

Supported by the US Department of Energy

- 1. Motivations
- 2. Introduction to einselection and the toy model
- 3. Einselection in equilibrium (technical explorations and overall assessment)
- 4. Eigenstate Einselection Hypothesis (if there is time)

5. Conclusions

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With A. Arrasmith

5/21/19

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Comments on the arrow of time in cosmology (at board)

Q: What features of the universe are correlated with classicality?

- Arrow of time?
- Locality?
- Etc.

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→ Explore the process of einselection in a toy model, relate to AoT, etc.



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Related to the emergence of

classical

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Traditionally connected with arrow of time

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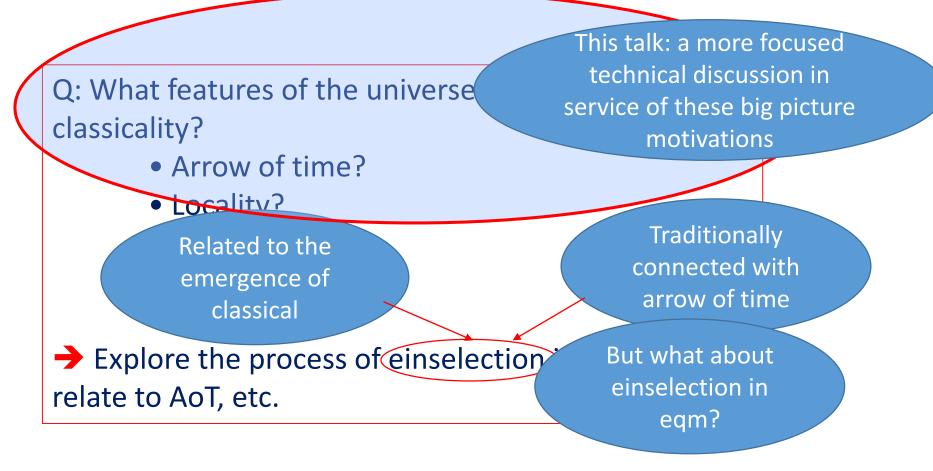
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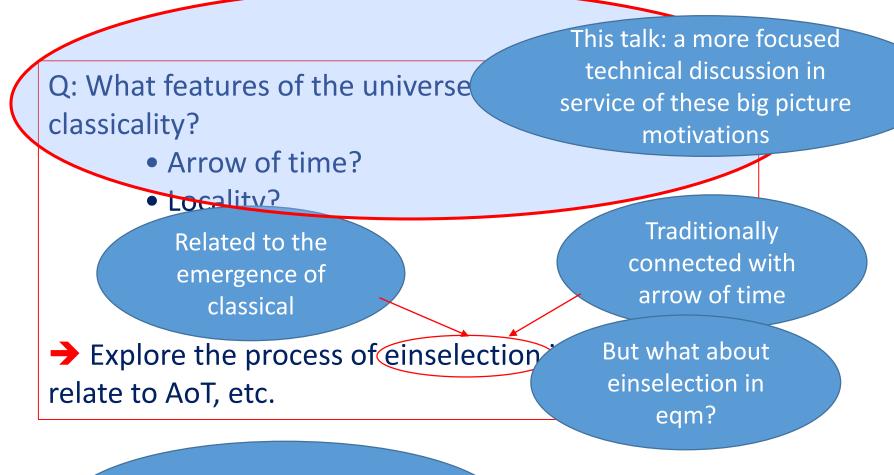
- Arrow of time?
- Locality?

Related to the emergence of classical

Explore the process of einselection relate to AoT, etc. Traditionally connected with arrow of time

But what about einselection in eqm?





If you are handed a theory, what are the classical degrees of freedom?

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The preference of special "pointer" states of a system due to interactions with the environment

- Stability of pointer states
- Destruction of non-pointer states (including "Schrödinger cat" superpositions of pointer states)
- Pure non-pointer states → mixtures of pointer states via entanglement with the environment.

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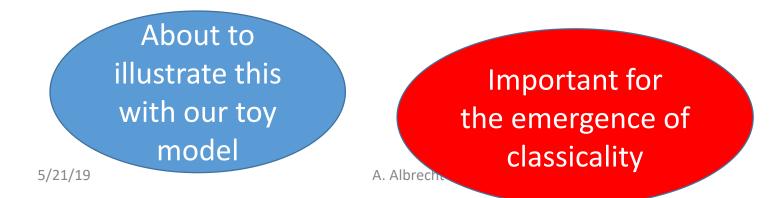
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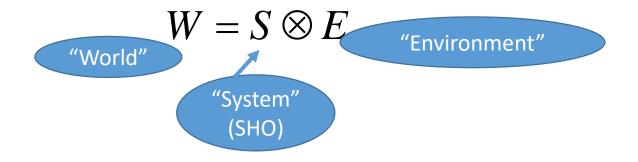
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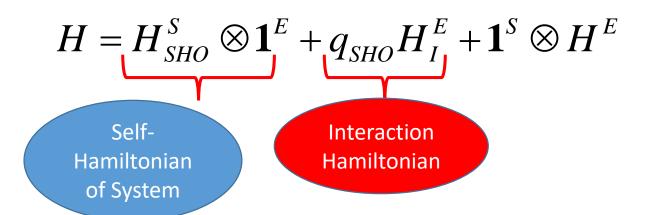


$$W = S \otimes E$$

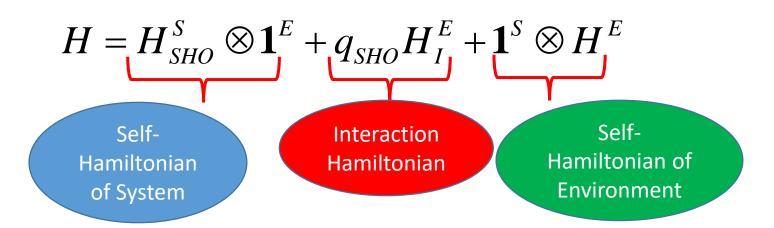
$$H = H_{SHO}^{S} \otimes \mathbf{1}^{E} + q_{SHO} H_{I}^{E} + \mathbf{1}^{S} \otimes H^{E}$$

Self-
Hamiltonian
of System

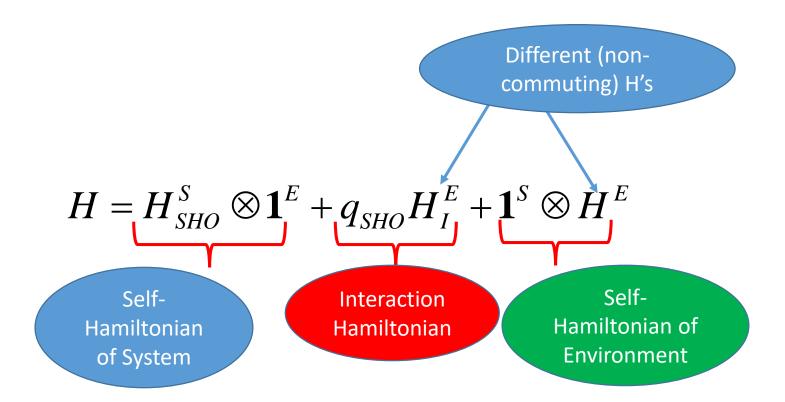
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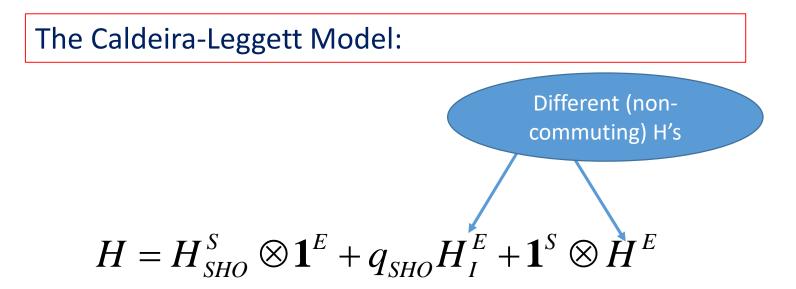


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<u>CL:</u>

- i) Model *E* as an infinite set of SHOs with different frequencies
- ii) Take special (order of) limits and parameter choices to get an <u>(irreversible) stochastic equation</u> that describes this (unitary) evolution under certain conditions (including AoT)

iii) Demonstrate einselection etc. (CL and others)

The Caldeira-Leggett Model:Different (non-commuting) H's
$$H = H_{SHO}^S \otimes \mathbf{1}^E + q_{SHO} H_{I_n}^E + \mathbf{1}^S \otimes H^E$$
Adapted CL:Random Hermitian
matricesi) Model E as finite systemii) Solve full unitary evolution in all regimes (numerical)

- iii) Demonstrate Einselection under certain conditions (AoT)
- iv) Explore scope of einselection (eqm?)

Introducing the toy model

- No interaction case $(H_I^E = 0)$
- Model SHO with d=30 Hilbert space

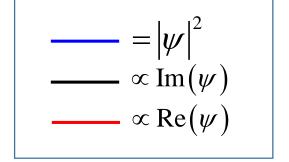
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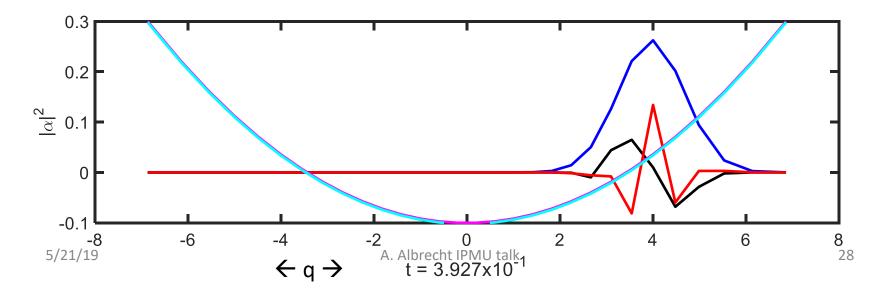
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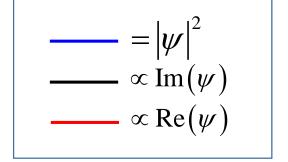
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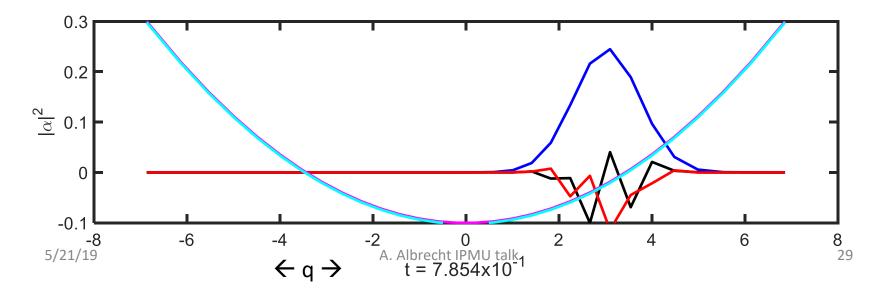
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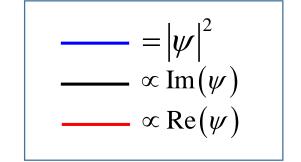
"Movie" A (Isolated SHO)

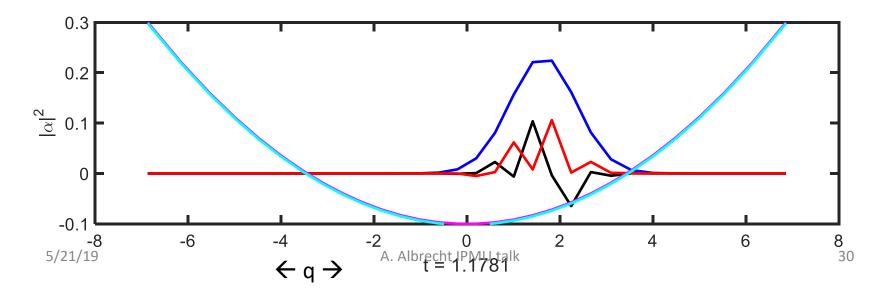


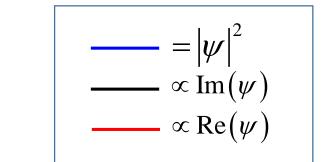


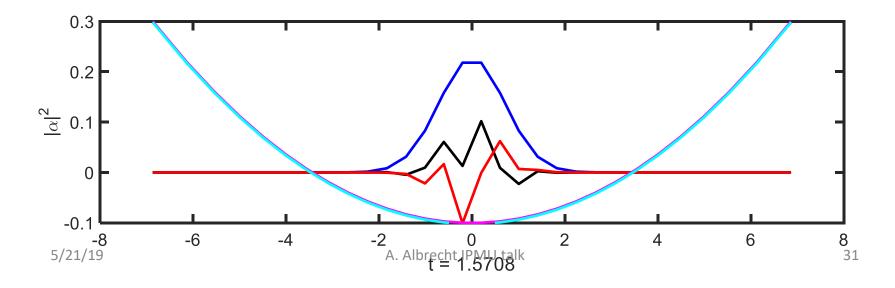


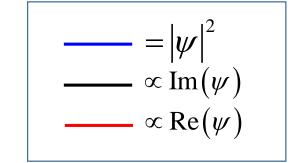


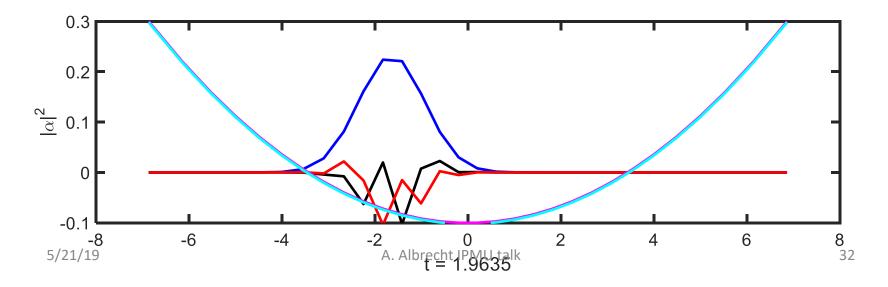


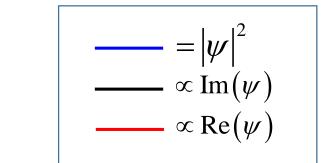


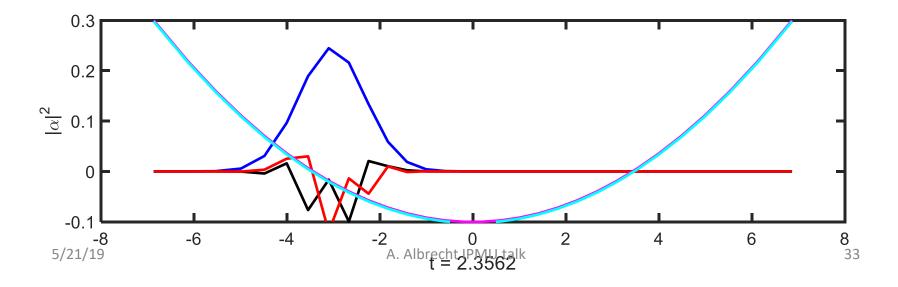


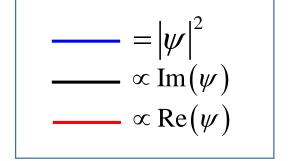


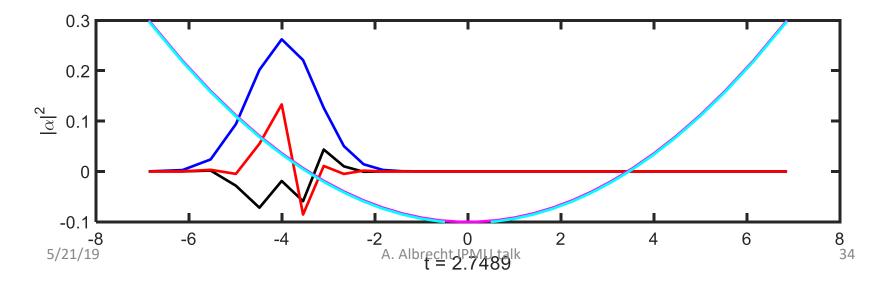


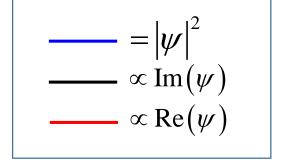


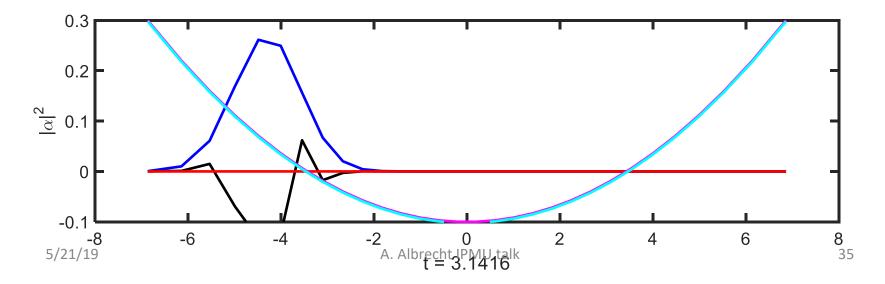


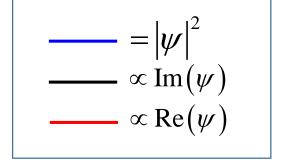


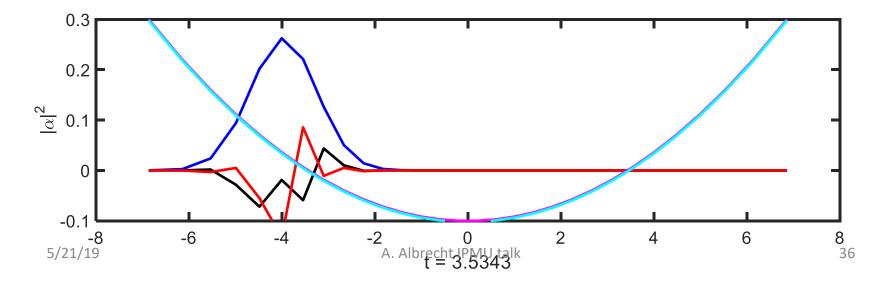


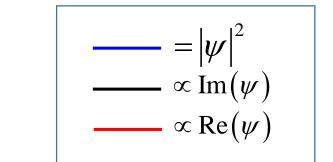


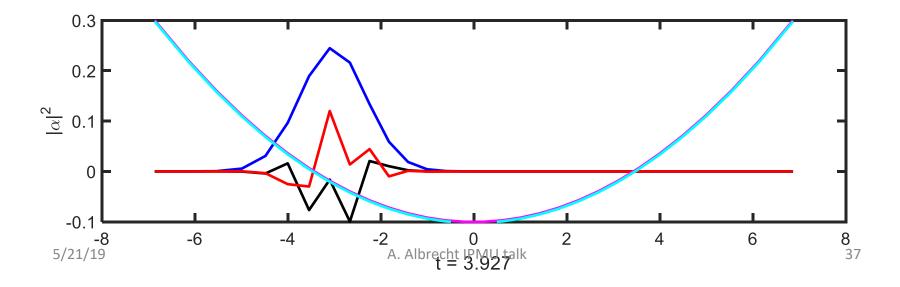


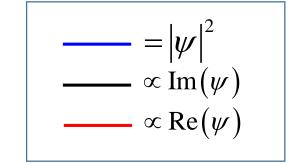


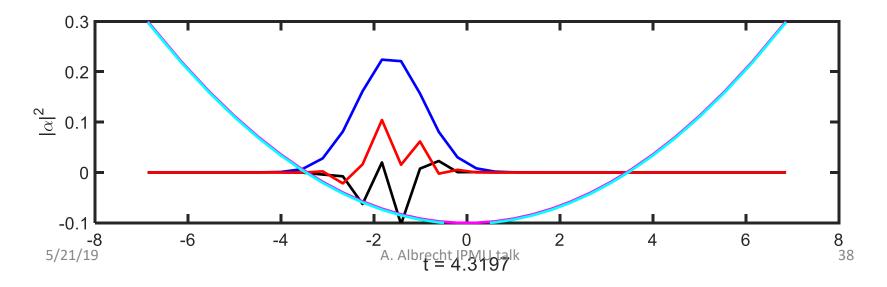


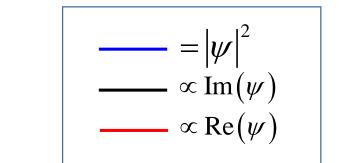


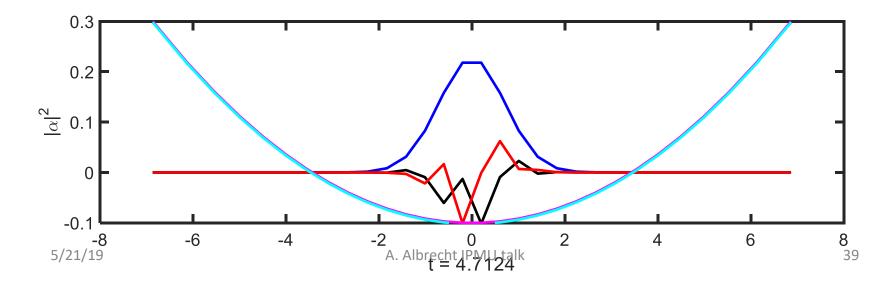


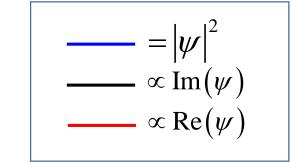


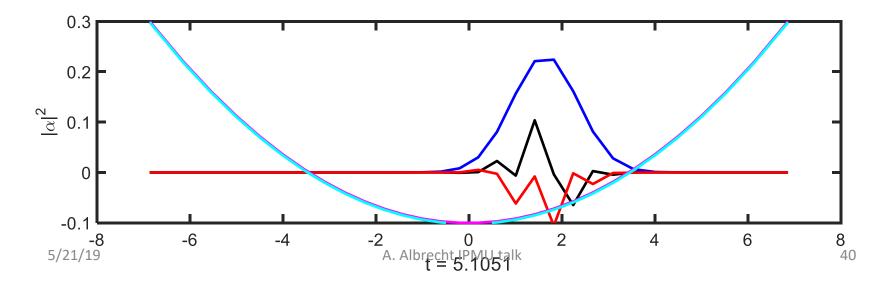


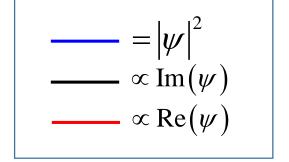


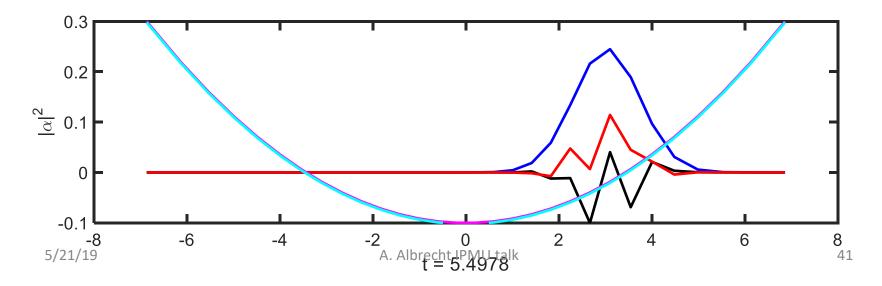


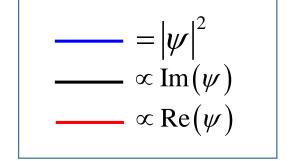


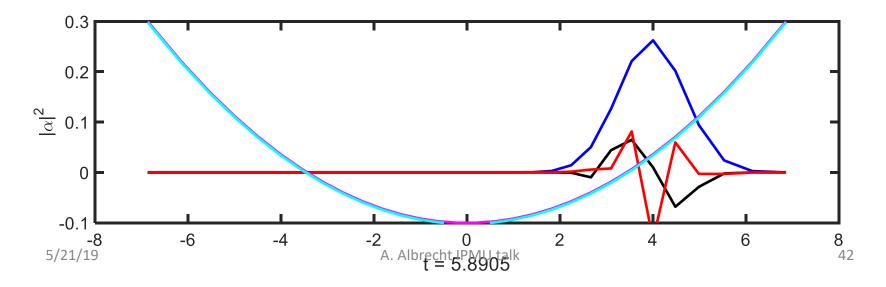


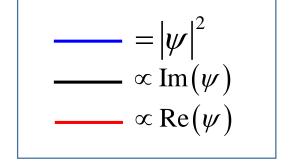


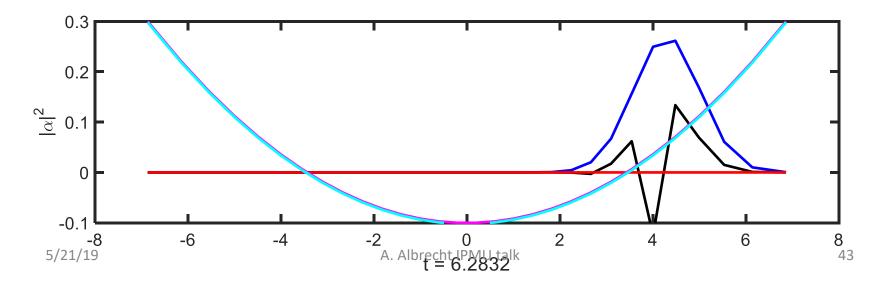






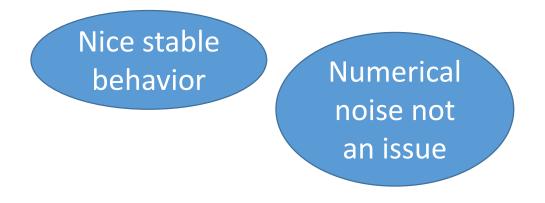






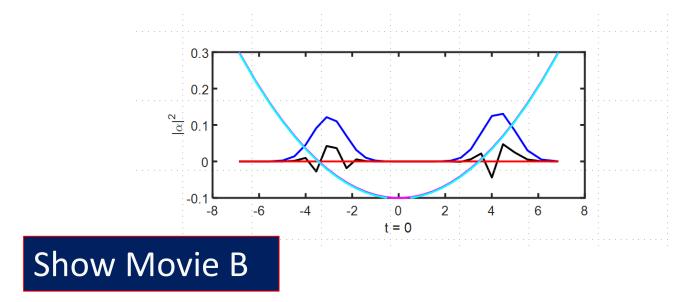
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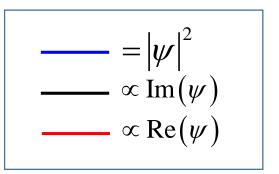


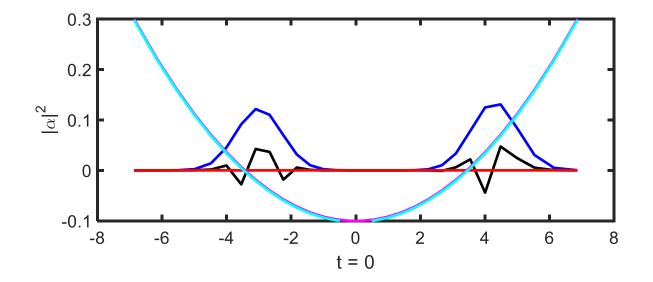
Introducing the <u>Schrödinger cat</u>

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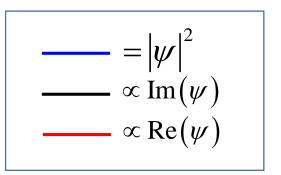


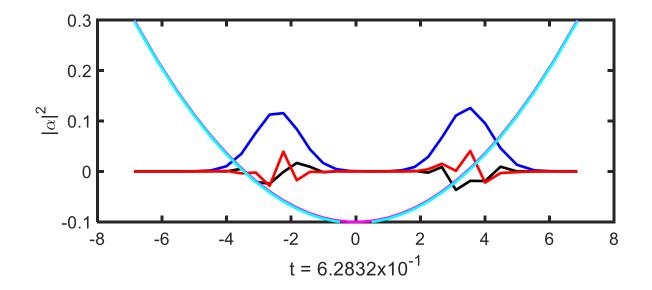




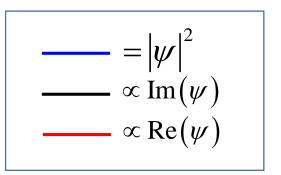


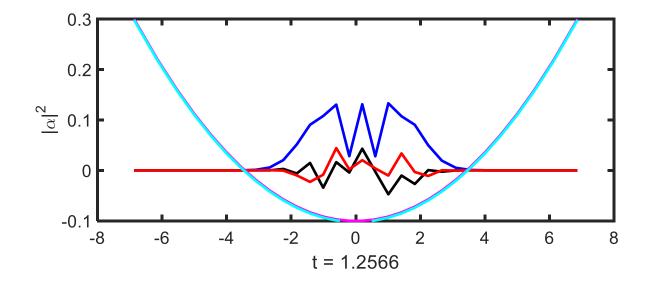




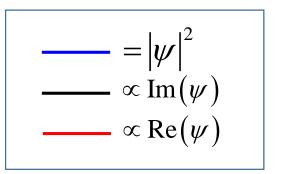


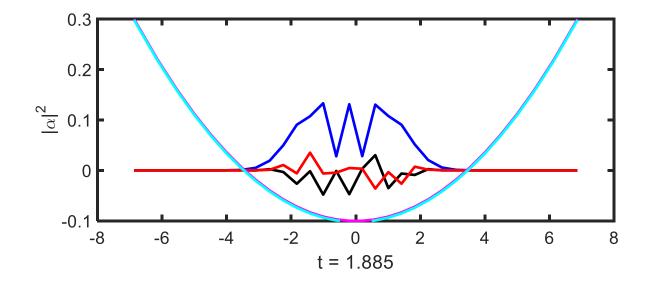




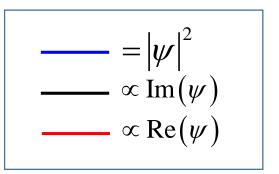


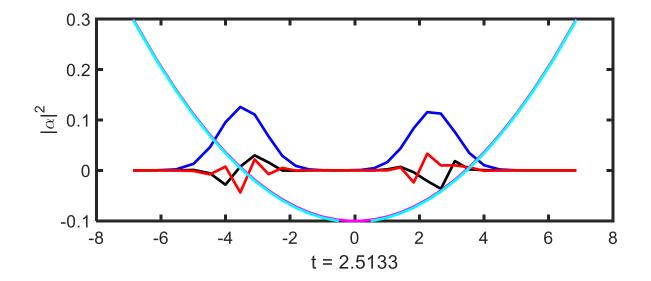




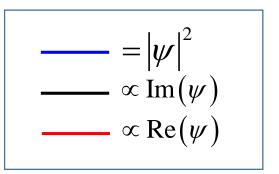


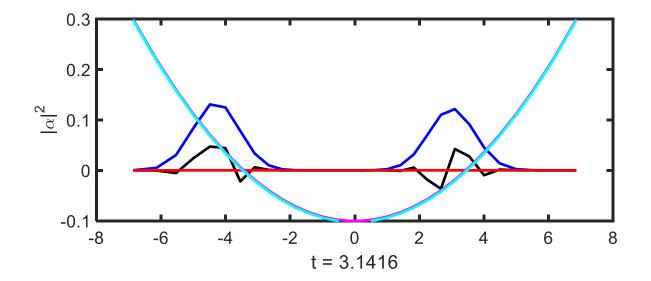




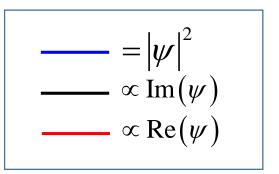


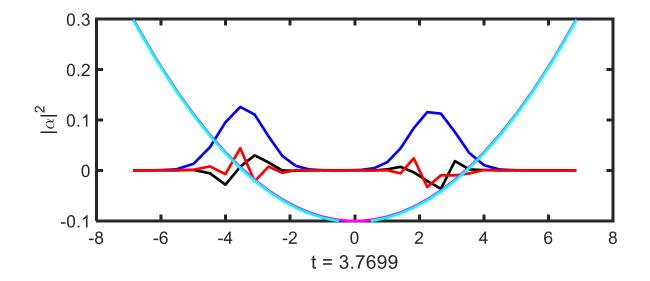




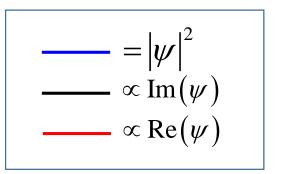


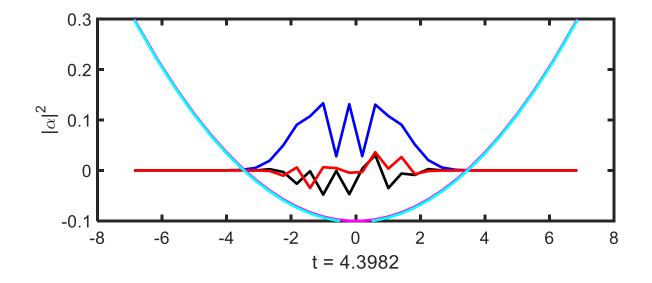




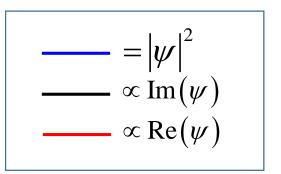


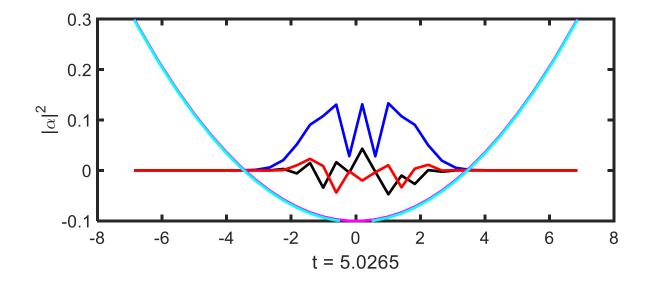




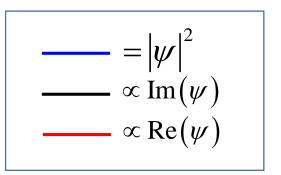


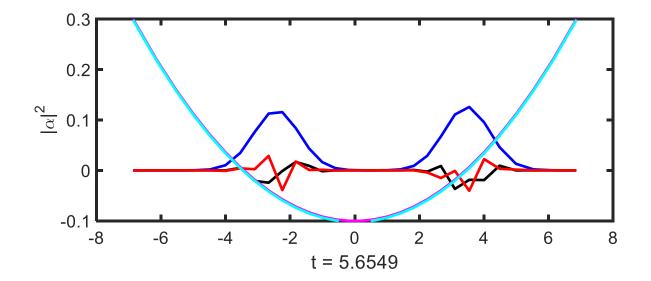




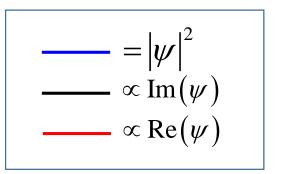


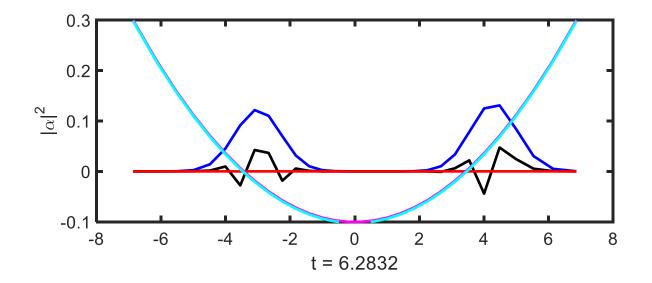












• Interactions turned on $(H_I^E \neq \mathbf{0})$

$$H = H_{SHO}^{S} \otimes \mathbf{1}^{E} + q_{SHO}H_{I}^{E} + \mathbf{1}^{S} \otimes H^{E}$$

This evolution will take an initial product state into a mixed state:

$$|\psi\rangle_{W} = |\psi\rangle_{S} |\psi\rangle_{E} \rightarrow \sum_{i,j} \alpha_{ij} |i\rangle_{S} |j\rangle_{E}$$

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m

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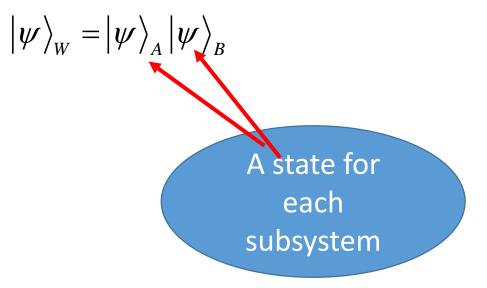
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$$\rho_{S} \equiv Tr_{E}(|\psi\rangle_{WW}\langle\psi|) = |\psi\rangle_{SS}\langle\psi| \Longrightarrow \text{more general } \rho_{S}$$

Some comments on entangled states

$$W = A \otimes B$$

Inclined to think:



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But the general case is:

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Which gives:

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A density
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Eigenvalues and eigenstates of ρ_A

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Eigenvalues and eigenstates of ρ_A

("Schmidt states")

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Eigenvalues and eigenstates of ρ_A

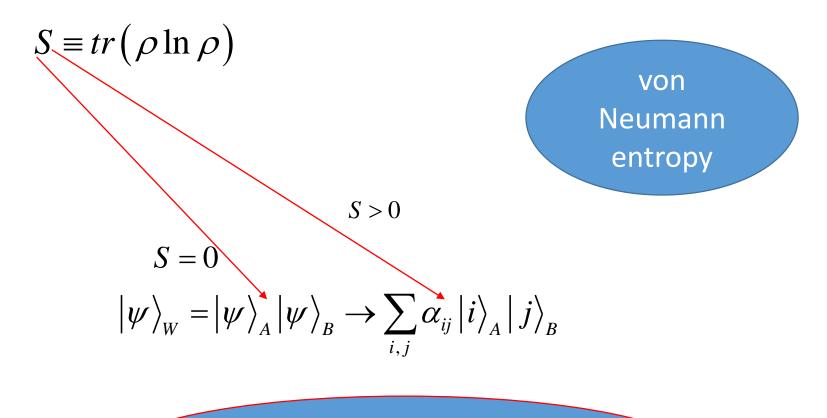
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Eigenvalues and eigenstates of ρ_A

A "classical mixture" of the eigenstates (no cross terms → no "quantum coherence")

Onset of entanglement → decoherence



Decoherence → increasing S → arrow of time In general, the density matrix eigenstates might vary significantly from one moment to the next, producing a "classical mixture of random stuff"

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- In general, the density matrix eigenstates might vary significantly from one moment to the next, producing a "classical mixture of random stuff"
- Special case: The nature of the interactions between A and B lead to a reliable preference for specific "pointer" eigenstates. This is <u>Einselection</u>

Discuss pendulum interacting with the air, leading to localized wave packed pointer states.

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An illustration of Einselection

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Show eigenstates and eigenvalues

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 $|\psi\rangle_{W} = |\psi\rangle_{S} |\psi\rangle_{E} \rightarrow \sum_{i,j} \alpha_{ij} |i\rangle_{S} |j\rangle_{E}$

Show First 2 eigenstates and eigenvalues

$$\rho_{S} \equiv Tr_{E} \left(|\psi\rangle_{WW} \langle \psi| \right) = |\psi\rangle_{SS} \langle \psi| \rightarrow \text{more general } \rho_{S}$$

• Interactions turned on $(H_I^E \neq 0)$

Show Movie C

$$H = H_{SHO}^{S} \otimes \mathbf{1}^{E} + q_{SHO}H_{I}^{E} + \mathbf{1}^{S} \otimes H^{E}$$

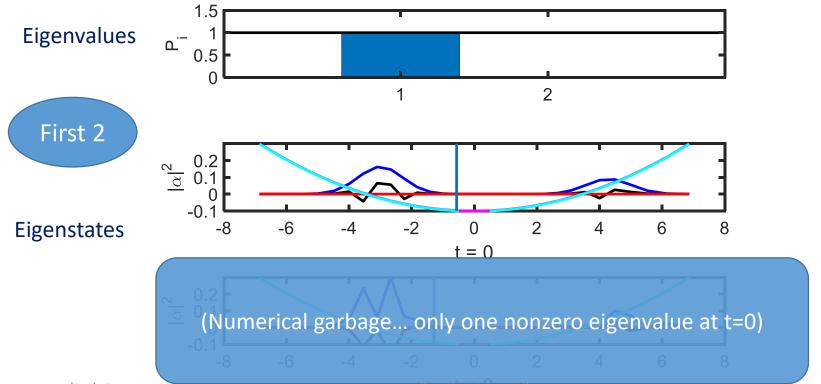
This evolution will take an initial product stat mixed state:

Show First 2 eigenstates and eigenvalues

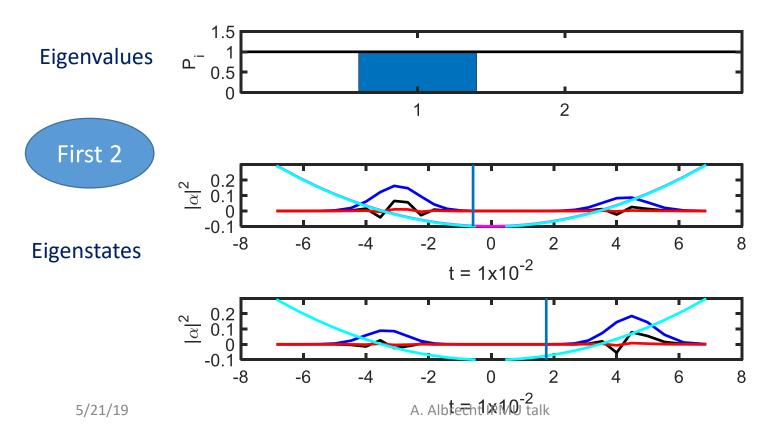
$$|\psi\rangle_{W} = |\psi\rangle_{S} |\psi\rangle_{E} \rightarrow \sum_{i,j} \alpha_{ij} |i\rangle_{S} |j\rangle_{E}$$

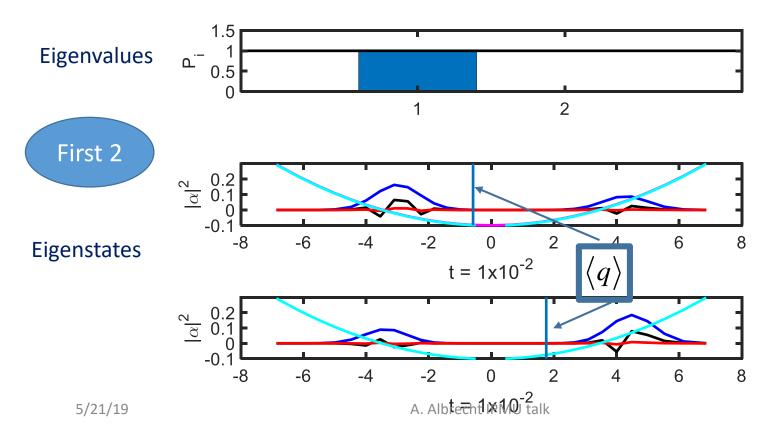
$$\rho_{S} \equiv Tr_{E}(|\psi\rangle_{WW}\langle\psi|) = |\psi\rangle_{SS}\langle\psi| \rightarrow \text{more general } \rho_{S}$$

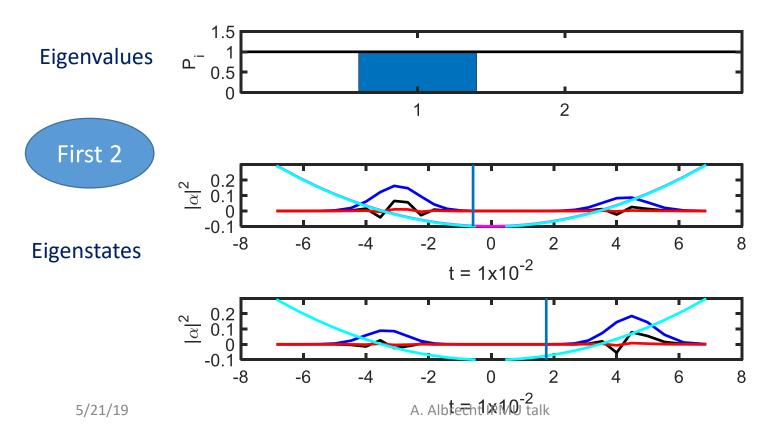


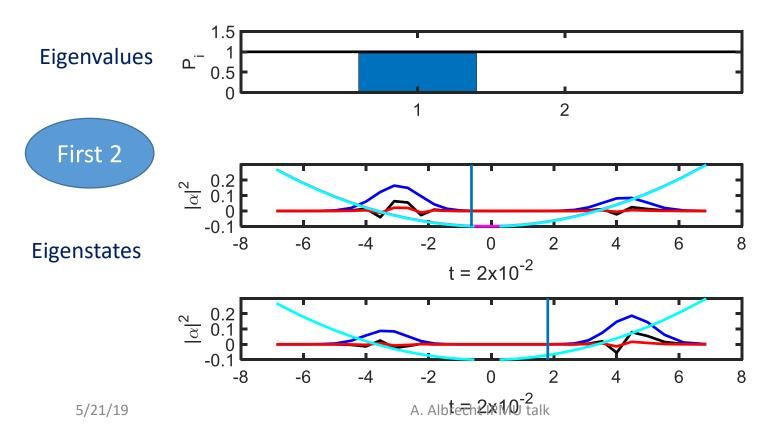


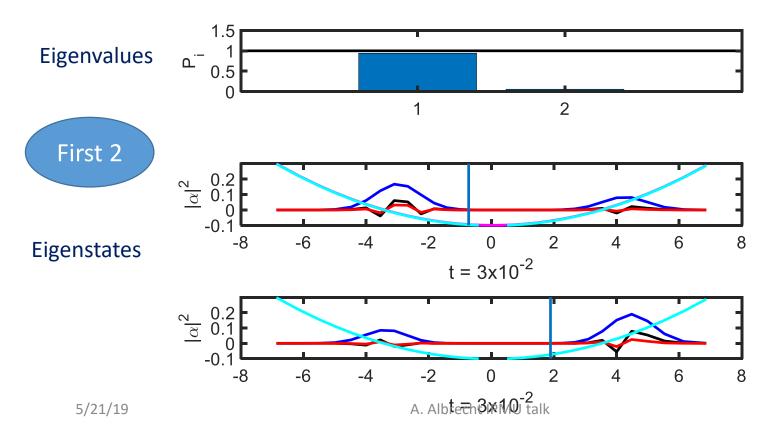
A. Albrecht PMU talk

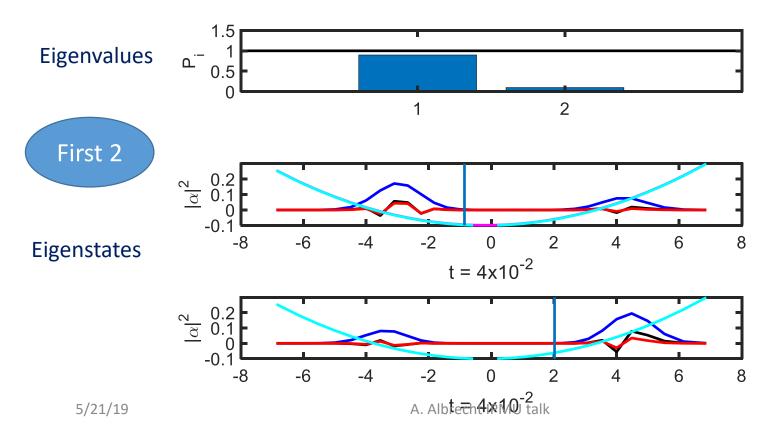


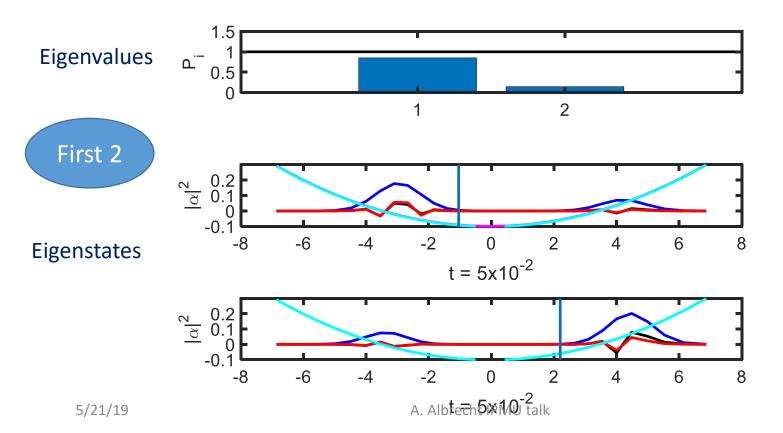


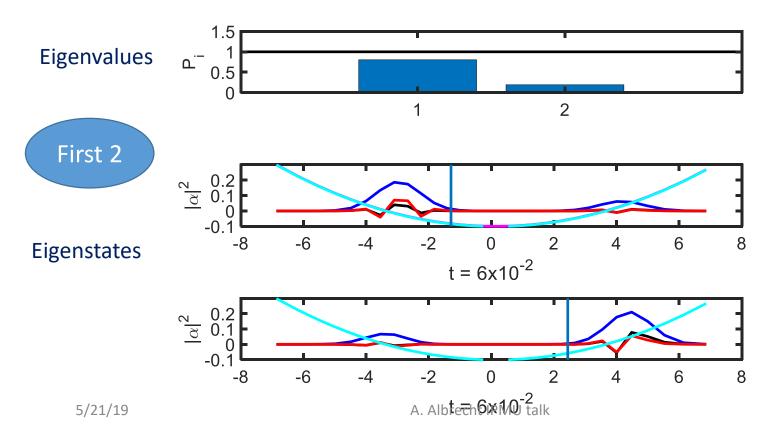


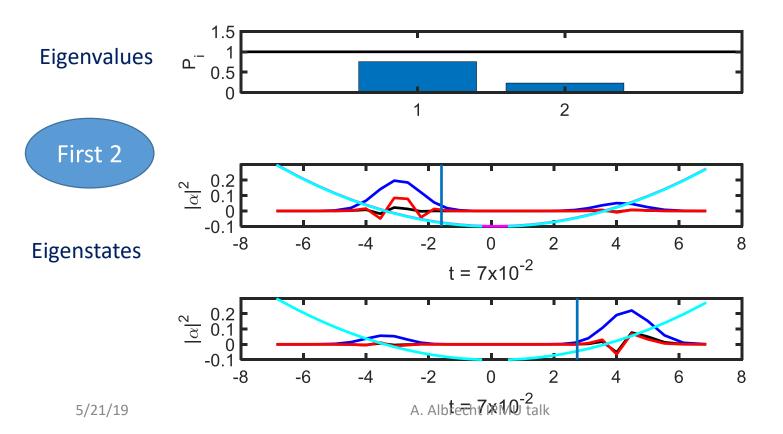


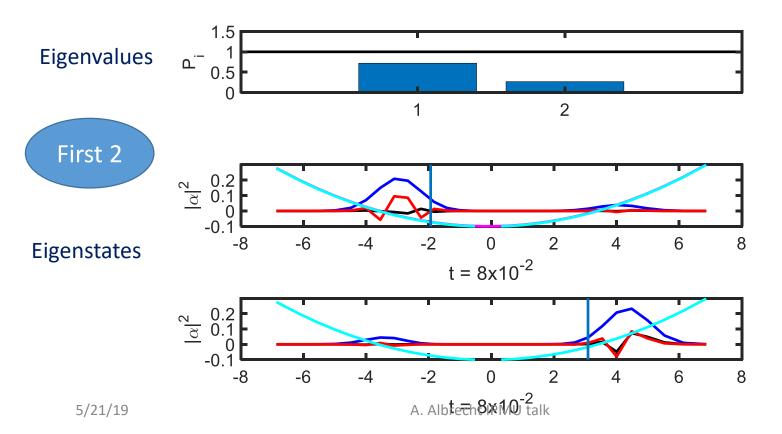


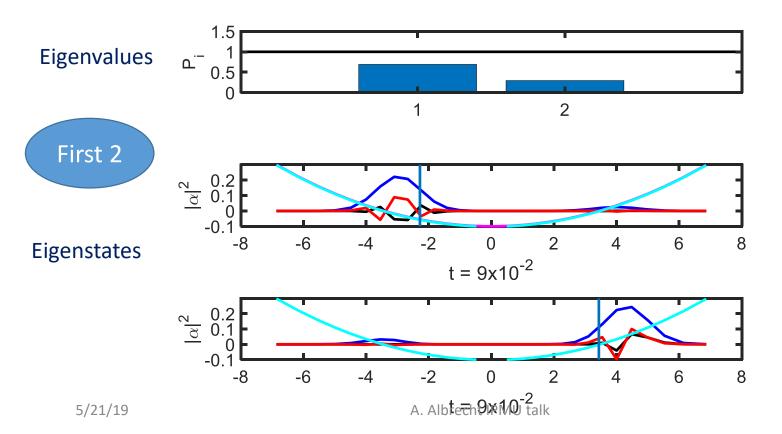


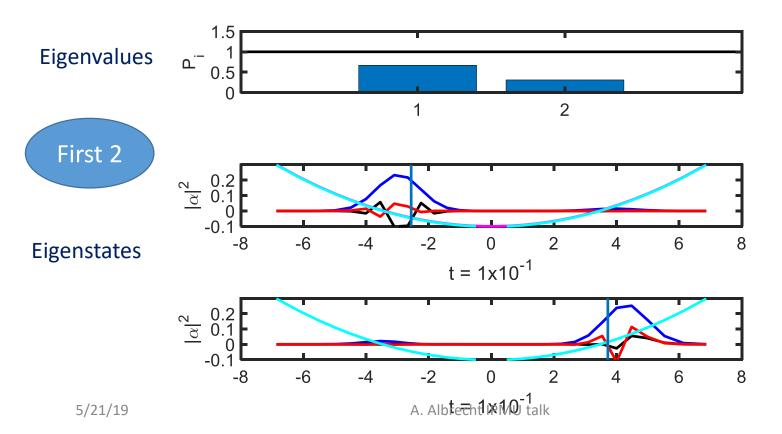


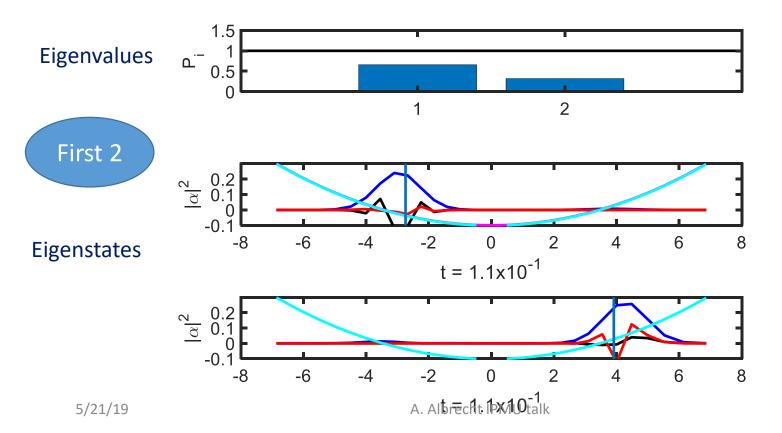


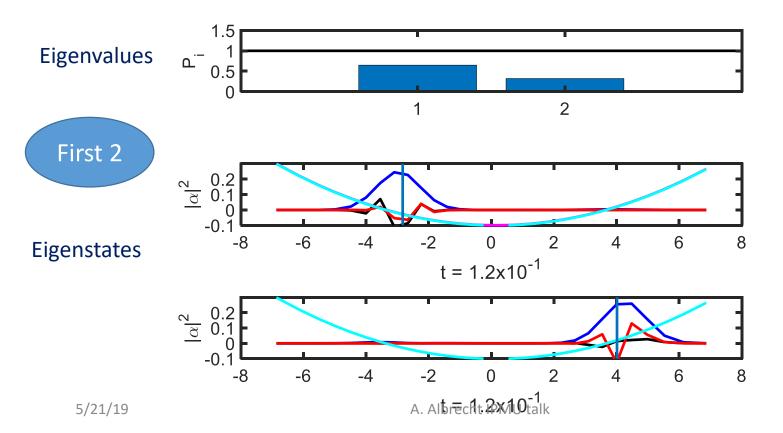


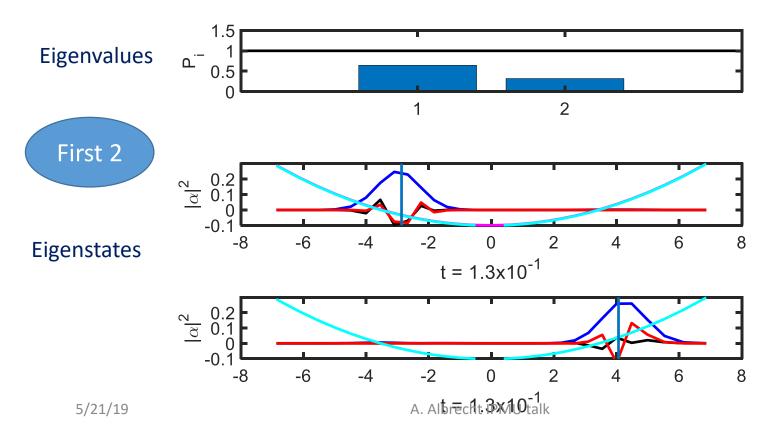


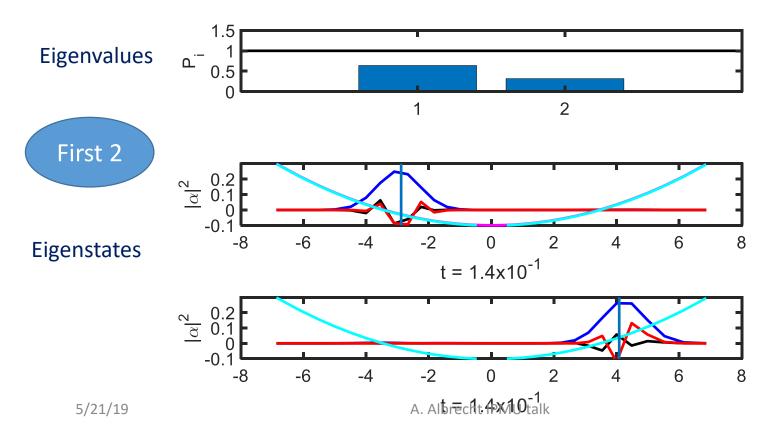




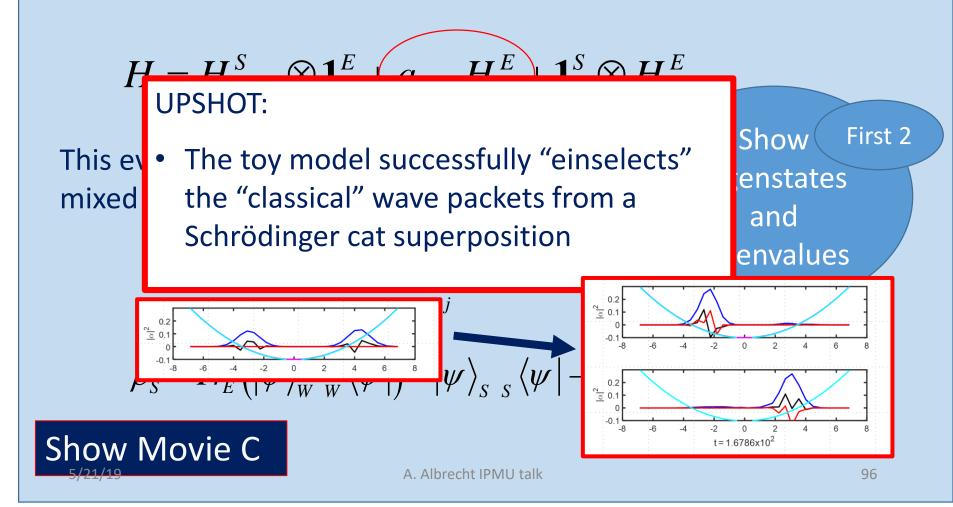








• Interactions turned on $(H_I^E \neq \mathbf{0})$



Outline

- 1. Motivations
- 2. Introduction to einselection and the toy model
- 3. Einselection in equilibrium (technical explorations and overall assessment)
- 4. Eigenstate Einselection Hypothesis (if there is time)

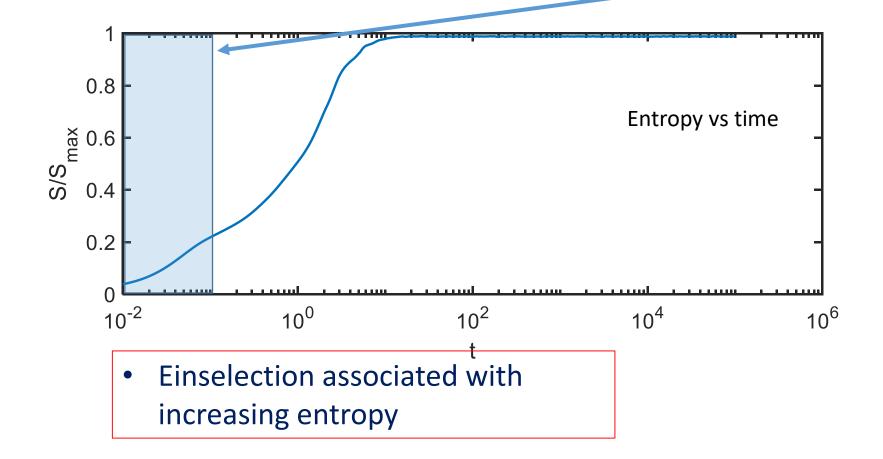
5. Conclusions

Outline

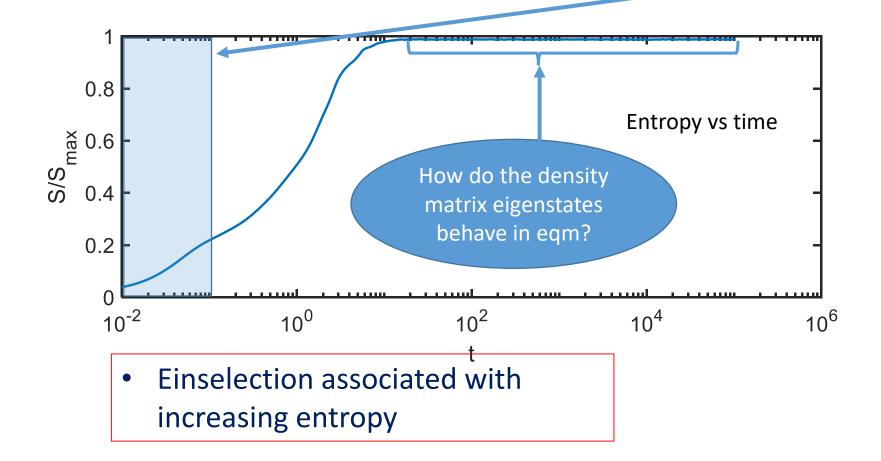
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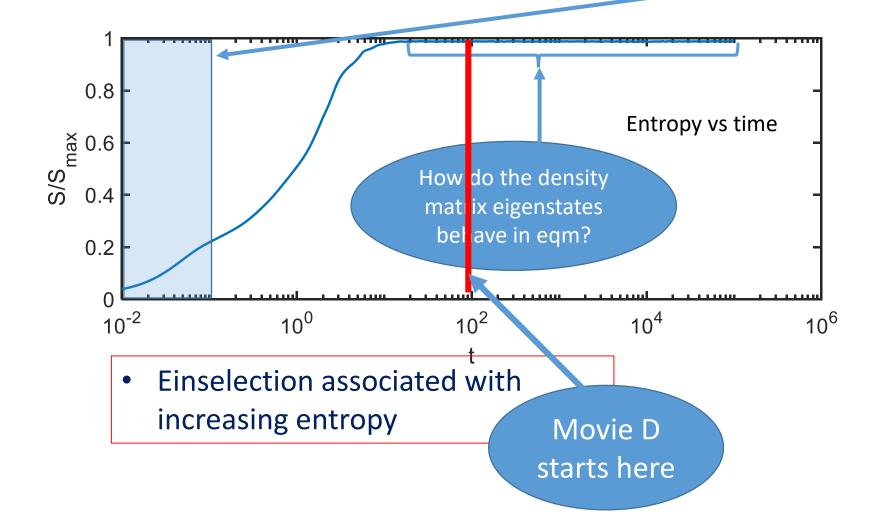
The collapsing Schrödinger cat (Movie C) was in this time window



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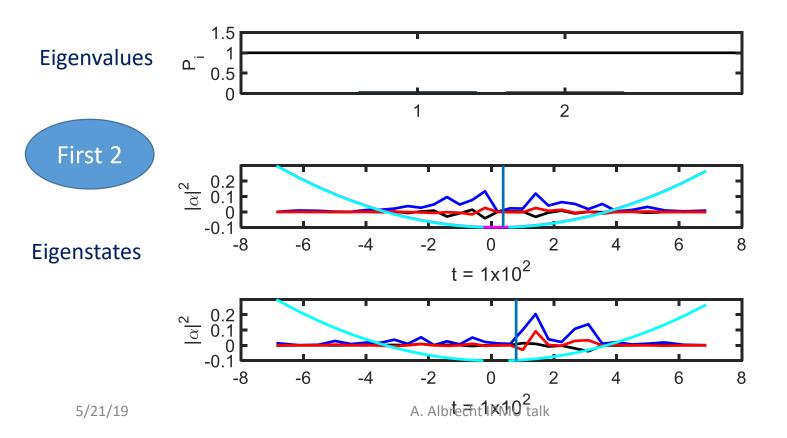


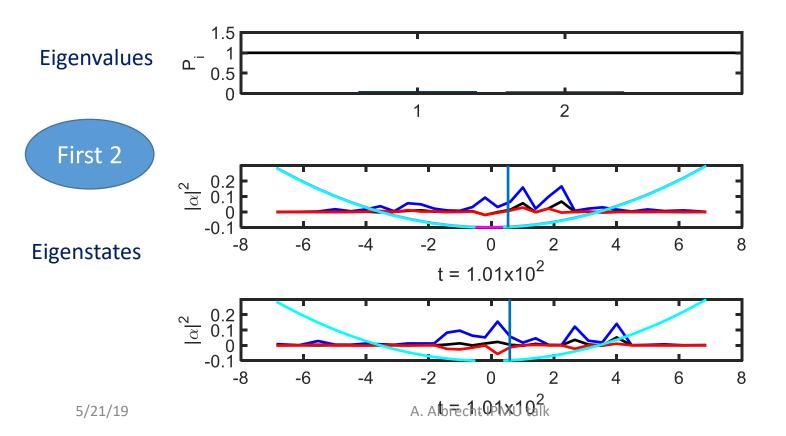
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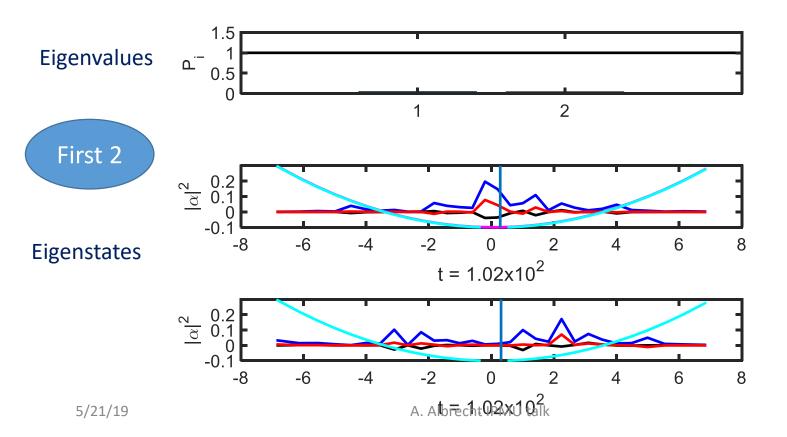


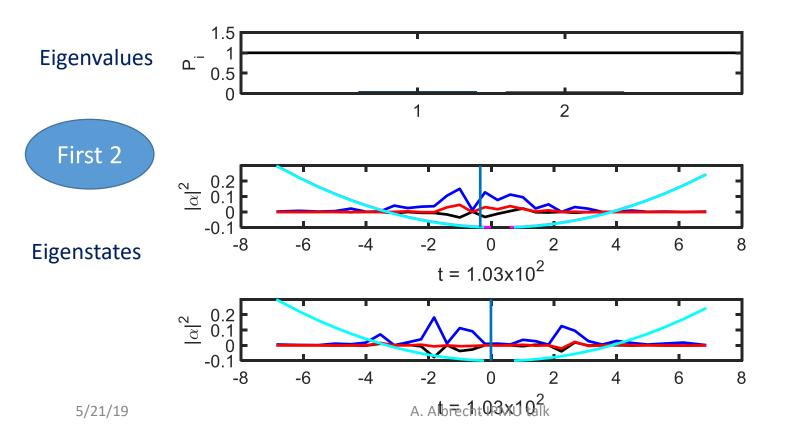
Show Movie D

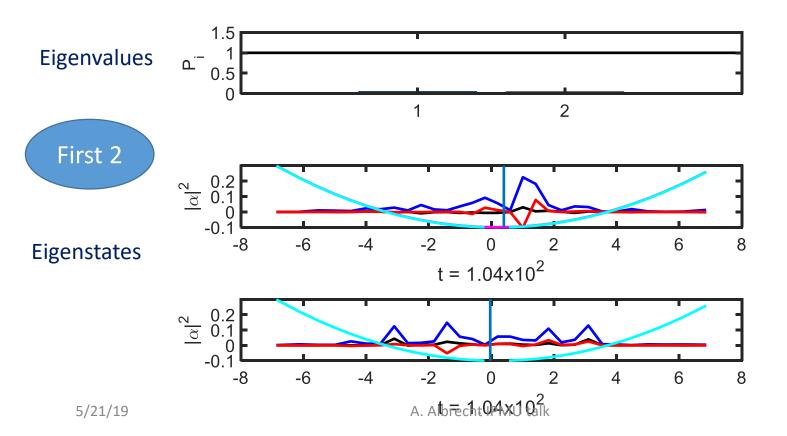
Discuss "detailed balance" (wallowing) of everett worlds at board

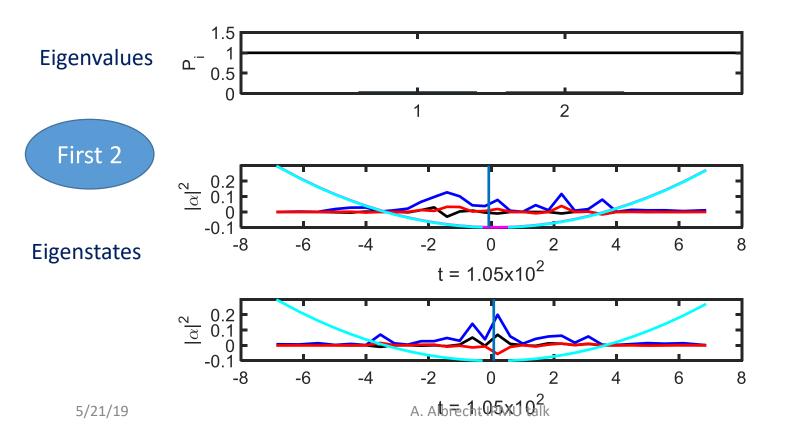


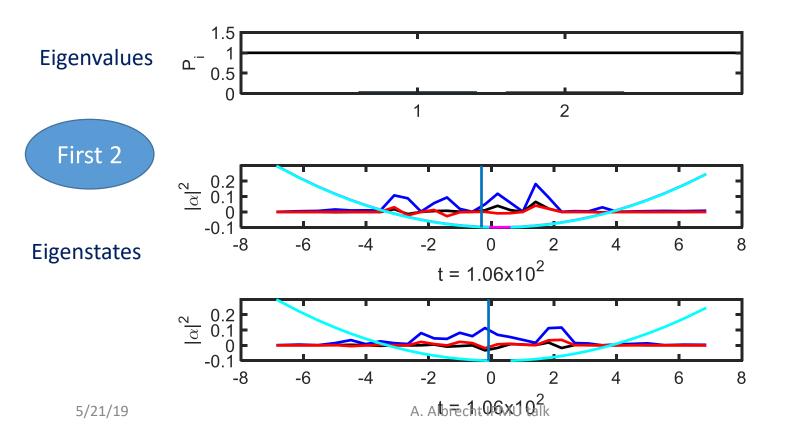


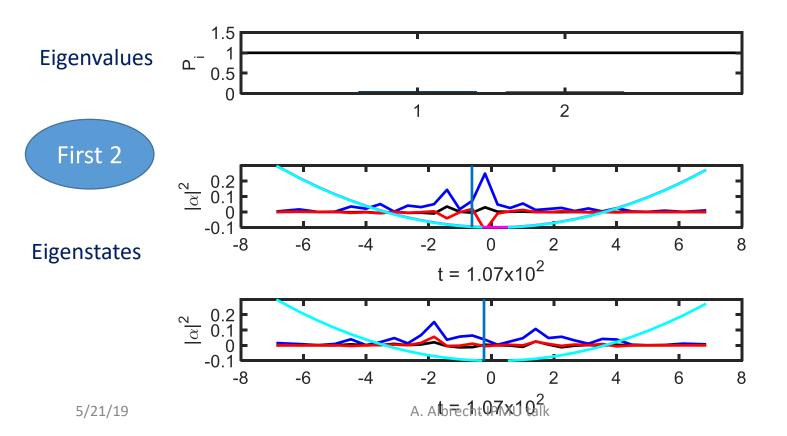


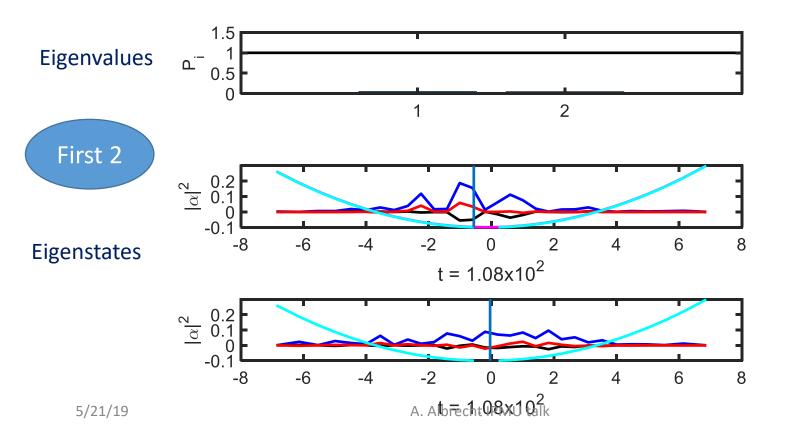


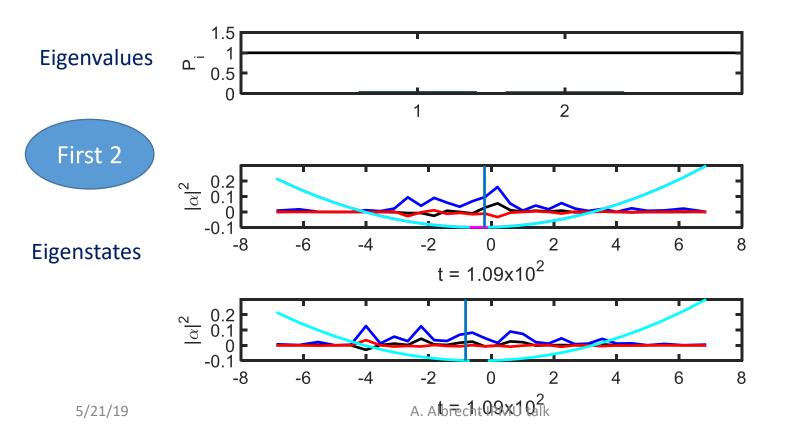


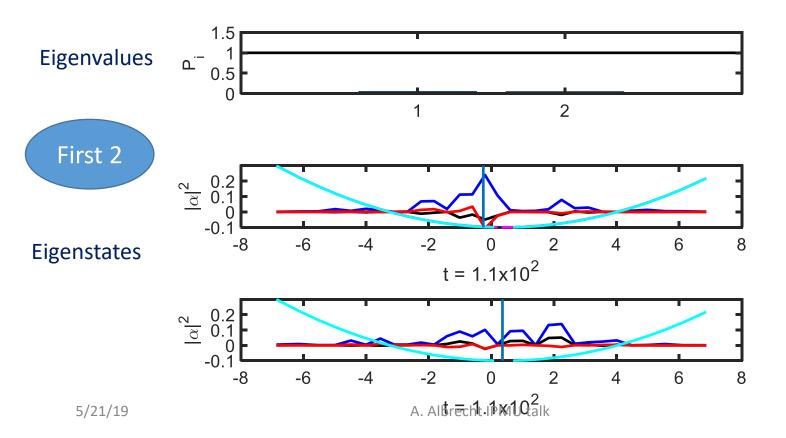


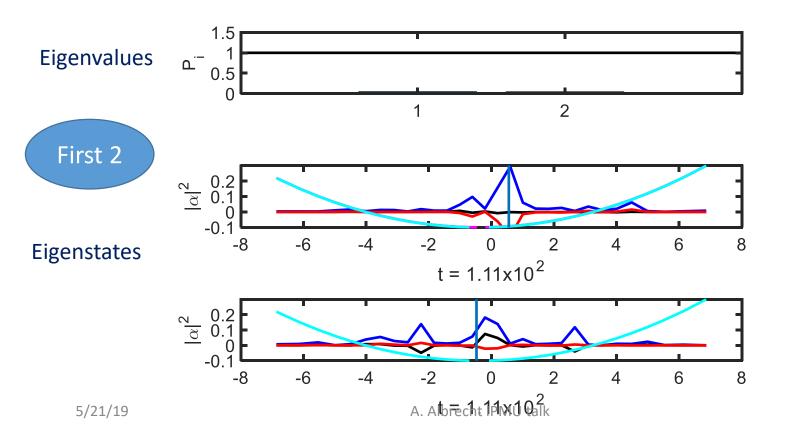


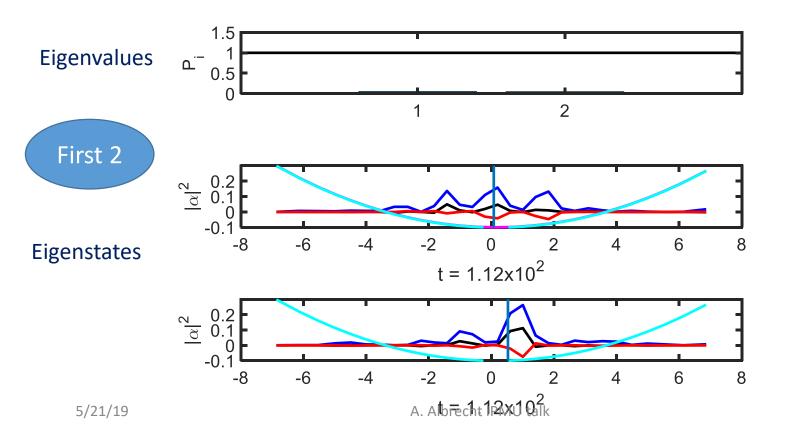


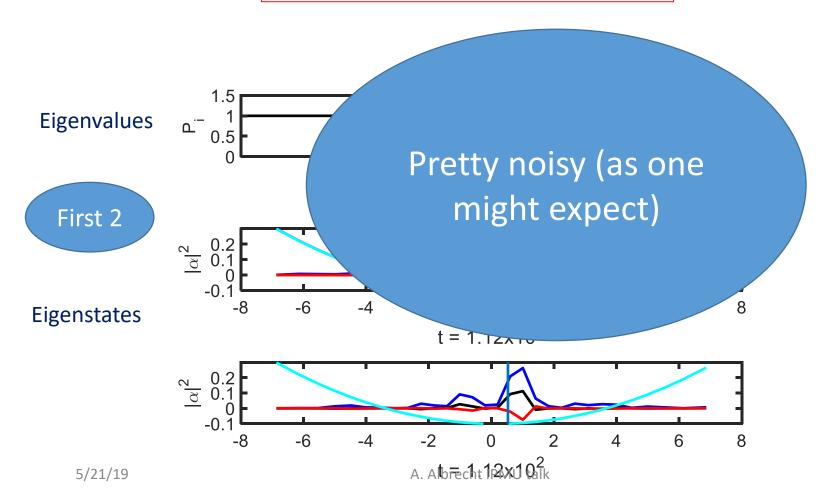


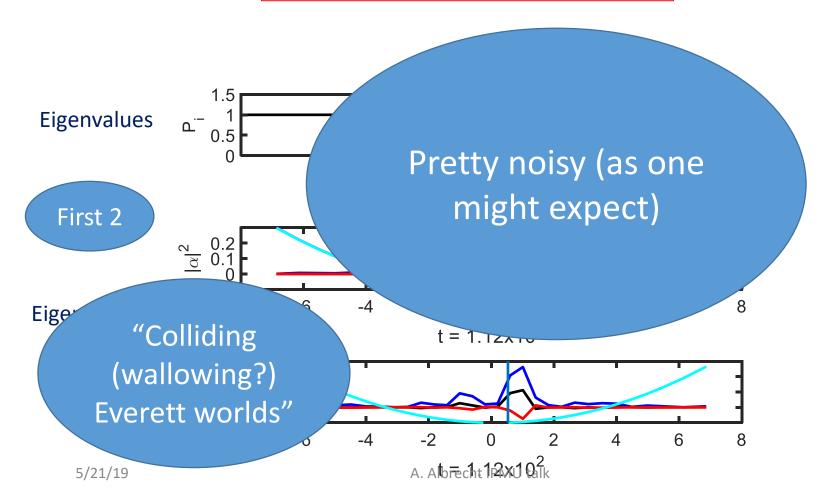












However, when the coupling strength is reduced, things get more interesting

-6

-4

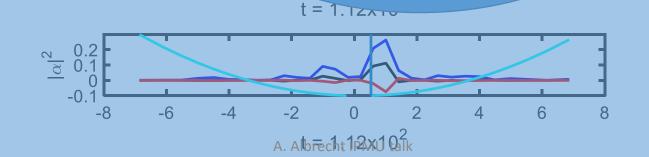
-0.1

-8

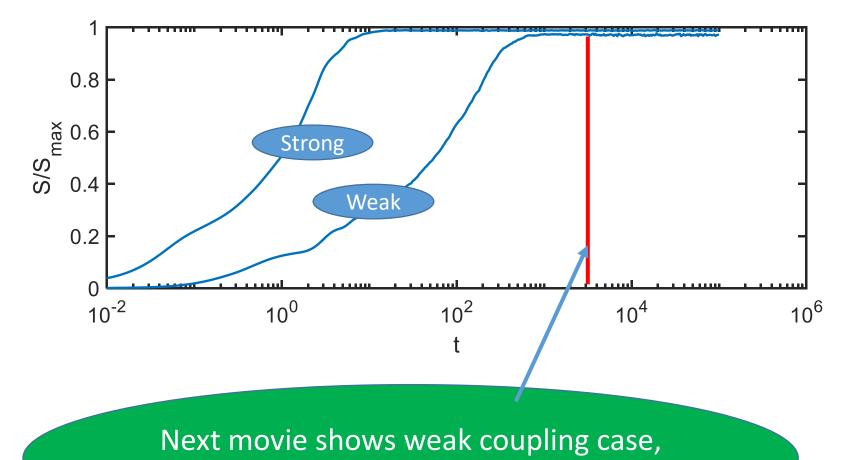
y noisy (as one night expect)

Eigenstates

Eigen



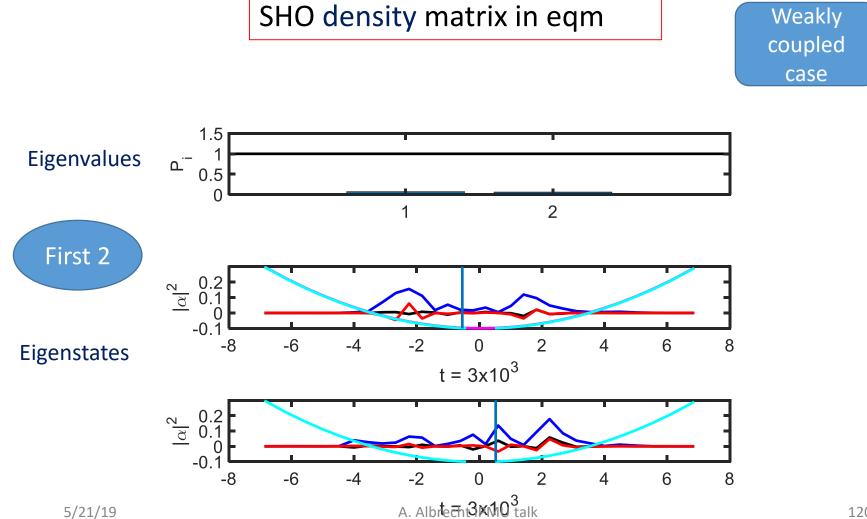
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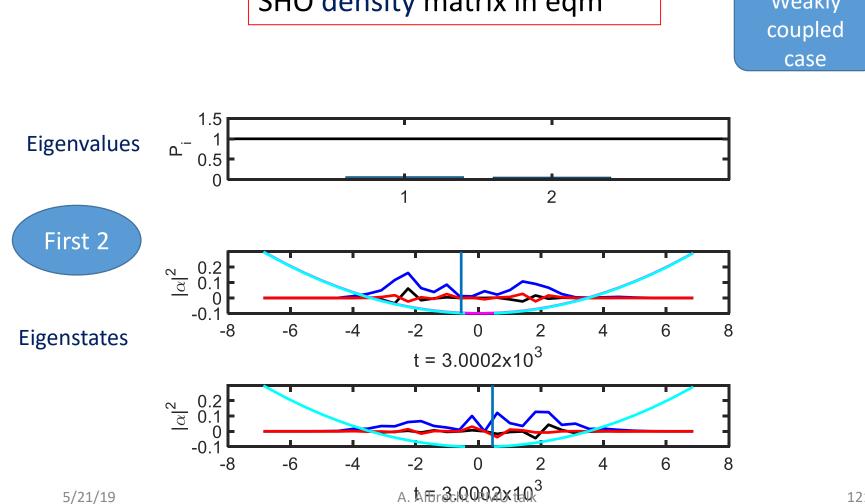


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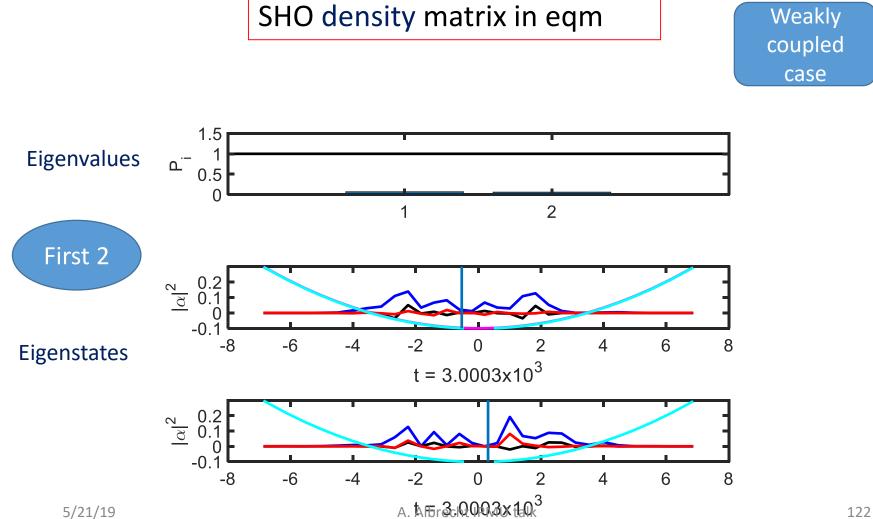


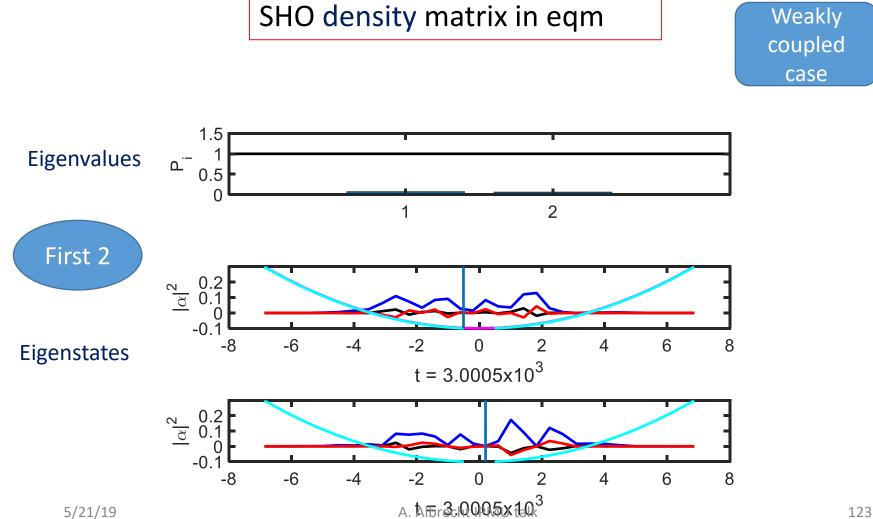
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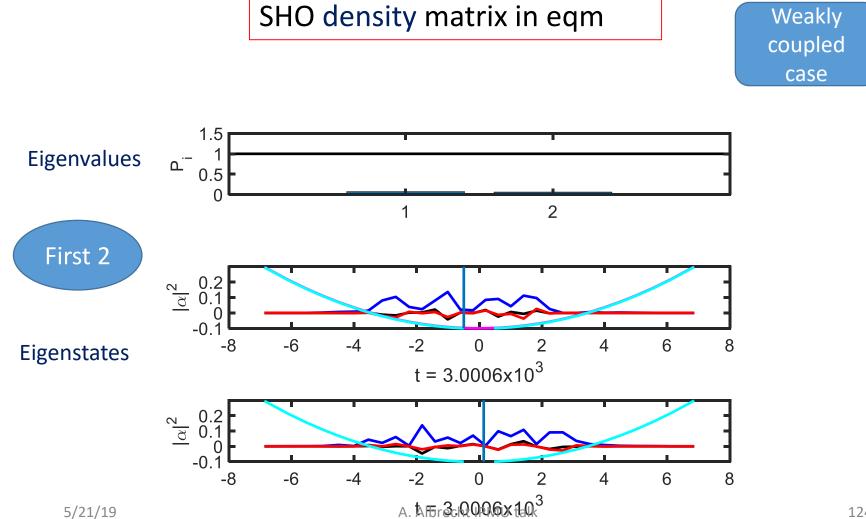


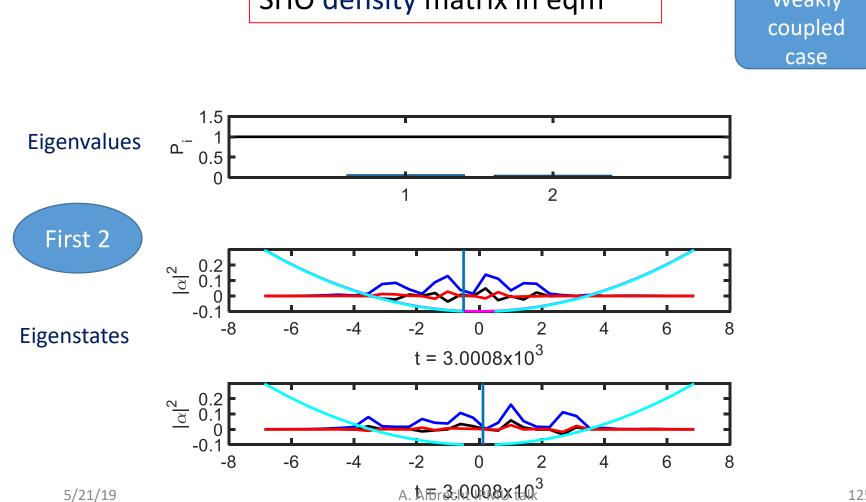


Weakly



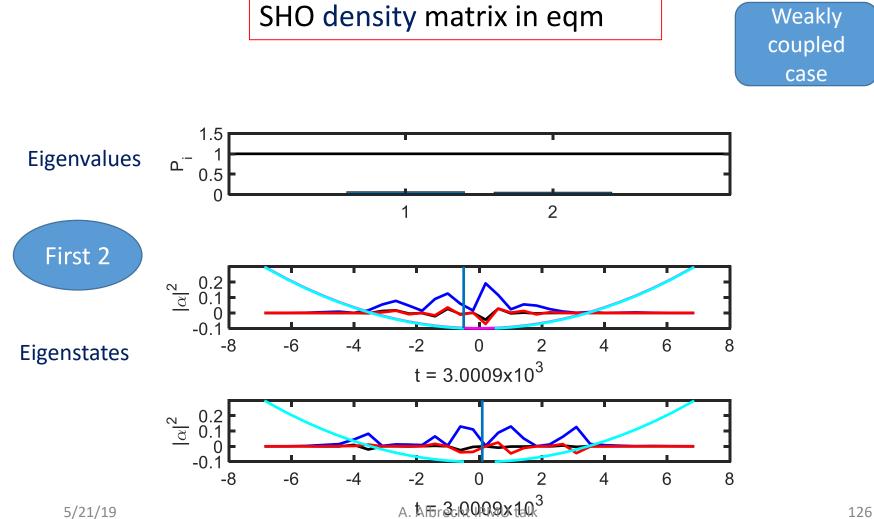


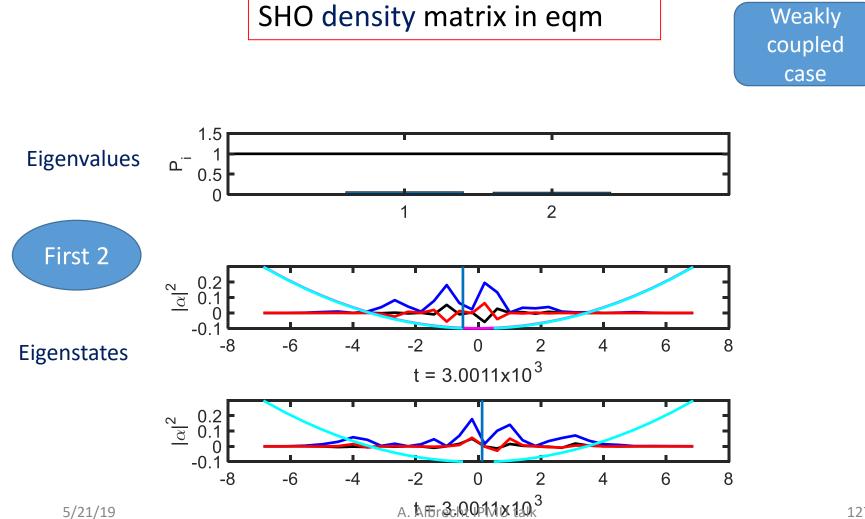


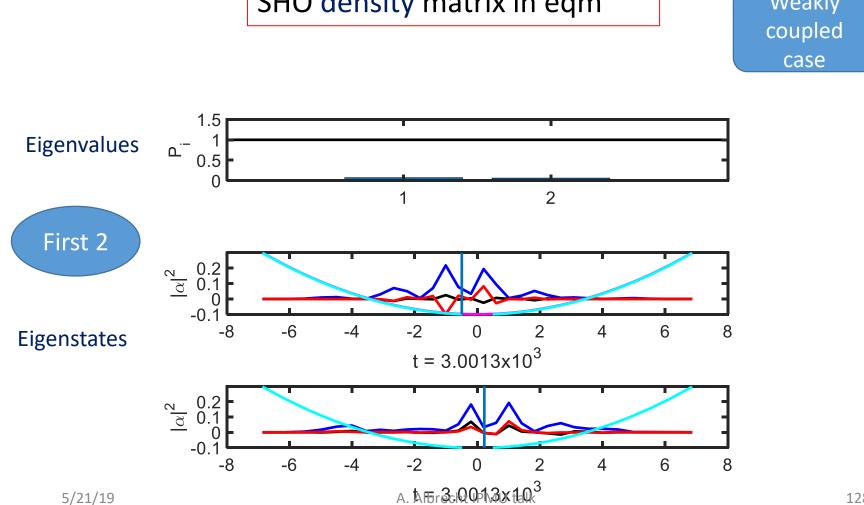


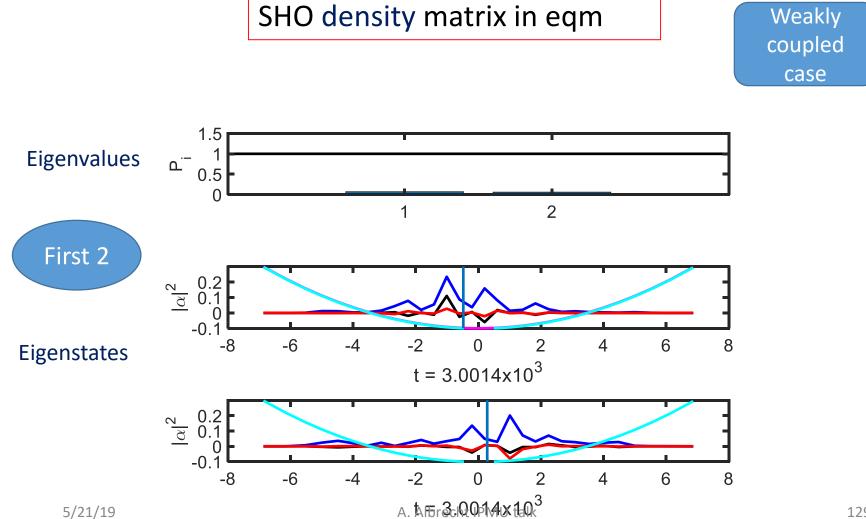
SHO density matrix in eqm

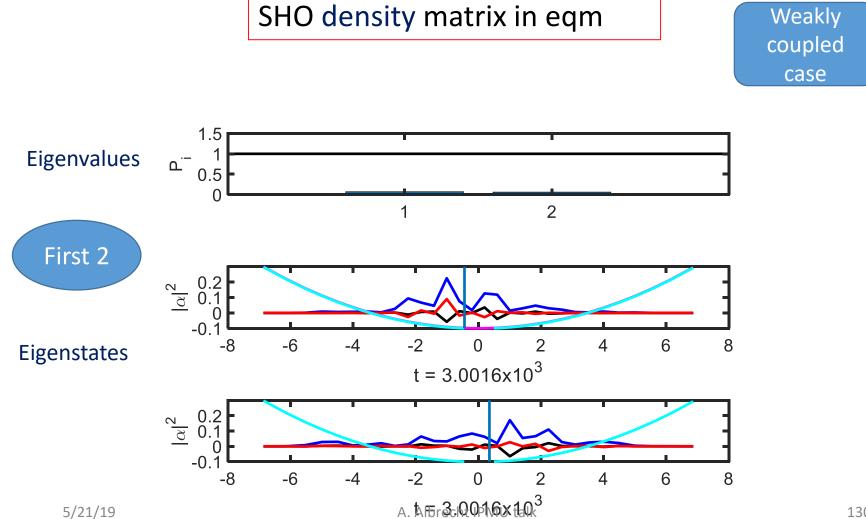
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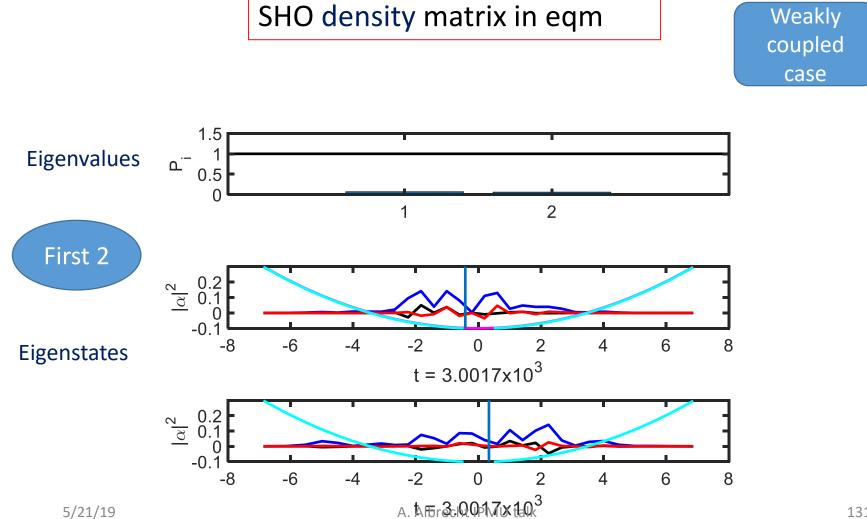


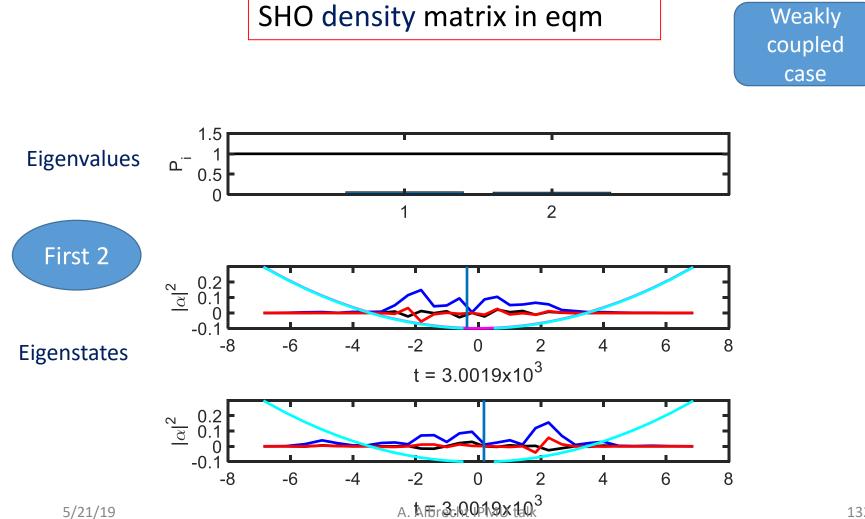


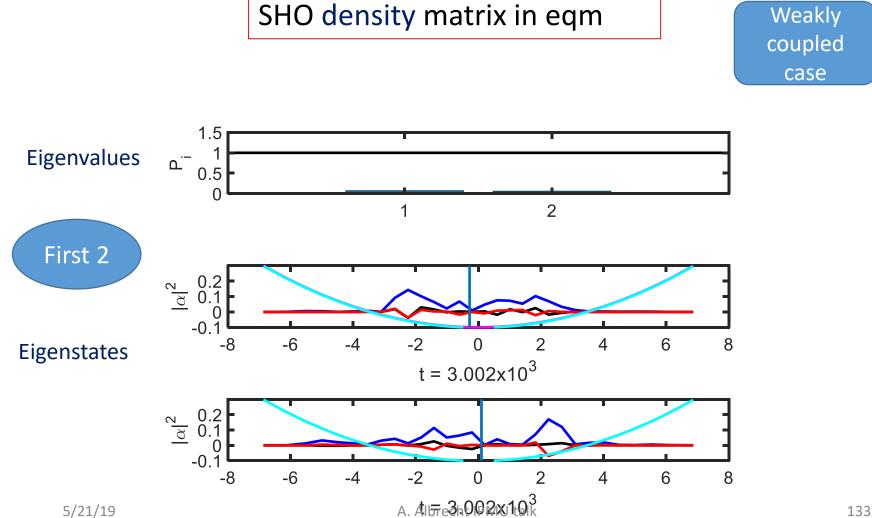


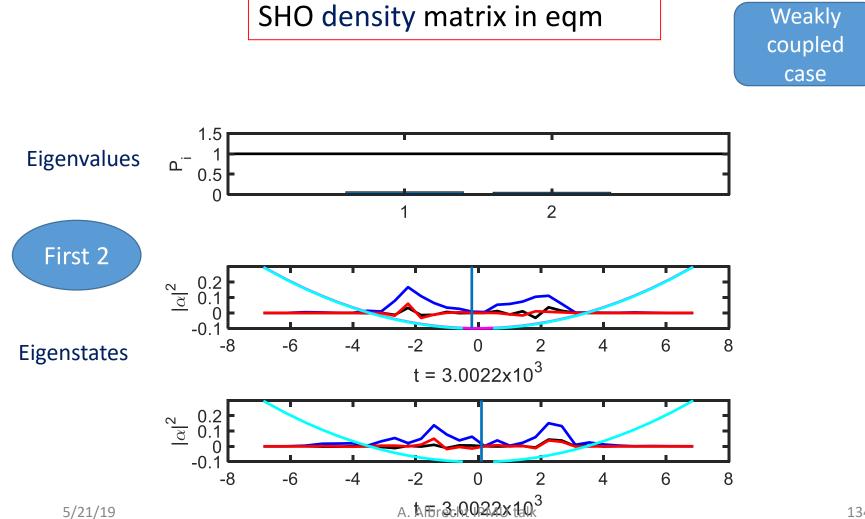


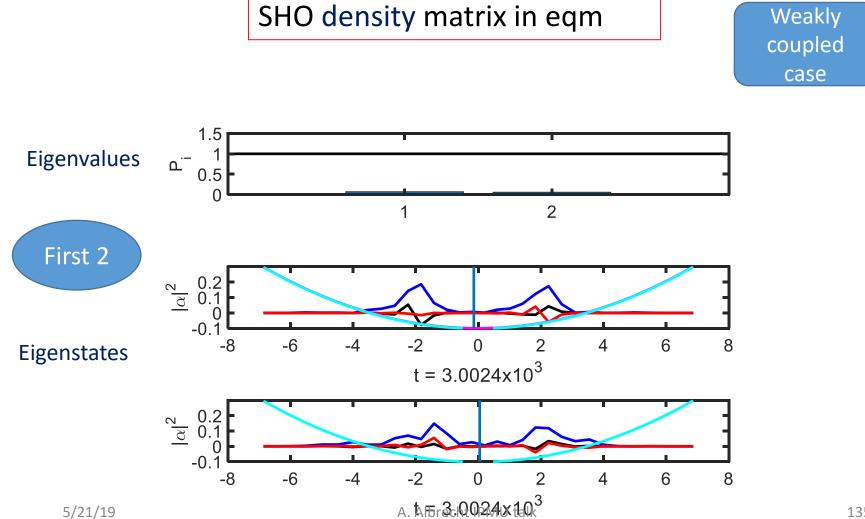




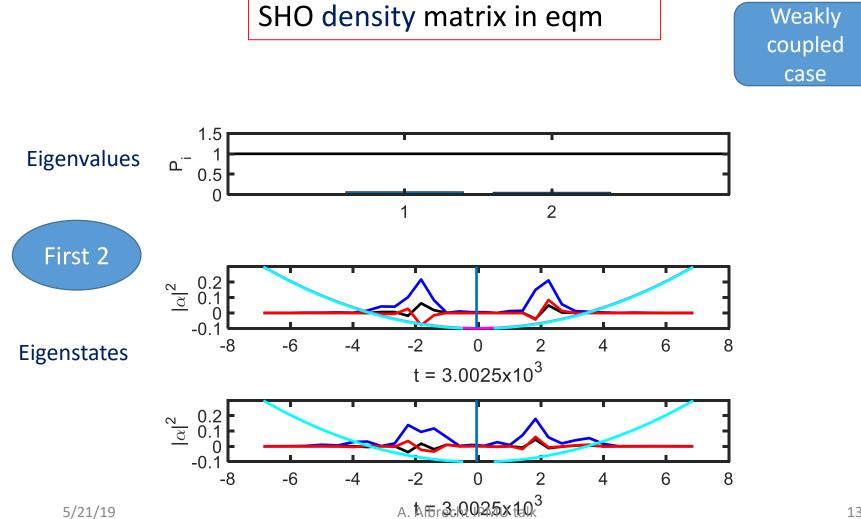


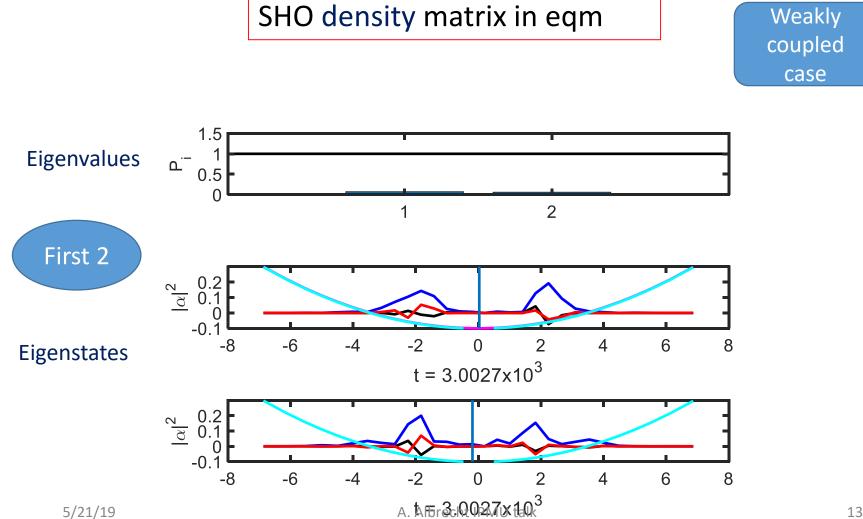


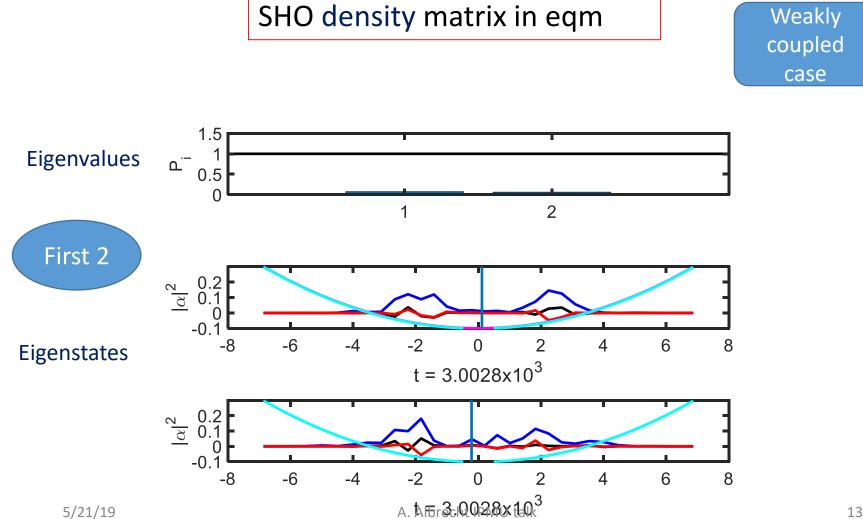


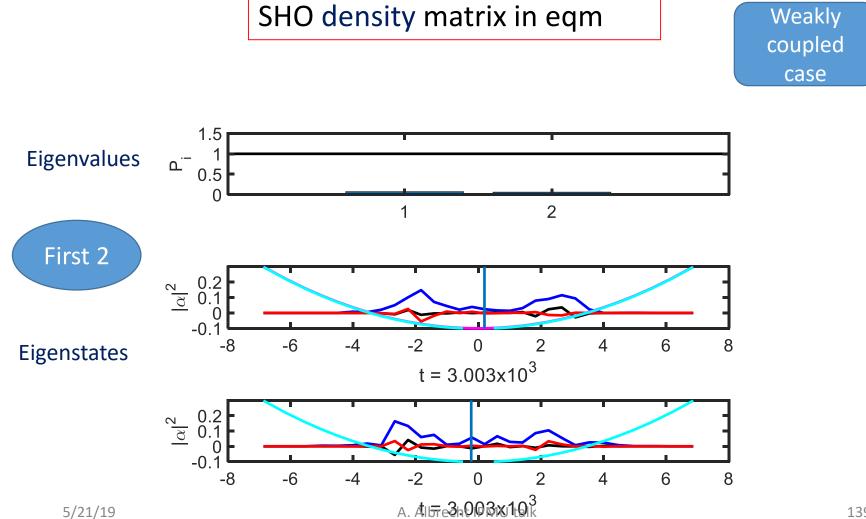


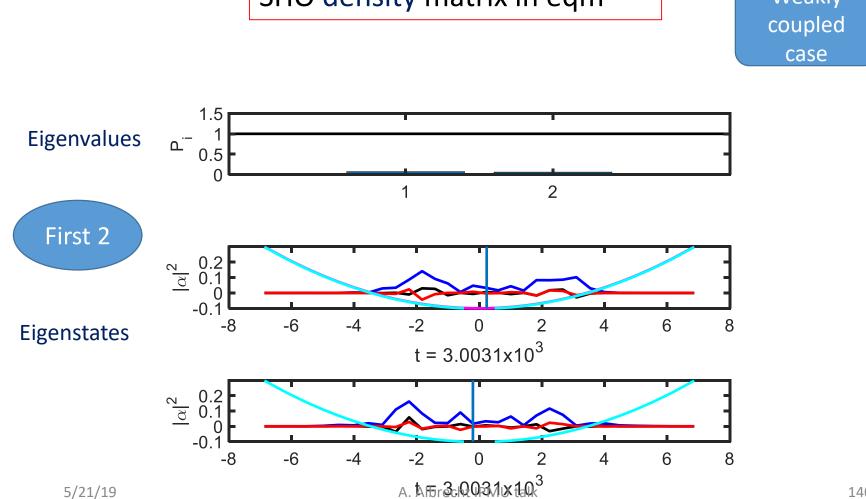






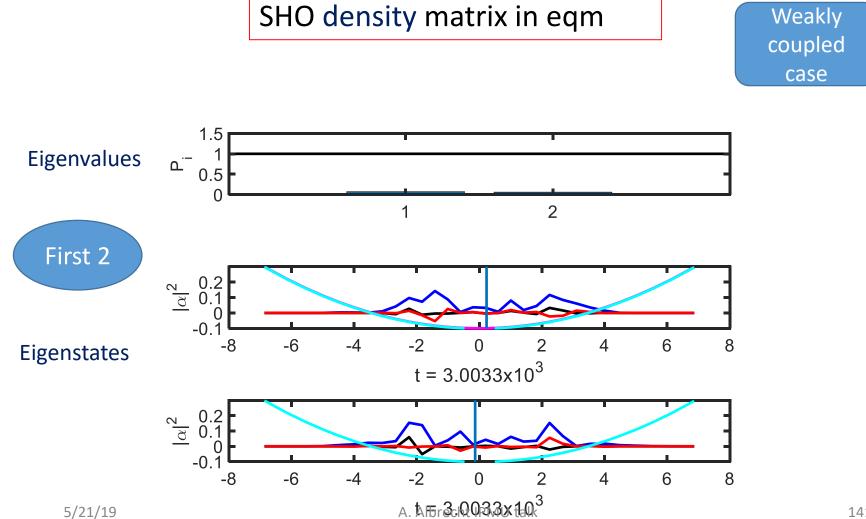


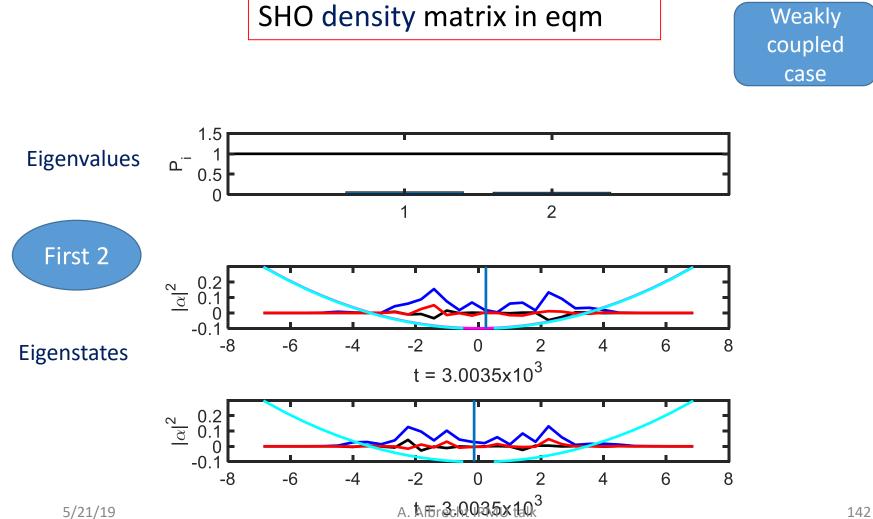


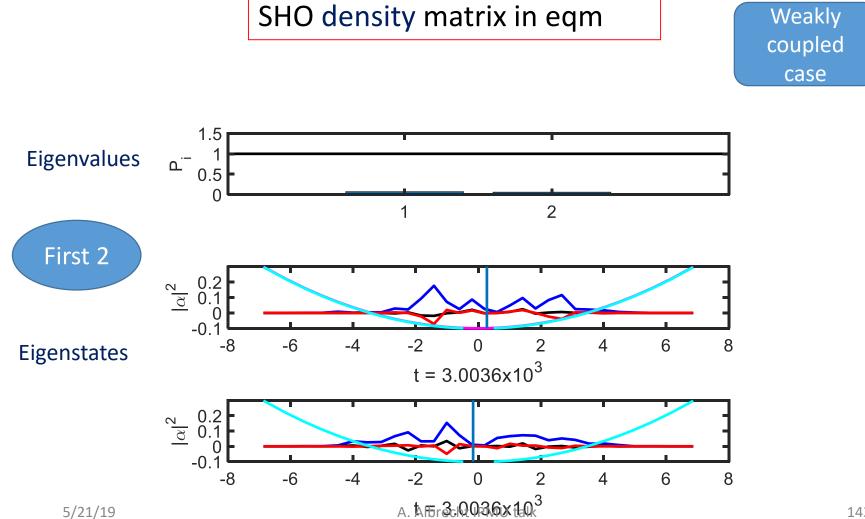


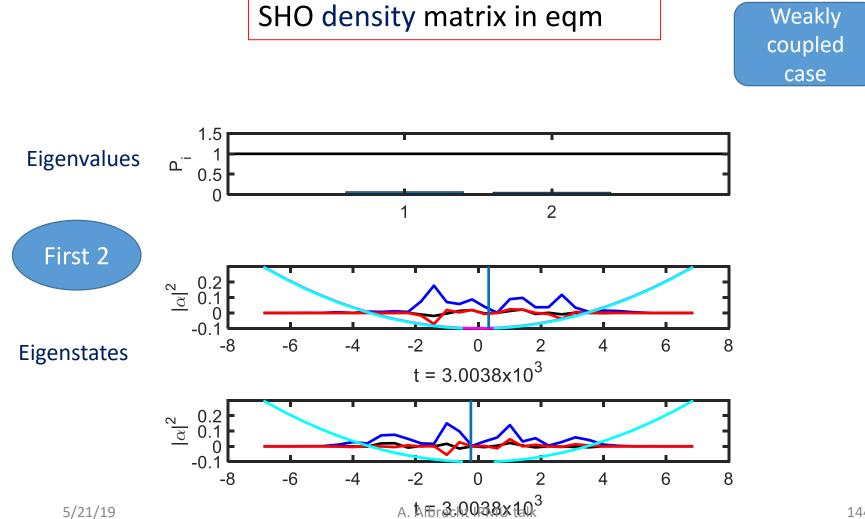
SHO density matrix in eqm

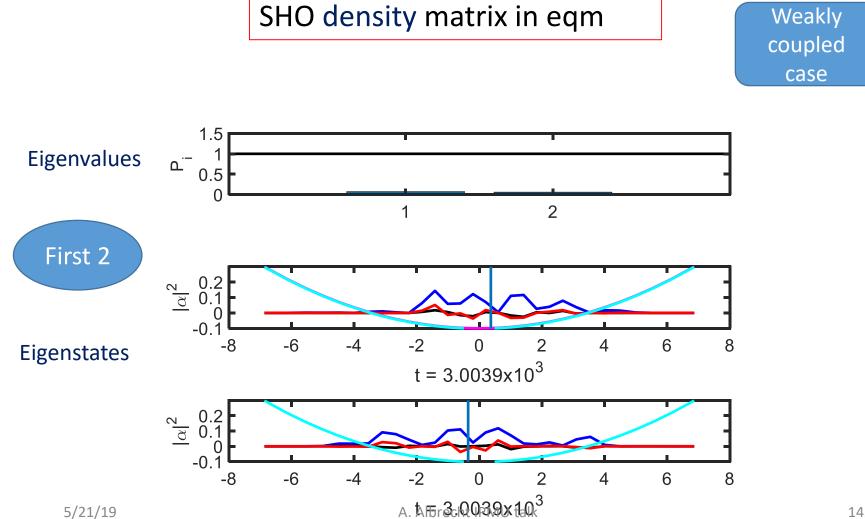
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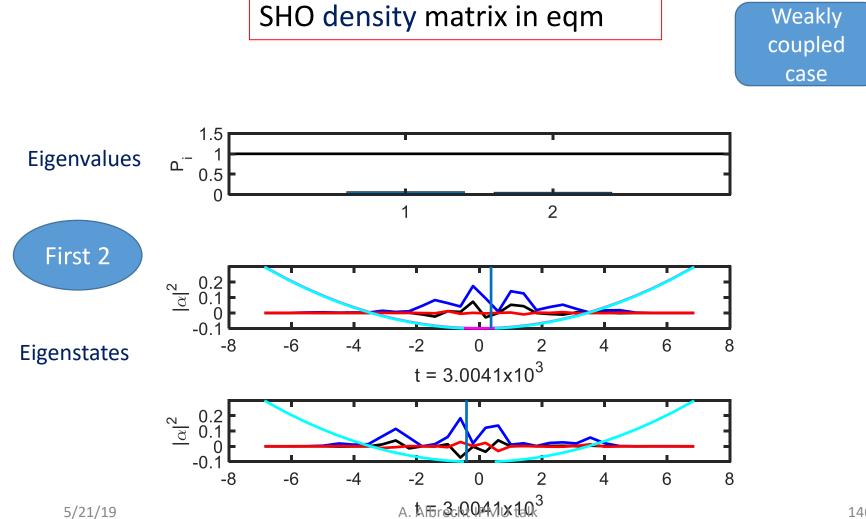


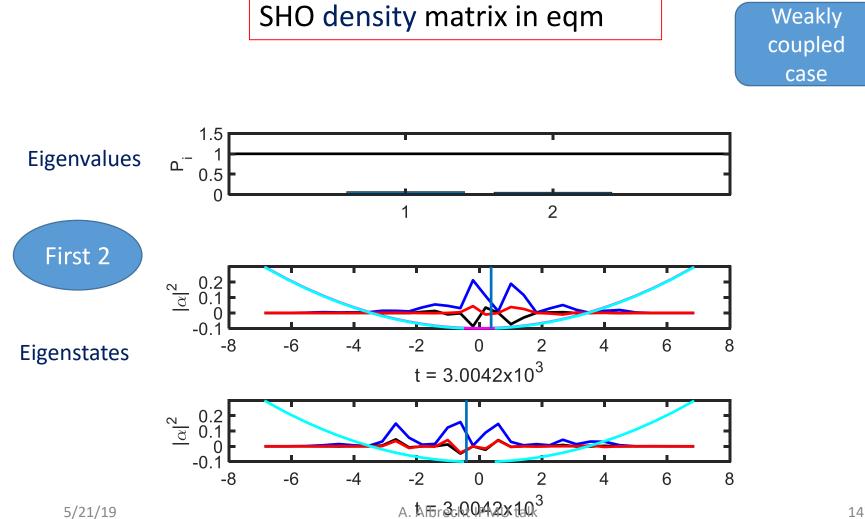


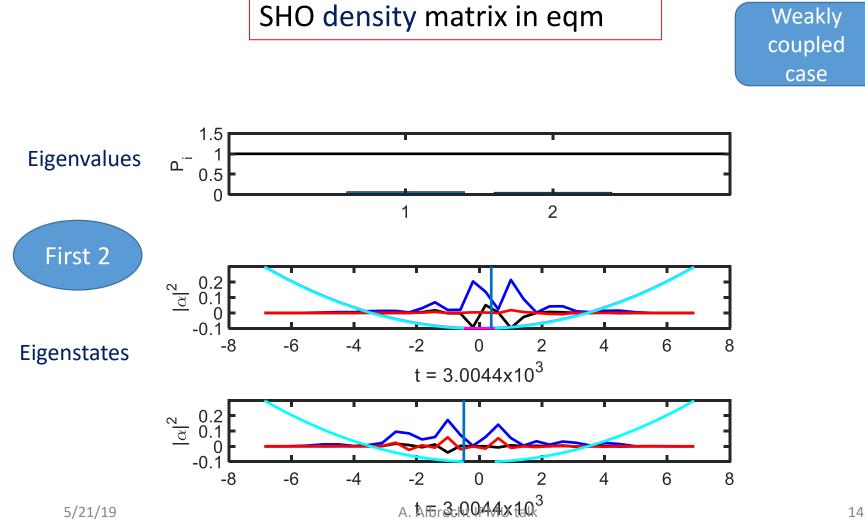


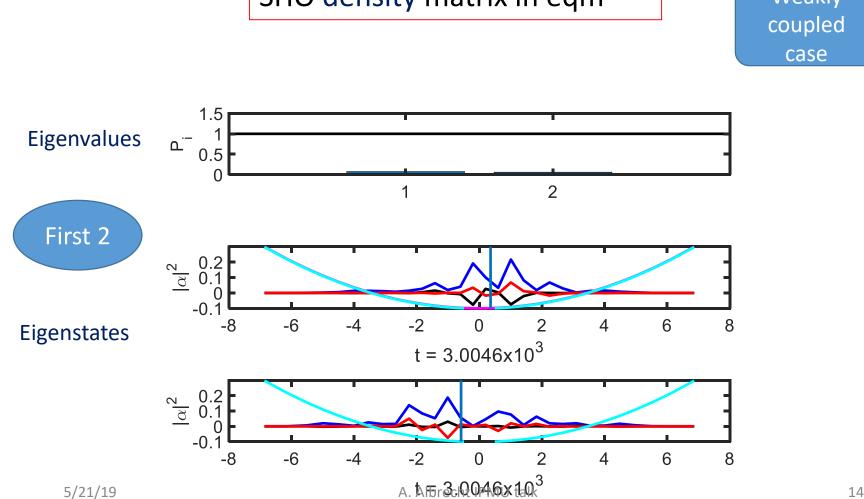




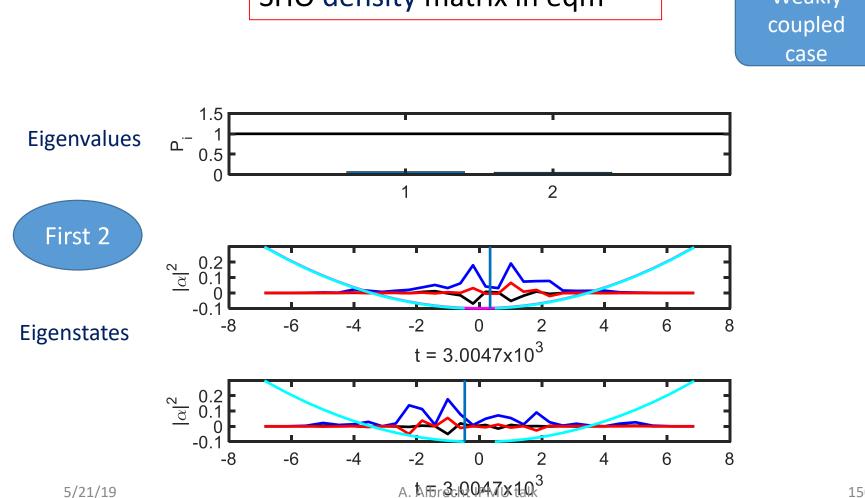






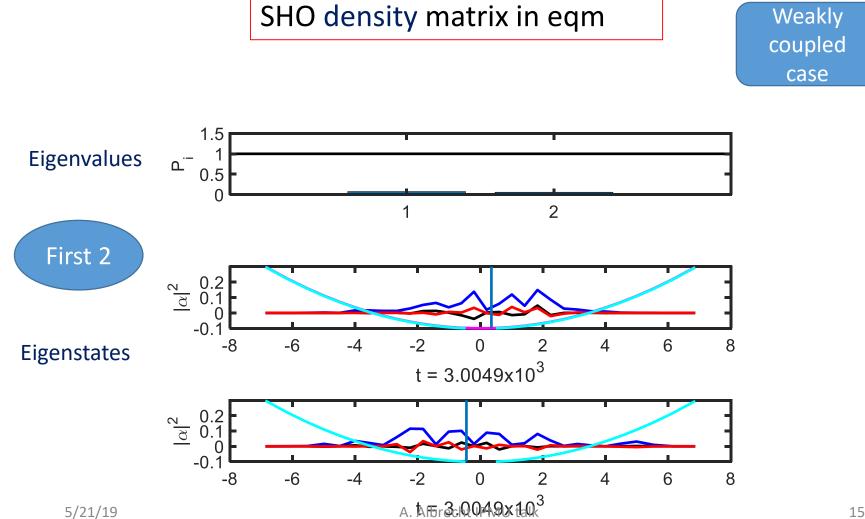


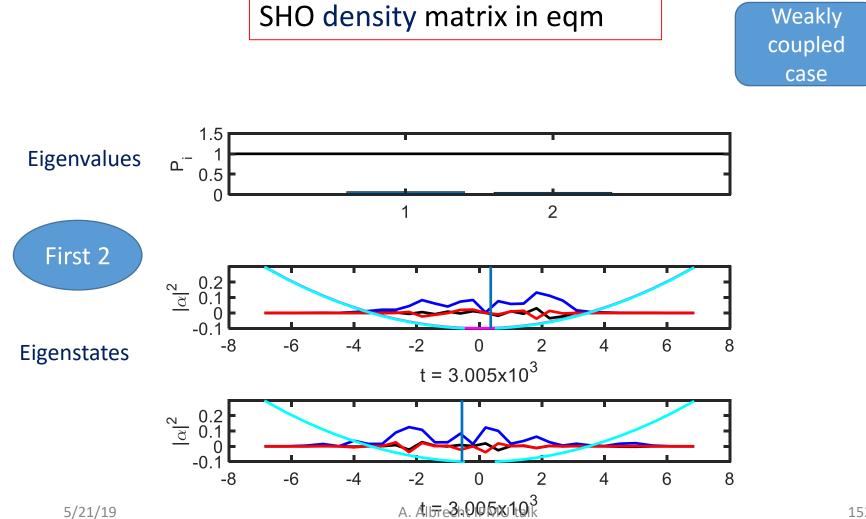
Weakly

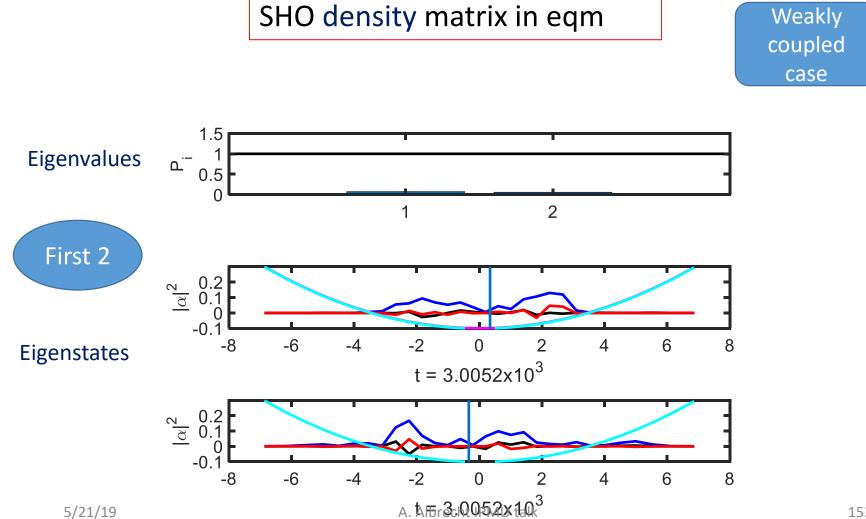


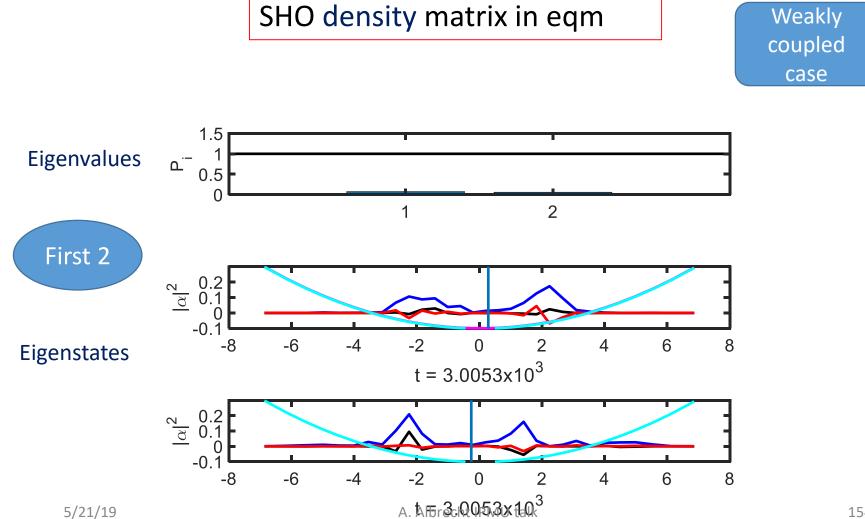
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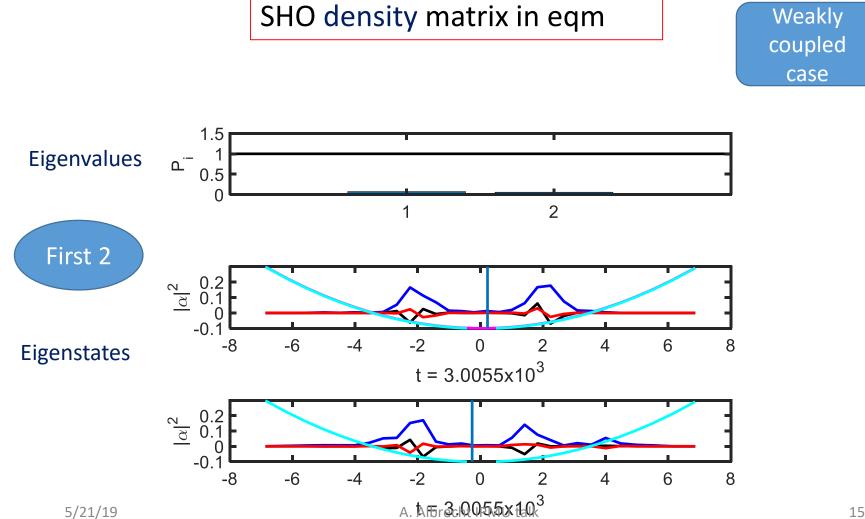
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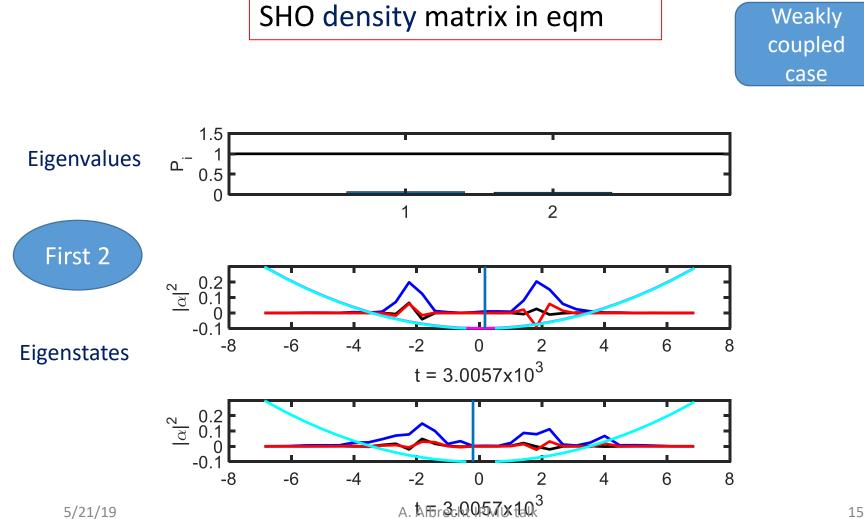


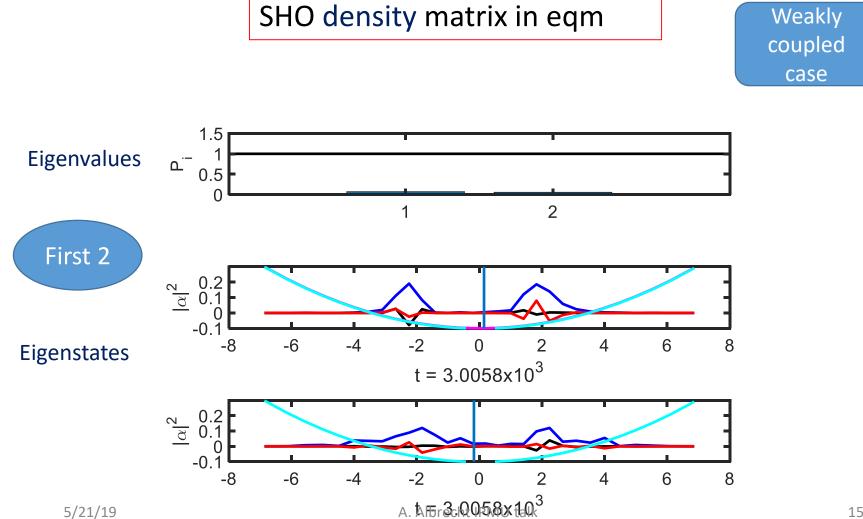


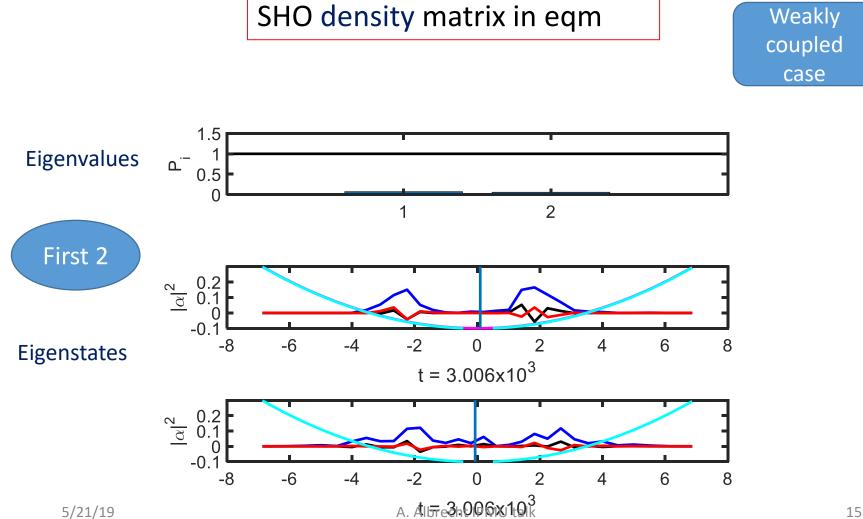


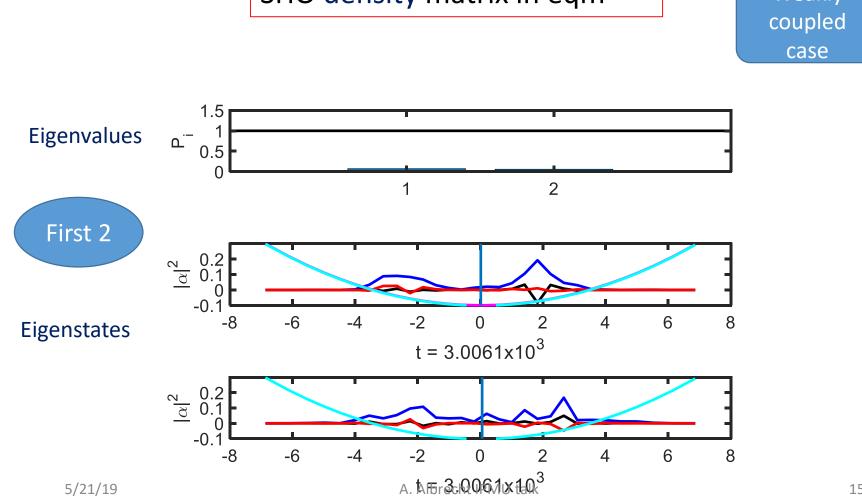


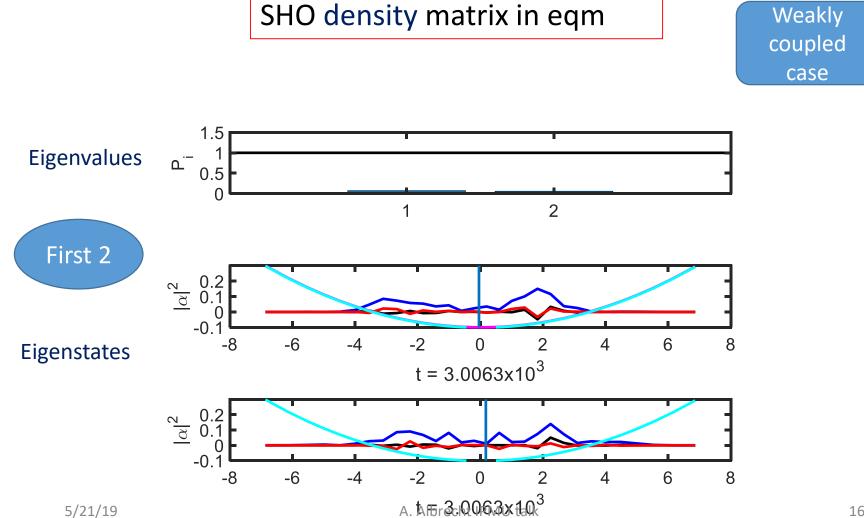


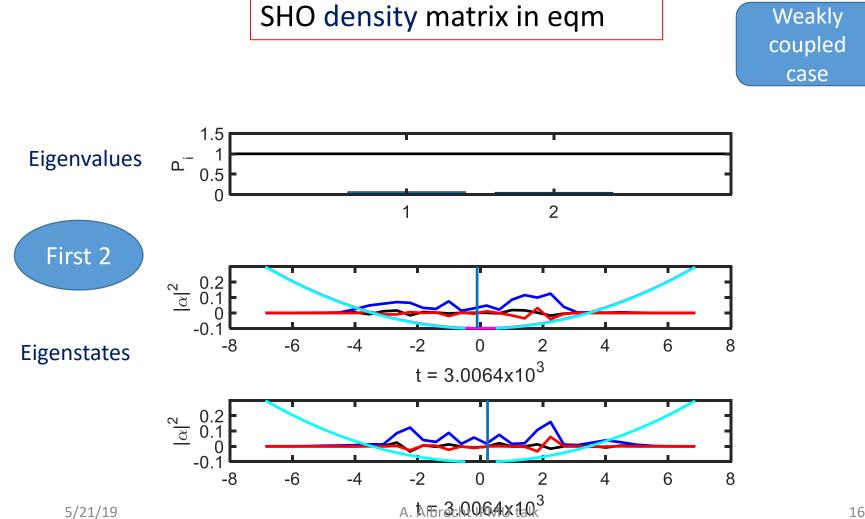


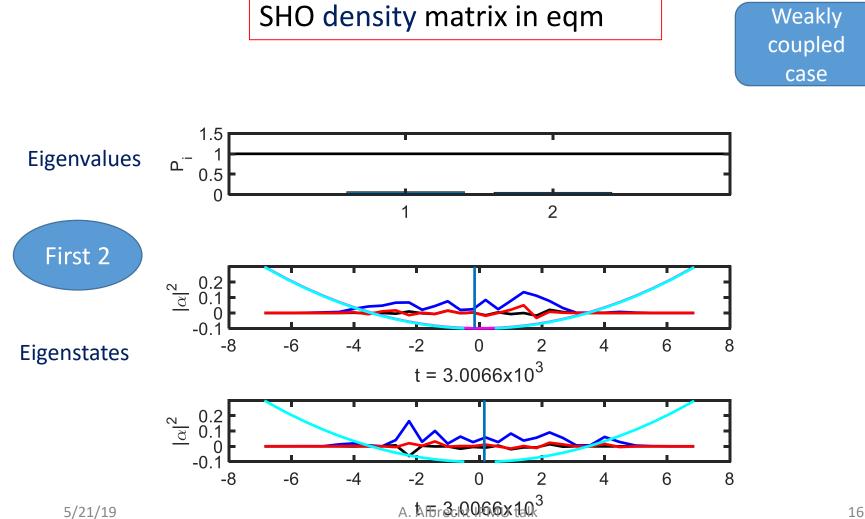


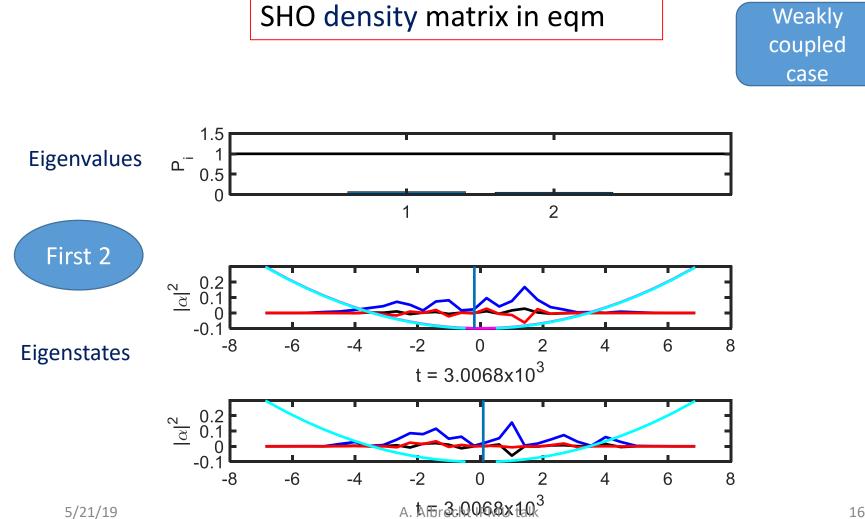


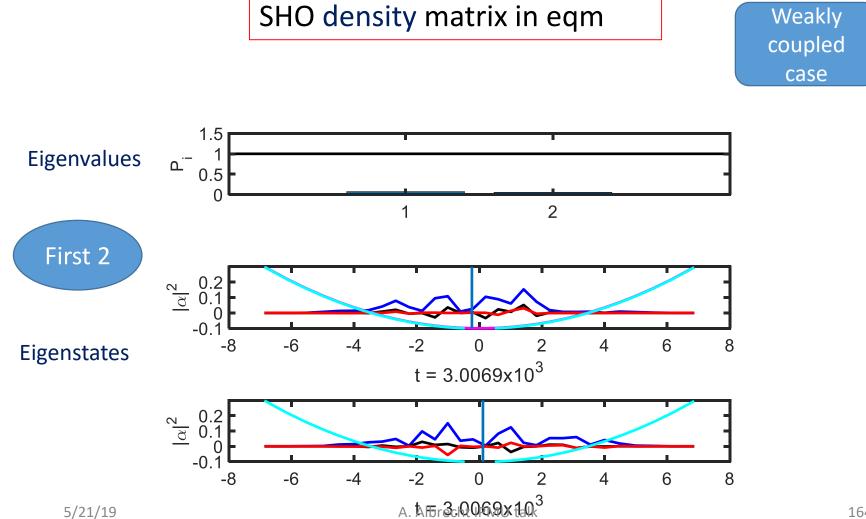


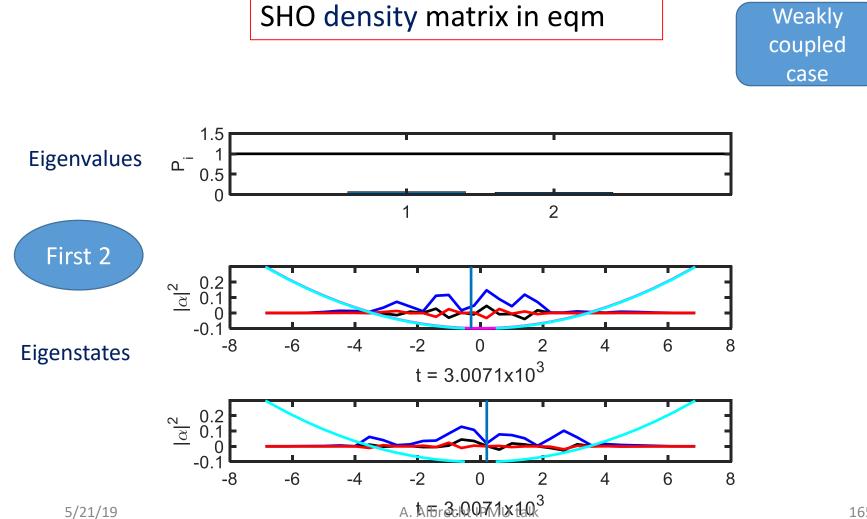


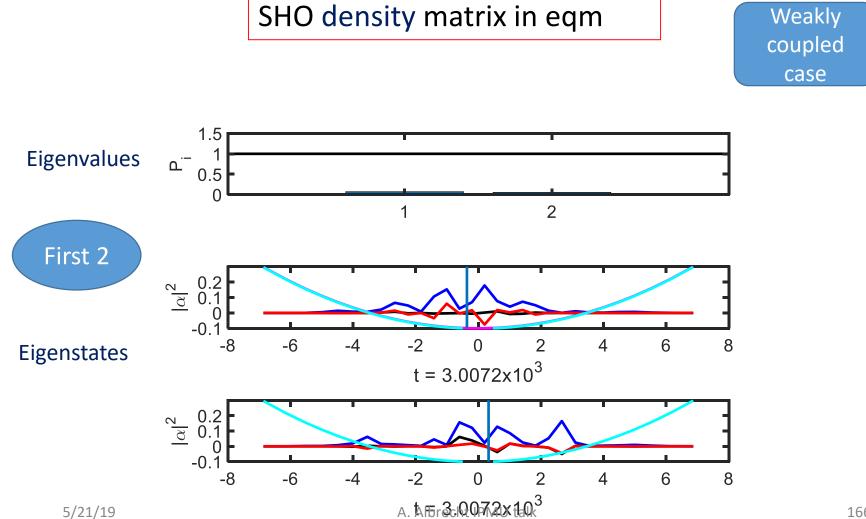


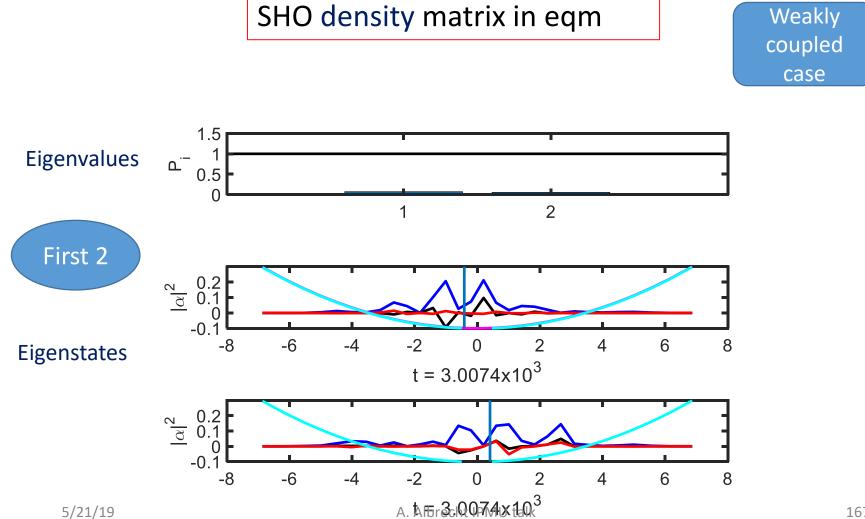


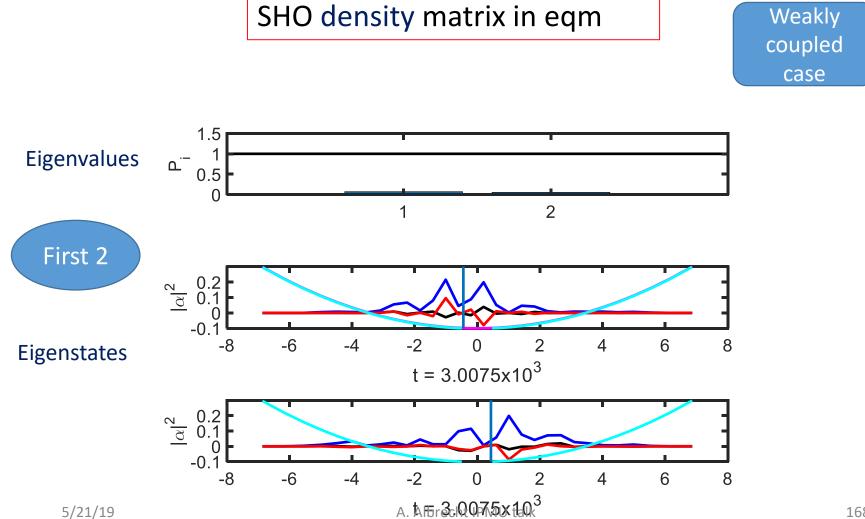


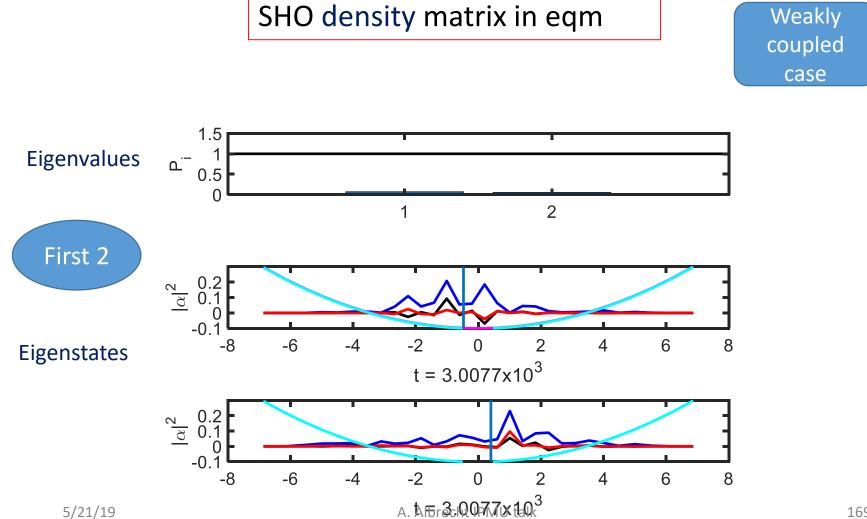


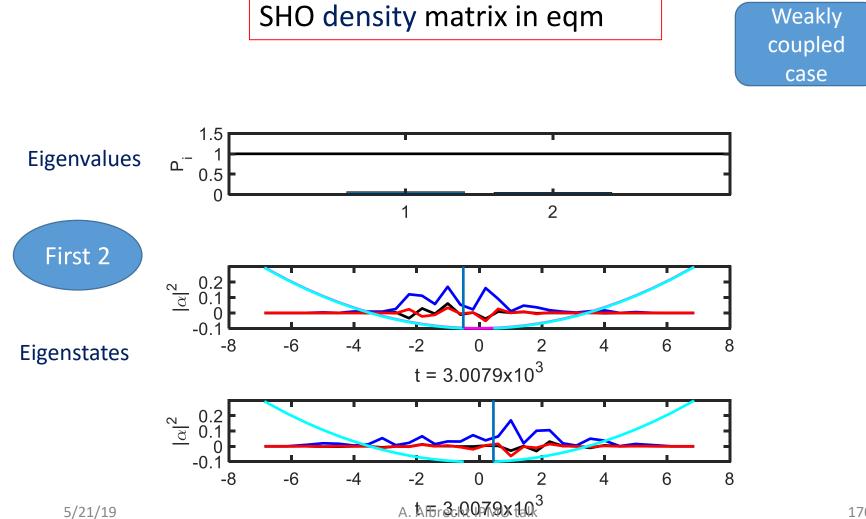


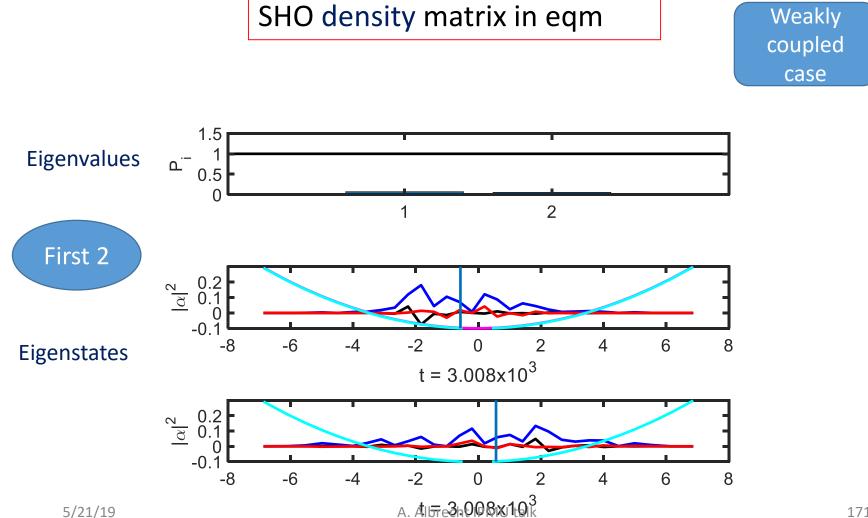




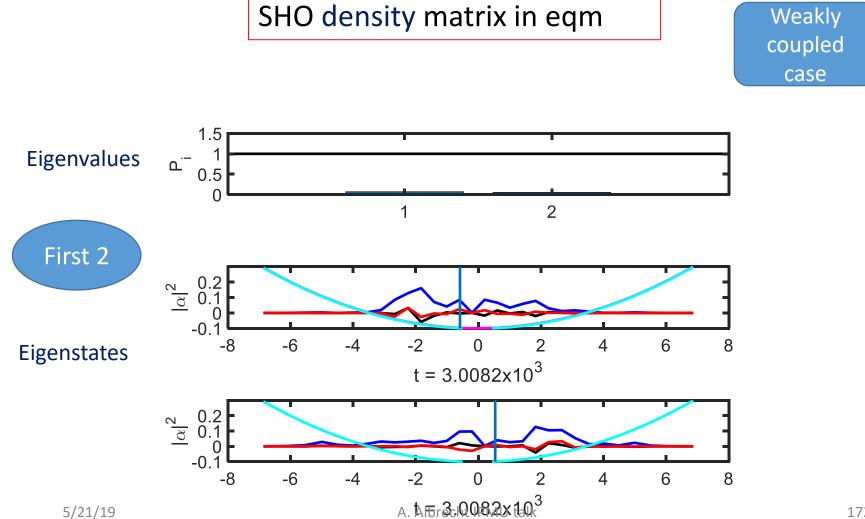


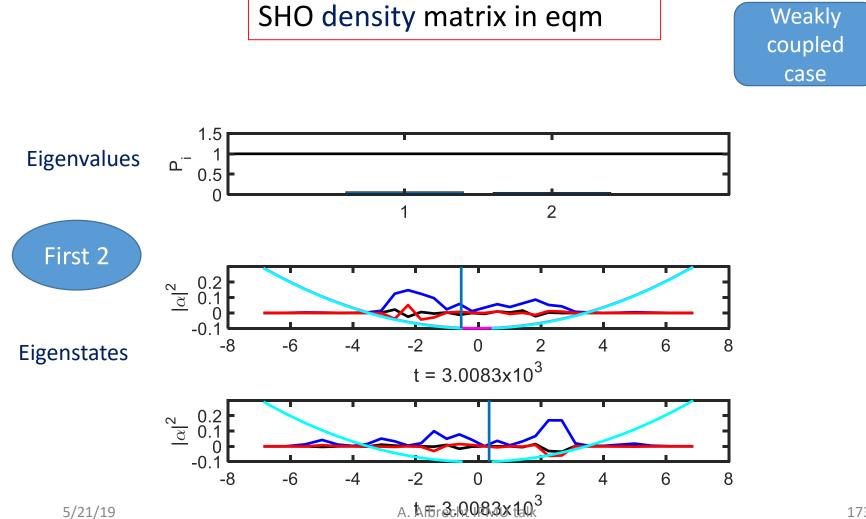


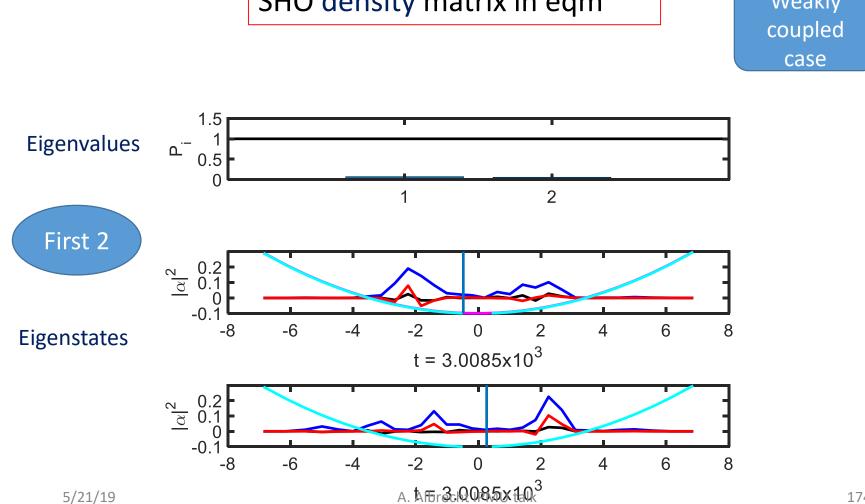




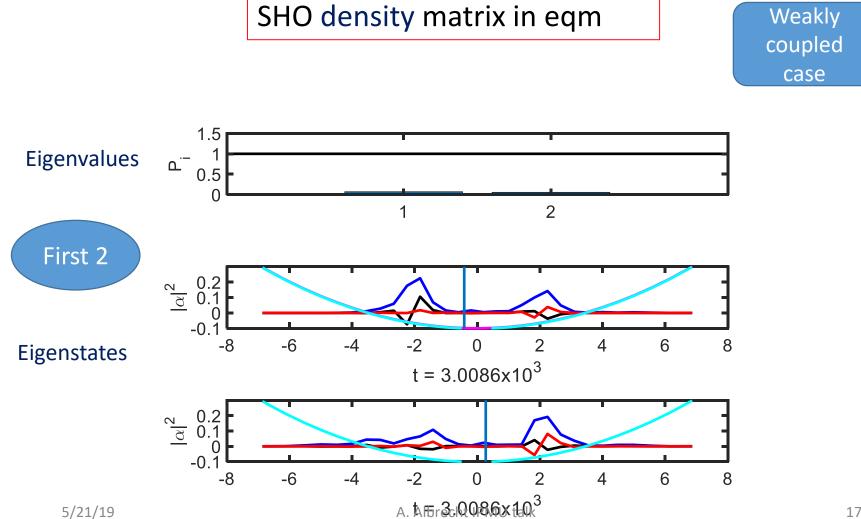
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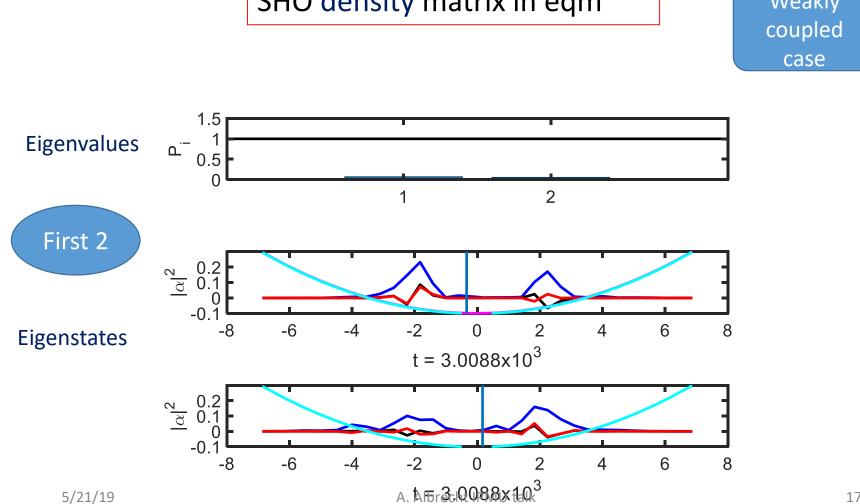


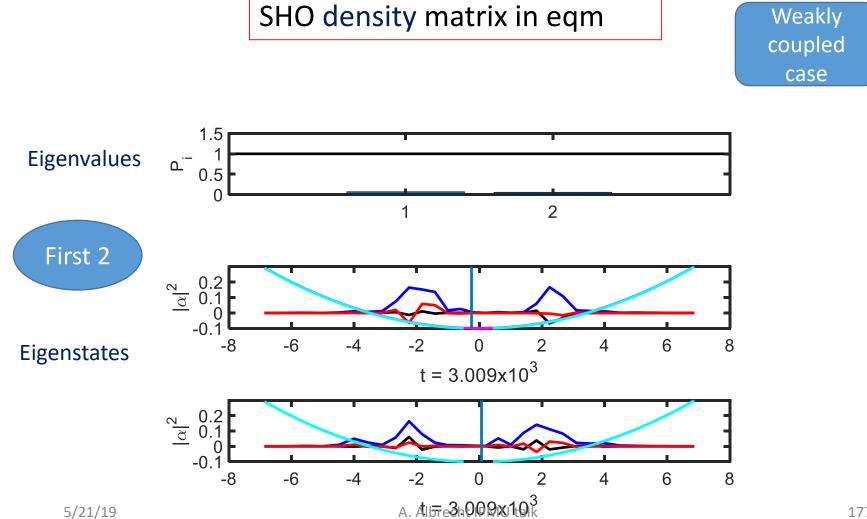


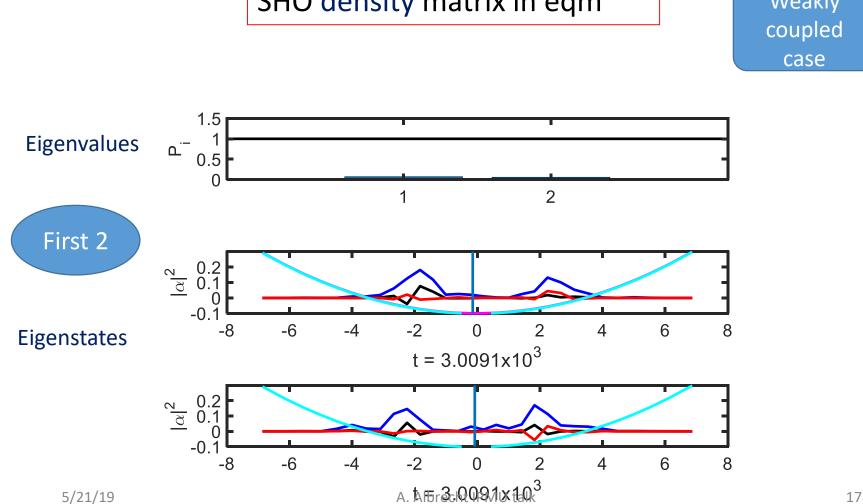


Weakly

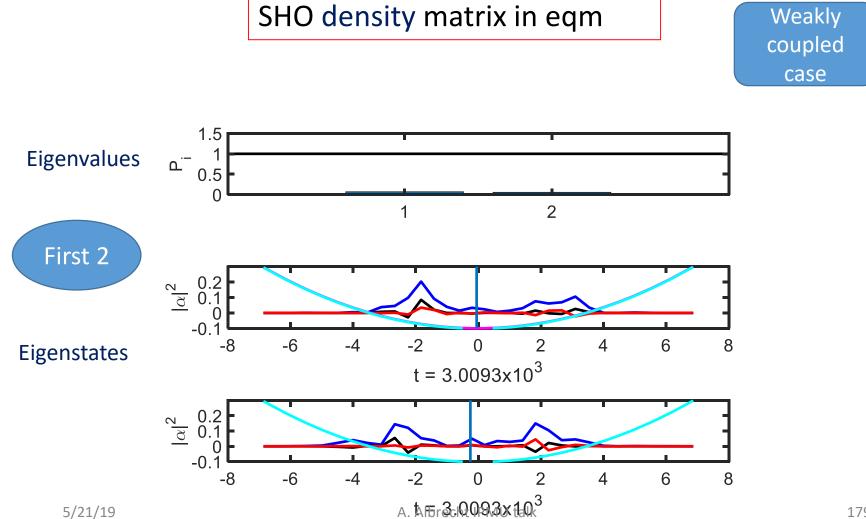


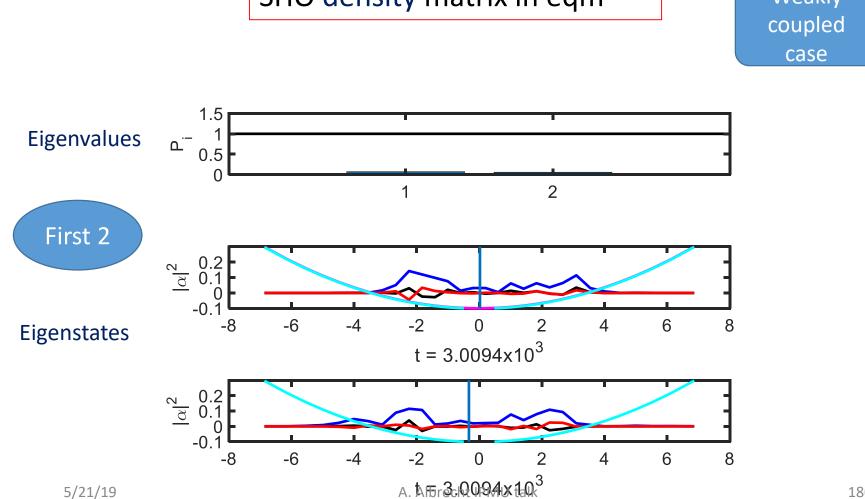






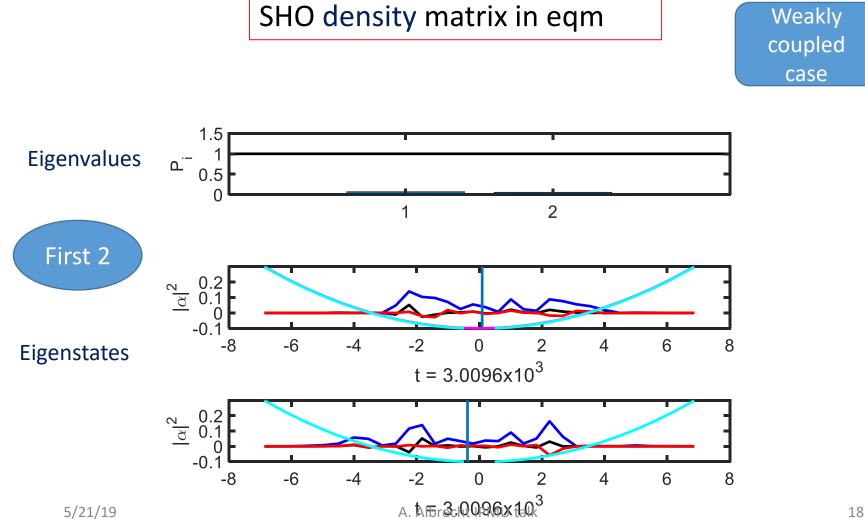
Weakly

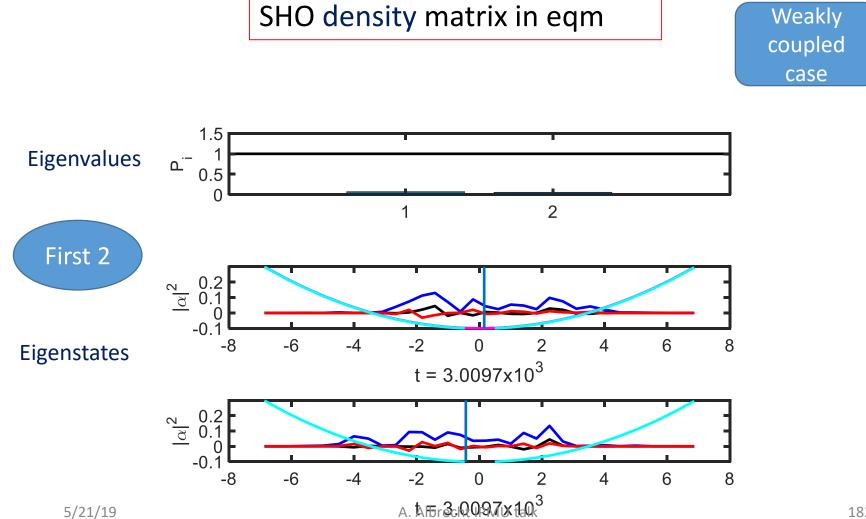


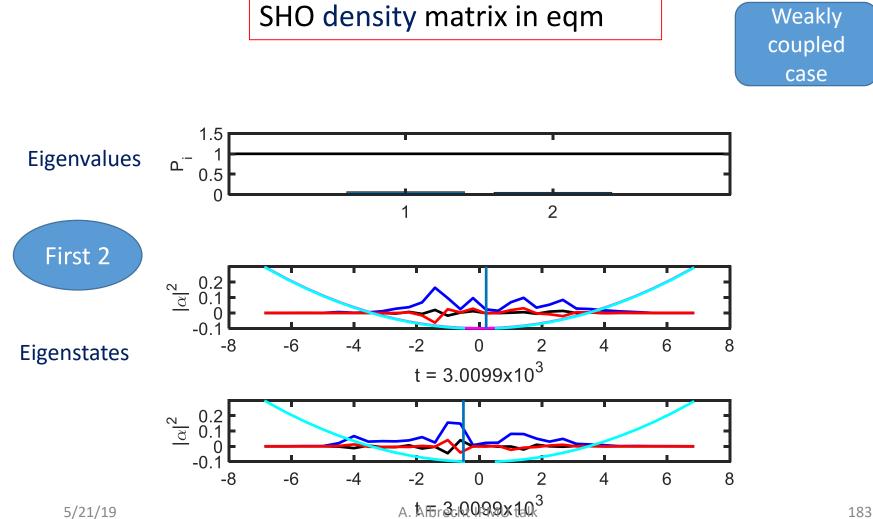


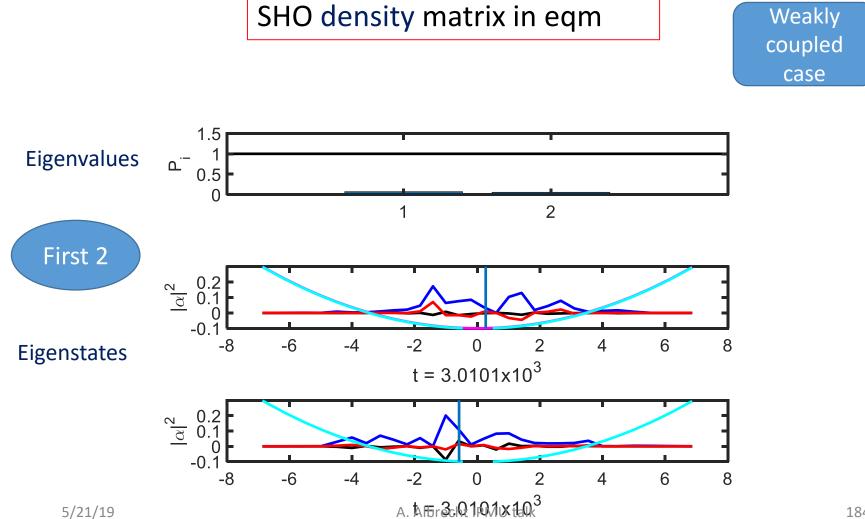
SHO density matrix in eqm

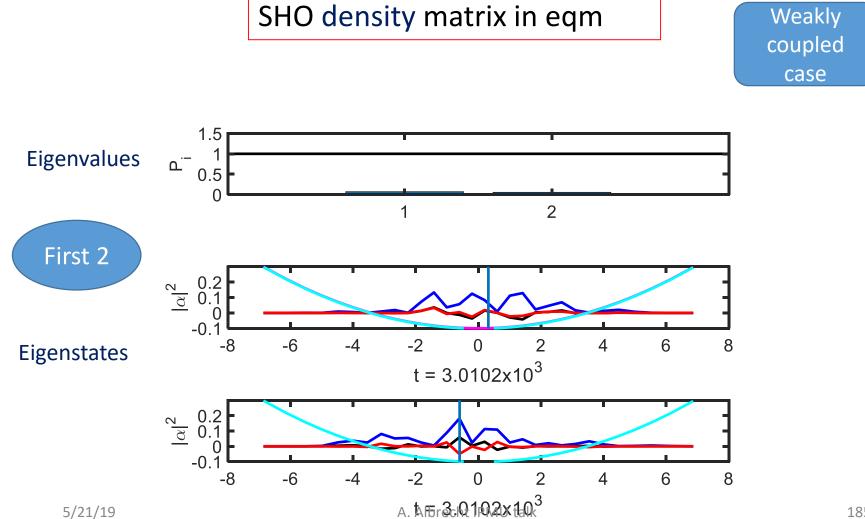
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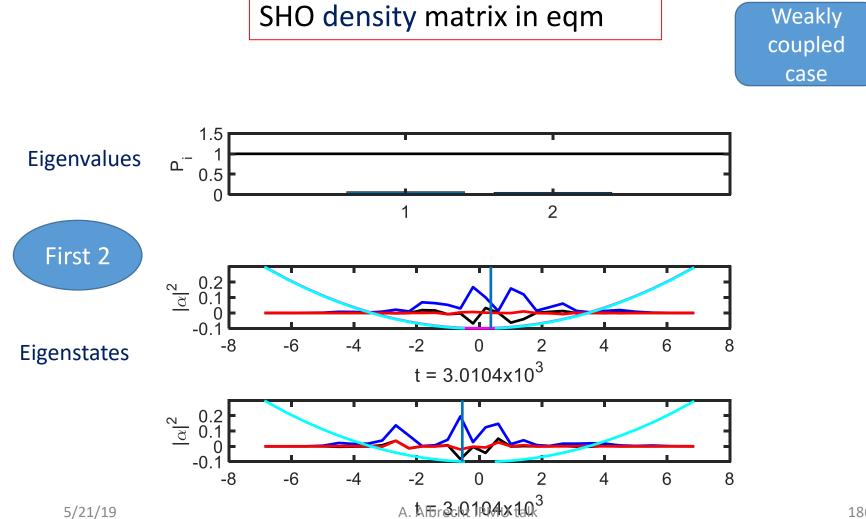


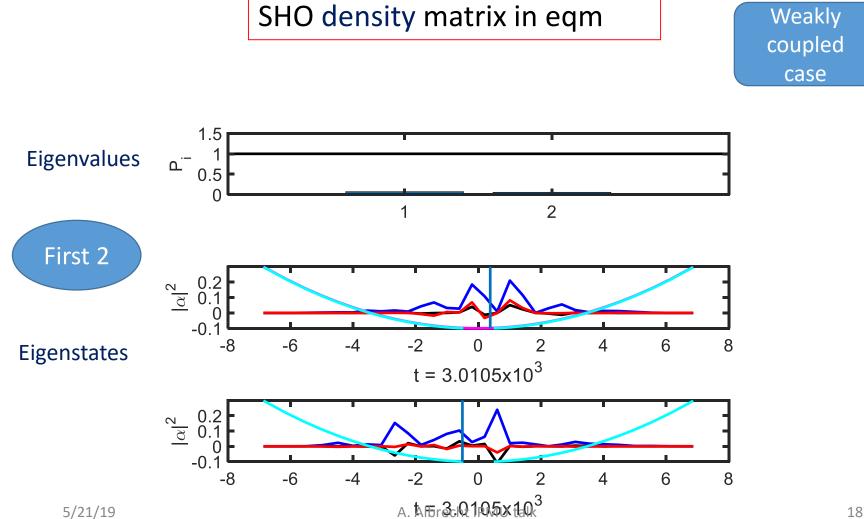


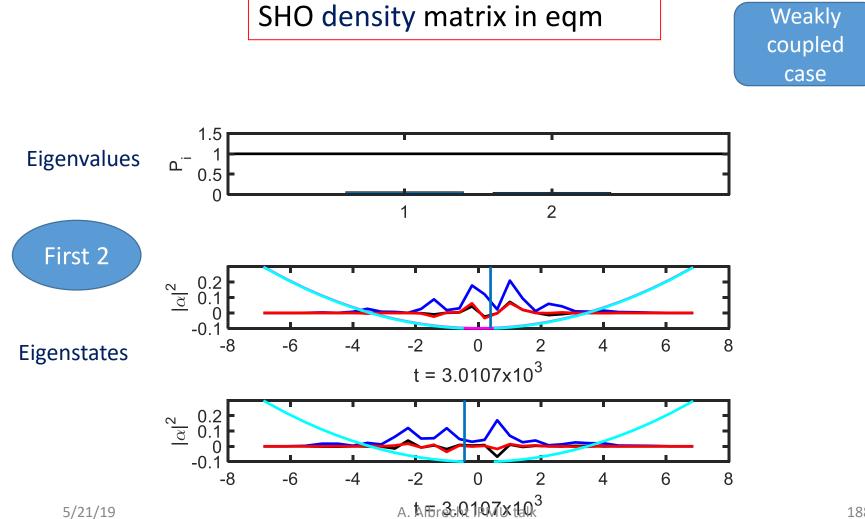


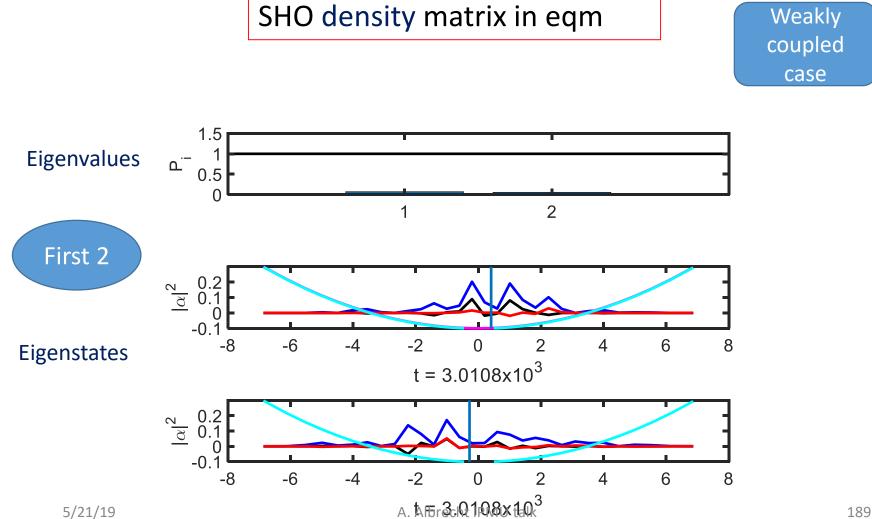


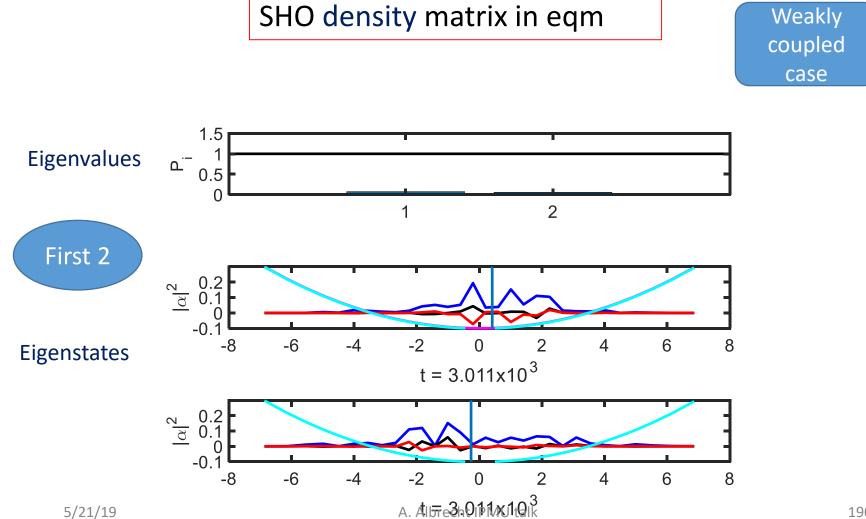


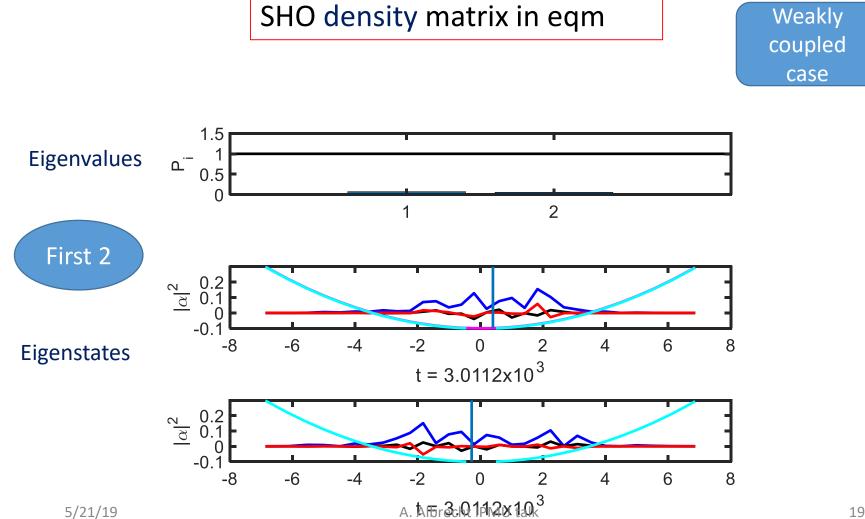


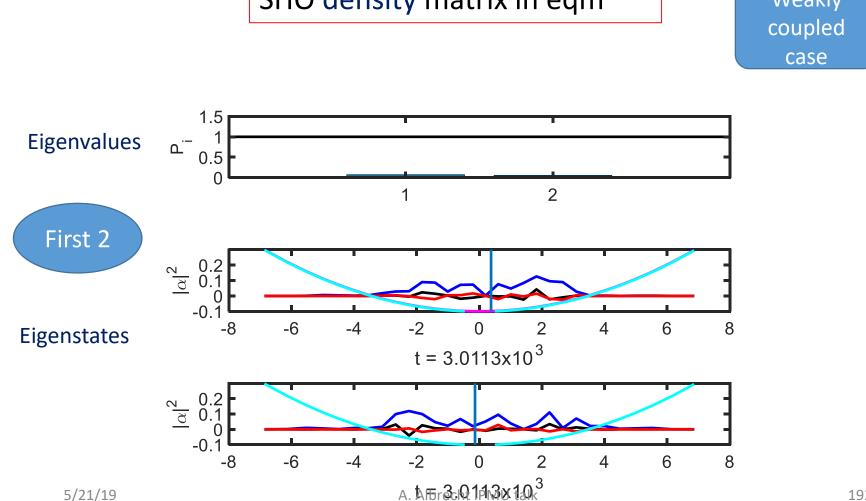






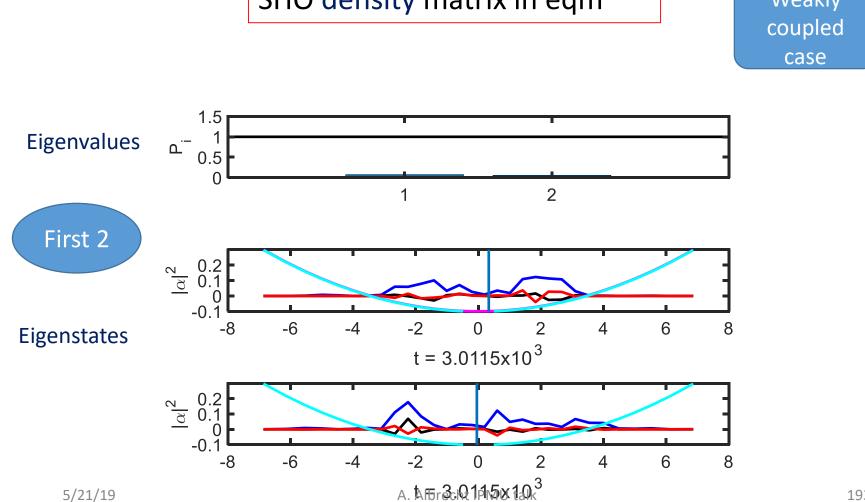




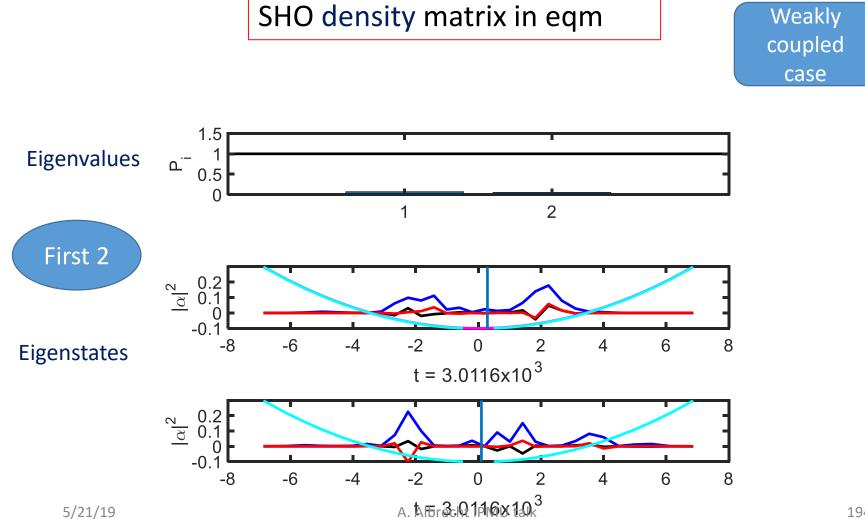


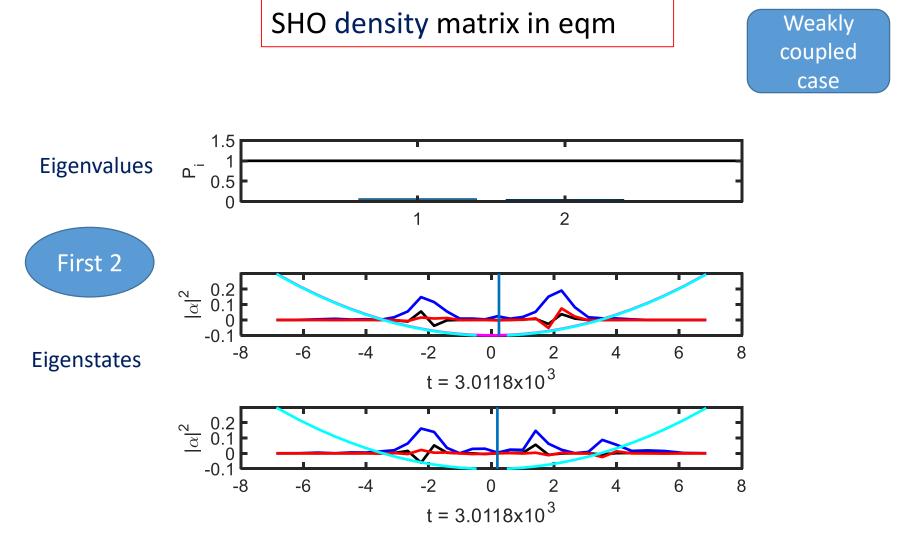
SHO density matrix in eqm

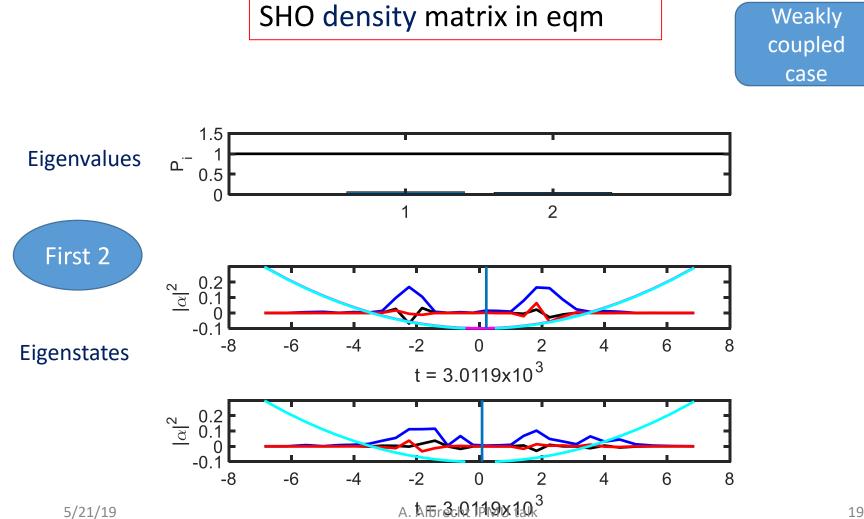
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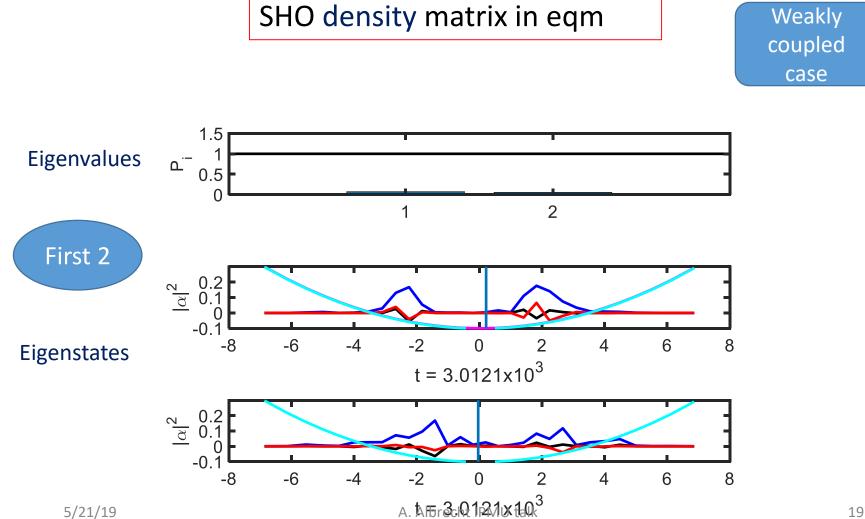


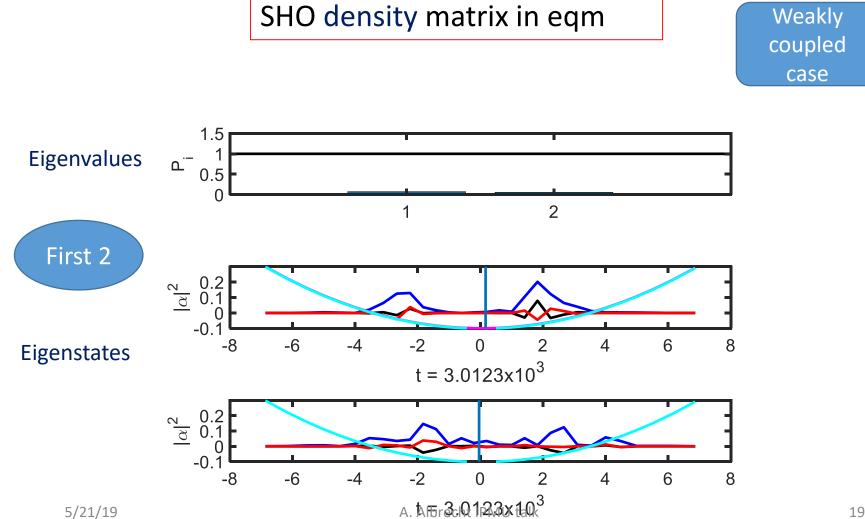
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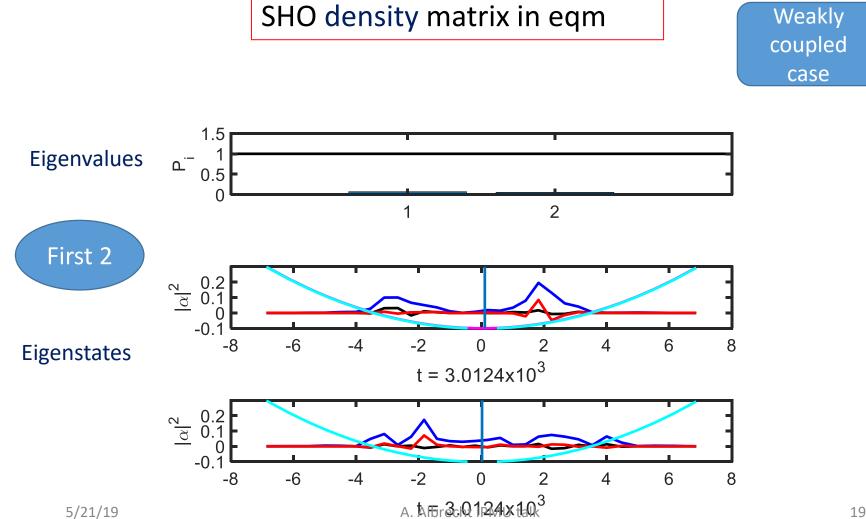


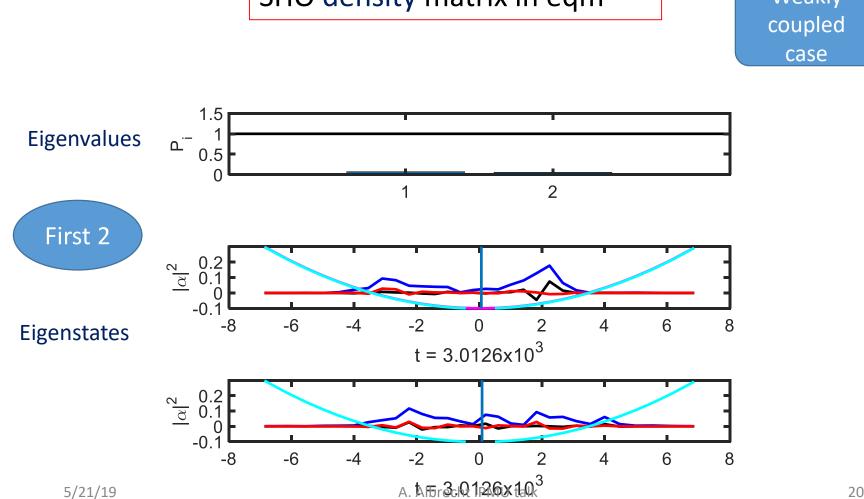




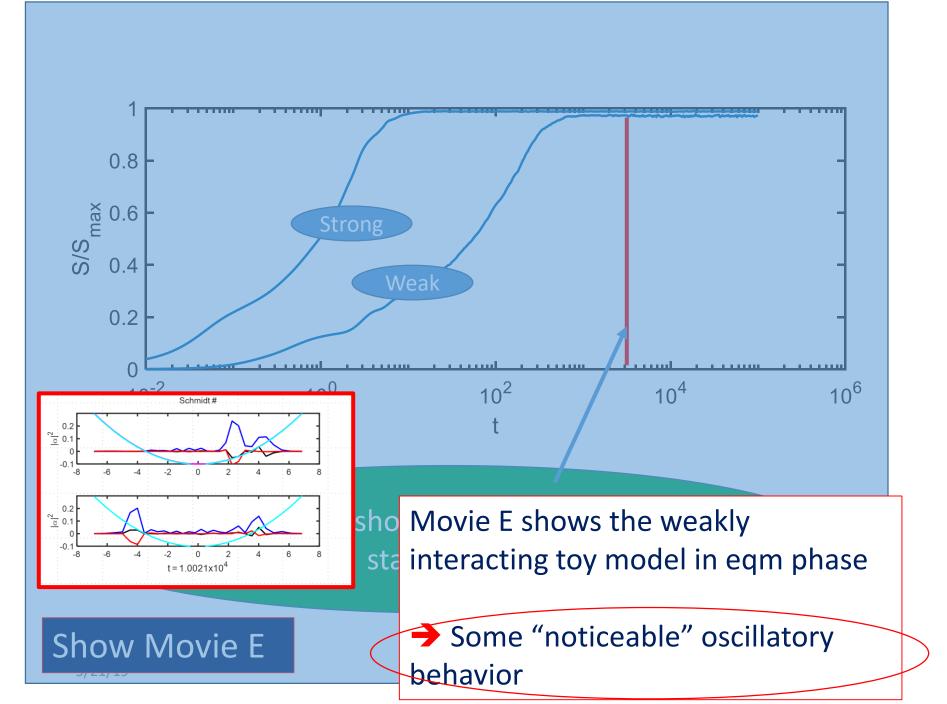


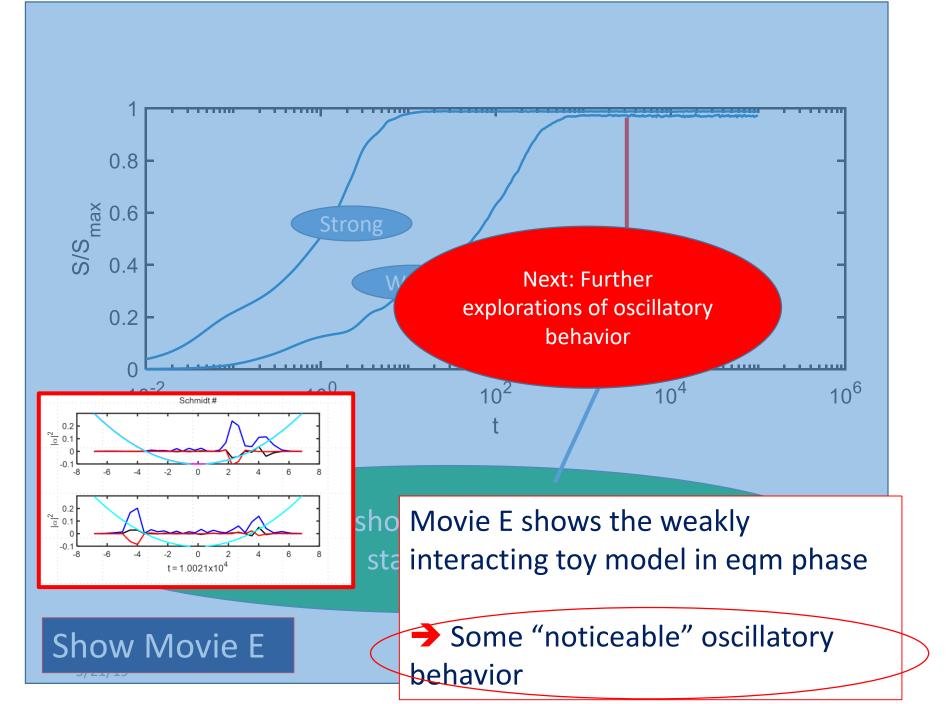




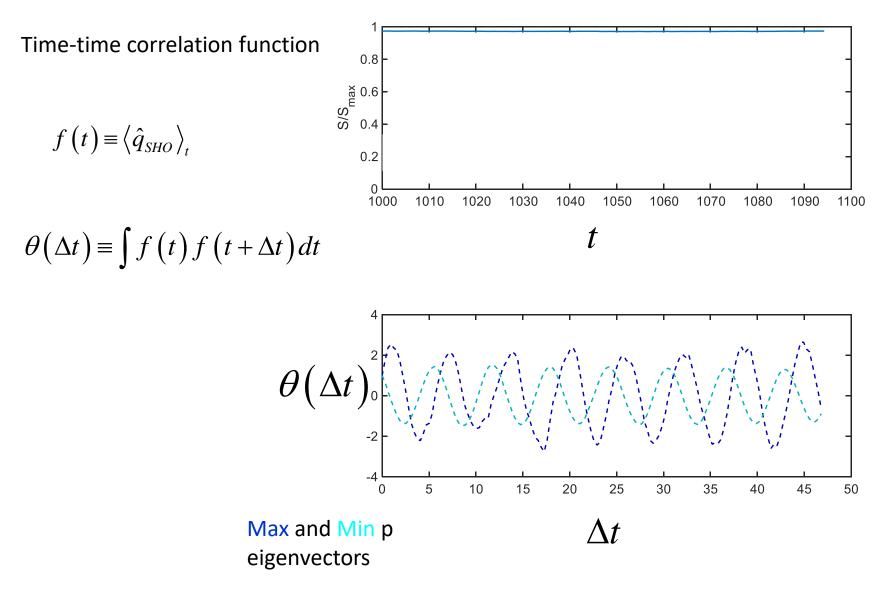


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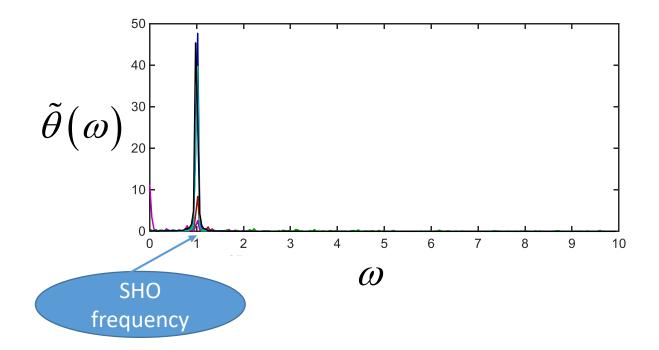


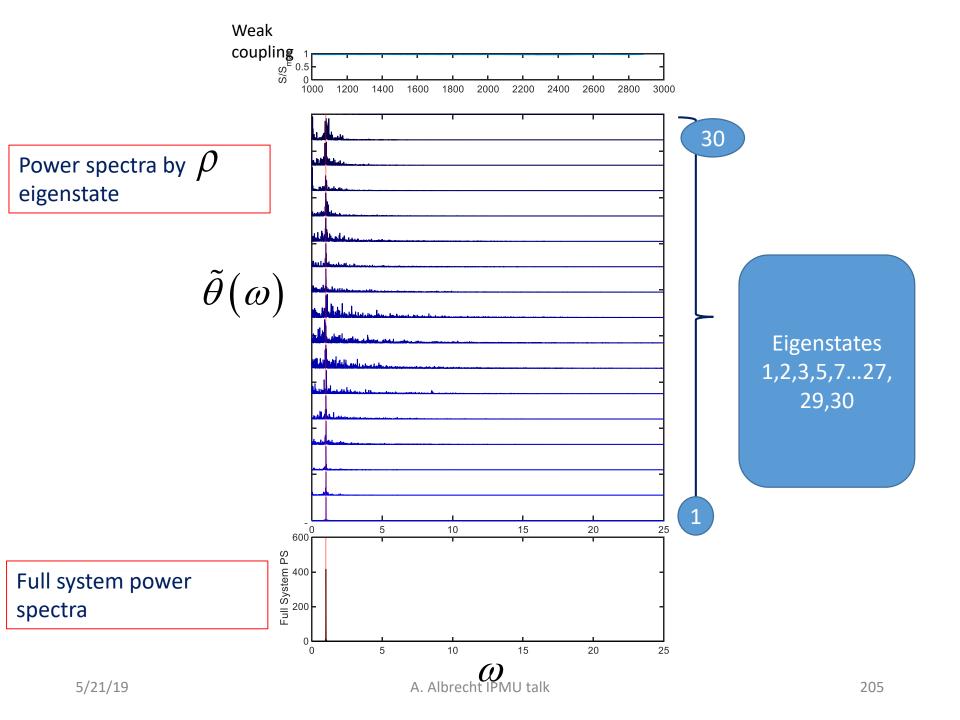


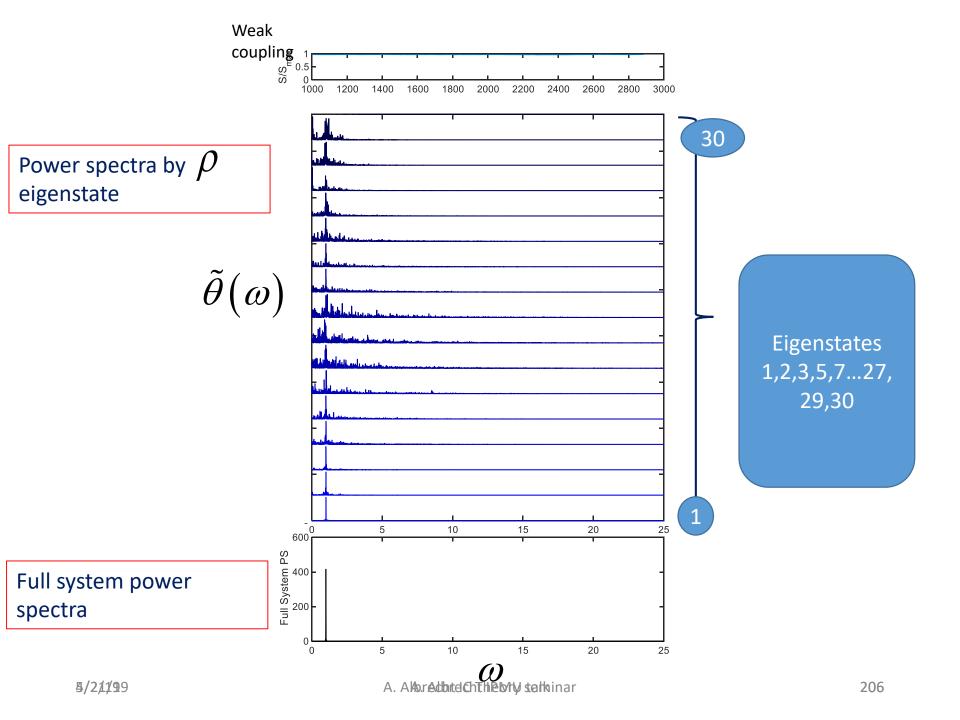
Further analysis:



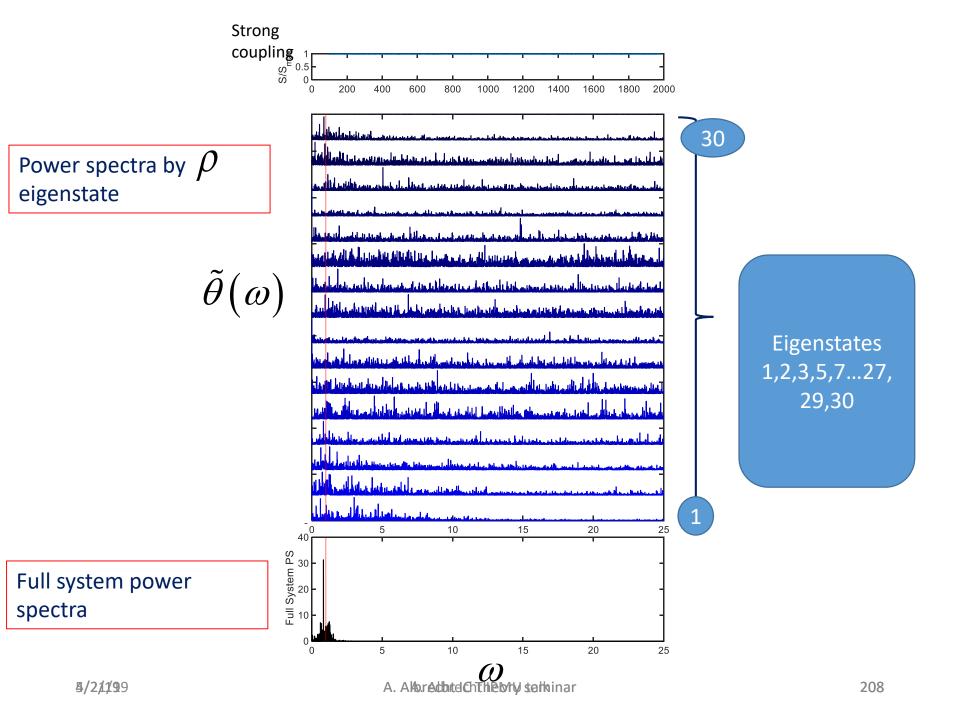
Power spectrum

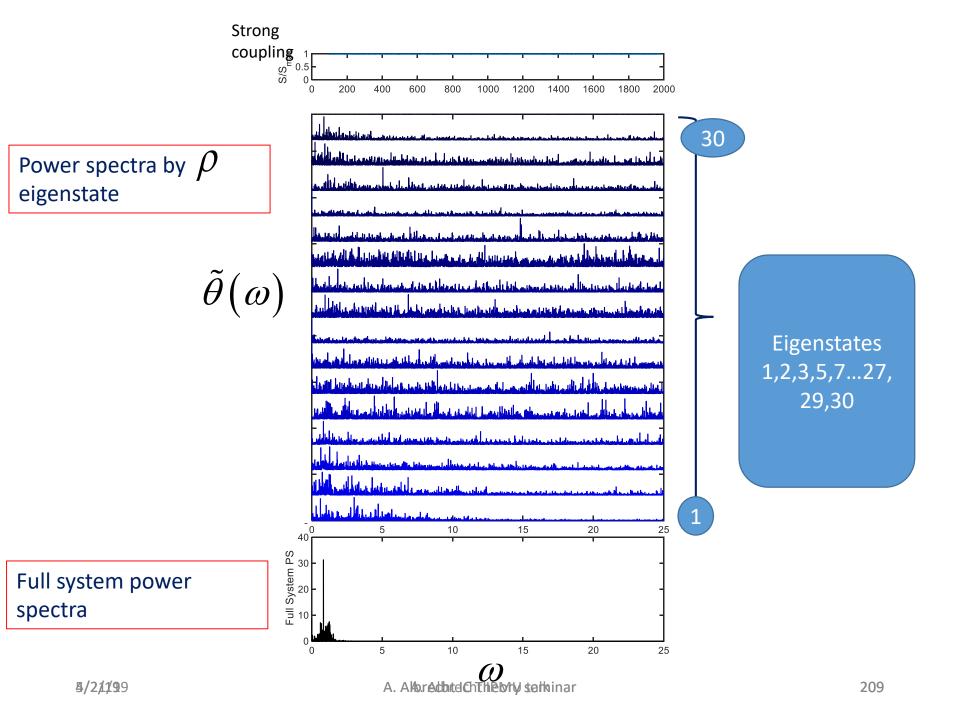


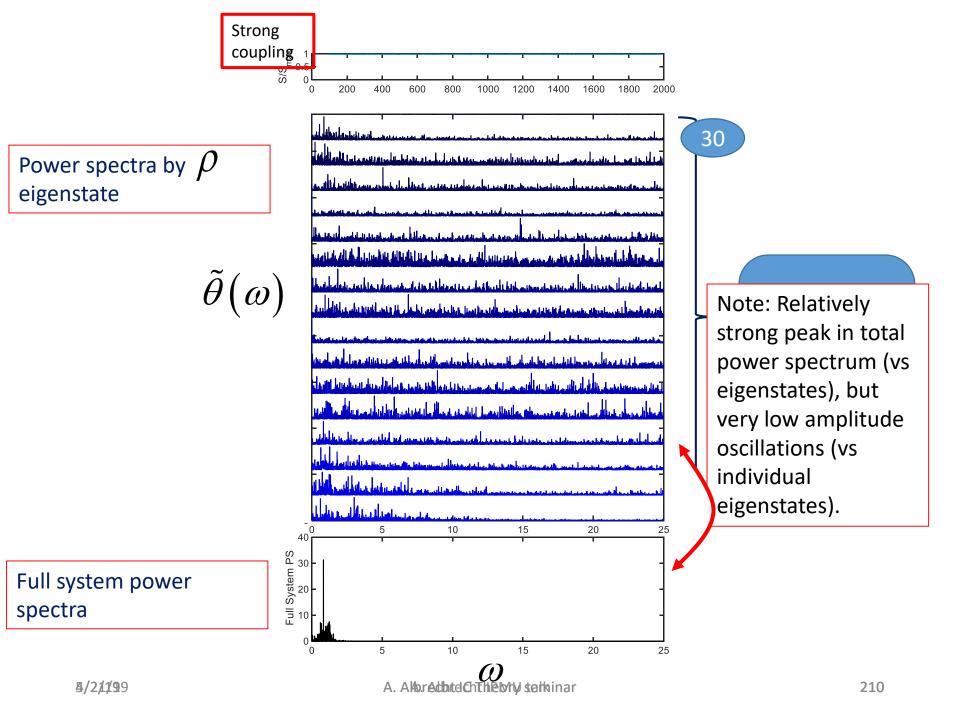




Now look at strong coupling, where Eqm seemed more noisy







Upshot:

- Strong oscillatory signal in <q> for ρ eigenstates for weakly coupled case (despite messy overall wavefunction shapes)
- → No such signal for strongly coupled case
- → Both cases show strong oscillatory signal for $\langle q\rho \rangle$ but amplitude is small.

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NEXT: A consistent histories approach

- → Generally, in the path integral formulation of QM interference among paths plays an important role
- → CH formalism identifies paths where interferences effects are NOT important. These are the paths to which probabilities can be assigned, and which are classical in that sense.
- → We have found that the messiness of the eqm physics of our toy model shows up as histories that degrade after a couple of SHO periods
- CH gives interesting test of "coherent state as most robust state" result from master equation work.

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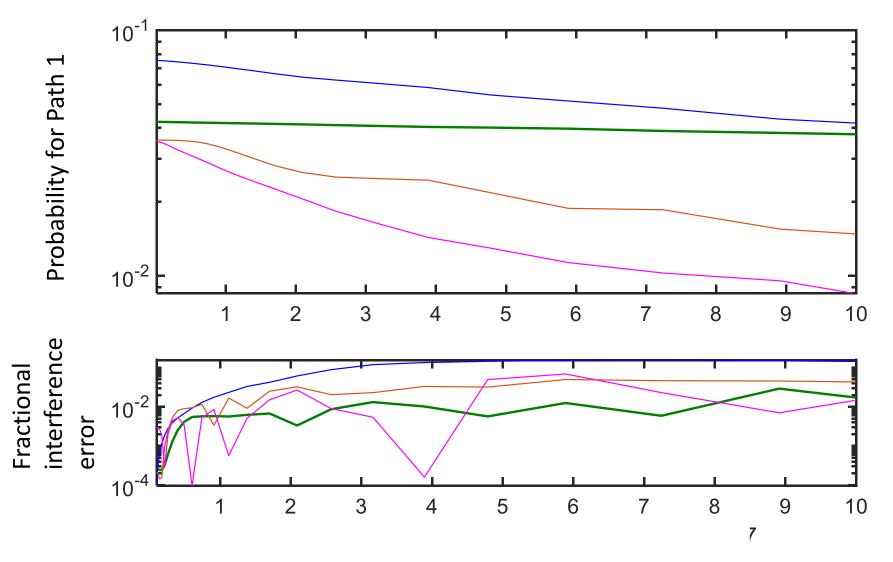
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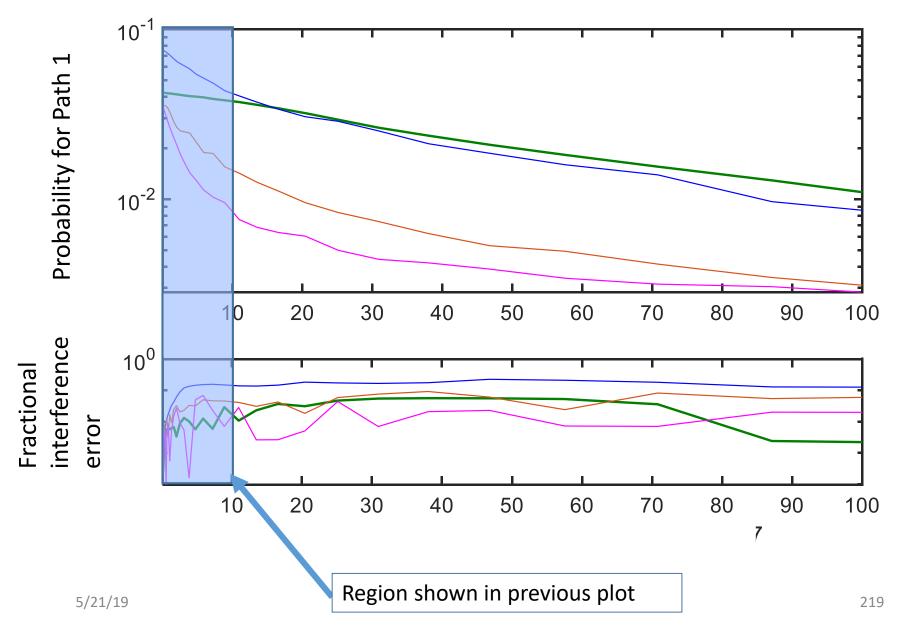
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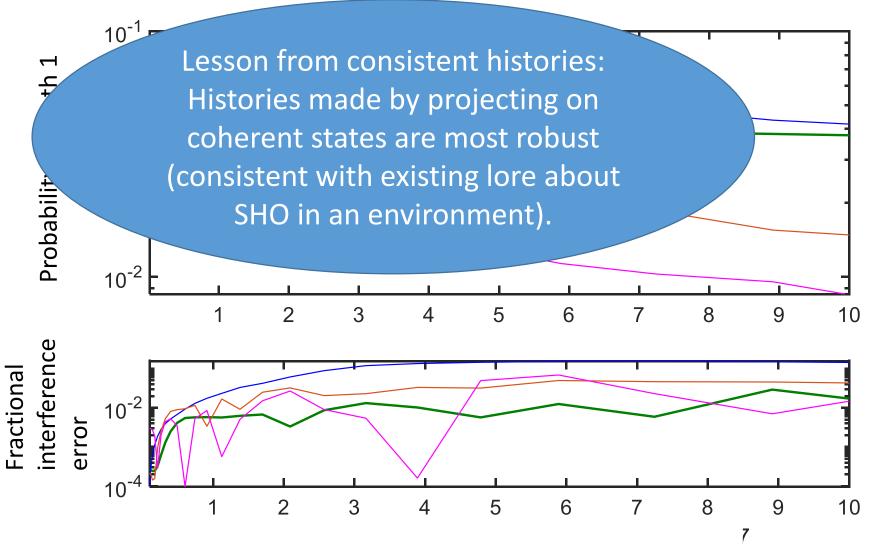
Histories built from coherent states (green) degrade more slowly than histories built from other states.



Histories built from coherent states (green) degrade more slowly than histories built from other states. <u>But eventually the coherent state paths degrade too.</u>



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Lesson from consistent histories: Histories made by projecting on coherent states are most robust (consistent with existing lore about SHO in an environment).

> But in EQM (detailed balance), even coherent state histories degenerate on the scale of a coherence time

error

10⁻¹

10⁻²

10⁻²

10

2

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Probabili

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Fractional

5

4

10

9

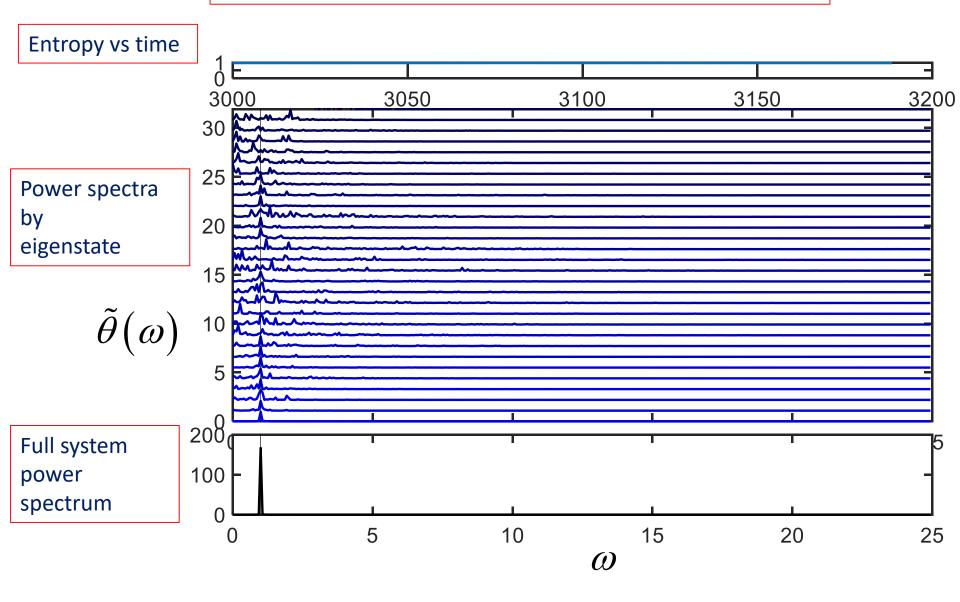
Lessons from eqm studies in the adapted CL model:

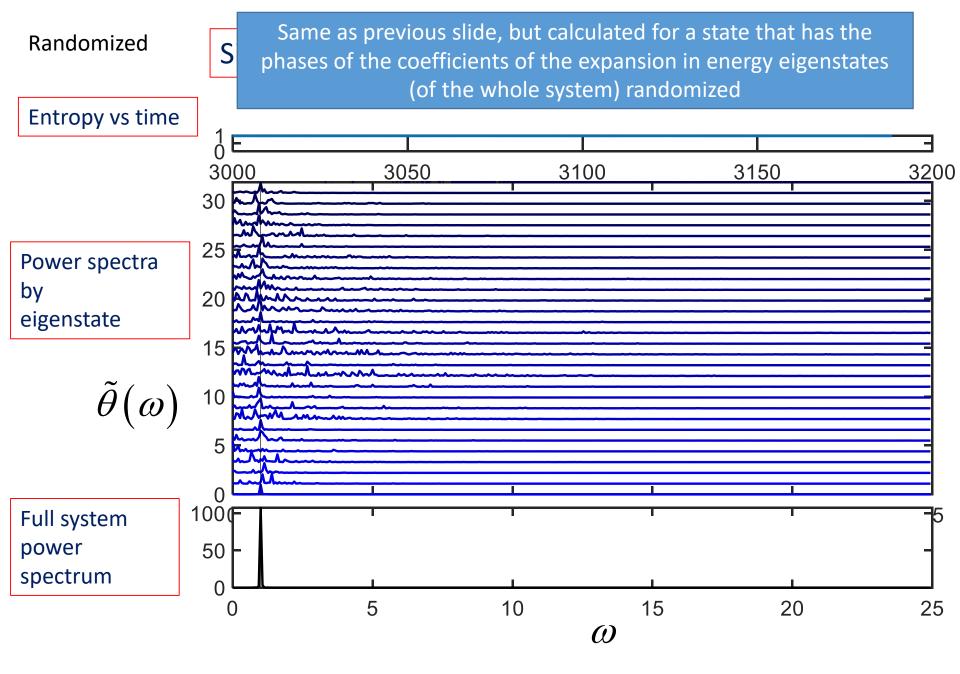
- Plenty of messiness in eqm ("wallowing or interfering Everett worlds")
- → Still, some intriguing sign of "classicality" show up in the power spectra of density matrix eigenstates.
- Consistent Histories formalism shows some classical behavior (which degrades after a couple of SHO periods)
- → Further discussion of implications at end of talk.

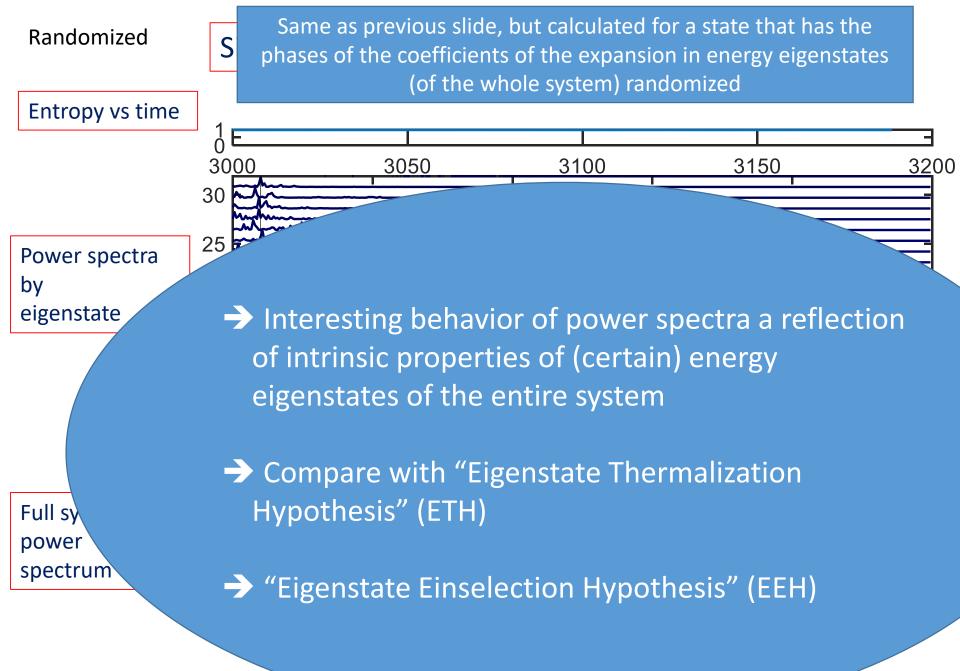
- 1. Motivations
- 2. Introduction to einselection and the toy model
- 3. Einselection in equilibrium (technical explorations and overall assessment)
- 4. Eigenstate Einselection Hypothesis (if there is time)

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Similar power spectra to those shown above







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- The ACL allows the exploration of these phenomena under conditions not accessible to the CL model (specifically eqm.)
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- A suggestion: Classical phenomena may persist in equilibrium on a time scale short compared to a decoherence time... interesting for cosmology.

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"classical" features emerged in eqm when the decoherence time was extended (weak coupling) to be longer than the SHO period Compare with "de Sitter Equilibrium" scenario where decoherence time is the age of the Universe