]	-6.4	-6.4	-6.4	-2.1	-6.4	-2.1	-2.1	-6.4
6	Little Higgs w. T-parity [39]	0.0	0.0	-6.1	-2.5	0.0	-2.5	-1.5	0.0
7	Little Higgs w. T-parity [40]	-7.8	-4.6	-3.5	-1.5	-7.8	-1.5	-1.0	-7.8
8	Higgs-Radion [41]	-1.5	- 1.5	10.	-1.5	-1.5	-1.5	-1.0	-1.5
9	Higgs Singlet [42]	-3.5	-3.5	-3.5	-3.5	-3.5	-3.5	-3.5	-3.5

Table 4: Deviations from the Standard Model predictions for the Higgs boson couplings, in %, for the set of new physics models described in the text. As in Table 1, the effective couplings g(hWW) and g(hZZ) are defined as proportional to the square roots of the corresponding partial widths.

--> quantitative assessment for models discrimination

model parameters (chosen as escaping direct search at HL-LHC)

- a PMSSM model with b squarks at 3.4 TeV, gluino at 4 TeV
- a Type II 2 Higgs doublet model with $m_A = 600 \text{ GeV}, \tan \beta = 7$
- a Type X 2 Higgs doublet model with $m_A = 450 \text{ GeV}, \tan \beta = 6$
- a Type Y 2 Higgs doublet model with $m_A = 600 \text{ GeV}, \tan \beta = 7$
- a composite Higgs model MCHM5 with $f = 1.2 \text{ TeV}, m_T = 1.7 \text{ TeV}$
- a Little Higgs model with T-parity with $f = 785 \text{ GeV}, m_T = 2 \text{ TeV}$
- A Little Higgs model with couplings to 1st and 2nd generation with $f=1.2 \text{ TeV}, m_T=1.7 \text{ TeV}$
- A Higgs-radion mixing model with $m_r = 500 \text{ GeV}$
- a model with a Higgs singlet at 2.8 TeVcreating a Higgs portal to dark matter and large λ for electroweak baryogenesis

BSM model discrimination at ILC250



once find deviation against SM —> can tell which BSM

then the next energy scales would be known: reachable 2~3 TeV

1.6 TeV

In general two Higgs doublet models (unitarity bound)						
[hZZ~0.38%]	1 TeV					
[hbb/hWW~0.64%] 2HDM Type II, Y	<u>3 TeV</u>					
[hττ/hWW~0.84%] 2HDM Type II, X	2.7 TeV					

In MSSM, NMSSM [hbb/hWW ~ 0.64%]

In the Higgs doublet and singlet model (unitarity bound) [hZZ~0.38%] <u>5 TeV</u>

In Minimal Composite Higgs Model (ξ=sin²(v/f)) [hZZ~0.38%] MCHM4 <u>2.8 TeV</u> [hbb/hWW~0.64%] MCHM5, MCHM10 <u>3.8 TeV</u>

Calculation done using the results from arXiv:1705.05399.

Endo, Kanemura, et al

determine Higgs CP (admixture)

• find CP-violating source in Higgs sector —> EW baryogenesis
• essential to understand structures of all Higgs couplings

through H—>T+T-
(or ttH)
$$L_{Hff} = -\frac{m_f}{v} H \bar{f} (\cos \Phi_{CP} + i\gamma^5 \sin \Phi_{CP}) f$$
Jeans et al, 1804.01241

through HZZ/HWW $L_{HVV} = 2C_V M_V^2 (\frac{1}{v} + \frac{a}{\Lambda}) HV_{\mu} V^{\mu} + C_V \frac{b}{\Lambda} HV_{\mu\nu} V^{\mu\nu} + C_V \frac{\tilde{b}}{\Lambda} HV_{\mu\nu} \tilde{V}_{\mu\nu}$ (CP-odd) $\Delta \tilde{b} \sim 0.016 \text{ (for } \Lambda = 1 \text{TeV} \text{)} \text{ Ogawa, 1712.09772}$

probing Hidden Sector: h->invisible / exotic decay



Higgs Physics beyond 250 GeV

the intrinsic advantage of a linear collider is its energy extendability

top mass measurement at ~350 GeV: vacuum stability

 λ runs < 0? top mass precision crucial for vacuum stability
 <p>
 at e+e-: top-pair threshold scan, much lower theory error
 Δm_t(MS-bar) ~ 50 MeV (Δm_H=14MeV)







Top-Yukawa coupling

- Iargest Yukawa coupling; crucial role in theory
- non-relativistic tt-bar bound state correction: enhancement by ~2 at 500 GeV
- Higgs CP measurement

$\Delta g_{ttH}/g_{ttH}$	500 GeV	550 GeV	1 TeV
H20	6.3%	2.7%	1.5%

Yonamine, et al., PRD84, 014033; Price, et al., Eur. Phys. J. C75 (2015) 309

Higgs self-coupling

$\Delta \lambda_{HHH} / \lambda_{HHH}$	500 GeV	+ 1 TeV
H20	27%	10%

Duerig, JT, et al, paper in preparation

Higgs self-coupling: when $\lambda_{\text{HHH}} \neq \lambda_{\text{SM}}$?

- ▷ constructive interference in ZHH, while destructive in vvHH (& LHC) —> complementarity between ILC & LHC, between √s ~500 GeV and >=1TeV
- If \u03c6 Interpretect in the second sec

Duerig, JT, et al, paper in preparation

references for large deviations

e.g.

Grojean, et al., PRD71, 036001; Kanemura, et al., 1508.03245; Kaori, Senaha, PHLTA, B747, 152; Perelstein, et al., JHEP 1407, 108

Status of ILC250

status of ILC250

- it's been a really long journey (my former two bosses at KEK & U.Tokyo both have been working on linear colliders for 30 years)
- scientific case is very strong and well justified -> Snowmass 2013, European Strategy 2013, AsiaHEP, and see most recently endorsement by ICFA, JAHEP
- technologically, entered from "matured stage" to "robust stage" -> extensive SCRF developments especially at European XFEL and LCSL-II
- under serious consideration by Japanese government -> expect some decision in this year, hopefully a green light

government

industry

#mylinearcollider

youtube

攻殻機動隊製作委員会. その他

鈴木 裕 Yu Suzuki

木村 優 Yu Kimura

backup

typical Higgs coupling deviation in BSM

Mixing with singlet

$$\frac{g_{hVV}}{g_{h_{\rm SM}VV}} = \frac{g_{hff}}{g_{h_{\rm SM}ff}} = \cos\theta \simeq 1 - \frac{\delta^2}{2}$$

Composite Higgs

$$\frac{g_{hVV}}{g_{h_{\rm SM}VV}} \simeq 1 - 3\% (1 \text{ TeV}/f)^2$$
$$\frac{g_{hff}}{g_{h_{\rm SM}ff}} \simeq \begin{cases} 1 - 3\% (1 \text{ TeV}/f)^2 & (\text{MCHM4})\\ 1 - 9\% (1 \text{ TeV}/f)^2 & (\text{MCHM5}) \end{cases}$$
SUSY

$$\frac{g_{hbb}}{g_{h_{\rm SM}bb}} = \frac{g_{h\tau\tau}}{g_{h_{\rm SM}\tau\tau}} \simeq 1 + 1.7\% \left(\frac{1 \text{ TeV}}{m_A}\right)^2$$

Higgs couplings in EFT

$$\begin{split} \Delta \mathcal{L}_{h} &= \frac{1}{2} \partial_{\mu} h \partial^{\mu} h - \frac{1}{2} m_{h}^{2} h^{2} - (1 + \eta_{h}) \overline{\lambda} v h^{3} + \frac{\theta_{h}}{v} h \partial_{\mu} h \partial^{\mu} h \\ &+ (1 + \eta_{W}) \frac{2m_{W}^{2}}{v} W_{\mu}^{+} W^{-\mu} h + (1 + \eta_{WW}) \frac{m_{W}^{2}}{v^{2}} W_{\mu}^{+} W^{-\mu} h^{2} \\ &+ (1 + \eta_{Z}) \frac{m_{Z}^{2}}{v} Z_{\mu} Z^{\mu} h + \frac{1}{2} (1 + \eta_{ZZ}) \frac{m_{Z}^{2}}{v^{2}} Z_{\mu} Z^{\mu} h^{2} \\ &+ \zeta_{W} \hat{W}_{\mu\nu}^{+} \hat{W}^{-\mu\nu} \left(\frac{h}{v} + \frac{1}{2} \frac{h^{2}}{v^{2}} \right) + \frac{1}{2} \zeta_{Z} \hat{Z}_{\mu\nu} \hat{Z}^{\mu\nu} \left(\frac{h}{v} + \frac{1}{2} \frac{h^{2}}{v^{2}} \right) \\ &+ \frac{1}{2} \zeta_{A} \hat{A}_{\mu\nu} \hat{A}^{\mu\nu} \left(\frac{h}{v} + \frac{1}{2} \frac{h^{2}}{v^{2}} \right) + \zeta_{AZ} \hat{A}_{\mu\nu} \hat{Z}^{\mu\nu} \left(\frac{h}{v} + \frac{1}{2} \frac{h^{2}}{v^{2}} \right) \,. \end{split}$$

$$\begin{split} \eta_h &= \delta \overline{\lambda} + \delta v - \frac{3}{2} c_H + c_6 & \theta_h = c_H \\ \eta_W &= 2 \delta m_W - \delta v - \frac{1}{2} c_H & \zeta_W = \delta Z_W \\ \eta_{WW} &= 2 \delta m_W - 2 \delta v - c_H & \zeta_Z = \delta Z_Z \\ \eta_Z &= 2 \delta m_Z - \delta v - \frac{1}{2} c_H - c_T & \zeta_A = \delta Z_A \\ \eta_{ZZ} &= 2 \delta m_Z - 2 \delta v - c_H - 5 c_T & \zeta_{AZ} = \delta Z_{AZ} \end{split}$$

comments on beam polarizations

	2 ab^{-1}	2 ab^{-1}	5 ab^{-1}	$+ 1.5 \text{ ab}^{-1}$	full ILC
	w. pol.	$350~{\rm GeV}$	no pol.	at 350 ${\rm GeV}$	$250{+}500~{\rm GeV}$
$g(hb\overline{b})$	1.04	1.08	0.98	0.66	0.55
$g(hc\overline{c})$	1.79	2.27	1.42	1.15	1.09
g(hgg)	1.60	1.65	1.31	0.99	0.89
g(hWW)	0.65	0.56	0.80	0.42	0.34
$g(h\tau\tau)$	1.16	1.35	1.06	0.75	0.71
g(hZZ)	0.66	0.57	0.80	0.42	0.34
$g(h\gamma\gamma)$	1.20	1.15	1.26	1.04	1.01
$g(h\mu\mu)$	5.53	5.71	5.10	4.87	4.95
g(hbb)/g(hWW)	0.82	0.90	0.58	0.51	0.43
g(hWW)/g(hZZ)	0.07	0.06	0.07	0.06	0.05
Γ_h	2.38	2.50	2.11	1.49	1.50
$\sigma(e^+e^- \to Zh)$	0.70	0.77	0.50	0.22	0.61
$BR(h \to inv)$	0.30	0.56	0.30	0.27	0.28
$BR(h \rightarrow other)$	1.50	1.63	1.09	0.94	1.15

comments on beam polarizations

	no pol.	80%/0%	80%/30%	
$g(hb\overline{b})$	1.33	1.13	1.04	
$g(hc\overline{c})$	2.09	1.97	1.79	
g(hgg)	1.90	1.77	1.60	
g(hWW)	0.98	0.68	0.65	
g(h au au)	1.45	1.27	1.16	
g(hZZ)	0.97	0.69	0.66	
$g(h\gamma\gamma)$	1.38	1.22	1.20	
$g(h\mu\mu)$	5.67	5.64	5.53	
$g(hb\overline{b})/g(hWW)$	0.91	0.91	0.82	
g(hWW)/g(hZZ)	0.07	0.07	0.07	
Γ_h	2.93	2.60	2.38	
$\sigma(e^+e^- \rightarrow Zh)$	0.78	0.78	0.70	
$BR(h \rightarrow inv)$	0.36	0.33	0.30	
$BR(h \rightarrow other)$	1.68	1.67	1.50	

Table 4: Projected relative errors for Higgs boson couplings and other Higgs observables with 2 ab⁻¹ of data at 250 GeV, comparing the cases of zero polarization, $80\% e^-$ polarization and zero positron polarization, and $80\% e^-$ polarization and 30% positron polarization. In each case, the running is equally divided into two samples with opposite beam polarization orientation.

discrimination between BSM models (ILC500 stage)

