

The Large Charge Sector of 3D parity violating CFTs

Umang Mehta
umangmehta@uchicago.edu

with Luca Delacretaz and Gabriel Cuomo

arXiv: 2102.05046

Outline

- EFT without parity
 - Euler term at $\mathcal{O}(\partial^1)$
 - Vortices on the sphere
 - Spectrum and OPEs
- Anyon Superfluids
 - Derivation of the leading order EFT on a plane
 - Monopole operators
 - Vortices on the sphere as holes
 - Beyond EFT - Chern-Simons terms, double scaling, rotons

Large $U(1)$ charge EFT

$$S \sim c_1 \int d^3x \sqrt{g} |\partial\phi|^3 + \mathcal{O}(\partial^2) \quad \Leftrightarrow \quad S \sim \frac{1}{\sqrt{c_1}} \int d^3x \sqrt{g} |f|^{3/2} + \mathcal{O}(\partial^2)$$

Currents	ϕ	a_μ
$U(1)$	$ \partial\phi ^2 \partial_\mu\phi$	$\epsilon_{\mu\nu\rho} f^{\nu\rho}$
$U(1)^{(1)}$	$\epsilon_{\mu\nu\rho} \partial^\rho\phi$	$f_{\mu\nu} / \sqrt{ f }$

EFT without parity

- J_{Euler}^0 integrated over space gives the Euler characteristic

$$S \sim \frac{1}{\sqrt{c_1}} \int d^3x \sqrt{g} |f|^{3/2} + \kappa \int d^3x \sqrt{g} a_\mu J_{Euler}^\mu$$

- $\kappa \in \mathbb{Z}$ is quantized

$$J_{Euler}^\mu = \frac{1}{8\pi} \epsilon^{\mu\nu\rho} \epsilon^{\alpha\beta\gamma} u^\alpha \left(\nabla_\nu u^\beta \nabla_\rho u^\gamma - \frac{1}{2} \mathcal{R}_{\nu\rho}{}^{\beta\gamma} \right)$$

- κ controls Hall viscosity

- Euler term breaks one-form symmetry

$$u^\alpha = \frac{\epsilon^{\alpha\beta\gamma} f_{\beta\gamma}}{\sqrt{2} |f|}$$

- Equation of motion: $\nabla_\mu J_{1-form}^{\mu\nu} = \kappa J_{Euler}^\nu$

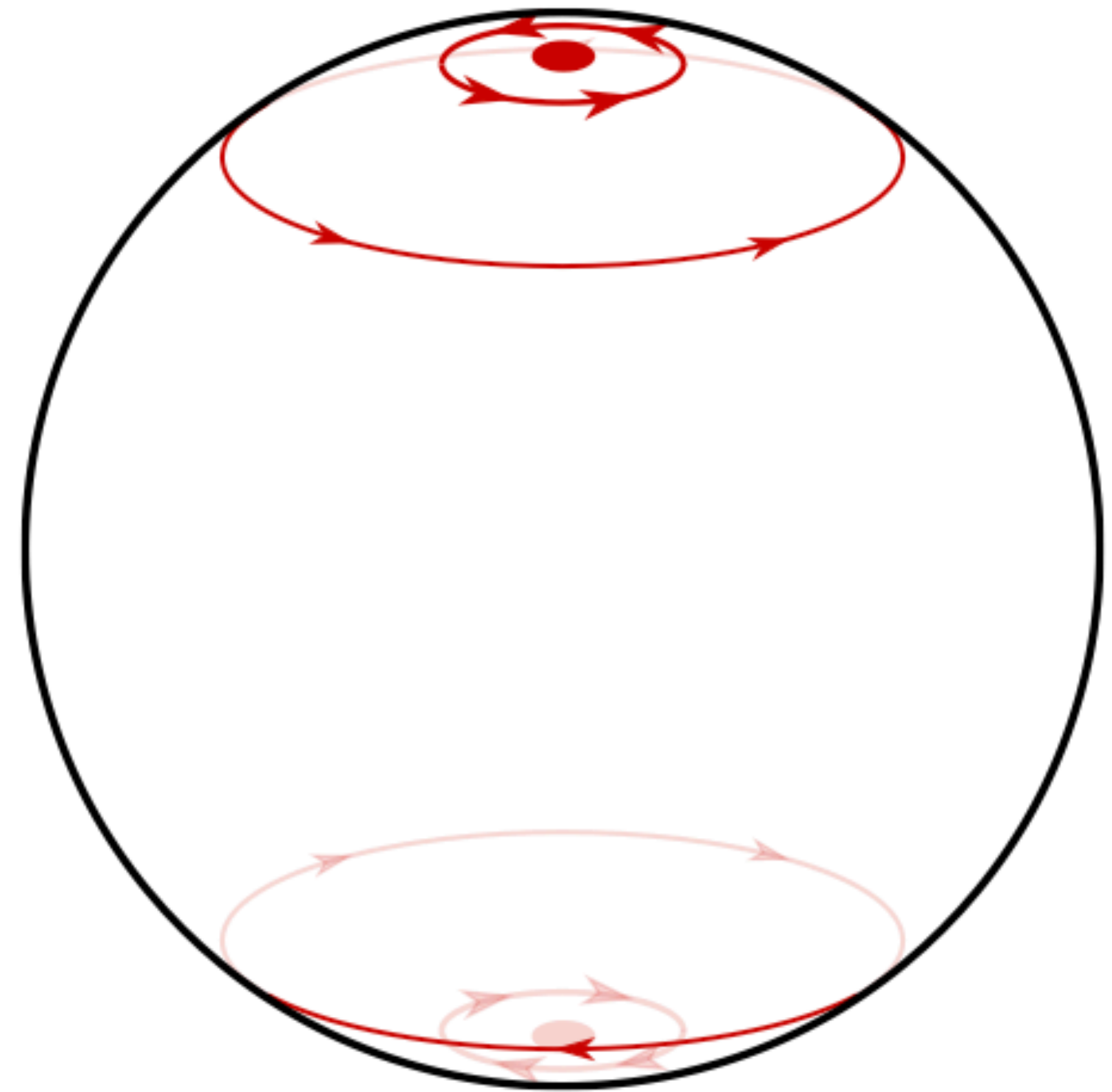
EFT on the sphere

- Integrated e.o.m. $\implies 0 = \kappa \chi_{S^2}/2 = \kappa$
need κ vortices \cong point charges

- $S_{vortices} = - \sum_{p=1}^{\kappa} \int_{X_p} a + \dots$

- Vortex angular momentum: $\vec{J} = \sum Q \vec{X}_p / 2$

- Coulomb interaction: $H = - \sum \log |\vec{X}_p - \vec{X}_{p'}|^2$



Spectrum: classical ground state

- Energy minimization = Whyte's problem - no general solution known except in special cases $\kappa = 1, 2, 3, 4, 5, 6, 12$
- Theorem: center of mass at sphere origin when $\kappa \neq 1$
- Main result: lightest operator has $J = 0$ for $\kappa \neq 1$ and $J = Q/2$ for $\kappa = 1$ with dimension

$$\Delta_{min}(Q) = \frac{1}{\sqrt{c_1}} Q^{3/2} + \kappa \sqrt{c_1} \sqrt{Q} \log Q + \dots$$

- Inhomogeneous ground state, large degeneracy

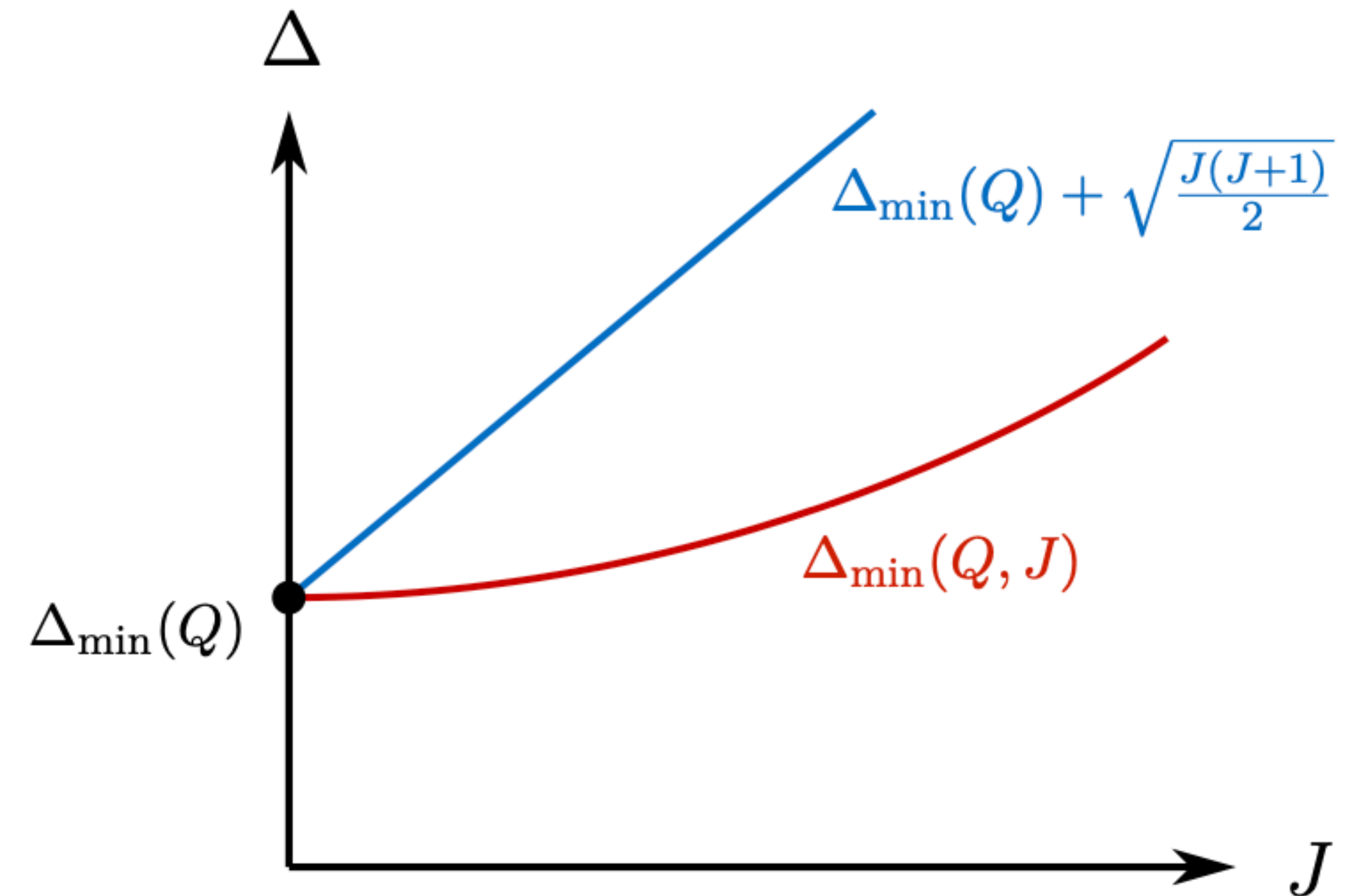
Spectrum: quantum ground state and vortex excitations

- Quantization \implies Landau levels on sphere (monopole harmonics)

- Quantization lifts degeneracy, eg $\kappa = 2$

$$\Delta_{min}(Q, J) = \Delta_{min}(Q) + \sqrt{c_1} \frac{J(J+1)}{Q^{3/2}}$$

- Much softer than phonon excitations with $\Delta_{Q,J} - \Delta_{min} = \sqrt{J(J+1)/2}$



Hellerman, Orlando, Reffert, Watanabe '15
Monin, Pirtskhalava, Rattazzi, Seibold '16
Jafferis, Mukhametzhanov, Zhiboedov '17

OPE coeffs and Hall viscosity

- Kubo formula: $\eta_H \sim \langle \bar{Q} T_{xy} T_{xx} Q \rangle$. The stress tensor is linear in Goldstone fluctuations

$$T \sim 1 + c_1 \partial\pi + \kappa \partial\partial\pi$$

- Therefore, we expect $\langle \bar{Q} T Q' \rangle_{odd} \sim \eta_H$ where $Q' = Q + \text{single phonon}$

$$\langle \bar{Q} T Q_J \rangle_{even} \sim \frac{Q^{3/2}}{c_1^{1/4}} J^{1/2}$$

$$\langle \bar{Q} T Q_J \rangle_{odd} \sim \kappa c_1^{1/4} Q^{3/2} J^{3/2}$$

Candidate theories

- $U(N)_{k_1, k_2}$ or $SU(N)_k$ Chern-Simons matter theories?
- Global $U(1)$ - topological or fermion number respectively
- Seem to be a Fermi liquid at strict $N = \infty$, possibly superfluid for finite N
- Superfluid for $N = 1$, $k \rightarrow \infty$ - “anyon superfluids”

Anyon Superfluids - flat space

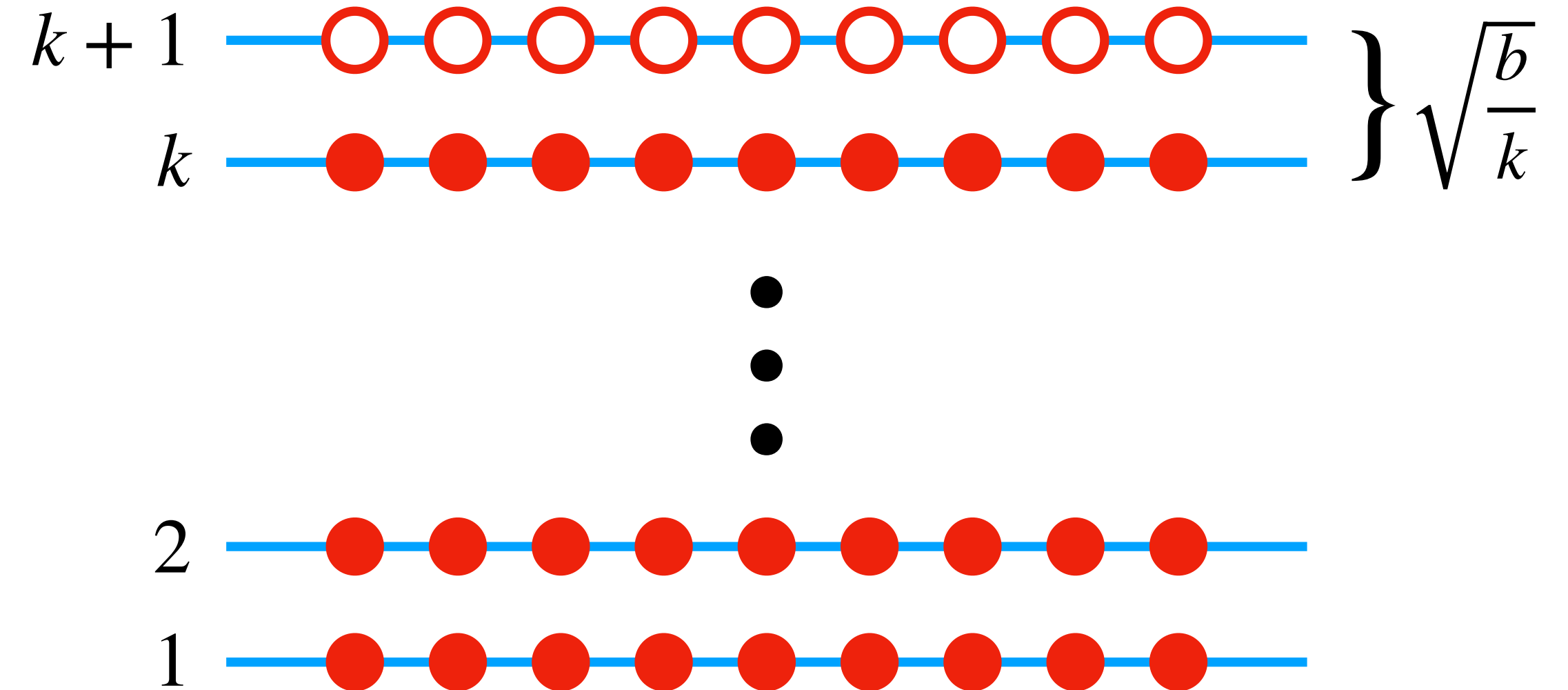
- Global $U(1)$ current $j^\mu = \frac{1}{4\pi} \epsilon^{\mu\nu\rho} f_{\nu\rho}$

$$S = \int d^3x \left(\bar{\psi} i \gamma^\mu D_\mu \psi - \frac{k}{4\pi} \epsilon^{\mu\nu\rho} a_\mu \partial_\nu a_\rho \right)$$

- Gauss law: $n_\psi = \bar{\psi} \gamma^0 \psi = k j^0 = \frac{kb}{2\pi}$

- k fully filled Landau levels

- Hall conductivity $\sigma_{xy} = \frac{k}{2\pi}$



Leading order EFT

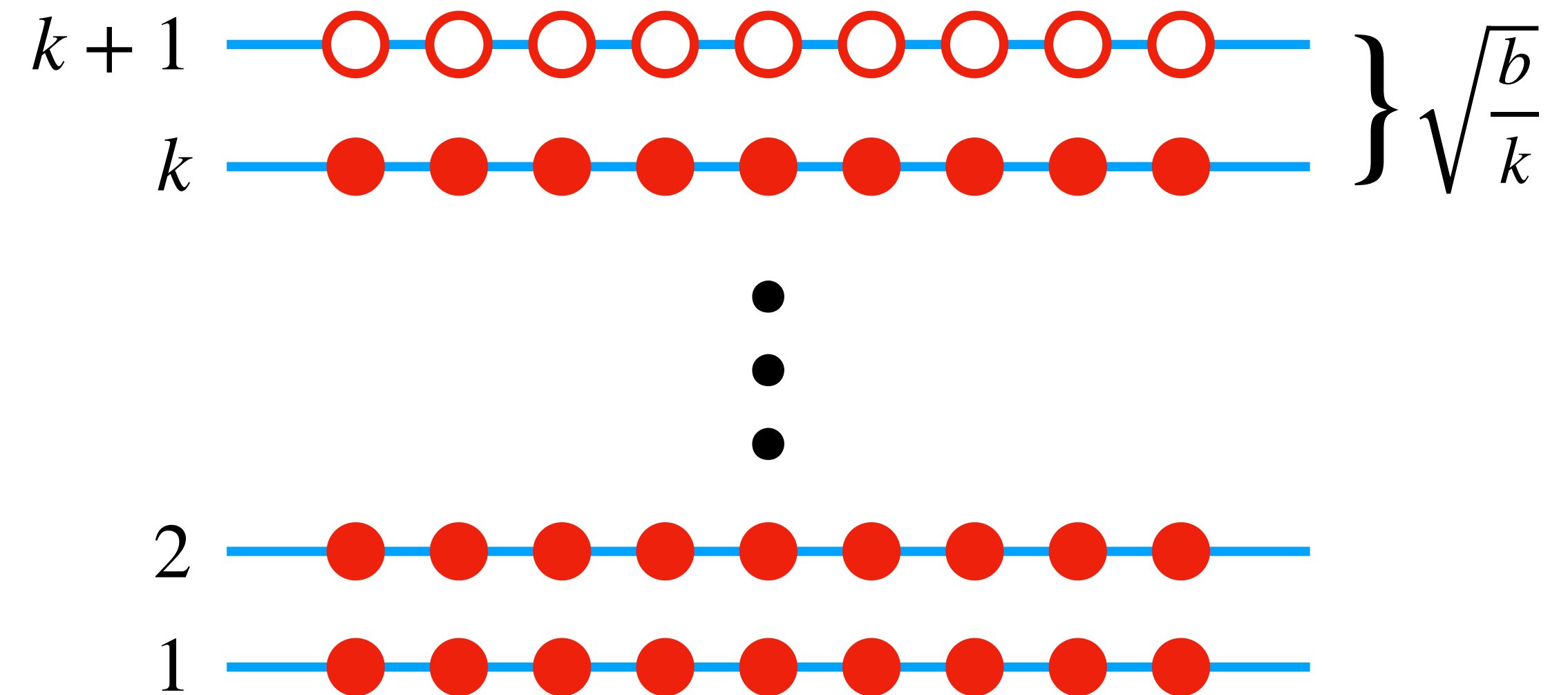
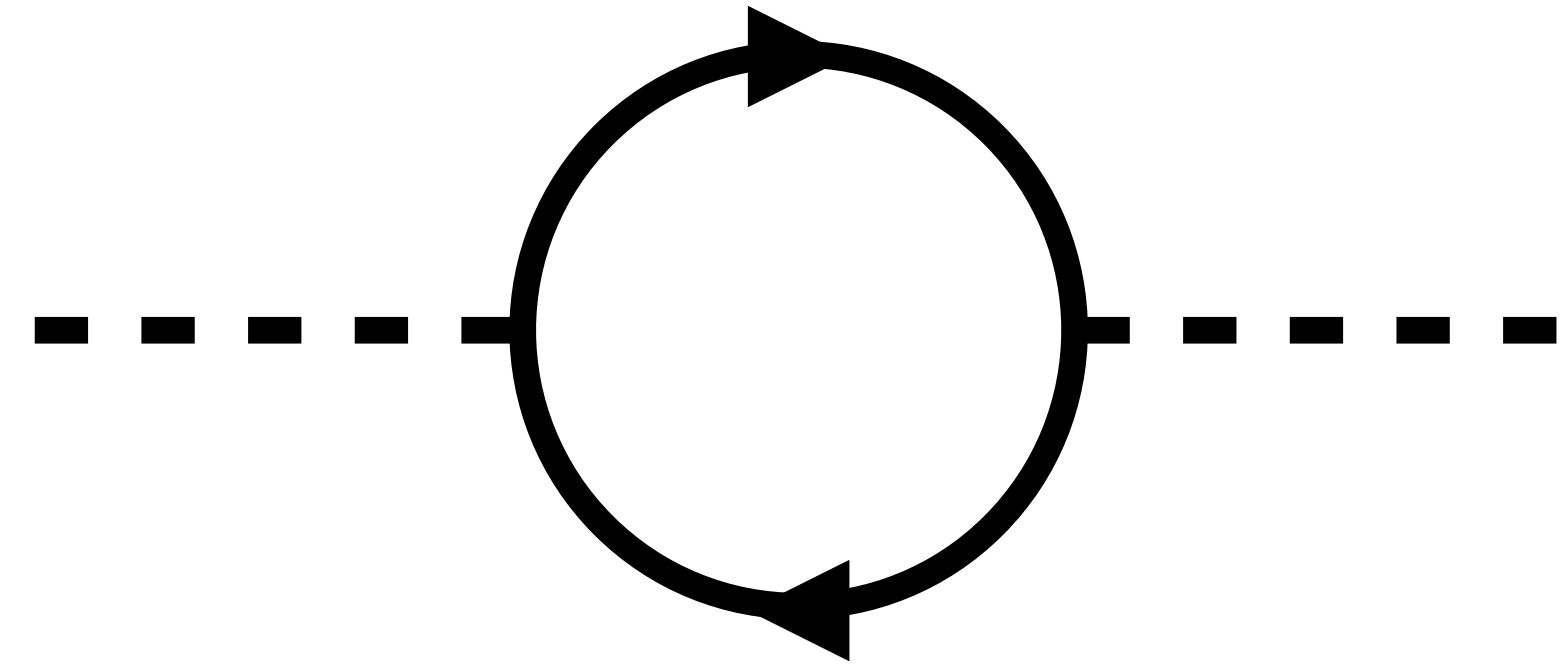
- Integrate out fermions to get

$$\delta S = \frac{k}{4\pi} \int a da$$

- Cancels CS term and generates higher order terms

$$S_{EFT} = \frac{k^{3/2}}{\sqrt{2\pi\rho}} \int \left(e^2 - \frac{1}{2} b^2 \right) + \dots \sim \int |f|^{3/2}$$

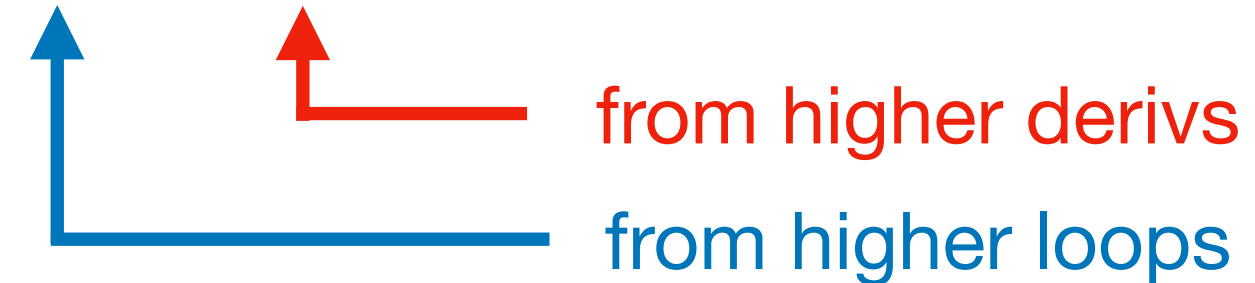
- Read off $c_1 = 1/2\pi k^3$



Monopole operators and double scaling

- Charged operators are monopoles. Operator dimension from EFT ($1 \ll k \ll Q$) is given by

$$\Delta_{min}(Q) = \frac{2}{3}(kQ)^{3/2} + \mathcal{O}(1/k, k/Q)$$



- Opposite limit $Q \ll k$

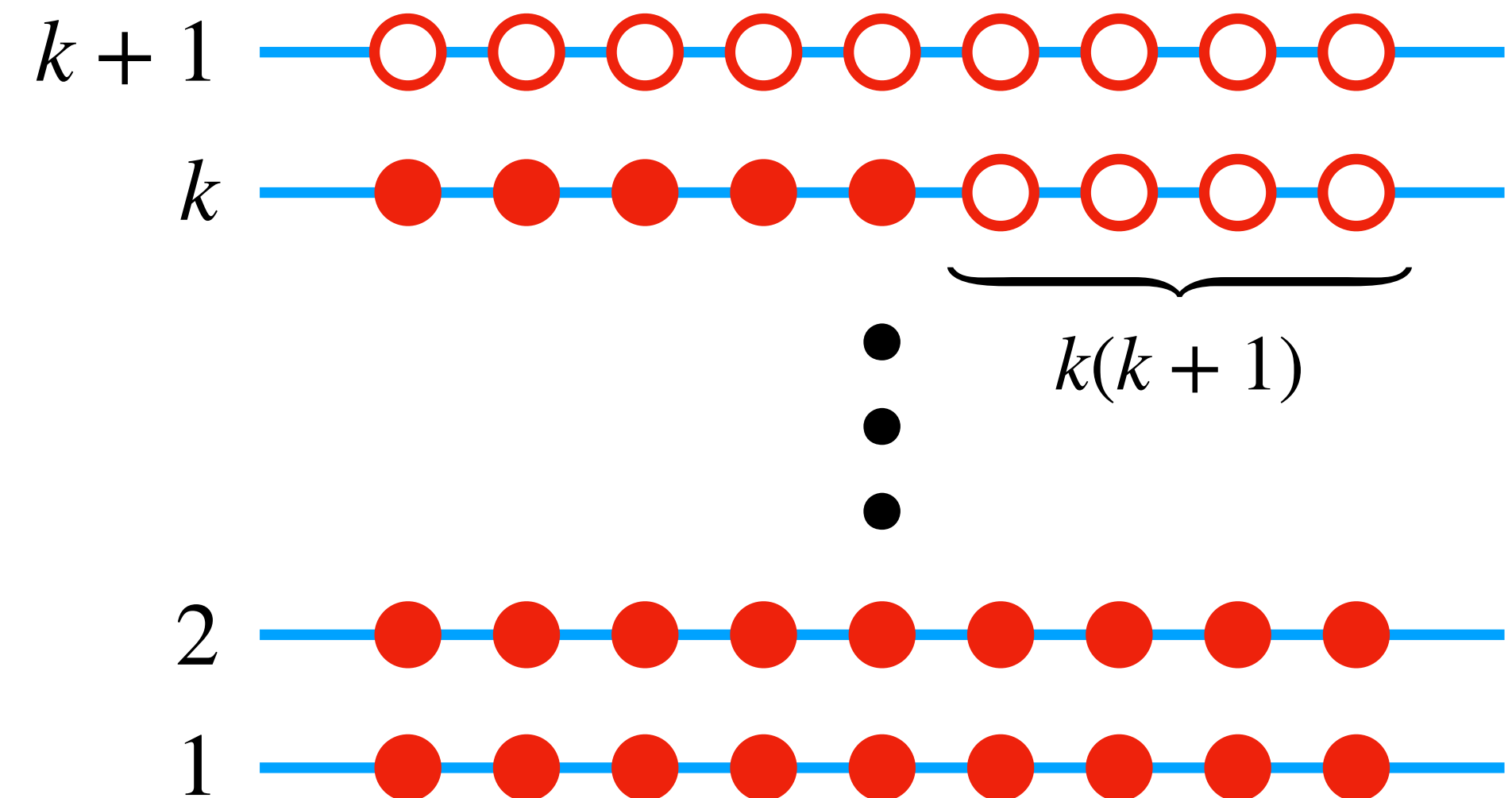
$$\Delta_{min}(Q) = \frac{2}{3}(kQ)^{3/2} + \dots$$

- Double scaling limit: $Q, k \rightarrow \infty$, Q/k fixed - organize Dirac particles in Landau levels obtained from the monopole background at leading order

$$\Delta_{min}(Q) = \frac{2}{3}k^3 \left(\frac{Q}{k}\right)^{3/2} \left[1 + \mathcal{O}\left(\frac{1}{k}\right)\right]$$

Anyon superfluids on the sphere

- kQ Dirac fermions in sphere Landau levels
- Degeneracy $Q + 2|p|$, $p \in \mathbb{Z}$. k fully filled LLs requires $kQ + k(k + 1)$ fermions
- $k(k + 1)$ holes on the last one - superfluid vortices
- Identify $\kappa = k(k + 1) \gg 1$



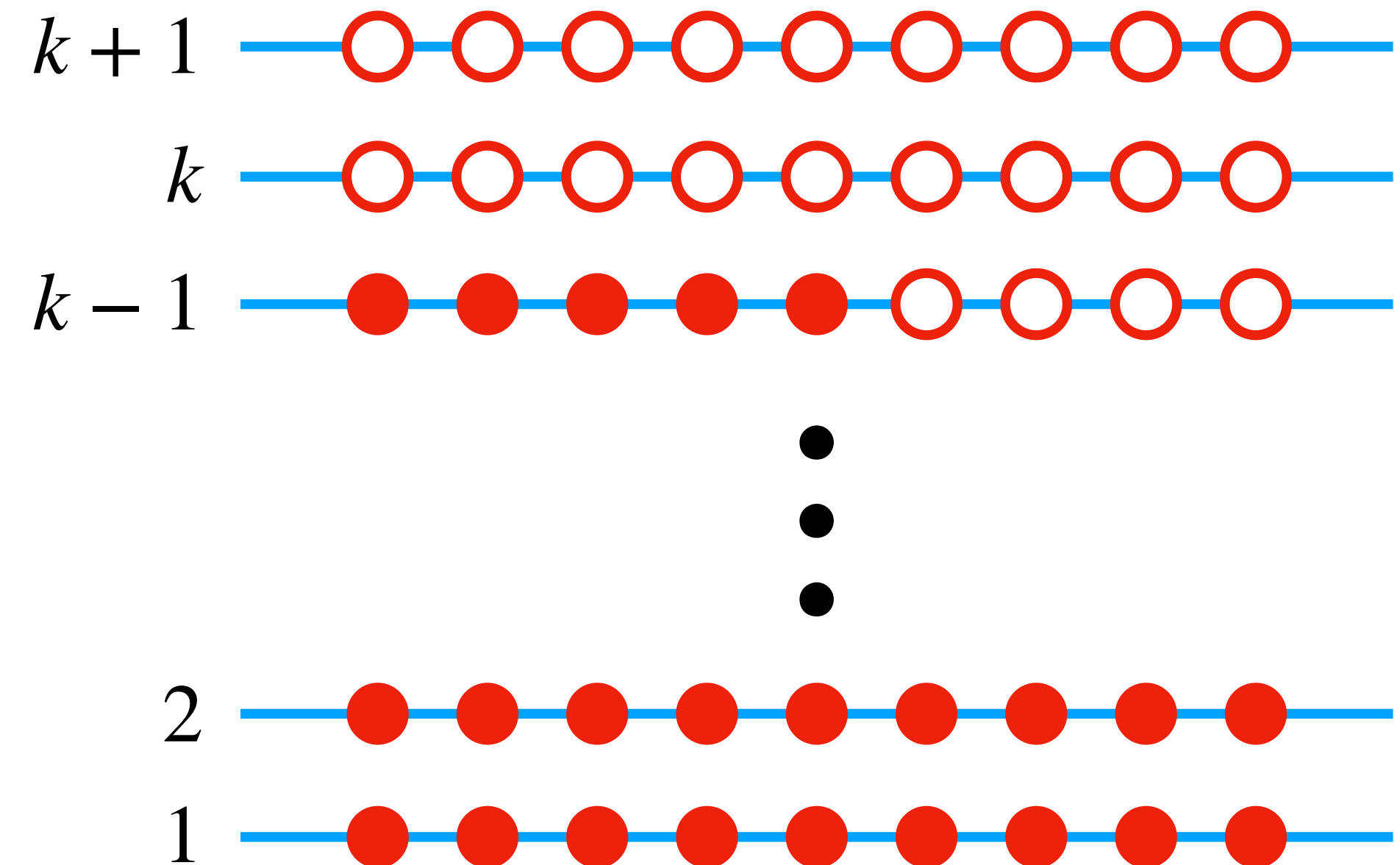
(assumed $k^2 \simeq \kappa \lesssim Q$)

Beyond EFT: Chern-Simons terms

- What happens when $k \ll Q \lesssim k^2$?
- Number of filled LLs $\sim k - [k^2/Q]$
- Integrating out ψ leaves residual Chern-Simons term

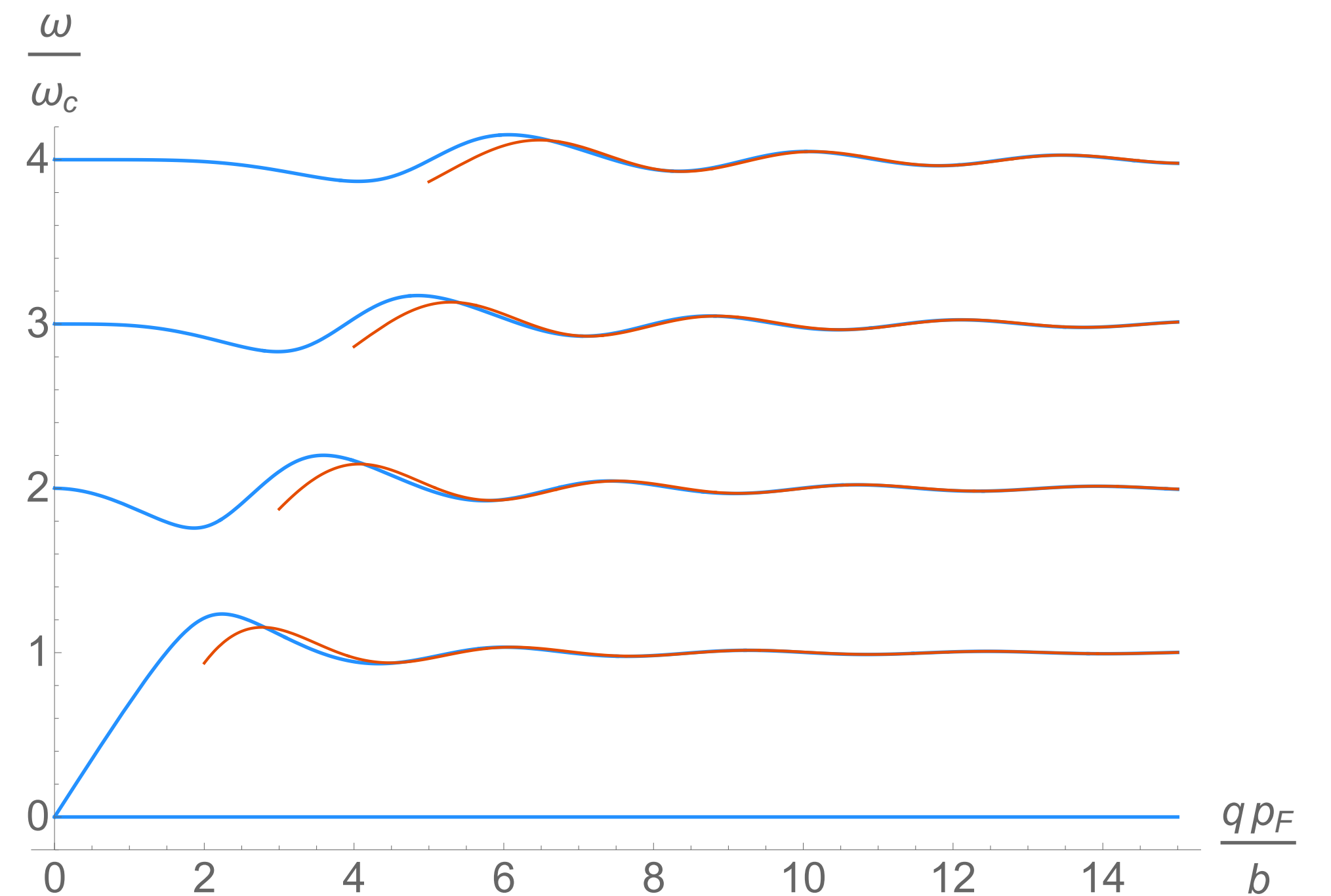
$$S_{EFT'} = S_{EFT} - \frac{[k^2/Q]}{4\pi} \int ada$$

- $\kappa \neq N_{vortices} \simeq k^2 - [k^2/Q]Q$; coefficient of subleading term in $\Delta_{min}(Q)$ changed from κ to $N_{vortices}$
- Superfluid phonons are gapped;
 $m_{phonon} \sim \sqrt{k/\rho} \ll 1/R_{S^2}$; gap invisible on sphere



Beyond EFT: rotons?

- Particle-hole excitations set the cutoff scale for EFT (and EFT') to $\sqrt{\rho/k}$
- Can compute flat space spectrum beyond EFT at weak coupling using a bosonized Fermi surface approach - "rotons"
- Rotons on the sphere? What CFT operators do they correspond to?



Thank you!