

Implications of the cosmic birefringence measurement for the axion dark matter search

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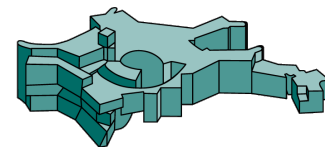
ダークマターの正体は何か？

広大なディスカバリースペースの網羅的研究

What is dark matter? - Comprehensive study of the huge discovery space in dark matter



文部科学省
科学研究費助成事業
学術変革領域研究
(2020-2024)



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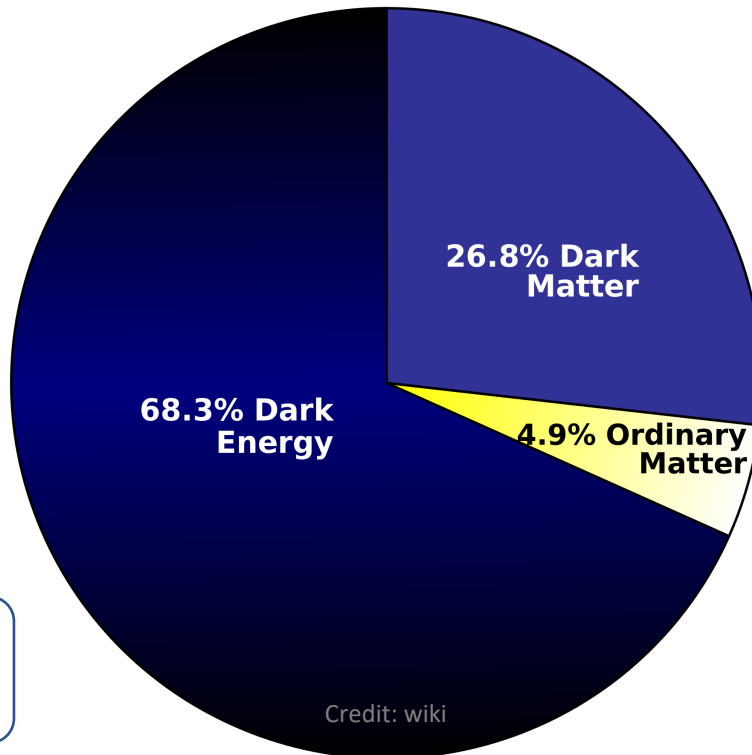
2022.3.29 FY2021 学術変革領域研究「ダークマター」シンポジウム

Dark sectors in our universe



Credit: higgstan.com

Cosmological Constant?
Quintessence?



Credit: higgstan.com

WIMP?
Axion?
MACHO?
PBHs?

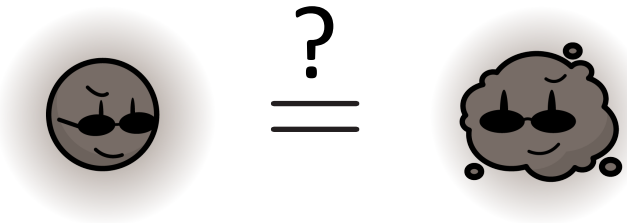
(standard model)

We know only 5 % in our universe!

Question

Are these sectors independent components?

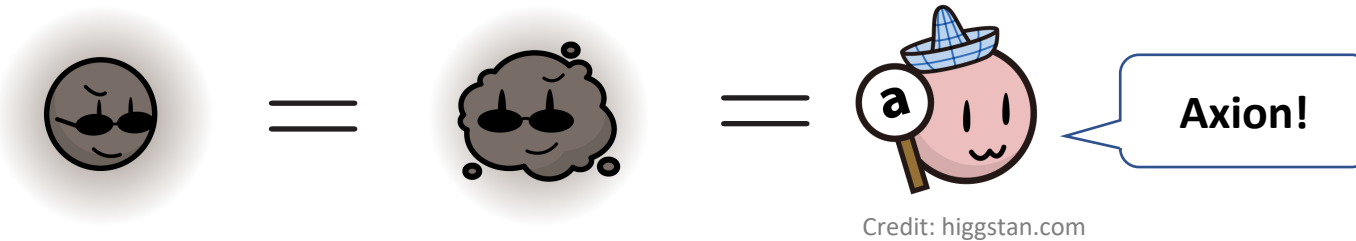
What if they are related by a common origin?



Unified models for dark matter and dark energy have been developed

- **Generalized Chaplygin gas:** Bento, Bertolami & Sen (2002); Makler, Oliveira, Waga (2003); ...
- **k-essence:** Scherrer (2004); Giannakis & Hu (2005); ...
- **Fast transition models:** Bruni, Lazkoz & Fernandez (2013); Leanizbarrutia, Fernandez & Tereno (2017); ...
- ...

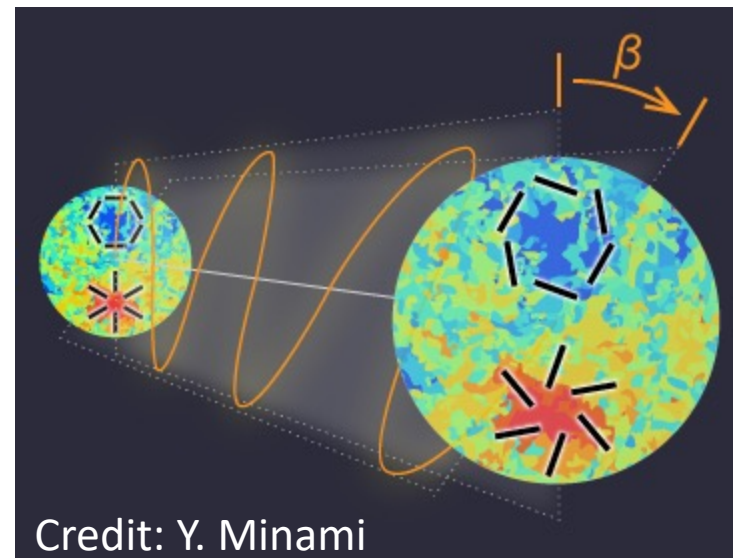
Overview of this talk



- Axion can be a candidate for both dark matter and dark energy

Motivation

- ✓ The constraints on this scenario are potentially connected by the measurement of cosmic birefringence effect in CMB!
- ✓ We can do a complementary search between CMB observations (axion dark energy) and resonant cavity experiments (axion dark matter)



Photon's birefringence by axion

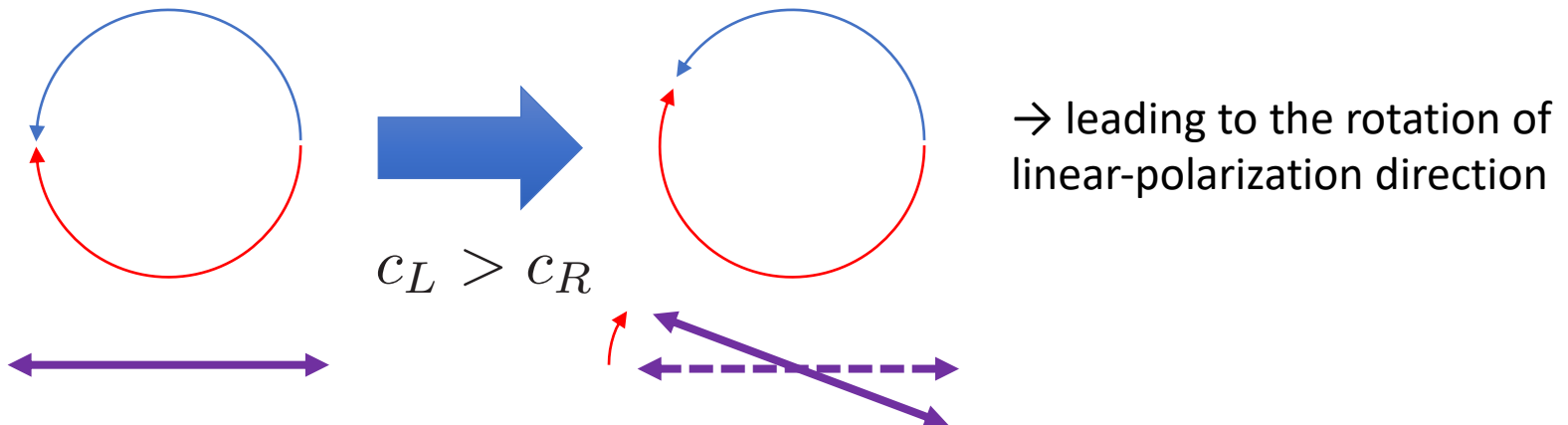
Carroll, Field & Jackiw (1990); Harari & Sikivie (1992); Carroll (1998); ...

Axion behaves as a birefringent material in our universe

➤ Axion differentiates the phase velocities of circular-polarized photon

$$\mathcal{L} \supset \frac{1}{4} g_{a\gamma} \varphi F_{\mu\nu} \tilde{F}^{\mu\nu}, \quad F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu$$

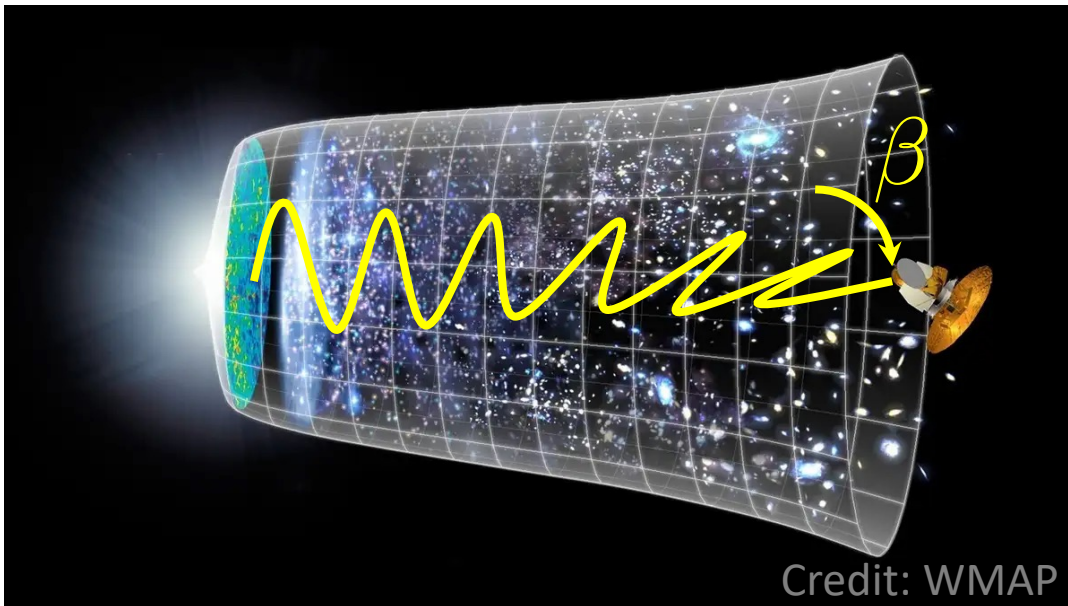
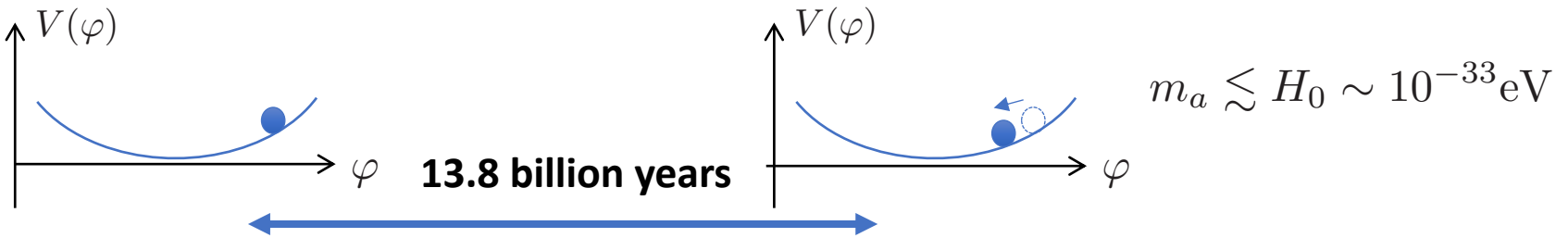
$$\text{Dispersion relation: } \ddot{A}_k^{L/R} + \omega_{L/R}^2 A_k^{L/R} = 0, \quad c_{L/R} \equiv \frac{\omega_{L/R}}{k} = \sqrt{1 \pm \frac{g_{a\gamma} \dot{\varphi}}{k}}$$



CMB Birefringence by axion DE

(Fukugita & Yanagida (1994); Friemann+ (1995); J.E.Kim+ (1999); ...)

- Axion with mass smaller than current Hubble scale behaves as a dark energy



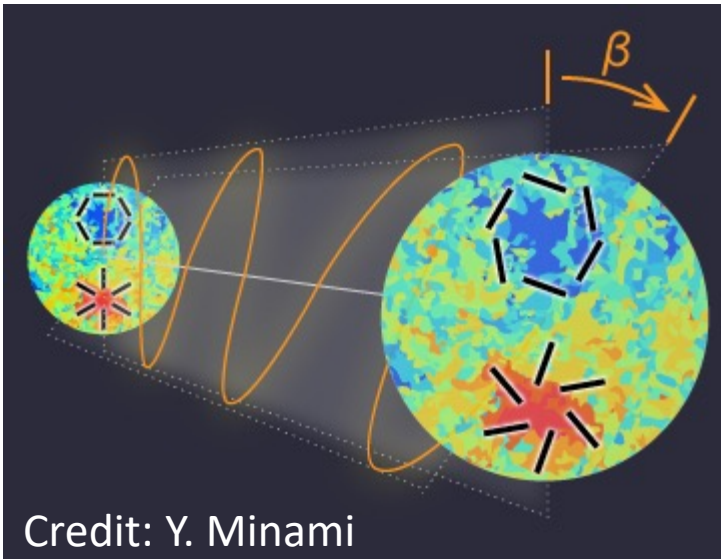
- If axion is responsible for dark energy, it makes the polarization plane of CMB rotate from the last-scattering-surface

Rotation angle:

$$\beta = \frac{g_{\phi\gamma}}{2} \Delta\phi \equiv \frac{g_{\phi\gamma}}{2} (\phi_0 - \langle \phi_{\text{LSS}} \rangle)$$

Generation EB correlation function

Lue, Wang & Kamionkowski (1999); Feng+ (2005,2006); Liu, Lee & Ng (2006); ...



Parity-violating interaction

$$\text{e.g. } \mathcal{L}_{\text{int}} = \frac{1}{4} g_{a\gamma} \varphi F_{\mu\nu} \tilde{F}^{\mu\nu}$$

produces the parity-odd EB correlation

$$C_{\ell}^{EB,o} = \frac{1}{2} \sin(4\beta) \left(C_{\ell}^{EE,\text{CMB}} - C_{\ell}^{BB,\text{CMB}} \right) + \cos(4\beta) C_{\ell}^{EB,\text{CMB}}$$

↑ measured value

↑ usually assume 0

History of measurements (WMAP, Planck, ACT,...)

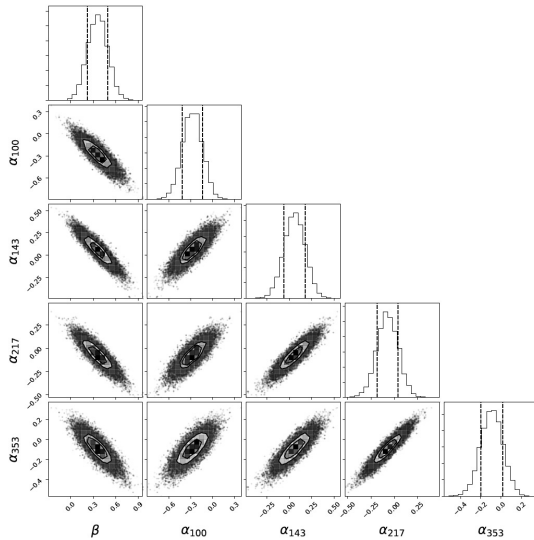
Non-zero $\langle EB \rangle$ has been detected.

But, not reliable estimates due to the miscalibration of instrumental angle α .

Foreground-based calibration

Minami & Komatsu (2020);

calibrate α by using the polarized emission from the galactic foregrounds and measures the intrinsic birefringence angle β by Planck 2018 PR3 data



Angles	Results (deg)
β	0.35 ± 0.14
α_{100}	-0.28 ± 0.13
α_{143}	0.07 ± 0.12
α_{217}	-0.07 ± 0.11
α_{353}	-0.09 ± 0.11

$$\beta = 0.35 \pm 0.14 \text{ deg } (2.4\sigma)$$

Diego-Palazuelos+ (2022)

$$\beta = 0.36 \pm 0.11 \text{ deg } (3.3\sigma)$$

Update! (PR4 data):

Implication for the axion search

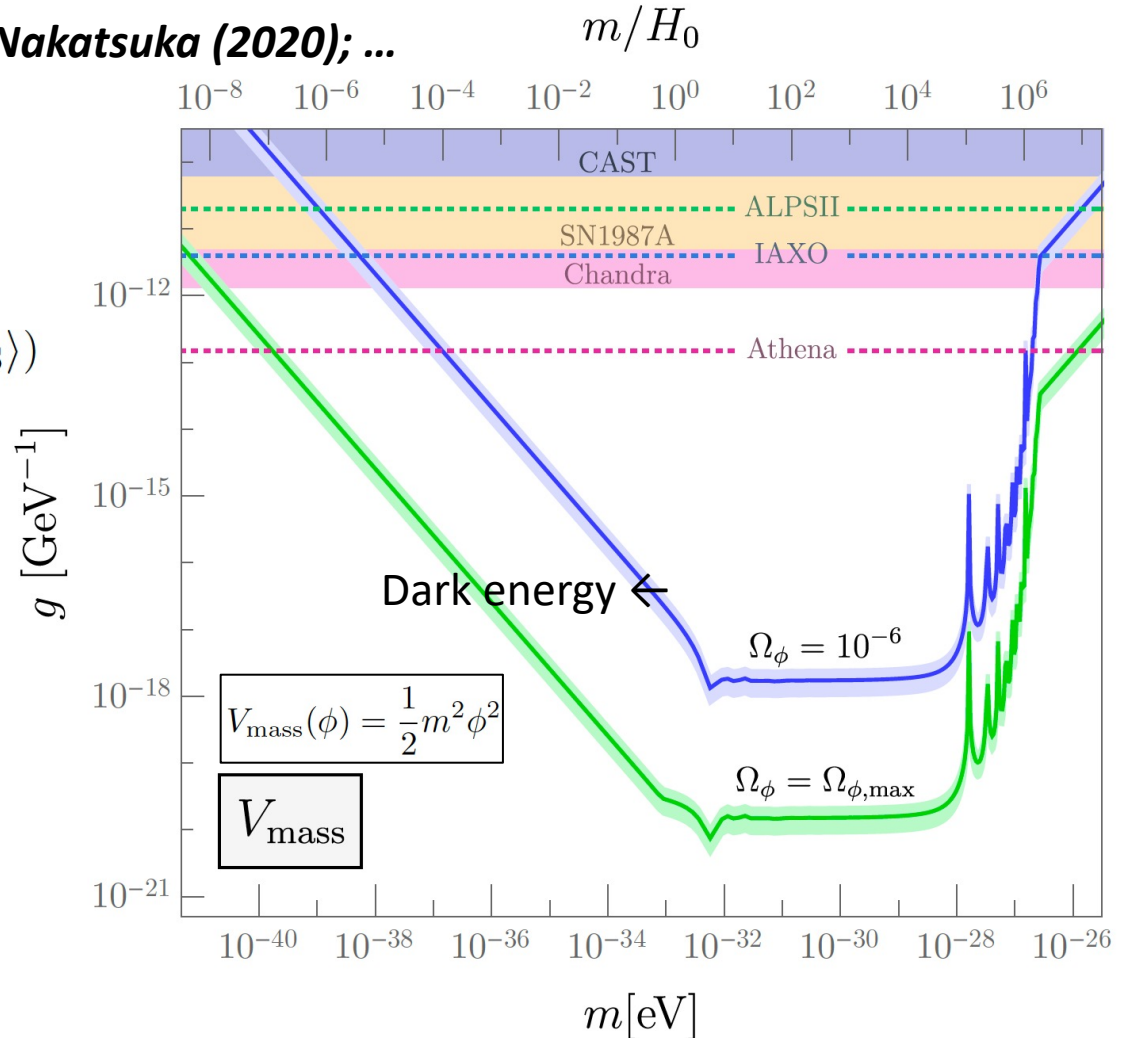
Fujita, Minami, Murai & Nakatsuka (2020); ...

- From this relationship

$$\beta = \frac{g_{\phi\gamma}}{2} \Delta\phi \equiv \frac{g_{\phi\gamma}}{2} (\phi_0 - \langle\phi_{\text{LSS}}\rangle)$$

we can constrain the parameter space of axion-photon coupling w.r.t. axion mass

- Support the presence of axion as dark energy

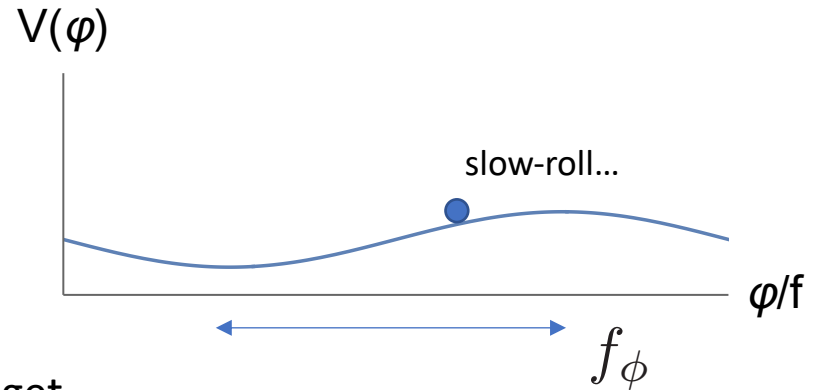


Some issues of single-field model

Friemann+ (1995); ...

- Consider a **nearly flat** axion cosine potential

$$V(\phi) = m_\phi^2 f_\phi^2 \left[1 - \cos \left(\frac{\phi}{f_\phi} \right) \right]$$



- To satisfy the constraint on EoS parameter, we get

$$f_\phi \simeq 14 M_{\text{Pl}} \left(\frac{\Omega_\phi}{0.69} \right)^{1/2} \left(\frac{m_\phi/H_0}{0.1} \right)^{-1} > M_{\text{Pl}}$$

requires a **super-Planckian** decay constant or a fine-tuning of initial axion displacement

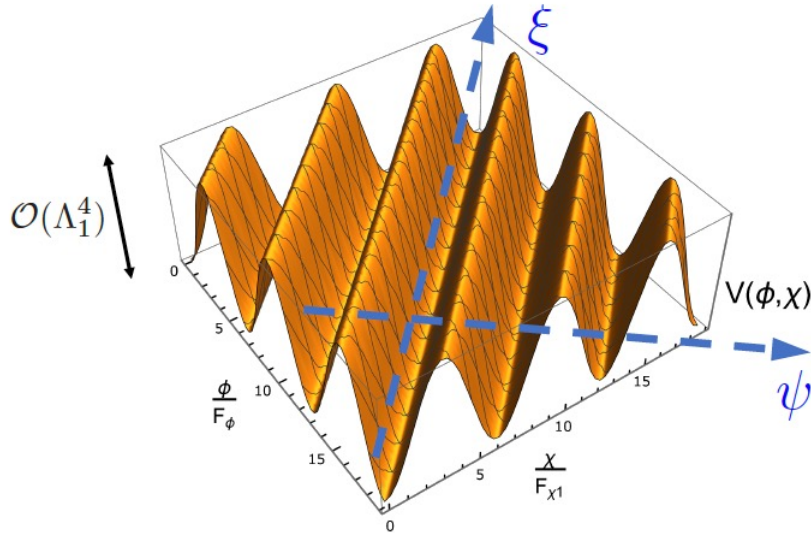
- To get the measured β , a **large anomaly coefficient** is required

$$g_{\phi\gamma} = \frac{\alpha}{2\pi} \frac{c_{\phi\gamma}}{f_\phi} \quad |c_{\phi\gamma}| \simeq 7.5 \times 10^3 \left(\frac{\beta}{0.35\text{deg}} \right) \left(\frac{m_\phi/H_0}{0.1} \right)^{-2} \gg 1$$

Two-fields axion model

Kim (1999)(2000), ...

Kim, Nilles & Peloso (2005), ...



Potential

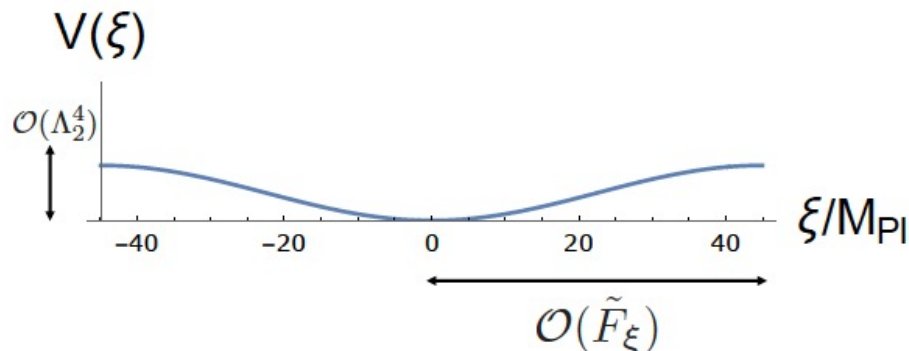
$$V(\phi, \chi) = \Lambda_1^4 \left[1 - \cos \left(\frac{\phi}{F_{\phi 1}} + \frac{\chi}{F_{\chi 1}} \right) \right] + \Lambda_2^4 \left[1 - \cos \left(\frac{\phi}{F_{\phi 2}} + \frac{\chi}{F_{\chi 2}} \right) \right]$$

$$(\Lambda_1^4 \gg \Lambda_2^4, F_i < M_{\text{Pl}})$$

$$\xi = \frac{F_\phi}{\sqrt{F_\phi^2 + F_{\chi 1}^2}} \phi - \frac{F_{\chi 1}}{\sqrt{F_\phi^2 + F_{\chi 1}^2}} \chi, \quad \psi = -\frac{F_{\chi 1}}{\sqrt{F_\phi^2 + F_{\chi 1}^2}} \phi - \frac{F_\phi}{\sqrt{F_\phi^2 + F_{\chi 1}^2}} \chi$$

- Nearly-flat direction can be realized by an alignment of the **multiple axion potentials**:

$$F_{\phi 1} = F_{\phi 2} \equiv F_\phi, \quad F_{\chi 2} = F_{\chi 1}(1 + \epsilon) \quad \epsilon \ll 1 : \text{misalignment parameter}$$



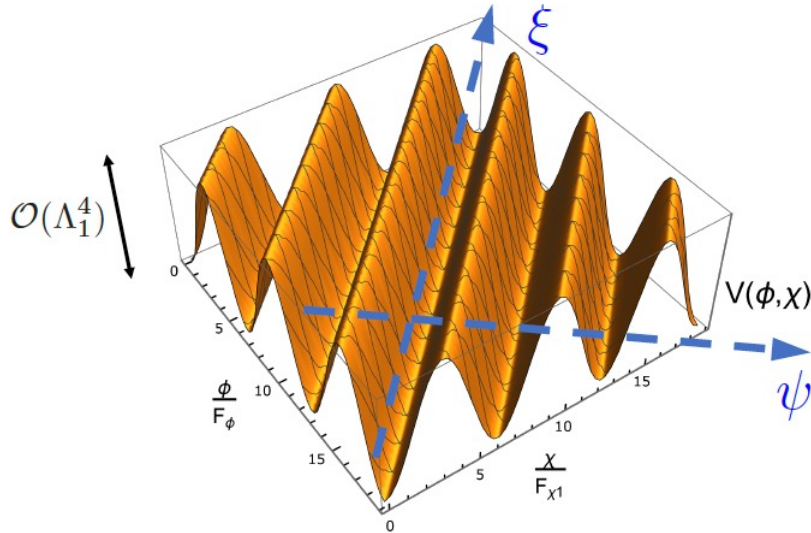
Effective field range of axion

$$\tilde{F}_\xi = \frac{\sqrt{F_\phi^2 + F_{\chi 1}^2}}{\epsilon} \gg M_{\text{Pl}}$$

Two-fields axion model

Kim (1999)(2000), ...

Kim, Nilles & Peloso (2005), ...



Potential

$$V(\phi, \chi) = \Lambda_1^4 \left[1 - \cos \left(\frac{\phi}{F_{\phi 1}} + \frac{\chi}{F_{\chi 1}} \right) \right] + \Lambda_2^4 \left[1 - \cos \left(\frac{\phi}{F_{\phi 2}} + \frac{\chi}{F_{\chi 2}} \right) \right]$$

$$(\Lambda_1^4 \gg \Lambda_2^4, F_i < M_{\text{Pl}})$$

$$\xi = \frac{F_\phi}{\sqrt{F_\phi^2 + F_{\chi 1}^2}} \phi - \frac{F_{\chi 1}}{\sqrt{F_\phi^2 + F_{\chi 1}^2}} \chi, \quad \psi = -\frac{F_{\chi 1}}{\sqrt{F_\phi^2 + F_{\chi 1}^2}} \phi - \frac{F_\phi}{\sqrt{F_\phi^2 + F_{\chi 1}^2}} \chi$$

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Linear combinations of two-fields provide two **(nearly) massless** & **massive** fields

Dark energy

Dark matter

ξ

ψ

Introduce axion-photon couplings

Obata (2021)

- The interactions of photon to the (original) axion fields are given by

$$\mathcal{L} \supset \frac{\alpha}{8\pi} \left(\frac{\phi}{F_{\phi\gamma}} + \frac{\chi}{F_{\chi\gamma}} \right) F_{\mu\nu} \tilde{F}^{\mu\nu}$$

- In terms of (ψ, ξ) , the effective coupling constants are obtained:

$$g_{\xi\gamma} = \frac{\alpha}{2\pi} \frac{c_{\xi\gamma}}{\tilde{F}_{\xi}}, \quad c_{\xi\gamma} \equiv \frac{1}{\epsilon} \left(\frac{F_{\phi}}{F_{\phi\gamma}} - \frac{F_{\chi 1}}{F_{\chi\gamma}} \right),$$
$$g_{\psi\gamma} = \frac{\alpha}{2\pi} \frac{c_{\psi\gamma}}{\tilde{F}_{\psi}}, \quad c_{\psi\gamma} \equiv - \left(\frac{F_{\phi}}{F_{\chi\gamma}} + \frac{F_{\chi 1}}{F_{\phi\gamma}} \right) \frac{F_{\phi} F_{\chi 1}}{F_{\phi}^2 + F_{\chi 1}^2}$$

(release 3)

- $g_{\xi\gamma}$ (dark energy) is fixed by the measured birefringence angle $\beta = 0.35 \pm 0.14$

→ the parameter space of $g_{\psi\gamma}$ (dark matter) is also constrained

Parameter Search (1)

- (Effective) axion DE field range: $\tilde{F}_\xi \simeq 14M_{\text{Pl}} \left(\frac{\Omega_\phi}{0.69} \right)^{1/2} \left(\frac{m_\phi/H_0}{0.1} \right)^{-1}$
leads to the tuning of the small misalignment parameter:

$$\epsilon \simeq 7.0 \times 10^{-2} \frac{\sqrt{F_\phi^2 + F_{\chi 1}^2}}{M_{\text{Pl}}} \left(\frac{\Omega_\xi}{0.69} \right)^{-1/2} \left(\frac{m_\xi/H_0}{0.1} \right)$$

- Birefringence condition: $|c_{\xi\gamma}| \simeq 7.5 \times 10^3 \left(\frac{\beta}{0.35\text{deg}} \right) \left(\frac{m_\xi/H_0}{0.1} \right)^{-2}$

leads to the condition of the ratio of decay constants:

$$\left| \frac{F_\phi}{F_{\phi\gamma}} - \frac{F_{\chi 1}}{F_{\chi\gamma}} \right| \simeq 5.2 \times 10^2 \frac{\sqrt{F_\phi^2 + F_{\chi 1}^2}}{M_{\text{Pl}}} \left(\frac{\beta}{0.35\text{deg}} \right) \left(\frac{\Omega_\xi}{0.69} \right)^{-1/2} \left(\frac{m_\xi/H_0}{0.1} \right)^{-1}$$

Parameter Search (2)

- Axion DM abundance by misalignment production: $\Omega_\psi \simeq \frac{1}{6} \left(\frac{\tilde{F}_\psi}{M_{\text{Pl}}} \right)^2 (9\Omega_r)^{3/4} \left(\frac{m_\psi}{H_0} \right)^{1/2}$

leads to the condition of axion DM decay constant:

Marsh & Ferreira (2010);

$$\tilde{F}_\psi \simeq 3.8 \times 10^{-2} M_{\text{Pl}} \left(\frac{\Omega_\psi}{0.31} \right)^{1/2} \left(\frac{m_\psi}{10^{-22} \text{eV}} \right)^{-1/4}$$

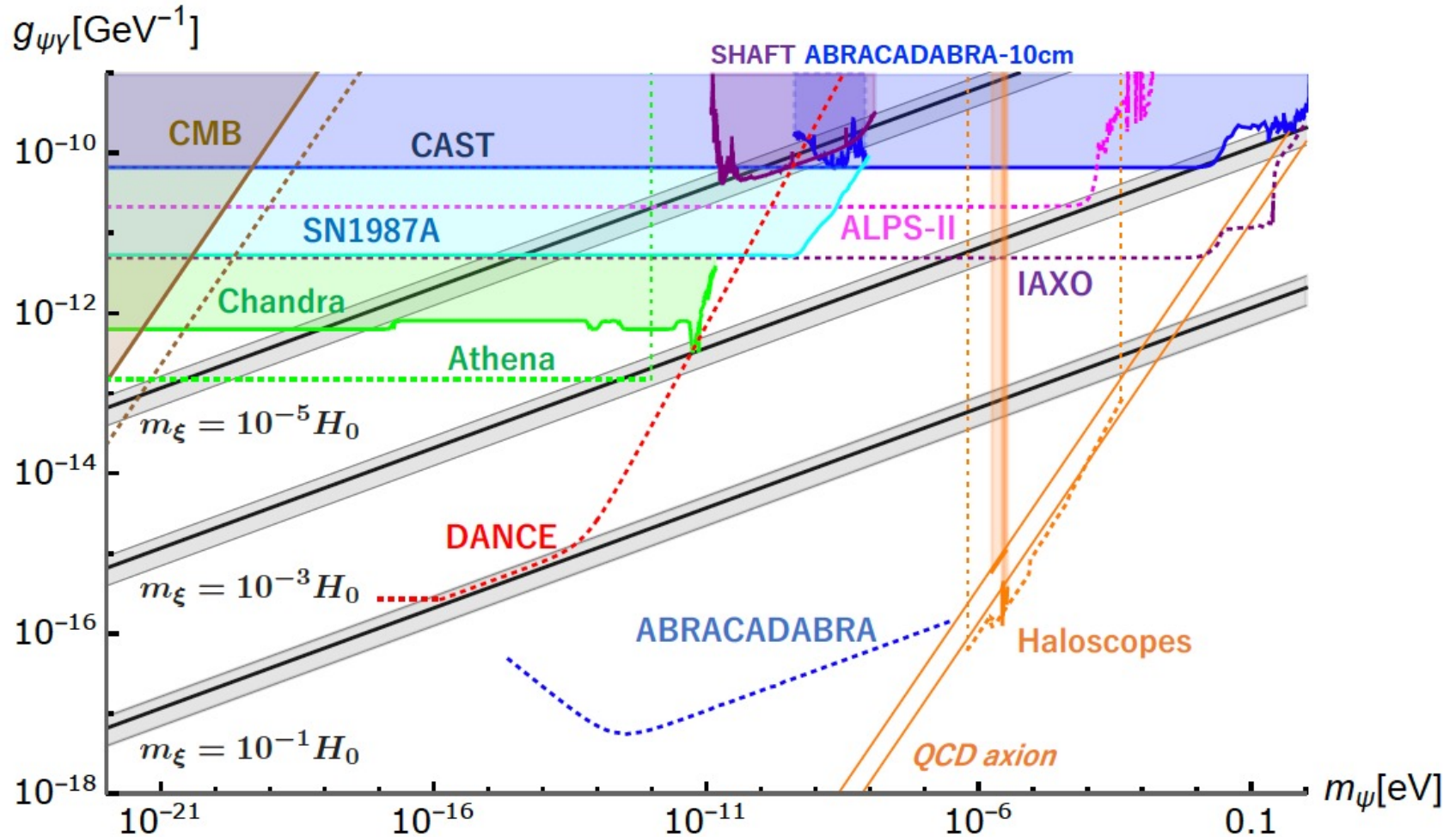
- The axion DM-photon coupling constant is obtained by:

$$g_{\psi\gamma} = \frac{\alpha}{2\pi} \frac{c_{\psi\gamma}}{\tilde{F}_\psi}, \quad c_{\psi\gamma} \equiv - \left(\frac{F_\phi}{F_{\chi\gamma}} + \frac{F_{\chi 1}}{F_{\phi\gamma}} \right) \frac{F_\phi F_{\chi 1}}{F_\phi^2 + F_{\chi 1}^2}$$

$$|g_{\psi\gamma}| \simeq 6.6 \times 10^{-18} \text{ GeV}^{-1} \frac{F_{\chi 1}}{M_{\text{Pl}}} \left(\frac{\Omega_\xi}{0.69} \right)^{-1/2} \left(\frac{\Omega_\psi}{0.31} \right)^{-1/2} \\ \times \left(\frac{\beta}{0.35 \text{deg}} \right) \left(\frac{m_\xi/H_0}{0.1} \right)^{-1} \left(\frac{m_\psi}{10^{-22} \text{eV}} \right)^{1/4} .$$

(assuming the region $F_\phi \ll F_{\chi 1}, F_{\chi 1} \ll F_{\chi\gamma}$)

Parameter space of axion DM



Summary & Outlook

- Axion is one of the promising candidates for the dark sector of our universe.
- Photon's birefringence measurements potentially develop a new frontier of the axion search! (b01, b06)
- A recent measurement of CMB birefringence gives us a tantalizing hint for the axion physics, especially for the axion as dark energy.
- Based on a multiple axion scenario, this observable can connect the constraints on axion as dark energy and dark matter.
- Extensions of this scenario to more generic ones? (e.g. N-field, kinetic mixings, ...)