Dark Matter early kinetic decoupling and finite temperature corrections to long-range force systems

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based on [1706.07433] in collaboration with T.Bringmann, M.Gustafsson, A.Hryczuk,

[1808.06472] in collaboration with L. Covi, K. Mukaida.

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Standard description of DM chemical decoupling process:

$$\dot{n} + 3Hn = -\langle \sigma v \rangle_{T_{SM}} \left[n^2 - n_{eq}^2 (T_{SM}) \right]$$

Lee&Weinberg 77', Gondolo&Gelmini 91'

 implemented in DarkSUSY and MicrOMEGAs, efficiently and accurately solving for DM relic density.

Captures special cases like co-annihilation, threshold and resonances.

- relies on certain assumptions and has limitations.
- developed refined description accounting for i) early kinetic decoupling ii) self-consistent treatment of thermal corrections.

Planck precision era: $\Omega_{DM} h^2 = 0.1198 \pm 0.0015$ (!), refinement of DM overlosure bound.

Early kinetic decoupling of dark matter: when the standard way of calculating the thermal relic density fails $_{\rm Binder+\ 17'}$

- ► Early kd if $\Gamma_{sc}|_{cd} \sim \Gamma_{an}|_{cd}$. Not only a theoretical possibility!
- ► DM annihilation ↔ temperature
- developed two methods to accurately compute the DM relic abundance for such case

Two methods to deal with early kd

Full phase-space density solution (full BE)

- ► Numerical solution of the non-linear integro PDE: $E(\partial_t - 2Hp\partial_p)f(t,p) = C_{an}[f, f_{eq}] + C_{sc}[f, f_{SM}]$
- accounting for phase-space density distortion caused by velocity-dependent annihilation.

Solving coupled system of momentum moments (cBE)

► Coupled system of ordinary differential equations: $\dot{n} + 3Hn = g \int \frac{d^3p}{(2\pi)^3 E} C_{an}[f, f_{eq}]$ $\dot{T} + \frac{\dot{n}}{n}T + 2HTw(T) = \frac{g}{3n} \int \frac{d^3p}{(2\pi)^3 E} \frac{\mathbf{p}^2}{E} \left(C_{an}[f, f_{eq}] + C_{sc}[f, f_{SM}] \right)$

► Closer of Boltzmann hierarchy by form Ansatz $f = \frac{n}{n_{eq}(T)}e^{-\frac{E}{T}}$.

Simplest example: Scalar singlet dark matter



- Two parameter model: m_S , λ_S .
- Breit-Wigner enhanced annihilation for $m_S \sim m_H/2$.
- Higgs couples preferably to heavy SM particles
 - → Boltzmann suppressed scattering



- Already in simplest scalar singlet DM model, the relic abundance can be affected by more than one order of magnitude!
- Public phase space density solver soon available

T.Binder, T.Bringmann, M.Gustafsson, A.Hryczuk in prep. 2018, provide many more examples: Higgsand Z-portals, forbidden DM, Sommerfeld enhanced annihilation,...

Dark matter Sommerfeld-enhanced annihilation and Bound-state decay at finite temperature Binder+ 18'



- DM (χ) longe-range force enhanced annihilation in the presence of hot and dense plasma background.
- Vacuum computation well understood.

 $\mathscr{L} \supset g_{\gamma} \bar{\chi} \gamma^{\mu} \chi A_{\mu} + g_{\psi} \bar{\psi} \gamma^{\mu} \psi A_{\mu}$

- Finite temp. effects: at first place a conceptional question.
 Charge screening? Dissociation ? No (correct) formal description available, claims about strong finite temperature corrections based on linear response theory
- comprehensive ab initio derivation from non-eq. QFT mostly new research in quantum statistical mechanics
- \blacktriangleright main result: new number density equation \checkmark

Master equation for s-wave annihilation:

$$\begin{split} \dot{h}_{\eta} + 3Hn_{\eta} &= -2(\sigma v)_{0} G_{\eta\xi}^{++--}(x, x, x, x)|_{eq} \Big[e^{2\beta \mu} - 1 \Big], \\ G_{\eta\xi}^{++--}|_{eq} &= e^{-2\beta M} \int_{-\infty}^{\infty} \frac{\mathrm{d}^{3} \mathbf{P}}{(2\pi)^{3}} e^{-\beta \mathbf{P}^{2}/4M} \int_{-\infty}^{\infty} \frac{\mathrm{d} E}{(2\pi)} e^{-\beta E} G_{\eta\xi}^{\rho}(\mathbf{0}, \mathbf{0}; E)|_{I=0}. \end{split}$$

- General equation, accounting for non-ideal and finite temperature corrections.
- Chemical potential $\mu[n_{\eta}]$ sets functional dependence of r.h.s on n_{η}
- Four point correlator G⁺⁺⁻⁻_{ηξ} (x, x, x, x)|_{eq} in chemical equilibrium includes Sommerfeld enhanced annihilation and bound state decay simultaneously via two-particle spectral function G^ρ_{ηξ}. Form invariant under interpretation of bound or scattering state at finite temperature.
- ▶ Self-consistent solution of $\mu[n_{\eta}]$ and $G_{\eta\xi}^{++-}(x,x,x,x)|_{eq}$ is important!
- Checked various limits and compared to existing literature: vacuum limit √ (Lee-Weinberg eq.), linear regime close to chemical equilibrium √ (linear response theory)

Vacuum spectral function

Vacuum relations between s-wave two-particle spectral function $G_{\eta\xi}^{\rho}$. Sommerfeld enhancement factor *S*, and bound-state decay width Γ :

$$(\sigma v_{\rm rel})_0 G^{\rho}_{\eta\xi}(\mathbf{0},\mathbf{0};E)|_{E>0,l=0} = \frac{1}{2\pi} {\rm Tr}[\mathbf{1}_{2\times 2}] M(ME)^{1/2} (\sigma v_{\rm rel})_0 S(E),$$

$$(\sigma v_{\rm rel})_0 G^{\rho}_{\eta\xi}(\mathbf{0},\mathbf{0};E)|_{E<0,l=0} = \frac{\pi}{2} {\rm Tr}[\mathbf{1}_{2\times 2}] \sum_n \delta(E-E_{B_n}) \Gamma_n.$$



Spectral function at finite temperature

Effective in-medium potential (in HTL resummed effective theory):

$$V_{\rm eff}(\mathbf{r}) = -\alpha_{\chi} m_D - \frac{\alpha_{\chi}}{r} e^{-m_D r} - i\alpha_{\chi} T \phi(m_D r),$$

where $\phi(0) = 0$ and $\phi(\infty) = 1$, and Debye $m_D^2 = g_\psi^2 T^2/3$.



Notation of spectral function more general than using S or Γ . Shift towards lower energies affects total rate exponentially.

Summary, Conclusion, Outlook

- Questioning the assumptions entering the standard description of the DM freeze-out can lead to surprising new phenomena.
- Early kinetic decoupling can have huge impact on DM overclosure bound.
- Public phase-space density code with many more examples soon available.
- Developed rigorous theoretical understanding of long-range force systems at finite temperature.
- ► If bound state solution exist $(m_V \leq \alpha_{\chi} M)$ they are important to include in the relic density computation.
- New number density equation allows to include finite temperature effects self-consistently.
- Impact of finite temperature effects on relic density needs careful investigation, work in progress.