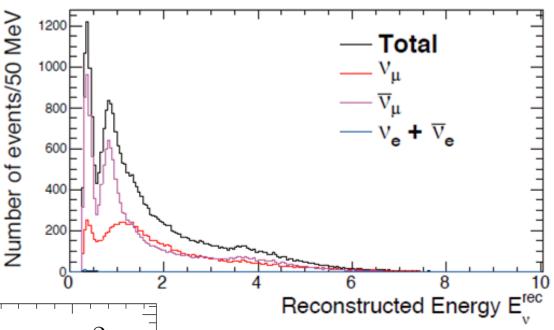
Charge Reconstruction with a Magnetised Muon Range Detector in TITUS

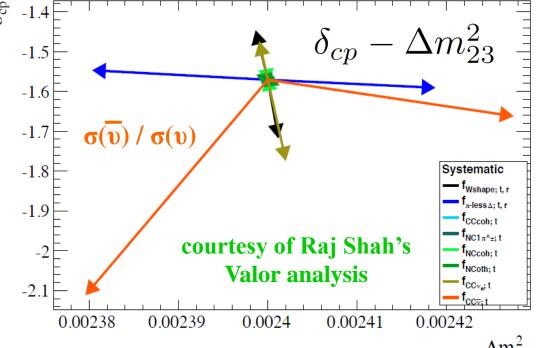
Mark A. Rayner – *Université de Genève*5th open Hyper-Kamiokande meeting, Vancouver
19th July 2014, Near Detector <u>pre-meeting</u>



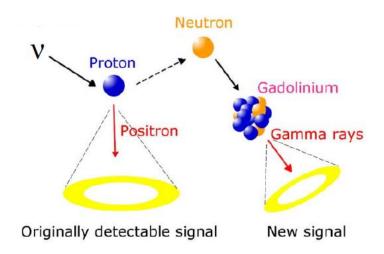
Motivation

Significant wrong-sign component in anti-neutrino mode





The anti-neutrino crosssection is the biggest unconstrained model systematic



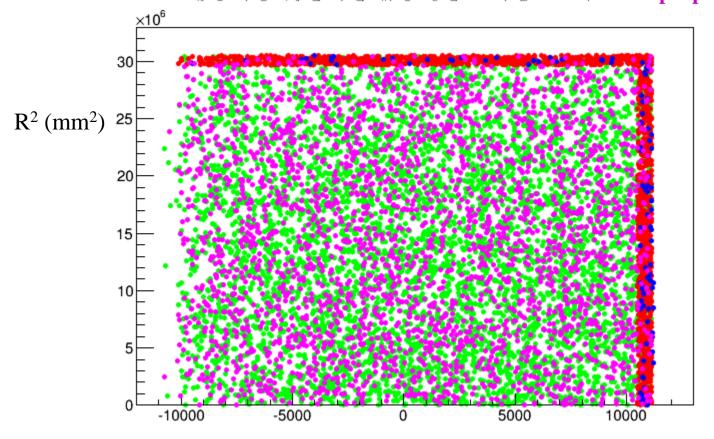
$$\upsilon n \rightarrow \ell p$$

Exciting, but somewhat untested

18% of muons escape the tank

((part_xEnd*part_xEnd)+(part_yEnd*part_yEnd)):part_zEnd (part_pid==13 && part_processEnd==0)

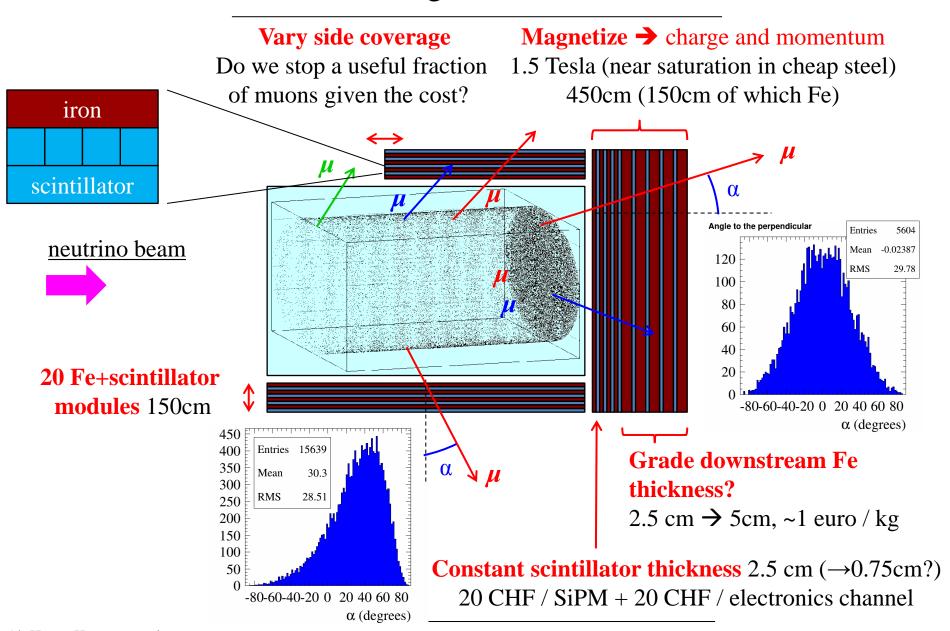
red: mu- leave tank
blue: mu+ leave tank
green: mu- stop in tank
purple: mu+ stop in tank



N.B. The tank size could be re-optimized with the MRD in mind

courtesy of Matthew Malek

MRD design considerations

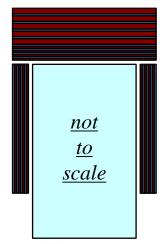


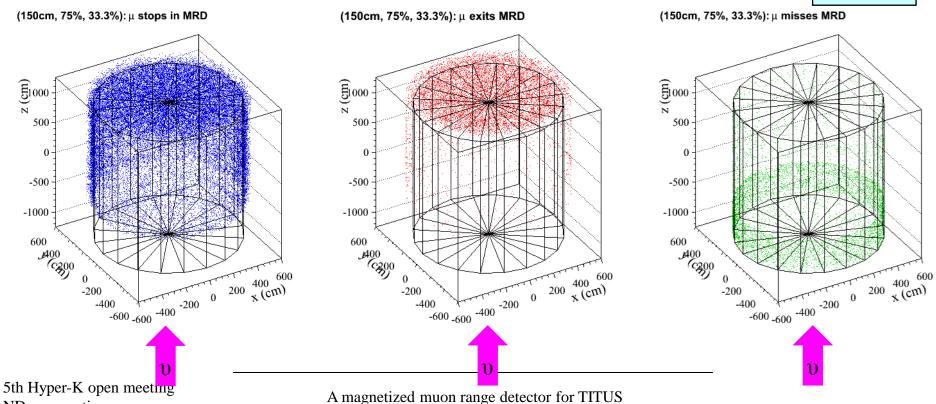
Muons tracks in the TITUS MRD

A simulation with 150cm end Fe and 75% side coverage of 50cm of Fe

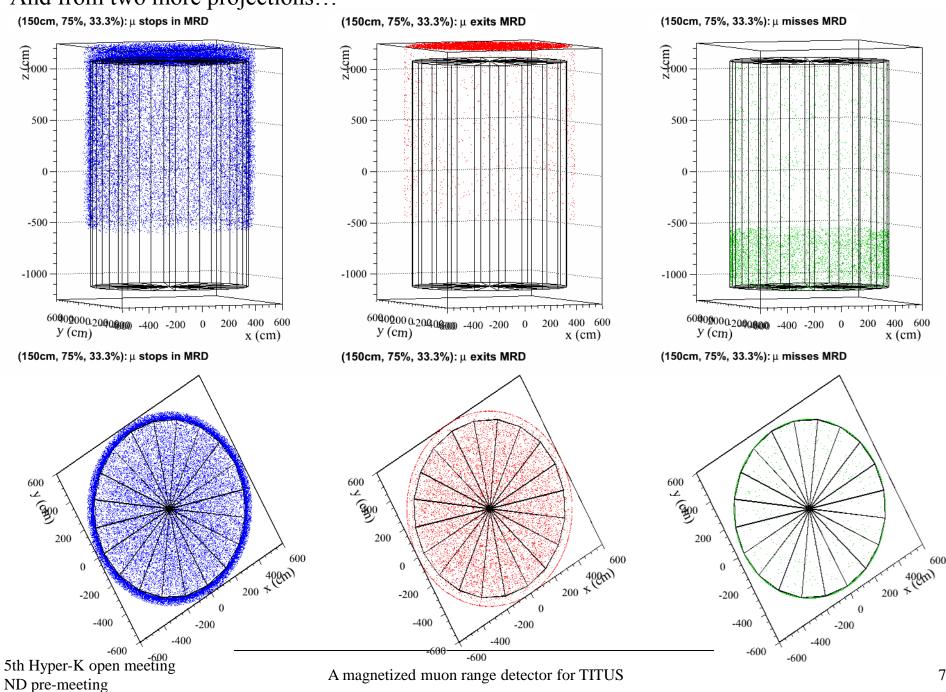
- range-out and stop in the MRD
- penetrate through the MRD
- miss the MRD

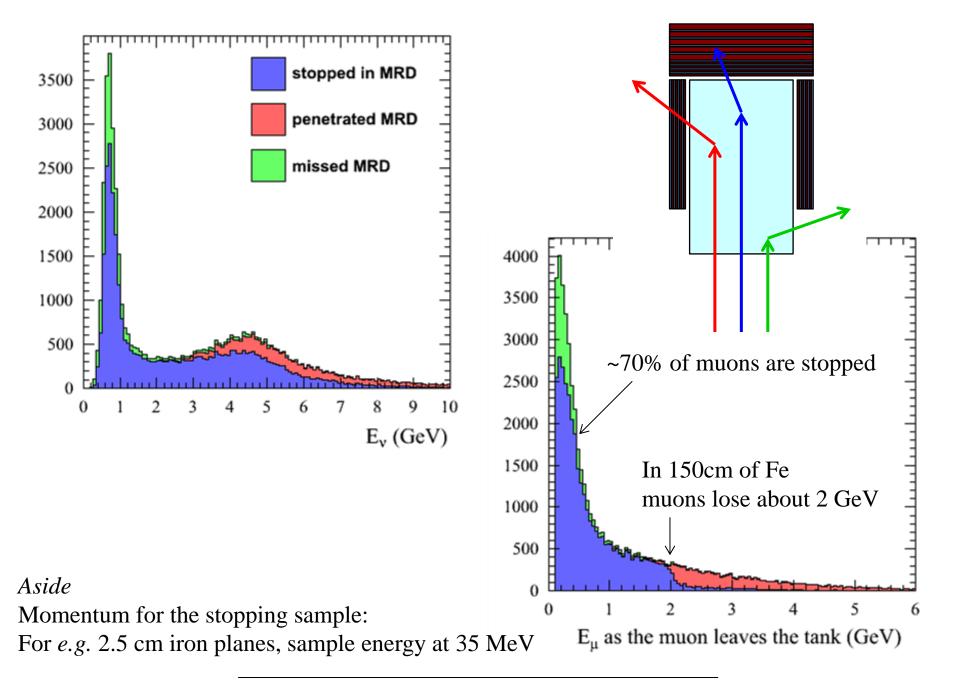
ND pre-meeting



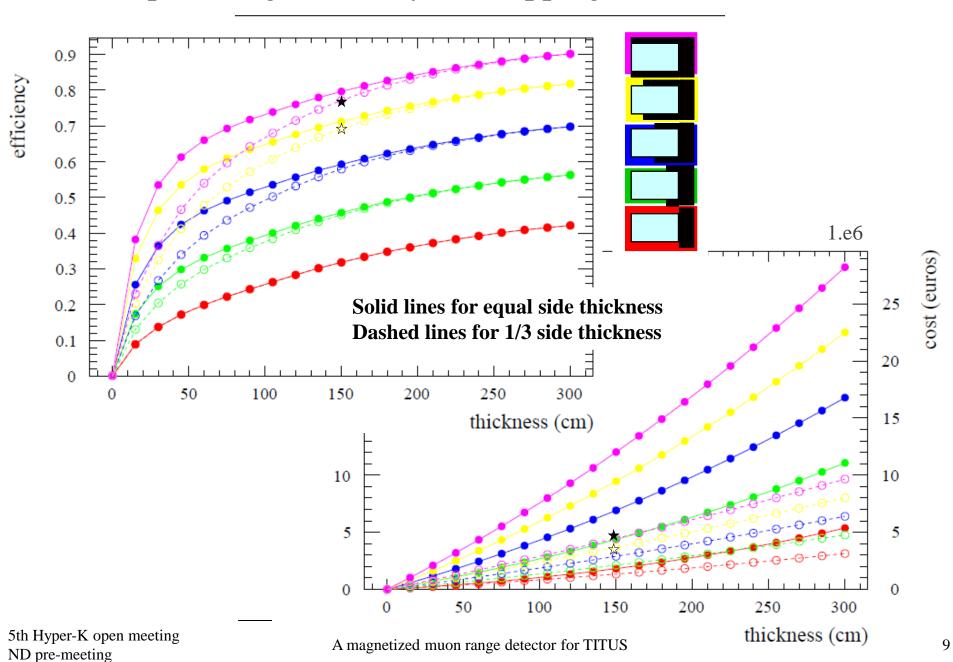


And from two more projections...

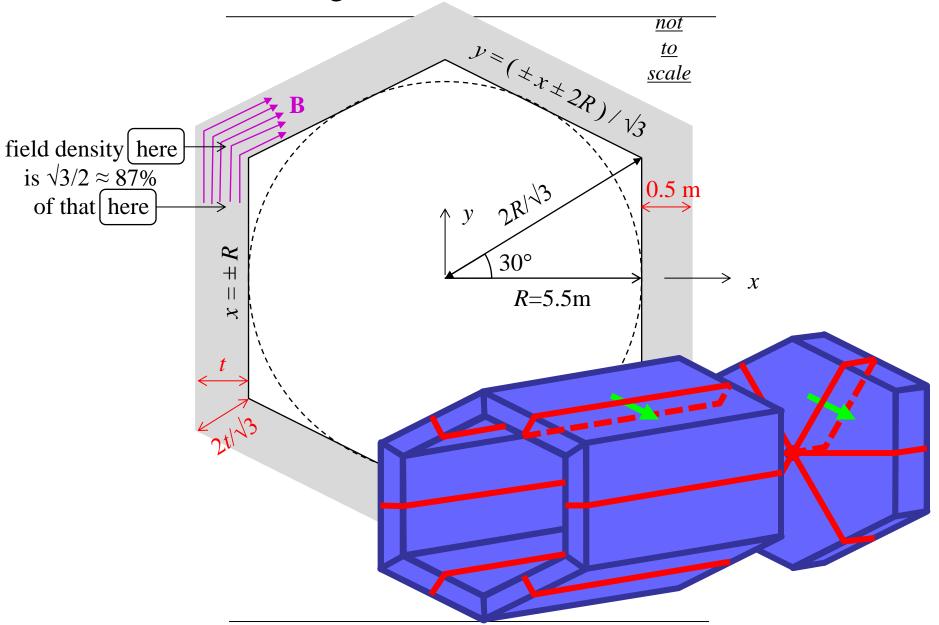




Optimizing efficiency for stopping muons, and cost

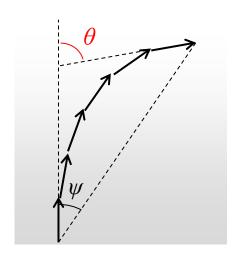


Magnetization of the MRD



Multiple Scattering in the iron is the biggest obstacle to charge reconstruction

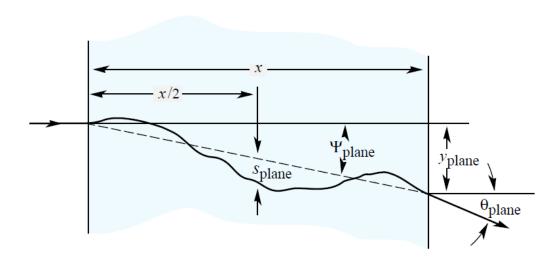
 $X_0 = 1.757$ cm in Fe $X_0 = 50.31$ cm in polyethylene



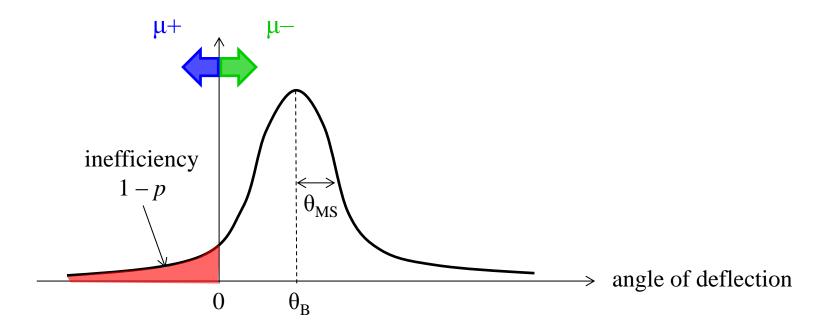
$$(X_0 / X_0)^{1/2} = 1.9\%$$

$$\theta_0 = \frac{13.6 \text{ MeV}}{\beta cp} z \sqrt{x/X_0} \Big[1 + 0.038 \ln(x/X_0) \Big]$$

$$\psi_{\text{plane}}^{\text{rms}} = \frac{1}{\sqrt{3}} \theta_{\text{plane}}^{\text{rms}} = \frac{1}{\sqrt{3}} \theta_0$$

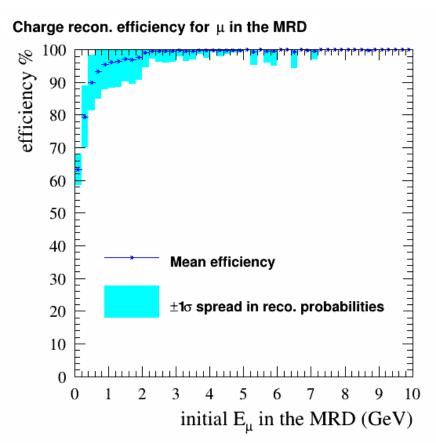


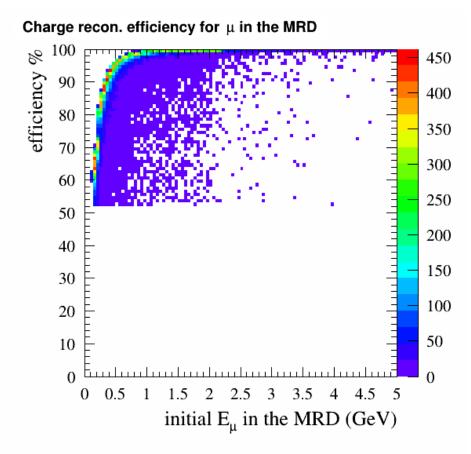
To start with, a probabilistic, back of a (fairly big) envelope calculation



TITUS MRD charge recon. efficiency vs. muon energy

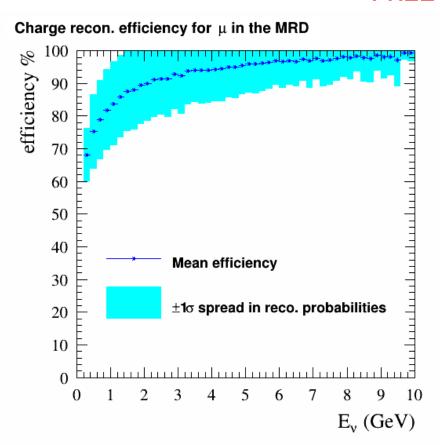
PRELIMINARY

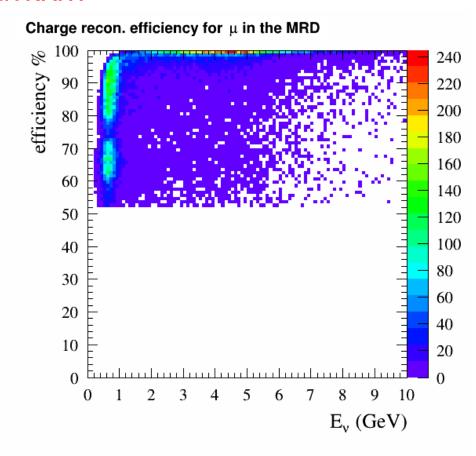




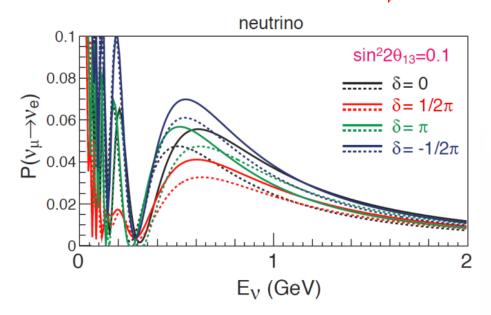
TITUS MRD charge recon. efficiency vs. neutrino energy

PRELIMINARY



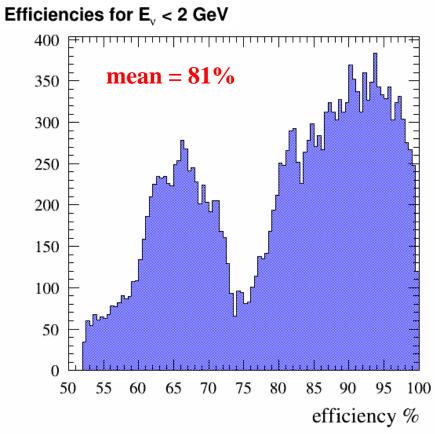


but of course $E_v < 2$ GeV is of particular interest



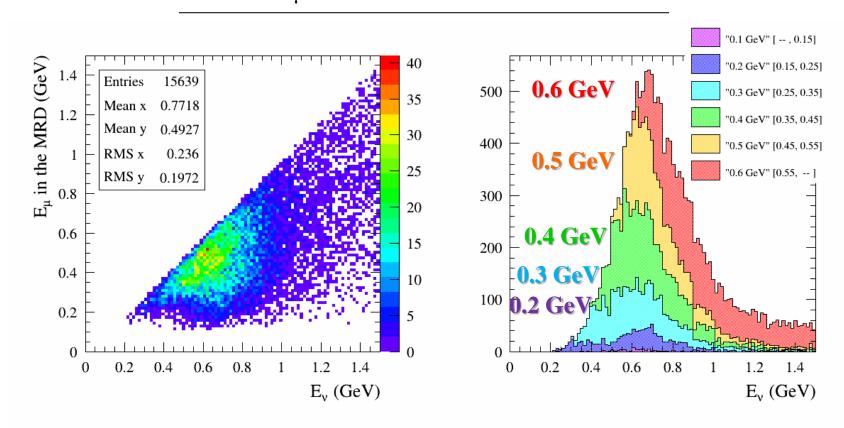
Here we expect ~80% efficiency

PRELIMINARY



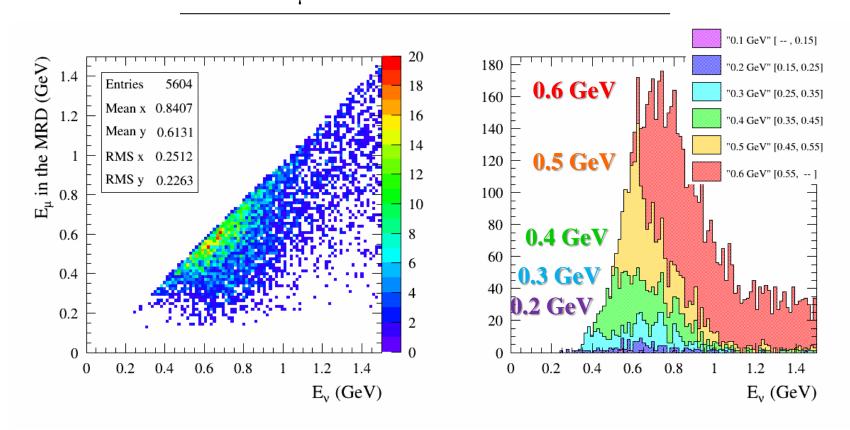
Let's think more carefully about this region...

What is the E_{μ} composition of the <u>side</u> MRD?



34 of events with muons which leave the tank

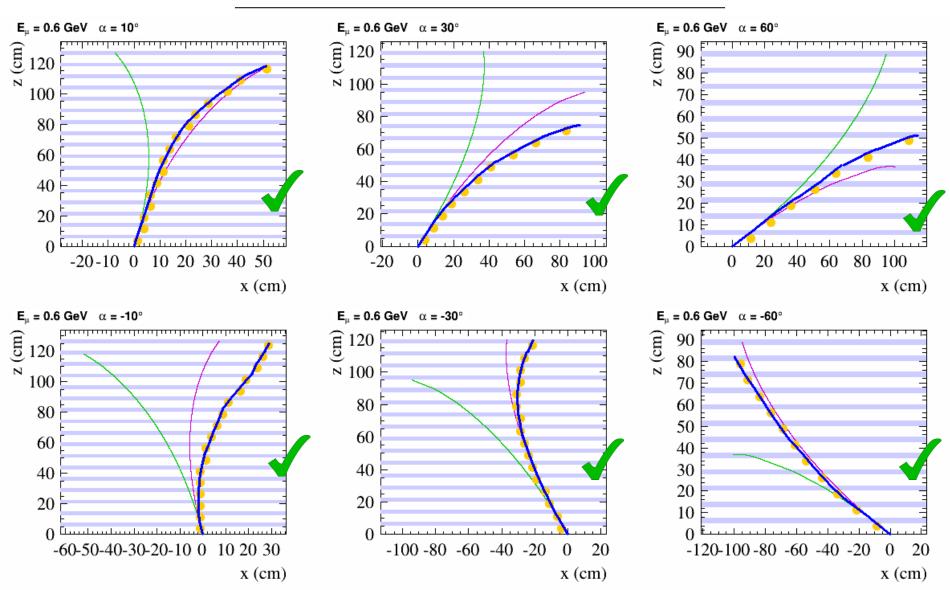
What is the E_{μ} composition of the end MRD?



More forward muons have slightly higher energies

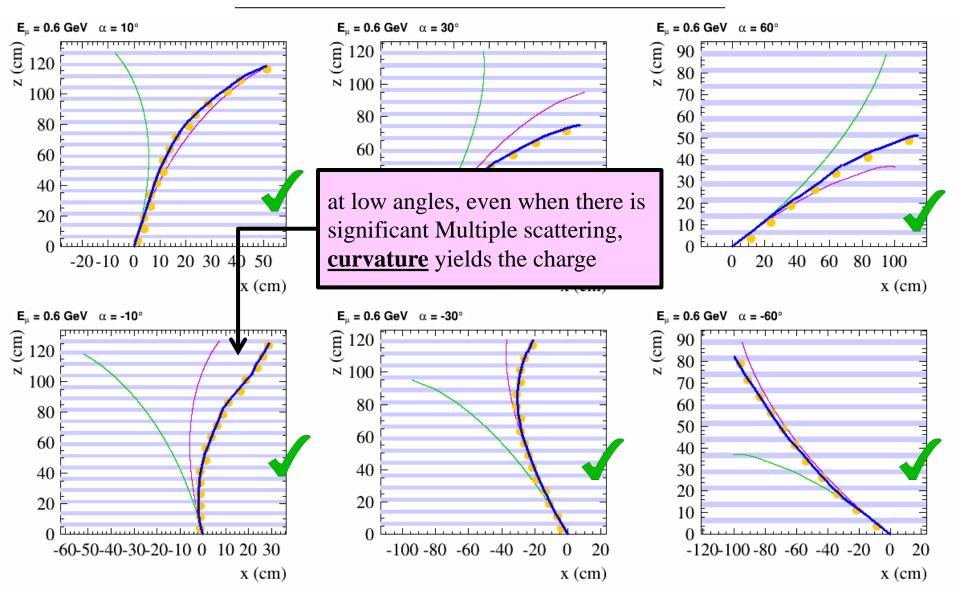
$$E_{\mu} = 0.6 \text{ GeV}$$

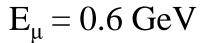
56% of END muons 32% of SIDE muons



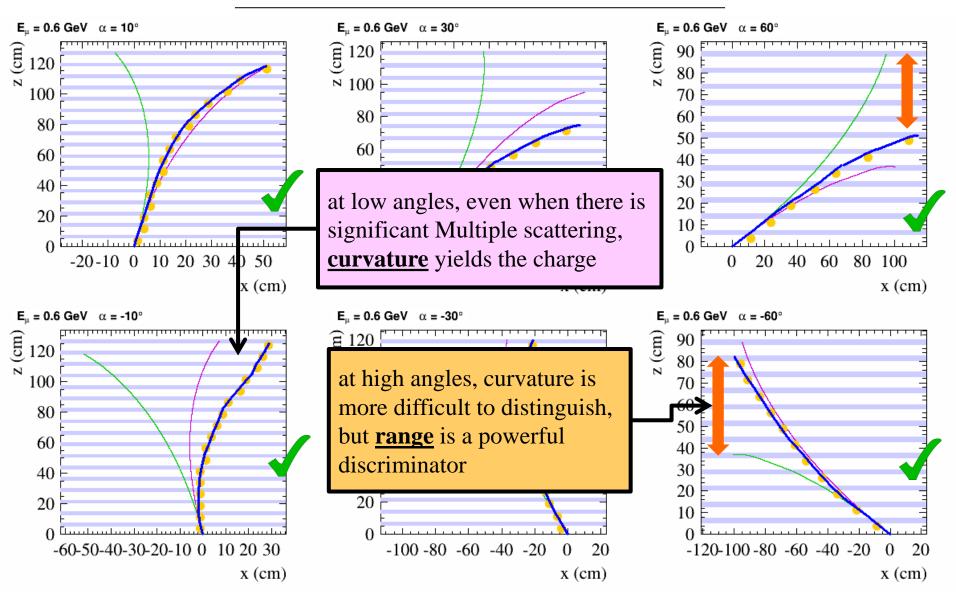


56% of END muons 32% of SIDE muons



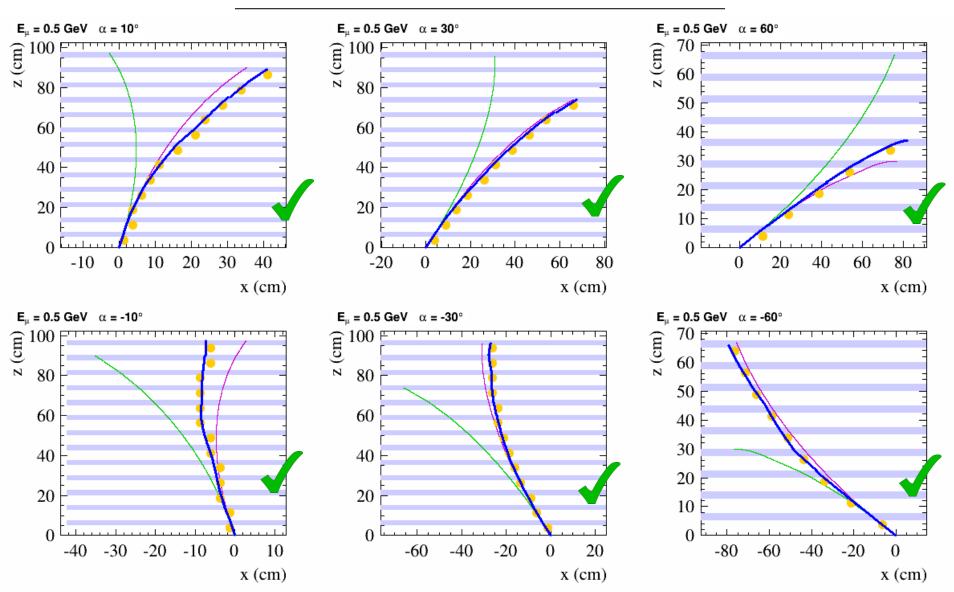


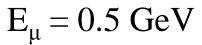
56% of END muons 32% of SIDE muons



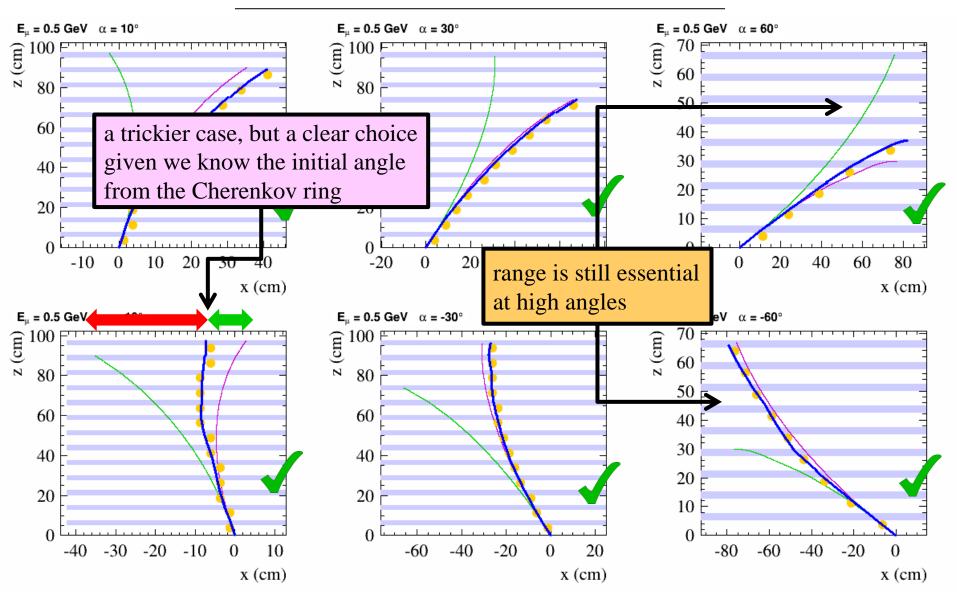
$$E_{\mu} = 0.5 \text{ GeV}$$

20% of END muons 21% of SIDE muons



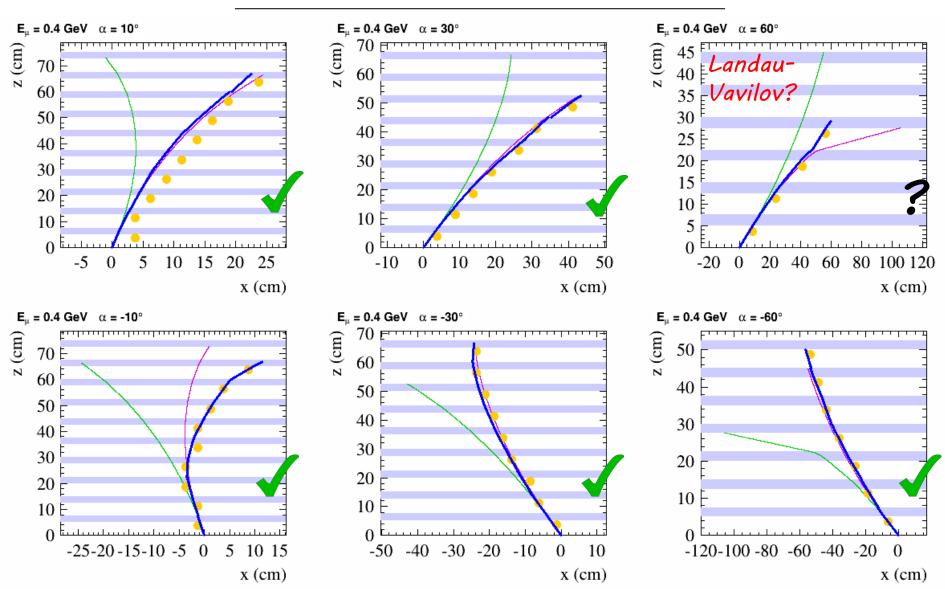


20% of END muons 21% of SIDE muons



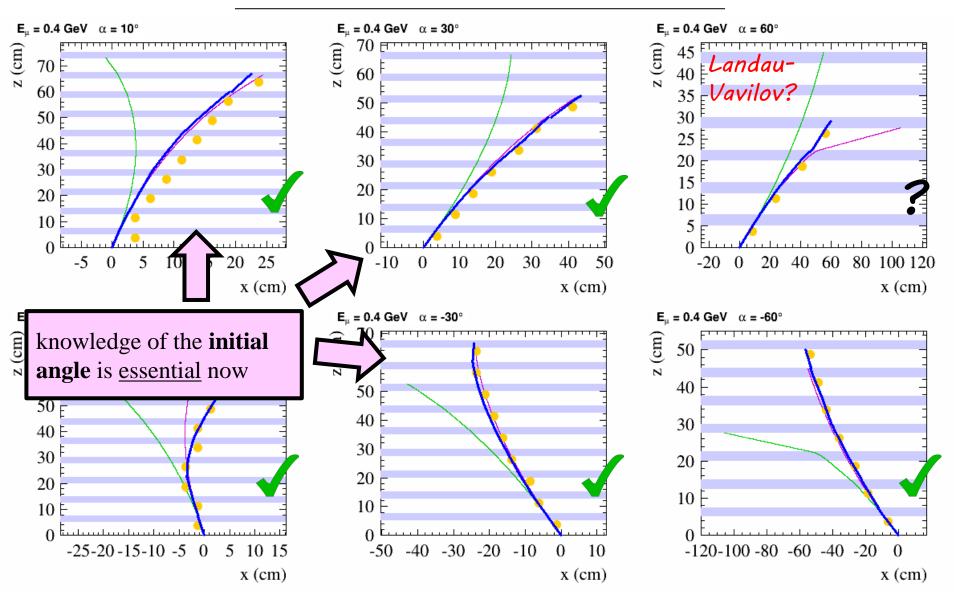
$$E_{\mu} = 0.4 \text{ GeV}$$

14% of END muons 24% of SIDE muons



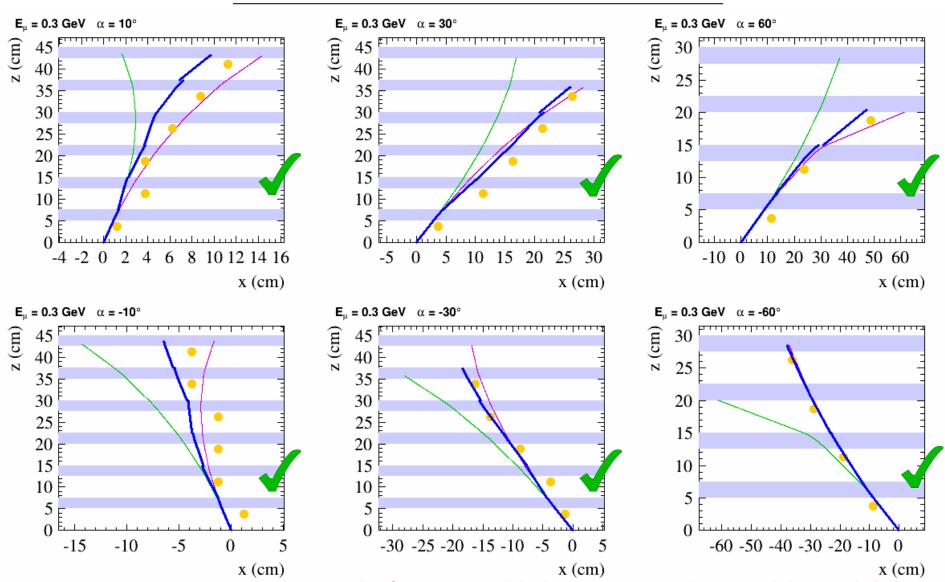


14% of END muons 24% of SIDE muons





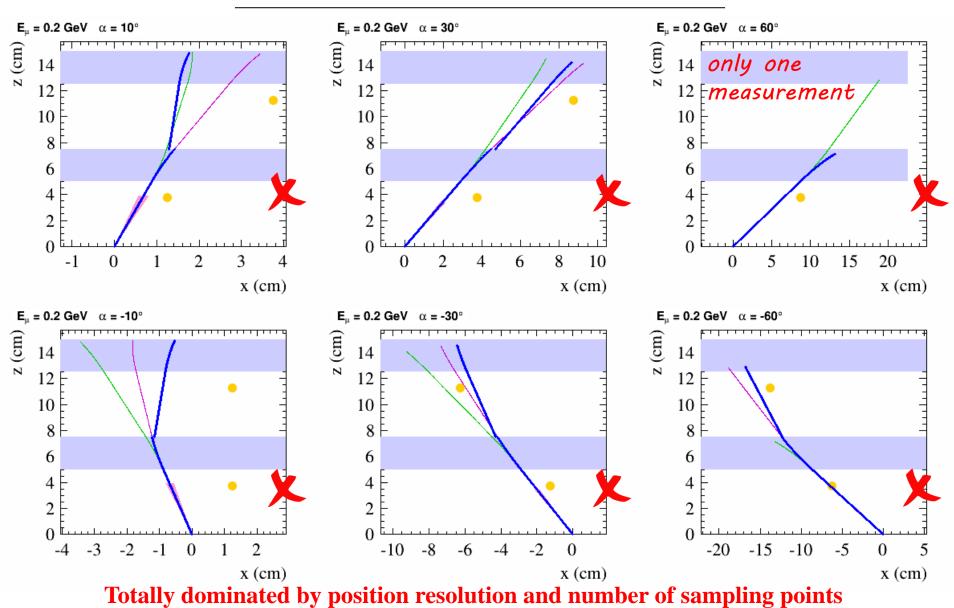
7% of END muons 17% of SIDE muons



Here I suspect we got a bit lucky! It's a probabilistic game, but still promising at 0.3 GeV

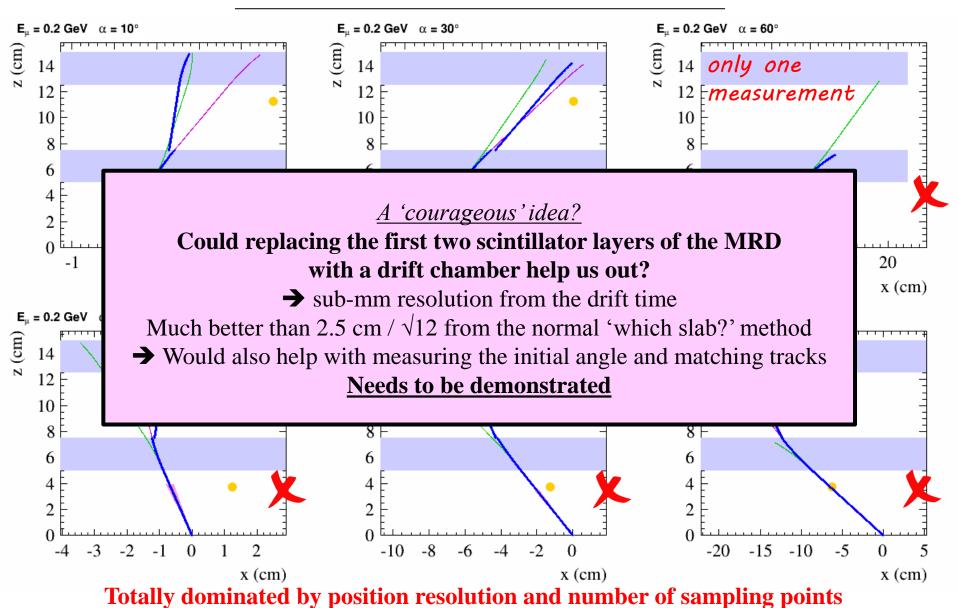


2% of END muons 7% of SIDE muons

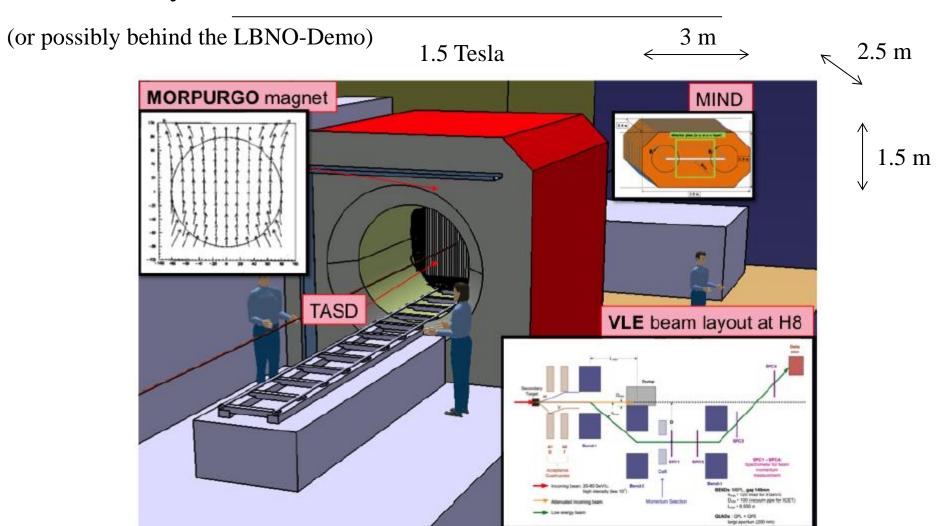




2% of END muons 7% of SIDE muons



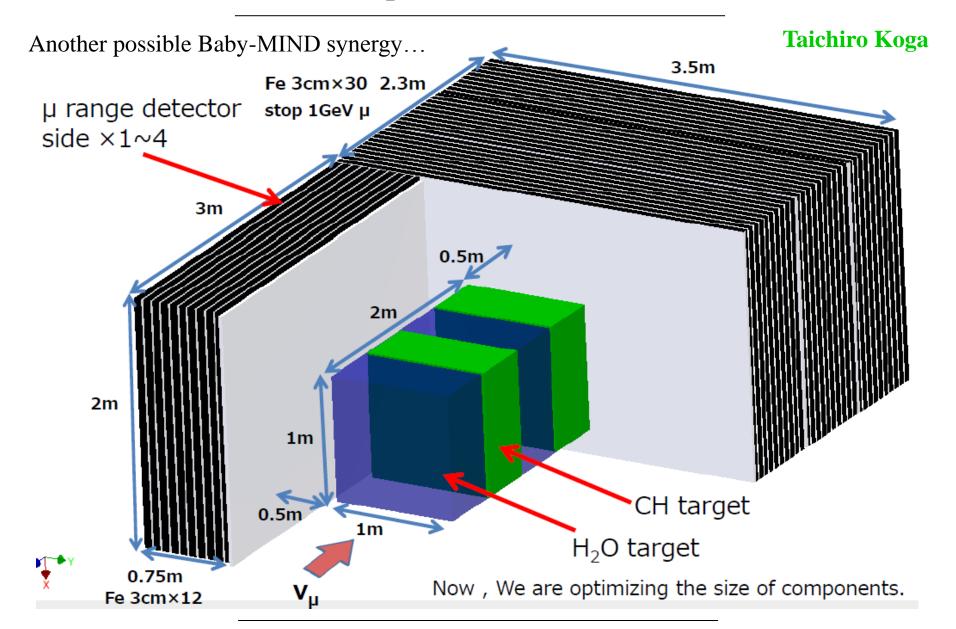
Baby-MIND and TASD: H8 beamline in North Area



Could also be a test of the TITUS MRD charge reconstruction

Contact: Etam Noah, University of Geneva

The B2 experiment / 'WAGASCI'



Summary

- 18% of muons escape the 22 m long, 11 m diameter TITUS tank 75% through the sides, 25% through the end
- With 150 cm of iron at the end, 50 cm of iron at the sides:
 - \rightarrow 75% of muons which escape the tank are stopped
 - \rightarrow 25% of muons which escape the tank penetrate through the MRD
- Preliminary studies show promising charge reconstruction in the oscillation region and impeccable resolution in the high energy tail (1.5 Tesla)

Work in progress

- Find the effect on δ_{CP} sensitivity
- Optimization of scintillator and irons layer thicknesses
- Answers to practical questions, such as PMT shielding
- The last lever: consider re-optimising the tank size and MRD size simultaneously

Backup slides

TITUS tank angle reconstruction?

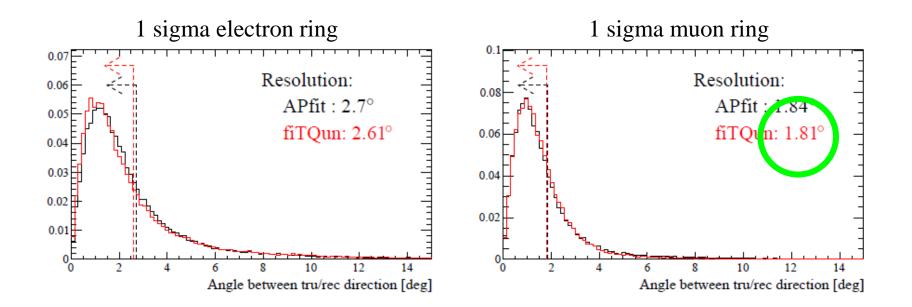
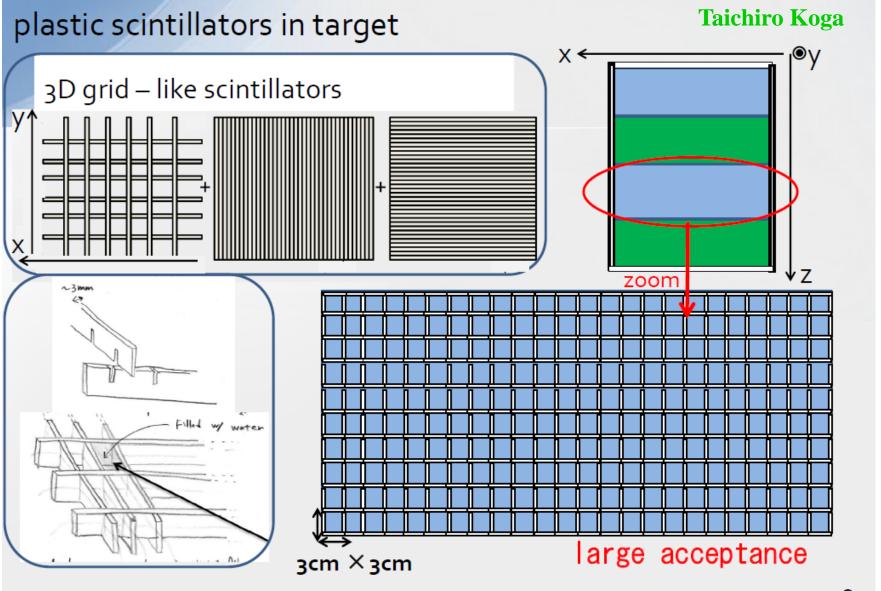


Figure 22: Distributions of the angle between the true and the reconstructed particle directions, for single-ring electron(left) and muon(right) particle gun events. The red histograms are the distributions for fiTQun, and the black histograms are for APfit. The resolutions are defined as the 68.3 percentiles, which are indicated by the dashed arrows.

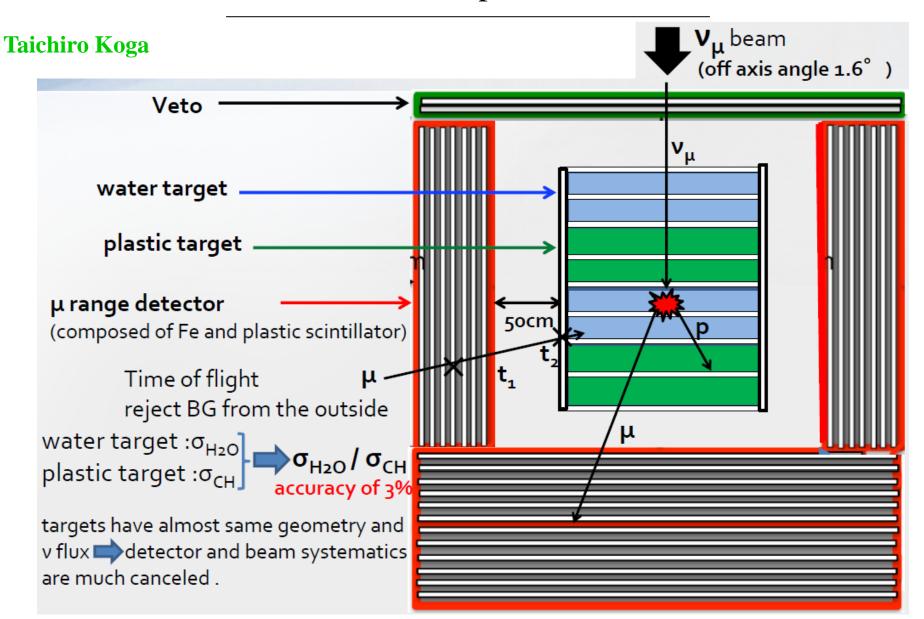
from the fitQun technical note



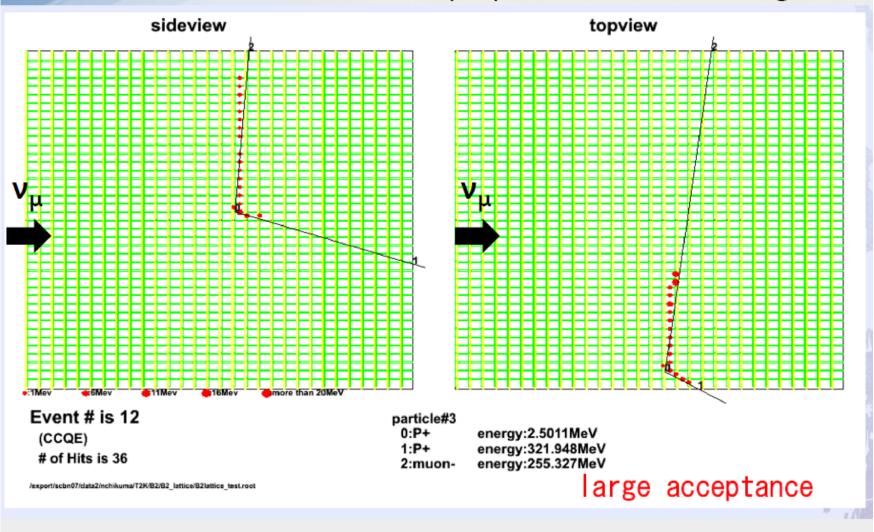
• We use thin plastic scintillators(~3mm) to increase water ratio in target.

Now **H2O:CH=70:30.** If the size of grid is changed to 5cm × 5cm, **H2O:CH=80:20.**

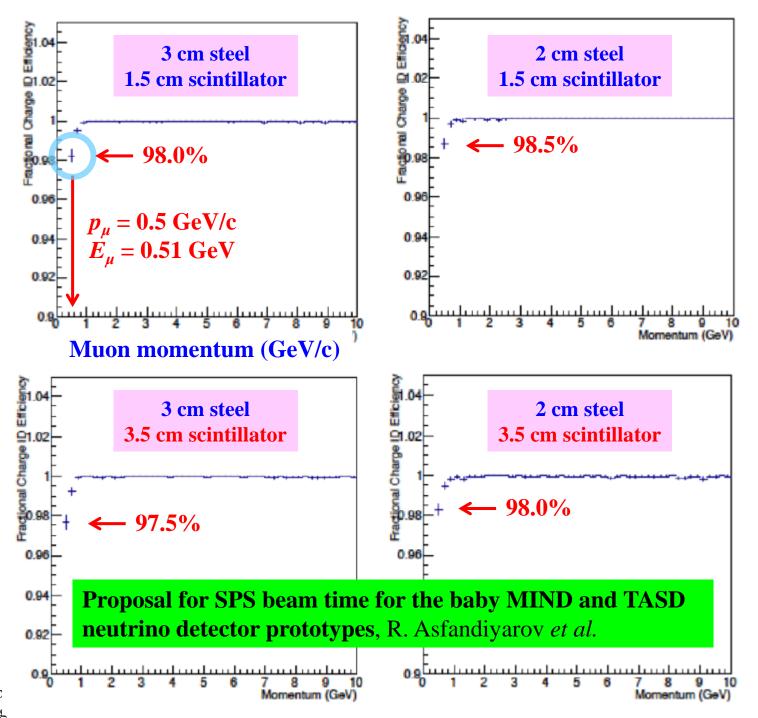
The B2 experiment

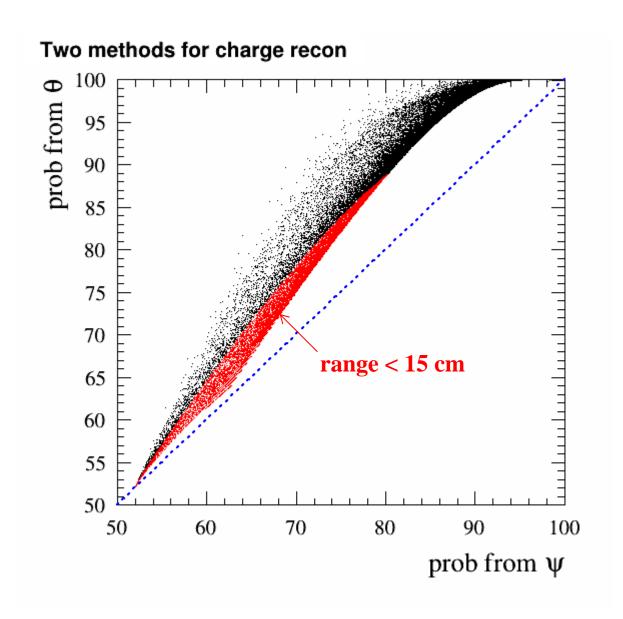


plastic scintillators in target event display(*Xchikuma san's figure)

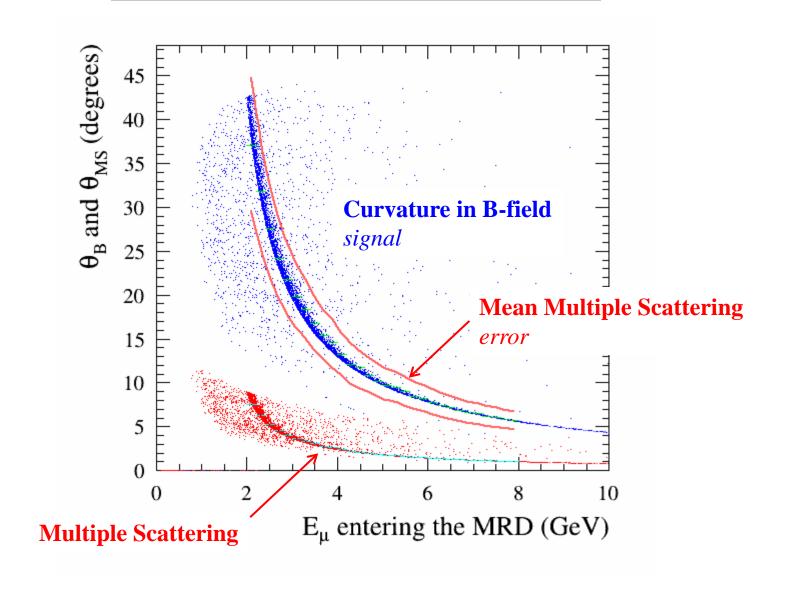


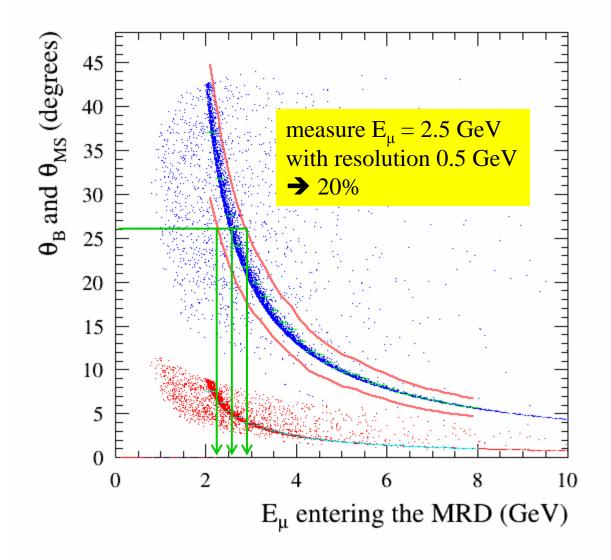
Taichiro Koga

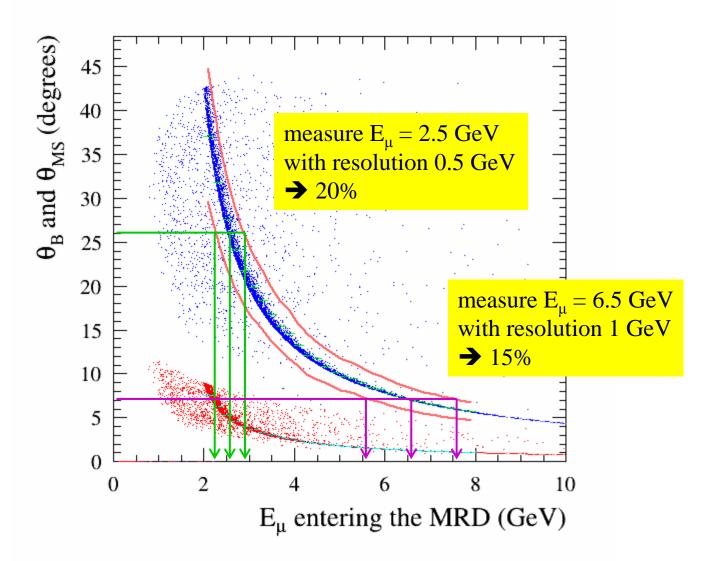




Aside: Momentum for the penetrating sample







Very conservative estimates

Landau-Vavilov most probable energy loss in iron

$$\int density = 7.87 g cm^{-2}$$

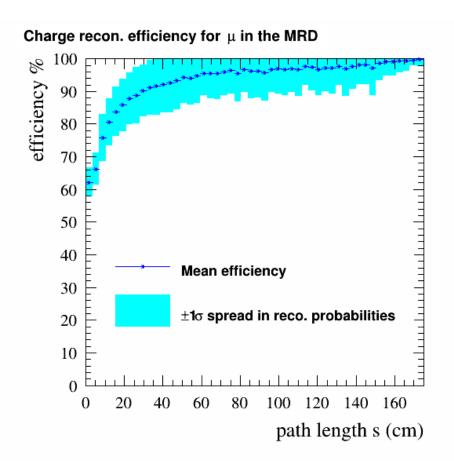
$$\xi = (K/2) \langle Z/A \rangle (x/\beta^2) \text{ MeV } \sim \text{1.13 MeV/cm (ultra-relativistic)}$$

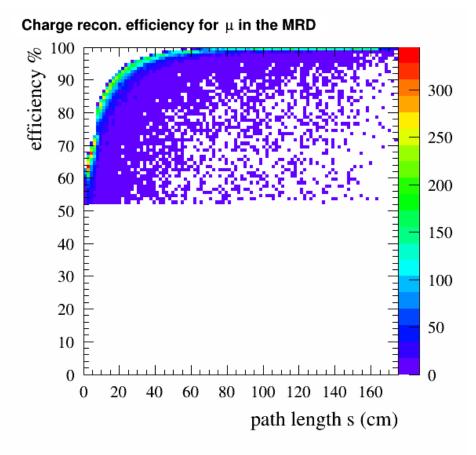
$$Z/A = 26/55.845 = 0.466$$

$$Z/A > \rho \text{ ratio (\simenergy loss / cm) = 1.4\%}$$

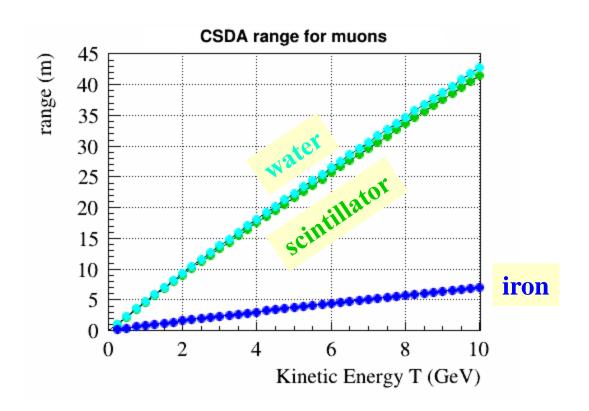
$$0.511 \text{ MeV}$$
 neglect density effect}
$$\Delta_p = \xi \left[\ln \frac{2mc^2\beta^2\gamma^2}{I} + \ln \frac{\xi}{I} + j - \beta^2 - \delta(\beta\gamma) \right]$$

$$0.200 \text{ all materials}$$
 Mean excitation energy
$$I = 286.0 \text{ eV in iron}$$





Range of muons in iron



Fiducial volume cut = 1 m with LAPPDs = 0.5 m

PDG 32.11. Measurement of particle momenta in a uniform magnetic field

The trajectory of a particle with momentum p (in GeV/c) and charge ze in a constant magnetic field \overline{B} is a helix, with radius of curvature R and pitch angle λ . The radius of curvature and momentum component perpendicular to \overrightarrow{B} are related by

where
$$B$$
 is in tesla and R is in meters. (32.49)

The distribution of measurements of the curvature $k \equiv 1/R$ is approximately Gaussian. The curvature error for a large number of uniformly spaced measurements on the trajectory of a charged particle in a uniform magnetic field can be approximated by

trajectory of a charged particle in a uniform magnetic field can be approximated by
$$(\delta k)^2 = (\delta k_{\rm res})^2 + (\delta k_{\rm ms})^2 , \qquad (32.50)$$

 $\delta k_{\rm ms} = {\rm curvature\ error\ due\ to\ multiple\ scattering.}$ If many (≥ 10) uniformly spaced position measurements are made along a trajectory

in a uniform medium,
$$\delta k_{\rm res} = \frac{\epsilon}{L'^2} \sqrt{\frac{720}{N+4}} \;, \tag{32.51}$$

(32.51)

where
$$N =$$
 number of points measured along track

 $\delta k = \text{curvature error}$

L' = the projected length of the track onto the bending plane $\epsilon = \text{measurement error for each point, perpendicular to the trajectory.}$

 $\delta k_{\rm res} = {\rm curvature\ error\ due\ to\ finite\ measurement\ resolution}$

If a vertex constraint is applied at the origin of the track, the coefficient under the radical becomes 320.

For arbitrary spacing of coordinates s_i measured along the projected trajectory and with variable measurement errors ϵ_i the curvature error $\delta k_{\rm res}$ is calculated from:

$$(\delta k_{\rm res})^2 = \frac{4}{w} \frac{V_{ss}}{V_{ss}V_{s^2s^2} - (V_{ss^2})^2} , \qquad (32.52)$$

where V are covariances defined as $V_s m_s n = \langle s^m s^n \rangle - \langle s^m \rangle \langle s^n \rangle$ with $\langle s^m \rangle = w^{-1} \sum (s_i^m / \epsilon_i^2)$ and $w = \sum \epsilon_i^{-2}$.

The contribution due to multiple Coulomb scattering is approximately

$$\delta k_{\rm ms} \approx \frac{(0.016)({\rm GeV}/c)z}{Lp\beta\cos^2\lambda}\sqrt{\frac{L}{X_0}}$$
, (32.53)

where p = momentum (GeV/c)

z =charge of incident particle in units of e

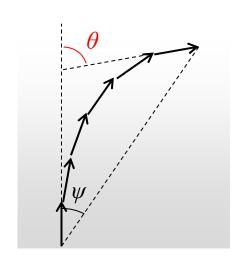
L = the total track length

 X_0 = radiation length of the scattering medium (in units of length; the X_0 defined elsewhere must be multiplied by density)

 β = the kinematic variable v/c.

More accurate approximations for multiple scattering may be found in the section on Passage of Particles Through Matter (Sec. 31 of this Review). The contribution to the curvature error is given approximately by $\delta k_{\rm ms} \approx 8 s_{\rm plane}^{\rm rms}/L^2$, where $s_{\rm plane}^{\rm rms}$ is defined there.





The uniform magnetic field B = 1.5T is in the z direction The particle moves along a curve of length s in the (x,y) plane

$$\mathrm{d}p_{\perp}/\mathrm{d}t = B \ q \ \mathrm{d}s/\mathrm{d}t$$

$$\Delta p_{\perp} = B \ q \ \Delta s$$

Take uniform steps of $\Delta s = 1$ cm

$$\Delta p_{\perp} = 4.5 \text{ MeV/c (for every cm)}$$

And hence the angle curved, depending on E at the time

 ΔE using most probable Landau-Vavilov value (Bethe overestimates due to long tails)

Charge identification for the muon if

 θ > Multiple Scattering

Muon path length in the iron of the MRD

