

The light dark matter search—

What's going on at  
an MeV?

Tom Melia, Kavli IPMU

Cosmic Acceleration, Kavli IPMU 17th Feb 2020

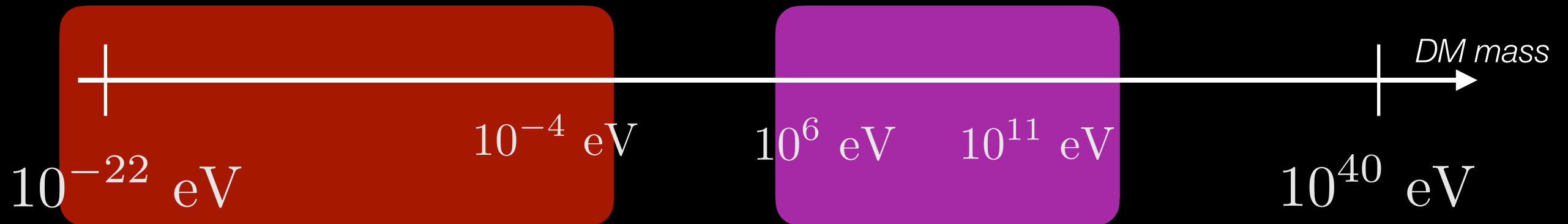
The search for dark matter interactions in earth-based detection experiments must confront a possible mass range that spans more than sixty orders of magnitude



*This necessarily makes search strategies diverse*

*Axion-like experiments: exploit coherence with resonance*

*WIMP-like experiments: energy deposits*



*Sub-halo structure*

GeV DM mass  
→

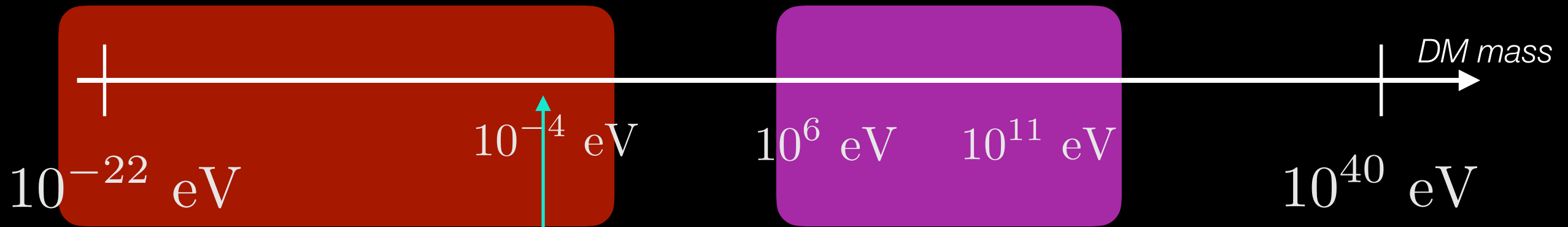
*1 DM through earth in a year*

*ADMX/CAPP*      Cavity  
*CASPEr*      NMR  
*Atomic Clocks*  
*LIGO*      Interferometers  
...

*XENON*  
*LUX*      Liquid Xe (recoil)  
*PANDAX*  
*CDMS*      Ge / Si  
*PICO*      C3F8  
...

*Axion-like experiments: exploit coherence with resonance*

*WIMP-like experiments: energy deposits*



*Sub-halo structure*

GeV DM mass  
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*1 DM through earth in a year*

*ADMX/CAPP*

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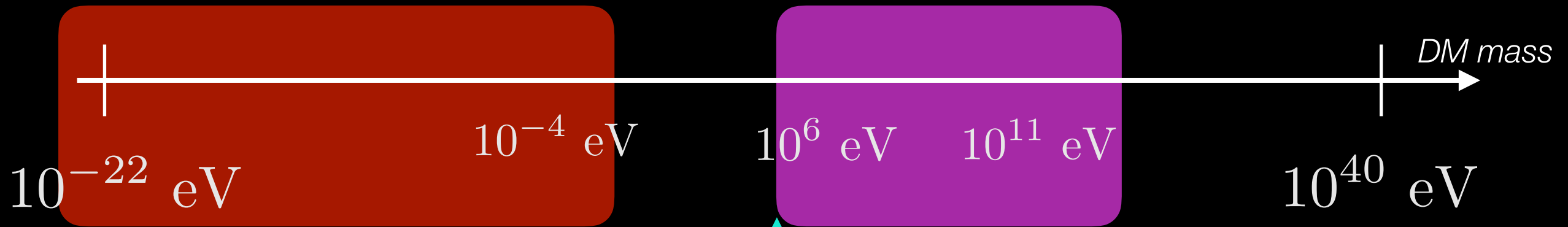
Interferometers

...

→  
*Spatial and temporal coherence becomes unusable*

*Axion-like experiments: exploit coherence with resonance*

*WIMP-like experiments: energy deposits*



*Sub-halo structure*

*1 DM through earth in a year*

MeV DM mass  
→

sub-GeV mass  
being probed

*CRESST*

*superCDMS*

*DarkSide (electron)*

*Xenon10 (electron)*

*SENSEI*

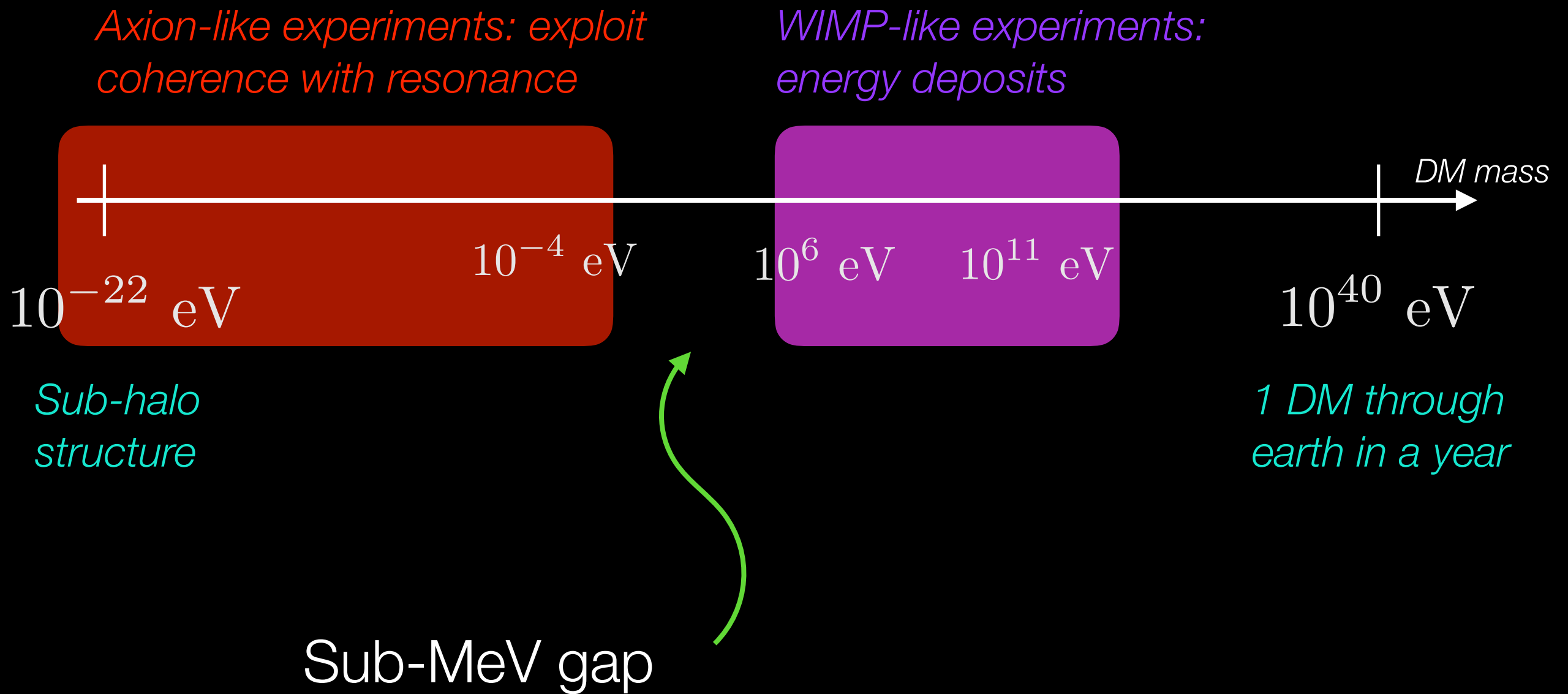
Phonons/  
ionization

Ionization of Ar,  
Xe electrons

Electrons in  
semi-conductor

$$\frac{1}{2}m_{DM}v^2 \sim 10^{-6}m_{DM}$$

→  
*Few-10s eV  
energy deposit*



What's going on at an  
MeV?

What's going  
on at an  
MeV?



```
graph TD; A((What's going on at an MeV?)) --> B[1]; A --> C[2]; A --> D[3]; A --> E[4]; A --> F[5];
```

1

2

3

4

5

What's going  
on at an  
MeV?

Sensitivity gap in direct  
detection meV-eV energy  
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$$\frac{1}{2}m_{DM}v^2 \sim 10^{-6}m_{DM}$$

i.e. keV-MeV mass

2

3

4

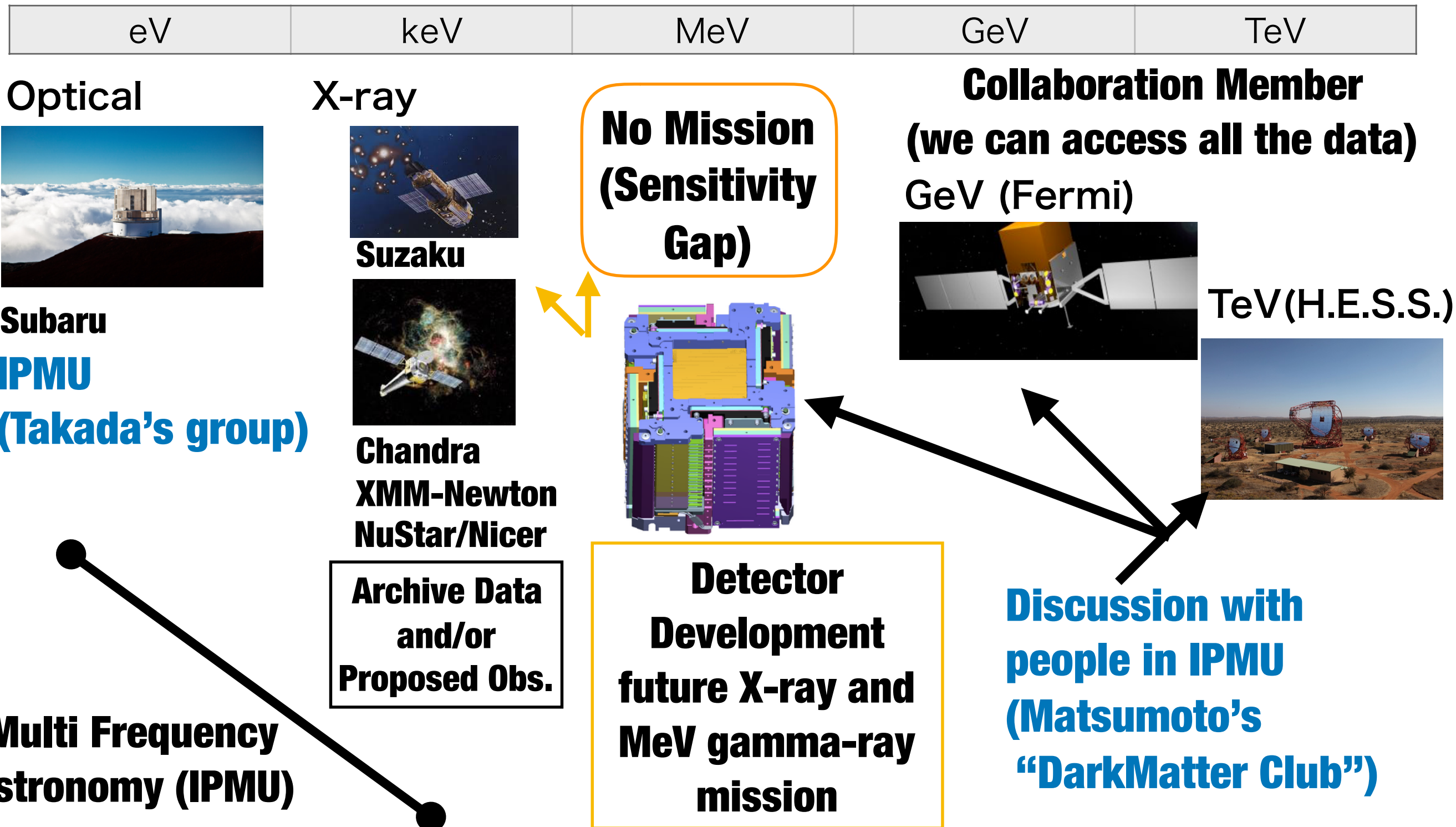
5



# High Energy Astrophysics

## 1&3: Analysis and Detector Development

### Our Activity (High Energy Astrophysics at IPMU)



## High Energy Astrophysics

**1&3: Analysis and Detector Development****Our Activity (High Energy Astrophysics at IPMU)**

## High Energy Astrophysics

**1&3: Analysis and Detector Development****Our Activity (High Energy Astrophysics at IPMU)**

Optical



Subaru

**IPMU**  
**(Takada's group)**

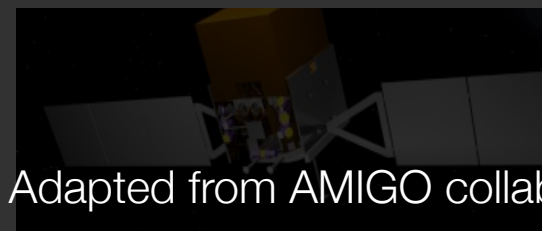
X-ray



Suzaku

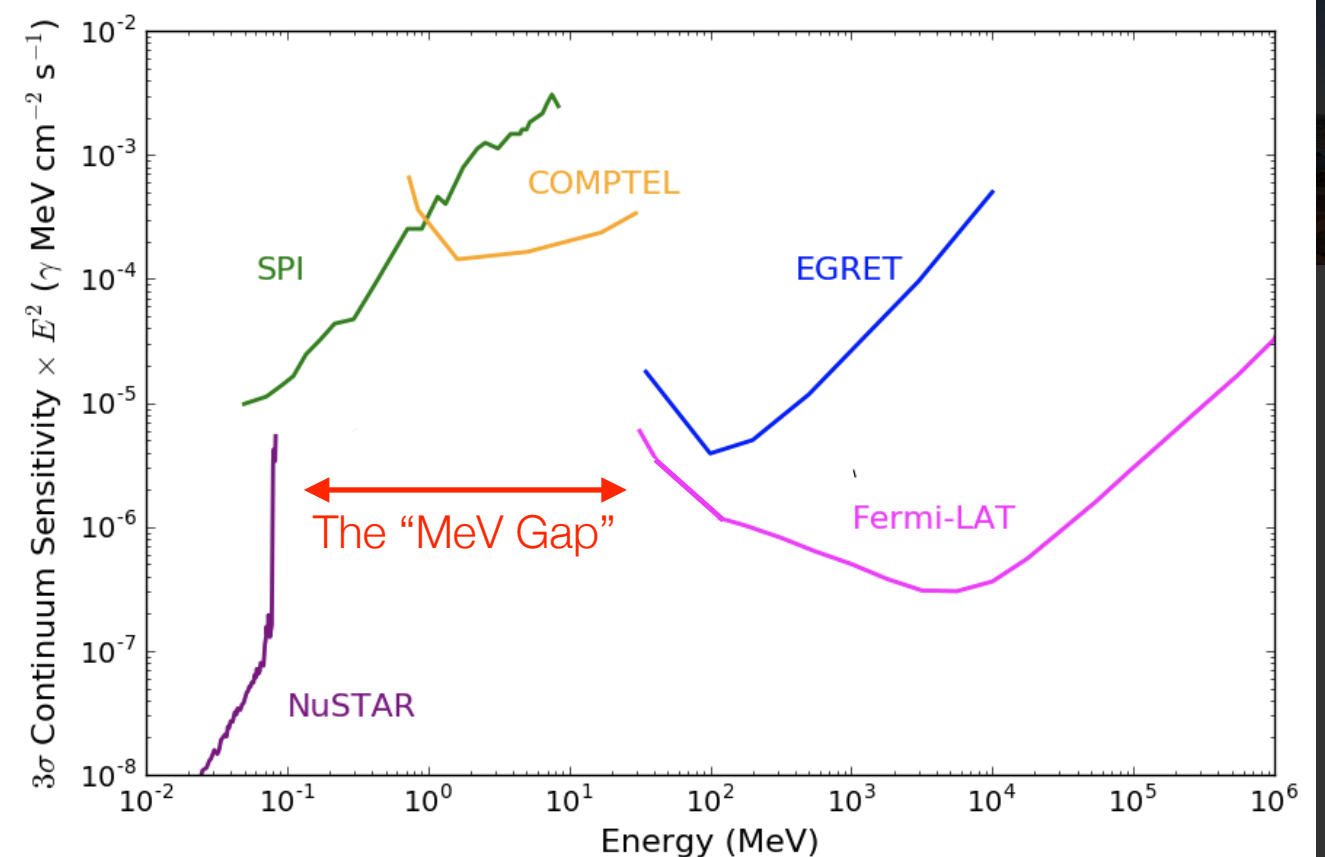

**Chandra**  
**XMM-Newton**  
**NuStar/Nicer**
**Archive Data**  
**and/or**  
**Proposed Obs.**
**No Mission**  
**(Sensitivity**  
**Gap)**
**Collaboration Member****(we can access all the data)**

GeV (Fermi)



TeV (H.E.S.S.)

Adapted from AMIGO collaboration 1907.07558

**Multi Frequency**  
**Astronomy (IPMU)**


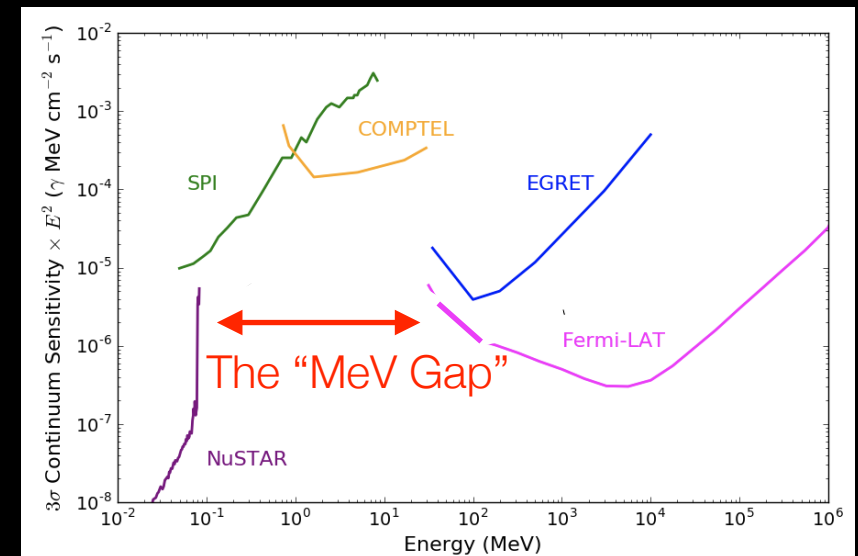
# What's going on at an MeV?

Sensitivity gap in direct detection meV-eV energy deposit

$$\frac{1}{2}m_{DM}v^2 \sim 10^{-6}m_{DM}$$

i.e. keV-MeV mass

Telescope sensitivity gap



3

4

5

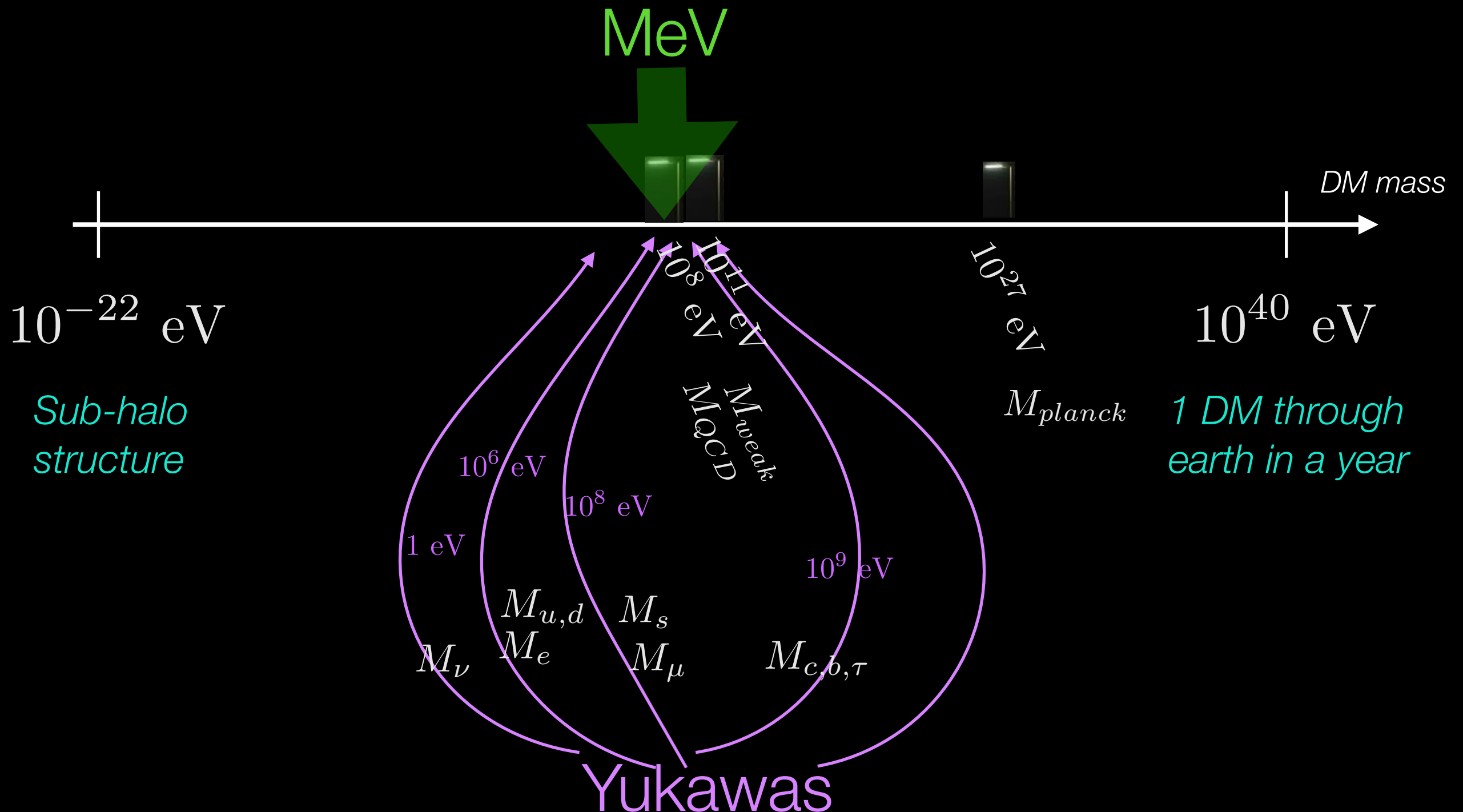
There exist models of keV-MeV DM that give 100% of DM abundance, not ruled out by astrophysics/cosmology, and are detectable in earth based experiments

e.g. [Green, Rajendran 1701.08750](#); [Knapen, Lin, Zurek 1709.07882](#)

Relatively un-developed/explored theory space, e.g. using relatively (past decade) new mechanism of freeze in for relic abundance

Plus, we are being forced to reconsider our ideas about naturalness & the weak scale, and what yardstick we should judge models by

(perhaps worth saying:  $\sim \text{MeV}$  is a pretty well motivated place to search, given we know particles exist there)



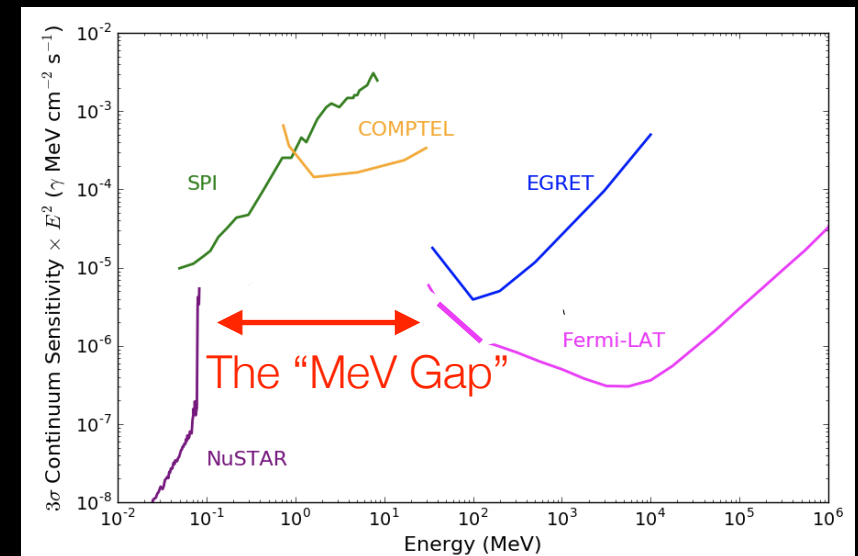
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Theory  
motivation

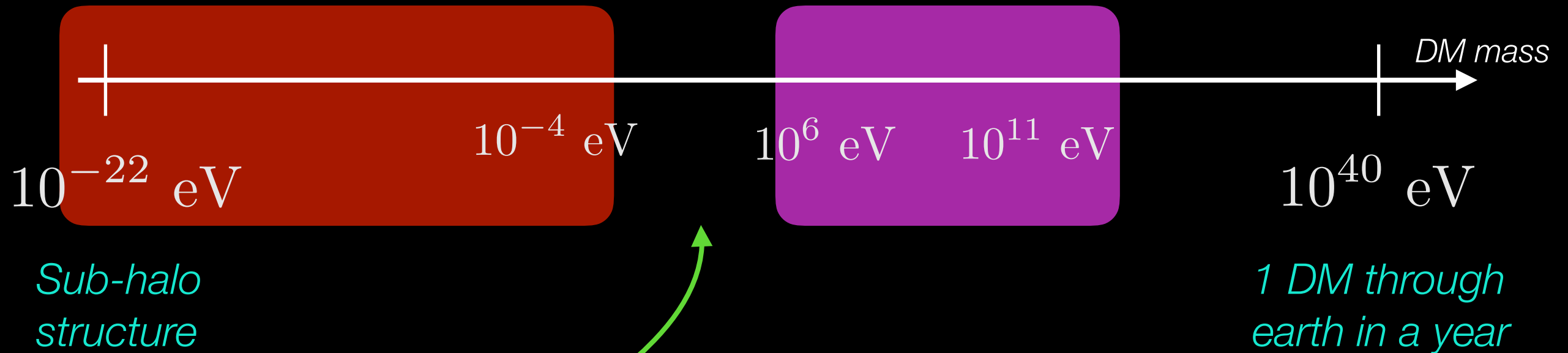
4

5



*Axion-like experiments: exploit coherence with resonance*

*WIMP-like experiments: energy deposits*



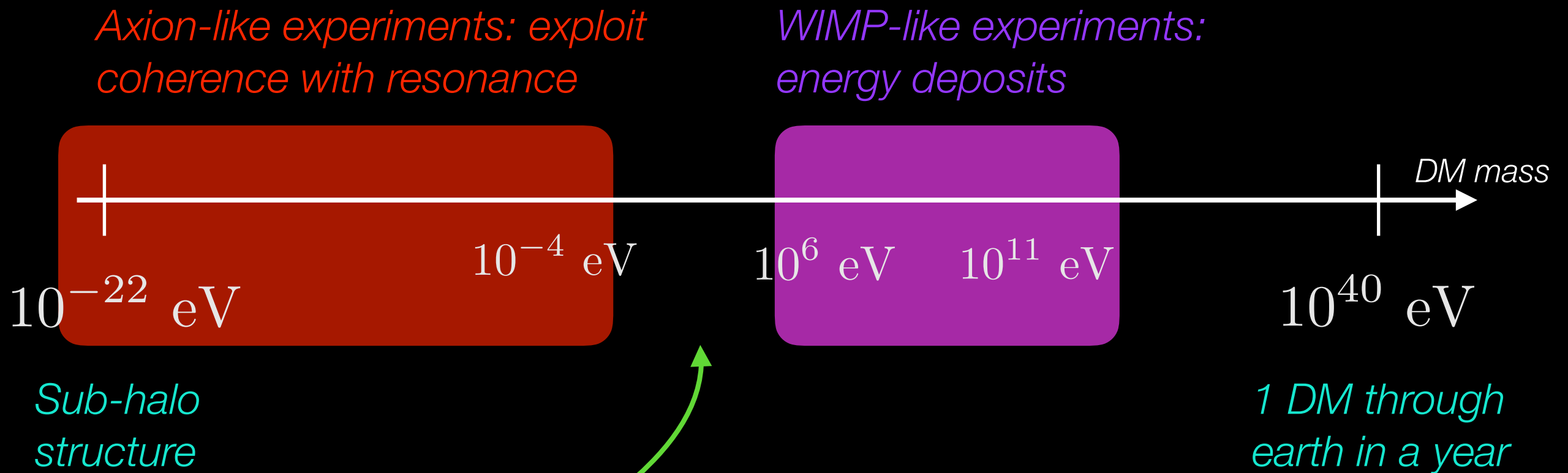
Sub-MeV window

*Energy deposits < eV*

This is the typical energy of quantum excitations of condensed matter systems, and is the energy scale of chemistry

Matching of energy scales suggests such systems should be good targets





## Sub-MeV window

*Energy deposits < eV at proposal/R&D stage*

### *Superfluid He*

Schutz and Zurek, Phys.Rev.Lett. 117 (2016) no.12, 121302

### *Superconductors*

Hochberg, Zhao, and Zurek, Phys.Rev.Lett. 116 (2016) no.1, 011301

### *Magnetic bubble chambers*

Bunting, Gratta, TM, and Rajendran, Phys.Rev. D95 (2017) no.9, 095001

### *Sapphire/ polar crystals*

Knapen, Lin, Pyle, and Zurek, Phys.Lett. B785 (2018) 386-390

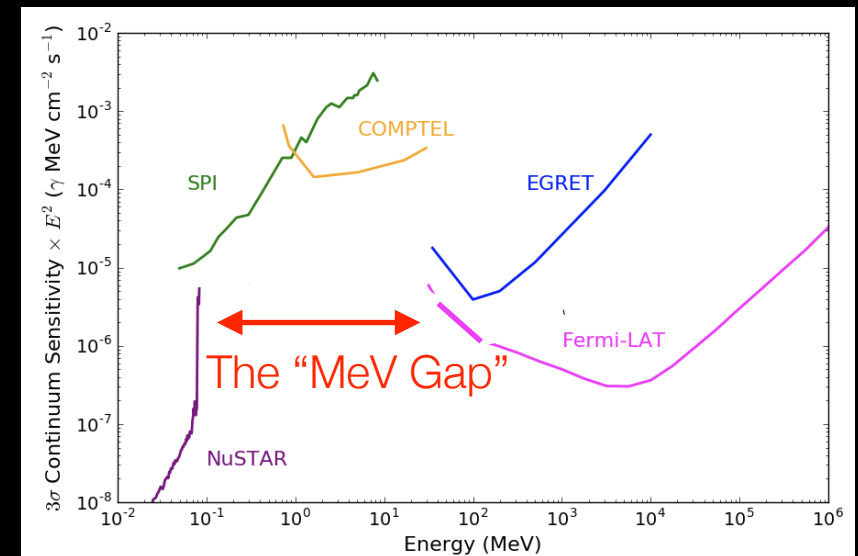
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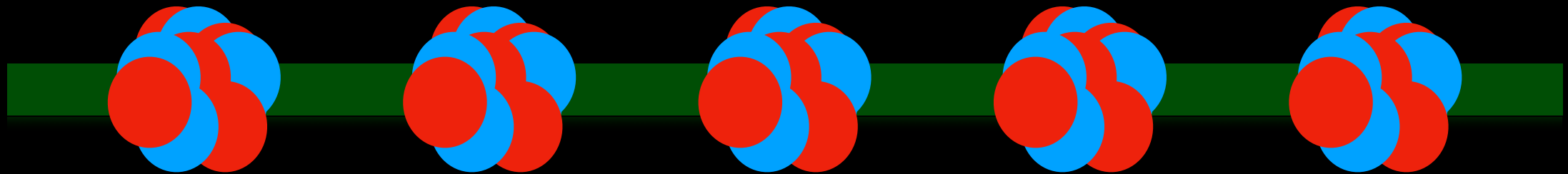


Theory  
motivation

Condensed matter,  
quantum devices,  
chemistry

*New technologies  
under R&D, with  
broad applications*

DM de Broglie wavelength  $>$  interatomic spacing



We are used to characterising DM via its interaction with **particles**. Here, DM doesn't resolve individual nuclei/atoms...

At these energies in materials, **phonons** (crystal/molecule vibrations) are the relevant **particle** description

Symmetries perhaps less familiar, so can expect novel features in DM interactions

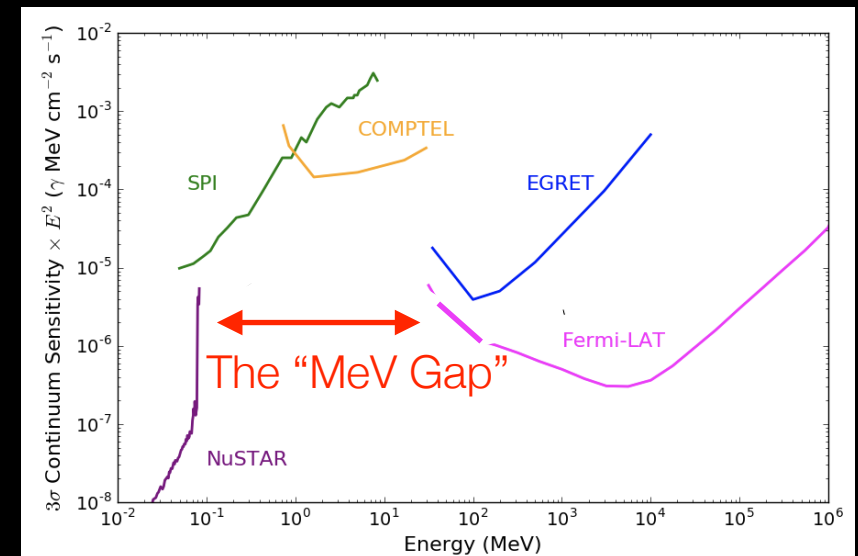
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*New technologies  
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*New theoretical input needed to determine  
rate and nature of DM interaction*

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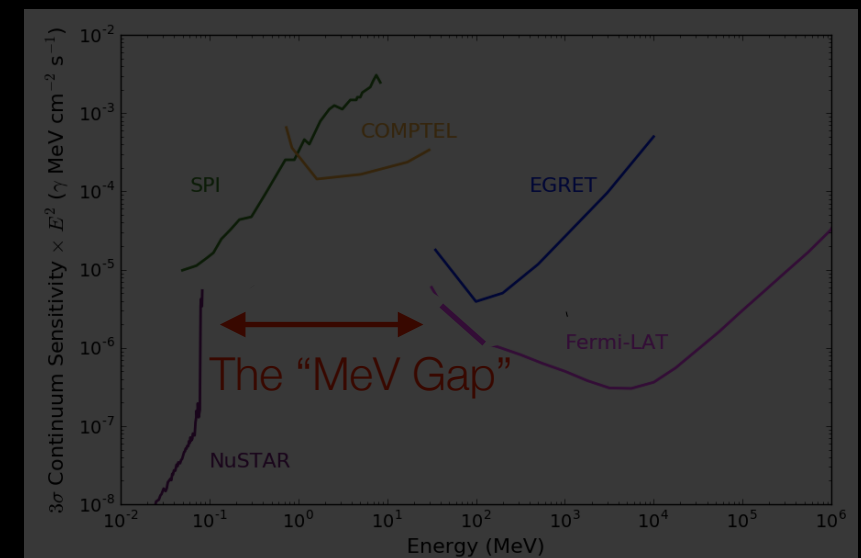
*Remainder of this talk I'll focus on this*

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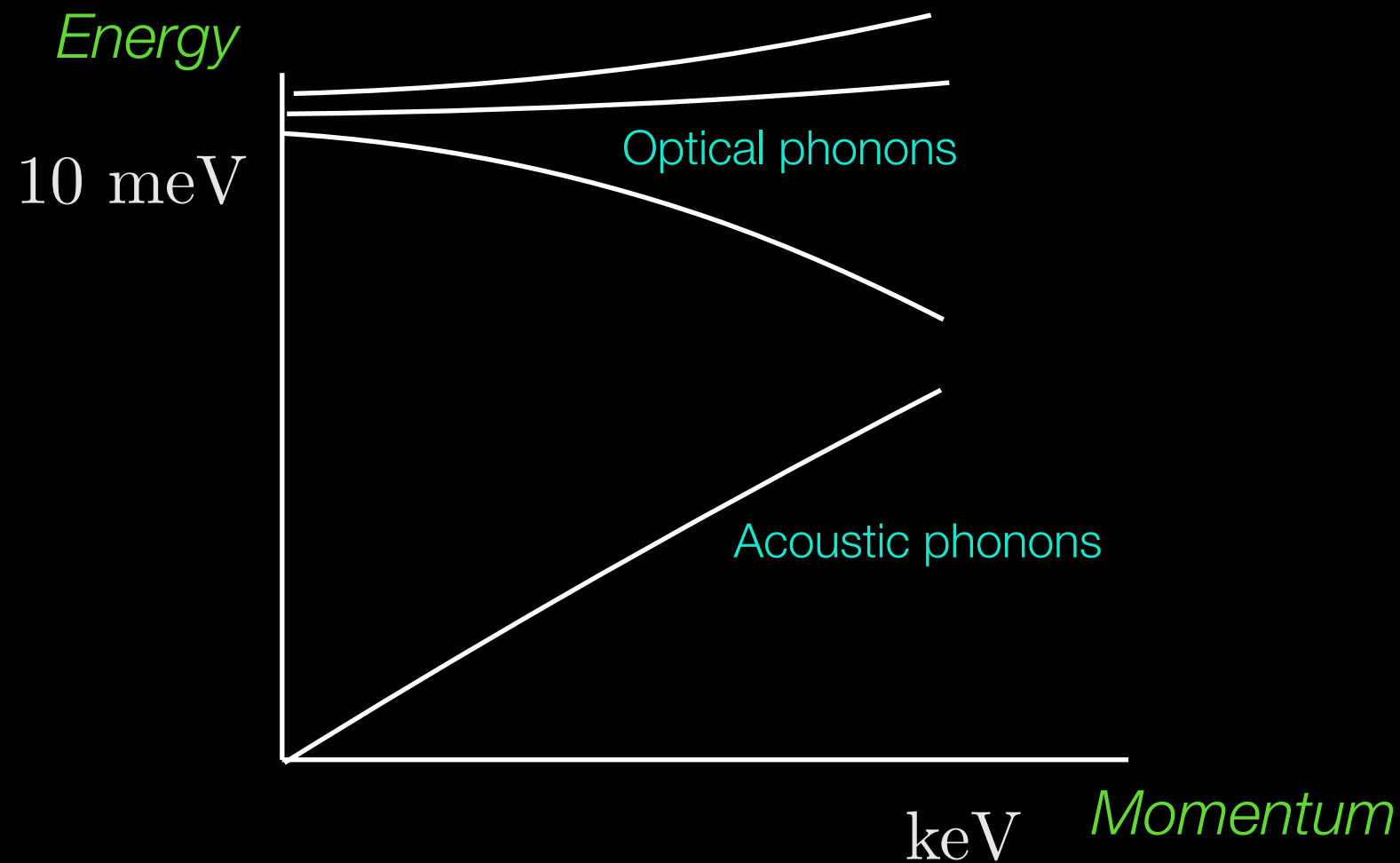
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rate and nature of DM interaction*

# The Dark Matter - Phonon Coupling

Phonons are quanta of vibrations, and are prototypical excitations in condensed matter systems

In crystals and molecules, it turns out that the dark matter-phonon coupling has a curious feature, with implications for direct detection experiments

# Phonon band-structure of crystals



DM energy scales

$$M_{DM} \quad q_{DM} = M_{DM}v \quad E_{DM} = \frac{1}{2}M_{DM}v^2$$

MeV

keV

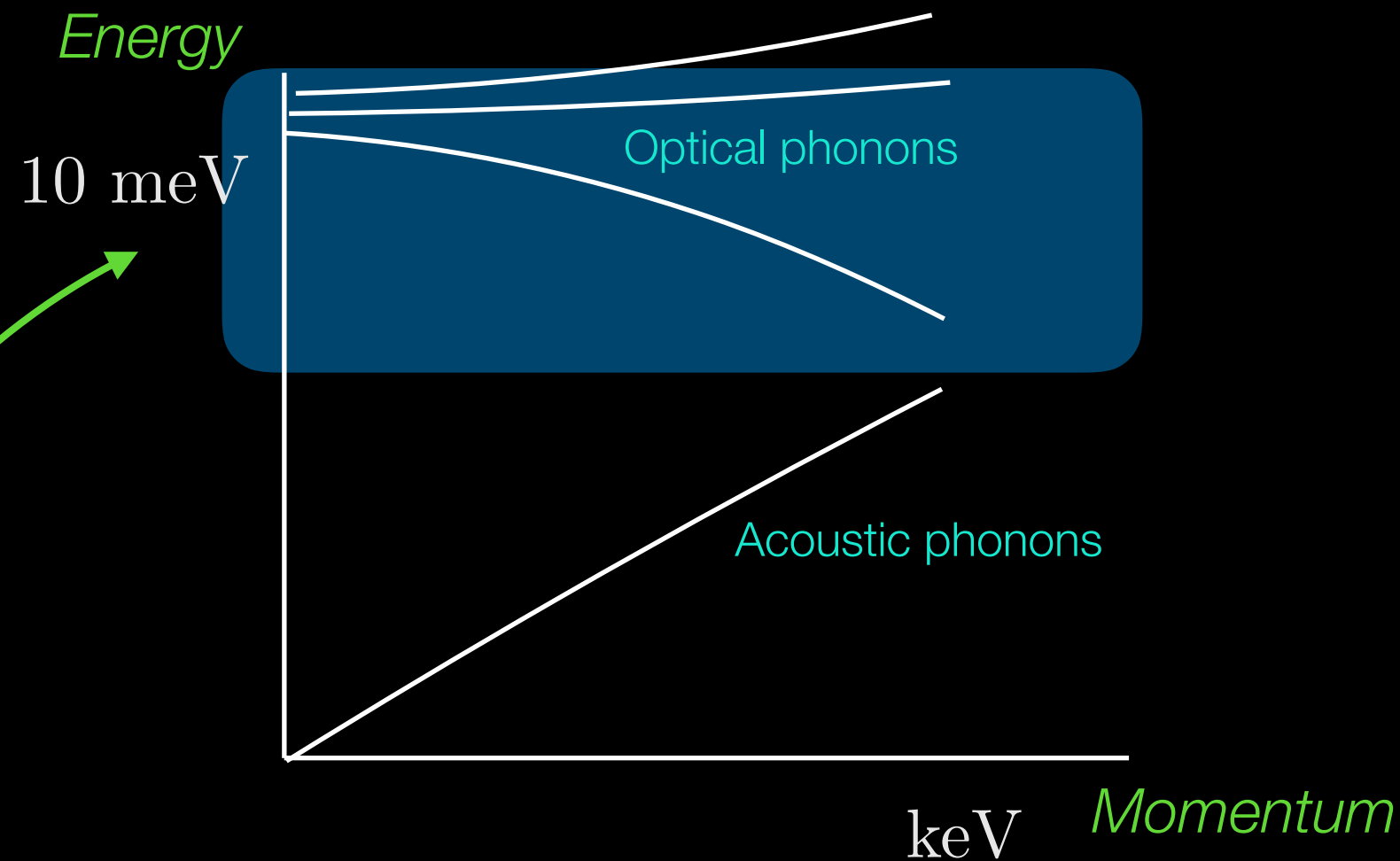
eV

keV

eV

meV

# Phonon band-structure of crystals



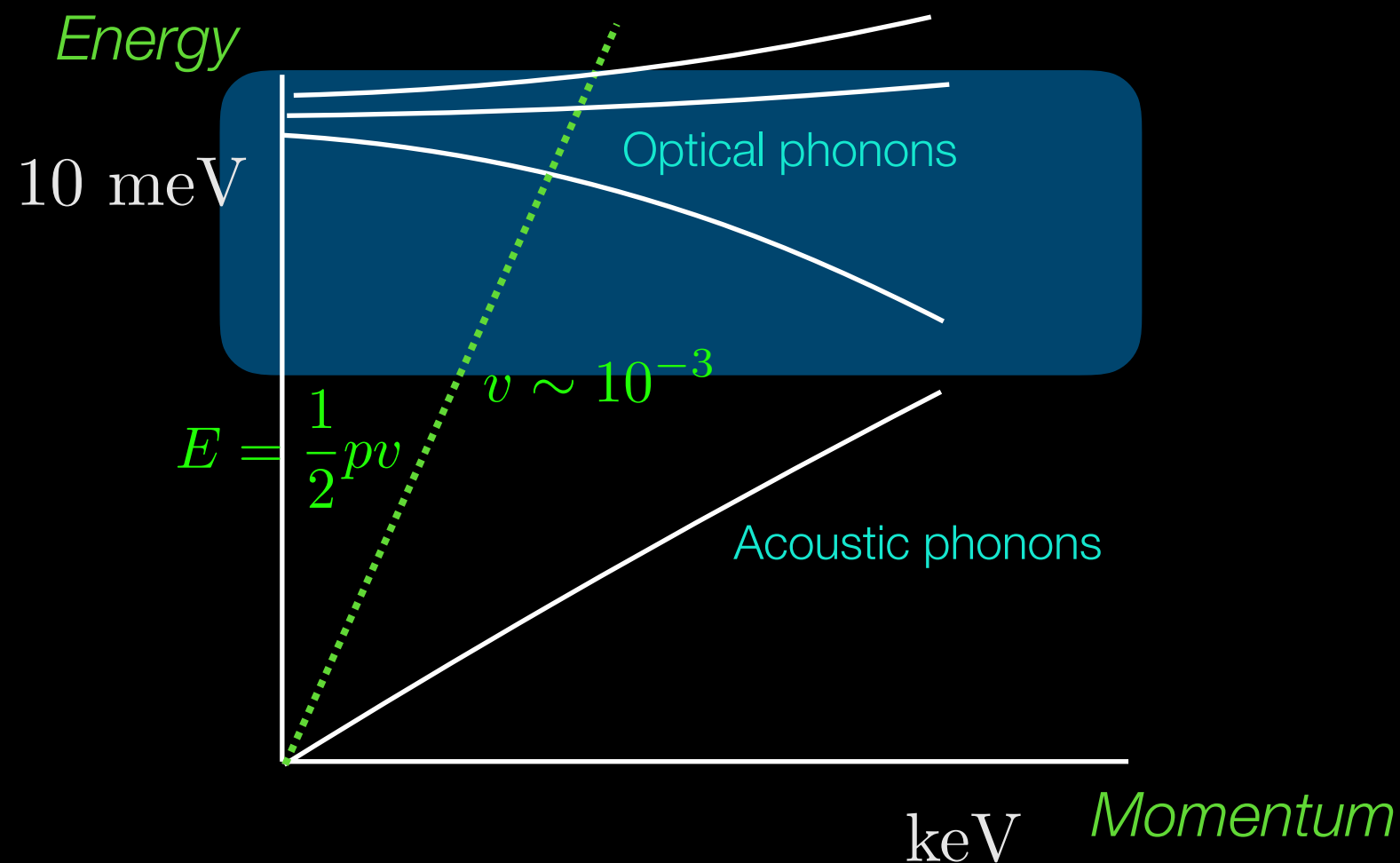
Optical phonons have: [Knapen, Lin, Pyle, and Zurek; Phys. Lett. B785 2018 386-390](#)

Higher threshold energy

Can have large density of states



# Phonon band-structure of crystals



Optical phonons have: [Knapen, Lin, Pyle, and Zurek; Phys. Lett. B785 2018 386-390](#)

Higher threshold energy

Can have large density of states

Kinematics to match DM kinematics

# The coupling-to-mass effect in DM-phonon scattering

Cox, TM, and Rajendran; [arxiv 1905.05575](#)

When the coupling of the DM to each atom is proportional to the atom's mass, the leading order scattering off optical phonons vanishes at small momentum transfer

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Cox, TM, and Rajendran; arxiv 1905.05575

When the coupling of the DM to each atom is proportional to the atom's mass, the leading order scattering off optical phonons vanishes at small momentum transfer

Coupling proportional to mass is generic:

Low momentum transfer  $\rightarrow$  coherence over atom

Couplings away from the unitarity limit  $\rightarrow$  effective coupling scales with number of constituents

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Cox, TM, and Rajendran; arxiv 1905.05575

When the coupling of the DM to each atom is proportional to the atom's mass, the leading order scattering off optical phonons vanishes at small momentum transfer

It holds for a general DM scattering potential of the form

$$V = \sum_{l=1}^N g_l V(\mathbf{r} - \mathbf{r}_l)$$

Here  $N$  is the number of atoms in the molecule

Or (by Bloch's Theorem) the number of atoms in the unit cell of a crystal

Simple proof

Scattering matrix element is, in the Born approximation

$$\langle \Phi_f | V(\mathbf{q}) \sum_l g_l e^{i\mathbf{q} \cdot \mathbf{r}_l} | \Phi_i \rangle \quad f \neq i$$

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Expand this at low  $q$       $q < 1/(\text{size molecule or unit cell}) \sim \text{keV}$

$$= iV(\mathbf{q}) \mathbf{q} \cdot \langle \Phi_f | \sum_l g_l \mathbf{r}_l | \Phi_i \rangle$$

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Now setting  $g_i = g m_i$

$$= iV(\mathbf{q})g \mathbf{q} \cdot \langle \Phi_f | \sum_l m_l \mathbf{r}_l | \Phi_i \rangle$$



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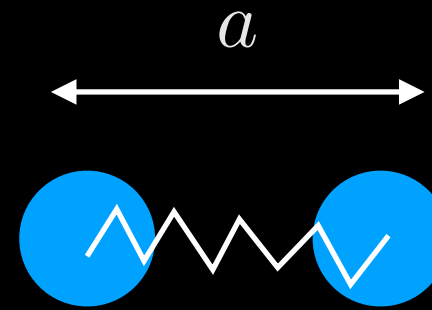
and identifying

$$\sum_{l=1}^N m_l \mathbf{r}_l = \mathbf{X}_{CM}$$

This can never induce  
internal changes

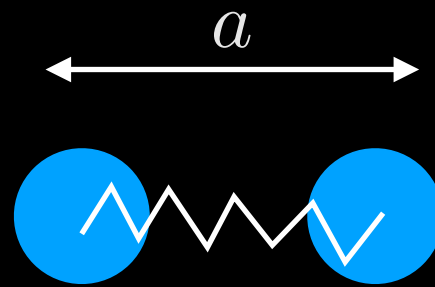
# Simple example

Di-molecule



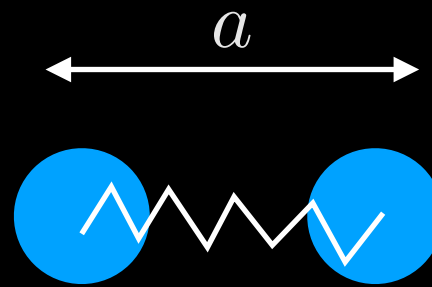
(Unit cell of a crystal with more than one atom works in essentially the same way)

Di-molecule



$$H = \frac{P^2}{2M} + \frac{p^2}{2\mu} + \frac{1}{2}\mu w^2 (x')^2$$

Di-molecule



$$H = \frac{P^2}{2M} + \frac{p^2}{2\mu} + \frac{1}{2}\mu\omega^2(x')^2$$

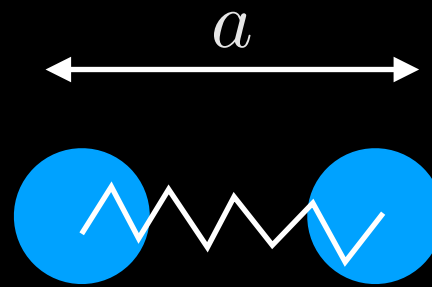
Matrix element for phonon creation

$$\mathcal{M} \sim \langle 1 | (g_1 e^{iq \cdot r_1} + g_2 e^{iq \cdot r_2}) | 0 \rangle$$

Rate proportional to

$$|\mathcal{M}|^2 \propto \frac{q^2}{M\omega} (g_1^2 m_2^2 + g_2^2 m_1^2 - 2g_1 g_2 m_1 m_2 \cos qa)$$

Di-molecule



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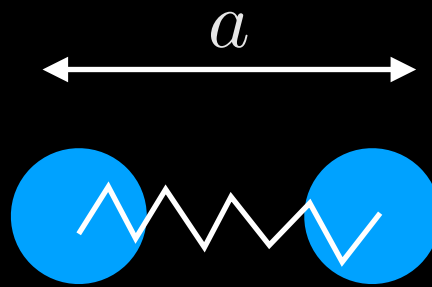
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Rate proportional to

This factor  
per phonon

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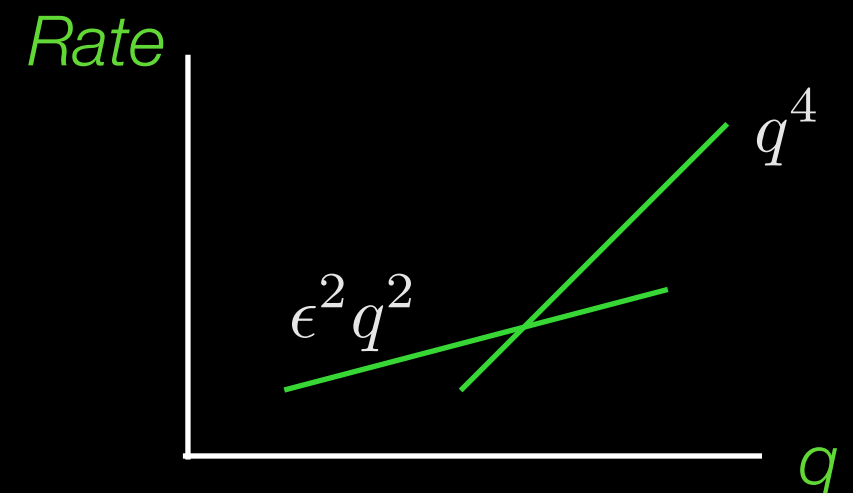
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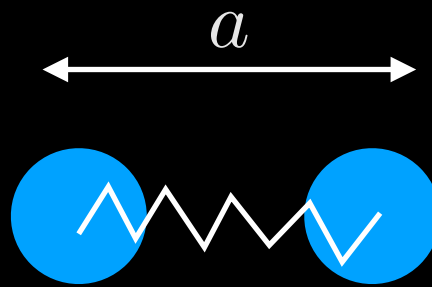
$$|\mathcal{M}|^2 \propto \frac{q^2}{Mw} (g_1^2 m_2^2 + g_2^2 m_1^2 - 2g_1 g_2 m_1 m_2 \cos qa)$$

Setting  $g_1 = gm_1(1 + \epsilon/2)$   $g_2 = gm_2(1 - \epsilon/2)$

$$|\mathcal{M}|^2 \propto q^2 g^2 \mu (\epsilon^2 + (qa)^2 + \dots)$$



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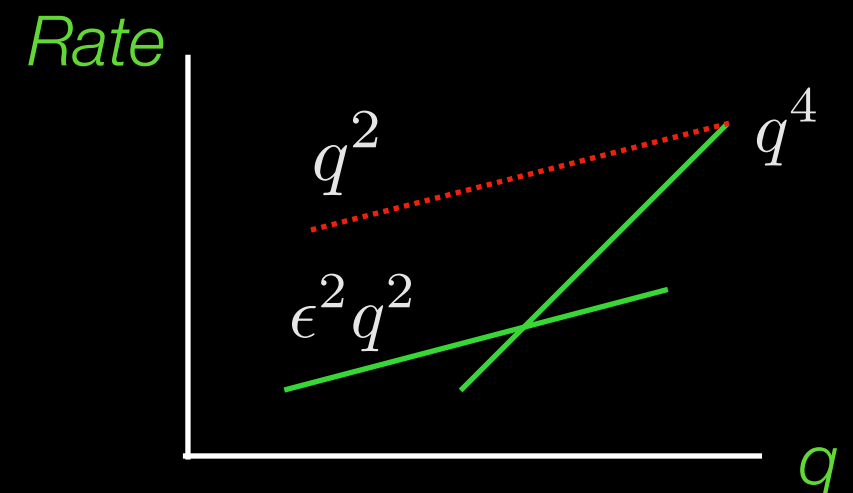
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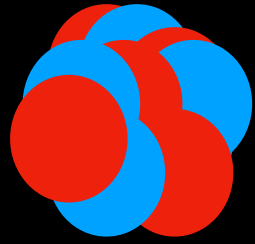
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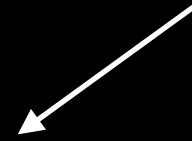
# Material consequences



# Characterisation of materials

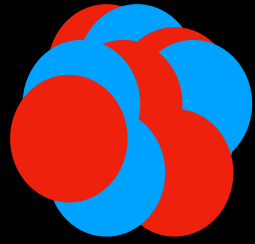


$$g_i = g_p Z_i + g_n (A_i - Z_i) = A_i \left( (g_p - g_n) \frac{Z_i}{A_i} + g_n \right)$$

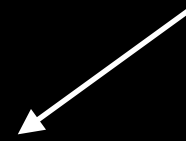


Atomic number and atomic mass only differ due to binding energy effects

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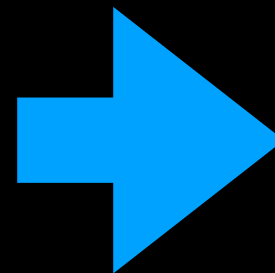
Atomic number and atomic mass only differ due to binding energy effects

*Suppression of LO term by  $\epsilon^2$*

$$g_p = g_n$$

*Always in the coupling to mass limit*

*Deviations of  $\epsilon \sim \mathcal{O}(10^{-3})$  due to binding energy*

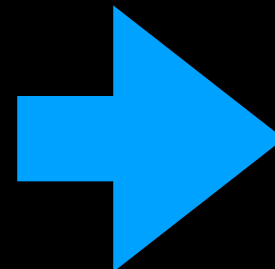


**Must calculate higher order terms to reliably estimate rate**

$$g_p \neq g_n$$

*Deviations determined by difference in the **ratio**  $Z/A$*

*Expect  $\epsilon \sim \mathcal{O}(0.01 - 0.1)$*



**Select materials with mix of light ( $Z/A=1/2$ ) and heavy ( $Z/A \sim 0.4$ ) elements, or contain hydrogens ( $Z/A=1$ )**

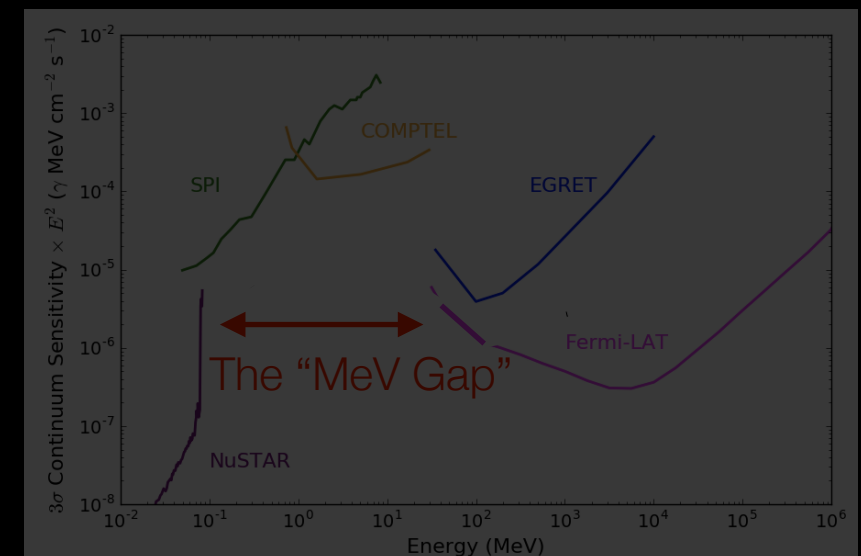
Simplest of systems  
already provides  
curiosities with  
implications for DM  
detector material

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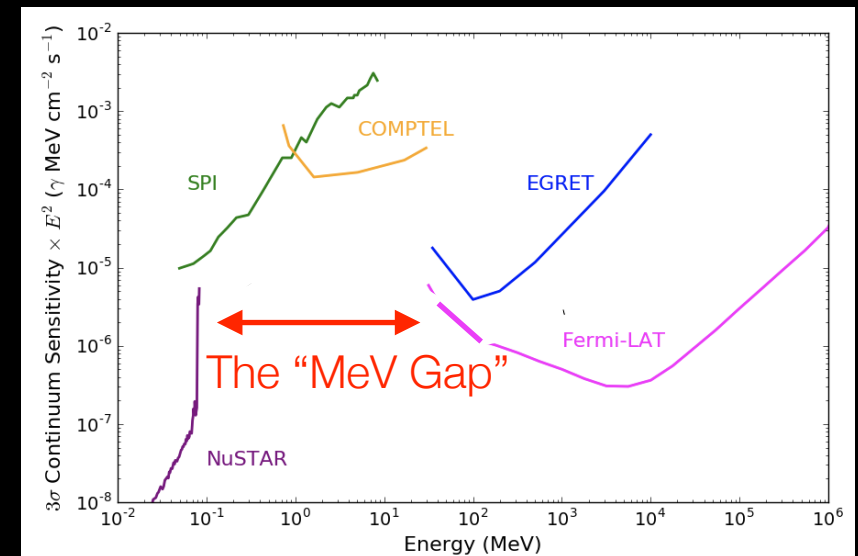
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