Determination of the Systematics Related to Polarization Angle Uncertainty and Its Impact on *r*

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Overview

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[Polarization Angle Uncertainty](#page-2-0)

- \triangleright One of the major systematics to be controlled is the uncertainty of the polarization sensitive angle of a polarimeter.
- Miscalibrated polarization angles induce a mixing of *E* and *B*-modes which obscures the primordial *B*-mode signal.
- ► In-flight polarization angle calibration strategies only reach a $\sim 0.5^\circ$ uncertainty.¹
- lacktriangleright Weintroduce an iterative angular power-spectra maximum Likelihood-based method to calculate the polarization angles ($\bar{\alpha}$) from the multi-frequency signal by nulling the *EB* power-spectra.

^{1&}lt;br>This uncertainty might be lower by a dedicated observational campaign for a better characterization of some astrophysical sources such as the Crab Nebula.

[Formalism](#page-4-0)

The observed *E* and *B* spectra of the *i*-th channels are rotated as follows:

$$
\begin{pmatrix} E_i \\ B_i \end{pmatrix} = \begin{pmatrix} \cos(2\alpha_i) & -\sin(2\alpha_i) \\ \sin(2\alpha_i) & \cos(2\alpha_i) \end{pmatrix} \begin{pmatrix} E_i^* \\ B_i^* \end{pmatrix},
$$

The full log-likelihood (assumed to be Gaussian) of all possible *EB*-spectra is:

$$
-2\log \mathcal{L} = \sum_{i=1}^{N_{ch}} \sum_{j=i}^{N_{ch}} \sum_{i'=1}^{N_{ch}} \sum_{j'=i'}^{N_{ch}} \Delta_{\ell}^{y} (\mathbf{C}^{-1})_{\ell, \vec{y}, i'j'} \Delta_{\ell}^{i'j'},
$$

where $\Delta_{\ell}^{\vec{y}} = (\chi_{\ell}^{\vec{y}} - \xi_{\ell}^{\vec{y}} \tan(2\alpha_i + 2\alpha_j))$ and,

$$
\chi_{\ell}^{\vec{y}} = \frac{C_{\ell}^{E_i B_j} + C_{\ell}^{B_i E_j}}{2}, \qquad \xi_{\ell}^{\vec{y}} = \frac{C_{\ell}^{E_i E_j} - C_{\ell}^{B_i B_j}}{2}
$$

It is a modification of the Minami et al. 2019 methodology.

Assuming:

- **If** Small angle approximation: $\alpha \ll 1 \rightarrow \tan(\alpha) \sim \alpha$.
- **In** The covariance matrix **C** does not depend on α . To correct the mismatch induced by this approximation, we perform an iterative approach that updates the polarization angle in the covariance matrix with the one estimated in the previous step.

With these approximations, we achieve a linear system which enable us to obtain analytical equations to calculate the rotation angles. Moreover, the uncertainties can be evaluated from the Fisher matrix. This results in a very fast computational methodology.

Formalism III. Covariance Matrix Implementation

- \blacktriangleright The covariance matrix depends on $\widetilde{C}_{\ell}^{XY}$ which is the ensemble average *XY* power-spectrum $(X, Y \in \{E, B\})$
- \triangle Due to the lack of a reliable model of the foregrounds cross-spectra, we require a good estimator of it from the observed power-spectra.
- \triangleright We have tried several estimators:
	- The observed power-spectra.
	- The observed power-spectra binned.
	- \bullet The observed power-spectra smooth with a square window ℓ -function convolution.
- \blacktriangleright Any of these previous approximations for the foregrounds cross-spectra model produce biases when a full covariance matrix including all the possible cross power-spectra are considered. Therefore in this work we use only the auto spectra.

Formalism IV. Rotation Angles Equations

The rotation angles are obtained solving the following system:

$$
\Omega\bar{\alpha}=\frac{1}{4}\bar{\eta}\,,
$$

where Ω and $\bar{\eta}$ elements are

$$
\Omega_{ij} = \sum_{\ell=\ell_{min}}^{\ell_{max}} \xi_{\ell}^{i} \left(\mathbf{C}^{-1} \right)_{\ell, jj} \xi_{\ell}^{j},
$$

$$
\eta_{i} = \sum_{j=1}^{N_{ch}} \sum_{\ell=\ell_{min}}^{\ell_{max}} \xi_{\ell}^{i} \left(\mathbf{C}^{-1} \right)_{\ell, jj} C_{\ell}^{EB_{j}}.
$$

And the Fisher matrix is calculated with:

$$
F_{ij} = \frac{1}{2} \frac{\partial^2 (-2 \log \mathcal{L})}{\partial \alpha_i \partial \alpha_j} = 16 \sum_{\ell = \ell_{min}}^{\ell_{max}} \xi_{\ell}^i \left(\mathbf{C}^{-1}\right)_{\ell, j} \xi_{\ell}^j.
$$

[Simulations](#page-9-0)

Simulations with Systematic

- Rotation Angles.
- \blacktriangleright Sky signal
	- CMB.
	- Polarized Foregrounds: synchrotron and dust.
- \blacktriangleright Instrumental Noise.

$$
\begin{pmatrix} \bar{Q}^{rot} \\ \bar{U}^{rot} \end{pmatrix} = \begin{pmatrix} \cos(2\bar{\alpha}) & -\sin(2\bar{\alpha}) \\ \sin(2\bar{\alpha}) & \cos(2\bar{\alpha}) \end{pmatrix} \begin{pmatrix} \bar{Q} \\ \bar{U} \end{pmatrix} + \begin{pmatrix} \bar{n}^{\mathcal{Q}} \\ \bar{n}^{\mathcal{U}} \end{pmatrix}
$$

Rotation Angles

Requirements with channels correlations.

The covariance among channels:

$$
\mathbf{C}_{\alpha} = \begin{pmatrix} M_{LFT} & 0 & 0 \\ 0 & M_{MFT} & 0 \\ 0 & 0 & M_{HFT} \end{pmatrix} \,,
$$

Using **L** the Cholesky decomposition of **C**_α, the rotation angles:

 $\bar{\alpha}_{ch} = \mathbf{L}\bar{x}$,

where \bar{x} random samples from the standard Gaussian distribution.

Based on the requirements shown by Enrique Martínez-González on Monday.

Sky Signal. CMB and foregrounds

CAMB. Lewis et al., astro-ph/9911177/ Planck 2018 results. VI. Cosmological parameters, 1807.06209 Thorne et al., 1608.02841 (PySM)

Instrumental Noise. LiteBIRD-like Experiment

[Validation with simulations](#page-14-0)

Convergence

Multipole Range

Multipole Range

Covariance Model. Estimator

- \blacktriangleright all ℓ : original power-spectra. $(\ell_{min} = 3 - \ell_{max} = 600)$
- binned (uniform nbin=50): power-spectra binned in ℓ chunks of equal length (*n*[|]*ell* ∼ 12 per bin).
- binned (log nbin=10): power-spectra binned in ℓ chunks with logarithmically increasing length.
- smooth: power-spectra smoothed by two subsequent convolutions with square window functions ($n_\ell = 5$, $n_e = 10$).

Uncertainty and Bias

Independent Telescope Study

[Impact on](#page-25-0) *r*

We perform component separation on the following set of maps:

- **In Non-rotated maps**. The original signal maps without any rotation.
- **Rotated maps**. The previous maps with the corresponding rotation applied.
- \triangleright De-rotated maps. The resulting maps after de-rotating the rotated maps with the solutions obtained with the described methodologies.

To study the impact introduced by the systematic.

Parametric Component Separation

The likelihood for a given pixel is:

$$
\mathcal{L}(\theta_p|\bar{d}_p) = \frac{1}{\sqrt{(2\pi)^{2N_{ch}}\det(\mathbf{C})}}\exp\left(-\frac{1}{2}\left(\bar{d}_p - m(\bar{\nu};\theta_p)\right)^T\mathbf{C}^{-1}\left(\bar{d}_p - m(\bar{\nu};\theta_p)\right)\right),\,
$$

The sky model is:

The amplitudes and spectral parameters are updated in different steps. Convergence is obtained quickly.

Component separation methodology used is an improved version of the one presented in de la Hoz, E., et al. JCAP 2020.06 (2020): 006

Map Processing before component separation

Power-spectra residuals comparison

[One minute summary](#page-30-0)

One Minute Summary

- ► Current polarization angles calibration strategies have uncertainties of $\sim 0.5^\circ$. Primordial B-mode search requires more stringent limits.
- I We introduce an iterative angular power-spectra maximum likelihood-based method to calculate the polarization angles $(\bar{\alpha})$ from the multi-frequency signal by nulling the C_{ℓ}^{EB} .
- \blacktriangleright Two major assumptions are made: i) the rotation angles are small (\leq 6deg), and, ii) the covariance matrix does not depend on $\bar{\alpha}$.
- \blacktriangleright We obtain an analytical linear system which leads to a very fast computational implementation.
- \triangleright With this methodology we reach uncertainties on the order of a few arc minutes.
- \triangleright We show that this accuracy is enough to remove the systematics by applying a parametric component separation technique to recover the CMB in three scenarios: i) rotated signal, ii) non-rotated signal, and iii) de-rotated signal.
- \triangleright We find that the systematic introduced by leftover polarization angles in r is removed after the signal is corrected.