

Axion clouds may survive the perturbative tidal interaction over the early inspiral phase of black hole binaries

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Based on : T. Takahashi & T. Tanaka 2106.08836[gr-qc]

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Introduction

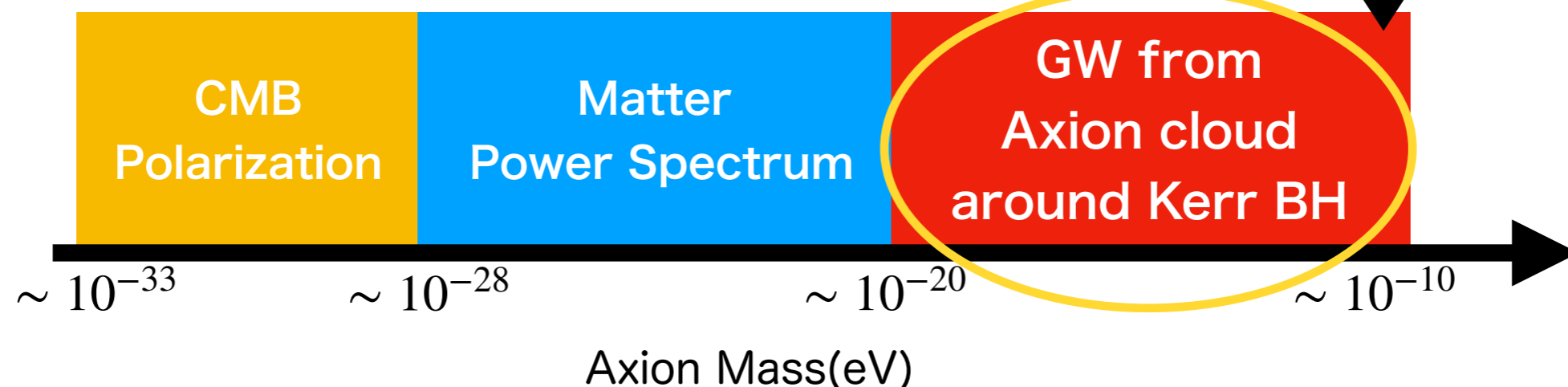
Axion & Astrophysics

The best-motivated extension of the Standard Model

- QCD/String axion
- Dark matter candidate

This talk

Cosmic phenomena used in axion exploration by axion mass

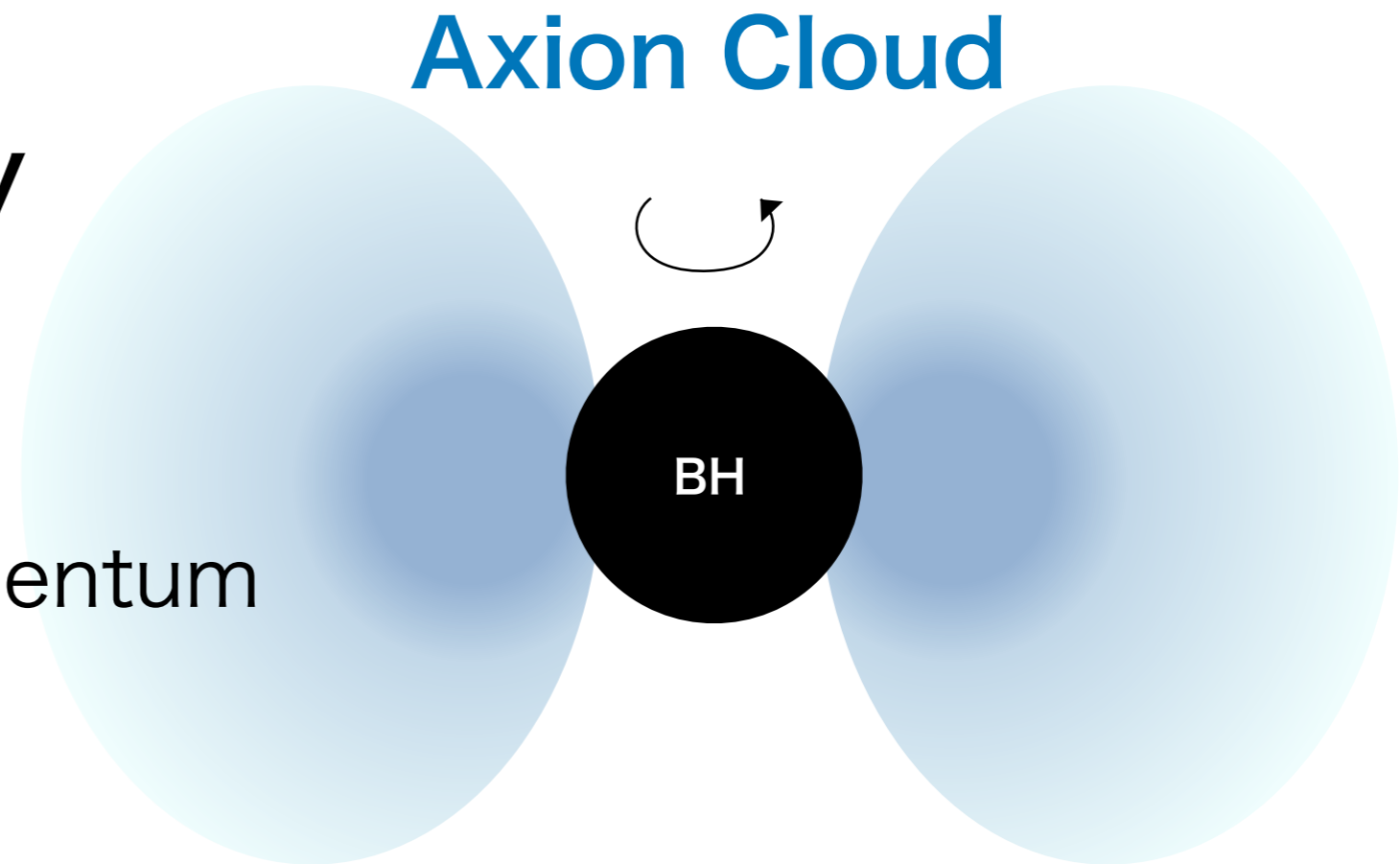


Rotating BHs and Axion

Superradiant instability

- This system admits quasi-bound solutions
- Energy and angular momentum extraction from BH

Axion from a cloud



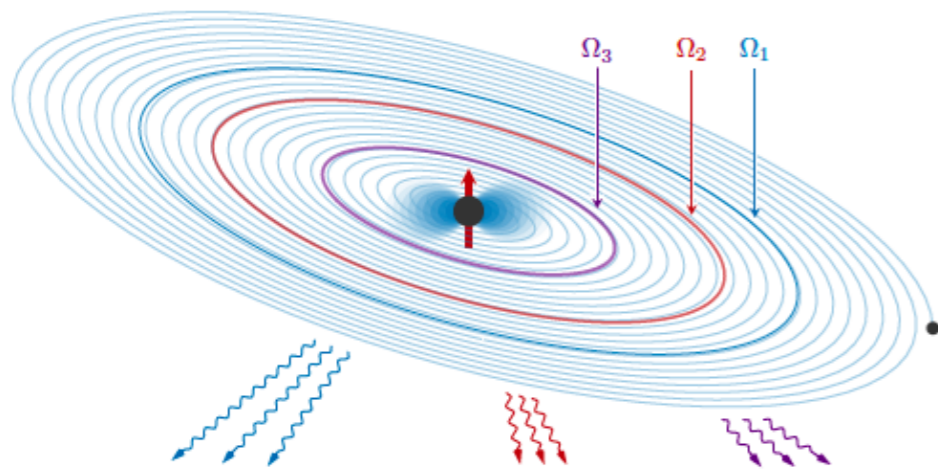
(BH radius M) \sim (Compton wavelength $1/\mu$)

$10^{-20} \sim 10^{-10} \text{eV}$

BHs and GWs observation
has the potential of probing axion

Axion Clouds in Binary systems

D.Baumann et al, 1912.04932



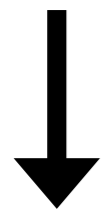
- Can we test axion with GW from binary coalescence?
- Does the axion cloud disappear due to binary's tidal interaction?

We made the exhaustive study of cloud depletion in a wide parameter region numerically for equal mass binaries.

Dynamics of the Axion Clouds

Isolated Axion Clouds

$$(g^{\mu\nu} \nabla_\mu \nabla_\nu - \mu^2) \phi = 0$$



non-relativistic approx. $M\mu \ll 1$

$$i \frac{\partial}{\partial t} \psi = \left[-\frac{1}{2\mu} \nabla^2 - \frac{M\mu}{r} \right] \psi$$

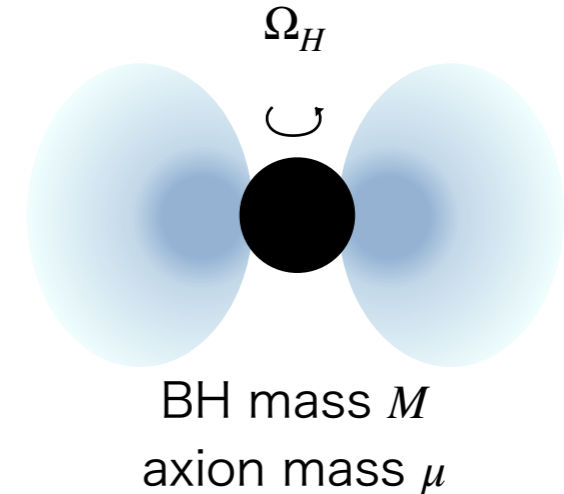
This is formally equivalent to the hydrogen atom in QM.

$$\omega_{nlm} = \underbrace{(\omega_R)_{nlm}}_{\text{Energy}} + i \underbrace{(\omega_I)_{nlm}}_{\text{Growth(+)/Decay(-)}}$$

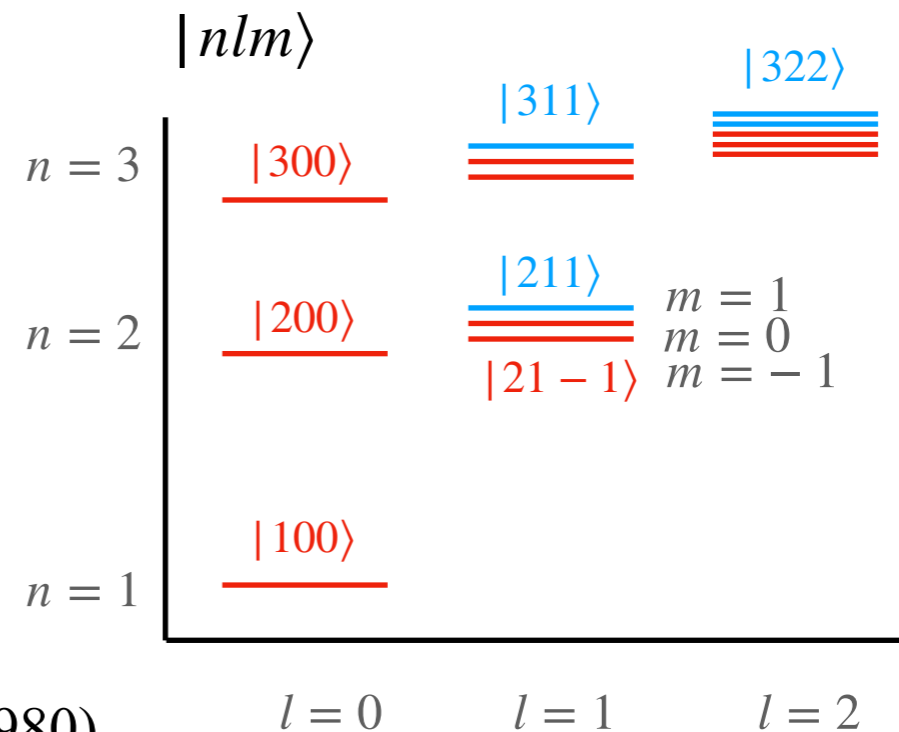
Energy **Growth(+)/Decay(-)**

$$(\omega_I)_{nlm} \propto \underbrace{(m\Omega_H - \mu)(M\mu)^{4l+5}}_{\text{superradiance condition}} \quad (\text{S.Detweiler, 1980})$$

superradiance condition



Energy level



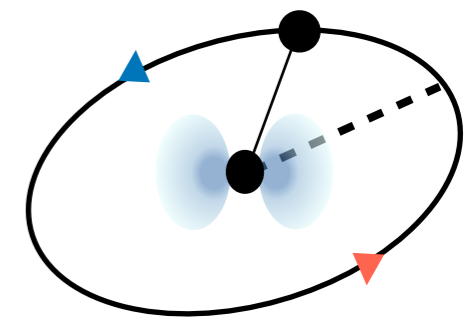
Binary's tidal perturbation

$$i \frac{\partial}{\partial t} \psi = \left[-\frac{1}{2\mu} \nabla^2 - \frac{M\mu}{r} + V_* \right] \psi$$

tidal potential

$$H = H_0 + V(t), \quad \underline{V(t) = \eta e^{-i\Omega t}}$$

Orbital freq. Ω



$|nlm\rangle$

311

310

31-1

211

210

21-1



counter-rotating

co-rotating

When orbital freq. coincides with the energy gap, **resonant level transition** occurs.

(D.Baumann et al, 1912.04932)

e.g.) co-rot. : $|211\rangle \rightarrow |21-1\rangle$

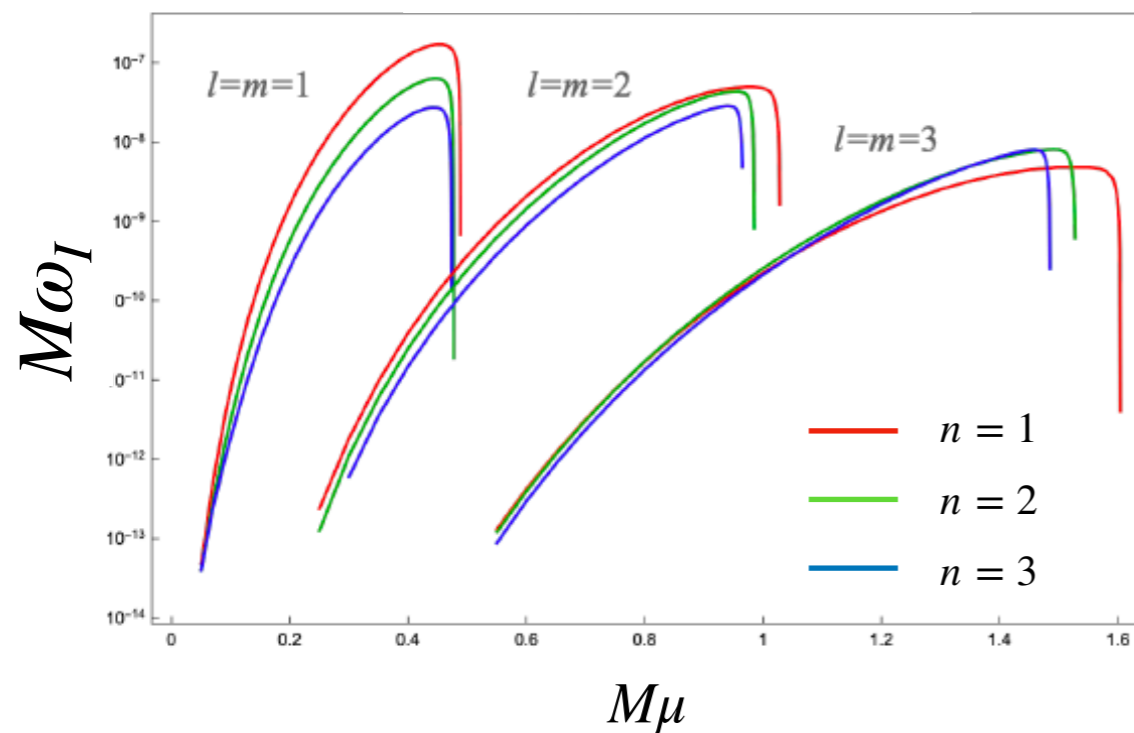
counter-rot. : $|211\rangle \rightarrow |31-1\rangle$

Axions in **the growing mode** are transferred to **the decaying mode**.

Is there any chance the Cloud won't disappear?

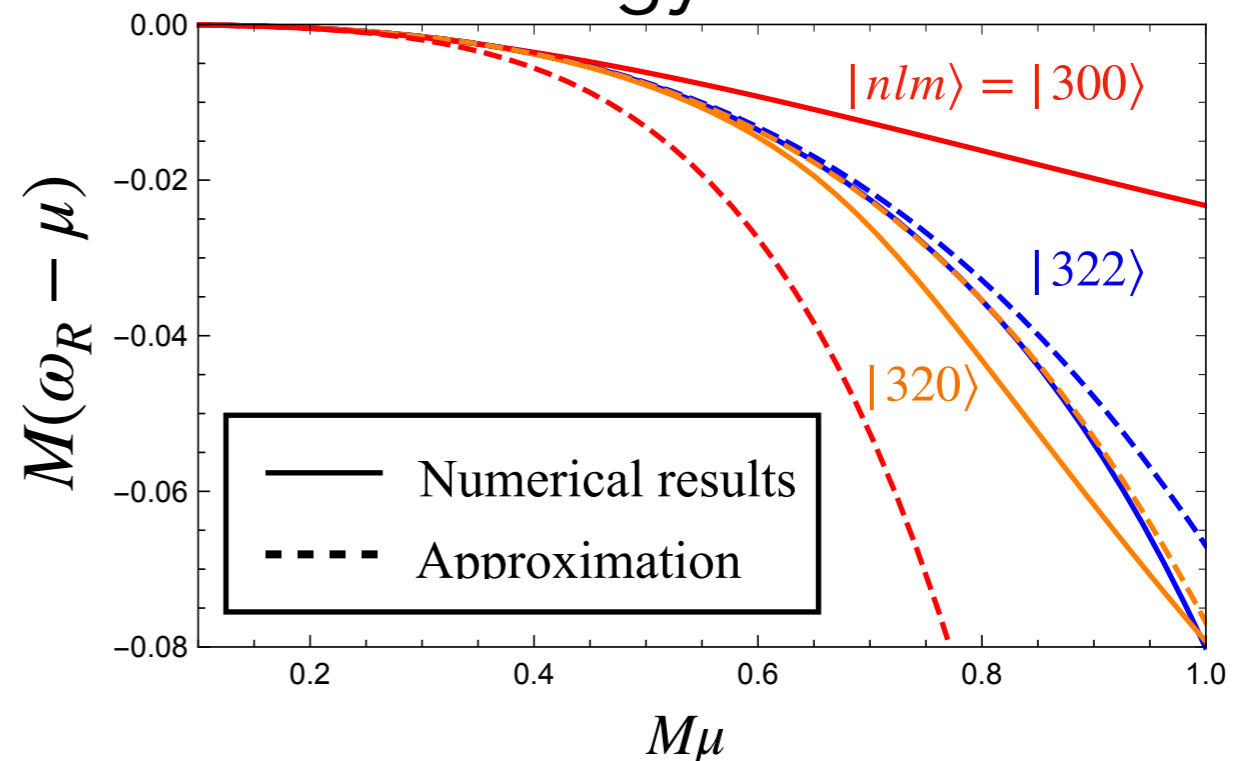
- We are interested in the case $M\mu \gtrsim 0.1$.
- Resonance freq. is determined by the energy gap.
- We need to calculate the eigenfrequencies numerically.

Growth rate



higher- l mode

Energy levels



non-trivial transition

Parameter : coupling and BH spin $(M\mu, a/M)$

Story

① Transition destination

Initially, axions occupy the fastest growing mode.

$$|n_0 l_0 m_0\rangle \rightarrow |? \rangle$$



- quadruplolar tidal perturbation V_*
- Selection rules

e.g.) non-rela approxi.

$$\text{co-rot. : } |211\rangle \rightarrow |21 - 1\rangle$$

$$|322\rangle \rightarrow |320\rangle$$

② How much fraction of the cloud will disappear

If transition is adiabatic, all axions are transferred to another mode.

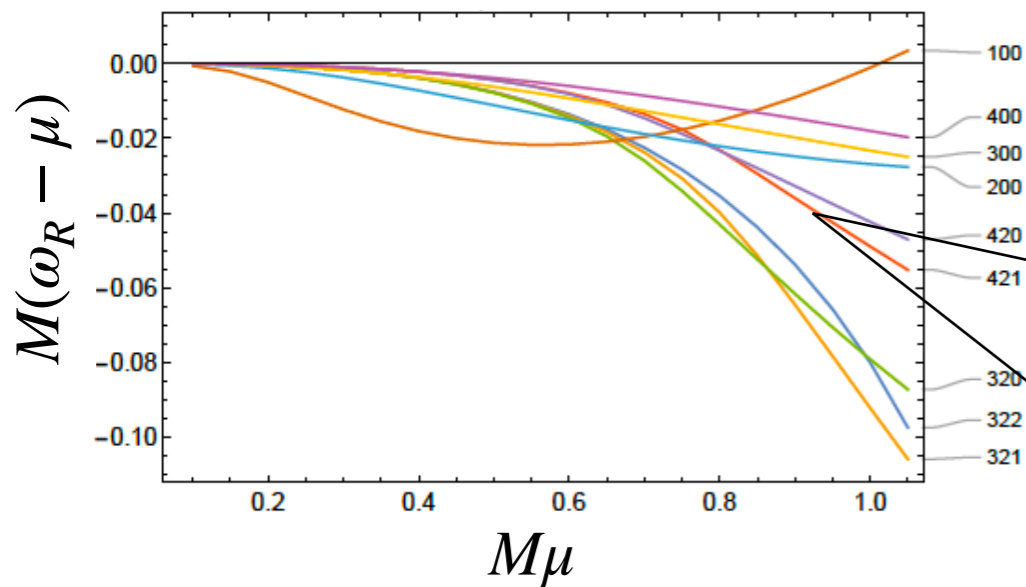
$$M_{cl} \simeq M_0 \exp \left(\underline{\omega_I} \underline{t_m} \right)$$

Decay width Time to merger

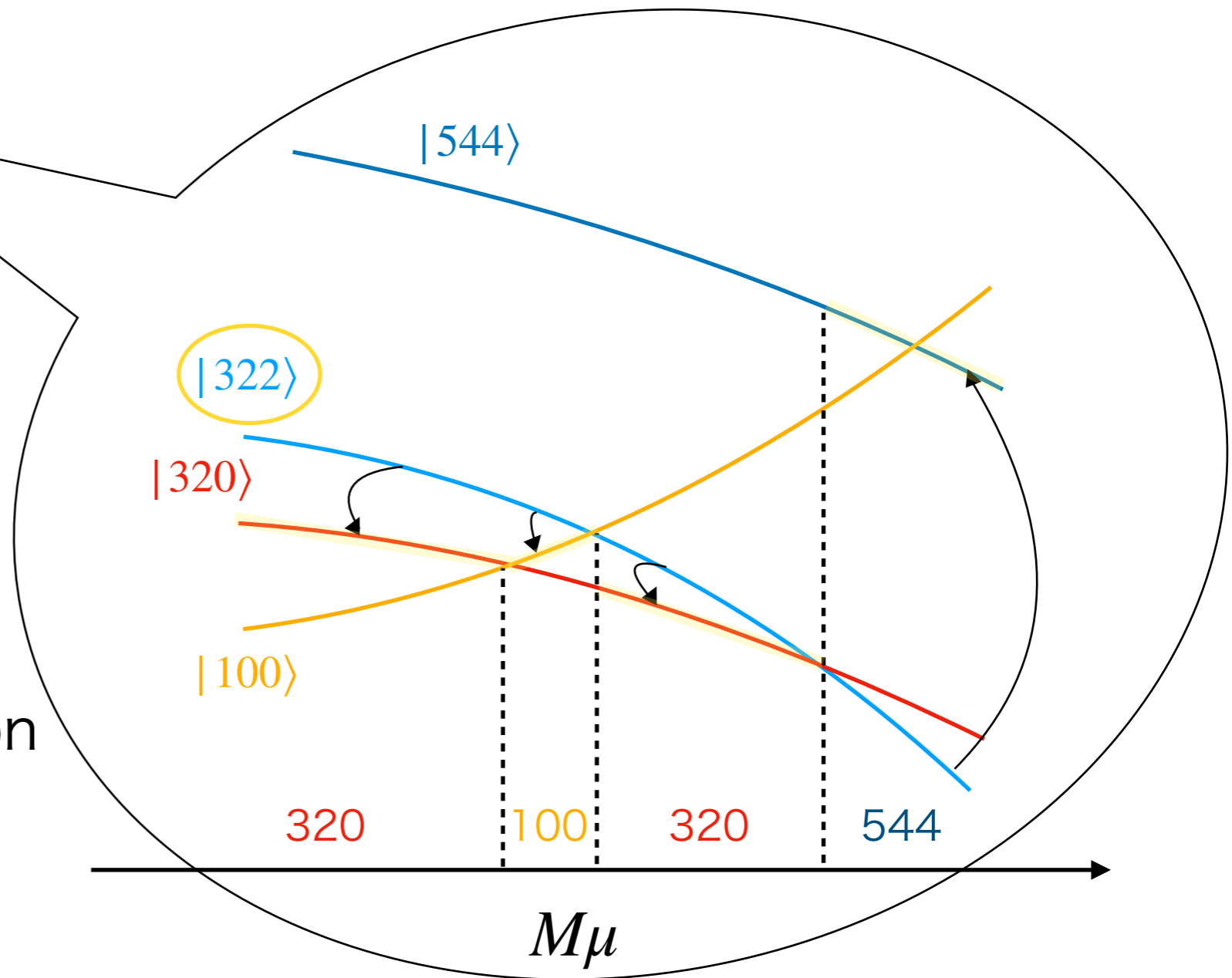
Transition destination

Example) $a/M = 0.998$, $l = m = 2$ is the fastest

for co-rotating $\frac{\Delta E}{\Delta m} > 0$

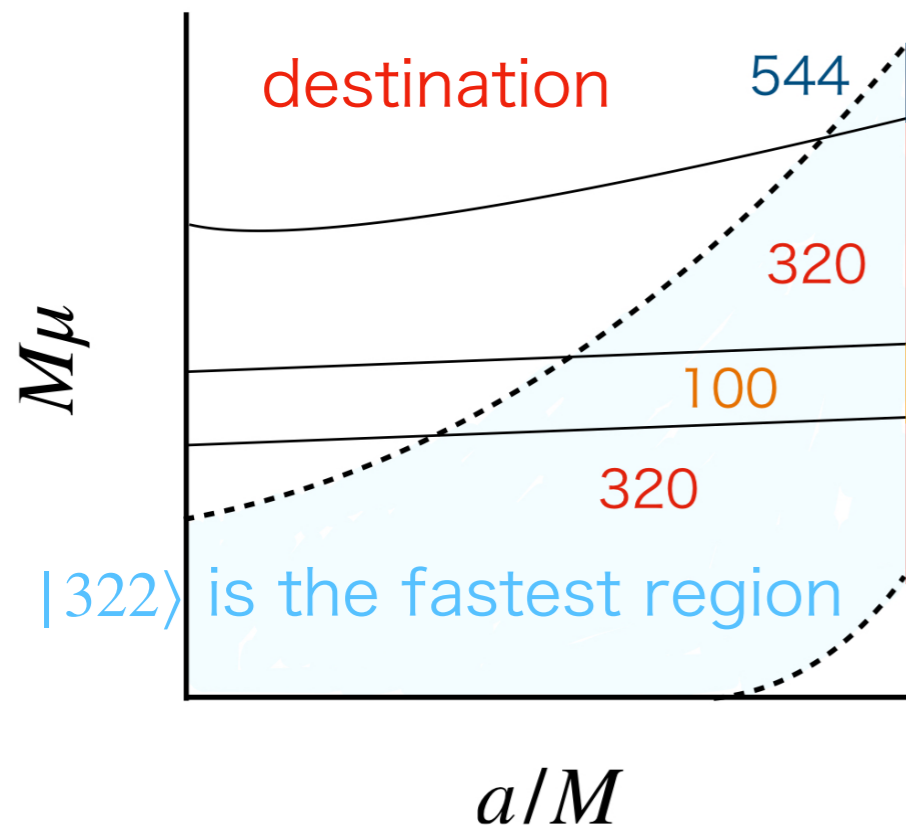


By comparing the energy gap between levels, we can read which level to transition to.

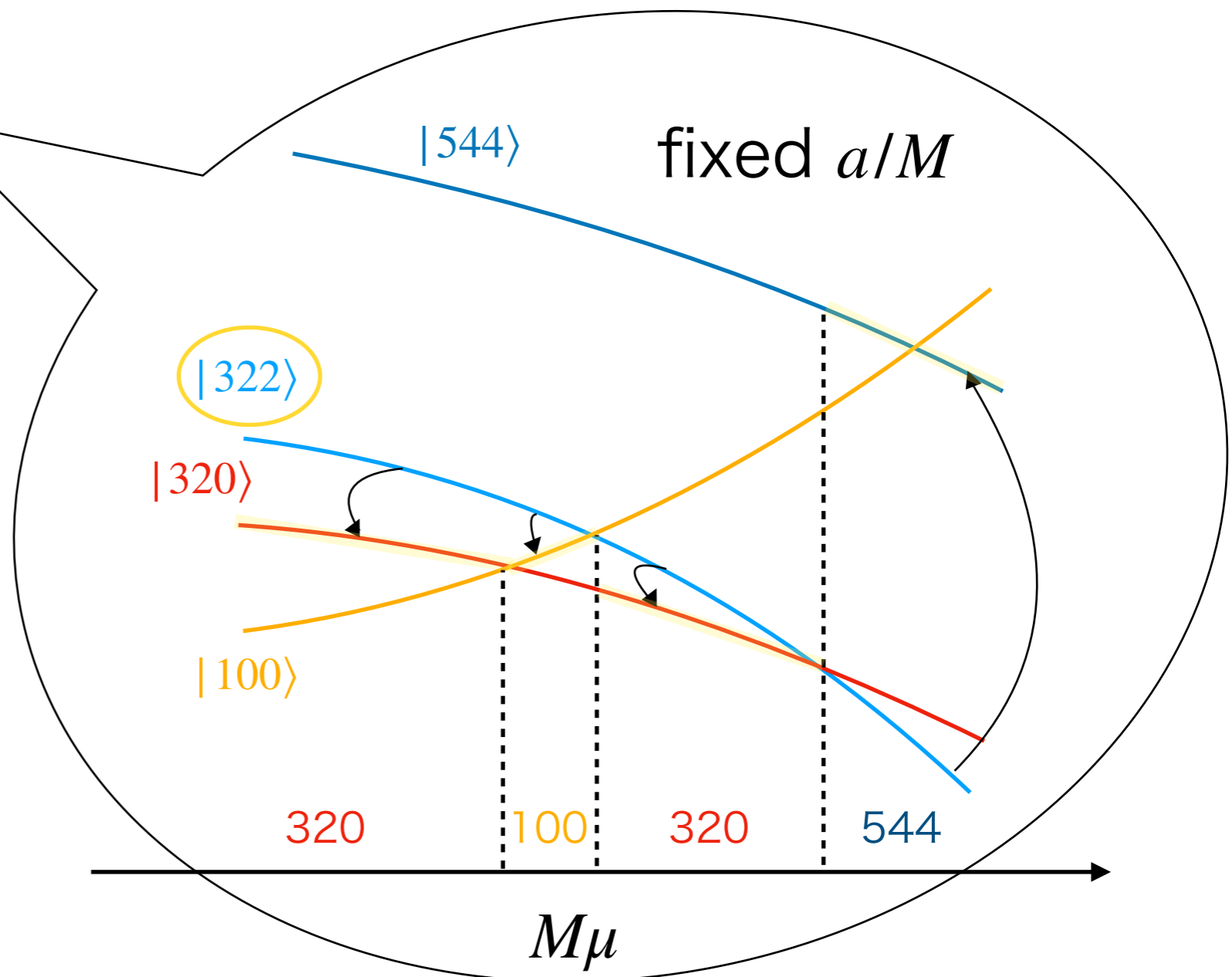


Transition destination

Parameter $(M\mu, a/M)$ plane



Transition destination changes where level crossing occur.



Transition Map

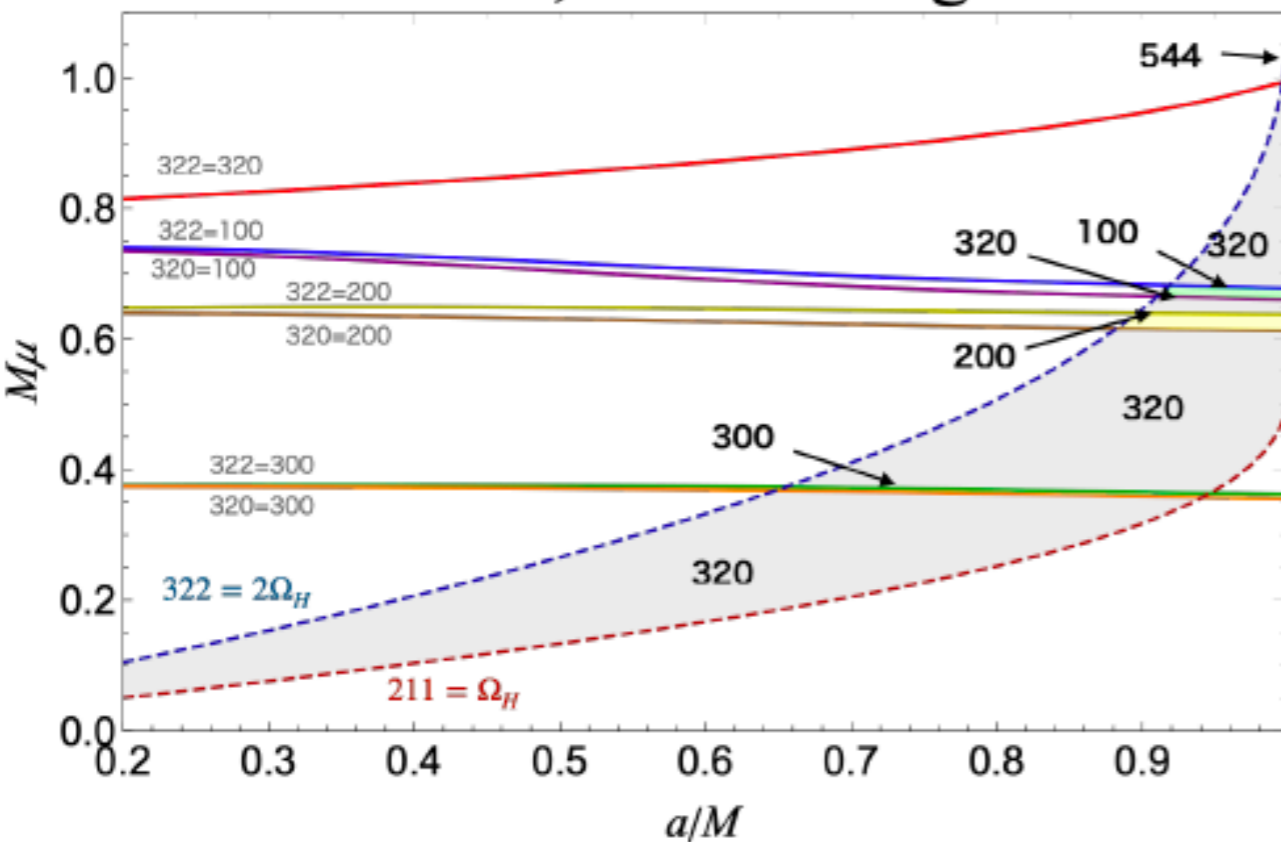
$$l = 1$$

Transition destinations are the same the non-rela approx.

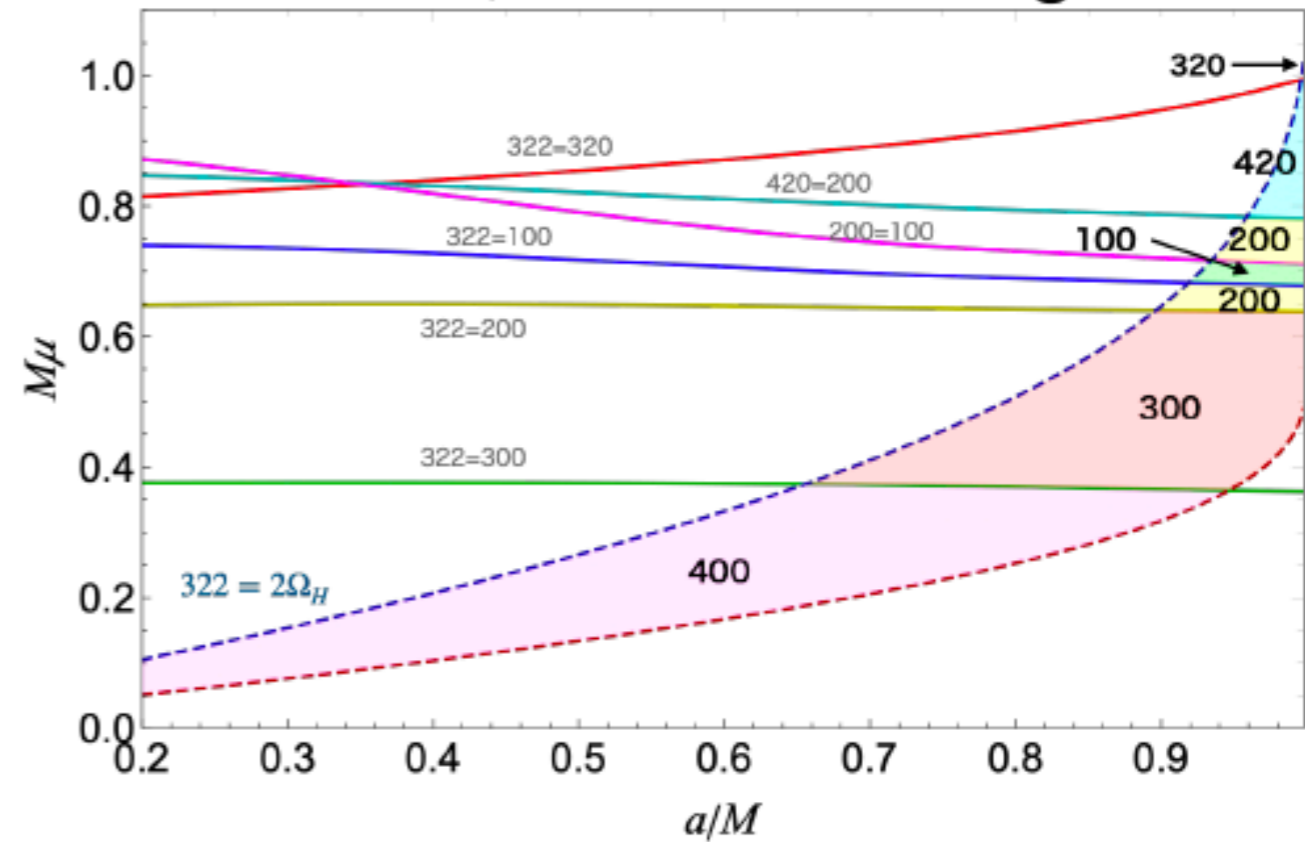
i.e. co-rot. $|211\rangle \rightarrow |21 - 1\rangle$, conter-rot. $|211\rangle \rightarrow |31 - 1\rangle$

$$l = 2$$

$l=2$, co-rotating



$l=2$, counter-rotating



Cloud depletion

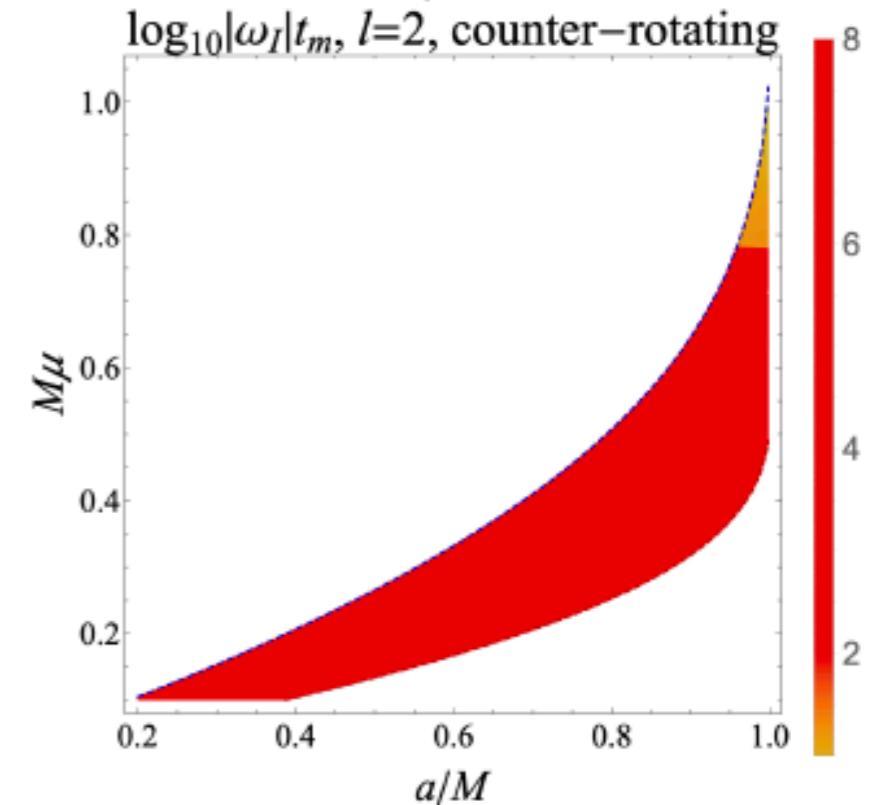
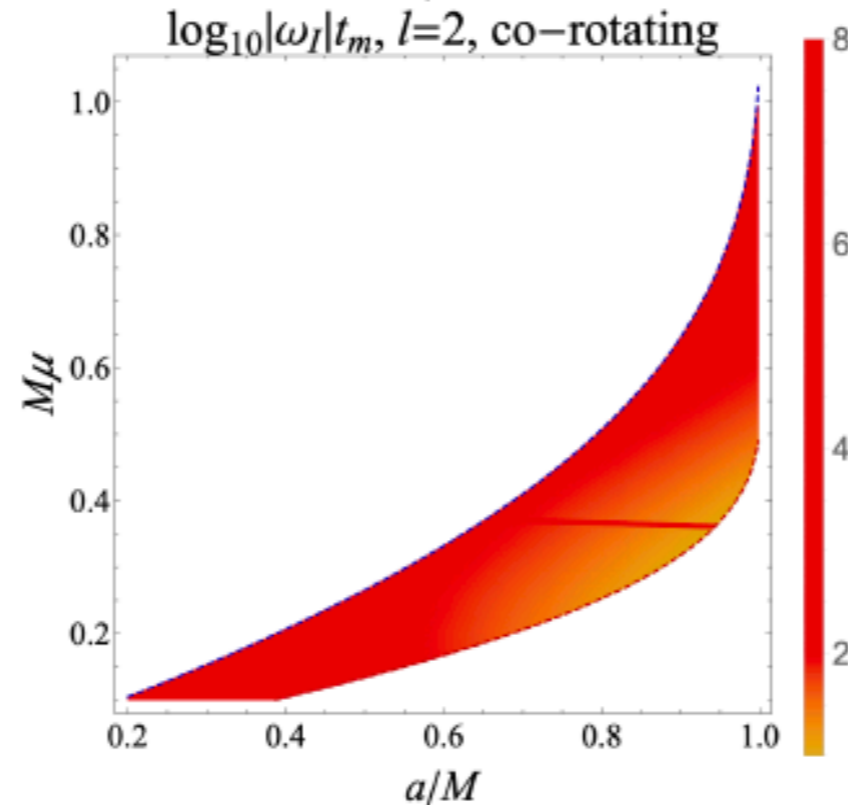
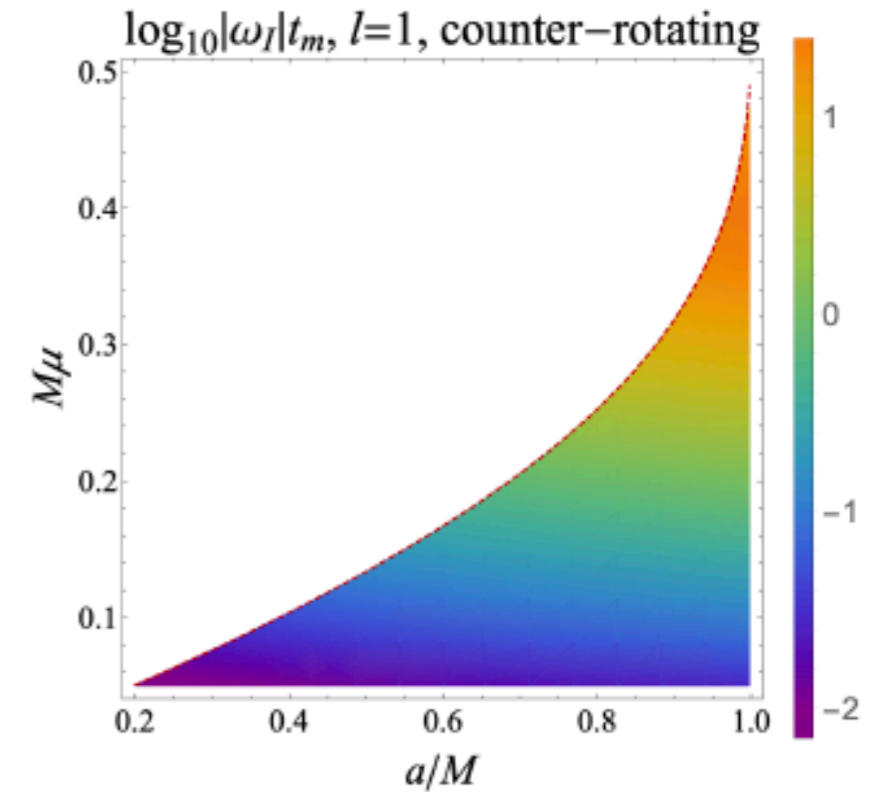
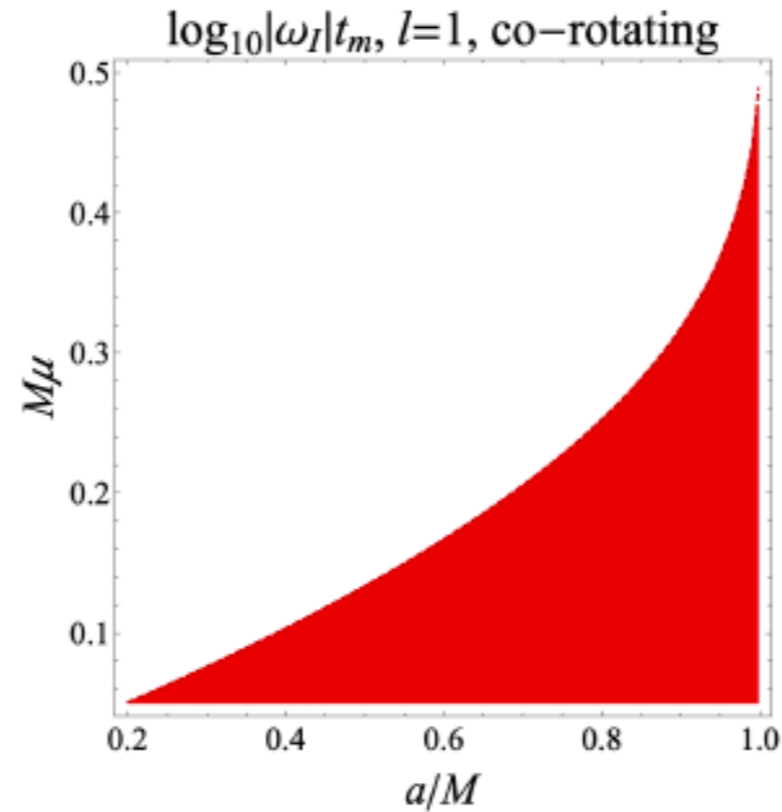
$$M_{cl} \simeq M_0 \exp(\underbrace{\omega_I}_{\text{Decay width}} \underbrace{t_m}_{\text{Time to merger}})$$

Decay width ↑
Time to merger

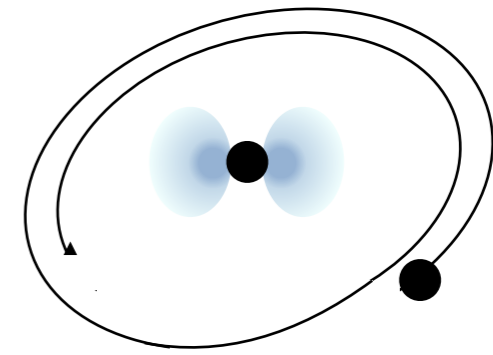
$$|\omega_I| t_m \gg 1$$

→ disappear

axions in $l \geq 2$ mode can be transferred to smaller l mode.



Summary



We studied the axion cloud depletion due to binary's tidal interaction in a wide parameter region including the cloud is in the relativistic regime.

- assuming • only leading quadrupolar tidal perturbation is at work
- transition is adiabatic

This is the first step!

Ongoing work

- Effect of higher multipole moment of tidal potential
- Study this system in more detail

Backup

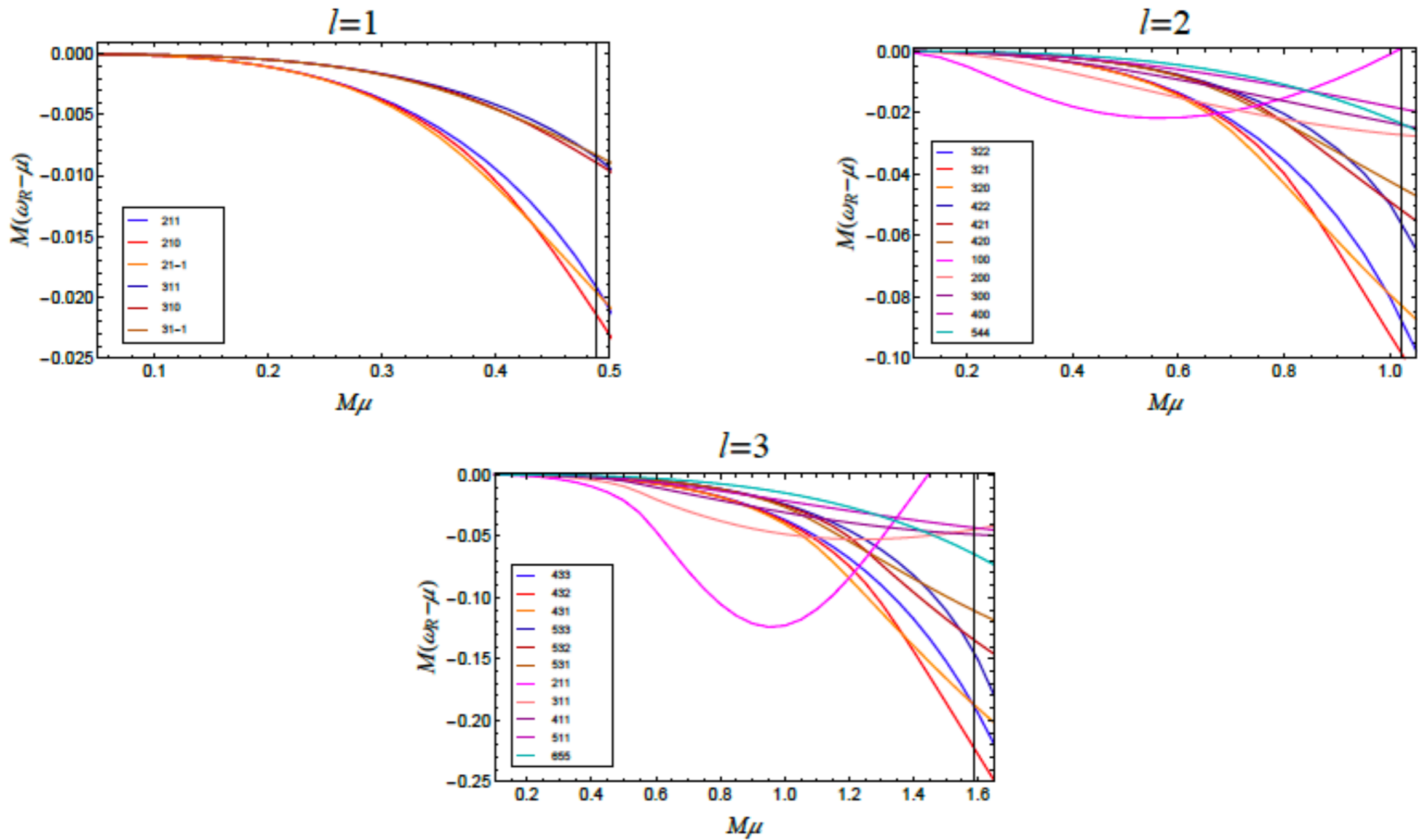


Figure 1. Energy spectra of bound states with $l = 1, 2$ and 3 relevant in discussing the transitions for $a/M=0.998$. Here, the vertical lines in the respective panels show the range of $M\mu$ beyond which $|nlm\rangle = |211\rangle, |322\rangle, |433\rangle$ are not the fastest growing modes of interest.