

## Abstract

We study the inflationary expansion in the early Universe within the framework of non linear electrodynamics(NLED). We propose a NLED lagrangian which is characterised by two parameters  $\alpha$ (dimensionless) and  $\beta$ (dimensionful). The classical stability and causality aspects of the model requires the speed of the sound wave  $c_s^2 > 0$  and  $0 \leq c_s \leq 1$  provided  $1.828 \leq \beta B^2 \leq 1.469$  with  $\alpha = 1$ . The magnetic field necessary to trigger the inflation is found to be  $B \simeq \sqrt{\frac{0.4\rho_B^{(max)}}{0.65}} = 4 \times 10^{51}$  Gauss, where  $\rho_B^{(max)} (\sim \rho_{inf}) = 10^{64}$  GeV<sup>4</sup>. The model also predicts the desired e-fold number  $N = 71$  in the parameter space of  $\alpha$  and  $\beta$ . With  $\alpha = 0.3$  and  $\beta B^2 = 0.3974$ , we find the scalar spectral index  $n_s = 0.9649$ , which is consistent with the PLANCK 2018 data and also predict the tensor-to-scalar ratio  $r = 0.1417$  and the tensor spectral index  $n_T = -0.0177$  corresponding to  $n_s = 0.9649$ .

## Introduction

The cosmic evolution of the early Universe was largely controlled by the non-linear electromagnetic(NLED) field. This NLED field, which acts as a source of gravity, can lead to negative pressure and thus drive the inflationary expansion of the Universe [1].

## A model of non-linear electrodynamics

The NLED Lagrangian, we propose, is given by

$$\mathcal{L} = -\frac{\mathcal{F}e^{-\beta\mathcal{F}}}{(\mathcal{F}\beta + \alpha)^2}$$

where  $\mathcal{F} = \frac{1}{4}F_{\mu\nu}F^{\mu\nu} = \frac{1}{2}(B^2 - E^2)$ . One finds the energy-momentum tensor of the NLED field,

$$T_{\mu\nu} = \frac{\partial\mathcal{L}}{\partial\mathcal{F}}F_{\mu\alpha}F_{\nu}^{\alpha} - g_{\mu\nu}\mathcal{L}$$

The early Universe being magnetic(where we set  $E = 0$  as the electric field is screened off due to the charged plasma), we find the energy density( $\rho(= \rho_B)$ ) and pressure( $P(= P_B)$ )

$$\rho_B = \frac{2B^2e^{-\beta B^2/2}}{(\beta B^2 + 2\alpha)^2}, P_B = -\frac{4B^2e^{-\beta B^2/2}}{3(\beta B^2 + 2\alpha)^3} \left[ \beta^2 B^4 + \beta B^2 \left( \frac{7}{2} + 2\alpha \right) - \alpha \right]$$

## Inflation in non-linear electrodynamics

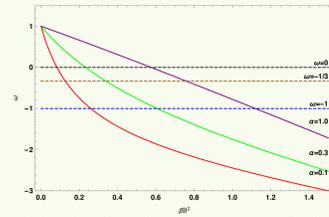
The Friedmann and the Roychaudhury equation are given by,

$$\left( \frac{\dot{a}}{a} \right)^2 = \frac{1}{3}\rho_B, \quad \frac{\ddot{a}}{a} = -\frac{1}{6}(\rho + 3P_B)$$

where we set  $8\pi G = 1$ .

- The EoS parameter  $\omega$  of the NLED field with  $B = B_0/a^2$ , is given by

$$\omega = \frac{P_B}{\rho_B} = -\frac{2}{3} \left( \frac{\beta B_0^2}{a^4} + 2\alpha \right)^{-1} \left[ \frac{\beta^2 B_0^4}{a^8} + \frac{\beta B_0^2}{a^4} \left( \frac{7}{2} + 2\alpha \right) - \alpha \right]$$



- $\omega$  is plotted against  $\beta B^2$  for different  $\alpha$ .
- $\omega$  crosses  $-1$ (negative pressure  $\rightarrow$  acceleration) at  $\beta B^2 = 0.25(1.112)$  corresponding to  $\alpha = 0.1(1.0)$ .

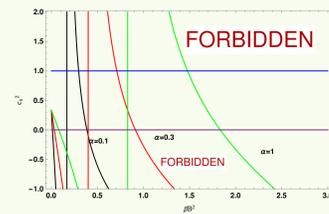
- The deceleration parameter  $q$  is defined as

$$q = -\frac{a\ddot{a}}{\dot{a}^2} = \frac{1}{2}(1+3\omega) = \frac{2\alpha - 3\beta B^2 - 2\alpha\beta B^2 - \beta^2 B^4}{2\alpha + \beta B^2} \rightarrow 1 \text{ as } B \rightarrow 0$$

So,  $q \rightarrow 1 \rightarrow \ddot{a} < 0 \implies$  *No acceleration without the magnetic field!!*

- Causality and Classical stability of  $c_s$ : The speed of the sound wave  $c_s$  is given by

$$c_s^2 = \frac{dP_B}{d\rho_B} = -1 - \frac{2}{3}\beta B^2 + \frac{8\alpha}{2\alpha + \beta B^2} + \frac{16\alpha - 2\beta B^2(2\alpha + 2)}{6\alpha \left( \frac{\beta B^2}{2} - 1 \right) + 3\beta B^2 \left( \frac{\beta B^2}{2} + 1 \right)}$$



- $c_s^2$  is plotted against  $\beta B^2$  for different  $\alpha$ .
- For  $\alpha = 0.1(1.0)$ , we see that the classically stability condition  $0 \leq c_s^2 \leq 1$  requires  $0.382(1.828) \leq \beta B^2 \leq 0.288(1.469)$ .
- The regions corresponding to  $c_s^2 < 0$  and  $c_s^2 > 1$  are forbidden.

## Effective potential $V_{eff}$ , the inflation energy density $\rho_B$ and the e-fold number( $N$ )

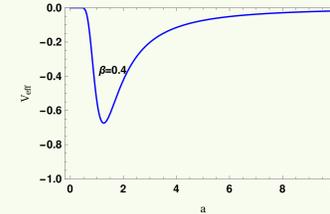
- The Friedmann equation reads

$$H^2 = \left( \frac{\dot{a}}{a} \right)^2 = \frac{\rho_B}{3} \rightarrow \dot{a}^2 + V_{eff}(a) = 0$$

where the effective potential  $V_{eff}(a) = -\frac{B_0^2}{6a^2}e^{-\beta B_0^2/2a^4} \left( \frac{\beta B_0^2}{2a^4} + \alpha \right)^{-2}$ .

- Taking a derivative of the Friedman eqn, we find

$$2\dot{a}\ddot{a} + \dot{a}\frac{dV_{eff}}{da} = 0 \rightarrow \ddot{a} + \frac{1}{2}\frac{dV_{eff}}{da} = 0$$



- $V_{eff}$  is plotted against  $a$  for  $\alpha = 0.3, \beta = 0.4$  and  $B_0 = 1$  and it has a minima at  $a_c = 1.69$  (the point where the inflation stops).
- On the left of  $a_c$ , we see  $\frac{dV_{eff}}{da} < 0$  i.e.  $\ddot{a} > 0 \rightarrow$  *acceleration*, whereas, on the right of  $a_c$ ,  $\frac{dV_{eff}}{da} > 0$  i.e.  $\ddot{a} < 0 \rightarrow$  *deceleration*.

- Inflation requires  $\dot{\rho}_B = 0 \rightarrow \rho_B + P_B = 0$  and one finds

$$\beta^2 B^4 + \beta B^2(2 + 2\alpha) - 4\alpha = 0$$

This gives  $\beta B^2 = 0.4$  for  $\alpha = 0.3$ .

- The energy density  $\rho_B(= \rho_B^{(max)})$  during inflation is calculated as

$$\rho_B = \frac{2\beta B^2 e^{-\beta B^2/2}}{\beta(2\alpha + \beta B^2)^2} = 0.65/\beta = \rho_B^{max} \rightarrow \beta \rho_B^{max} = 0.65$$

- Now  $\beta B^2 = 0.4 \rightarrow B \simeq \sqrt{\frac{0.4}{\beta}} = \sqrt{\frac{0.4\rho_B^{max}}{0.65}} = 4 \times 10^{51}$  G with  $\rho_B^{max} (\rho_{inf}) = 10^{64}$  GeV<sup>4</sup>.

- The e-fold number( $N$ ) is defined as

$$N = \int_t^{t_{end}} H dt = \ln \frac{a_{end}}{a} = \ln \sqrt{\frac{B_{start}}{B_{end}}}$$

With  $B_{end} = 10^{-10}$  G and  $B_{start} = 4 \times 10^{51}$  G, we find

$$N = \frac{1}{2} \ln(4 \times 10^{61}) = 71$$

## Spectral index analysis

The scalar spectral index( $n_s$ ), tensor-to-scalar ratio ( $r$ ) and the tensor spectral index ( $n_T$ ) in NLED are defined as

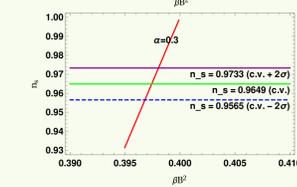
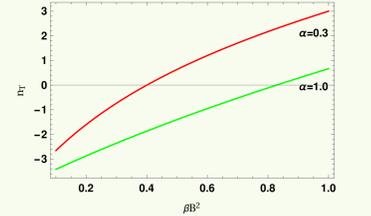
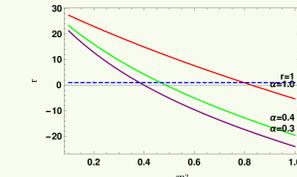
$$n_s = 1 - \frac{4(4\alpha - 2\beta B^2 - \beta B^2(\beta B^2 + 2\alpha))}{\beta B^2 + 2\alpha},$$

$$r = \frac{16(4\alpha - 2\beta B^2 - \beta B^2(\beta B^2 + 2\alpha))}{(\beta B^2 + 2\alpha)}$$

and

$$n_T = -\frac{r}{8} = -2 \left[ \frac{4\alpha - 2\beta B^2 - \beta B^2(\beta B^2 + 2\alpha)}{(\beta B^2 + 2\alpha)} \right]$$

Results:



- $n_s, r$  and  $n_T$  are plotted as a function of  $\beta B^2$  for different  $\alpha$  values.

- Also shown are the scalar spectral index  $n_s = 0.9649 \pm 0.0042$  (Planck 2018 CMB data)[2].

- We find that  $\alpha = 0.3$  and  $\beta B^2 = 0.3974$  are consistent with the scalar spectral index data  $n_s = 0.9649 \pm 0.0042$ (Planck 2018).

- With these  $\alpha, \beta$ , we predict the tensor-to-scalar ratio  $r = 0.1417, 0.1089, 0.1744$  and the tensor spectral index  $n_T = -0.0177, -0.0136, -0.0218$  corresponding to c.v. and c.v. $\pm 2\sigma$  values of  $n_s$ .

## Results & Conclusions

- We find that a highly non-linear magnetic field  $B = 4 \times 10^{51}$  G can trigger the inflationary expansion of the early Universe.

- Our NLED model predicts the e-fold number  $N = 71$  in the space of nonlinear parameters  $\alpha$  and  $\beta$ .

- We find that the nonlinear parameters  $\alpha = 0.3$  and  $\beta B^2 = 0.3974$  are consistent with the scalar spectral index  $n_s = 0.9649 \pm 0.0042$ (Planck 2018) and we also predict the tensor-to-scalar ratio  $r = 0.1417$  and tensor spectral index  $n_T = -0.0177$ .

## Acknowledgements

PS thanks to Department of Science and Technology(DST), Government of India for the Inspire fellowship (No. DST/INSPIRE Fellowship/2017/IF170807). The work of PKD and GCS is supported by CSIR Grant No.25(0260)/17/EMR-II .

## References

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