



**Part 1: Cosmic Birefringence from  
Domain Wall without a string**

**Part 2: Dark Matter stability from Pauli-  
exclusion principle**

**Wen Yin (Tohoku University)**

@29 March 2022

“What is dark matter? - Comprehensive study of the huge discovery space in dark matter”

# Part 1

## **Cosmic Birefringence from Domain Wall without a string**

Based on Takahashi, WY, 2012.11576

and ongoing project in collaboration with Takahashi, Kitajima and Kozai

# What does the measured parity-violating cosmic birefringence (CB) suggest?

★ An ALP can alter the Maxwell equation

Carroll, Field, Jackiw, 1990; Harari, Sikivie, 1992; Carroll, 1998;

$$\mathcal{L} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} - c_\gamma \frac{\alpha}{4\pi} \frac{\phi}{f_\phi} F_{\mu\nu} \tilde{F}^{\mu\nu}$$

$$\Phi(\Omega) = 0.42 \text{ deg} \times c_\gamma \left( \frac{\phi_{\text{Earth}} - \phi_{\text{LSS}}(\Omega)}{2\pi f_\phi} \right),$$

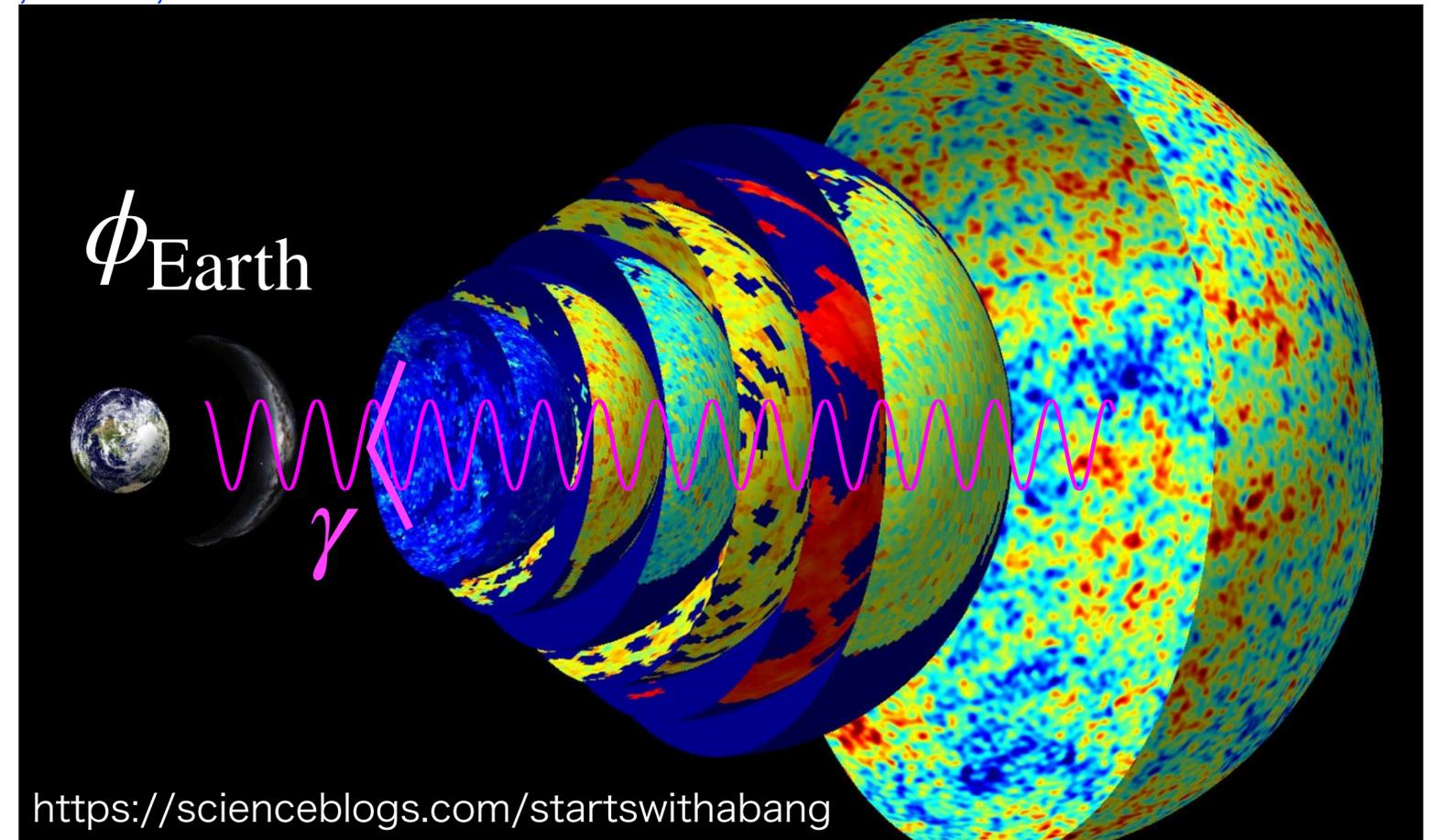
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Diego-Palazuelos et al, 2201.07682

see also Minami and Komatsu, 2006.15982, and B06's session.

$$\longrightarrow c_\gamma \left( \frac{\phi_{\text{Earth}} - \bar{\phi}_{\text{LSS}}}{f_\phi} \right) \sim (0.9\pi - 1.9\pi)$$



Averaged  $\phi$  on last scattering surface

$$\bar{\phi}_{\text{LSS}}$$

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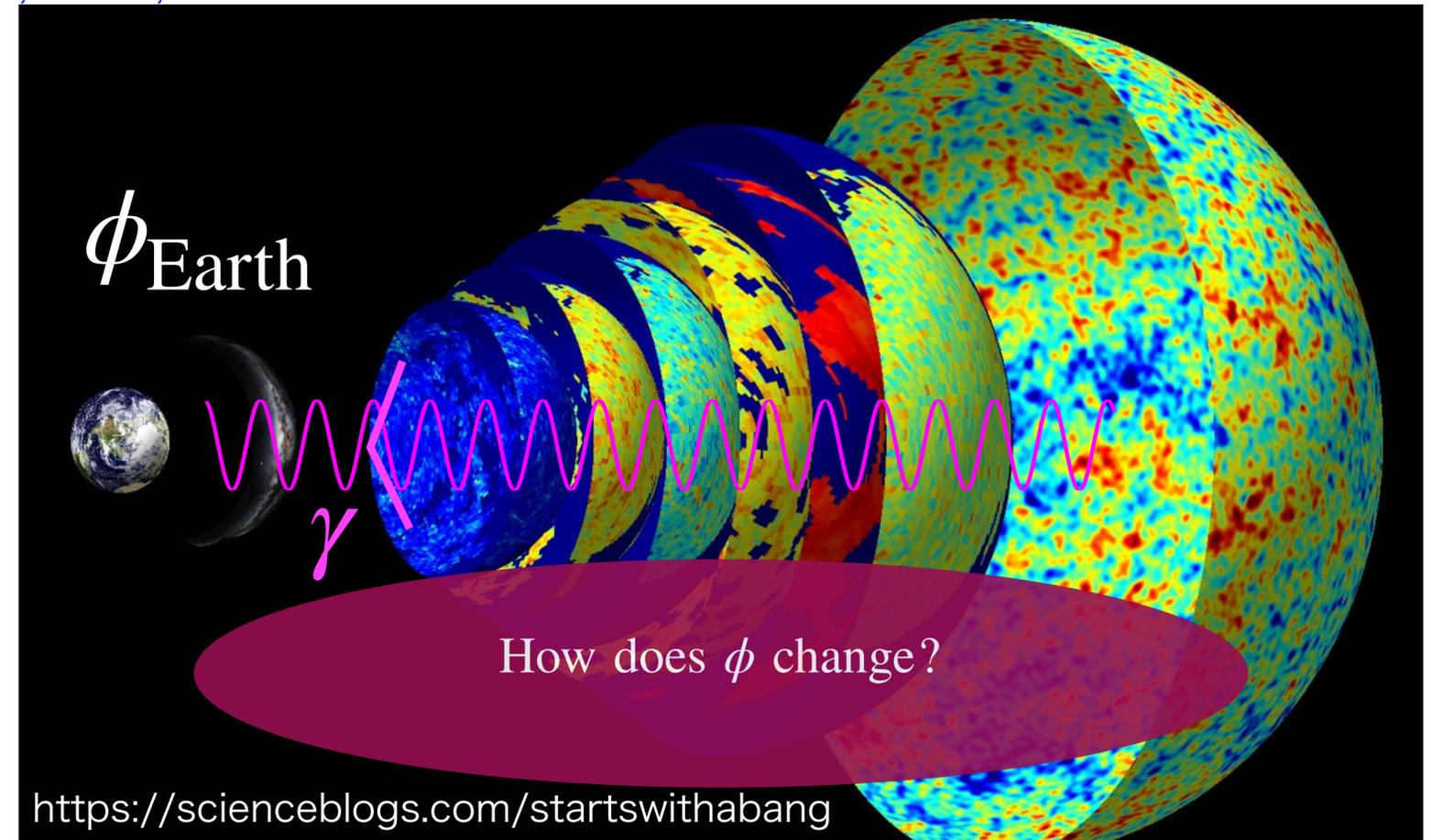
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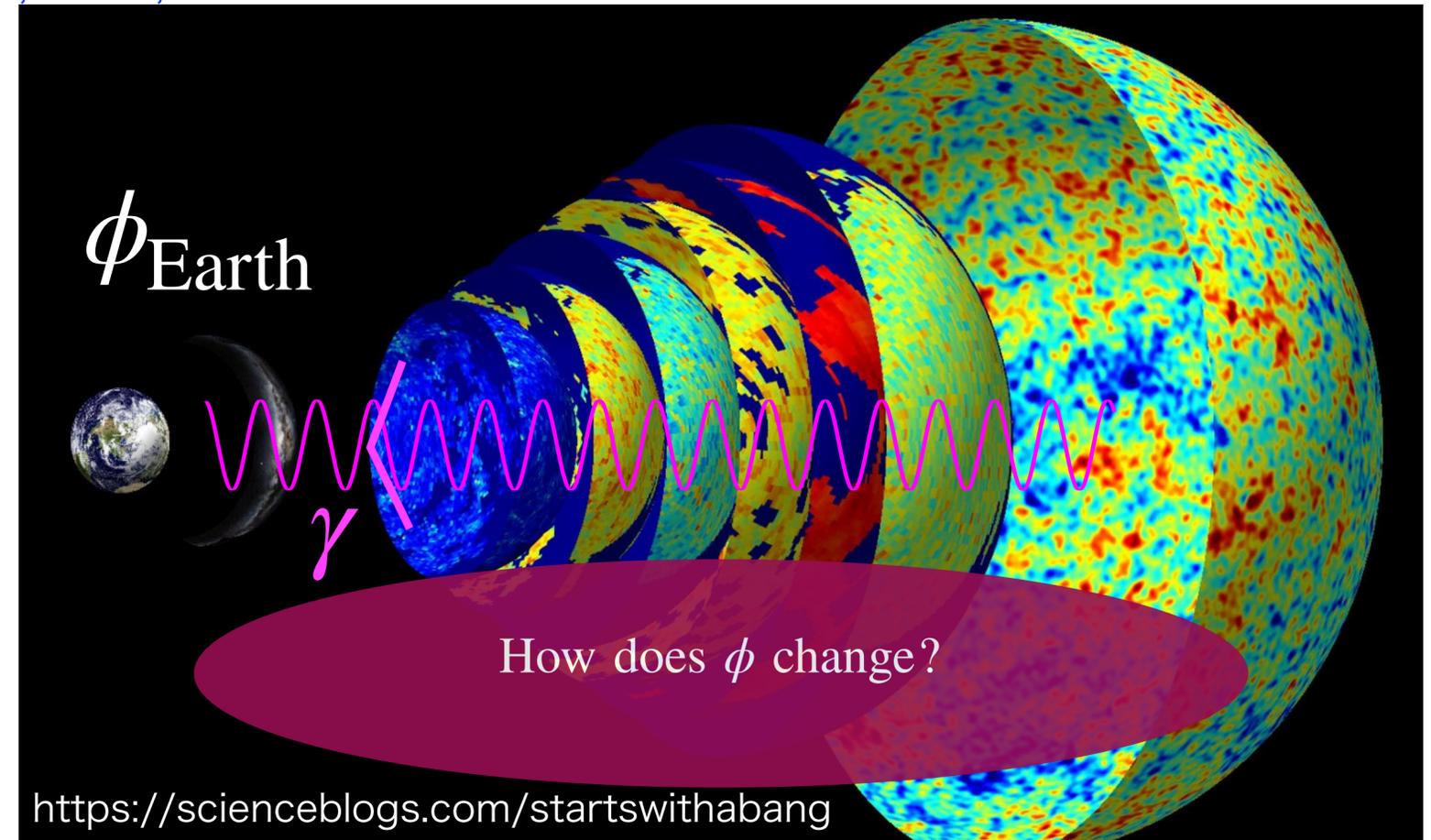
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Minami and Komatsu, 2006.15982, Fujita et al, 2011.11894, (CB and  $H_0$  tension)

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Takahashi, WY, 2012.11576 **my focus**

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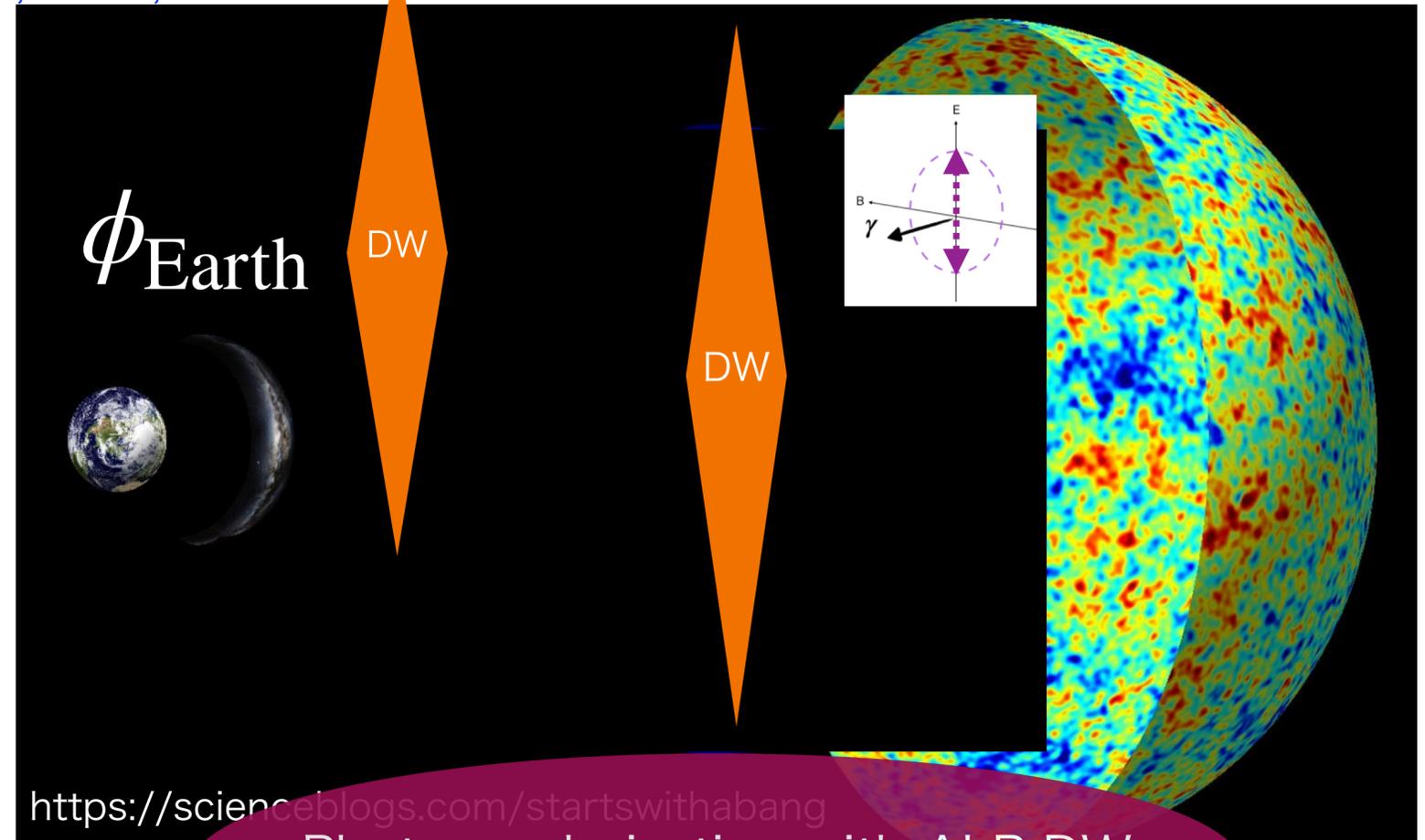
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Photon polarization with ALP DW

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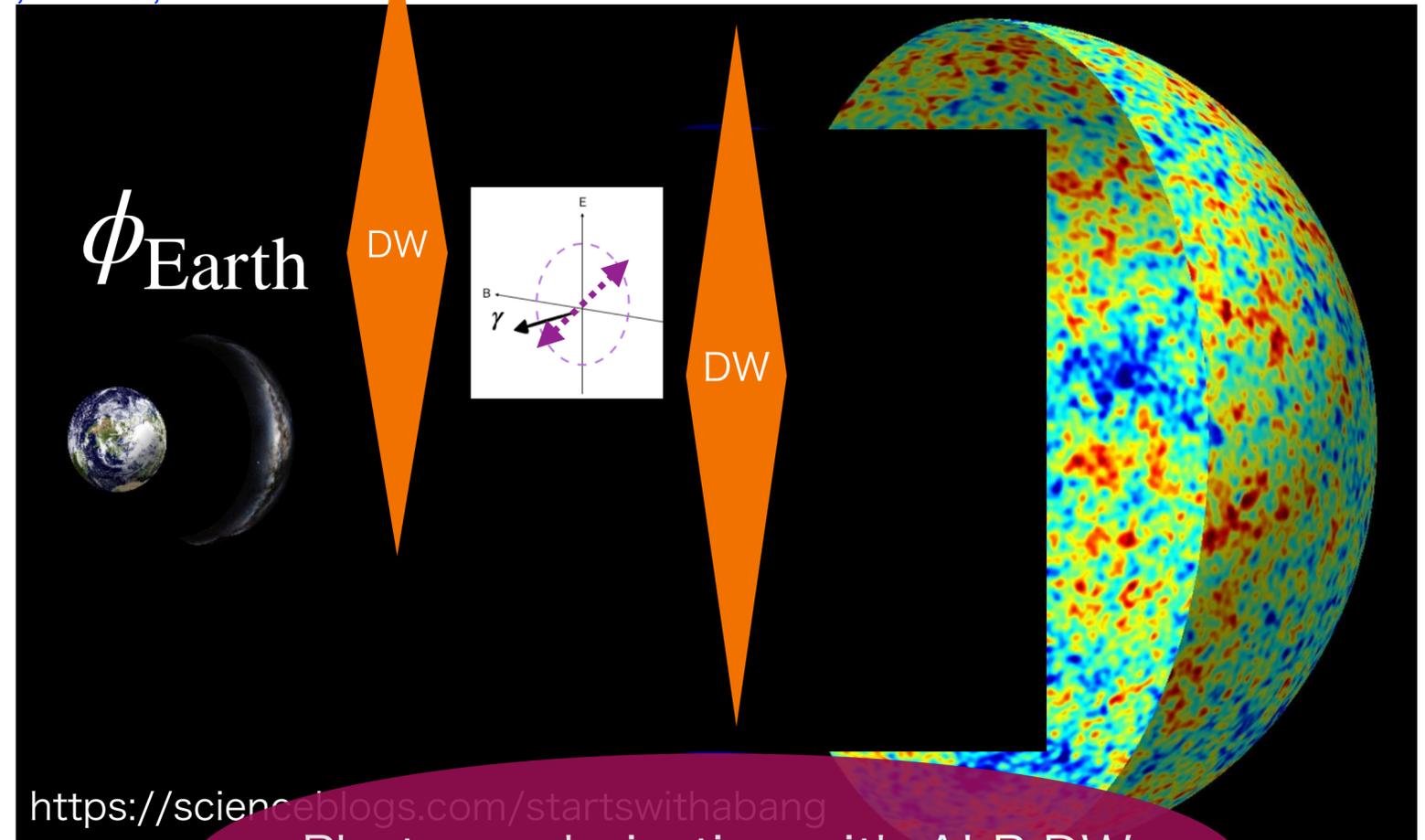
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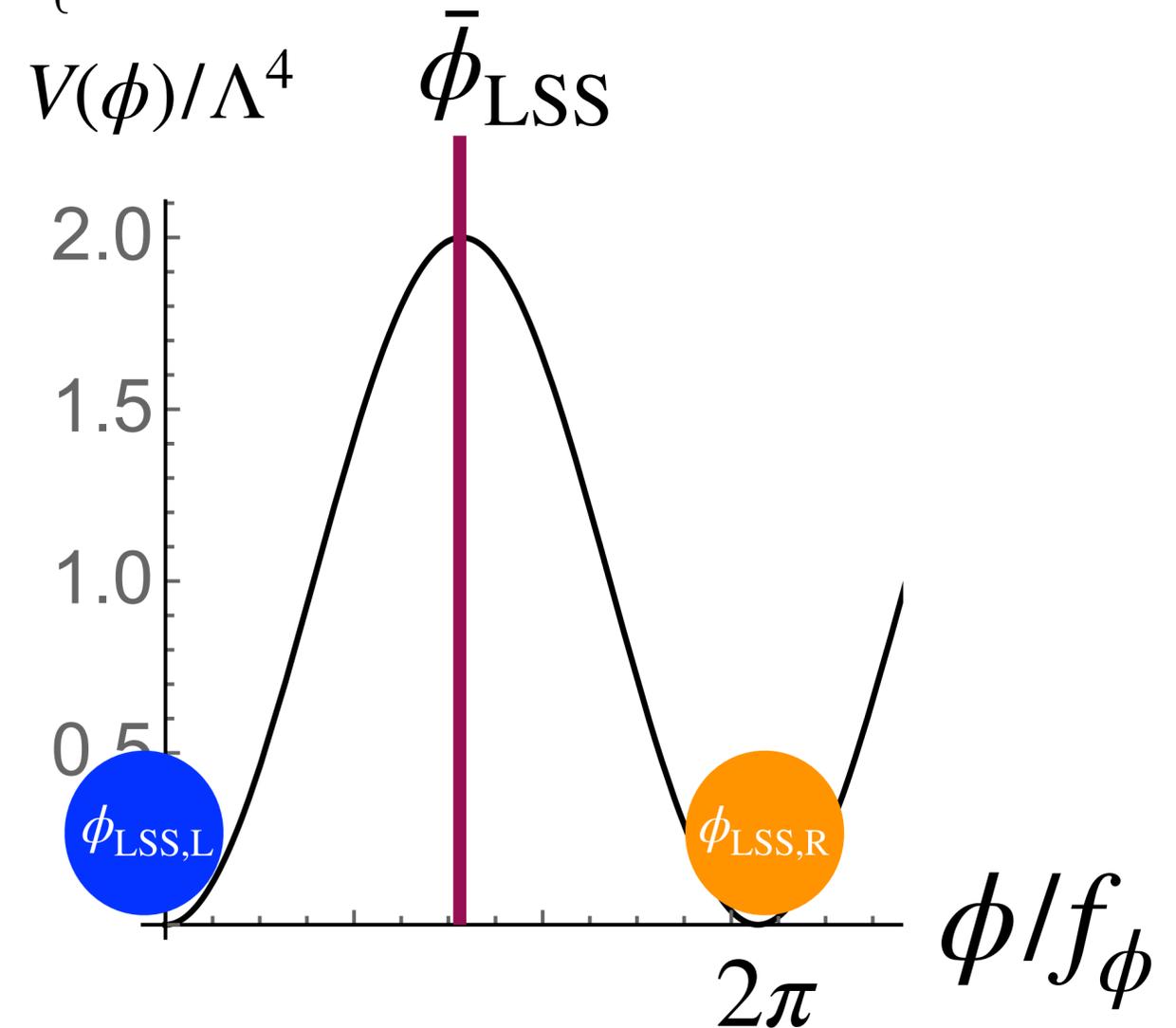
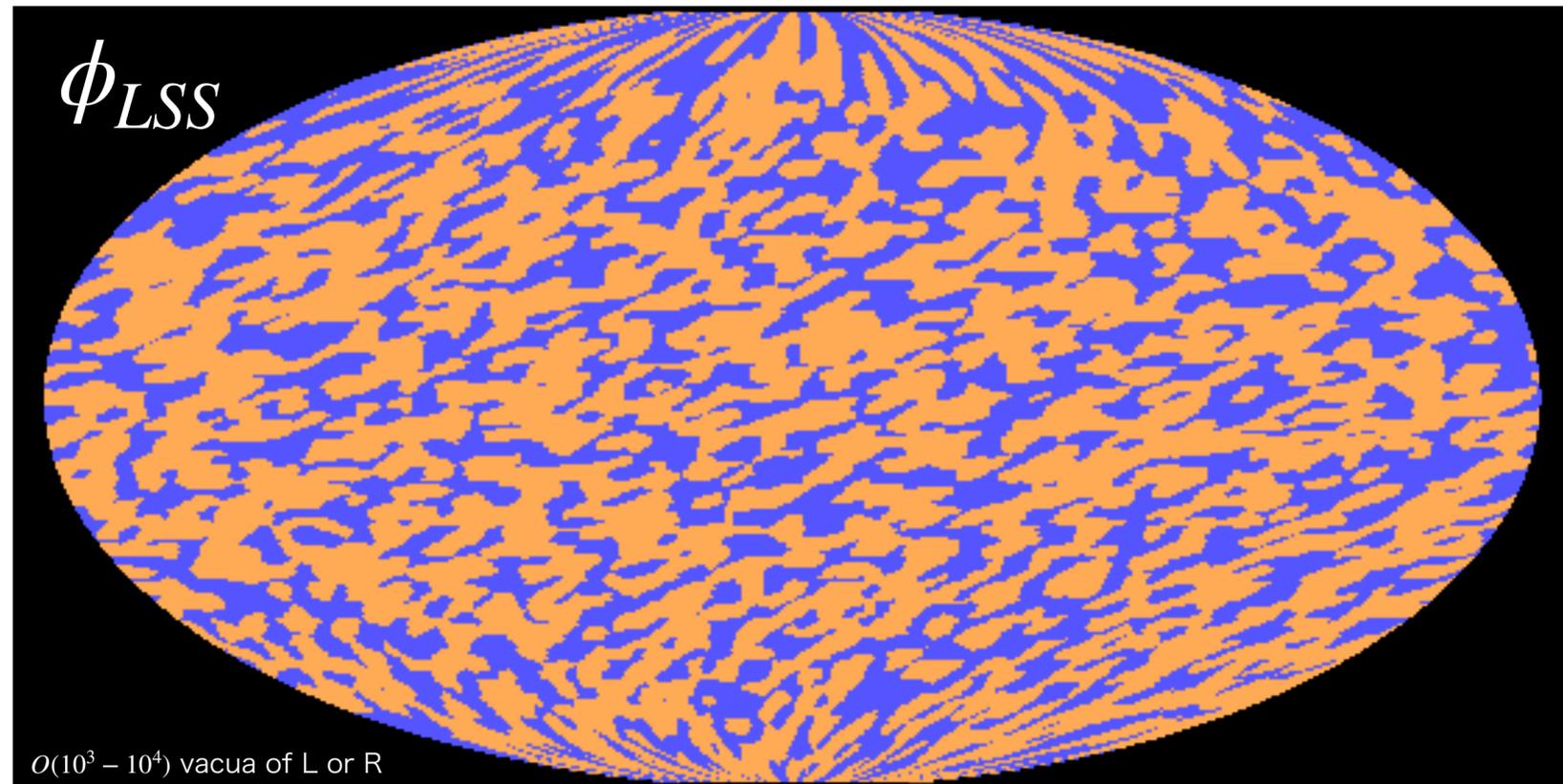
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ALP DW without a string following scaling solution ( $\equiv$  attractor solution with  $O(1)$  domains within a Hubble horizon) explains the isotropic CB.

$$|\beta| = 0.21 \text{deg} |c_\gamma|$$

c.f.  $\begin{cases} \beta = 0.42 \text{deg} \frac{(\phi_{\text{Earth}} - \bar{\phi}_{\text{LSS}})}{2\pi f_\phi} c_\gamma \\ \beta_{\text{obs}} = 0.30 \pm 0.11 \text{ deg}, \end{cases}$  [Takahashi, WY, 2012.11576](#)

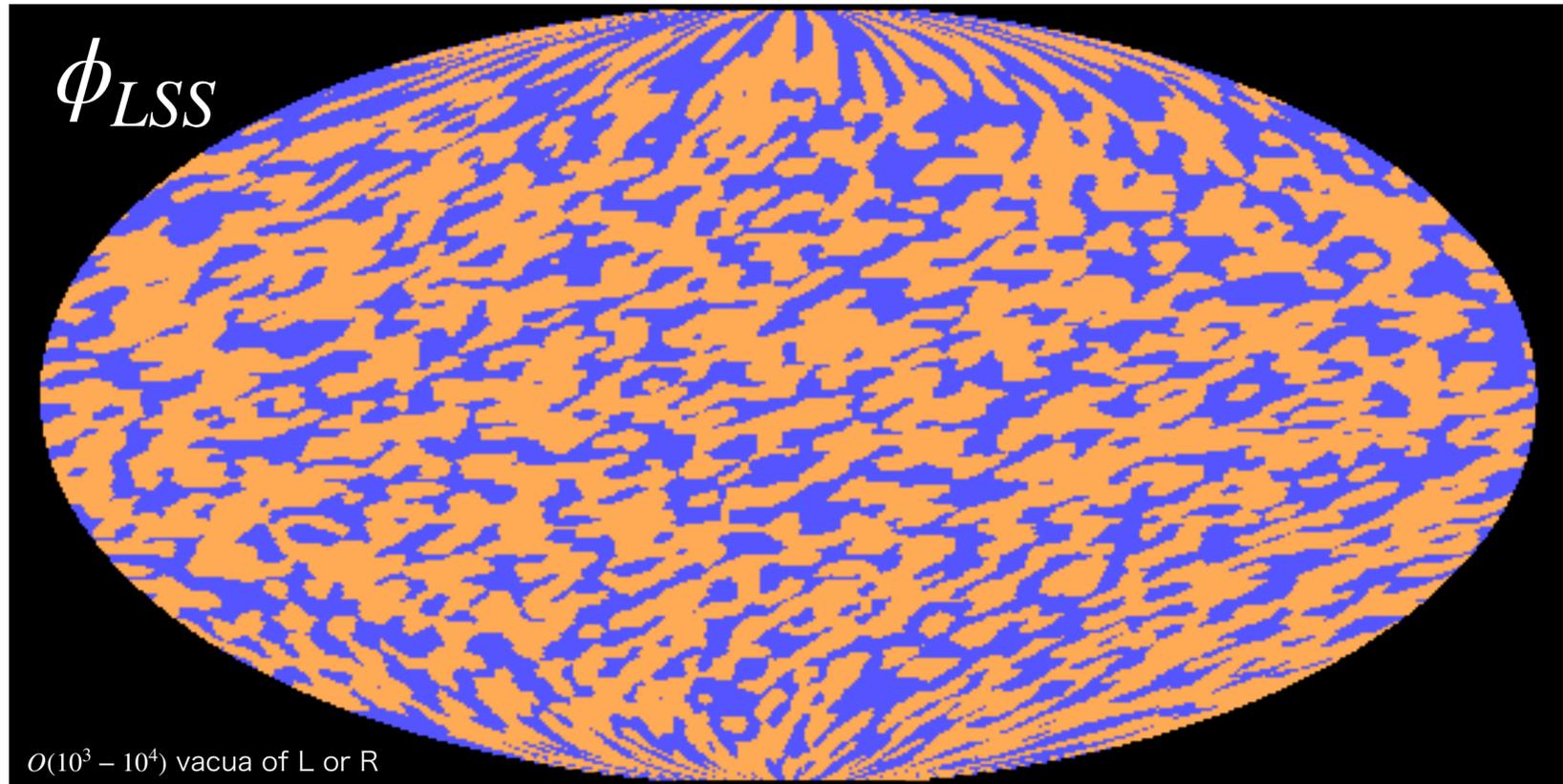
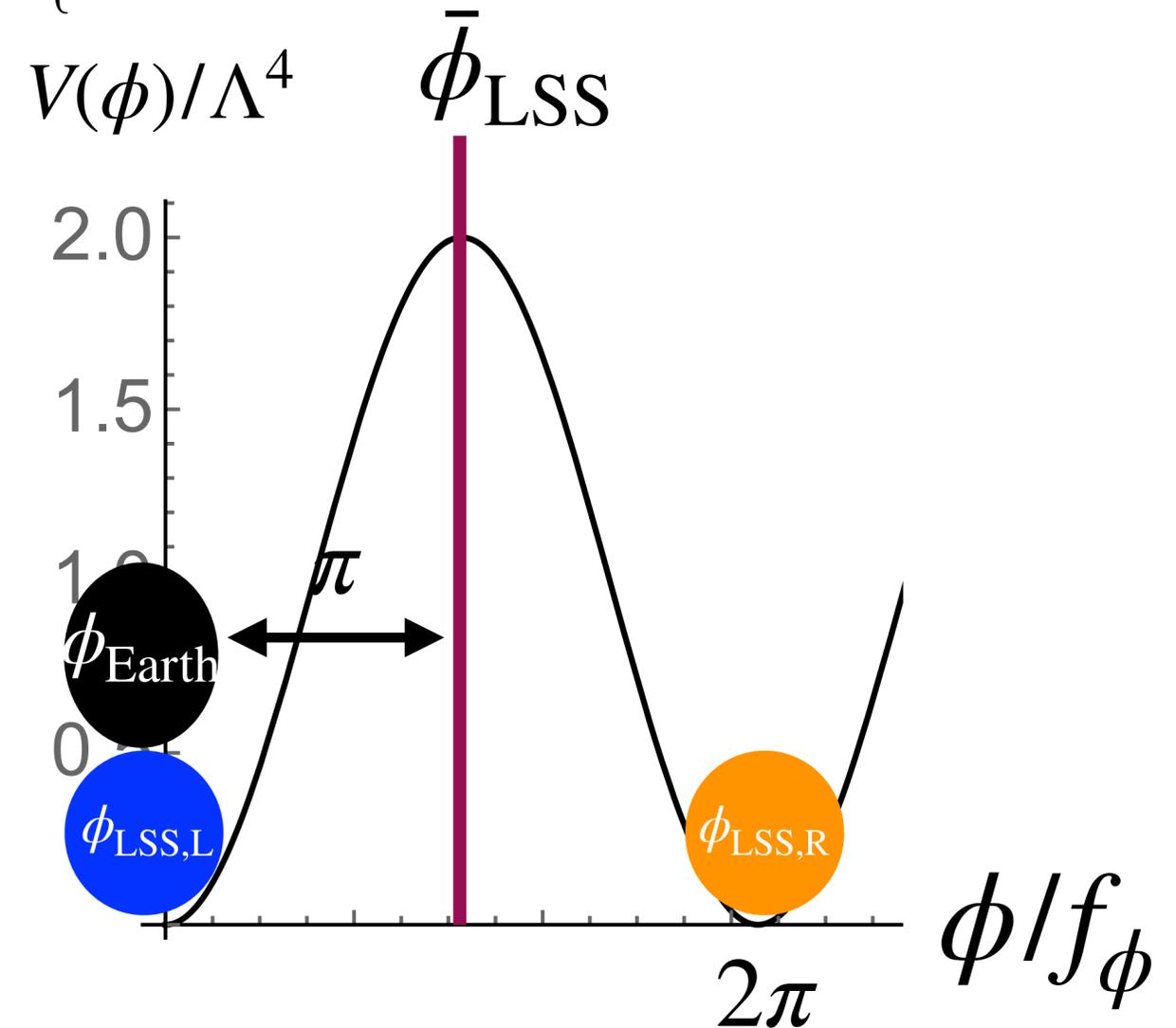


Strings with DWs cannot induce isotropic CB [Agrawal et al, 1912.02823;](#)

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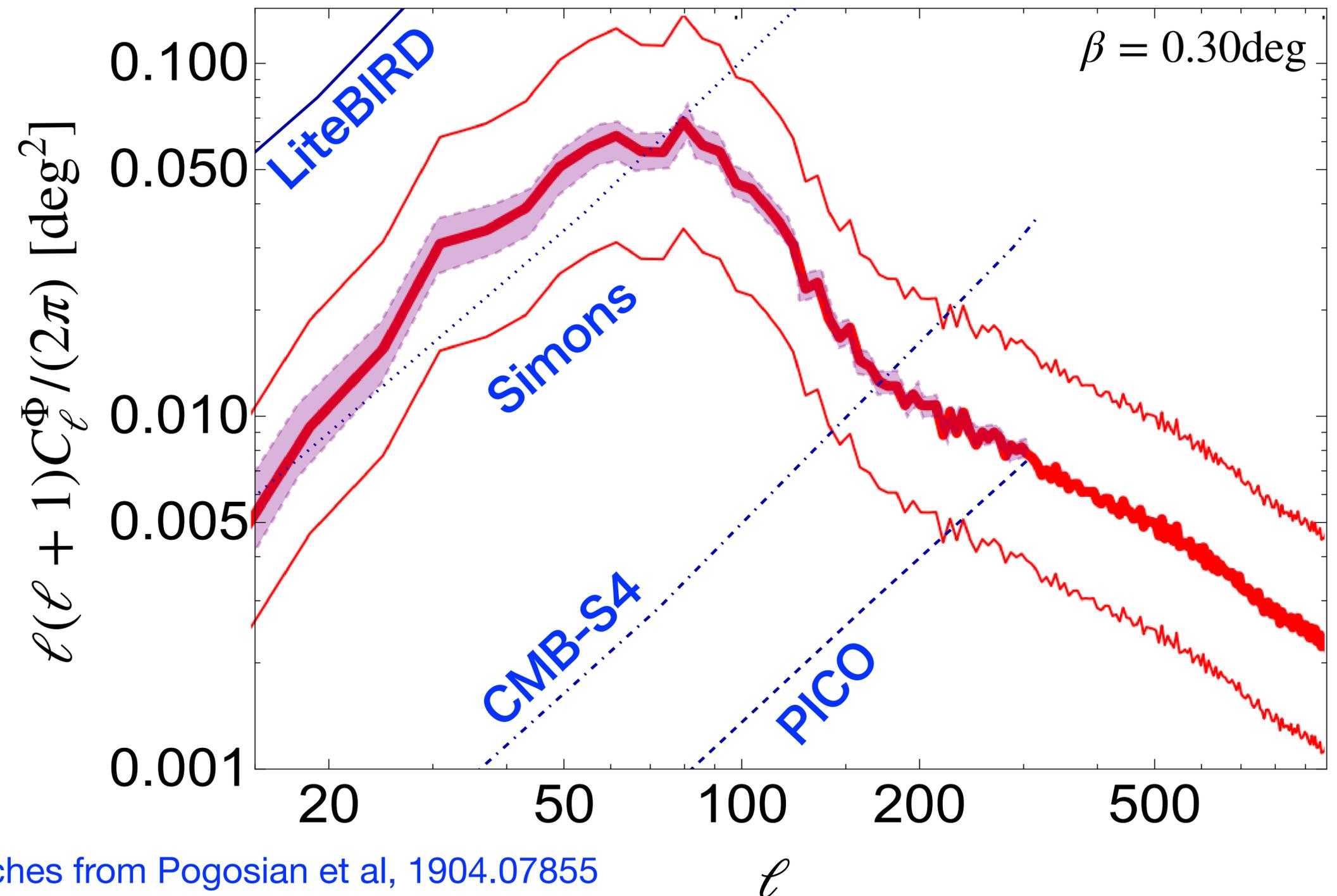
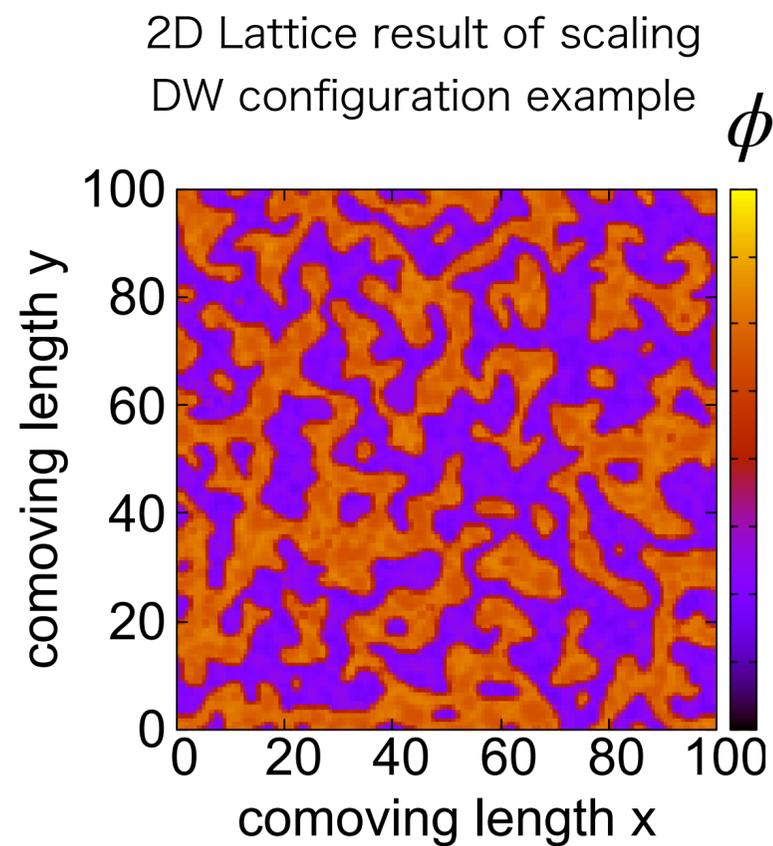


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# Predicted anisotropic birefringence can be tested.

We use 4096\*4096 Lattice simulation to estimate the power spectrum of anisotropic cosmic-birefringence.

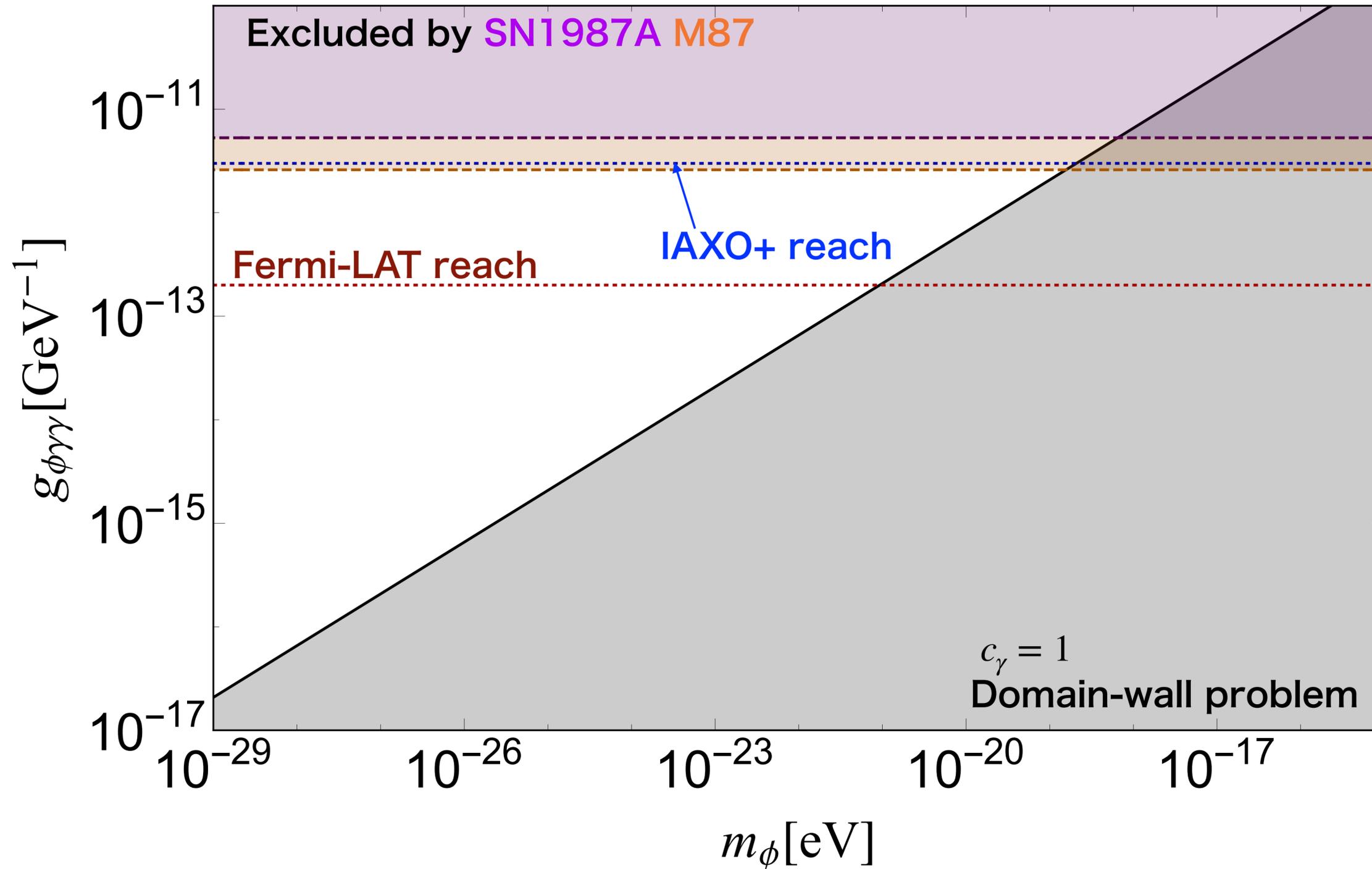
Takahashi, Kitajima, Kozai, WY, preliminary



Future reaches from Pogosian et al, 1904.07855

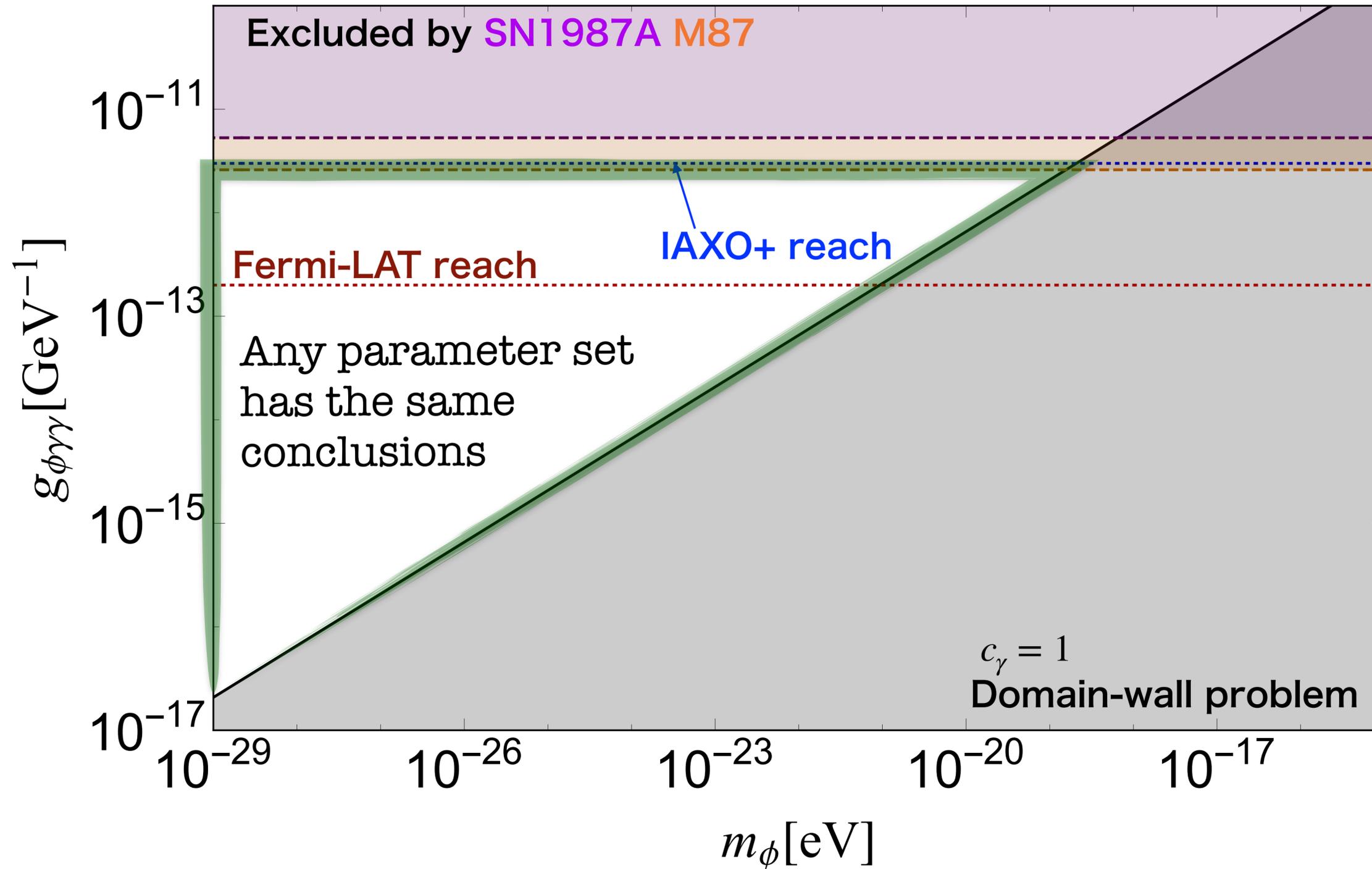
# Parameter region and other testability

Takahashi, WY, 2012.11576



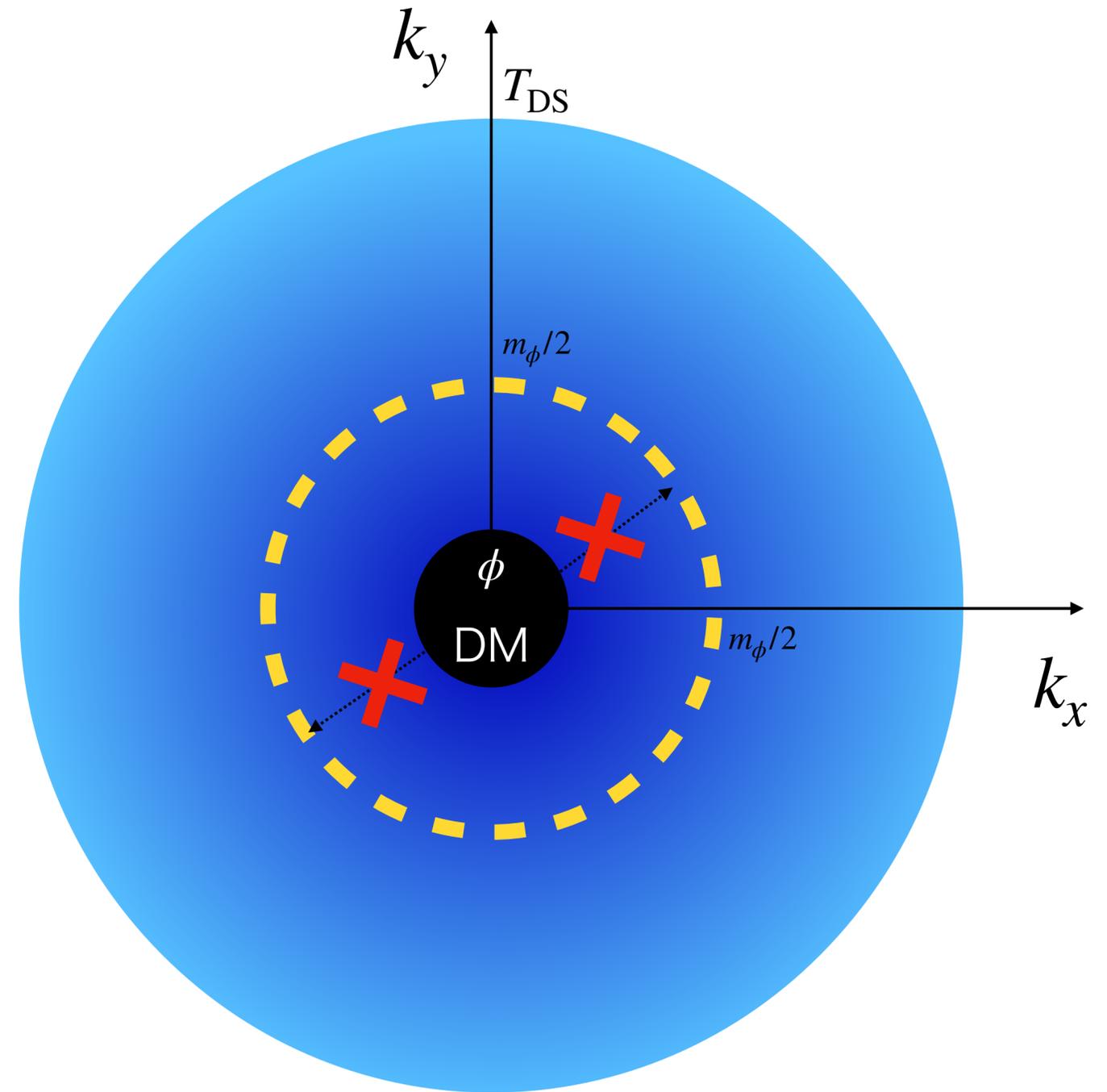
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# Part 2

## DM stability from Pauli-exclusion principle



# Dark matter

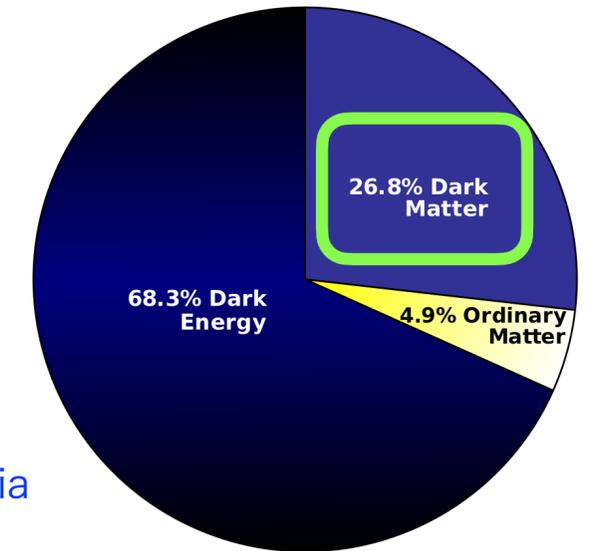
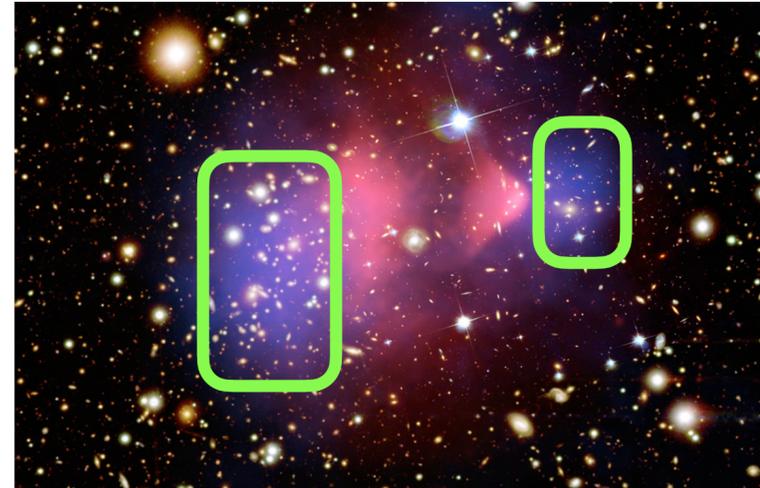
- What is dark matter?

Very stable

Neutral

Cold

$$\rho_{\text{DM}} \sim \text{keVcm}^{-3}$$



wikipedia

- Why is dark matter very stable?

“Charge” conservation, e.g. WIMP

Small mass/coupling,

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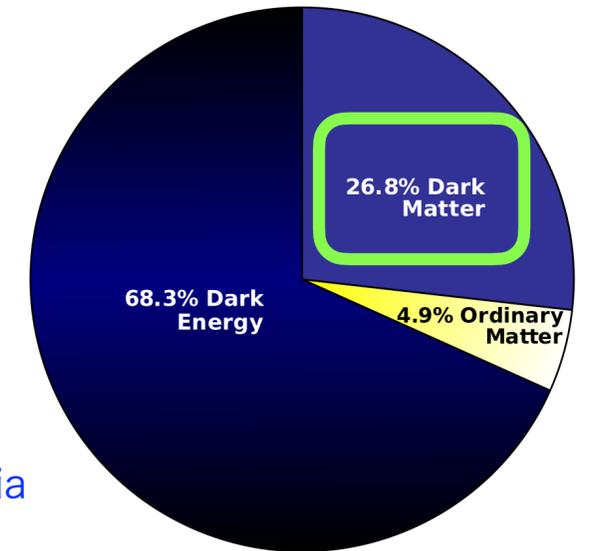
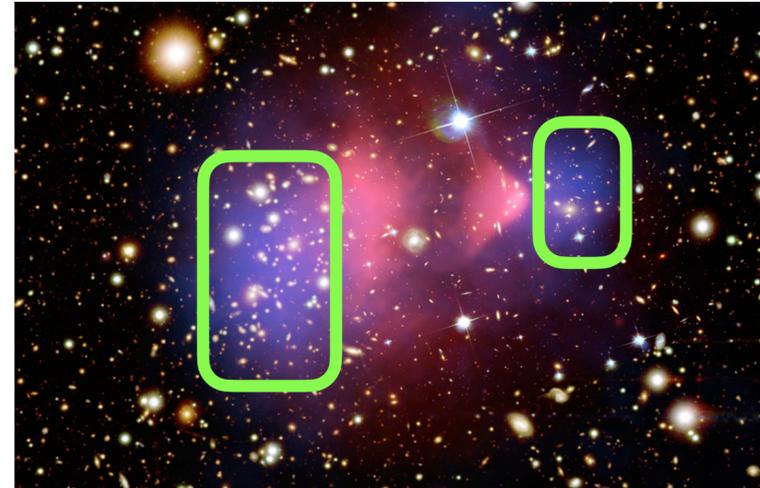
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## DM on Sea of Fermions

Batell, WY, in progress

$$1/\Gamma_{\text{DM}}^{\text{decay}} \ll 13.8\text{Gyr}$$

# What I will be talking about

## setup

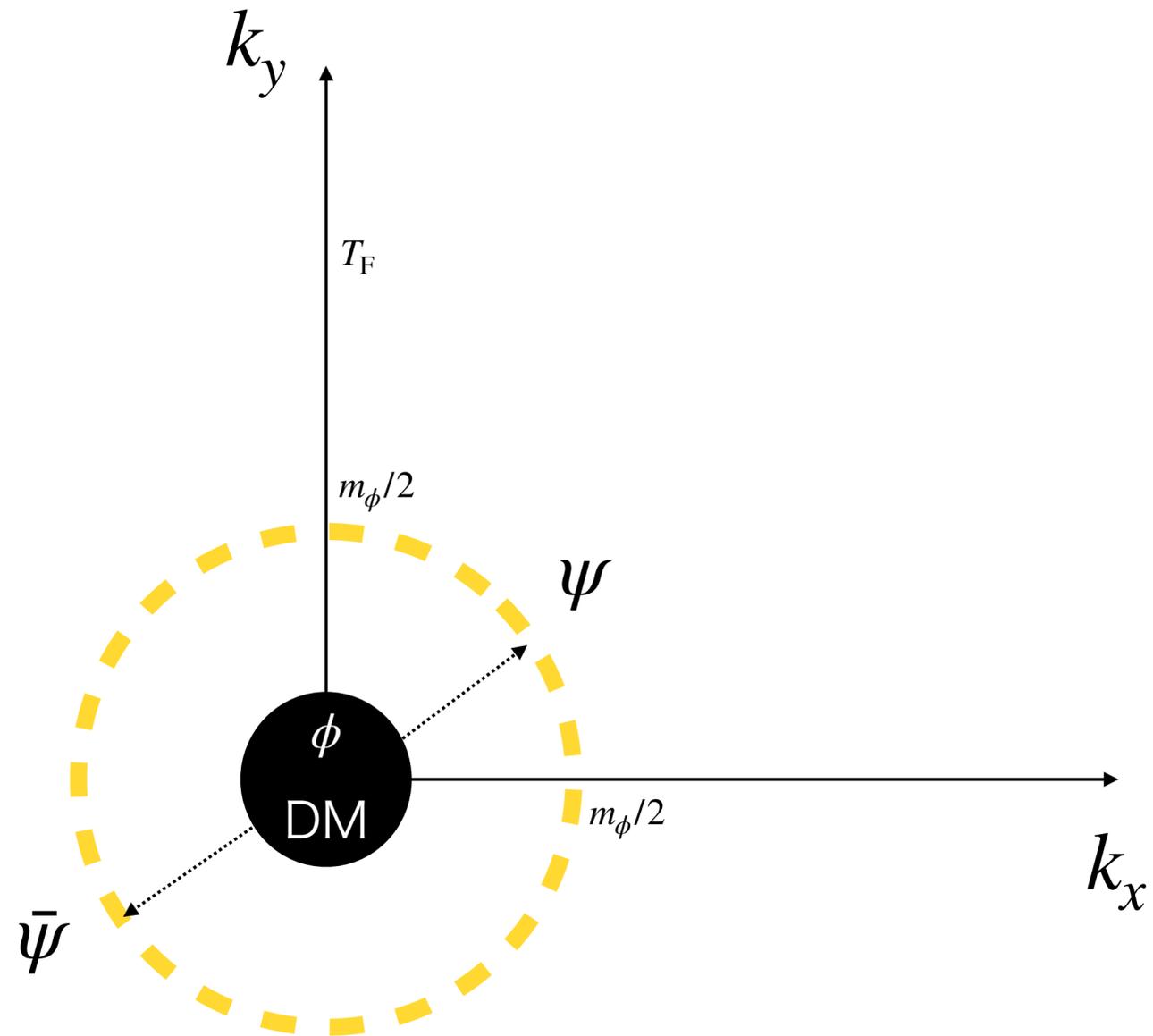
$$\mathcal{L} \supset -\frac{y}{2}\phi\bar{\psi}^c\psi \quad \phi : DM \quad \psi : \text{light fermion}$$

$$V = m_\phi^2\phi^2/2 \quad 1/\Gamma_{\phi\rightarrow\psi\psi} \ll 13.8\text{Gyr}$$

## Conclusions

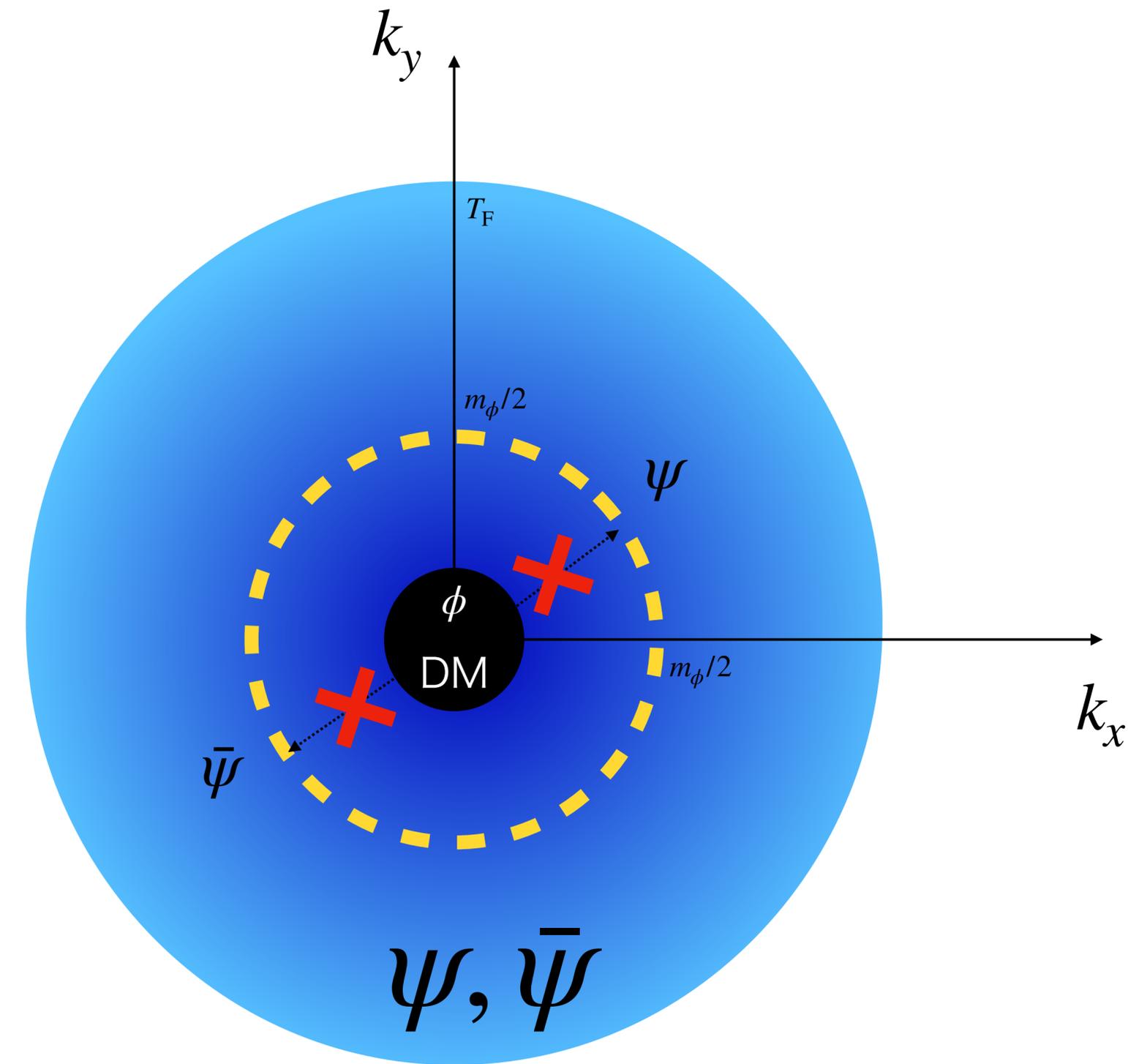
1. DM is stabilized due to Pauli-blocking from the produced fermi-gas.
2. Dark radiation of the fermi gas is self-interacting and may alleviate the Hubble tensions.

# Generic conditions for DM on Sea of Fermions



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$$T_F \gtrsim m_\phi/2$$

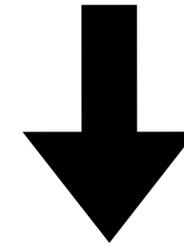


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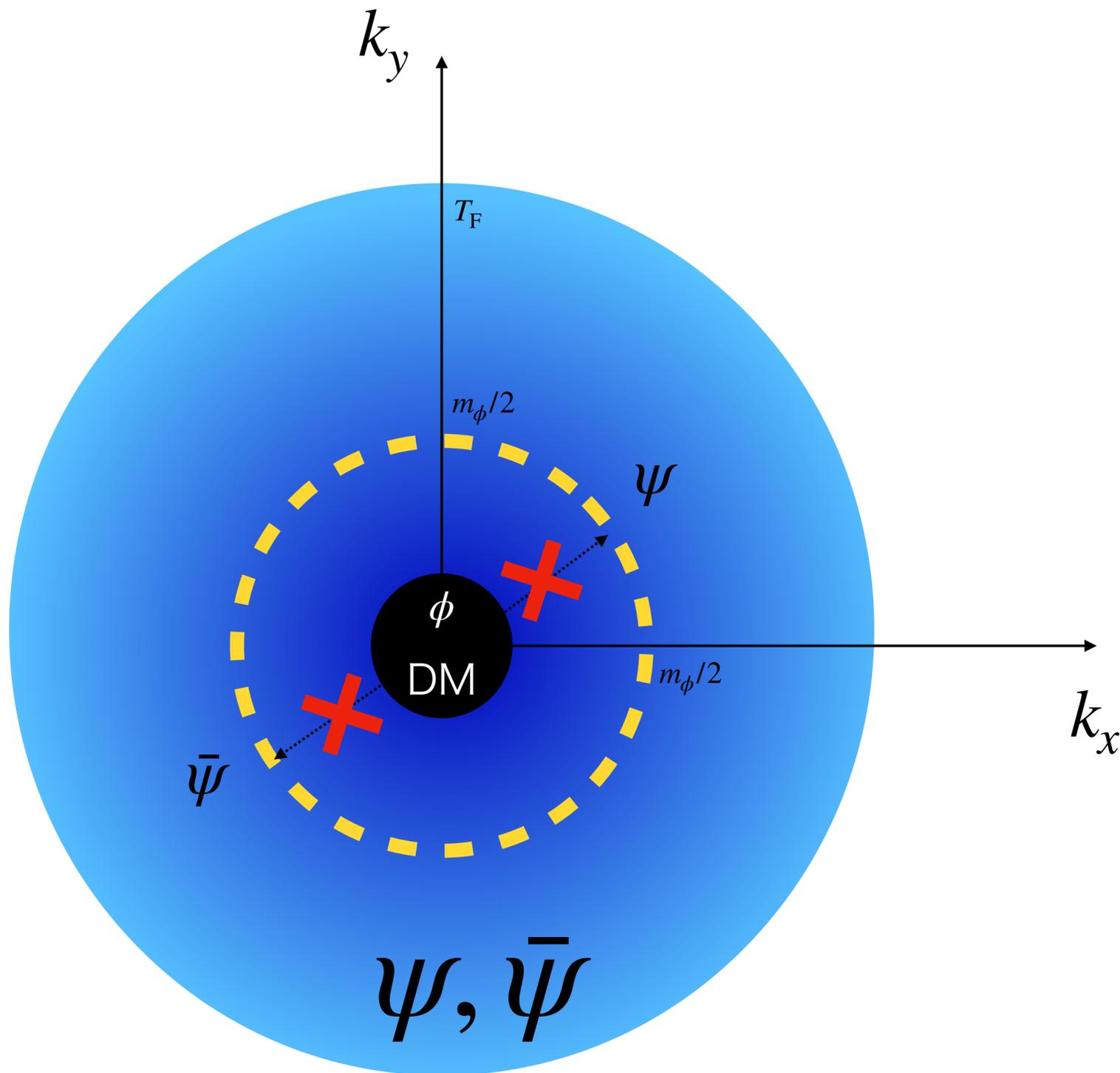
$$\rho_\phi \sim \text{keV cm}^{-3} \gg T_F^4$$



$$m_\phi \lesssim 0.01 \text{eV}$$

Thus I will talk about a system

- DM is light
- Number density is large



# Cosmology of DM on Sea of Fermions

# DM production and very early stage

setup

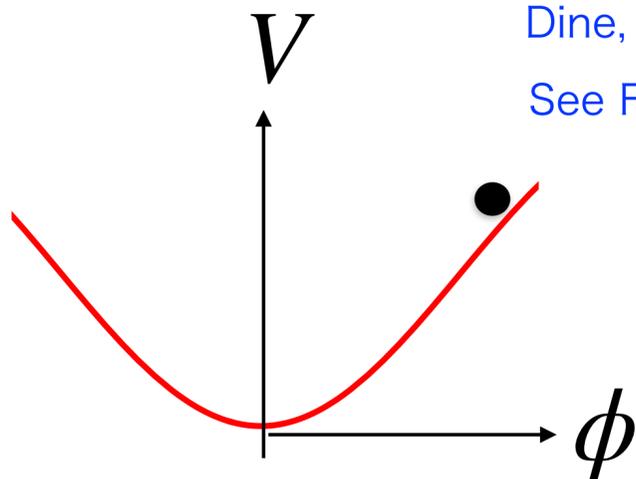
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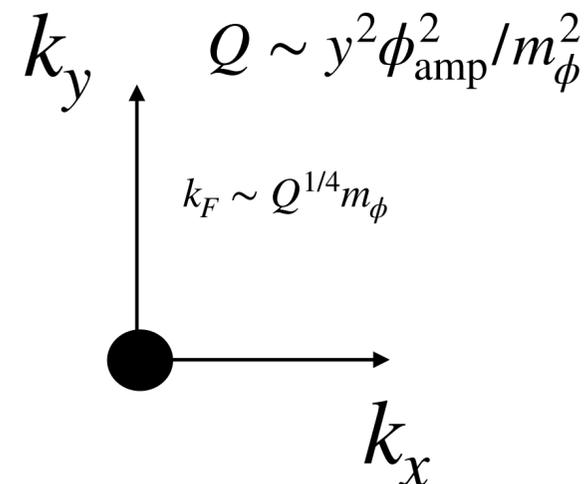
1. Misalignment mechanism to produce  $\phi$  condensate.

Preskill et al, 1983;  
Abbott, Sikivie, 1983;  
Dine, Fishler, 1983;  
See Fumi's talk



2. Initial Fermi sea produced via parametric resonance

Greene:1998nh, Baacke:1998di, Greene:2000ew



3. DM is stable until the  $Q \sim 1$

$$\rho_\psi \sim \frac{g_\psi}{2\pi^2} Q^{5/4} m_\phi^4 = \frac{g_\psi}{\pi^2} y^2 Q^{1/4} \rho_\phi \sim m_\phi^4 \quad (Q \sim 1)$$

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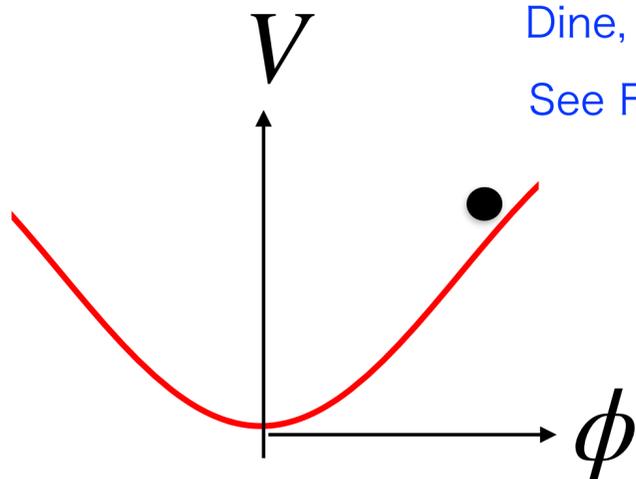
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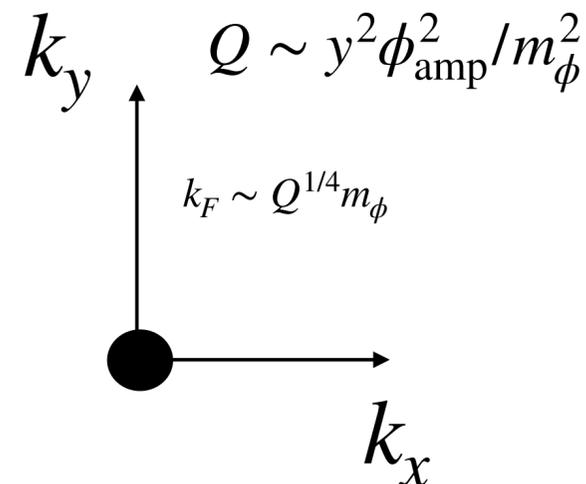
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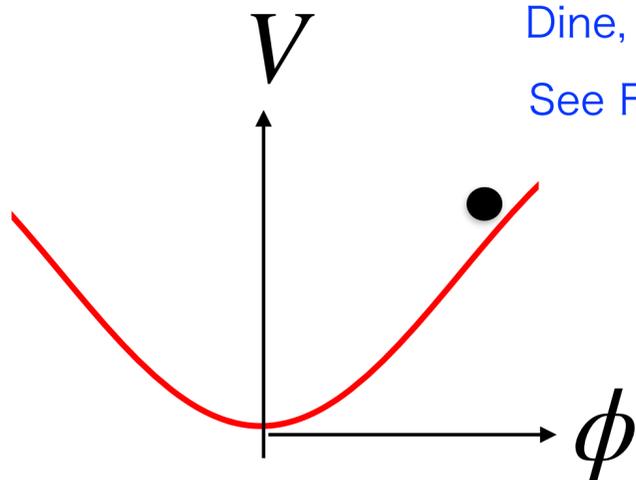
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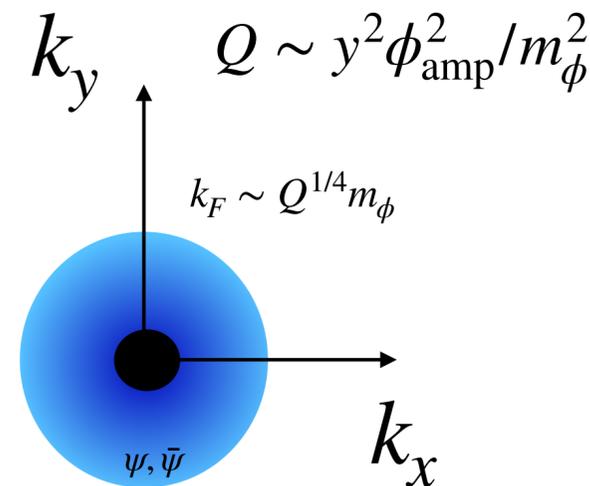
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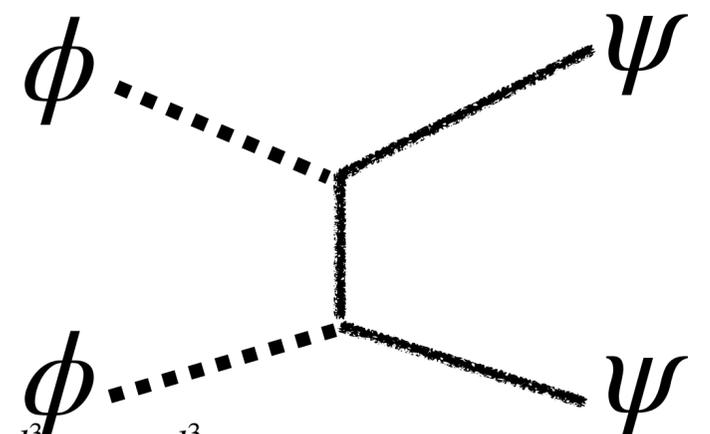
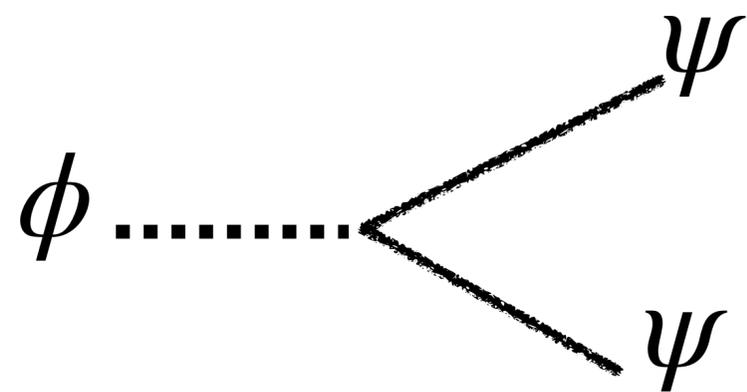
# DM stability in perturbative regime

Q redshifts to  $Q \ll 1$ , we can use the Boltzmann equation to study the stability.

Boltzmann equation

$$\frac{\partial f_\phi(p_\phi, t)}{\partial t} - p_\phi H \frac{\partial f_\phi(p_\phi, t)}{\partial p_\phi} = C^\phi(p_\phi, t),$$

$$\frac{\partial f_\psi(p_\psi, t)}{\partial t} - p_\psi H \frac{\partial f_\psi(p_\psi, t)}{\partial p_\psi} = C^\psi(p_\psi, t).$$



e.g.  $C^\phi \supset$

$$C_{\phi \leftrightarrow \psi\psi}^\phi = -\frac{1}{S_\psi} \frac{1}{g_\phi} \frac{1}{2E_\phi} \sum_{\text{spins}} \int \frac{d^3 p_1}{(2\pi)^3 2E_1} \frac{d^3 p_2}{(2\pi)^3 2E_2} (2\pi)^4 \delta^4(p_\phi - p_1 - p_2) |\mathcal{M}|^2$$

$$\times \left\{ f_\phi(p_\phi) [1 - f_\psi(p_1)] [1 - f_\psi(p_2)] - [1 + f_\phi(p_\phi)] f_\psi(p_1) f_\psi(p_2) \right\}.$$

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$$\times \left\{ f_\phi(p_\phi) f_\phi(p_2) [1 - f_\psi(p_3)] [1 - f_\psi(p_4)] - [1 + f_\phi(p_\phi)] [1 + f_\phi(p_2)] f_\psi(p_3) f_\psi(p_4) \right\}.$$

.....

Pauli-blocking, Bose-enhancement factors are very important

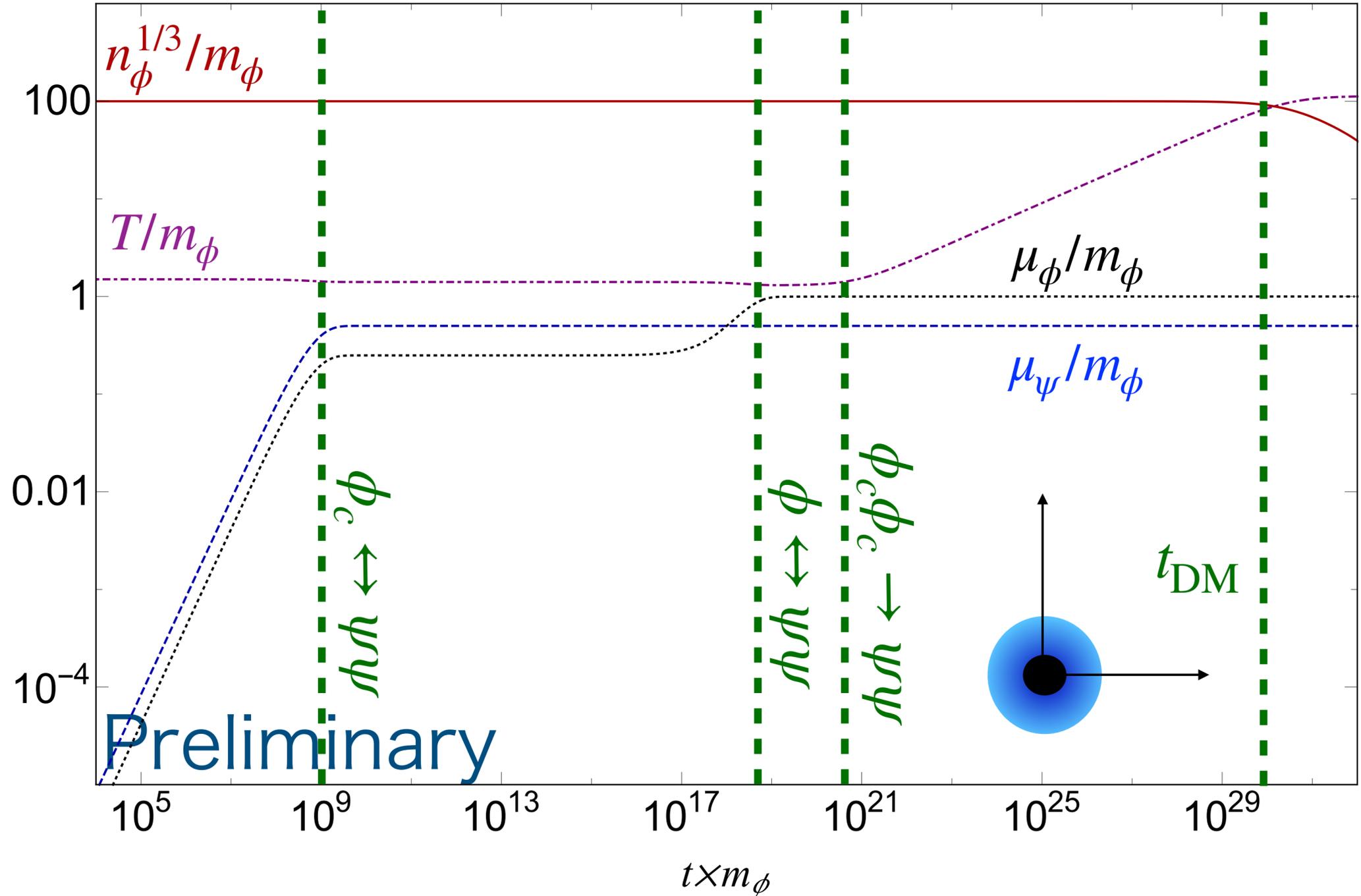
# Life-time of the DM

We solve the Boltzmann equation by neglecting Hubble expansion with some approximations.

Initial conditions:

$n_\phi[0]/m_\phi^3 = 10^8,$   
 $T[0] = 3/2 m_\phi,$   
 $\mu_\psi[0] = \mu_\phi[0] = 0,$   
 $y = 10^{-8}$   
 $M_\psi = m_\phi/50$

lifetime in the vacuum    Slow-thermalization    DM lifetime



Preliminary

# Life-time of the DM

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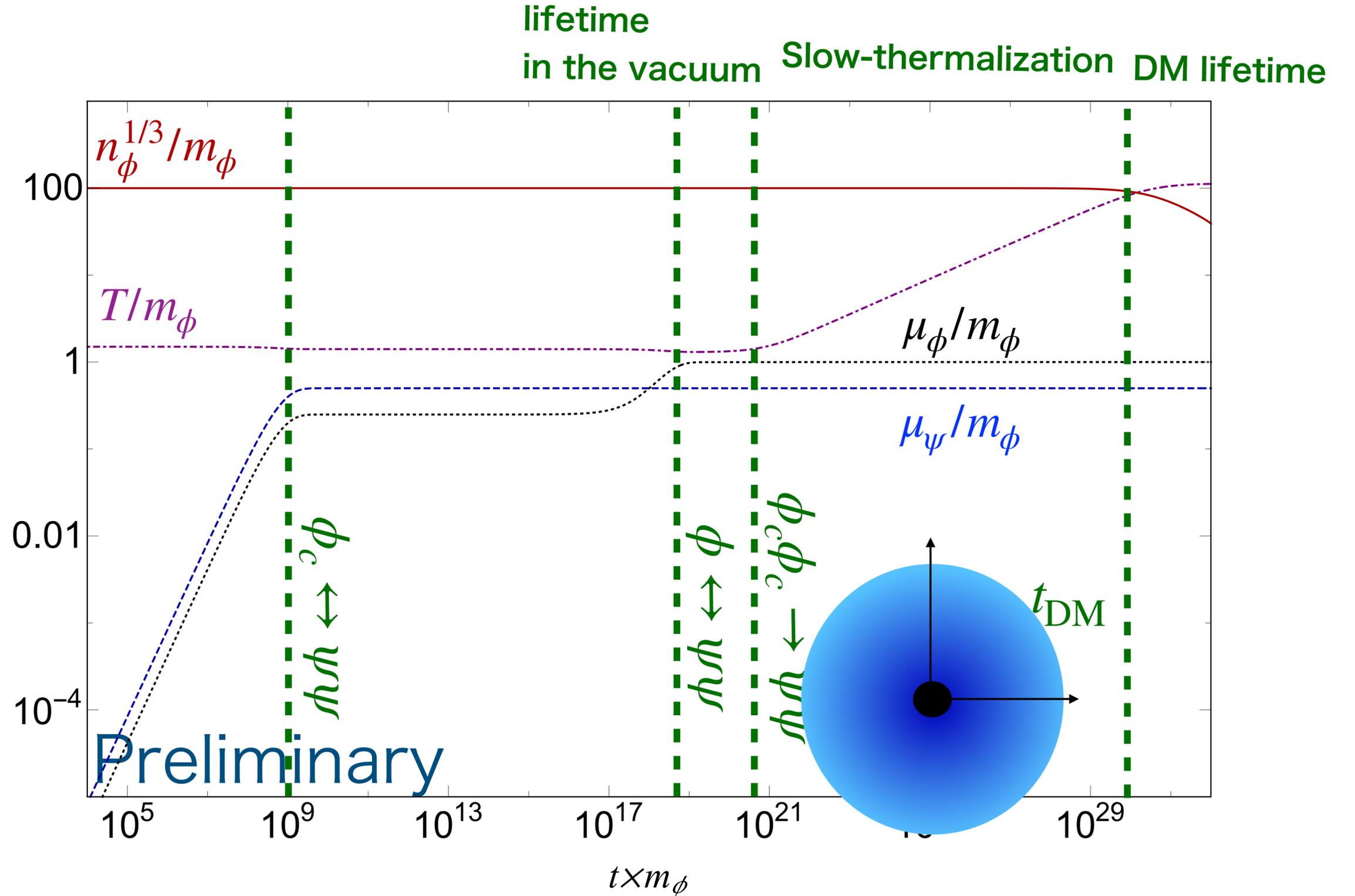
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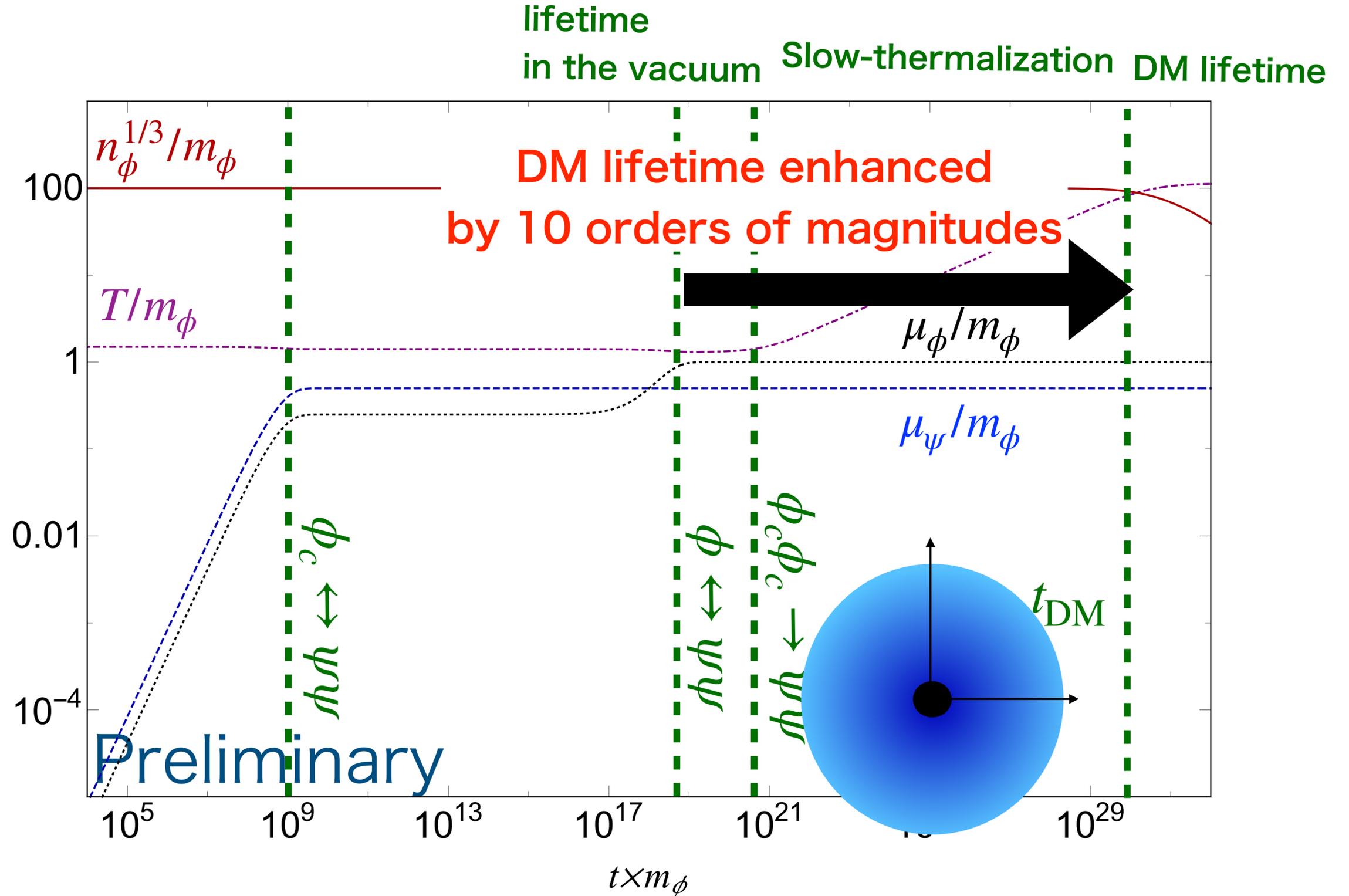
$$n_\phi[0]/m_\phi^3 = 10^8,$$

$$T[0] = 3/2 m_\phi,$$

$$\mu_\psi[0] = \mu_\phi[0] = 0,$$

$$y = 10^{-8}$$

$$M_\psi = m_\phi/50$$



# Dark radiation self-interaction prevents DM from decay.

Initial conditions:

conditions:

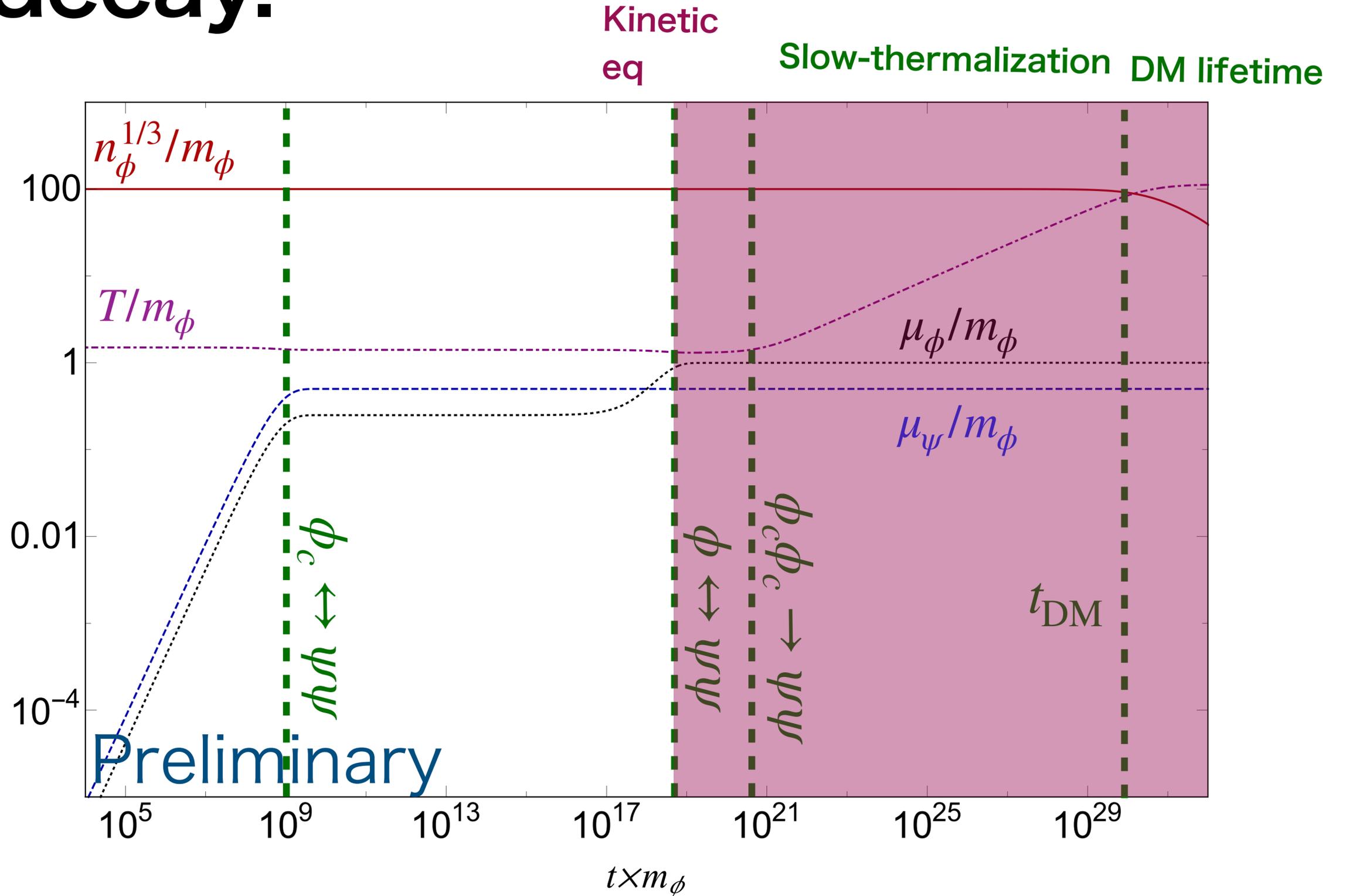
$$n_\phi[0]/m_\phi^3 = 10^8,$$

$$T[0] = 3/2 m_\phi,$$

$$\mu_\psi[0] = \mu_\phi[0] = 0,$$

$$y = 10^{-8}$$

$$M_\psi = m_\phi/50$$



Produced Dirac sea is self-interacting!

# Dark radiation and deviation from $\Lambda$ CDM

We solve (integrated) Boltzmann equation with Hubble expansion

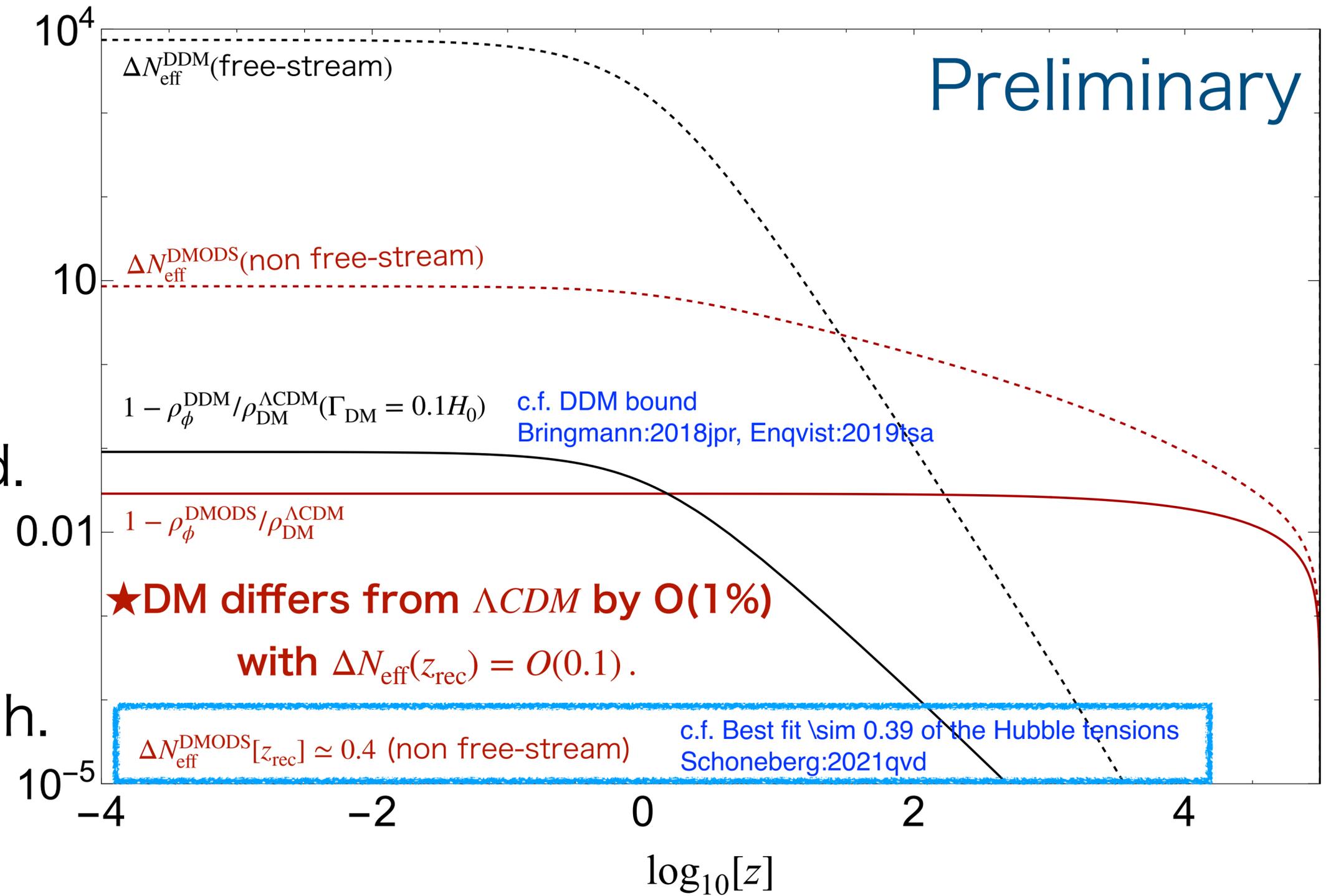
$$y = 7.4 \times 10^{-10}, m_\phi = 10^{-4} eV \text{ and } M_\psi = m_\phi/5.$$

Initial conditions:

$$\rho_\phi(z = 10^5) = \rho_{\text{DM}}^{\Lambda\text{CDM}} \text{ and } T[z = 10^5] = m_\phi.$$

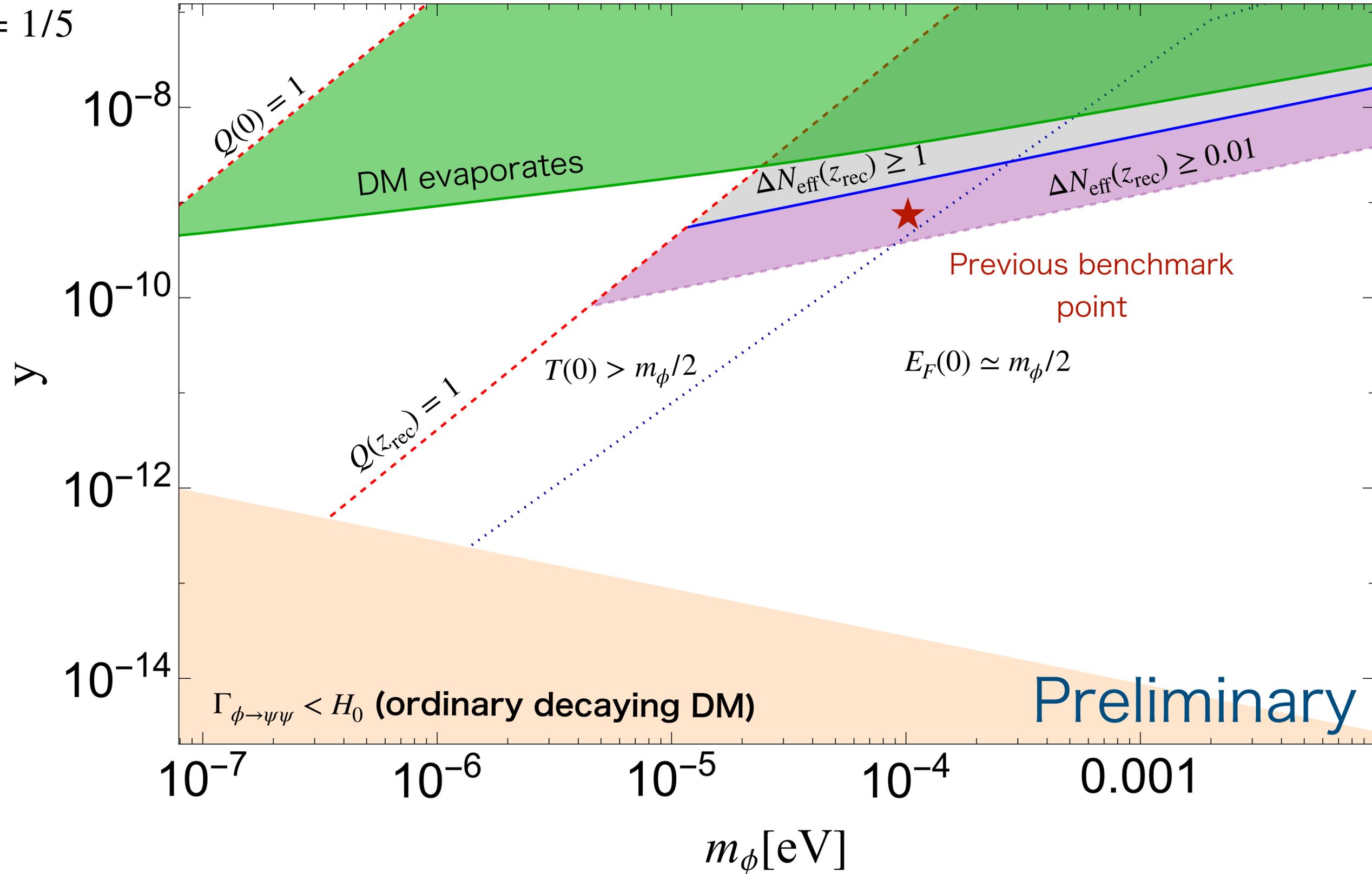
◦ Self-interacting DR significantly produced.

◦ DM comoving density does not change much.



# Parameter region

$$M_\psi/m_\phi = 1/5$$



# Conclusions

1. ALP DW without a string following scaling solution naturally explains the isotropic CB. [Takahashi, WY, 2012.11576](#)

2. The scenario can be fully tested in the future observation of anisotropic CB. [Takahashi, Kitajima, Kozai, WY, in progress](#)

1. DM can be stabilized if  $m_\phi < 0.01 eV, y \lesssim 10^{-7}$ , even if  $\Gamma_{decay,vac} \gg H_0$ . [Batell, WY, in progress.](#)

2. The resulting dark radiation of the Fermi gas is self-interacting and may alleviate the Hubble tension.

3. The present Universe DM density has an upper bound. DM stabilized by CnuB may predict the enhancement of  $\nu_e$  capturing rate by  $\sim 5$  in PTOLEMY.

(will appear in our paper)