



Realization of spontaneous gauge and supersymmetry breaking vacuum

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based on I705.00809
collaboration with T. Kobayashi, O. Seto, K. Ueda

Outline

- Introduction
- SUSY breaking associated with gauge symmetry breaking.
 - gauge $U(1)$ model*
 - Pati-Salam*
- Summary

Introduction

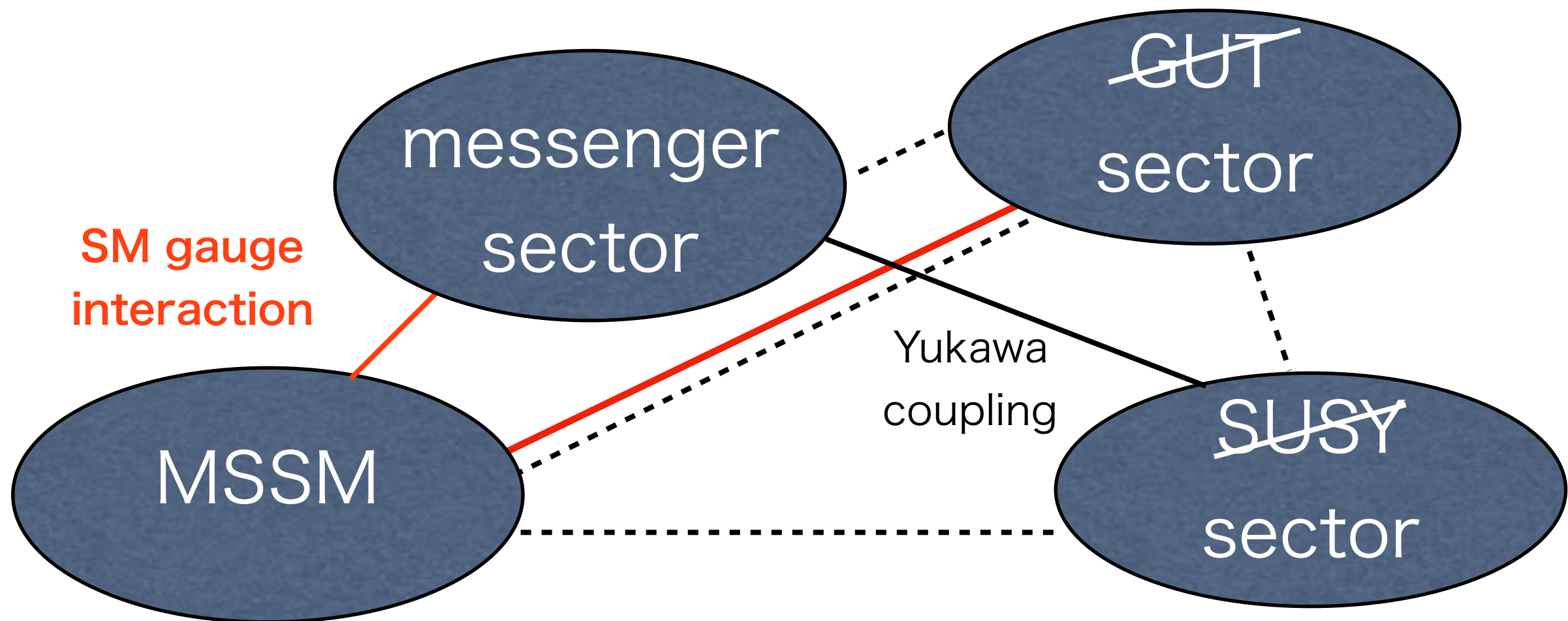
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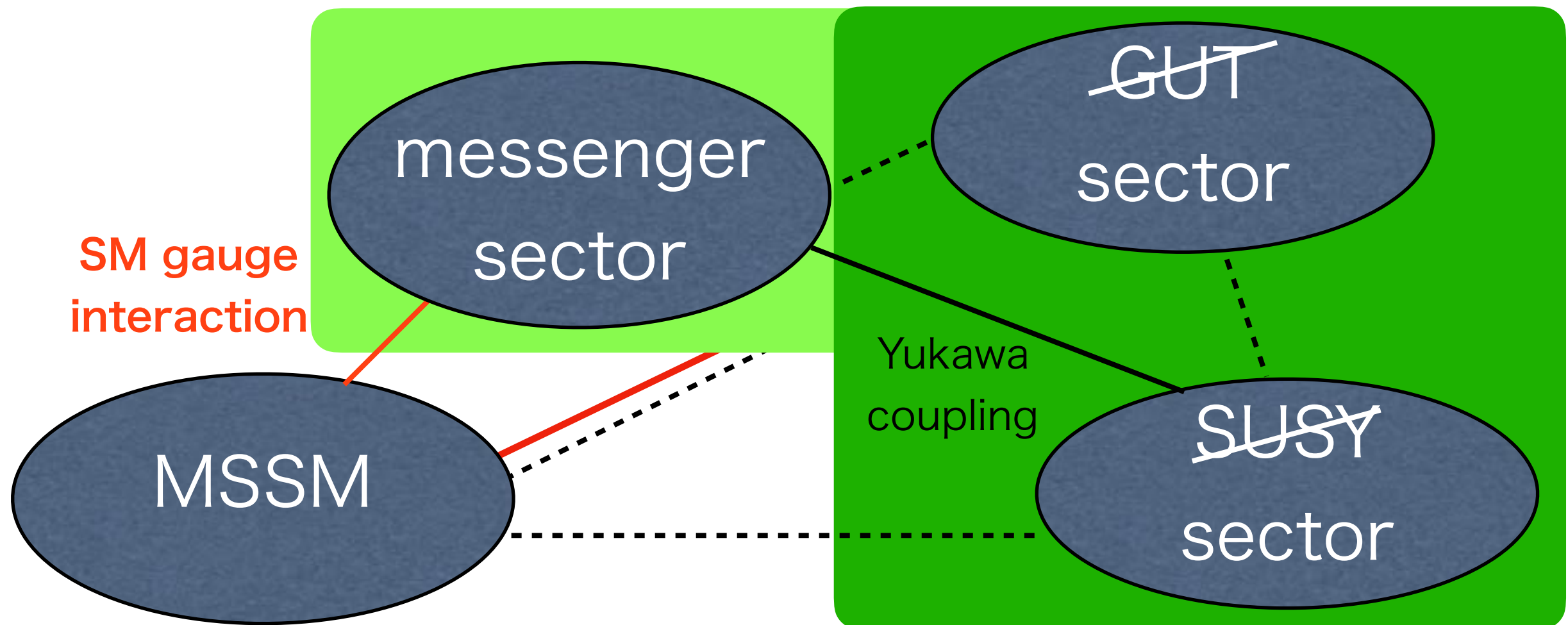


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If SUSY is conserved, $\partial_{\Phi_I} W = 0$ ($V = 0$)

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$$W = fX \quad \Rightarrow \quad \frac{\partial W}{\partial X} = f \neq 0 \quad \Rightarrow \quad V = f^2$$

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*Phase rotation of X is namely **R-symmetry**,
and it is crucial to SUSY breaking.*

X is flat

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Spontaneous SUSY breaking seems to imply the existence of R-symmetry in the models.

$$R_{charge}[W] = 2$$

(Nelson,Seiberg,[hep-th/9309299](#))

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Problem 2

*The scalar comp. of X may be stabilized by the radiative correction,
but not so simple. (Often stabilized at $X=0$.)*

(Shih,hep-th/0703196;Sun,0810.0477)

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One illustrative R-symmetric model

$$W = fX + \lambda X \phi_+ \phi_- + m_1 X_+ \phi_- + m_2 X_- \phi_+$$

$X = v_X + \theta^2 F$

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Motivated by those problems, I propose a model where ϕ_+ , ϕ_- , X_+ , X_- are charged under **extra gauge symmetry**.

Then, we find a realistic vacuum to cause SUSY breaking in association with gauge symmetry (GUT) breaking.

SUSY breaking
associated with gauge symmetry breaking

For the demonstration, let me consider the U(1) case:

(Kobayashi, Seto, YO, Ueda, I 705.00809)

Matter contents

Chiral superfield	SM	$U(1)$	$U(1)_R$
ϕ_{\pm}	singlet	± 1	0
X_{\pm}	singlet	± 1	2
X	singlet	0	2

superpotential

$$W = fX + \lambda X \phi_+ \phi_- + m_1 X_+ \phi_- + m_2 X_- \phi_+$$

U(1) D-term

$$V_D = \frac{g^2}{2} (|\phi_+|^2 - |\phi_-|^2 + |X_+|^2 - |X_-|^2)^2$$

Important point is that *there is no vacuum where all F-terms are vanishing in this model.*

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U(1) D-term

$$V_D = \frac{g^2}{2} (|\phi_+|^2 - |\phi_-|^2 + |X_+|^2 - |X_-|^2)^2$$

$$\partial_X W = f + \lambda \phi_+ \phi_-$$

$$\partial_{X_-} W = m_2 \phi_+$$

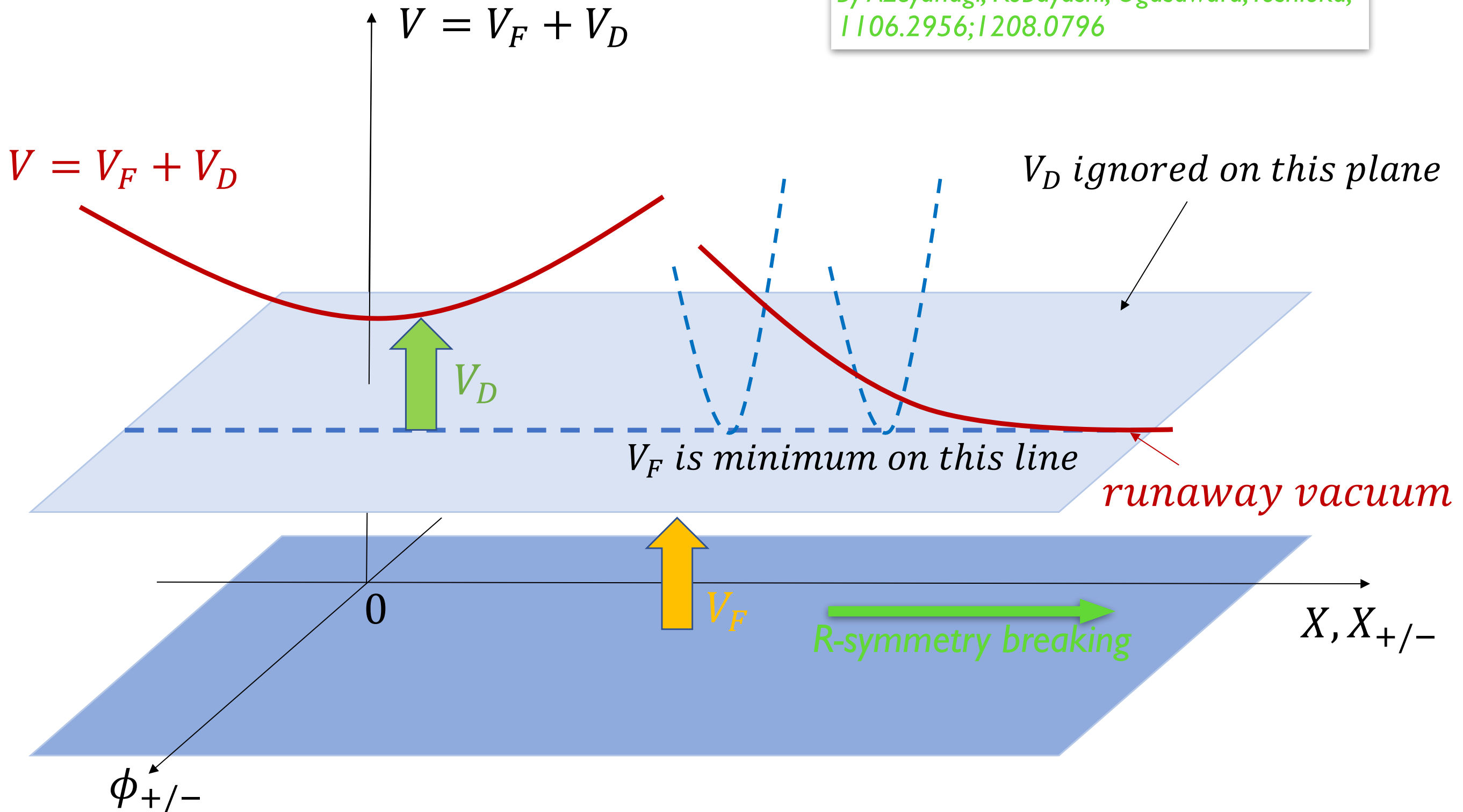
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sketch of the scalar potential

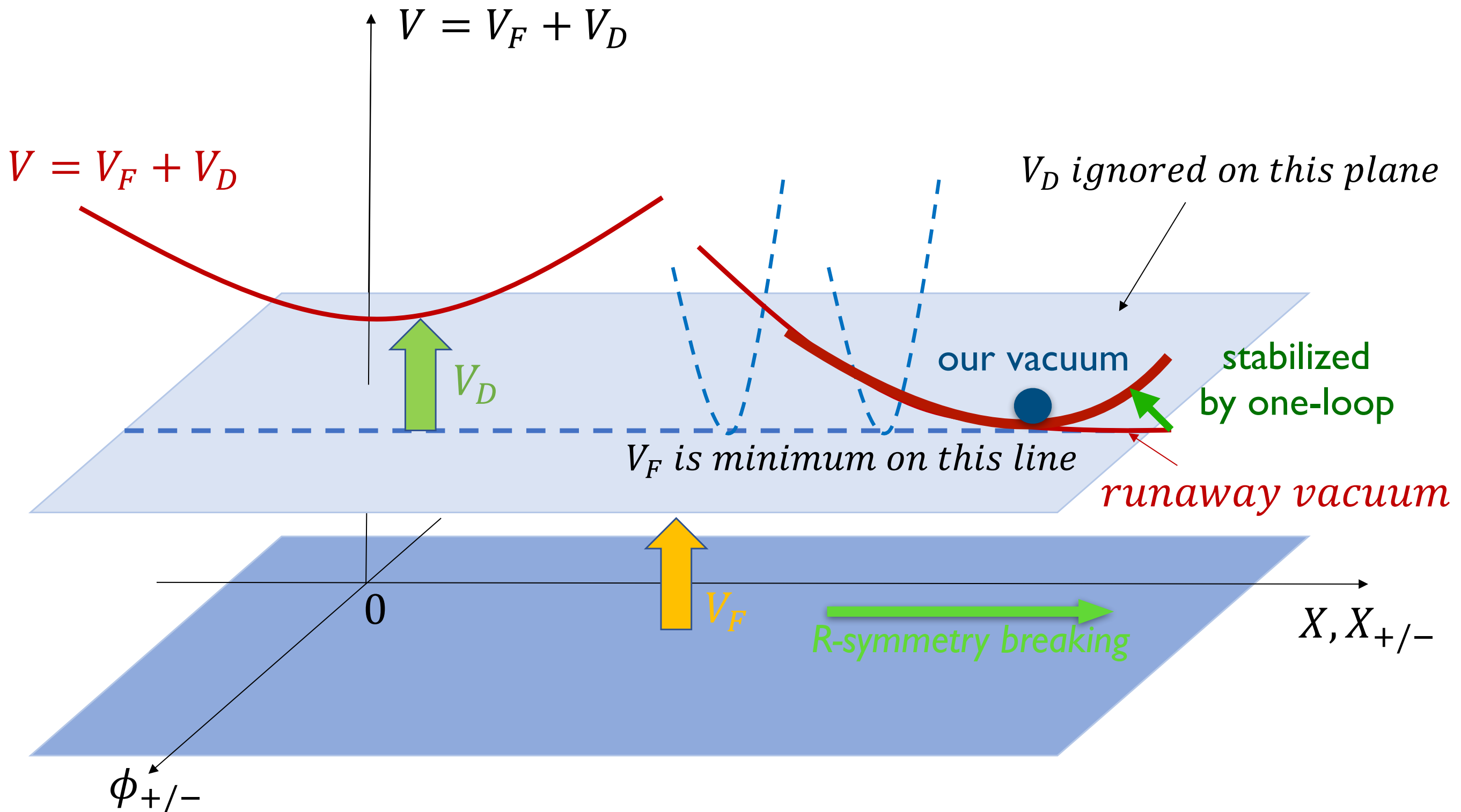
Note

This runaway vacuum is proposed
by Azeyanagi, Kobayashi, Ogasawara, Yoshioka,
1106.2956;1208.0796



sketch of the scalar potential

_Kobayashi,YO, Seto,Ueda,
1705.00809



Important points of this vacuum

- There is a runaway direction and stabilized by one-loop correction. Then, ***R-symmetry is also spontaneously broken.***

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- There is a runaway direction and stabilized by one-loop correction. Then, ***R-symmetry is also spontaneously broken.***
- The runaway behavior is one generic feature of this kind of model, where V_F is non-vanishing anywhere.

When $\langle \Phi_I \rangle = v_I$ satisfies $\partial_{\Phi_I} V_F = W_{IJ} \overline{W}^J = 0$

there is flat direction (z) in V_F : $\Phi_I = v_I + \overline{W}_I(v) z$

$$V_D = \frac{g_A^2}{2} \left(\Phi_I^\dagger q_{IJ}^A \Phi_L \right)^2 \xrightarrow{\hat{\Phi}_I = v_I + W_I(v)z + \frac{c_I}{z}} V_D|_{W_{IJ}\overline{W}^J=0} = \frac{g_A^2}{2} \left(v_I q_{IJ}^A c_J \frac{1}{z} + c.c. + \mathcal{O}(z^{-2}) \right)$$

because of the gauge symmetry

- We have discussed a simple U(1) model:

$$U(1) \rightarrow \textit{nothing}$$

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- We can extend this to GUT symmetry:

If the fields, $\phi_{+/-}$ and $X_{+/-}$, are charged under GUT symmetry (e.g., Pati-Salam, flipped $SU(5)$ etc.),

$$SU(2)_R \times SU(2)_L \times SU(4) \rightarrow SM$$

$$SU(5) \times U(1) \rightarrow SM$$

GUT breaking and SUSY breaking are unified.

Let's consider the application to "Pati-Salam" ($SU(4) \times SU(2)_R \times SU(2)_L$).

(Kobayashi, YO, Seto, Ueda, [I 705.00809](#))

The charge assignment of $\phi_{+/-}$ and $X_{+/-}$

$$X_+, \phi_+ : (4, 2, 1), \quad X_-, \phi_- : (\bar{4}, 2, 1)$$

Superpotential

$$W = fX + \lambda X \text{Tr}(\phi_+ \phi_-) + m_1 \text{Tr}(X_+ \phi_-) + m_2 \text{Tr}(X_- \phi_+)$$

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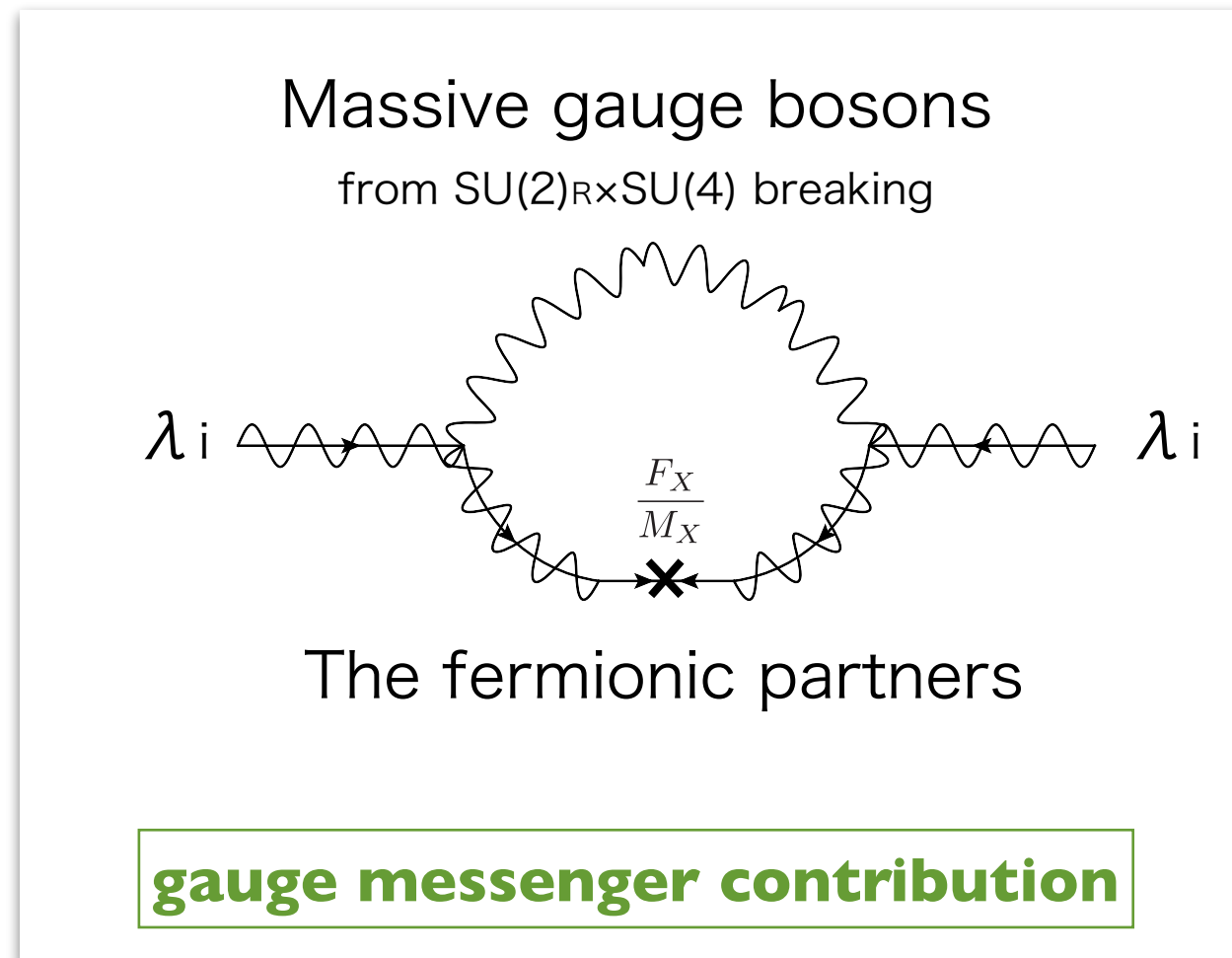
$$W = fX + \lambda X \text{Tr}(\phi_+ \phi_-) + m_1 \text{Tr}(X_+ \phi_-) + m_2 \text{Tr}(X_- \phi_+)$$

We can find the runaway direction that breaks R-symmetry, and we obtain the SUSY breaking vacuum at the one-loop.

$$\langle \phi_{\pm} \rangle = \begin{pmatrix} \mp \frac{F}{m_{1,2}} & 0 \\ 0 & 0 \\ 0 & 0 \\ 0 & 0 \end{pmatrix} \quad \langle X_{\pm} \rangle = \begin{pmatrix} \pm z_0 & 0 \\ 0 & 0 \\ 0 & 0 \\ 0 & 0 \end{pmatrix} \quad \langle F_{X_{\pm}} \rangle = \begin{pmatrix} \pm F & 0 \\ 0 & 0 \\ 0 & 0 \\ 0 & 0 \end{pmatrix}$$

$$SU(2)_R \times SU(4) \rightarrow SU(3) \times U(1)_Y$$

- Another interesting point of this setup is *gaugino mass is generated by the gauge messenger contributions.*



$$M_a = \frac{\alpha_a}{4\pi} \Delta b_a \frac{F}{z_0}$$

where $(\Delta b_Y \Delta b_2, \Delta b_3) = (-10/3, 0, -1)$

Summary

- How to break SUSY spontaneously is one of the crucial issues.
- SUSY breaking model generally has R-symmetry.
 - *Both R-symmetry and SUSY should break down at our vacuum.*
- I introduce one type of SUSY breaking models:

Gauge symmetry causes a runaway direction to break R-symmetry. Our vacuum is induced by stabilizing the runaway according to the one-loop correction.
- I gave some comments on the applications to GUT models in our paper. The detail will be shown near future.

Backup

Comments (Kobayashi,YO, Seto,Ueda,1705.00809)

- *Wino mass is vanishing at LO in the Pati-Salam. Other contributions like gravity mediation may be required.*
- *This mechanism can be applied to other GUTs, e.g., flipped $SU(5)$.*
- *One issue of this mechanism is about tachyonic masses of sfermions.*

one-loop contribution: almost vanishing at our vacuum.

two-loop contribution: we need more careful study.

- *Spontaneous R -symmetry breaking leads massless R -axion.*

Gauged $U(1)_A$ instead of R -symmetry does work well.

N. Maekawa, Y. Shigekami, M. Yoshida, 1712.05107

Gauged $U(1)_A$ Model

(Maekawa,YO,Shigekami,Yoshida,1712.05107)

Matter contents

Chiral superfield	SM	$U(1)_A$
S	singlet	$s \gg 1$
Θ	singlet	-1

Superpotential

$$W = a_0 + a_1 S \left(\frac{\Theta}{\Lambda} \right)^s + \frac{a_2}{2} \left(S \left(\frac{\Theta}{\Lambda} \right)^s \right)^2 + \frac{a_3}{6} \left(S \left(\frac{\Theta}{\Lambda} \right)^s \right)^3 + \dots$$



$W_{eff} \approx a_0 + (a_1 \lambda^s) S$ is obtained when $\lambda \ll 1$.

$$\left(\lambda \equiv \left\langle \frac{\Theta}{\Lambda} \right\rangle \right)$$

$U(1)_A$ D-term potential

$$V_D = \frac{g^2}{2} \left(\xi^2 - |\Theta|^2 + s|S|^2 \right)^2$$

Effective Potential for S in **SUGRA**

$$V_{eff} = f^2 + |W_{\Theta S}| |S|^2 + g^2 s D_0 |S|^2 + \frac{1}{2} g^2 s^2 |S|^4$$

$$-3 \frac{a_0^2}{M_p^2} + \boxed{(S + S^\dagger) \frac{a_0 f}{M_p^2}} + \frac{a_0^2}{M_p^2} |S|^2 + \dots$$

(originated from constant term, a_0 , in W)

$$V = V_F + V_D$$

SUGRA

(Maekawa, YO, Shigekami, Yoshida, I712.05107)

SUSY is broken.

R symmetry is broken by SUGRA effect

