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UV-IR BOTTLENECKS FOR INFLATION AND THE HIGGS Alexander Westphal (DESY) with M. Dias, J. Frazer [1810.05199] IPMU, Tokyo, Feb 18, 2020 and N. Kaloper [1907.05837]











and future: Simons Array with $\sigma(r) = 0.003$; LiteBIRD with $\sigma(r) = 6 \times 10^{-4}$ -> see talks yesterday afternoon!



<u>the string theory landscape:</u> many isolated *vacua*, connected by tunneling some mountain slopes drive <u>inflation</u>

string theory's 6 compact dimensions: strings, branes & fluxes





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string theory's 6 compact dimensions: strings , branes & fluxes

D3's





BB in I~80 bandbower I(I+1)C / 2π [$_{
m II}$ K²]

inflation in string theory ...



inflation



axion monodromy - a summary

[Silverstein & AW; McAllister, Silverstein & AW; Kaloper & Sorbo '08] [Kaloper, Lawrence & Sorbo '11]

erc

$$\int \mathrm{d}^{10}x \left(\frac{|dB|^2}{g_s^2} + |F_1|^2 + |F_3|^2 + \left| \tilde{F}_5 \right|^2 \right)$$

• example - B₂ - axion monodromy from flux:

$$|\tilde{F}_p|^2 = |dC_{p-1} + B_2 \wedge F_{p-2}|^2$$
 mass term
$$flux \ F_{p-2} = Nf_{p-2} \ , \ \int_{\Sigma_{p-2}} f_{p-2} = 1$$

$$B_2 \to B_2 + d\Lambda_1 \implies C_{p-1} \to C_{p-1} - N\Lambda_1 \wedge f_{p-2}$$

many models in '10-'18: Stanford/Cornell/Hamburg, Madrid, Madison, Heidelberg, ...

Flattening 1: moduli backreact in axion monodromy

• bare bones monodromy: $\int d^{10}x \left(\frac{|dB|^2}{g_s^2} + |F_1|^2 + |F_3|^2 + |\tilde{F}_5|^2 \right)$

$$V = \frac{C_1}{\phi} + C_2 \phi^2 (\mu^2 + b^2) \implies \langle \phi \rangle = \langle \phi \rangle_0 (1 + b^2 / \mu^2)^{-1/3}$$

• 2 types of flattening — additive & multiplicative:

$$V_{eff.}(b) = V|_{\langle \phi \rangle} \sim \langle \phi \rangle_0^2 \frac{b^2}{(1+b^2/\mu^2)^{2/3}} \sim \begin{cases} b^2 - \frac{2}{3} \frac{b^4}{\mu^2} &, \quad \mu \gg 1\\ b^{2/3} &, \quad \mu \ll 1 \end{cases}$$

• other powers as well: ϕ , $\phi^{4/3}$, ϕ^2 [McAllister, Silverstein, AW & Wrase '14] many models in '10-'18: Stanford/Cornell/Hamburg, [Hebecker et al. '14] Madrid, Madison, Heidelberg, ... [Buchmüller, Dudas, Heurtier, AW, Wieck & Winkler '15]

4D effective axion monodromy inflation

[Kaloper & Sorbo '08] [Kaloper, Lawrence & Sorbo '11; Kaloper & Lawrence ...]

$$\mathcal{L} = \frac{1}{2} (\partial \phi)^2 + \frac{1}{48} F_{(4)}^2 - \frac{\mu}{24} \phi \star F_{(4)} - \sum_{n \ge 2} \chi_n C_n^{\text{eff.}} \frac{(F_{(4)}^2)^n}{M_P^{4n-4}} - V_{\text{np}}$$

$$V_{\rm p}^{\rm eff.} = \sum_{n \ge 2}^{N} \chi_n \ C_n^{\rm eff.} \frac{\left(V^{(0)}(\phi)\right)^n}{M_{\rm p}^{4n-4}} \quad , \quad V_{\rm np} = \Lambda_{\rm UV}^4 \sum_{m \ge 1}^{M} D_m e^{-mS} \cos\left(\frac{m\phi}{f_{\phi}}\right)$$

$$S = \mathcal{C} \mathcal{V}^k \qquad \Lambda_{UV}^4 \sim \frac{1}{\mathcal{V}^2} \qquad \chi_n C_n^{\text{eff.}} \equiv C_n^{\text{tree}} + \chi_n C_n$$

$$V_{\rm p}^{\rm tree} = \sum_{n \ge 2}^{N} C_n^{\rm tree} \frac{\left(V^{(0)}(\phi)\right)^n}{M_{\rm p}^{4n-4}} = V_0 \left[\left(1 + \frac{\phi^2}{\mu^2}\right)^{p/2} - 1 \right] \quad , \quad V_{\rm p} = \sum_{n \ge 2}^{N} \chi_n C_n \frac{\left(V^{(0)}(\phi)\right)^n}{M_{\rm p}^{4n-4}}$$



EXAMPLE: AXION MONODROMY

Steps:

- 1. Identify relevant scales (class of models)
- 2. Learn the mapping from parameters to observables
- 3. Study how predictions change according to prior choice

Use numerical methods developed in previous work to generate a large sample assuming

 $\mu \sim \mathcal{U}(0.1, 1) \quad p \sim \mathcal{U}(0.1, 2)$

- Take a random draw of μ and p
- Solve background equations of motion to get total number of e-folds
- Solve equations of motion for the perturbations and compute n_s at pivot scale
- Repeat many times
- Use machine learning to get $n_s(\mu,p)$



$$\mathcal{L} = \frac{1}{2} (\partial \phi)^2 - V_0 \left[\left(1 + \left(\frac{\phi}{\mu}\right)^2 \right)^{p/2} - 1 \right]$$

erc







From 4-forms to Goldilocks ...



string dS construction have many dS vacua from fluxes
 <u>accommodate</u> small CC a la Bousso-Polchinski:

[Bousso & Polchinski '00]

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$$S = \int d^4x \sqrt{-g} \left(\frac{M_{\rm P}^2}{2} R - \Lambda_0 - \sum_i \left(F_{(4)}^{(i)} \right)^2 \right) + S_{bound.}$$

$$F_{(4)}^{(i)} = F_{\mu\nu\rho\sigma}^{(i)} = \partial_{[\mu}A_{\nu\rho\sigma}^{(i)} , \qquad S_{membr.} = q_i \int A_{(3)}^{(i)}$$



string dS construction have many dS vacua from fluxes
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for simplicity, U(1) Higgs:

$$V_{0} = \frac{\lambda}{4} |\phi|^{4} - \frac{v^{2}}{2} |\phi|^{2} + \Lambda$$

couple one *F*₍₄₎ to Higgs
$$\Delta V = \frac{c}{24} \epsilon_{\mu\nu\lambda\sigma} F^{\mu\nu\lambda\sigma} |\phi|^{2} - \frac{1}{48} F_{\mu\nu\lambda\sigma}^{2}$$















The Goal

- energy scale of inflation 10¹³ x LHC close to Planck Scale !
 - → unique window to quantum gravity !

 proposal: test string theory with inflation & CMB, or EW SSB!



A Roadmap into the Post -Naturalness Era (?) ... [Giudice '17]



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