

Dark Matter Heating vs. Rotochemical Heating in Old Neutron Stars

Koichi Hamaguchi (University of Tokyo)

Berkeley week at Kavli IPMU, January 16, 2020.

Based on

KH, N. Nagata, K. Yanagi, [[arXiv:1905.02991](#)] Phys.Lett. B795 (2019) 484

K. Yanagi, N. Nagata, KH, [[arXiv:1904.04667](#)] MNRAS (to be published)

cf. KH, N. Nagata, K. Yanagi, J. Zheng [[arXiv:1806.07151](#)], Phys.Rev. **D98** (2018) 103015. -> bound on **axion**.

Dark Matter Heating

vs.

Rotochemical Heating

in Old Neutron Stars

See also the talk by Keisuke Yanagi tomorrow.

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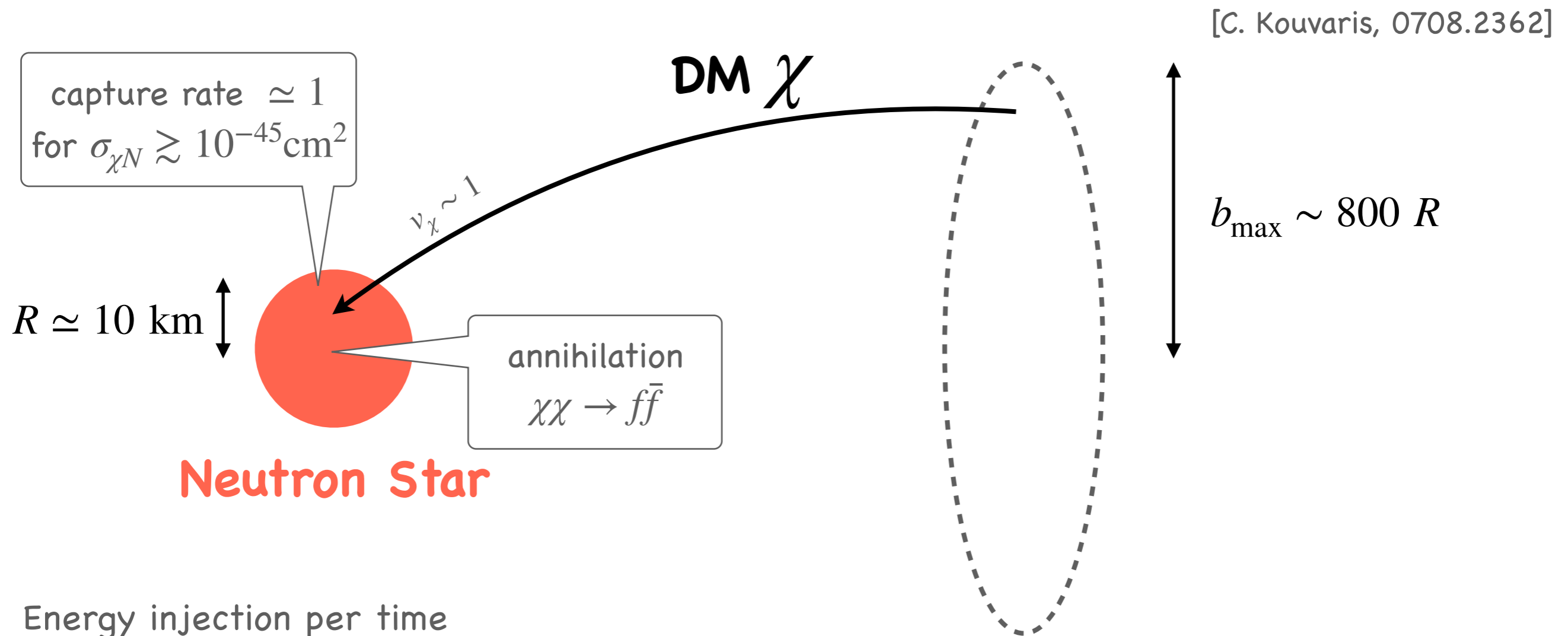
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Introduction (outline of today's talk)

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1. WIMP DMs hit a neutron star (NS), and annihilate inside the NS.



Energy injection per time

$$L_{\text{WIMP} \rightarrow \text{NS}} \sim \pi b_{\text{max}}^2 \rho_\chi v_\chi \simeq 3 \times 10^{22} \text{erg/s}$$

(independent of DM mass)

→ **DM heats NS !**

Introduction (outline of today's talk)

1. WIMP DMs hit a neutron star (NS), and annihilate inside the NS.
2. Old and warm ($\sim 2000K$) NS = DM signal?!

C. Kouvaris 0708.2362,
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...

+ many recent works: e.g.,
J. Bramante+ 1703.04043
M. Baryakhtar+ 1704.01577
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C.-S. Chen+ 1804.03409
N. F. Bell+ 1807.02840
D. A. Camargo+ 1901.05474
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KH, N.Nagata, K.Yanagi 1905.02991
R. Garani+ 1906.10145
J. F. Acevedo+, 1911.06334
A. Joglekar+, 1911.13293

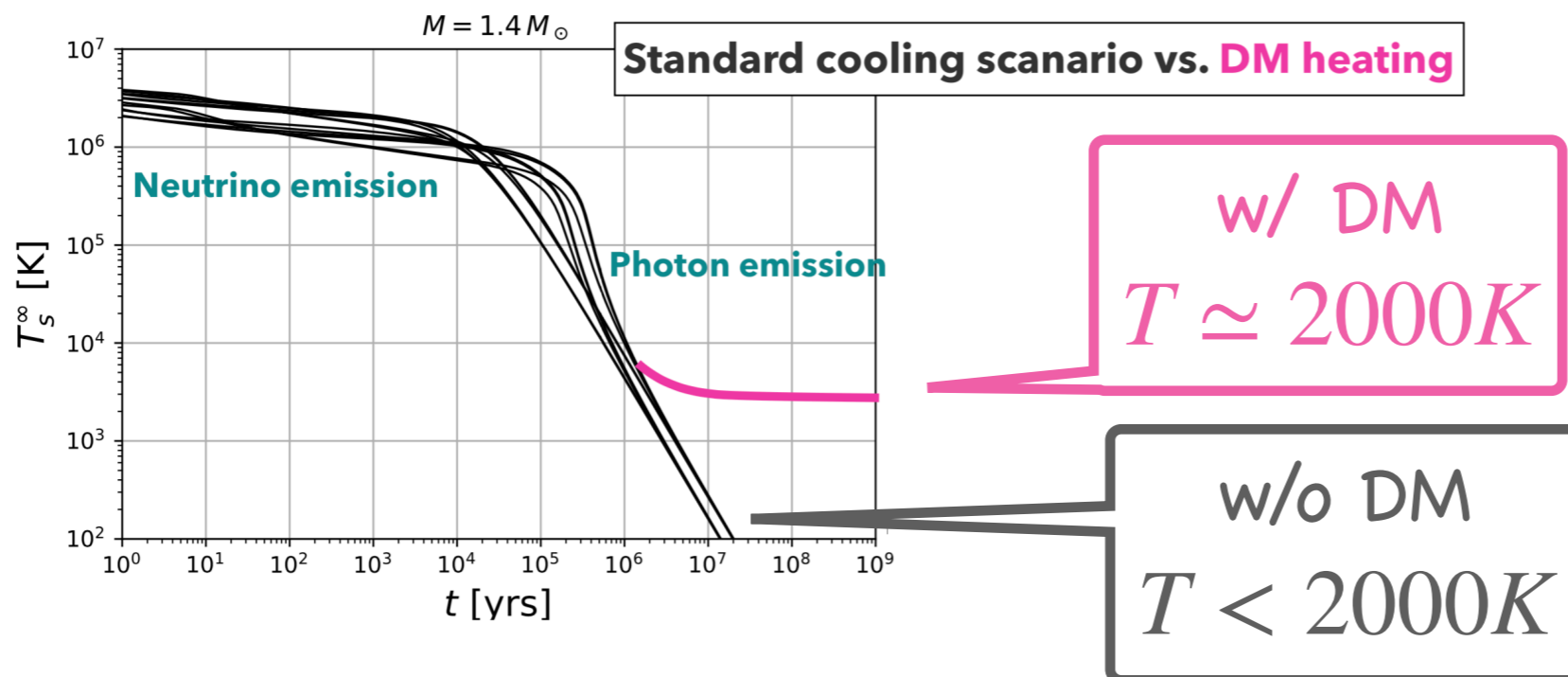
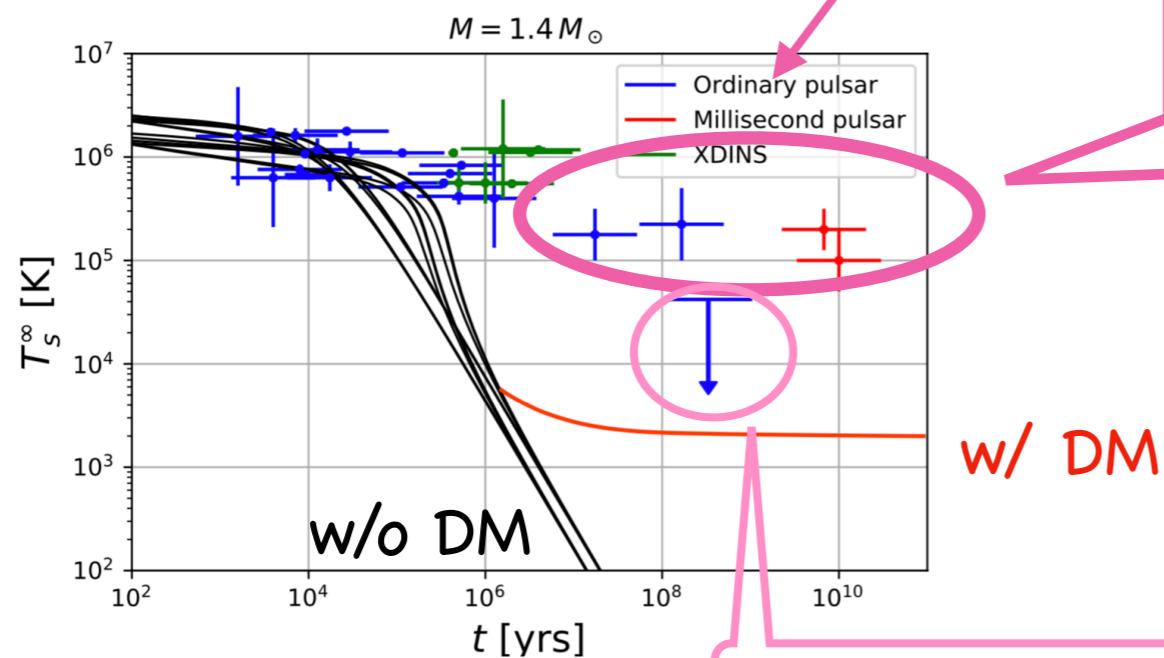


Fig. thanks to K.Yanagi.

Introduction (outline of today's talk)

1. WIMP DMs hit a neutron star (NS), and annihilate inside the NS.
2. Old and warm ($\sim 2000K$) NS = DM signal?!

3. But... old and warmer ($T \gg 2000K$) NSs are already observed!



Neither DM nor standard NS cooling can explain those old and warm NSs.

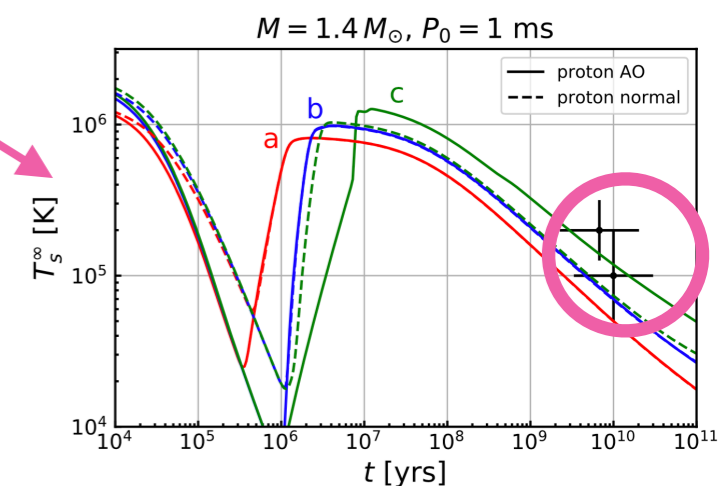
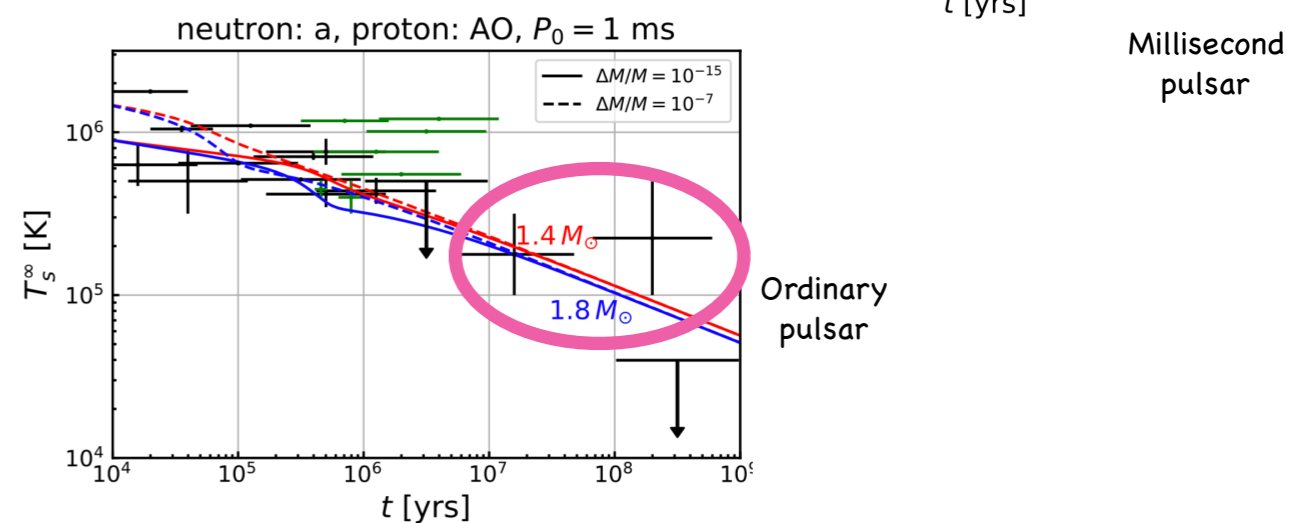
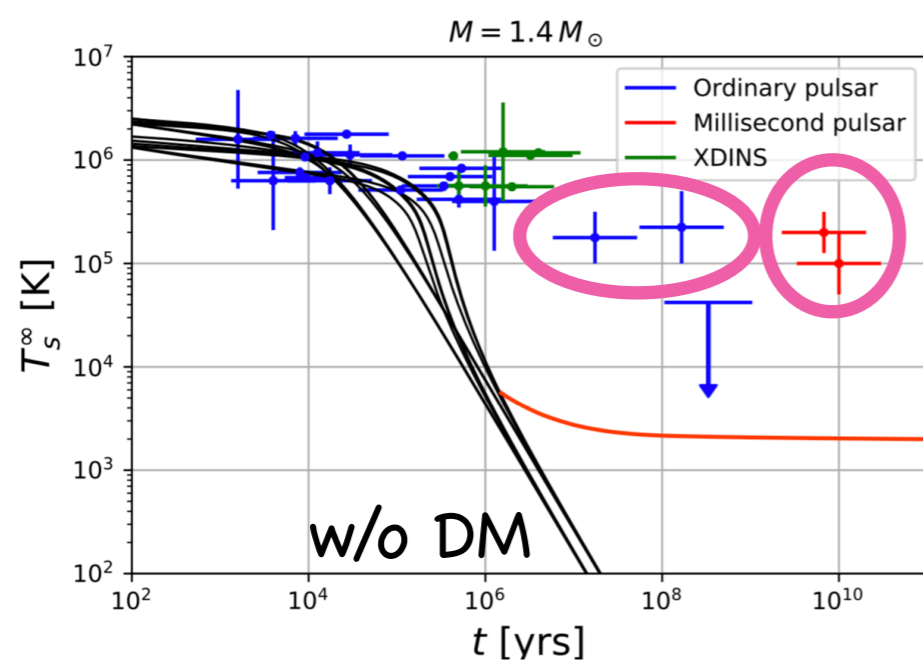
There are also cooler ones.

Introduction (outline of today's talk)

1. WIMP DMs hit a neutron star (NS), and annihilate inside the NS.
2. Old and warm ($\sim 2000K$) NS = DM signal?!

3. But... old and warmer ($T \gg 2000K$) NSs are already observed!

In fact, a mechanism inherent in NS ("rotochemical heating") can explain them.



Reisengger,'94, Haensel,'92, Gourgoulhon, Haensel,'93,
Fernandez, Reisenegger,'05,..... Yanagi, Nagata, KH, 1904.04667

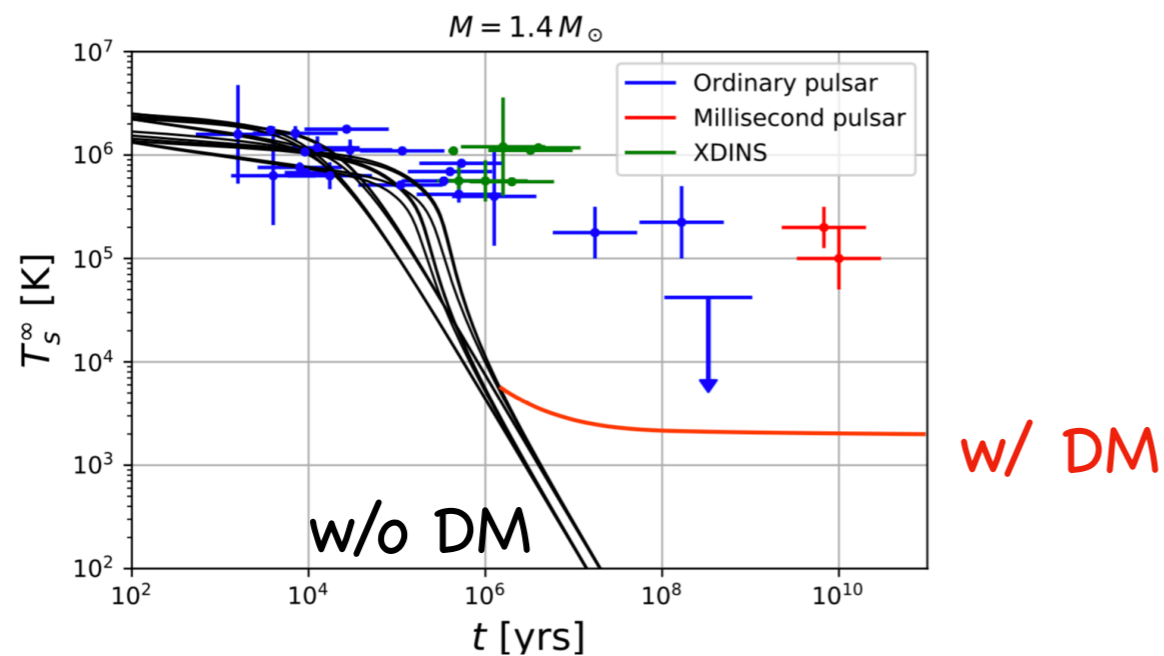
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2. Old and warm ($\sim 2000K$) NS = DM signal?!
3. **But...** old and warmer ($T \gg 2000K$) NSs are already observed!

In fact, a mechanism inherent in NS ("rotochemical heating") can explain them.

4. Question:

Can we really see the signal of the DM heating?
If so, what is the condition for that?



Plan

0. Introduction  done

1. Neutron Star

2. Neutron Star Cooling and Heating

3. Summary

Neutron Star

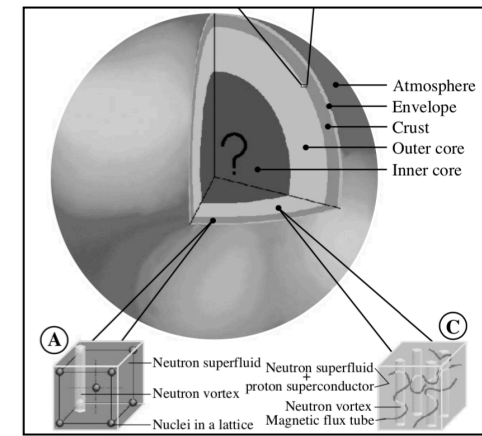


fig. from 1302.6626

- **Mass** : $M \sim (1 - 2)M_{\odot}$ (M_{\odot} = solar mass)

heaviest one found recently: $M \simeq 2.14M_{\odot}$ (pulsar MSP J0740+6620 [Nature, 2019 September])

Neutron Star

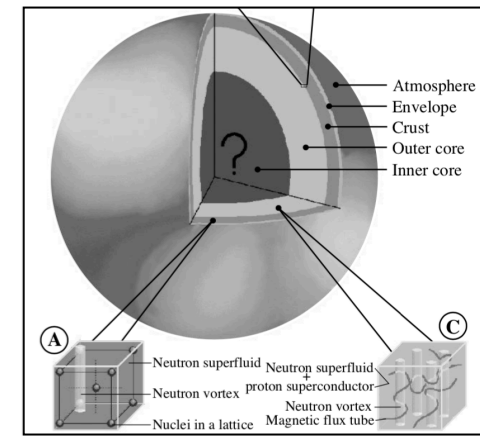
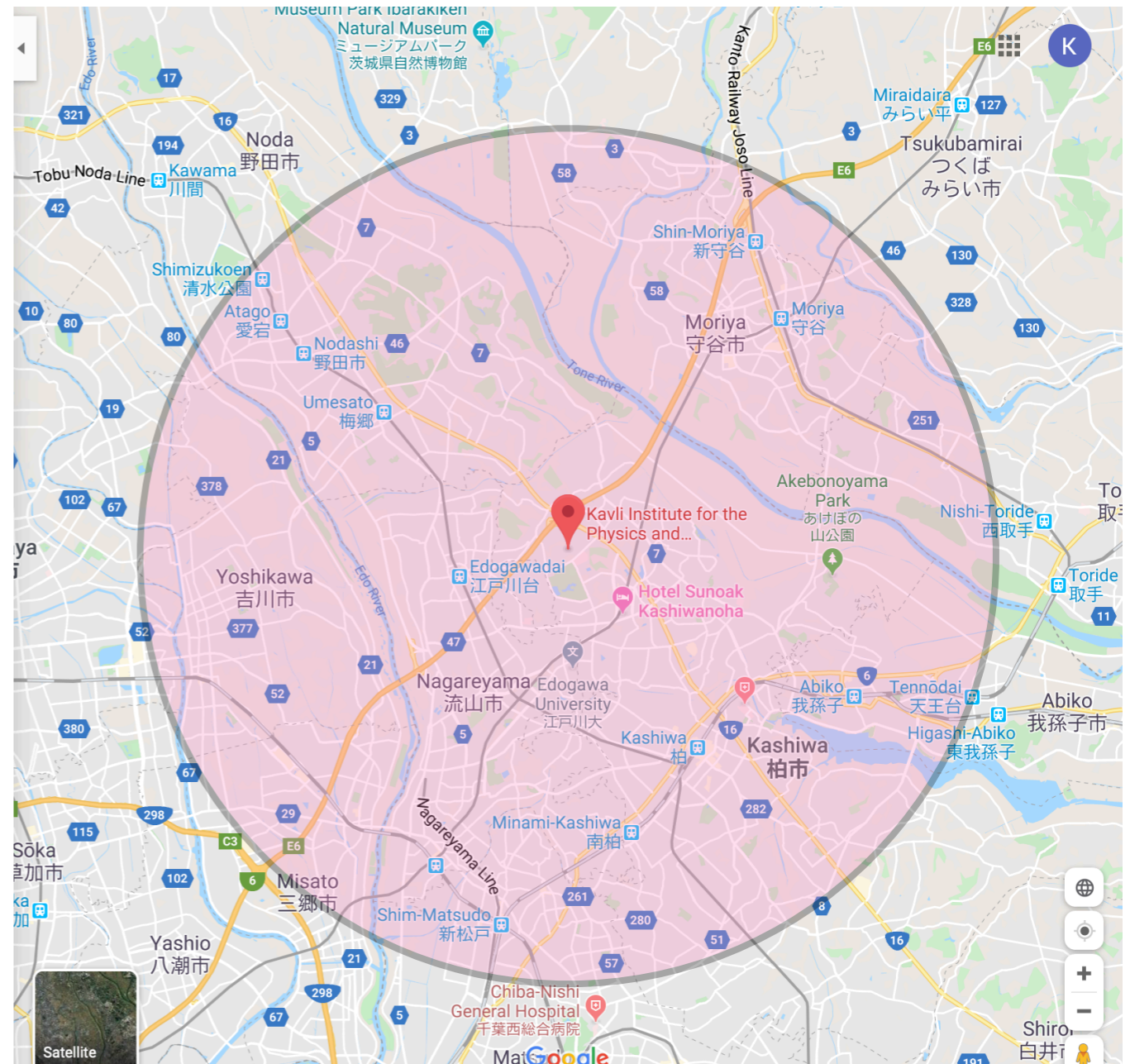


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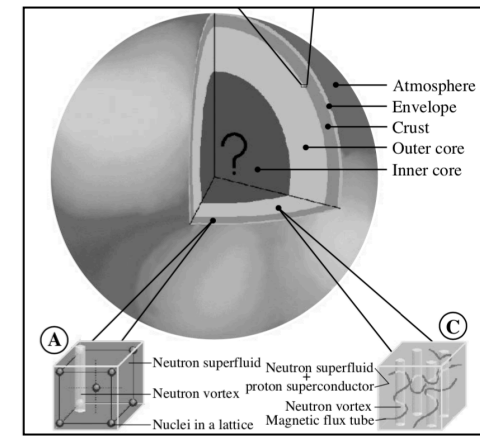


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- **Density** : $\bar{\rho} = \frac{M}{(4\pi/3)R^3} \simeq 7 \times 10^{14} \text{g/cm}^3$

cf. nuclear density $\sim 3 \times 10^{14} \text{g/cm}^3$

Neutron Star

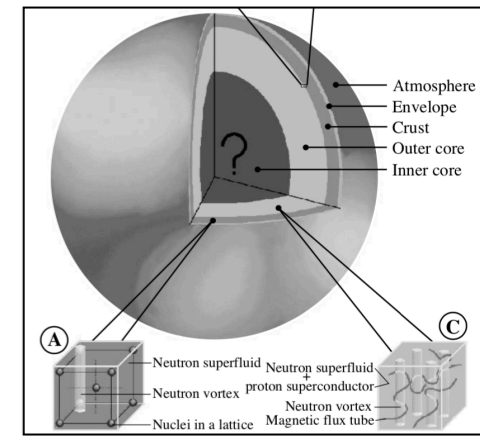


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Neutron Star

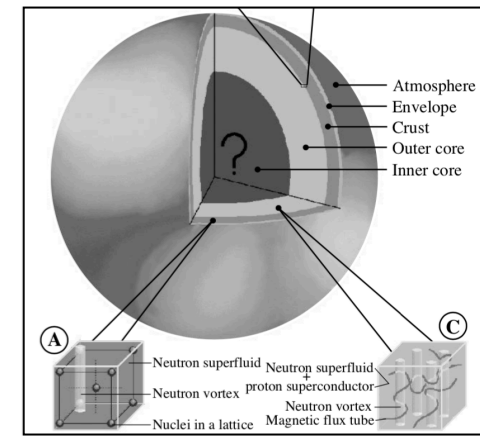


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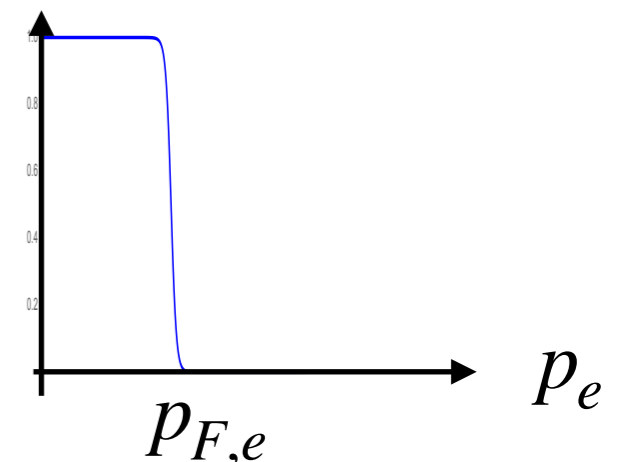
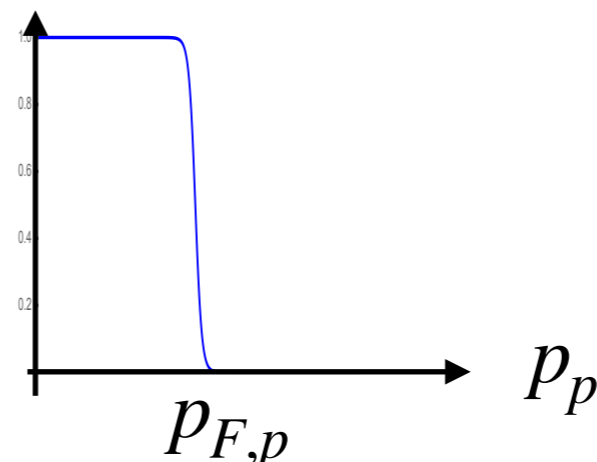
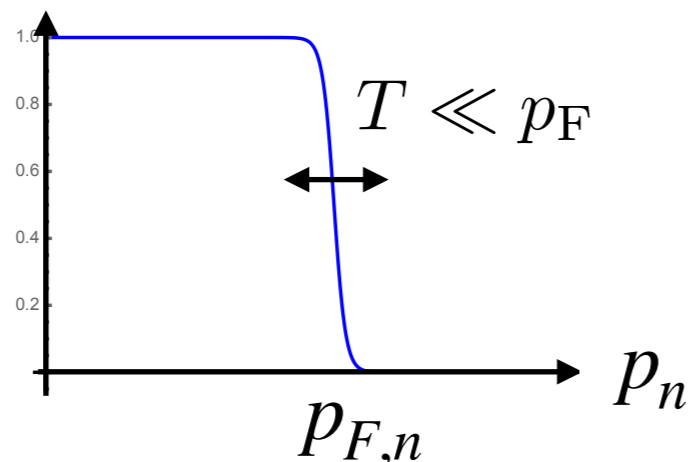
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- mostly composed of **neutrons**, + O(10%) **protons, electrons, muons**.

- **Neutrons, protons, and electrons are all Fermi degenerate.**



Neutron Star

- Most of NSs are found as **pulsars**.

Crab Pulsar



Neutron Star

Crab Pulsar

- Most of NSs are found as **pulsars**.

> 2700 pulsars found so far.

ATNF pulsar catalogue:



ATNF Pulsar Catalogue

[Catalogue Tutorial](#) | [Documentation](#) | [Expert](#) | [ATNF Pulsar Home](#) | [Pulsar Tutorial](#) | [Glitch table](#) | [Feedback](#) | [Download](#) | [History](#)

Catalogue version: 1.61

Predefined Variables [Display parameters](#)

<input type="checkbox"/> Name	<input type="checkbox"/> JName	<input type="checkbox"/> RaJ	<input type="checkbox"/> DecJ	<input type="checkbox"/> PMRA	<input type="checkbox"/> PMDec
<input type="checkbox"/> PX	<input type="checkbox"/> PosEpoch	<input type="checkbox"/> ELong	<input type="checkbox"/> ELat	<input type="checkbox"/> PMELong	<input type="checkbox"/> PMELat
<input type="checkbox"/> GL	<input type="checkbox"/> GB	<input type="checkbox"/> RaJD	<input type="checkbox"/> DecJD		
<input type="checkbox"/> P0	<input type="checkbox"/> P1	<input type="checkbox"/> F0	<input type="checkbox"/> F1	<input type="checkbox"/> F2	<input type="checkbox"/> F3
<input type="checkbox"/> PEpoch	<input type="checkbox"/> DM	<input type="checkbox"/> DM1	<input type="checkbox"/> RM	<input type="checkbox"/> W50	<input type="checkbox"/> W10
<input type="checkbox"/> Units	<input type="checkbox"/> Tau_sc	<input type="checkbox"/> S400	<input type="checkbox"/> S1400	<input type="checkbox"/> S2000	
<input type="checkbox"/> Binary	<input type="checkbox"/> T0	<input type="checkbox"/> PB	<input type="checkbox"/> A1	<input type="checkbox"/> OM	<input type="checkbox"/> Ecc
<input type="checkbox"/> Tasc	<input type="checkbox"/> Eps1	<input type="checkbox"/> Eps2	<input type="checkbox"/> MinMass	<input type="checkbox"/> MedMass	<input type="checkbox"/> BinComp

Neutron Star

Crab Pulsar

- Most of NSs are found as **pulsars**.

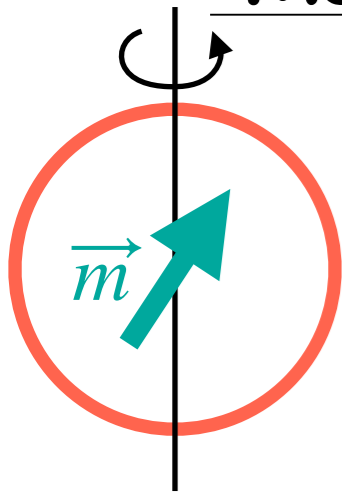
- **Magnetic Dipole Model**

Rotational energy loss \simeq magnetic dipole radiation

$$P(t) = \sqrt{P_0^2 + (P_{\text{now}} \dot{P}_{\text{now}}) t} \quad P_0 = \text{initial period}$$

$$\implies t \simeq \frac{P_{\text{now}}}{2\dot{P}_{\text{now}}} \equiv \tau_{\text{sd}}$$

spin down age / characteristic age



Neutron Star

- Most of NSs are found as **pulsars**.

• Magnetic Dipole Model

Rotational energy loss \simeq magnetic dipole radiation

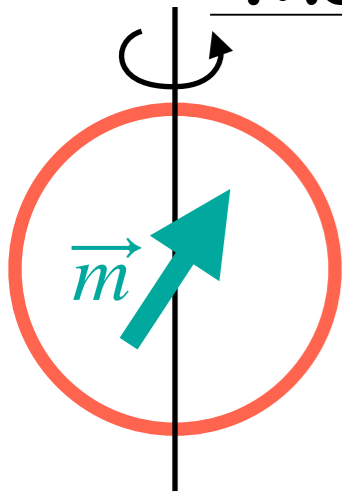
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spin down age / characteristic age

Example: Crab Pulsar

- actual age: $\tau = 965$ yrs (from historical records of SN1054.)
- spin down age: $P_{\text{now}} \simeq 0.033$ sec, $\dot{P}_{\text{now}} \simeq 4.2 \times 10^{-13} \implies \tau_{\text{sd}} \simeq 1200$ yrs



Plan

0. Introduction

1. Neutron Star  done

2. Neutron Star Cooling and Heating

3. Summary

Plan

0. Introduction

1. Neutron Star

2. Neutron Star Cooling and Heating

$$(i). C \frac{dT}{dt} = -L_\nu - L_\gamma \quad \text{Neutron Star Cooling}$$

$$(ii). C \frac{dT}{dt} = -L_\nu - L_\gamma + L_{DM}^{\text{heat}}$$

$$(iii). C \frac{dT}{dt} = -L_\nu - L_\gamma + L_{\text{rotochemical}}^{\text{heat}}$$

$$(iv). C \frac{dT}{dt} = -L_\nu - L_\gamma + L_{\text{rotochemical}}^{\text{heat}} + L_{DM}^{\text{heat}}$$

Previous works

Our works

K. Yanagi, N. Nagata, KH, [1904.04667]

KH, N. Nagata, K. Yanagi, [1905.02991]

3. Summary

(i). NS Cooling

$$C \frac{dT}{dt} = -L_\nu - L_\gamma$$

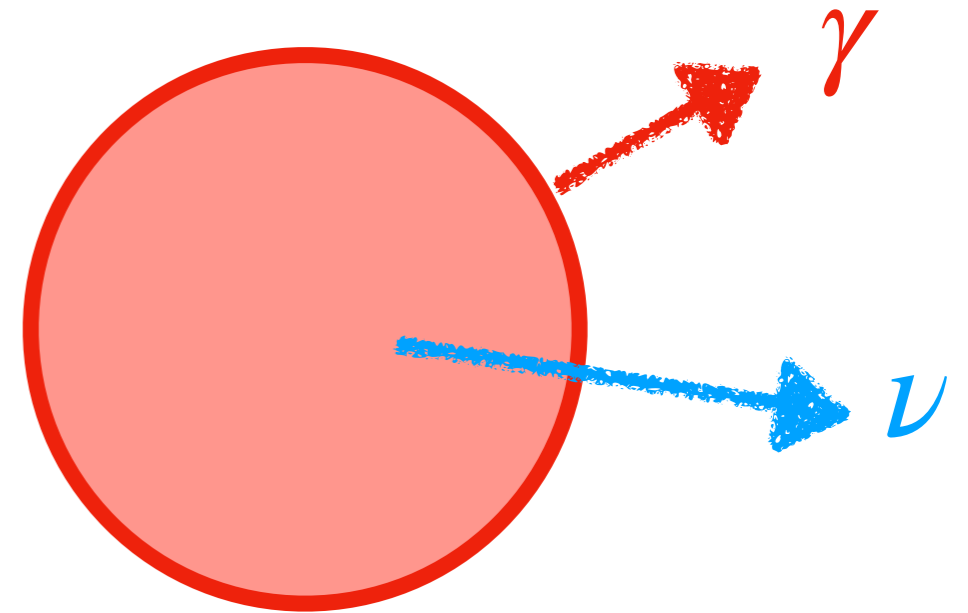
LHS = Temperature Evolution.

$$C = \frac{dE_{\text{thermal}}}{dT} \text{ (heat capacity)}$$

$$C = C_n + C_p + C_e + C_\mu$$

RHS = Cooling Luminosity.

$$-L = \frac{dE_{\text{thermal}}}{dt}$$



✂ assuming isothermal state $T(r) \propto e^{-\Phi(r)}$ for simplicity (valid for $t \gtrsim 100$ sec).

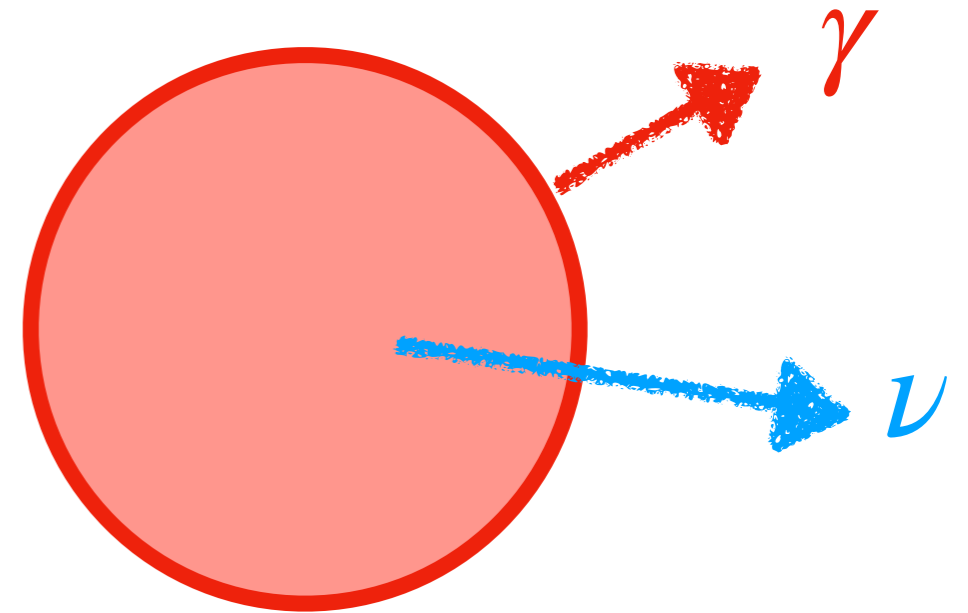
(i). NS Cooling

$$C \frac{dT}{dt} = -L_\nu - L_\gamma \leftarrow$$

Photon emission

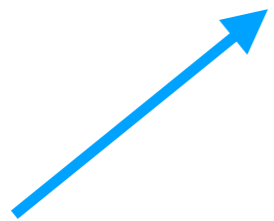
$$L_\gamma = 4\pi R^2 \sigma_{SB} T_s^4$$

dominant process at late time



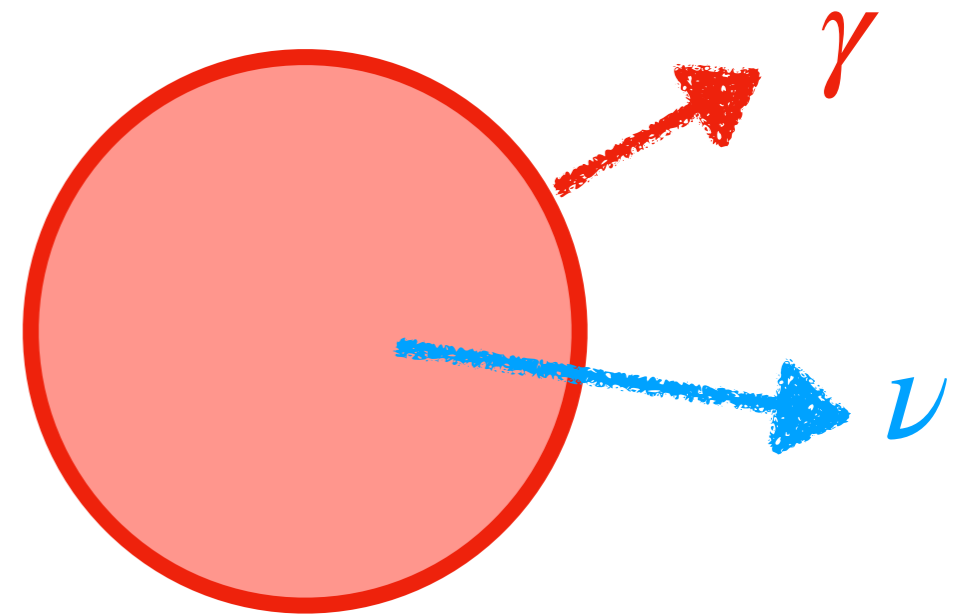
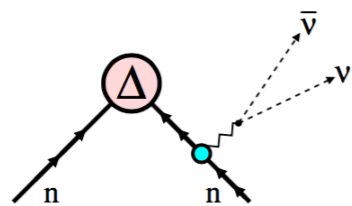
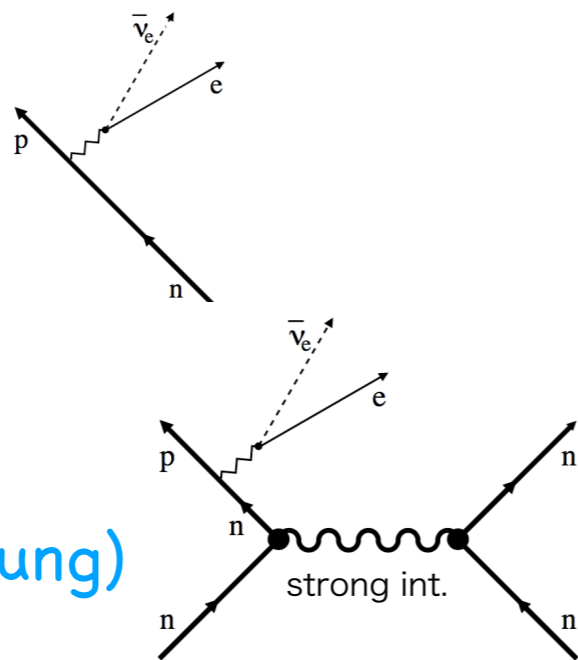
(i). NS Cooling

$$C \frac{dT}{dt} = -L_\nu - L_\gamma$$



Neutrino emission

- Direct Urca
- Modified Urca (& Bremsstrahlung)
- PBF



(i). NS Cooling

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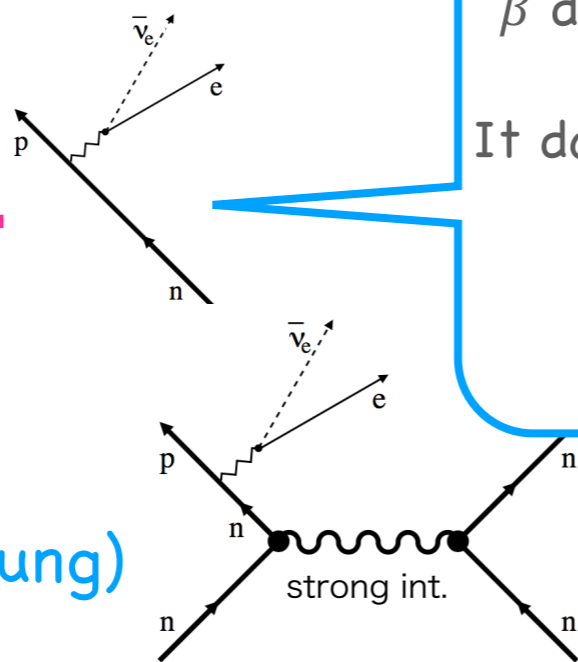
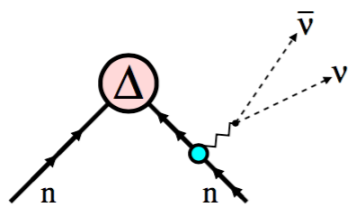


Neutrino emission

- ~~Direct Urca~~

- Modified Urca (& Bremsstrahlung)

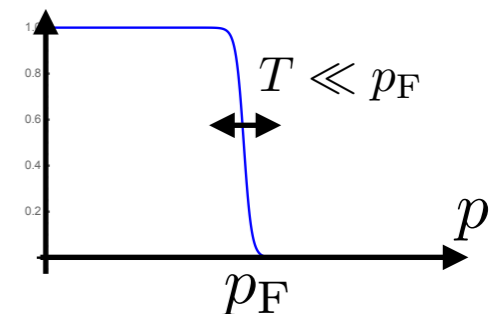
- PBF



β decay and its inverse: $\begin{cases} n \rightarrow p + e^- + \bar{\nu}_e \\ p + e^- \rightarrow n + \nu_e \end{cases}$

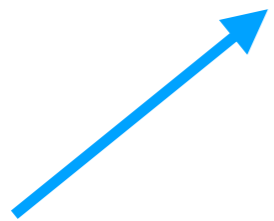
It does **NOT** work in typical NS because $p_p + p_e < p_n$.
Discarded in "minimal cooling" scenario.
D.Page+, astro-ph/0403657,
M.E.Gusakov+, astro-ph/0404002,
D.Page+, 0906.1621

⊗ Neutron, proton, electron are all Fermi degenerate.



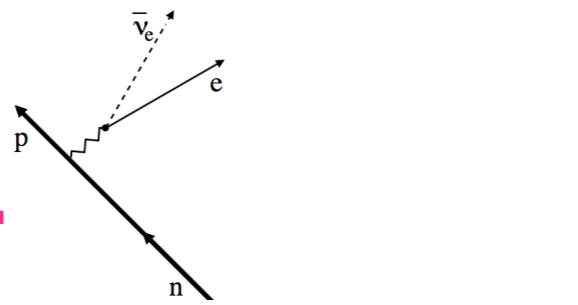
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$$C \frac{dT}{dt} = -L_\nu - L_\gamma$$



Neutrino emission

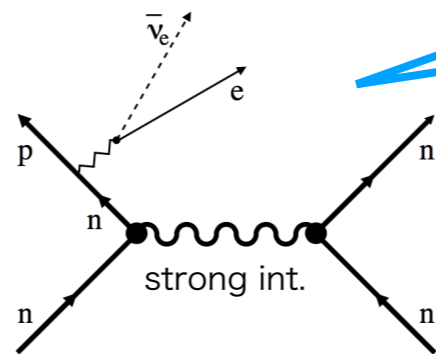
- ~~Direct Urca~~



$$\begin{cases} n + N \rightarrow p + e^- + N + \bar{\nu}_e \\ p + N + e^- \rightarrow n + N + \nu_e \end{cases} \quad (N = p \text{ or } n)$$

dominant process for $T > T_c$

- Modified Urca (& Bremsstrahlung)

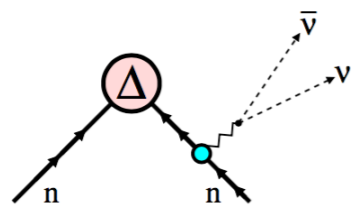


PBF (Cooper-pair breaking and formation)

$$\begin{cases} [\tilde{N}\tilde{N}] \rightarrow \tilde{N} + \tilde{N} \\ \tilde{N} + \tilde{N} \rightarrow [\tilde{N}\tilde{N}] + \nu + \bar{\nu} \end{cases} \quad (\tilde{N} : \text{quasi-particle}, [\tilde{N}\tilde{N}] : \text{Cooper-pair})$$

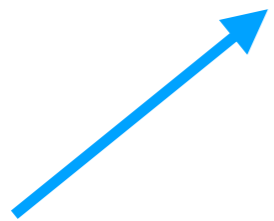
Important for $T < T_c$.

- PBF



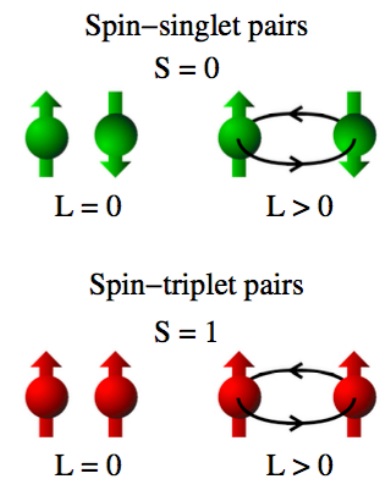
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Neutrino emission

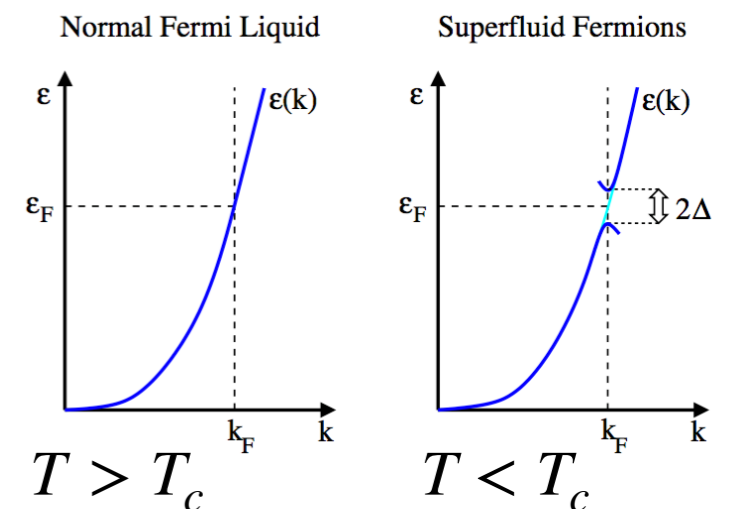
neutron singlet (1S_0)
neutron triplet (3P_2)
proton singlet (1S_0)



Superfluidity (pairing) plays important roles.

At $T < T_c$, Cooper pairing (p-p and n-n) occurs.

- Heat capacity C is suppressed.
- M.Urca luminosity $L_{\nu, MU}$ is suppressed.
- PBF occurs at $T < T_c$.
- It is also important for the "rotocemical heating" (see below).



(i). NS Cooling

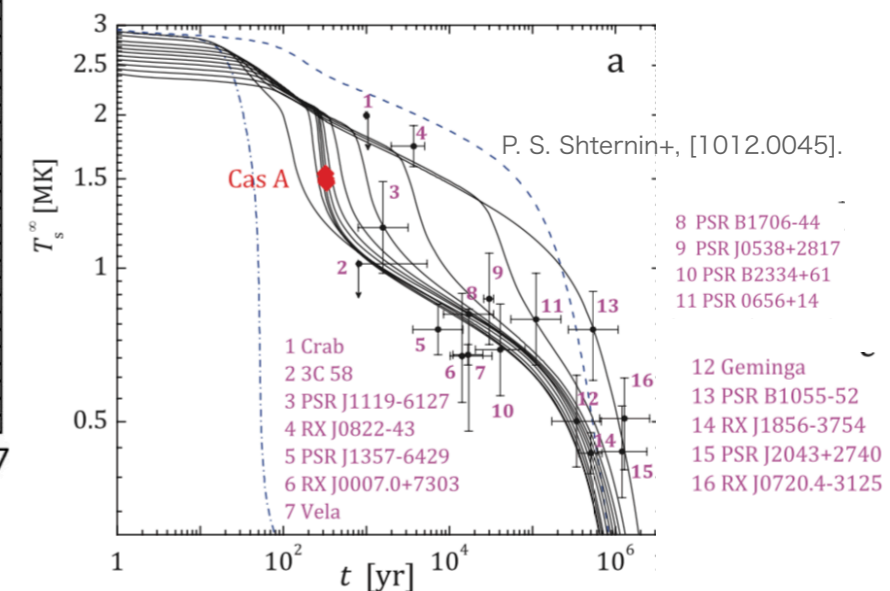
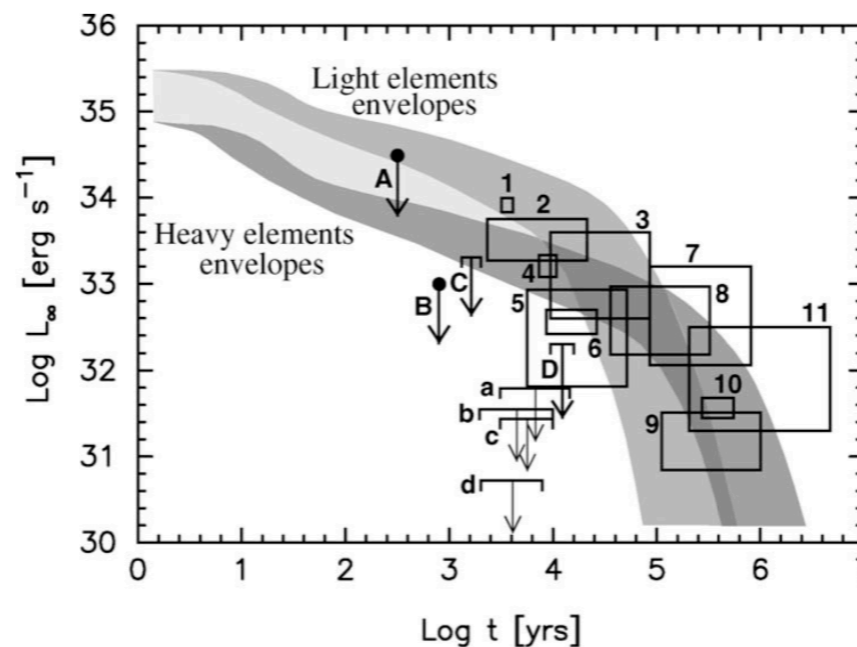
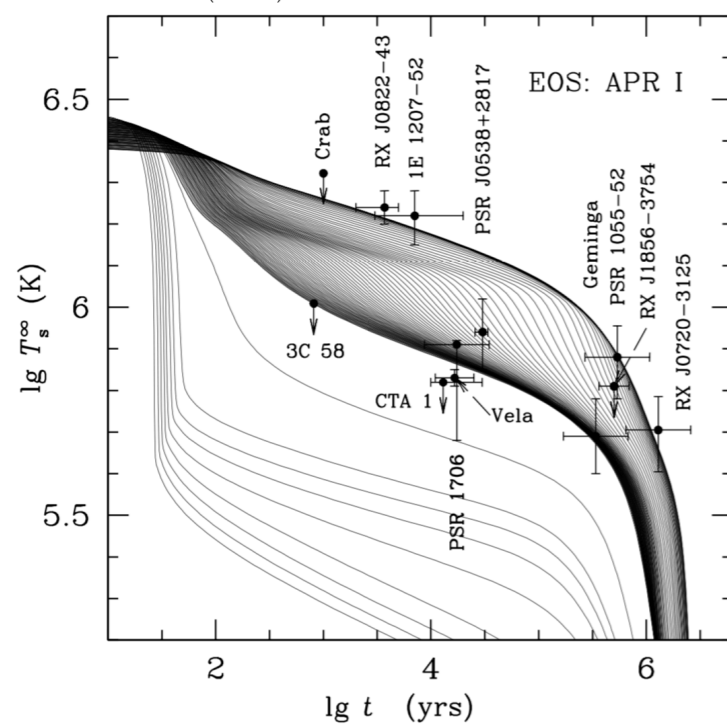
$$C \frac{dT}{dt} = -L_\nu - L_\gamma$$

The minimal cooling scenario can successfully explain many NS temperature observations.

D.Page+, astro-ph/0403657,
M.E.Gusakov+, astro-ph/0404002,
D.Page+, 0906.1621

D. Page et al. / Nuclear Physics A 777 (2006) 497–530

M. E. Gusakov, A. D. Kaminker, D. G. Yakovlev, O. Y. Gnedin
Mon.Not.Roy.Astron.Soc. **363** (2005) 555-562




Plan

0. Introduction

1. Neutron Star

2. Neutron Star Cooling and Heating

(i). $C \frac{dT}{dt} = -L_\nu - L_\gamma$  done

(ii). $C \frac{dT}{dt} = -L_\nu - L_\gamma + L_{DM}^{\text{heat}}$

(iii). $C \frac{dT}{dt} = -L_\nu - L_\gamma + L_{\text{rotochemical}}^{\text{heat}}$

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Previous works

Our works

K. Yanagi, N. Nagata, KH, [1904.04667]

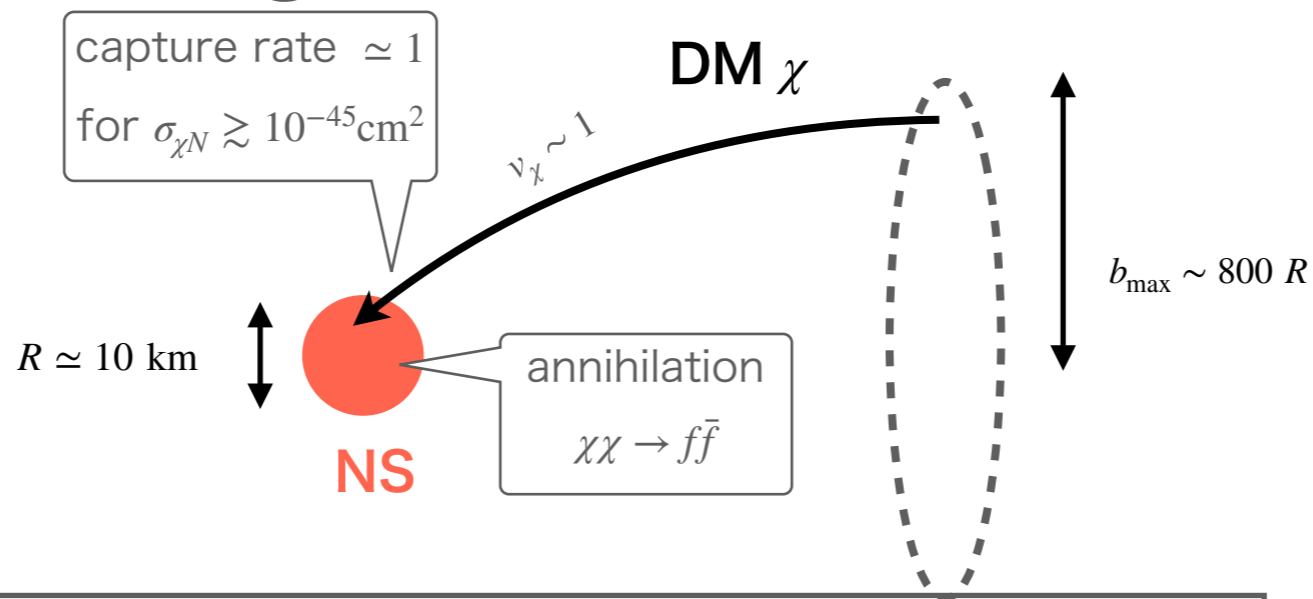
KH, N. Nagata, K. Yanagi, [1905.02991]

3. Summary

(ii). NS Heating by DM

C. Kouvaris 0708.2362,
 G. Bertone+ 0709.1485,
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 A. de Lavallaz+ 1004.0629
 . . .

+ many recent works: e.g.,
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$$C \frac{dT}{dt} = -L_\nu - L_\gamma + L_{\text{DM heat}}$$

Photon emission

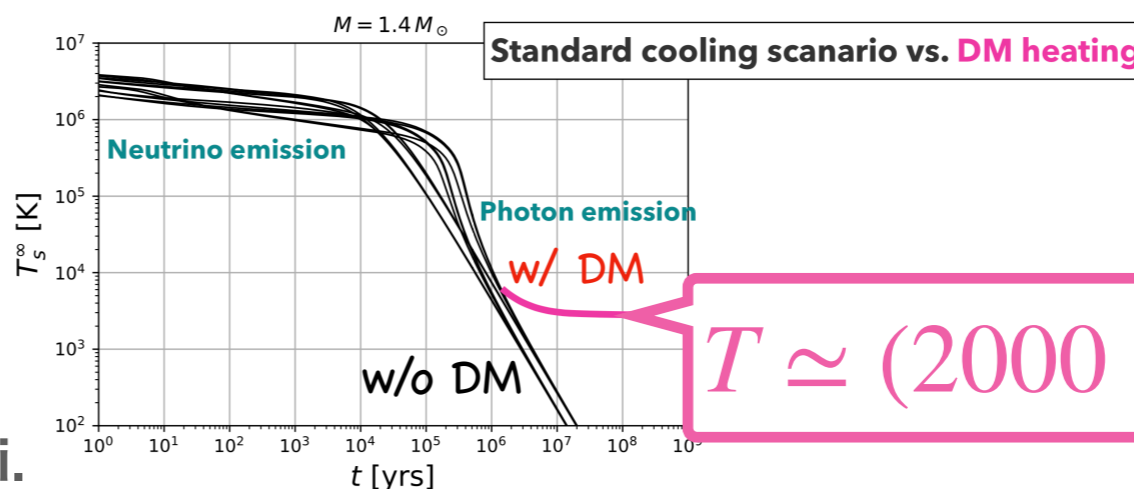
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dominant process at late time

DM heating

$$L_{\text{DM heat}} \sim \pi b_{\text{max}}^2 \rho_\chi v_\chi \simeq 3 \times 10^{22} \text{erg/s}$$

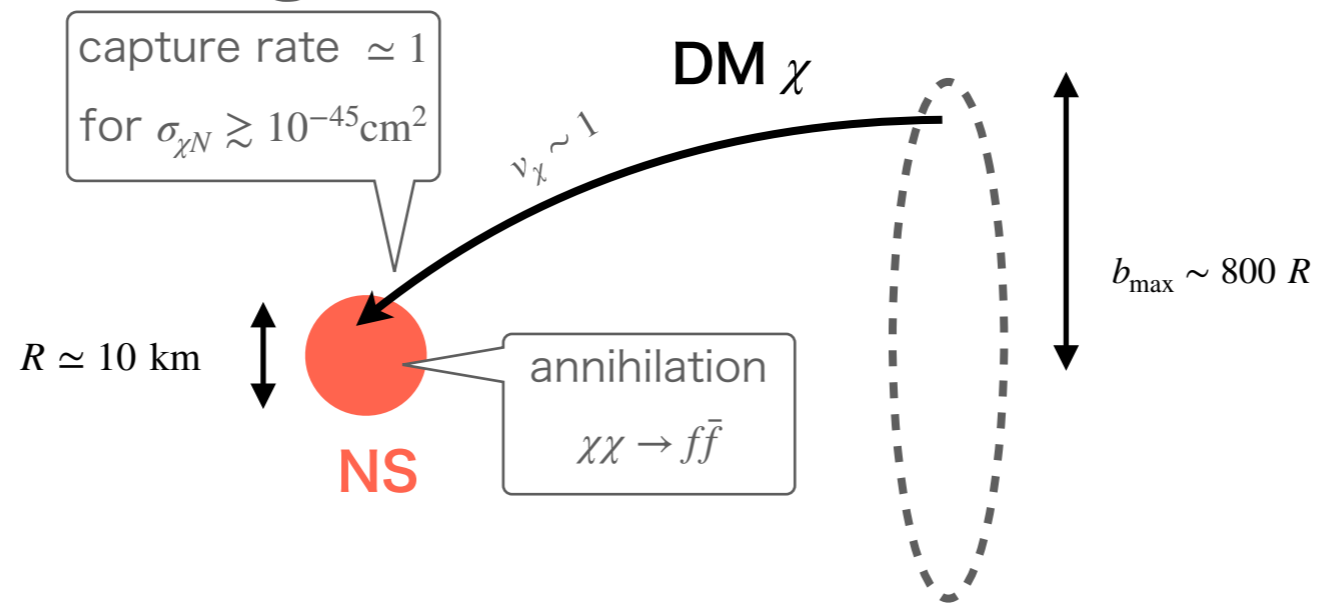
independent of the DM mass



$T \simeq (2000 - 3000)K$

Fig. by K.Yanagi.

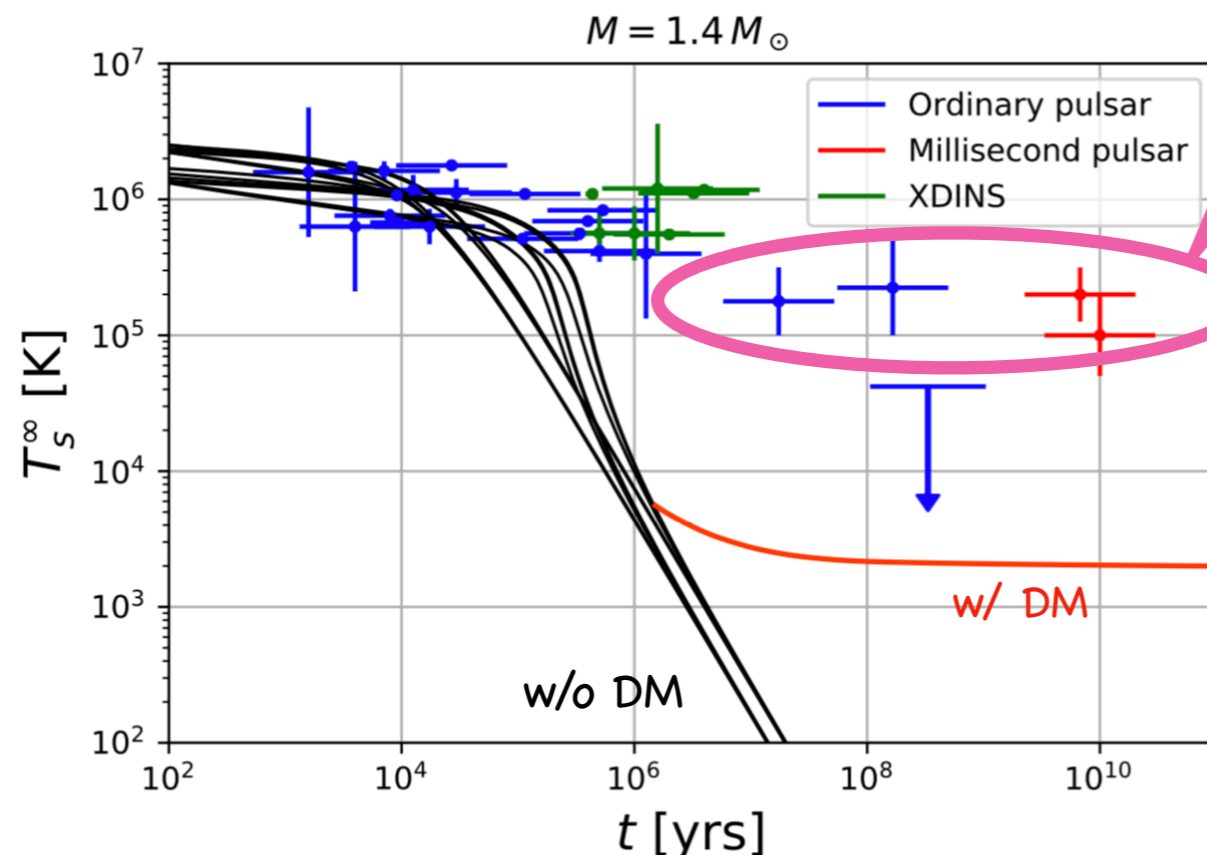
(ii). NS Heating by DM



C. Kouvaris 0708.2362,
G. Bertone+ 0709.1485,
C. Kouvaris+ 1004.0586,
A. de Lavallaz+ 1004.0629
...

+ many recent works: e.g.,
J. Bramante+ 1703.04043
M. Baryakhtar+ 1704.01577
N. Raj+ 1707.09442
C.-S. Chen+ 1804.03409
N. F. Bell+ 1807.02840
D. A. Camargo+ 1901.05474
N. F. Bell+ 1904.09803
KH, N.Nagata, K.Yanagi 1905.02991
R. Garani+ 1906.10145
J. F. Acevedo+, 1911.06334
A. Joglekar+, 1911.13293

But,...



Neither standard NS cooling nor DM can explain those old and warm NSs.

Plan

0. Introduction

1. Neutron Star

2. Neutron Star Cooling and Heating

$$(i). C \frac{dT}{dt} = -L_\nu - L_\gamma$$

$$(ii). C \frac{dT}{dt} = -L_\nu - L_\gamma + L_{DM}^{\text{heat}}$$

↑ done

$$(iii). C \frac{dT}{dt} = -L_\nu - L_\gamma + L_{\text{rotochemical}}^{\text{heat}}$$

$$(iv). C \frac{dT}{dt} = -L_\nu - L_\gamma + L_{\text{rotochemical}}^{\text{heat}} + L_{DM}^{\text{heat}}$$

Previous works

Our works

K. Yanagi, N. Nagata, KH, [1904.04667]

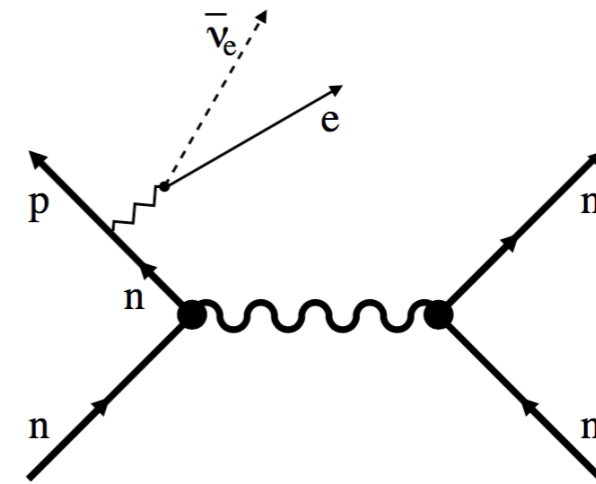
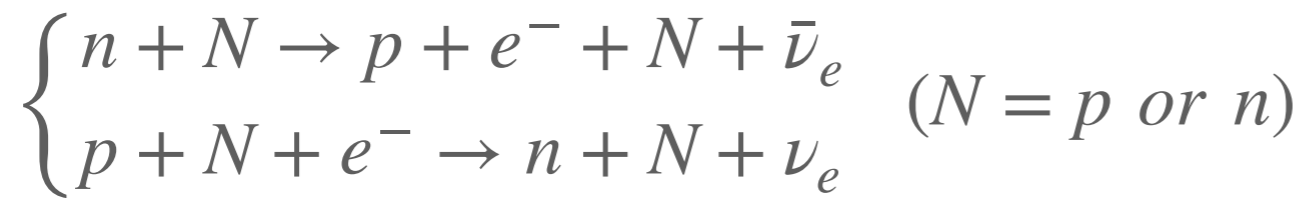
KH, N. Nagata, K. Yanagi, [1905.02991]

3. Summary

(iii). Rotochemical Heating $C \frac{dT}{dt} = -L_\nu - L_\gamma + L_{\text{rotochemical}}^{\text{heat}}$

(iii). Rotochemical Heating $C \frac{dT}{dt} = -L_\nu - L_\gamma + L_{\text{heat rotochemical}}$

- Modified Urca (dominant process at $T > T_c$)

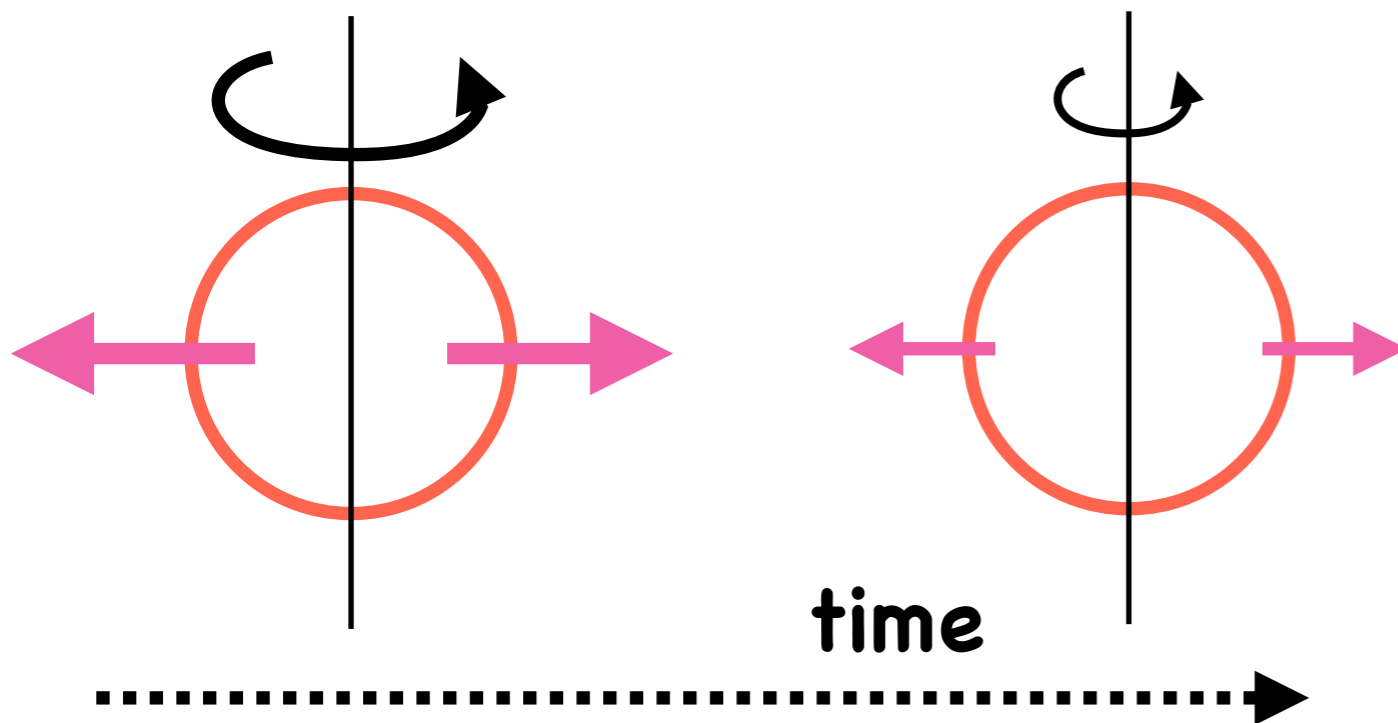


- In the minimal cooling, β -equilibrium is assumed.

$$\Gamma_{n \rightarrow p+e} = \Gamma_{p+e \rightarrow n}, \quad \mu_n = \mu_p + \mu_e.$$

- However, β -equilibrium is NOT maintained in rotating pulsars!

A.Reisenegger [astro-ph/9410035]

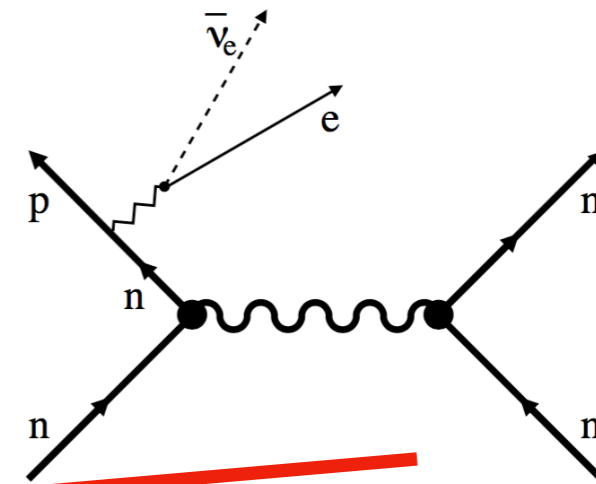


spin-down weakens the centrifugal force.
 \Rightarrow pressure changes.
 \Rightarrow chemical eq. condition changes
 \Rightarrow at low T ,
 the modified Urca process (slow, $\sim T^8$)
 can no longer maintain the equilibrium.

(iii). Rotochemical Heating $C \frac{dT}{dt} = -L_\nu - L_\gamma + L_{\text{rotochemical}}^{\text{heat}}$

- Modified Urca (dominant process at $T > T_c$)

$$\begin{cases} n + N \rightarrow p + e^- + N + \bar{\nu}_e \\ p + N + e^- \rightarrow n + N + \nu_e \end{cases} \quad (N = p \text{ or } n)$$



- In the minimal cooling, β -equilibrium is assumed.

~~$$\Gamma_{n \rightarrow p+e} = \Gamma_{p+e \rightarrow n}, \quad \mu_n = \mu_p + \mu_e.$$~~

- However, β -equilibrium is NOT maintained in rotating pulsars!

A.Reisenegger [astro-ph/9410035]

$$\Gamma_{n \rightarrow p+e} > \Gamma_{p+e \rightarrow n}, \quad \mu_n > \mu_p + \mu_e$$

- The deviation from β -equilibrium **heats the NS.**

$$L_{\text{rotochemical}}^{\text{heat}} = \int dV (\mu_n - \mu_p - \mu_e) (\Gamma_{n \rightarrow p+e} - \Gamma_{p+e \rightarrow n}) > 0.$$

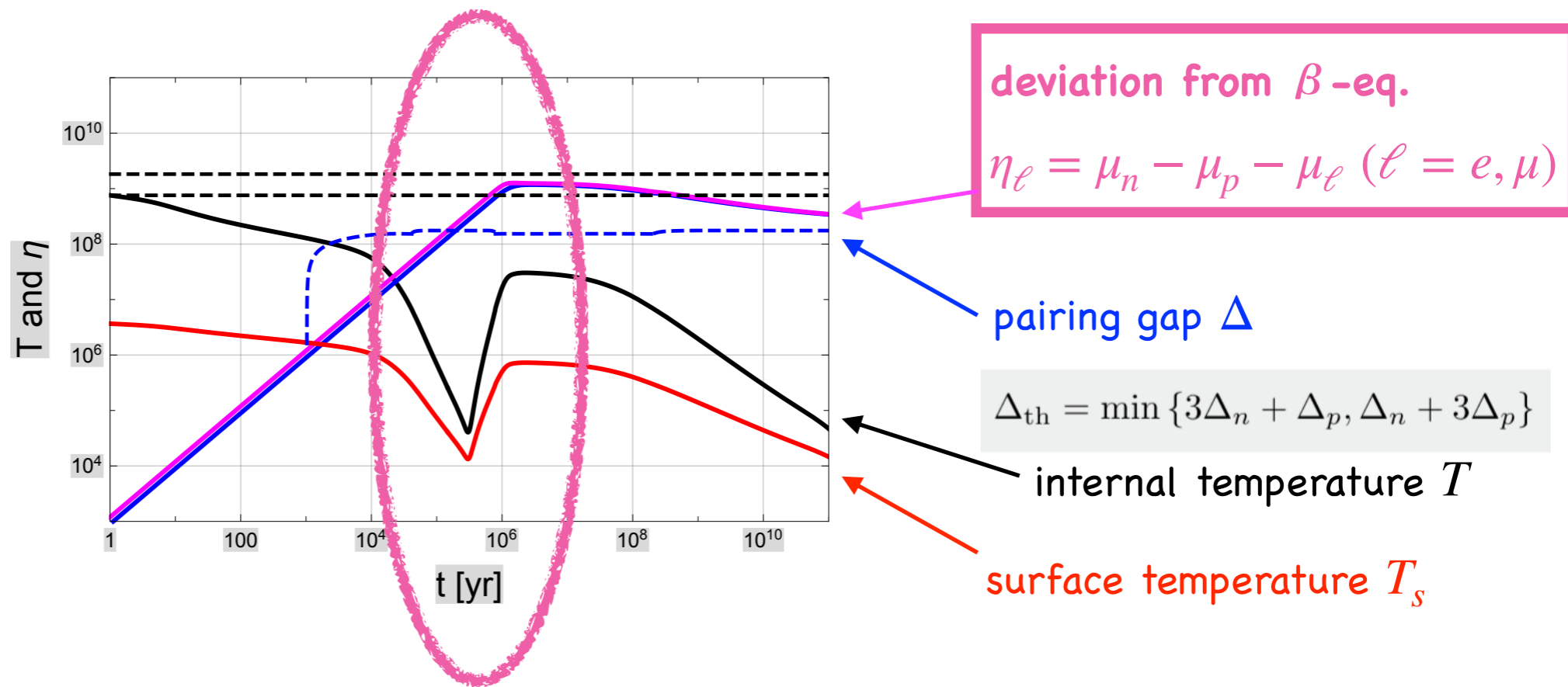
“Rotochemical heating” (nothing special, just normal physics!)

(iii). Rotochemical Heating $C \frac{dT}{dt} = -L_\nu - L_\gamma + L_{\text{rotochemical}}^{\text{heat}}$

Example

for millisecond pulsar

- neutron gap: "a"
- proton gap: AO
- $M = 1.4M_\odot$
- $P_0 = 1\text{ms}$
- $P\dot{P} = 3.3 \times 10^{-22}\text{s}$

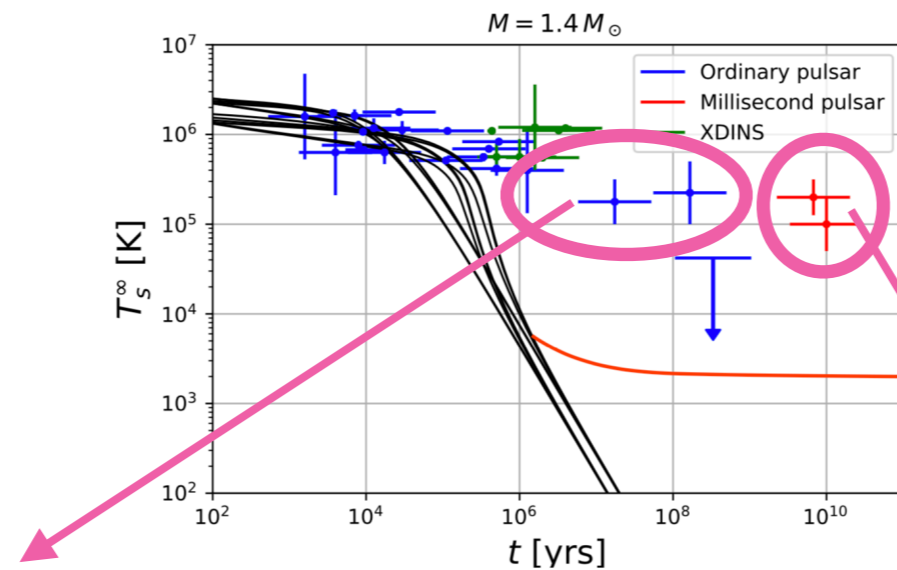


- The **superfluidity (pairing gap)** plays an important role.
Rotochemical heating begins when $\eta_\ell > \Delta$. [Petrovich & Reisenegger, 0912.2564]
- Recently, we have updated the calculation.
K. Yanagi, N. Nagata, KH, [arXiv:1904.04667] MNRAS (to be published)

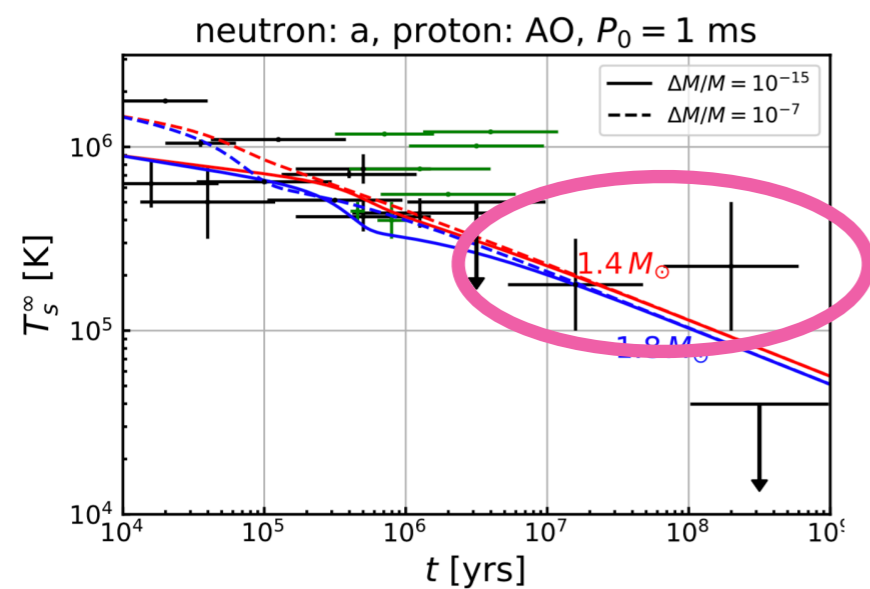
(iii). Rotochemical Heating $C \frac{dT}{dt} = -L_\nu - L_\gamma + L_{\text{rotochemical}}^{\text{heat}}$

The rotochemical heating can explain the old and warm NSs.

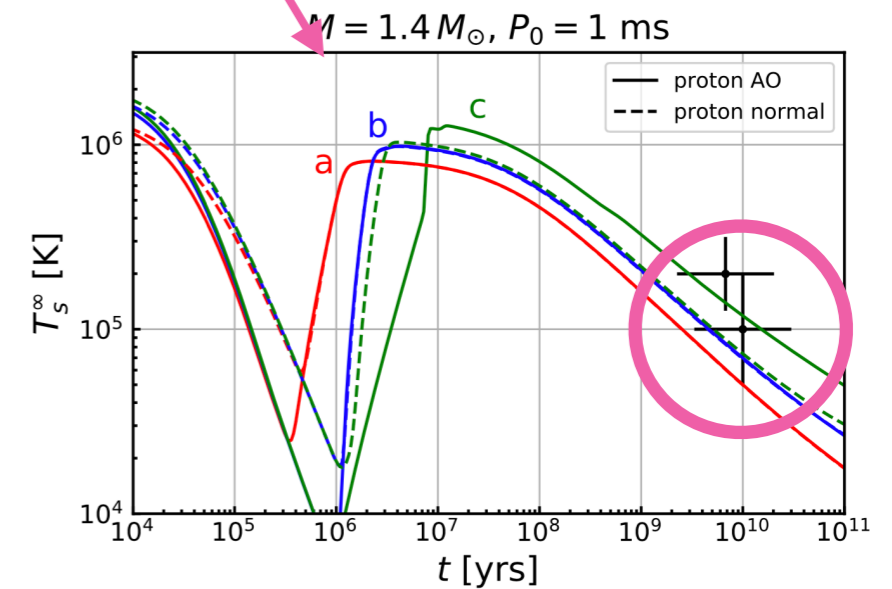
NOTE: No exotic physics assumed. The same setting as the Standard Cooling.



K. Yanagi, N. Nagata, KH
[arXiv:1904.04667]



Ordinary pulsar
(typically $P \sim 1\text{s}$, $\dot{P} \sim 10^{-14}$, $B \sim 10^{12}\text{G}$)



Millisecond pulsar
(typically $P \sim 1\text{ms}$, $\dot{P} \sim 10^{-20}$, $B \sim 10^8\text{G}$)

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Previous works

Our works

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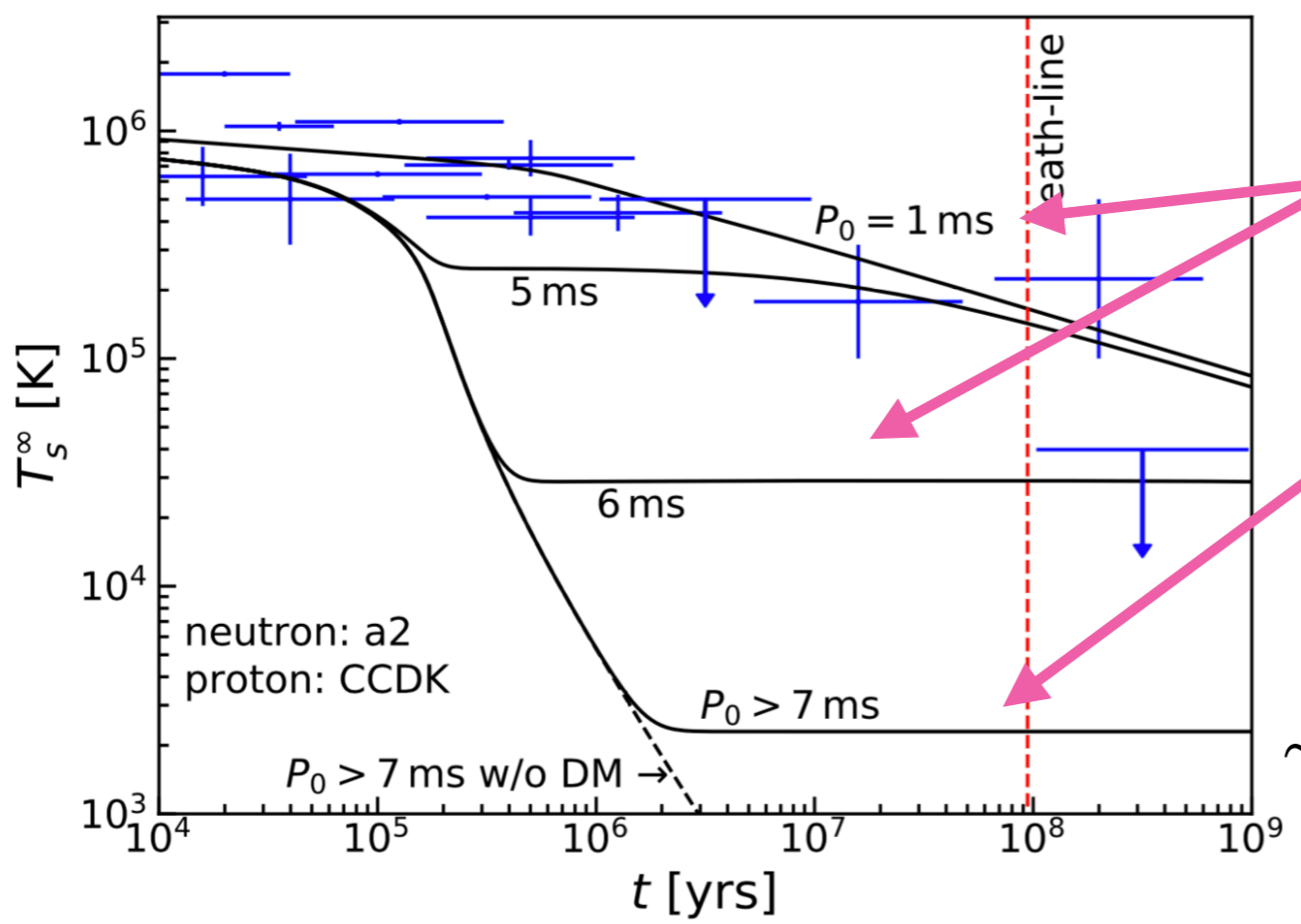
(iv). DM heating vs. Rotochemical heating

$$C \frac{dT}{dt} = -L_\nu - L_\gamma + L^{\text{heat rotochemical}} + L^{\text{heat DM}}$$

KH, N. Nagata, K. Yanagi, [1905.02991]

Result

$$P = 1\text{s}, \dot{P} = 10^{-15}$$



P_0 : initial rotation period is the key parameter.

• For a short P_0 , DM heating effect is invisible.

• For a long P_0 , DM heating effect is visible!

$\sim (2000 - 3000)K$

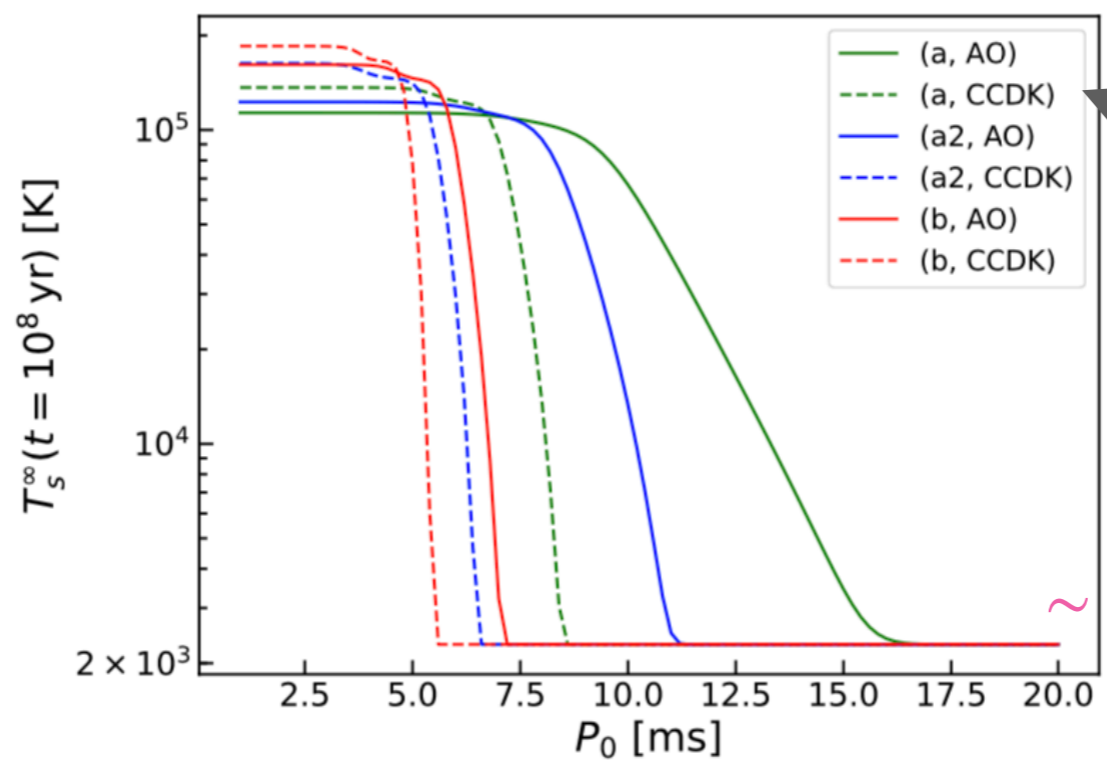
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KH, N. Nagata, K. Yanagi, [1905.02991]

Result

Late time temperature



initial rotation period (P_0) dependence

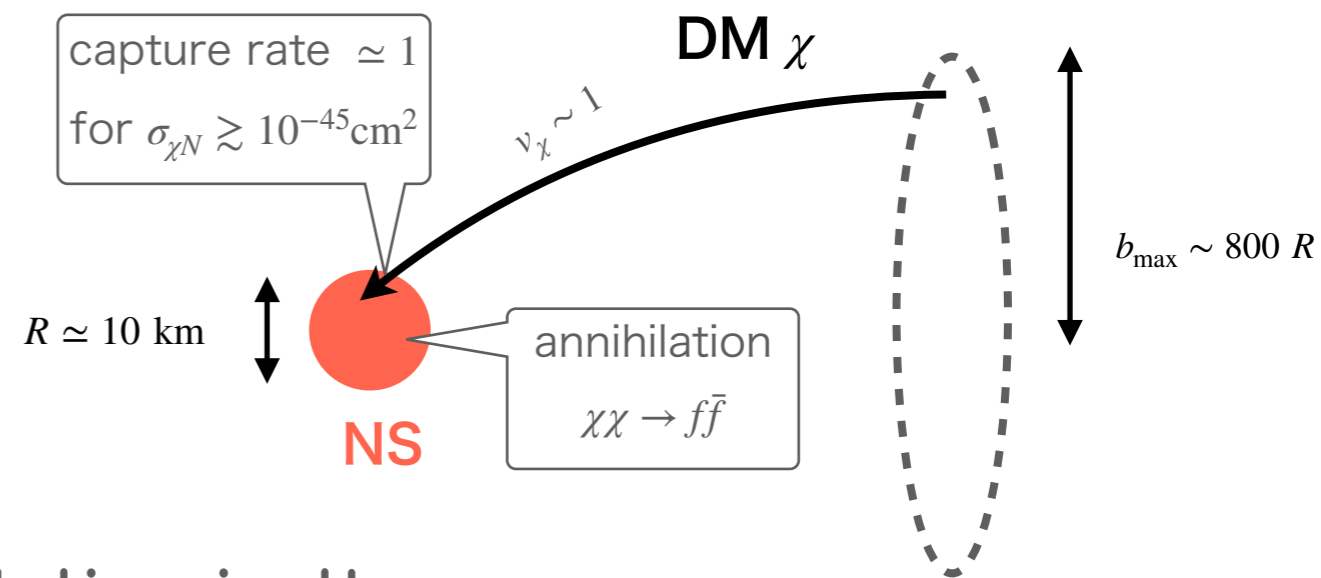
• For large enough P_0 , DM signal is visible.

Recent studies suggest P_0 can indeed be very large. ($> 100\text{ms}$). [cf. references in 1905.02991.]

Currently no NS with such a low T is observed.

• Conversely, discovery of a NS with $T < 2000\text{K}$ will exclude many DM models, such as Wino DM.

Summary



- We studied NS temperature evolution in the presence of both **rotochemical heating** and **DM heating**.

$$C \frac{dT}{dt} = -L_\nu - L_\gamma + L_{\text{rotochemical}}^{\text{heat}} + L_{\text{DM}}^{\text{heat}}$$

and found that DM heating effect is indeed visible for a large initial rotation period P_0 .

Future works

- application to concrete DM models.
- vs. other heating mechanisms. (cf. D.Gonzalez, 1005.5699)
- observational feasibility.

New Directions in Cosmology

24-27 March 2020

Hongo Campus

Asia/Tokyo timezone

Workshop at Hongo campus, March 24-27, 2020.

Overview

Timetable

Registration

Participant List

Accommodations

Transportation

Information

Contact

✉ newcosmo@hep-th.phys...

We are pleased to announce the international workshop "New Directions in Cosmology." The aim of the workshop is to bring together experts and to exchange ideas in cosmology, astroparticle physics, and related subjects in particle physics.

Dates and venue

Dates: March 24-27, 2020

Venue: Room 285/337A, Faculty of Science Bldg.1 East, Hongo Campus, The University of Tokyo.

Speakers include:

Tobias Binder (Kavli IPMU)
Kfir Blum (CERN; Weizmann Institute of Science)
Gongjun Choi (TDLI)
Nagisa Hiroshima (iTHEMS, RIKEN)
Kenta Hotokezaka (Princeton Univ.; RESCEU)
Alejandro Ibarra (TUM)
Kiyotomo Ichiki (Nagoya Univ.; KMI)
Koji Ishiwata (Kanazawa Univ.)
Akito Kusaka (Univ. of Tokyo)
Tongyan Lin (UC, San Diego)
Kohta Murase (Pennsylvania State Univ.)
Seong Chan Park (Yonsei University)
Nirmal Raj (TRIUMF)
Kenichi Saikawa (Kanazawa Univ.)
Yevgeny Stadnik (Kavli IPMU)
Tomo Takahashi (Saga Univ.)
Kenneth C. Wong (Kavli IPMU)
+ more to be confirmed...

Organizers:

Koichi Hamaguchi (co-chair, Physics Department, Tokyo U.)
Takeo Moroi (co-chair, Physics Department, Tokyo U.)
Natsumi Nagata (Physics Department, Tokyo U.)
Kazunori Nakayama (Physics Department, Tokyo U.)
Tom Melia (IPMU)
Shigeki Matsumoto (IPMU)
Tsutomu Yanagida (IPMU, TDLI)
Javier Menendez (University of Barcelona)
Kentaro Nagamine (Osaka)
Satoshi Iso (KEK)
Masahiro Yamaguchi (Tohoku)

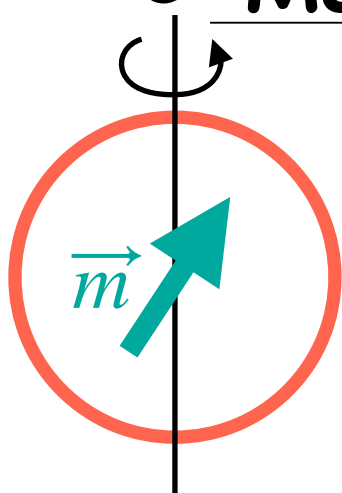
Backup

Neutron Star

- Most of NSs are found as **pulsars**.

● Magnetic Dipole Model

Rotational energy loss \simeq magnetic dipole radiation


$$\dot{E} = \frac{d}{dt} \left(\frac{1}{2} I \Omega^2 \right) = - \frac{\sin^2 \alpha}{6} R^6 \Omega^4 B_p^2 \quad \left\{ \begin{array}{l} I = \text{moment of inertia} \\ \Omega = 2\pi/P = \text{angular velocity} \\ P = \text{rotation period} \\ B_p = \text{magnetic field at the pole} \end{array} \right.$$

By solving this,

$$P(t) = \sqrt{P_0^2 + (P_{\text{now}} \dot{P}_{\text{now}}) t} \quad P_0 = \text{initial period}$$

$$\implies t \simeq \frac{P_{\text{now}}}{2\dot{P}_{\text{now}}} \equiv \tau_{\text{sd}}$$

spin down age / characteristic age

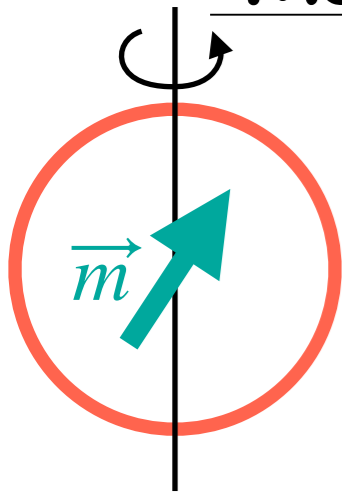


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spin down age / characteristic age

Example: Crab Pulsar

- actual age: $\tau = 965$ yrs (from historical records of SN1054.)
- spin down age: $P \simeq 0.033$ sec, $\dot{P} \simeq 4.2 \times 10^{-13} \implies \tau_{\text{sd}} \simeq 1200$ yrs



Neutron Star

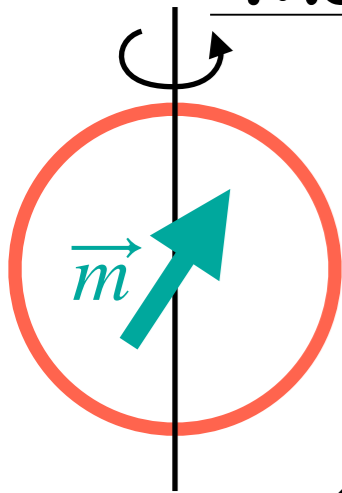
Crab Pulsar



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$$\implies t \simeq \frac{P_{\text{now}}}{2\dot{P}_{\text{now}}} \equiv \tau_{\text{sd}}$$

spin down age / characteristic age

The **magnetic field** can also be estimated from P_{now} and \dot{P}_{now} .

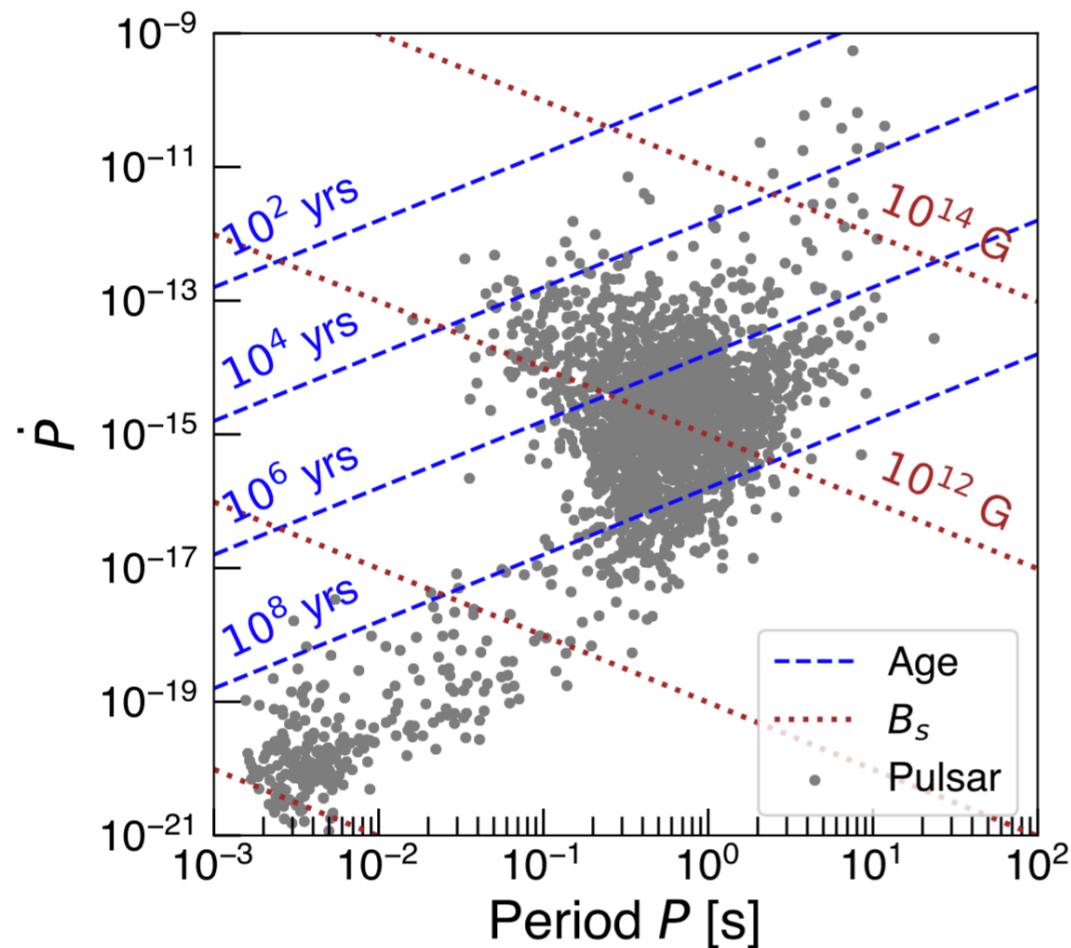
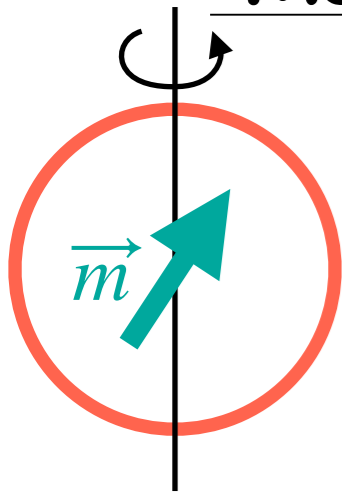
$$B_p \simeq 1 \times 10^{20} \text{G} \frac{1}{\sin \alpha} \left(\frac{P_{\text{now}} \dot{P}_{\text{now}}}{1 \text{ sec}} \right)^{1/2} \times \left(\frac{10 \text{ km}}{R} \right)^3 \left(\frac{I}{10^{45} \text{g} \cdot \text{cm}^2} \right)^{1/2}$$

Neutron Star

Crab Pulsar

- Most of NSs are found as **pulsars**.

● Magnetic Dipole Model



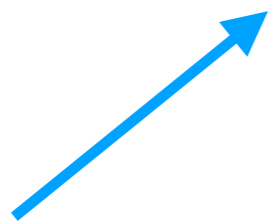
$P - \dot{P}$ diagram

\Leftrightarrow Age and magnetic field of pulsars.

Fig. thanks to N.Natsumi.

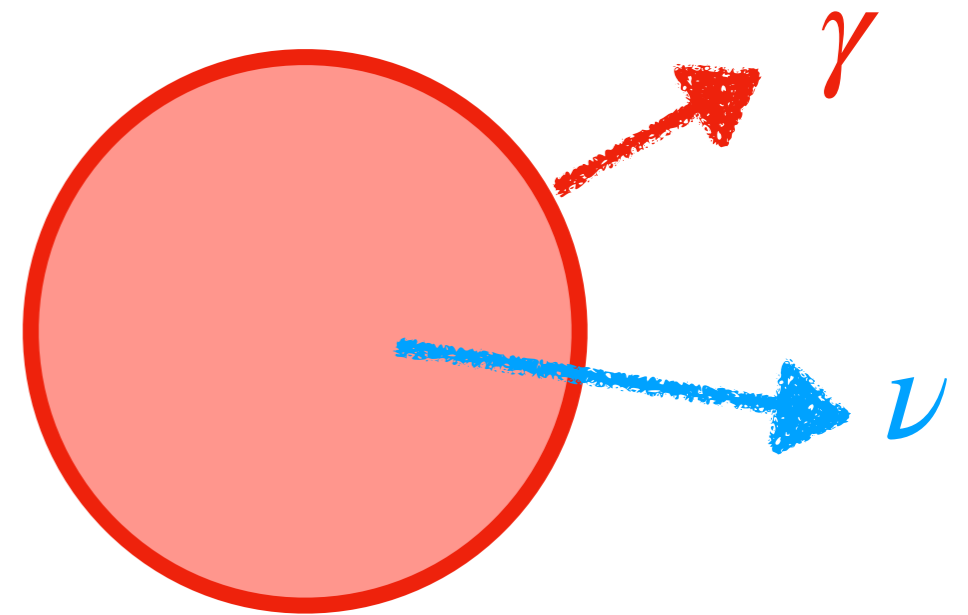
(i). NS Cooling

$$C \frac{dT}{dt} = -L_\nu - L_\gamma$$



Neutrino emission

- Direct Urca
- Modified Urca
- Bremsstrahlung
- PBF



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$$C \frac{dT}{dt} = -L_\nu - L_\gamma$$



Neutrino emission

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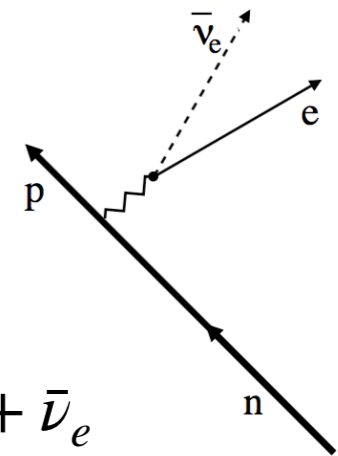
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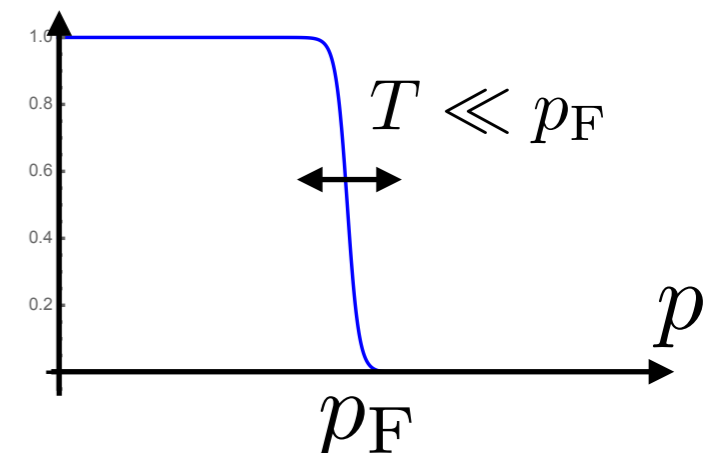
β decay and its inverse: $\begin{cases} n \rightarrow p + e^- + \bar{\nu}_e \\ p + e^- \rightarrow n + \nu_e \end{cases}$

It does **NOT** work in typical NS because $p_p + p_e < p_n$.

Discarded in "minimal cooling" scenario.
D.Page+, astro-ph/0403657,
M.E.Gusakov+, astro-ph/0404002,
D.Page+, 0906.1621

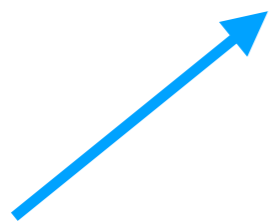


⊗ Neutron, proton, electron
are all Fermi degenerate.



(i). NS Cooling

$$C \frac{dT}{dt} = -L_\nu - L_\gamma$$



Neutrino emission

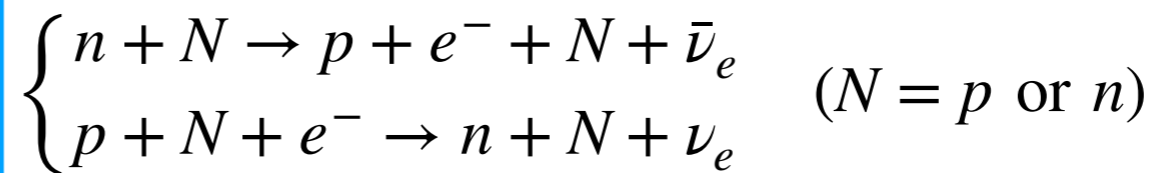
~~• Direct Urca~~

• **Modified Urca**

• Bremsstrahlung

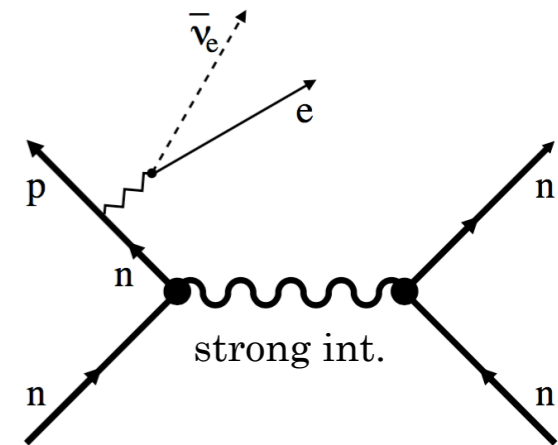
• PBF

- Dominant process (before the onset of Cooper pairing)

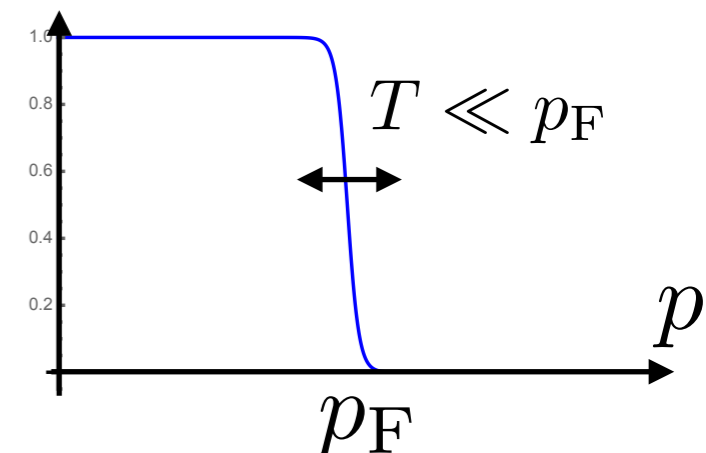


$$L_\nu^{\text{MU}} \sim T^8$$

$$L_\nu^{\text{MU}} \sim \underbrace{\int d^3 p_n}_T \underbrace{\int d^3 p_N}_T \cdot \underbrace{\int d^3 p_p}_T \underbrace{\int d^3 p_N}_T \underbrace{\int d^3 p_e}_T \cdot \underbrace{\int d^3 p_\nu \delta^4(p_i - p_f) E_\nu}_{T^3}$$

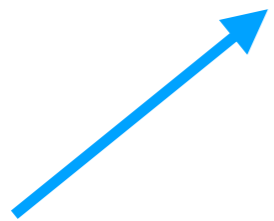


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(i). NS Cooling

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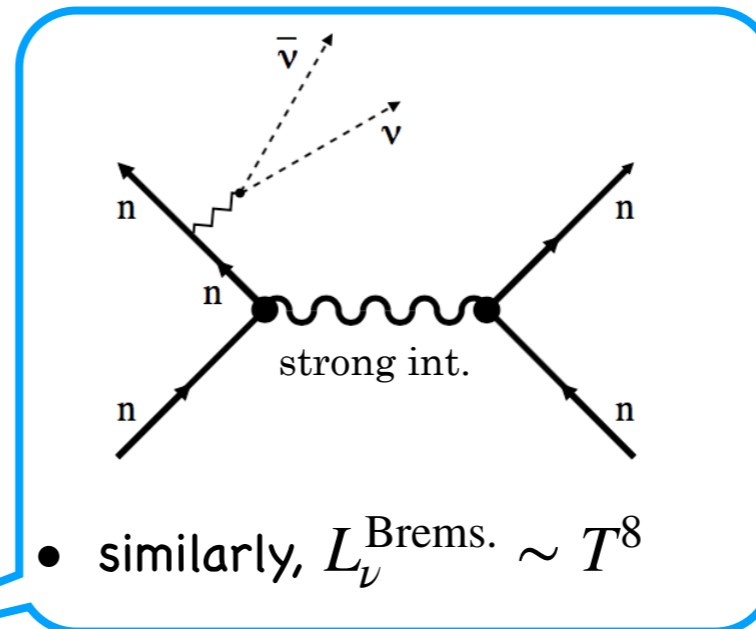
Neutrino emission

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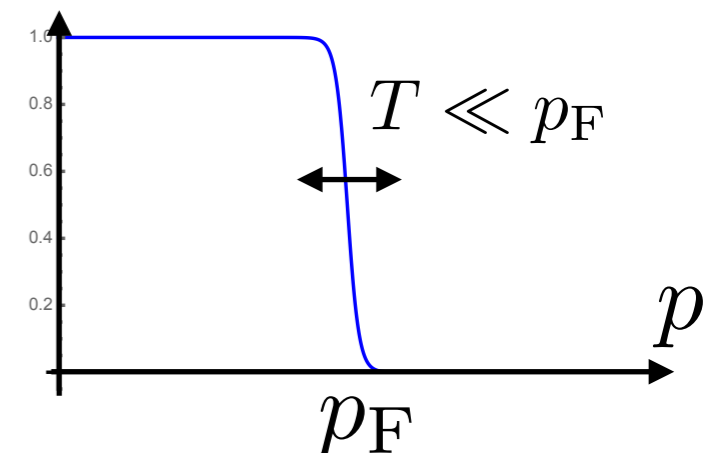
• Modified Urca

• **Bremsstrahlung**

• PBF

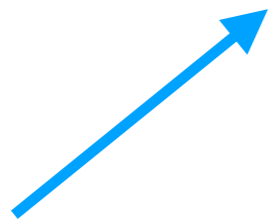


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(i). NS Cooling

$$C \frac{dT}{dt} = -L_\nu - L_\gamma$$



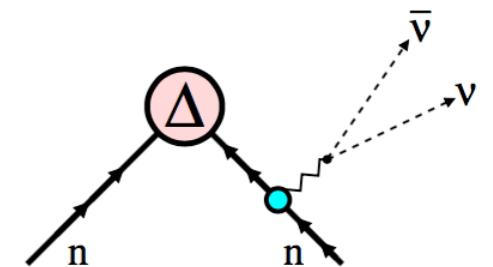
Neutrino emission

- ~~Direct Urca~~
- Modified Urca
- Bremsstrahlung
- **PBF**

PBF (Cooper-pair breaking and formation)

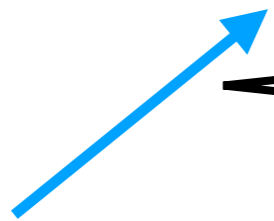
$$\begin{cases} [\tilde{N}\tilde{N}] \rightarrow \tilde{N} + \tilde{N} \\ \tilde{N} + \tilde{N} \rightarrow [\tilde{N}\tilde{N}] + \nu + \bar{\nu} \end{cases} \quad (\tilde{N} : \text{quasi-particle}, [\tilde{N}\tilde{N}] : \text{Cooper-pair})$$

Important for $T < T_c$.



(i). NS Cooling

$$C \frac{dT}{dt} = -L_\nu - L_\gamma$$



Neutrino emission

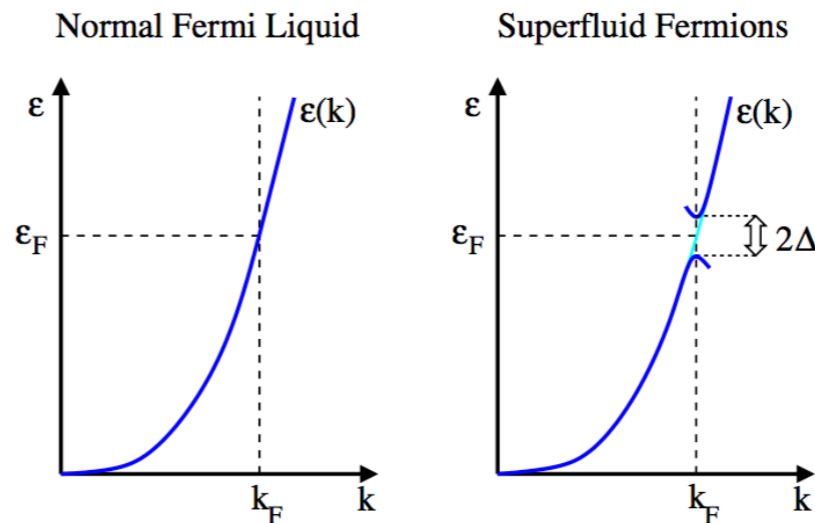
- ~~Direct Urca~~
- Modified Urca
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Superfluidity (pairing) plays important roles.

- At $T < T_c$, Cooper pairing occurs.

$T > T_c$

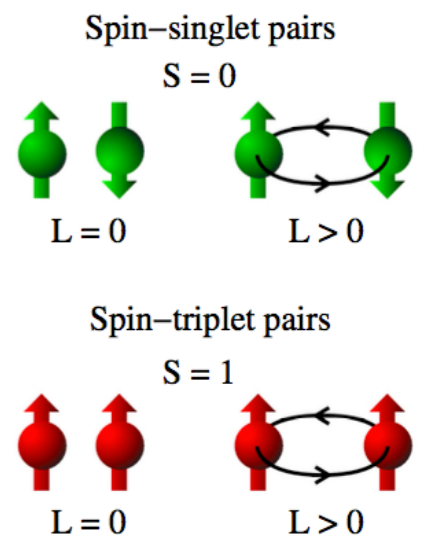
$T < T_c$



$$\epsilon(p) \simeq \sqrt{\Delta^2 + v_F^2(p - p_F)^2}$$

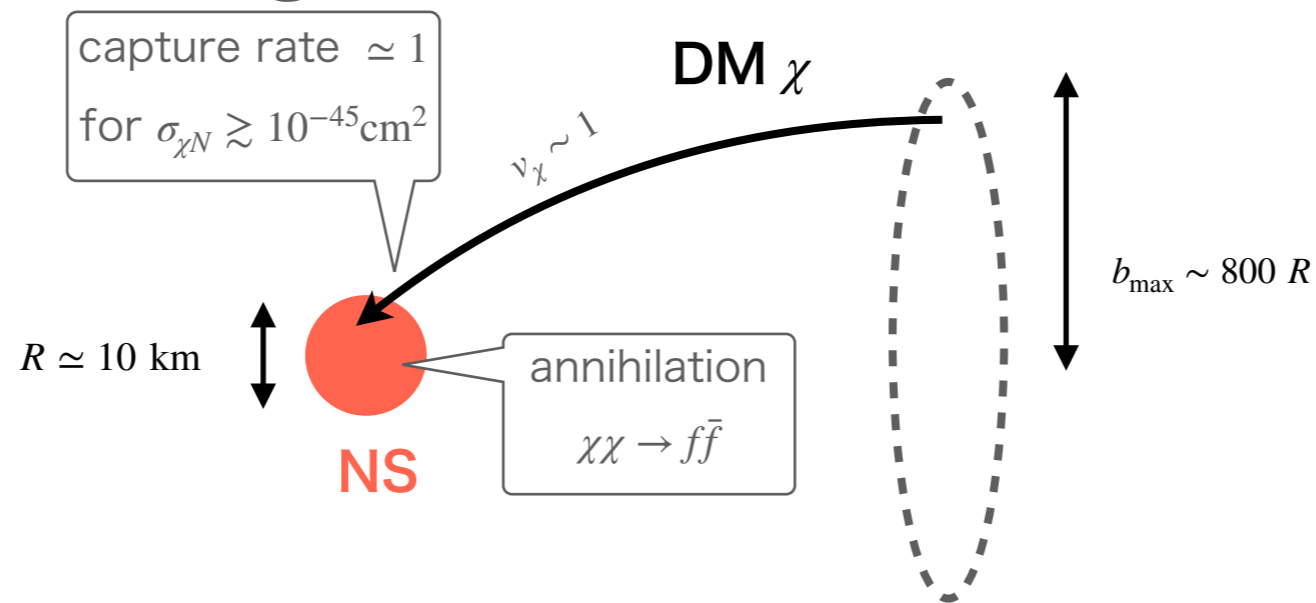
Figs. from Page et.al. 1302.6626

neutron singlet (1S_0)
neutron triplet (3P_2)
proton singlet (1S_0)



- Heat capacity C is suppressed. ($\sim e^{-\Delta/T}$)
- M.Urca luminosity $L_{\nu, MU}$ is suppressed. ($\sim e^{-\Delta/T}$)
- PBF occurs at $T < T_c$.
- It is also important for the "rotochemical heating" (see below).

(ii). NS Heating by DM



C. Kouvaris 0708.2362,
G. Bertone+ 0709.1485,
C. Kouvaris+ 1004.0586,
A. de Lavallaz+ 1004.0629
...

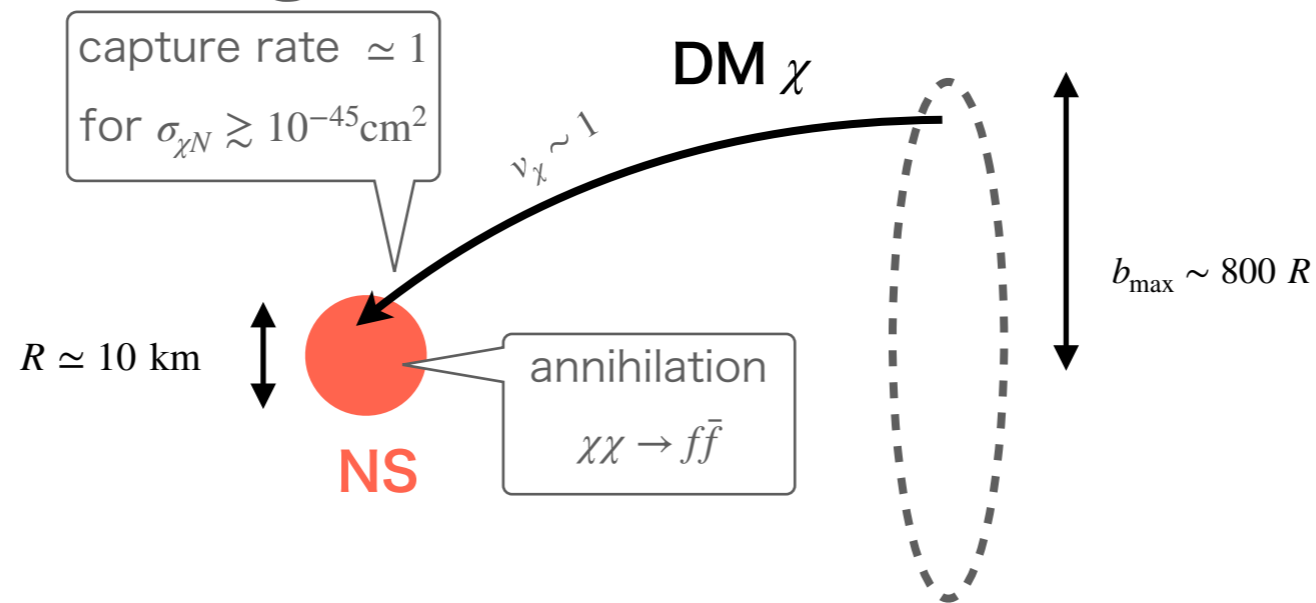
+ many recent works: e.g.,
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N. F. Bell+ 1807.02840
D. A. Camargo+ 1901.05474
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KH, N.Nagata, K.Yanagi 1905.02991
R. Garani+ 1906.10145
J. F. Acevedo+, 1911.06334
A. Joglekar+, 1911.13293

- velocity : $v_\chi^\infty \simeq 10^{-3} \rightarrow v_\chi^{\text{surface}} \sim 1$
- mean free path $L_{\text{mfp}} < R$ for $\sigma_{\chi N} \gtrsim 10^{-45} \text{cm}^2 \Rightarrow$ capture probability $\simeq 1$.
- For old NSs, accretion rate = annihilation rate.

$$\Rightarrow \text{energy injection } L_{\text{DM}}^{\text{heat}} \sim \dot{N} m_\chi \sim \pi b_{\text{max}}^2 \rho_\chi v_\chi^\infty \simeq 3 \times 10^{22} \text{erg/s} .$$

(independent of details of DM properties, such as mass, interactions,...)

(ii). NS Heating by DM



C. Kouvaris 0708.2362,
G. Bertone+ 0709.1485,
C. Kouvaris+ 1004.0586,
A. de Lavallaz+ 1004.0629
...

+ many recent works: e.g.,
J. Bramante+ 1703.04043
M. Baryakhtar+ 1704.01577
N. Raj+ 1707.09442
C.-S. Chen+ 1804.03409
N. F. Bell+ 1807.02840
D. A. Camargo+ 1901.05474
N. F. Bell+ 1904.09803
KH, N.Nagata, K.Yanagi 1905.02991
R. Garani+ 1906.10145
J. F. Acevedo+, 1911.06334
A. Joglekar+, 1911.13293

NS probe is advantageous in the sense that...

☑ $v_\chi \sim 1$ at NS surface

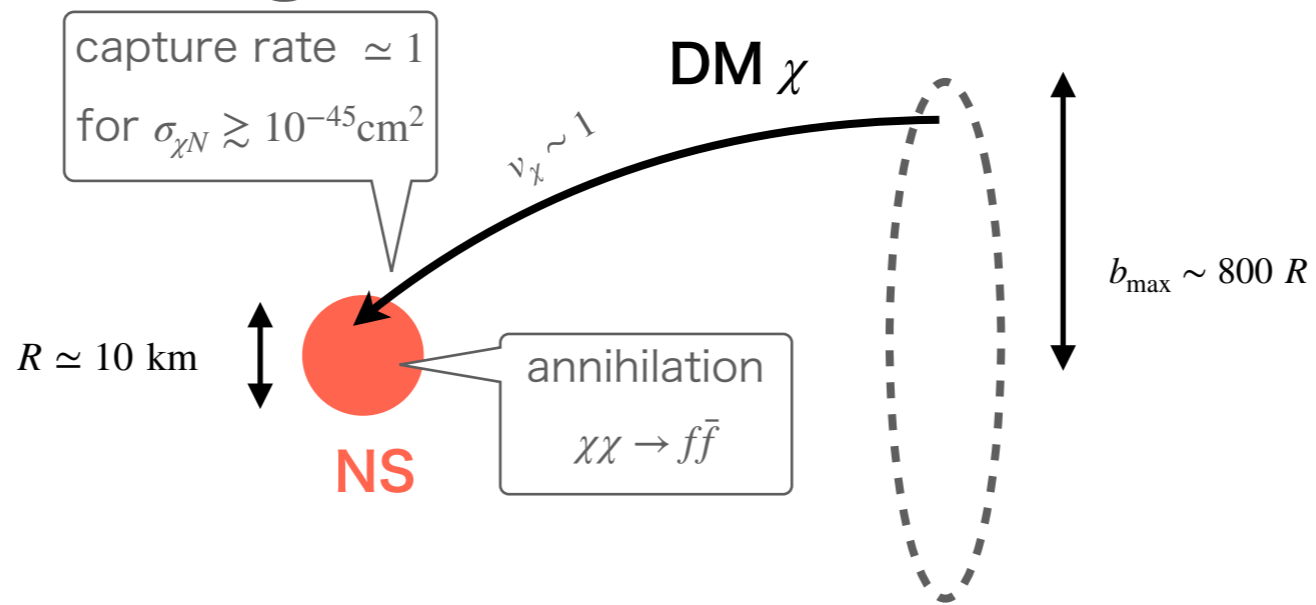
-> It is also sensitive to inelastic scattering (e.g., pure-Wino: $\chi^0 + N \rightarrow \chi^- + N'$) or other velocity-suppressed scatterings.

☑ It can also probe light DM ($\ll 1 \text{ GeV}$).

☑ It is also sensitive to lepton coupling (e and mu).

☑ In principle, it can go beyond the neutrino floor.

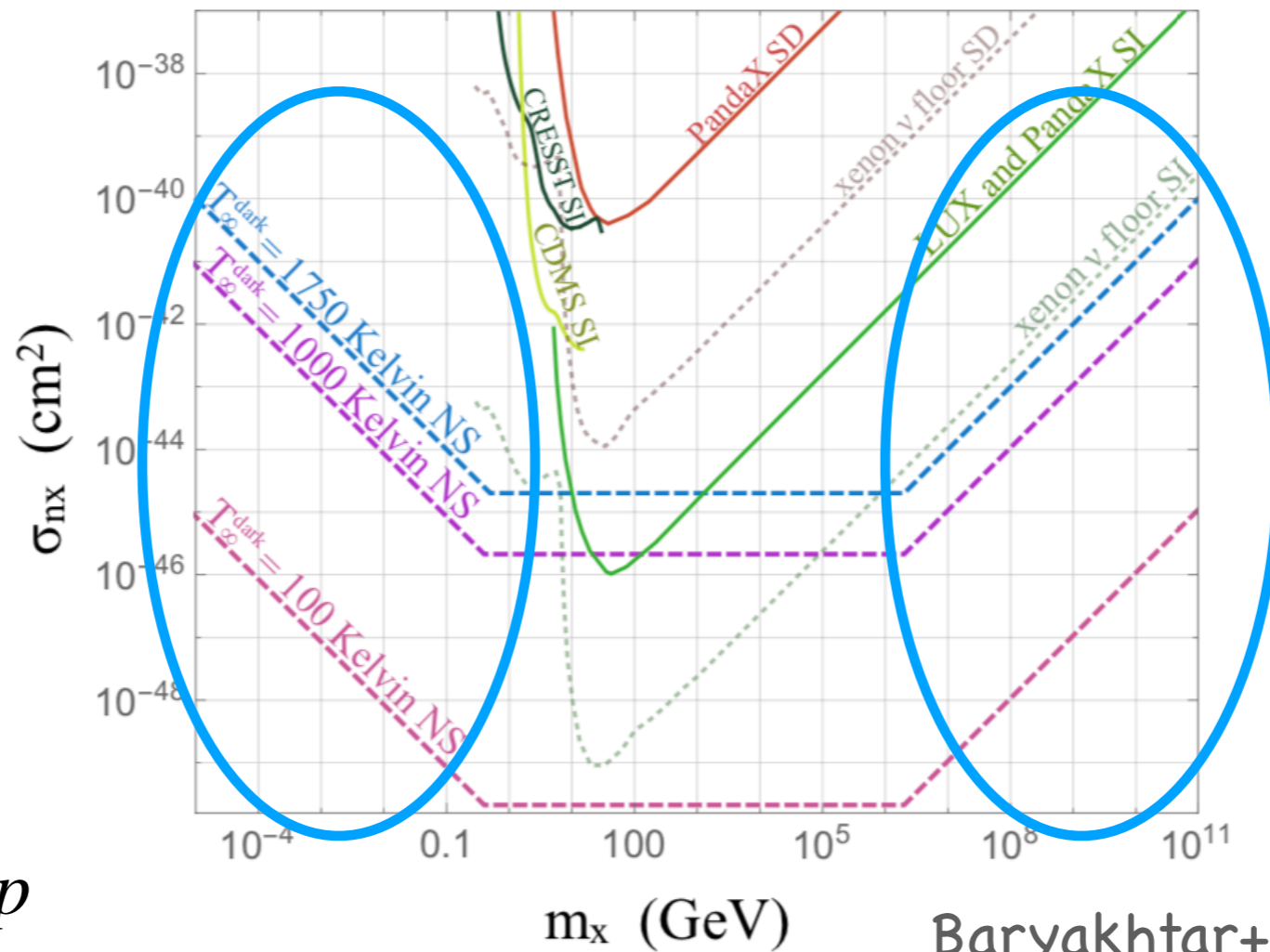
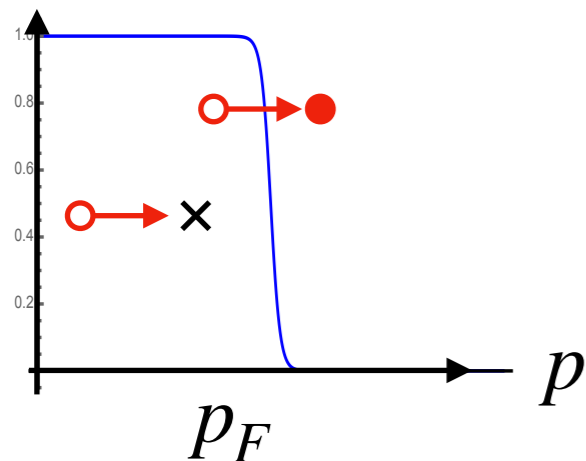
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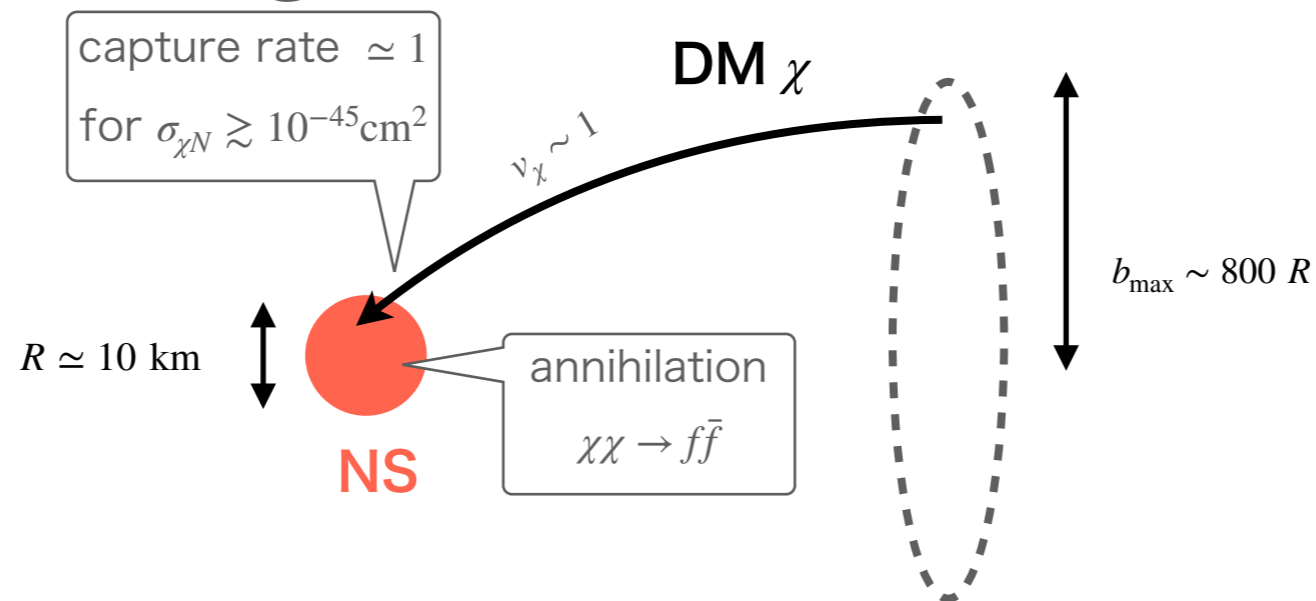
For $m_\chi < 1 \text{ GeV}$,
Pauli blocking prevents
scattering of low p
neutrons.



For $m_\chi > 100 \text{ TeV}$,
single scattering is not
enough to catch DM.
→ multiple scattering.

Baryakhtar+, 1704.01577

(ii). NS Heating by DM



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But,...

Old warm neutron stars?

Recently, “old but warm neutron stars” have been observed.

Milli-second pulsars

▶ J0437-4715: $t_{\text{sd}} = (6.7 \pm 0.2) \times 10^9$ years, $T_s^\infty = (1.25 - 3.5) \times 10^5$ K

O. Kargaltsev, G. G. Pavlov, and R. W. Romani, *Astrophys. J.* **602**, 327 (2004);
M. Durant, *et al.*, *Astrophys. J.* **746**, 6 (2012).

▶ J2124-3358: $t_{\text{sd}} = 11_{-3}^{+6} \times 10^9$ years, $T_s^\infty = (0.5 - 2.1) \times 10^5$ K

B. Rangelov, *et al.*, *Astrophys. J.* **835**, 264 (2017).

Ordinary pulsars

▶ J0108-1431: $t_{\text{sd}} = 2.0 \times 10^8$ years, $T_s^\infty = (1.1 - 5.3) \times 10^5$ K

R. P. Mignani, G. G. Pavlov, and O. Kargaltsev, *Astron. Astrophys.* **488**, 1027 (2008).

▶ B0950+08: $t_{\text{sd}} = 1.75 \times 10^7$ years, $T_s^\infty = (1 - 3) \times 10^5$ K

G. G. Pavlov, *et al.*, *Astrophys. J.* **850**, 79 (2017).

These observations cannot be explained in the standard cooling.

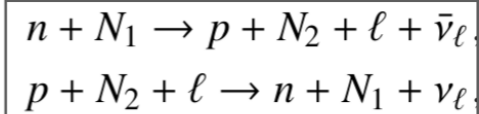
(iii). Rotochemical Heating $C \frac{dT}{dt} = -L_\nu - L_\gamma + L_{\text{heat rotochemical}}$

- Recently, we have updated the calculation, including for the first time

- ☑ both neutron superfluidity and proton superconductivity
- ☑ with radius dependence
- ☑ with temperature dependence
- ☑ with angular dependence (for neutron triplet pairing)

simultaneously. (K. Yanagi, N. Nagata, KH, [arXiv:1904.04667])

non-equilibrium



chemical potential

$$\eta_\ell \equiv \mu_n - \mu_p - \mu_\ell$$

$$C \frac{dT^\infty}{dt} = -L_\nu^\infty - L_\gamma^\infty + L_H^\infty$$

$$\begin{aligned} \frac{d\eta_e^\infty}{dt} &= - \sum_{N=n,p} \int dV (Z_{npe} \Delta\Gamma_{M,Ne} + Z_{np} \Delta\Gamma_{M,N\mu}) e^{\Phi(r)} + 2W_{npe} \Omega \dot{\Omega}, \\ \frac{d\eta_\mu^\infty}{dt} &= - \sum_{N=n,p} \int dV (Z_{np} \Delta\Gamma_{M,Ne} + Z_{np\mu} \Delta\Gamma_{M,N\mu}) e^{\Phi(r)} + 2W_{np\mu} \Omega \dot{\Omega} \end{aligned}$$

$$L_{\nu,M}^\infty = \sum_{\ell=e,\mu} \sum_{N=n,p} \int dV Q_{M,N\ell} e^{2\Phi(r)}$$

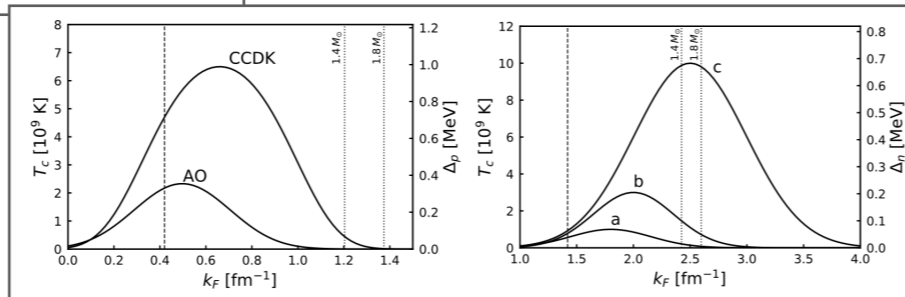
$$Q_{M,N\ell} = Q_{M,N\ell}^{(0)} I_{M,\epsilon}^N$$

$$L_H^\infty = \sum_{\ell=e,\mu} \sum_{N=n,p} \int dV \eta_\ell \cdot \Delta\Gamma_{M,N\ell} e^{2\Phi(r)}$$

$$\Delta\Gamma_{M,N\ell} = \frac{Q_{M,N\ell}^{(0)}}{T(r)} I_{M,\Gamma}^N$$

$$I_{M,\epsilon}^N = \frac{60480}{11513\pi^8} \frac{1}{A_0^N} \int \prod_{j=1}^5 \frac{d\Omega_j}{4\pi} \int_0^\infty dx_\nu \int_{-\infty}^\infty dx_n dx_p dx_{N_1} dx_{N_2} x_\nu^3 \cdot f(z_n) f(z_p) f(z_{N_1}) f(z_{N_2}) \times [f(x_\nu - \xi_\ell - z_n - z_p - z_{N_1} - z_{N_2}) + f(x_\nu + \xi_\ell - z_n - z_p - z_{N_1} - z_{N_2})] \delta^3\left(\sum_{j=1}^5 \mathbf{p}_j\right), \quad (18)$$

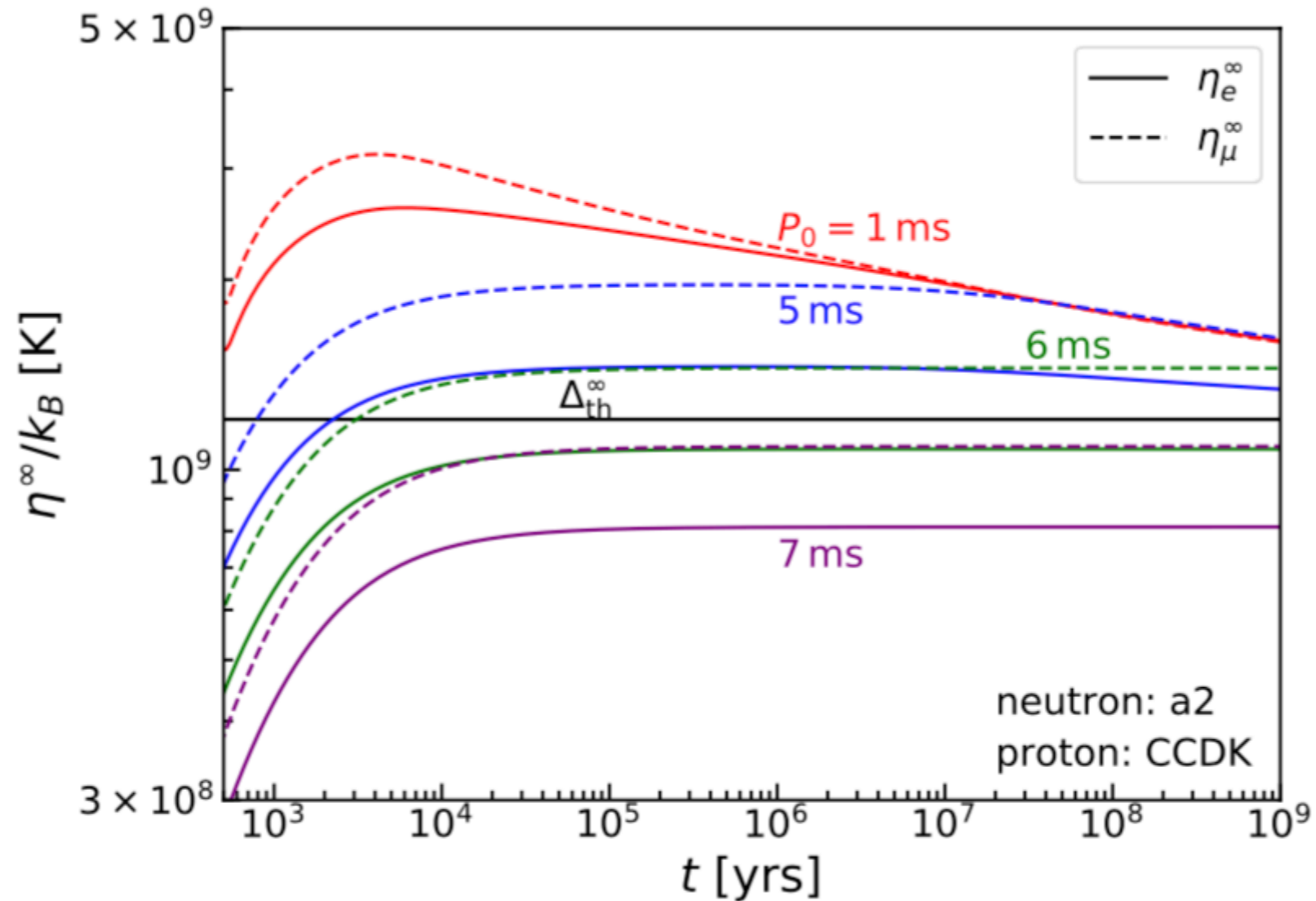
$$I_{M,\Gamma}^N = \frac{60480}{11513\pi^8} \frac{1}{A_0^N} \int \prod_{j=1}^5 \frac{d\Omega_j}{4\pi} \int_0^\infty dx_\nu \int_{-\infty}^\infty dx_n dx_p dx_{N_1} dx_{N_2} x_\nu^2 \cdot f(z_n) f(z_p) f(z_{N_1}) f(z_{N_2}) \times [f(x_\nu - \xi_\ell - z_n - z_p - z_{N_1} - z_{N_2}) - f(x_\nu + \xi_\ell - z_n - z_p - z_{N_1} - z_{N_2})] \delta^3\left(\sum_{j=1}^5 \mathbf{p}_j\right), \quad (19)$$



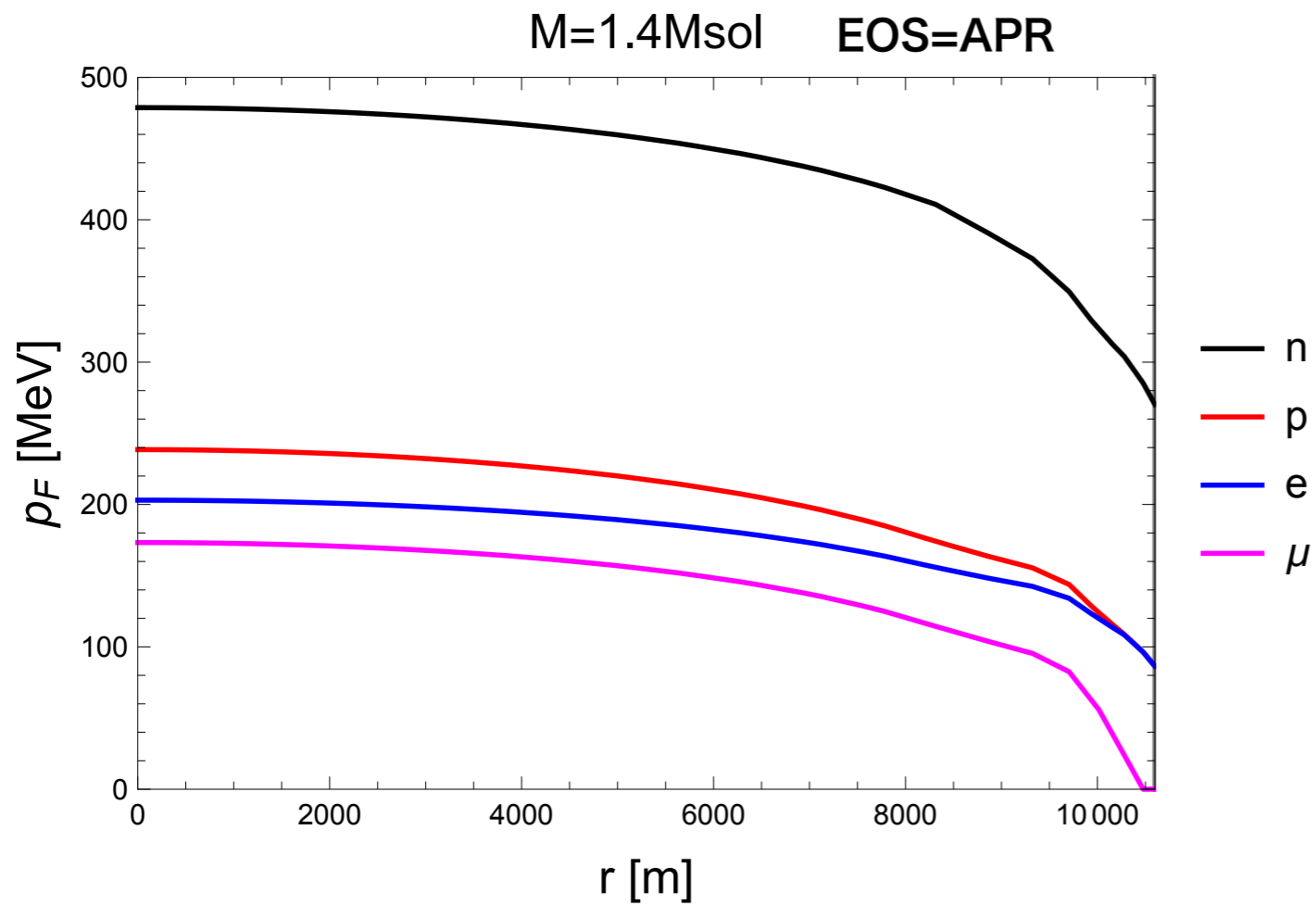
(iv). DM heating vs. Rotochemical heating

$$C \frac{dT}{dt} = -L_\nu - L_\gamma + L_{\text{heat rotochemical}} + L_{\text{heat DM}}$$

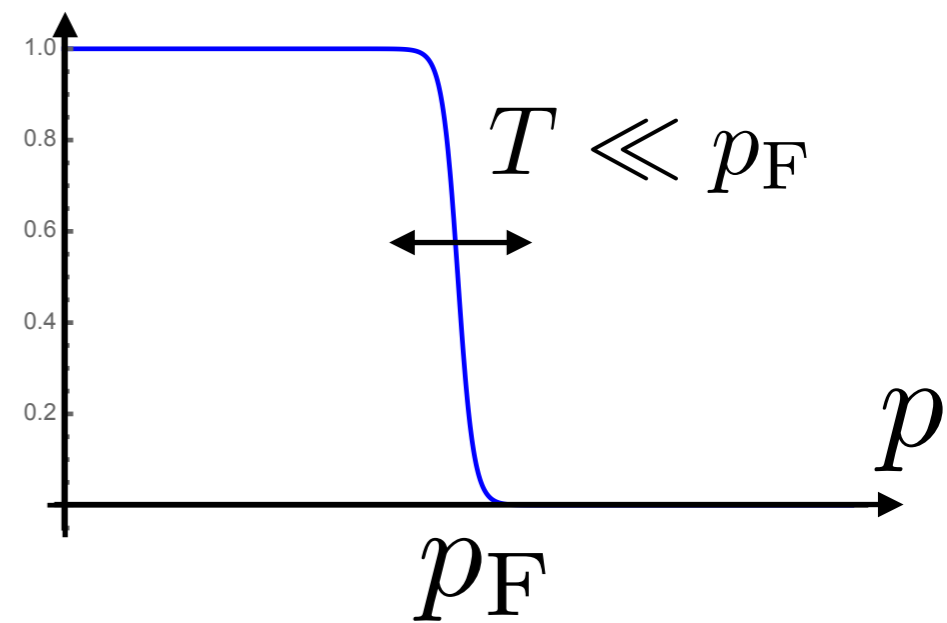
KH, N. Nagata, K. Yanagi, [1905.02991]



$M = 1.4M_\odot$
 $P = 1$ s
 $\dot{P} = 10^{-15}$



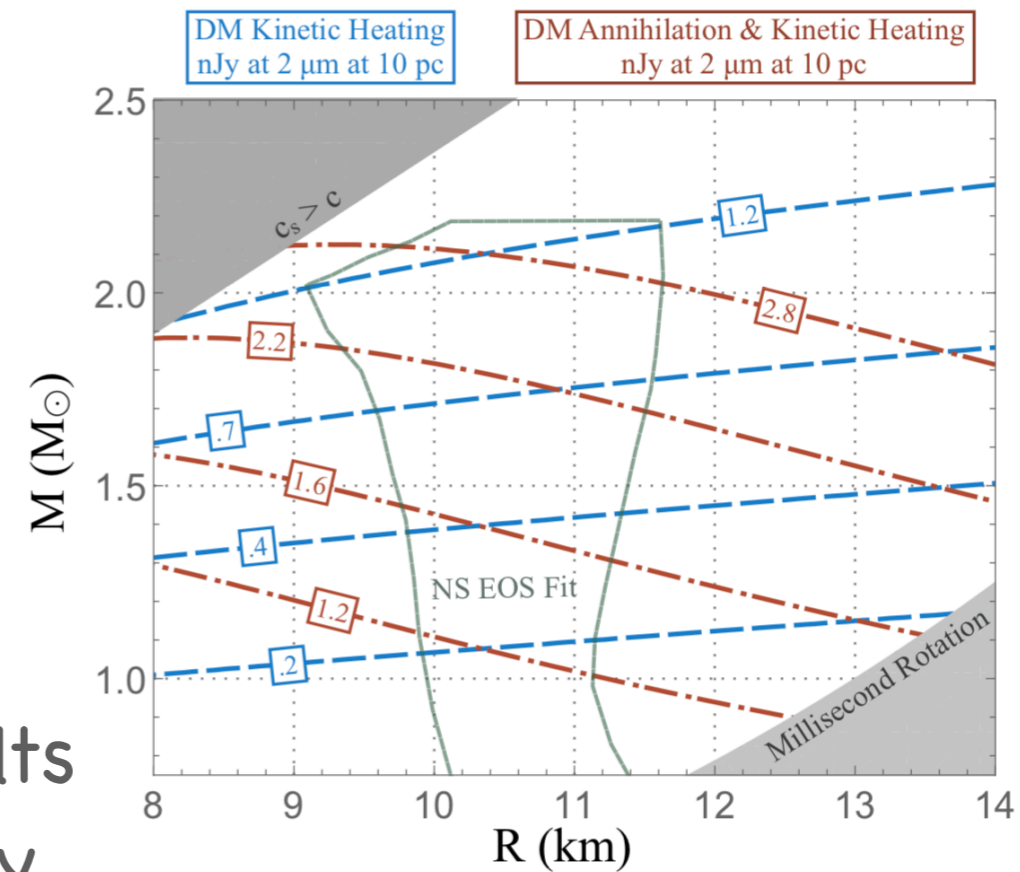
Fermi momentum $\sim \mathcal{O}(100)$ MeV



Fermi degenerate

observational feasibility

- See e.g., the discussion in M.Baryakhtar+, 1704.01577.
- $O(1)$ old and cold NSs can be at $d = 10\text{pc}$.
- Radiation from a DM-heated NS there results in a spectral flux density of $O(1)$ nanoJansky (nJy) at wavelength $\nu^{-1} = \mathcal{O}(1) \mu\text{m}$.
- Maybe within the sensitivity of the upcoming telescopes such as the JWST, TMT, and E-ELT.



M.Baryakhtar+, 1704.01577

$$T_s \sim T^\alpha \quad (\alpha \sim 0.5)$$

