Towards ending partial sky *E-B* ambiguity in CMB observations

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1 minute Summary

- Leakage of power from *E*-to-*B* modes on an incomplete sky patch is a critical problem for search for primordial B modes.
- We demonstrate that it is possible to remove this leakage contribution by completing the E-mode ٠ information in suborbital experiments with ancillary *E*-mode information from satellite experiments.
- For Planck E-modes which have considerably higher noise than current and upcoming ground-based ٠ mission, we can Wiener filter the E-modes.
- We combine the filtered *E*-mode only QU maps from Planck outside the observed sky patch and show ٠ that we effective remove the impact of leakage.
- We show that for AliCPT-like or LSPE-like cases we outperform pure-B PCL estimator giving unbiased, ٠ with optimal PCL error bars.
- For CMB-S4, with signal dominated LiteBIRD *E*-mode data, the *E*-to-*B* leakage can be nearly eliminated. ۲

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This talk is based on arXiv: 2007.09928

Work done in collaboration with:

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[Ref: Planck Picture Gallery]

100

Multipole

250 500

1000

0.0

10

30

2000

10 30

100 250 500 1000

Multipole

Introduction

- Planck mission is complete and LiteBIRD is still little far in the future
- In between there are large number of suborbital missions both in northern and southern hemisphere
- Goal is to observe the primordial *B* modes in a small sky patch with low noise and relatively free from foregrounds
- The Ali CMB Polarization Telescope will begin observing the northern sky sometime next year at 95 and 150 GHz



Ali CMB Polarization Telescope site



We observe Stokes Q and U fields in experiments

 $P_{\pm}(\hat{n}) = Q(\hat{n}) \pm i U(\hat{n})$

The Stokes Q and U fields are not scalars. P_{\pm} represents a spin-(\pm 2) field.

$$P_{\pm}(\hat{n}) = \sum_{\ell m} a_{\pm 2,\ell m \pm 2} Y_{\ell m}(\hat{n})$$

The *E-B* decomposition is unique for FULL SKY FIELDS ONLY!

Complication for incomplete sky.

The spherical harmonic coefficients for *E*- and *B*-fields are defined as:

$$\begin{aligned} a_{E,\ell m} &= -\frac{1}{2} \left(a_{2,\ell m} + a_{-2,\ell m} \right) = -\frac{1}{2} \left[\int P_+(\hat{n})_2 Y_{\ell m}^*(\hat{n}) d\Omega + \int P_-(\hat{n})_{-2} Y_{\ell m}^*(\hat{n}) d\Omega \right] \\ a_{B,\ell m} &= \frac{i}{2} \left(a_{2,\ell m} - a_{-2,\ell m} \right) = \frac{i}{2} \left[\int P_+(\hat{n})_2 Y_{\ell m}^*(\hat{n}) d\Omega - \int P_-(\hat{n})_{-2} Y_{\ell m}^*(\hat{n}) d\Omega \right]. \end{aligned}$$

The E-B ambiguity

For an incomplete sky observation the E-B mode definition becomes:

$$\begin{split} \tilde{a}_{E,\ell m} &= -\frac{1}{2} \left[\int P_+(\hat{n}) W(\hat{n})_2 Y_{\ell m}^*(\hat{n}) d\Omega + \int P_-(\hat{n}) W(\hat{n})_{-2} Y_{\ell m}^*(\hat{n}) d\Omega \right], \\ \tilde{a}_{B,\ell m} &= \frac{i}{2} \left[\int P_+(\hat{n}) W(\hat{n})_2 Y_{\ell m}^*(\hat{n}) d\Omega - \int P_-(\hat{n}) W(\hat{n})_{-2} Y_{\ell m}^*(\hat{n}) d\Omega \right]. \end{split}$$

$$\begin{split} & \text{Window} \\ \text{defining the sky} \\ \text{patch} \end{split}$$

Rewriting partial sky *E-B* harmonics to full sky *E-B* harmonics:

$$\begin{split} \tilde{a}_{E,\ell m} &= \sum_{\ell' m'} \left[\mathcal{K}_{\ell m\ell' m'}^{EE} a_{E,\ell' m'} + i \mathcal{K}_{\ell m\ell' m'}^{EB} a_{B,\ell' m'} \right], \\ \tilde{a}_{B,\ell m} &= \sum_{\ell' m'} \left[-i \mathcal{K}_{\ell m\ell' m'}^{BE} a_{E,\ell' m'} + \mathcal{K}_{\ell m\ell' m'}^{BB} a_{B,\ell' m'} \right]. \end{split} \qquad \mathcal{K}_{\ell m'}^{X}$$

 $K_{\ell m \ell' m'}^{XY} \longrightarrow$ Kernel matrix



5 μ K-arcmin N_{ℓ}^{P} (19')

10²

10¹

10³

that there are no foreground residuals or time stream filtering.

The E-to-B Leakage



E-to-B Leakage Control

Standard PCL method:

$$\langle \tilde{C}_{\ell}^{EE} \rangle = \sum_{\ell'} \left[M_{\ell\ell'}^{EE} \langle C_{\ell'}^{EE} \rangle + M_{\ell\ell'}^{EB} \langle C_{\ell'}^{BB} \rangle \right]$$

$$\langle \tilde{C}_{\ell}^{BB} \rangle = \sum_{\ell'} \left[M_{\ell\ell'}^{BE} \langle C_{\ell'}^{EE} \rangle + M_{\ell\ell'}^{BB} \langle C_{\ell'}^{BB} \rangle \right]$$

'pure' fields:

$$\begin{split} \mathcal{E}(\hat{n}) &= -\frac{1}{2} \left[\bar{\eth} \bar{\eth} P_{+}(\hat{n}) + \eth \eth P_{-}(\hat{n}) \right] \\ \mathcal{B}(\hat{n}) &= \frac{i}{2} \left[\bar{\eth} \bar{\eth} P_{+}(\hat{n}) - \eth \eth P_{-}(\hat{n}) \right] \end{split}$$

Pure-*B* method:

$$\tilde{a}_{B,\ell m} = \sqrt{\frac{(\ell-2)!}{(\ell+2)!}} \int W(\hat{n}) \mathcal{B}(\hat{n}) Y_{\ell m}^* d\Omega$$



Combined *E*-map method

The *E*-to-*B* leakage originates from completeness of the *E*-mode information.

So if we use full sky *E*-mode maps from Planck to fill-in the *E*-mode information, it can reduce and possibly remove leakage



However, Planck is too noisy for direct combination.

Wiener filtered E-maps



Combined *E*-map

- 1. Use Wiener filtered E-modes to generate QU maps at ground-based exp. resolution.
- 2. Combine the QU maps:

$$d_{\text{com}} = \widehat{d}_{\text{sat}} \left(1 - W_{\text{grd}} \right) + d_{\text{grd}} W_{\text{grd}}.$$

3. Get B-map for the combined QU map.











Optimal error band:

Our method results shown in red Pure-B method shown in blue Standard method in green

$$\Delta \mathcal{D}_{\bar{\ell};\text{optimal}} = \sqrt{\frac{2}{(2\bar{\ell}+1)f_{\text{sky}}\Delta\ell}} \left[\mathcal{D}_{\bar{\ell}} + \frac{\mathcal{N}_{\bar{\ell}}}{B_{\bar{\ell}}^2} \right] \left[\frac{w_{(4)}}{w_{(2)}^2} \right]^{1/2}$$

fsky = 8% patch, 10 uK arcmin noise; Full sky: isotropized 100-143-217 Planck noise



Optimal error band:

Our method results shown in red Pure-B method shown in blue

r = 0

$$\Delta \mathcal{D}_{\bar{\ell};\text{optimal}} = \sqrt{\frac{2}{(2\bar{\ell}+1)f_{\text{sky}}\Delta\ell}} \left[\mathcal{D}_{\bar{\ell}} + \frac{\mathcal{N}_{\bar{\ell}}}{B_{\bar{\ell}}^2} \right] \left[\frac{w_{(4)}}{w_{(2)}^2} \right]^{1/2}$$

r = 0.05



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fsky = 2% patch, 5 uK arcmin noise; Full sky: isotropized 100-143-217 Planck noise



Optimal error band:

Our method results shown in red Pure-B method shown in blue

r = 0

$$\Delta \mathcal{D}_{\bar{\ell};\text{optimal}} = \sqrt{\frac{2}{(2\bar{\ell}+1)f_{\text{sky}}\Delta\ell}} \left[\mathcal{D}_{\bar{\ell}} + \frac{\mathcal{N}_{\bar{\ell}}}{B_{\bar{\ell}}^2} \right] \left[\frac{w_{(4)}}{w_{(2)}^2} \right]^{1/2}$$

r = 0.05



fsky ~ 30% patch, combined SWIPE noise, 85' beam; Full sky: isotropized 100-143-217 Planck noise

μK -0.5 0.5



r = 0.05

How important is SNR?

The reconstruction ratio:

$$r_{\text{rec}} = \frac{1}{C_{\ell}^{EE} B_{E,\ell}^2} \left[\langle \widehat{d}_{E,\ell m} \mathcal{W}_{\ell'}^{EE} s_{E,\ell'm'}^* \rangle + \langle \widehat{d}_{E,\ell m} \mathcal{W}_{\ell'}^{ET} s_{T,\ell'm'}^* \rangle \right]$$

Quantifies how much of the E-mode signal is present after Wiener filtering.



Potential: LiteBIRD - CMB-S4



10²

10¹

10³

patch 95-145-155 GHz channel noise, 22.7' resolution

LiteBIRD - CMB-S4 Results

r = 0

r = 0.05



Optimal error band:

Our method results shown in red Pure-B method shown in blue

$$\Delta \mathcal{D}_{\bar{\ell};\text{optimal}} = \sqrt{\frac{2}{(2\bar{\ell}+1)f_{\text{sky}}\Delta\ell}} \left[\mathcal{D}_{\bar{\ell}} + \frac{\mathcal{N}_{\bar{\ell}}}{B_{\bar{\ell}}^2} \right] \left[\frac{w_{(4)}}{w_{(2)}^2} \right]^{1/2}$$

Summary and conclusions

- We demonstrate that it is possible to remove this leakage contribution by completing the *E*-mode information in suborbital experiments with ancillary *E*-mode information from satellite experiments.
- For Planck *E*-modes which have considerably higher noise than current and upcoming ground-based mission, we can Wiener filter the E-modes.
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- We show that for AliCPT-like or LSPE-like cases we outperform pure-*B* PCL estimator giving unbiased, with optimal PCL error bars.
- For CMB-S4, with signal dominated LiteBIRD *E*-mode data, the *E*-to-*B* leakage can be nearly eliminated.
- Next we need to update the method for inhomogeneous and correlated noise, with foreground residuals etc.