Update from group C02: **Cosmic structure formation**

Shin'ichiro Ando University of Amsterdam / University of Tokyo



GRavitation AstroParticle Physics Amsterdam

Dark Matter Symposium March 29, 2022

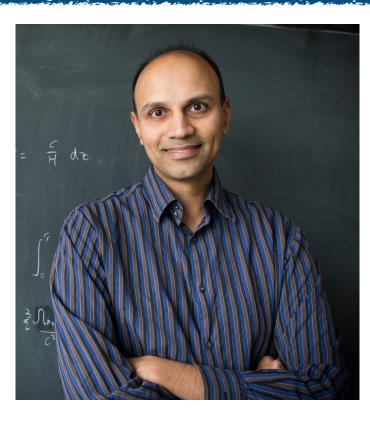


C02: Members

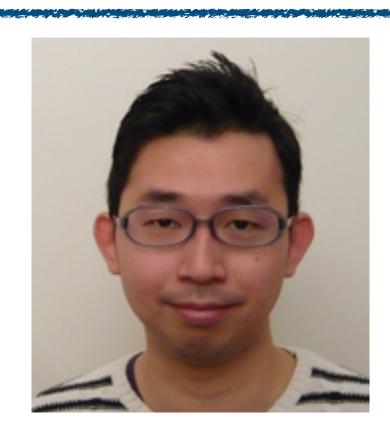
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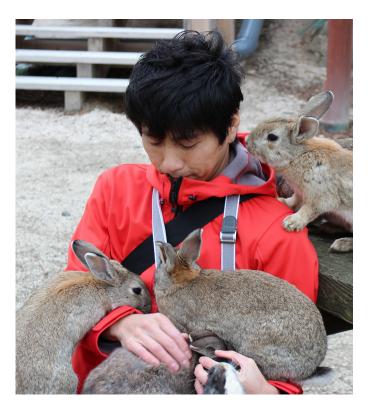
Shin'ichiro Ando



Neal Dalal



Takahiro Nishimichi



Takashi Okamoto



Masato Shirasaki



Atsushi Taruya





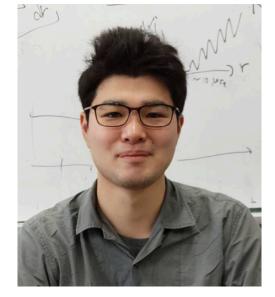


Shunichi Horigome

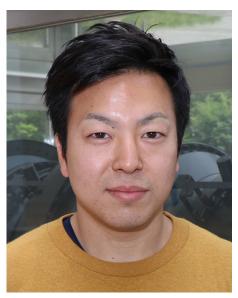


Shigeki Inoue

Collaborators



Yohsuke Enomoto



Kohei Hayashi

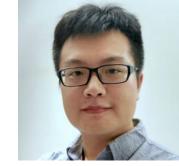




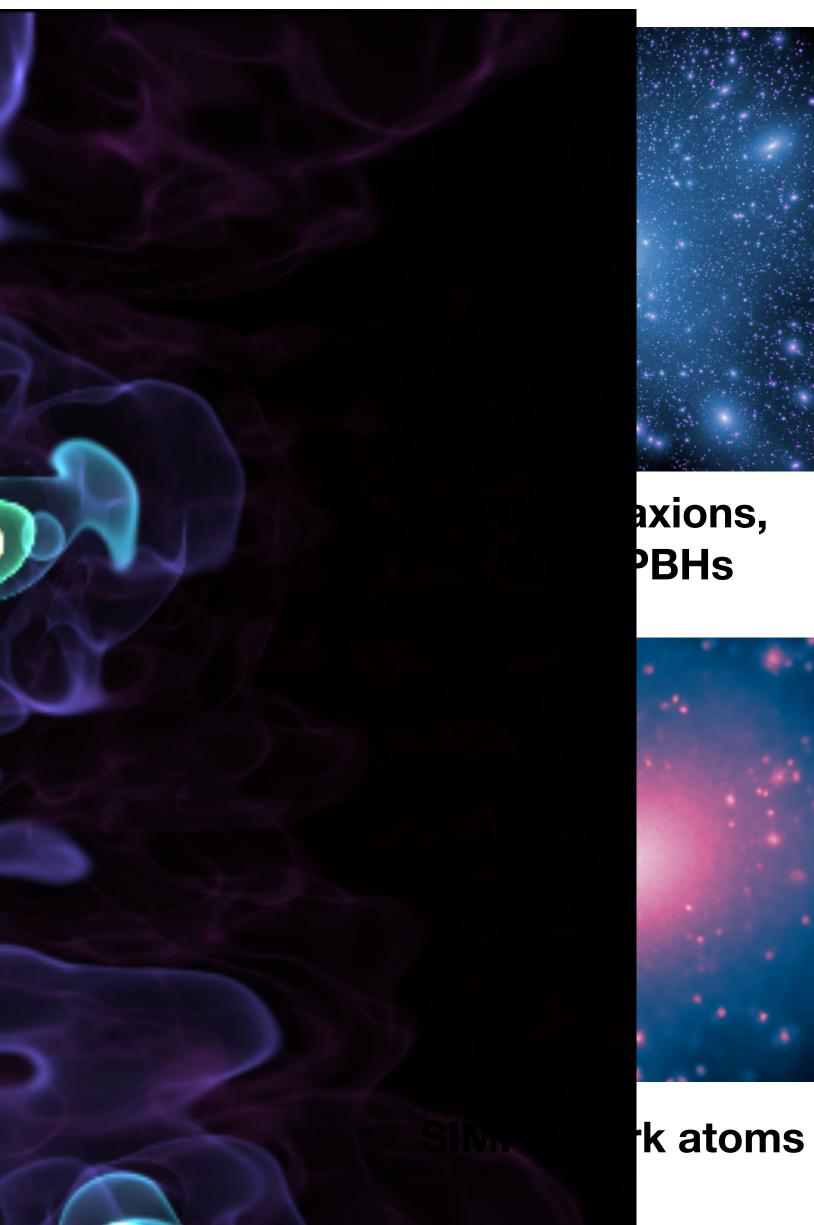


A. Dekker

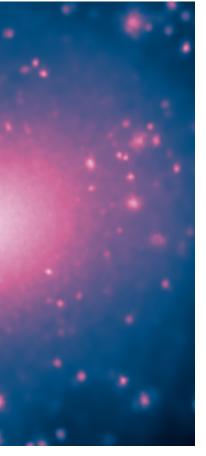




K.C.Y. Ng

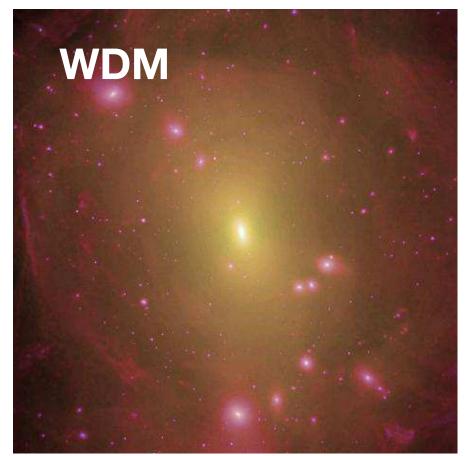






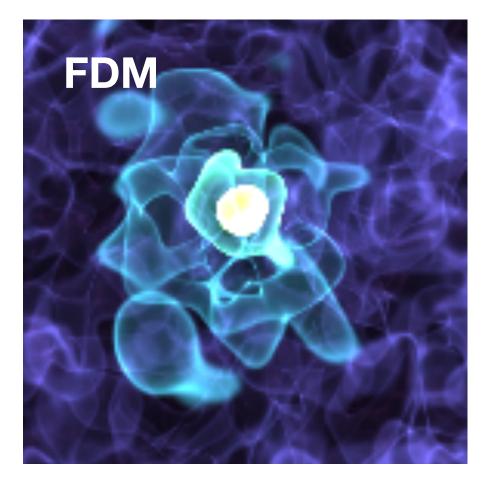
- Cusps in density profiles
- Very many small (sub)structures

Cores in density profiles induced by self scattering



• Cutoff at sub-galaxy scale in the power spectrum

Sterile neutrinos



Pattern induced by de Broglie length at sub-galactic scales

Ultralight bosons

output. Sman Scale Structure

Releve of public ndo et al.)

falling Shirasakie

et al.)

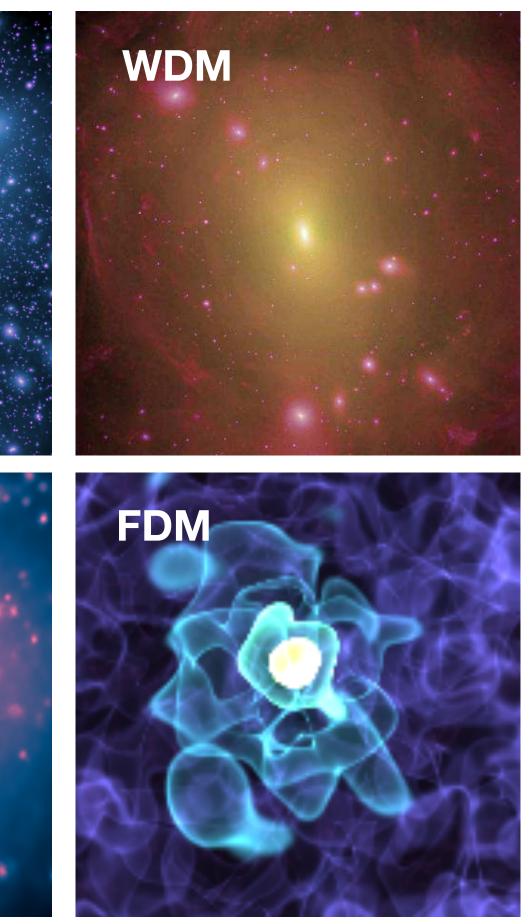
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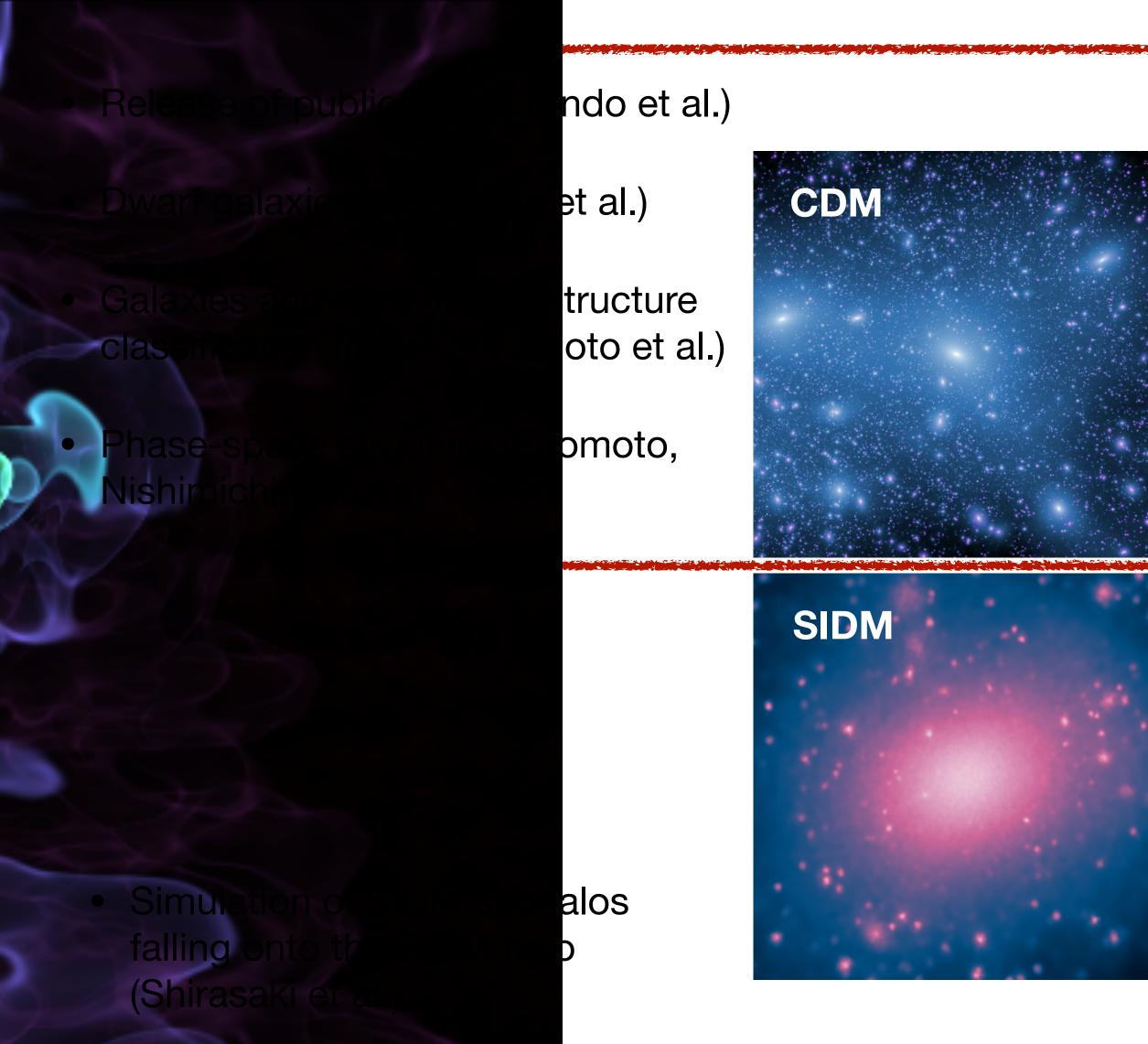


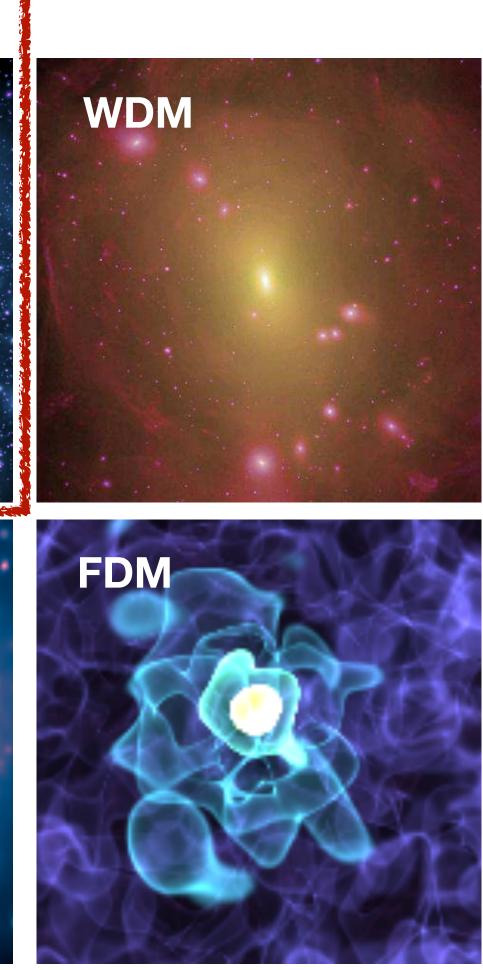
- Numerical simulations of WDM halos and subhalos (Okamoto, Inoue et al.)
- Developing semi-analytical models and constraints from satellite number counts (Ando et al.)

• Tight constraints using stellar motion in ultrafaint dwarf galaxies (Dalal et al.)



output. Sman Scale Structure





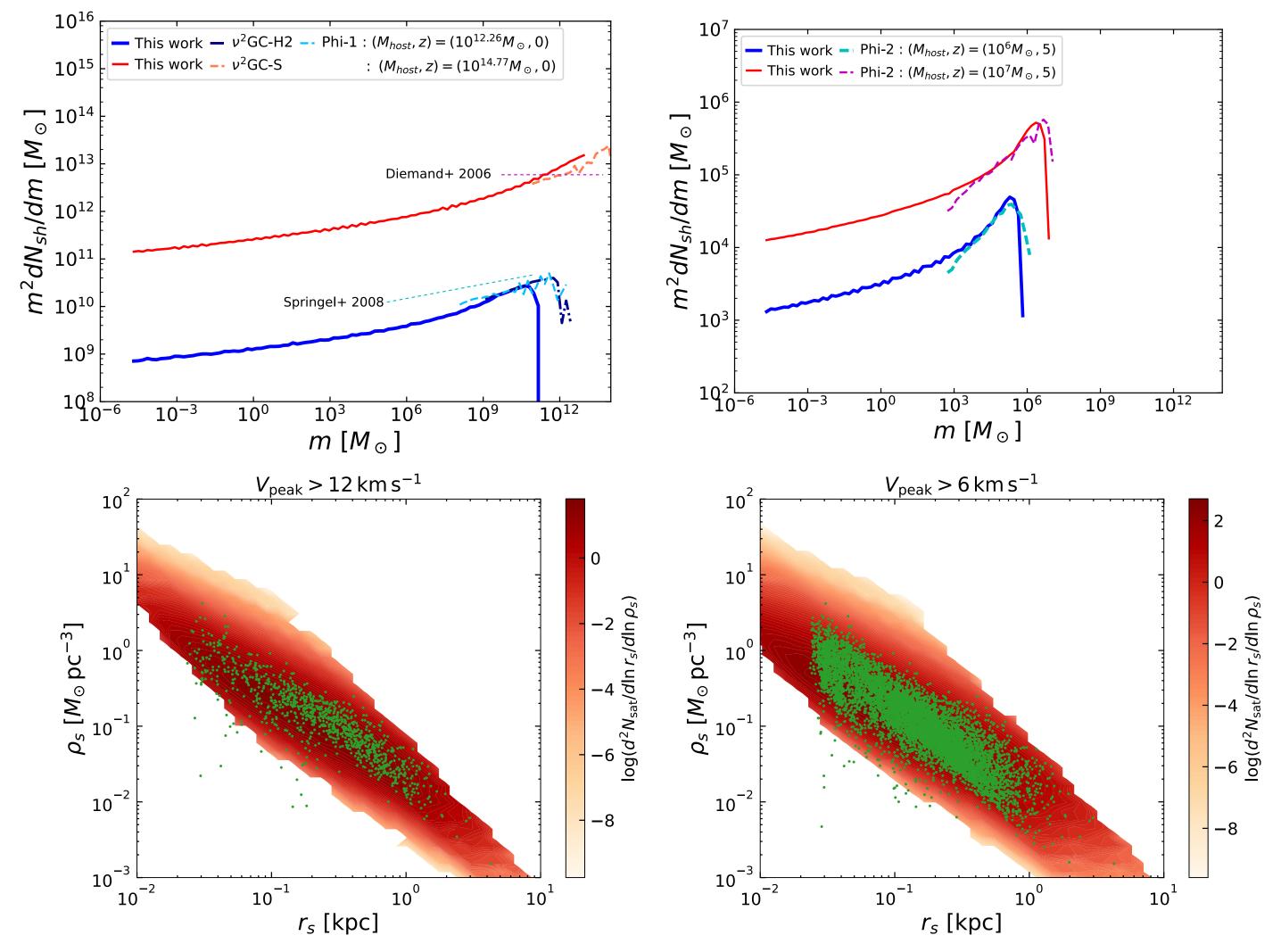
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Release of public codes for semianalytical subhalo models (CDM)

Hiroshima, Ando, Ishiyama, Phys. Rev. D 97, 123002 (2018)



Ando, Geringer-Sameth, Hiroshima, Hoof, Trotta, Walker, Phys. Rev. D 102, 061302 (2020)





N. Hiroshima

T. Ishiyama

- Recap: Semi-analytical models combining the extended Press-Schechter formalism with tidalevolution prescription
- Well recovers subhalo mass function and distribution of density profile parameters
- Cost-effective, free from numerical resolution and Poisson noise

Release of public codes for semianalytical subhalo models (CDM)

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Modeling for CDM (SASHIMI-C) arXiv 1803.07691 arXiv 1903.11427 The codes allow to calculate various subhalo properties efficiently using semi- analytical models for cold dark matter (CDM). The results are well in agreement with those from numerical N-body simulations. Authors					Packages No packages published Publish your first package Languages Jupyter Notebook 80.0% Python 20.0%	
	'ichiro Ando sa Hiroshima					
-	ne Dekker					
Special th		oaki Ishiyama, who providec used for calibration of mode	ed data of cosmological N-body el output.			
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- It can be used to quickly compute the subhalo mass function without making any assumptions such as power-law functional forms, etc. Only power law that we assume here is the one for primordial power spectrum predicted by inflation! Everything else is calculated theoretically.
- SASHIMI is not limited to numerical resolution which is often the most crucial





N. Hiroshima

A. Dekker

Semi-Analytical SubHalo Inference ModelIng

"Cold" SASHIMI: <u>github.com/shinichiroando/sashimi-c</u>

Only 760 lines of simple python codes, which enable to calculate (nearly) everything we did in Hiroshima et al. (2018)

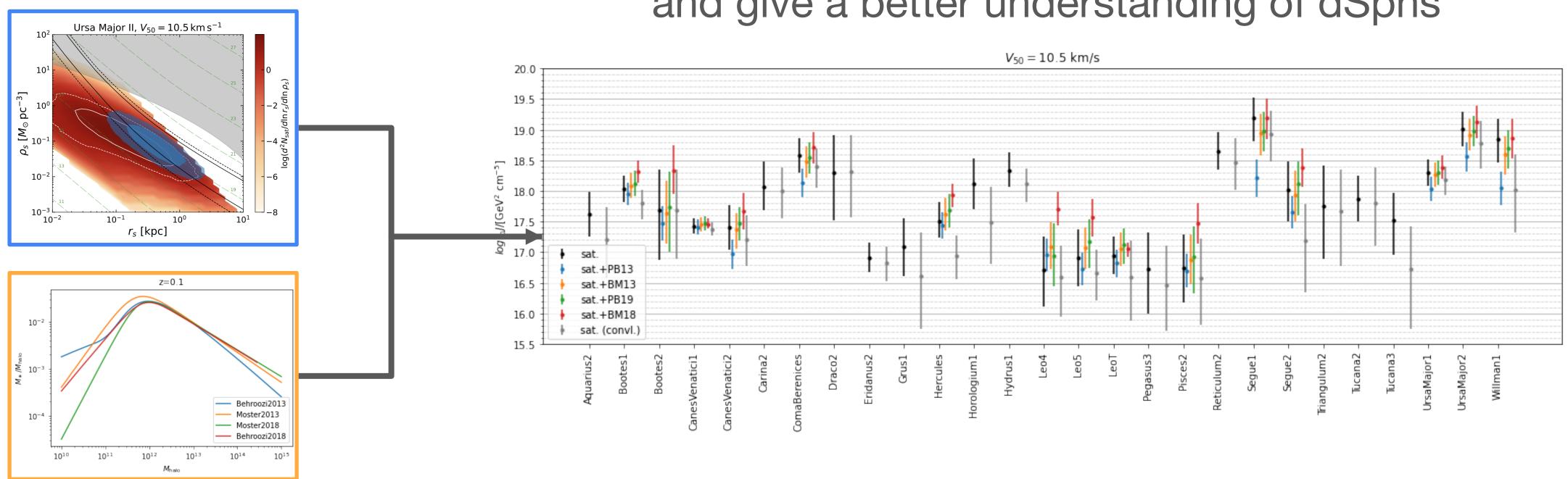
• Subhalo mass function, substructure boost of dark matter annihilation, etc.

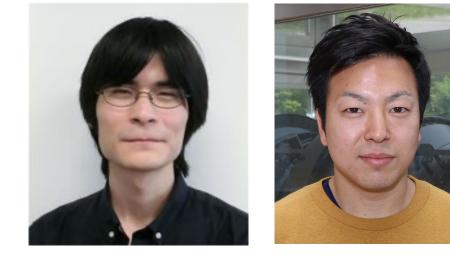
Well documented and useful sample codes provided



Cosmological prior for the J-factor estimation of dwarf spheroidal galaxies

- their dark matter halo profiles have large uncertainties
- - \bigcirc
 - Stellar-to-halo mass relation prior: empirical relation between stellar mass and halo mass \bigcirc







Dwarf spheroidal galaxies (dSph) play important roles for dark matter detection but

For the halo profile estimation of dSphs, we apply two cosmological priors:

Satellite prior: constraint distribution of halo parameter based on a structure formation model

The cosmological priors are useful to decrease the uncertainty in the estimation

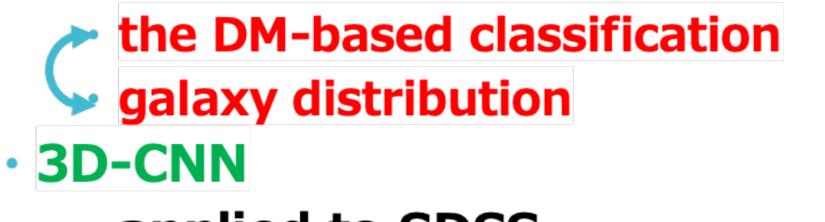
and give a better understanding of dSphs

Classification of cosmic structures for galaxies with machine learning: connecting cosmological simulations with observations

Shigeki Inoue, Xiaotian Si, Takashi Okamoto & Moka Nishigaki

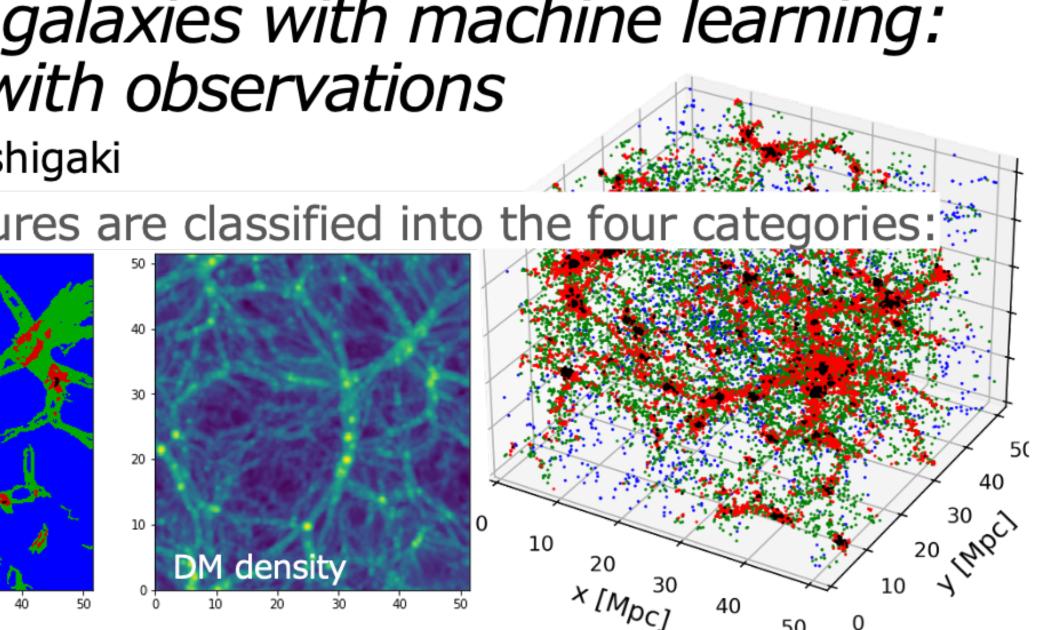
- In theory, using DM density, large-scale structures are classified into the four categories:
 - Knot, Filament, Sheet and Void
- However, since DM is unobservable Observations use galaxy distribution
- Thus, the structure classification is inconsistent between theory and observation.
- This study construct a machine-learning model to solve the problem.
 - Learning data of a cosmological simulation including baryons to learn the relationship between

[degree]



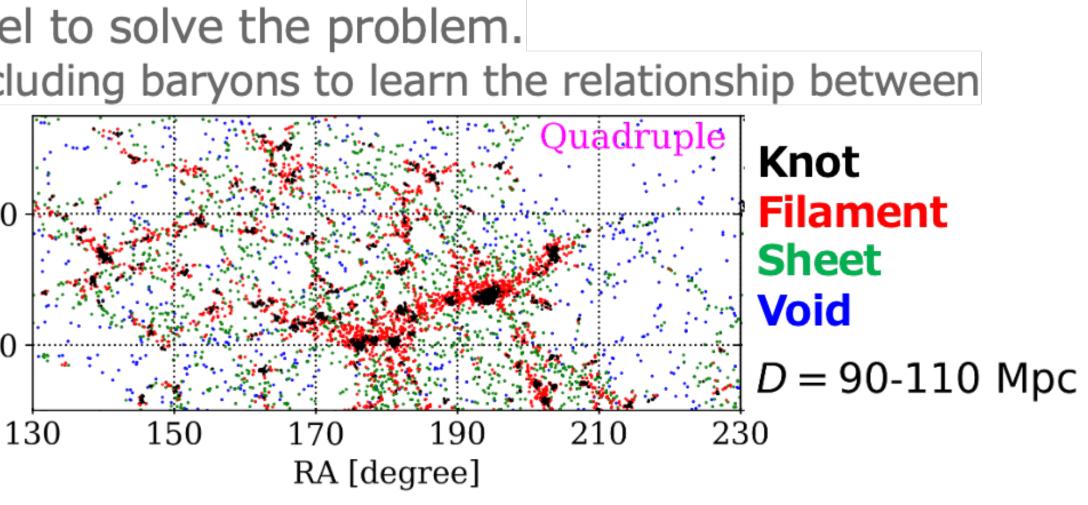
applied to SDSS

Observed galaxies are classified using DM!!









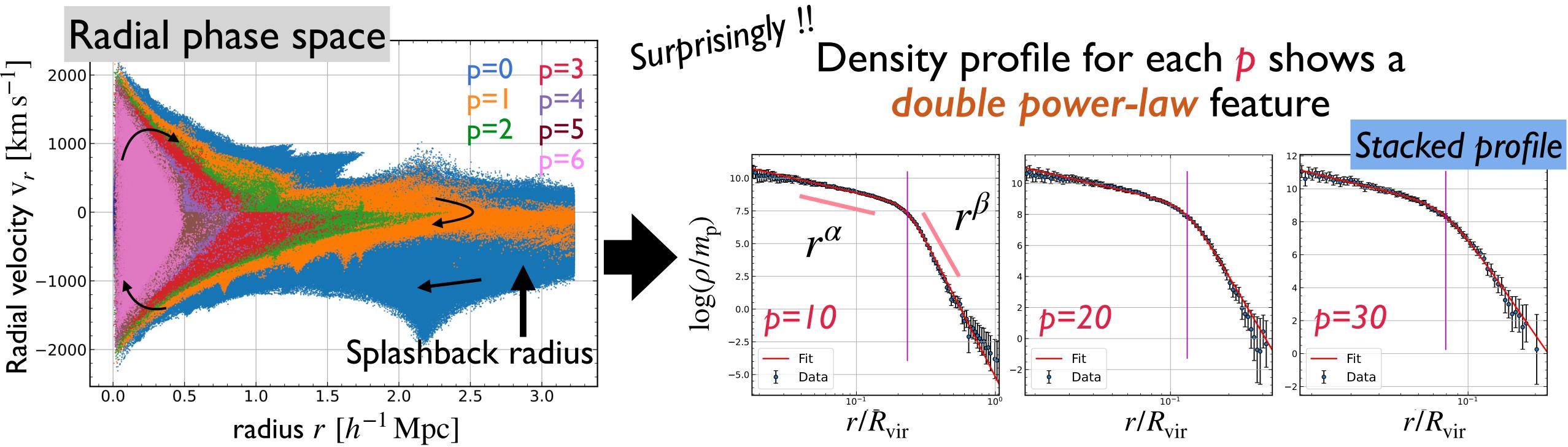
S. Inoue

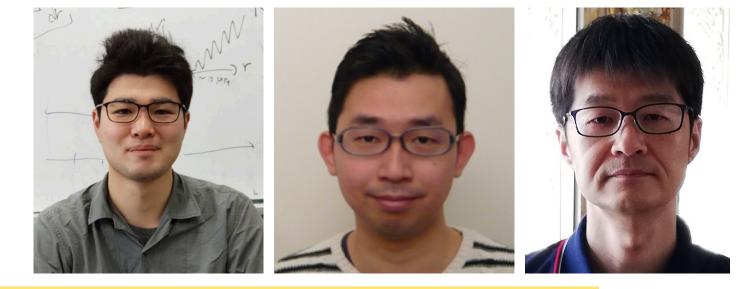


Inner structure of cold dark matter (CDM) halo

Enomoto, Nishimichi & Taruya ('22, in prep)

Applying the improved version of the method developed by Sugiura et al. ('20),

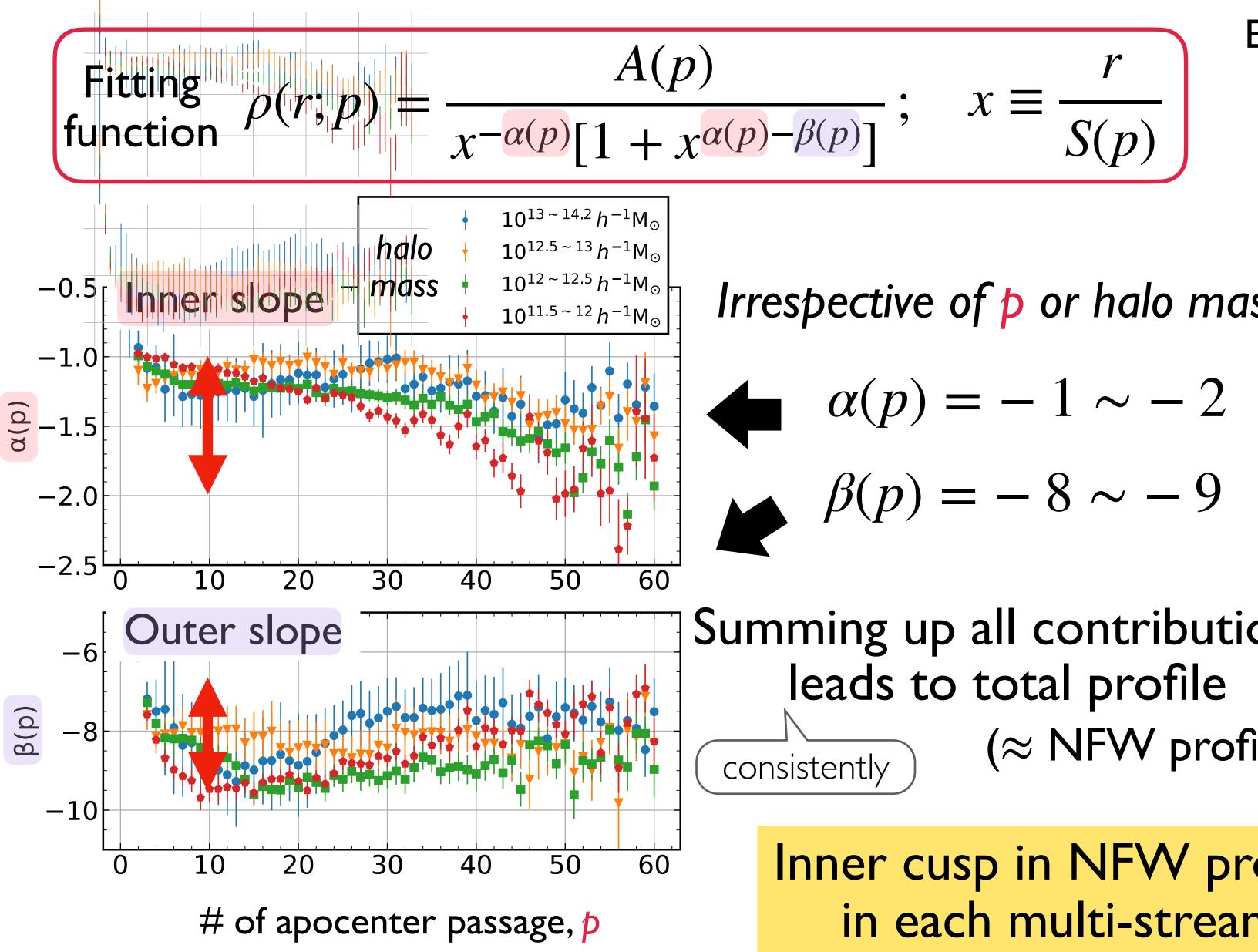


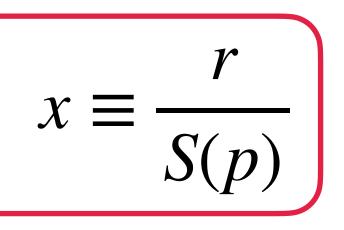


- Phase-space structure of dark matter halos by tracking particle trajectories in cosmological N-body simulations \rightarrow a clue to clarify nature of dark matter
 - Up to $p \sim 60$
 - N-body particles in CDM halos are classified by <u># of apocenter passages</u>, p



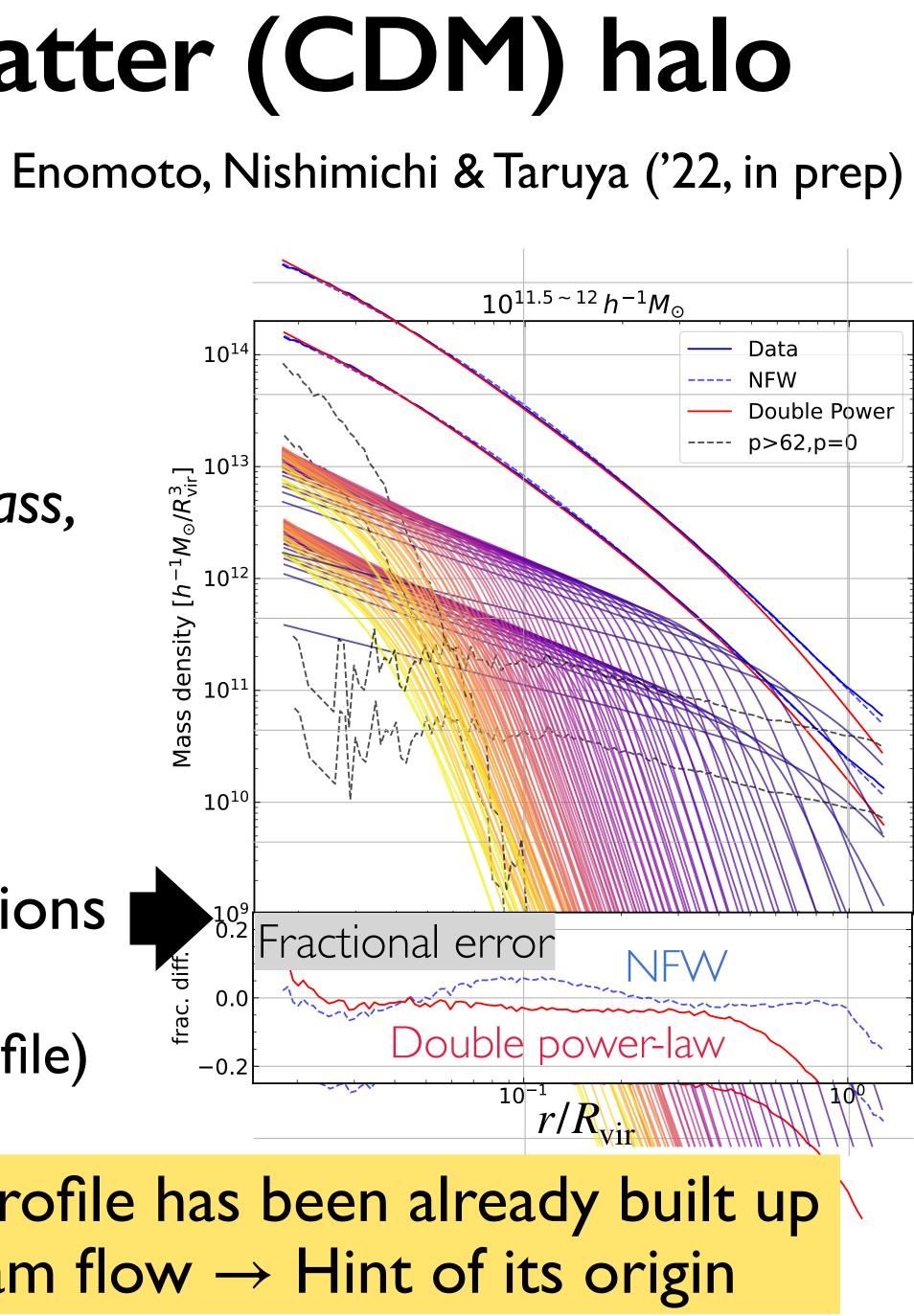
Inner structure of cold dark matter (CDM) halo





Irrespective of p or halo mass,

Summing up all contributions | leads to total profile $(\approx NFW \text{ profile})$



Inner cusp in NFW profile has been already built up in each multi-stream flow \rightarrow Hint of its origin

output. Sman Scale Structure

Release of public ndo et al.)

falling Shirasakie

et al.)

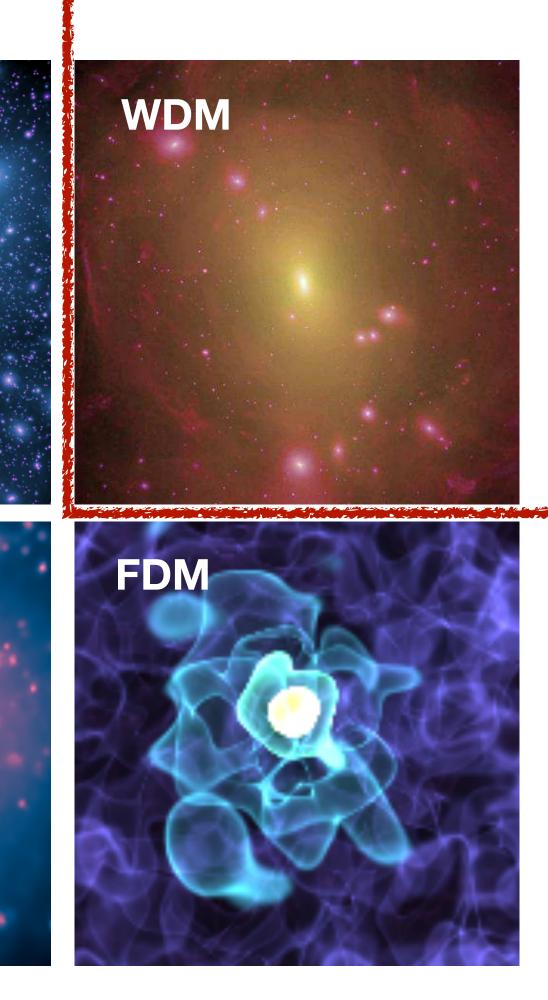
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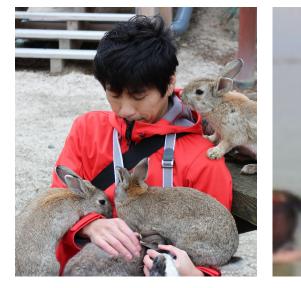
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Warm Dark Matter simulations

- A parameterization of the WDM mass function relative to that of the CDM is needed
- Lovell (2020) gives such a fitting function with WDM simulations of the half-mode mass of $M_{\rm hm} = 5 \times 10^8 3.5 \times 10^9 M_{\odot}$.
- We first extend their work to a wider range of the WDM mass $m_{WDM} = 1$ - 10 keV ($M_{\rm hm} = 6.4 \times 10^6 - 1.3 \times 10^{10} M_{\odot}$) and benchmark their fitting function
- Then construct a theoretical model that explains the simulation results.



T. Okamoto S.



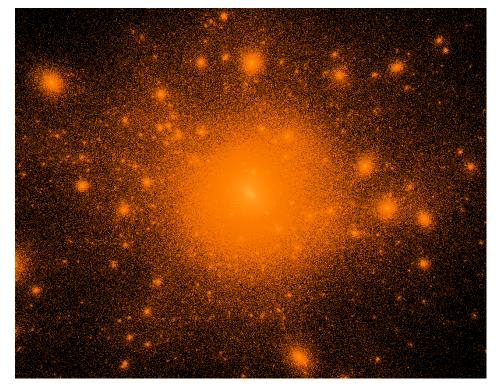
S. Inoue

Simulations

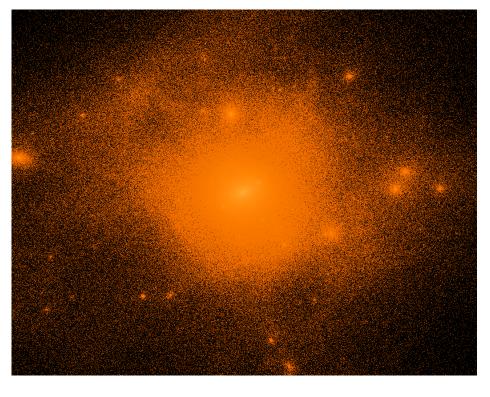
- Zoom simulations of a Milky Way mass halo lacksquare
 - $M_{200} = 1.18 \times 10^{12} M_{\odot}$
 - Mass resolution: $m_{\rm DM} = 2.491 \times 10^3 M_{\odot}$
 - The mean particle separation: $3.05 \times 10^{-3} h^{-1}$ cMpc ullet

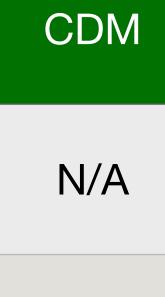
	1 keV	3 keV	10 keV
$M_{\rm hm}~(h^{-1}~M_{\odot})$	1.37x10 ¹⁰	3.52x10 ⁸	6.40x10 ⁶
$\lambda_{\rm fs}~(h^{-1}~Mpc)$	0.048	0.014	0.0037
Status	Running	if necessary	Done

CDM (low-resolution)

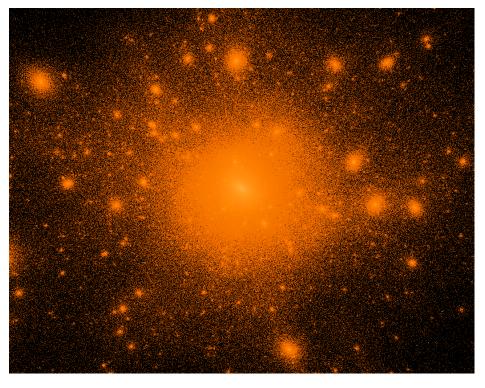


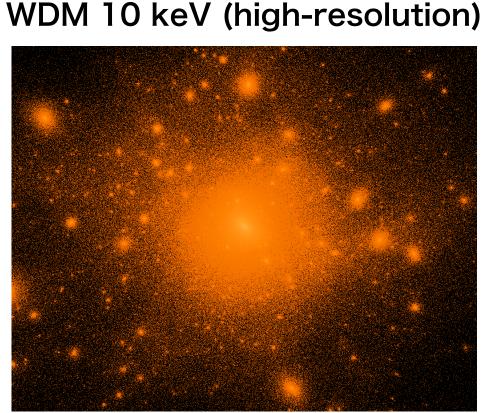
WDM 1 keV (low-resolution)





CDM (high-resolution)





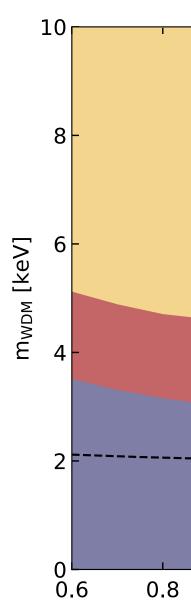
Done

N/A



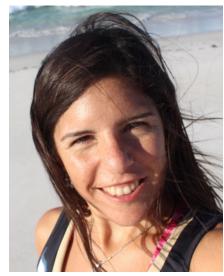
Semi-analytical models

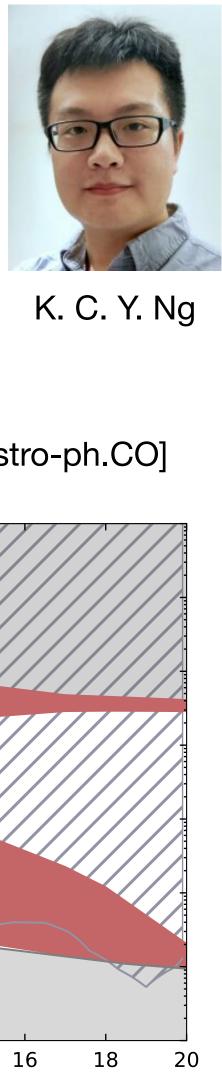
- "Warm" SASHIMI (*github.com*/ shinichiroando/sashimi-w)
 - Applied SASHIMI codes to the case of WDM by modifying power spectrum, etc.
- Compare with satellite number counts (DES+PanSTARRS1)
 - Excluding WDM mass of < 3.6-5.1 keV (without baryon physics) uncertainties)
 - Excluding sterile neutrino dark matter (combined with X rays)

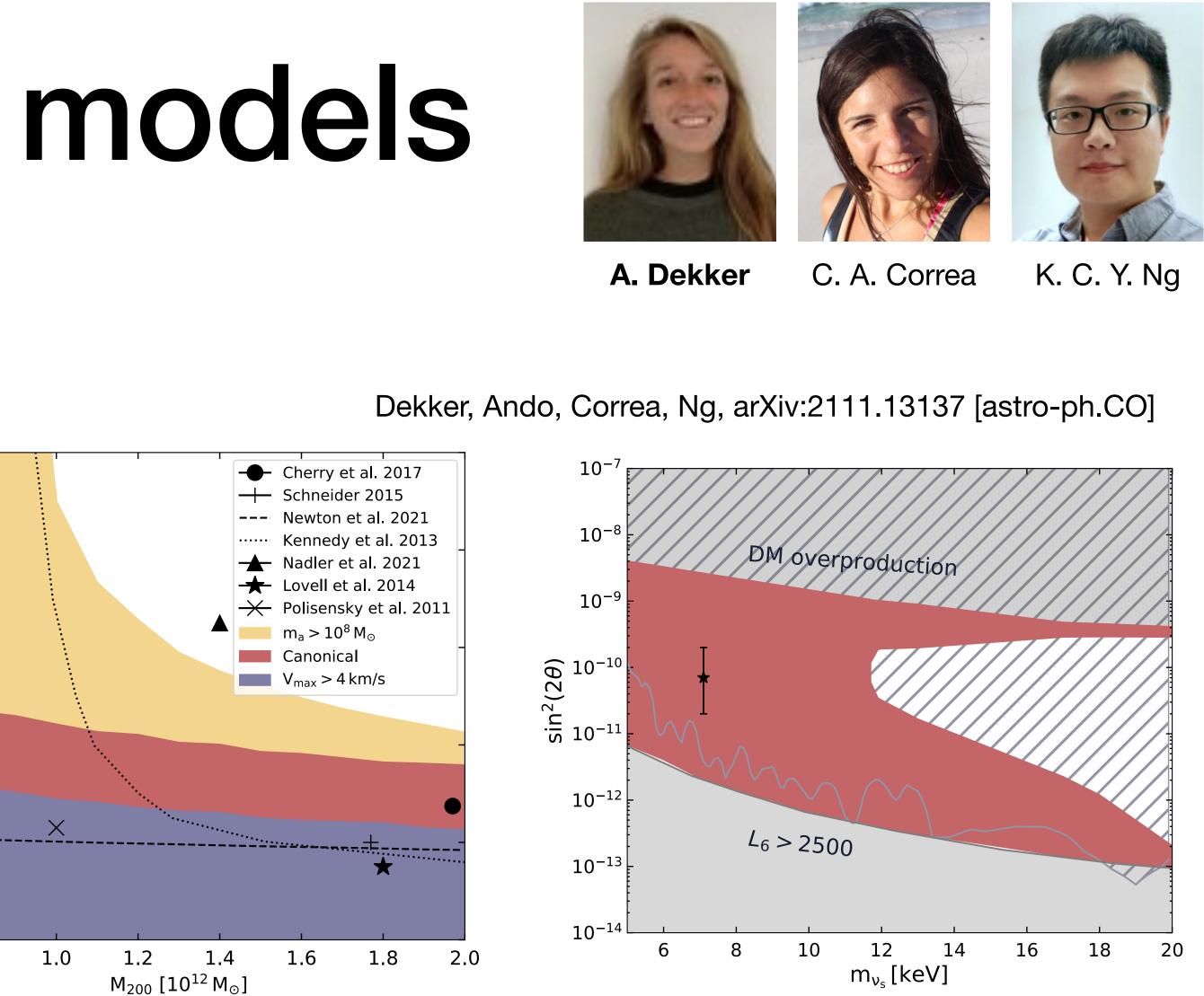


This needs to be calibrated against more numerical simulations!



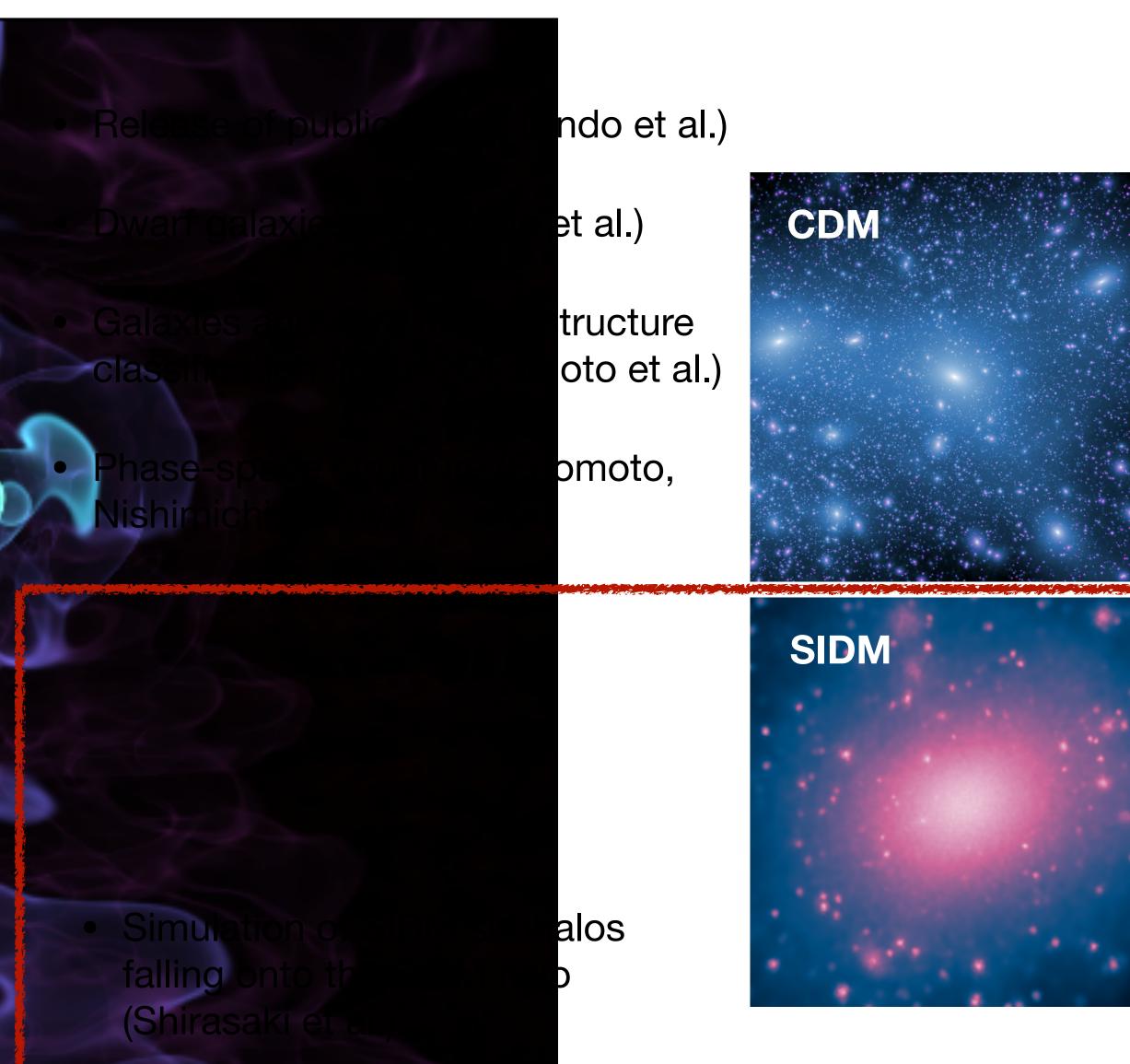


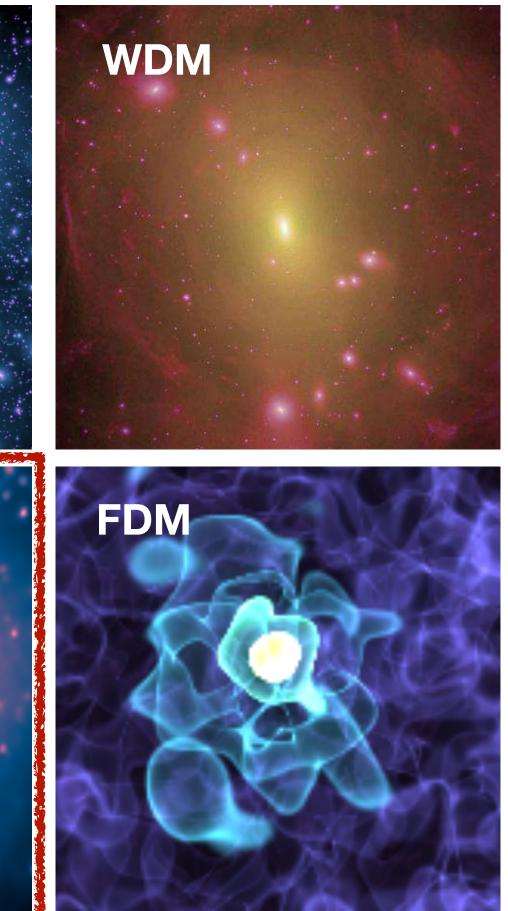






output. Sman Scale Structure





- Numerical simulations of WDM halos and subhalos (Okamoto, Inoue et al.)
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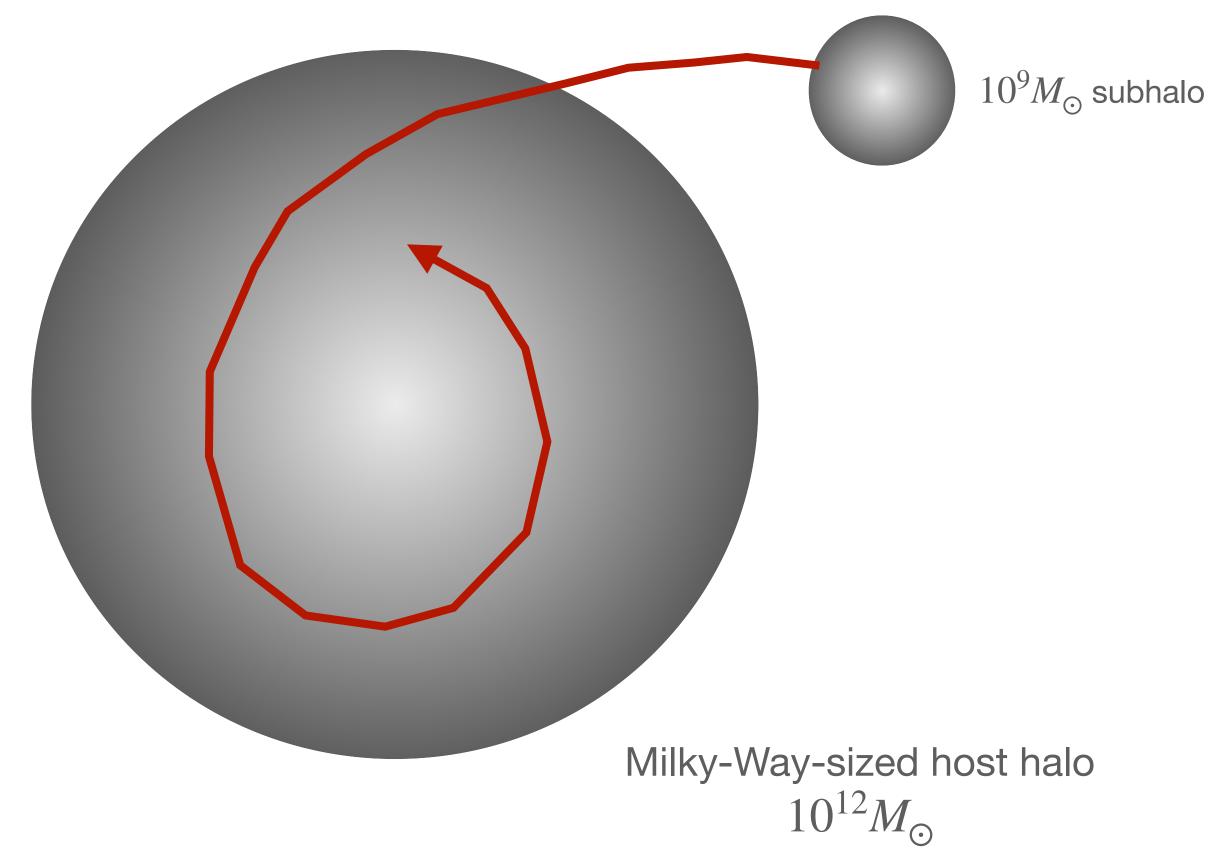
• Tight constraints using stellar motion in ultrafaint dwarf galaxies (Dalal et al.)



Semi-analytic model of SIDM subhaloes **Calibration with ideal N-body simulations of minor mergers**

- Testing self-interactions of DM particles would require a precise modeling of
 - thermalization of SIDM halo/ subhalo
 - Tidal stripping / Ram pressure
- Develop a semi-analytic model of infalling subhalos to a MW-sized halo and calibrate it with (isolated) N-body simulations





A brief summary of our model

Gravothermal fluid model (e.g. Balberg+2002)

Host halo density $\rho_{\rm h}(r,t)$

Mass conservation $dM/dr = 4\pi r^2 \rho_{\rm h}(r)$

Heat Flux = $-\kappa (m/k_B) \partial \sigma_v^2 / \partial r$ re-arranges ρ_h and σ_v

Subhalo

Density $\rho_{\rm sub}(r,t)$

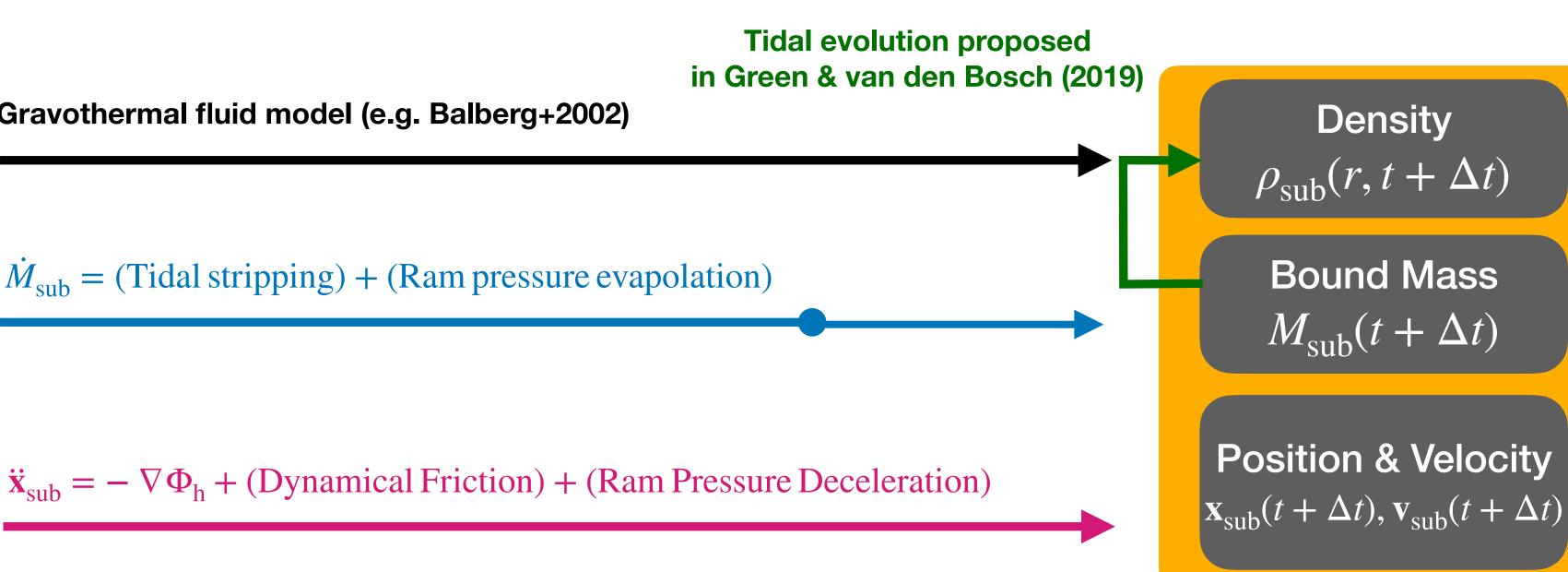
Bound Mass $M_{\rm sub}(t)$

Position & Velocity $\mathbf{X}_{sub}(t), \mathbf{V}_{sub}(t)$ Gravothermal fluid model (e.g. Balberg+2002)

 $\dot{M}_{sub} = (Tidal stripping) + (Ram pressure evapolation)$

Hydrostatic equilibrium $d(\rho \sigma_v^2)/dr = -GM\rho/r^2$

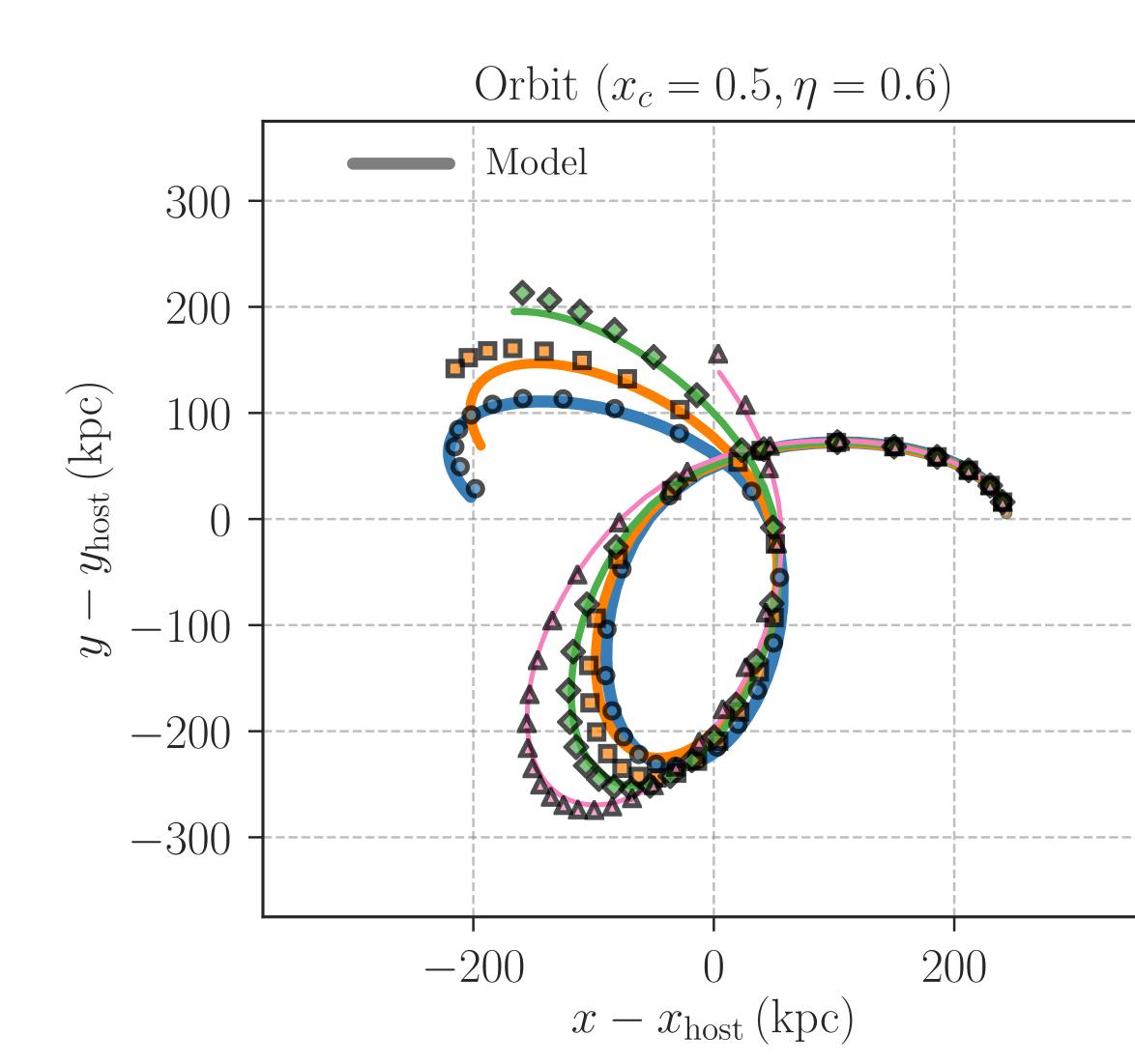


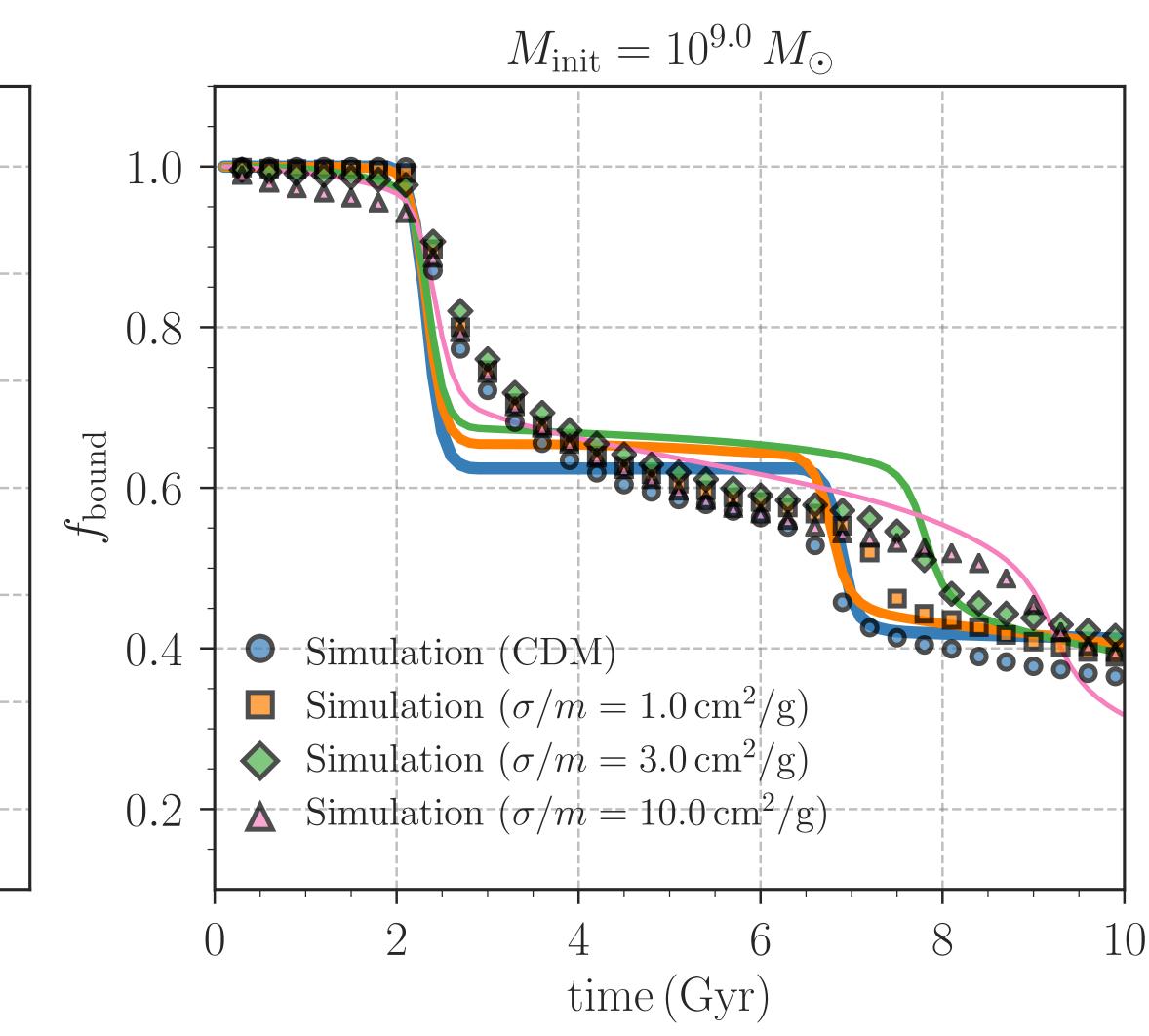


Note: We ignore possible changes of subhalo density profiles due to ram pressure effects



Comparison with our model and simulations







output. on an scale structure

Releve of public ndo et al.)

falling onto Shirasaki et

et al.)

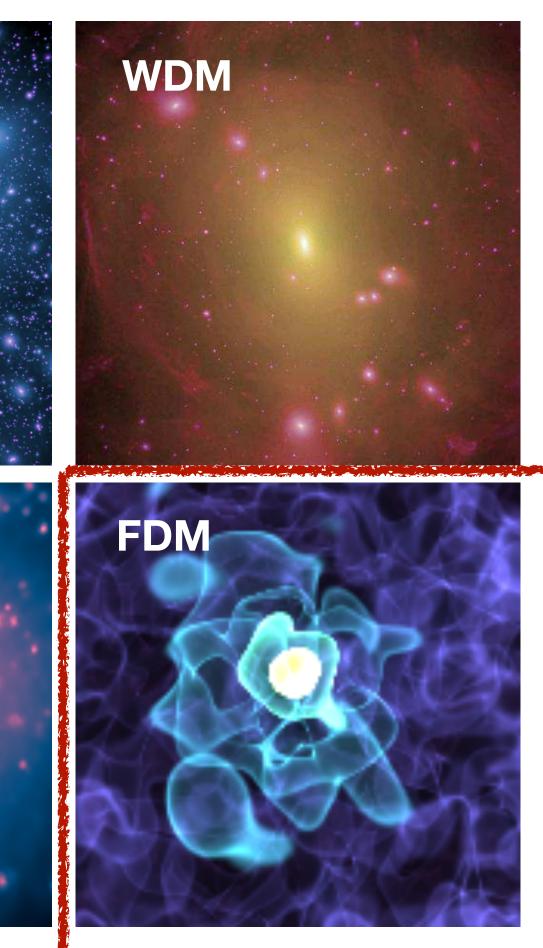
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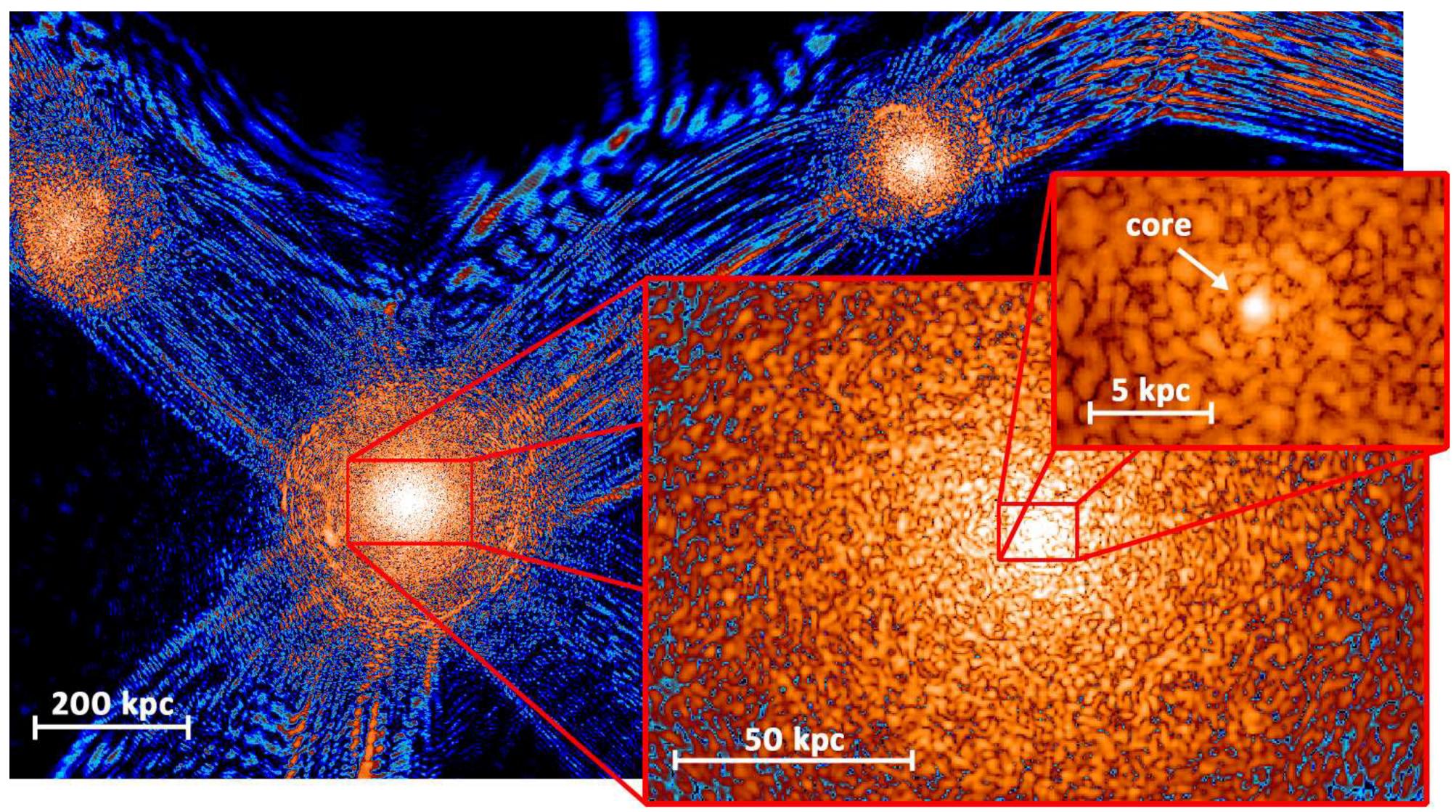
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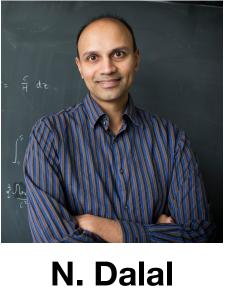




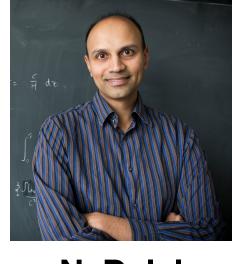
FDM wave interference



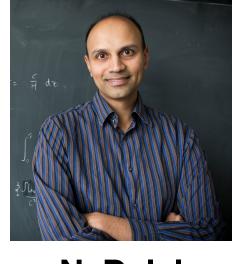
Schive et al., Nature Physics, 10, 496 (2014)



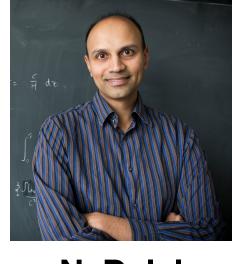
- Interference fringes have density contrast $\delta \rho \sim \rho$ everywhere all of the time
- These lead to fluctuating gravitational forces that can perturb stars



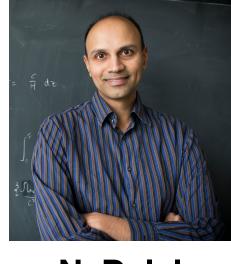
- Interference fringes have density contrast $\delta \rho \sim \rho$ everywhere all of the time
- These lead to fluctuating gravitational forces that can perturb stars
- Where to look for this signature of FDM? Crude estimate:
 - $\delta M \sim \delta \rho \lambda^3 \propto \rho / \sigma_v^3 \Rightarrow$ acceleration perturbation $\delta a \sim G \delta M / \lambda^2 \propto G \rho / \sigma_v$



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 - At that location, enclosed mass $M \sim \rho R^3$, so $a \sim GM/R^2 \propto G\rho R^3$

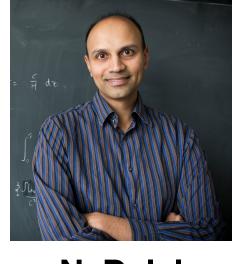


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 - At that location, enclosed mass N
 - So fractional effect $\delta a/a \propto (R \sigma_v)$



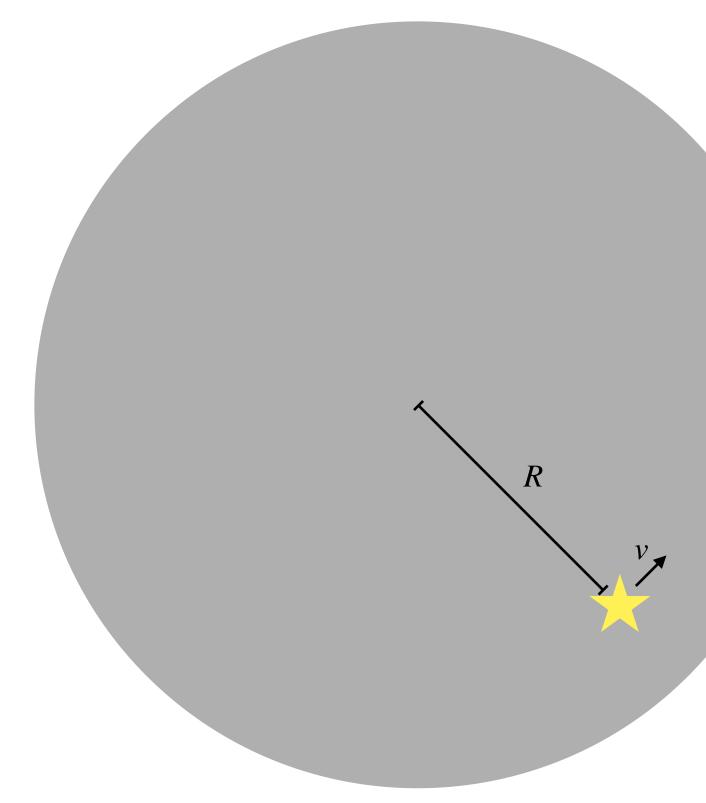
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 - At that location, enclosed mass N
 - So fractional effect $\delta a/a \propto (R \sigma_{\nu})$
- Biggest effect where R is small and σ_v is small, i.e. centres of smallest halos -> ultrafaint dwarf galaxies



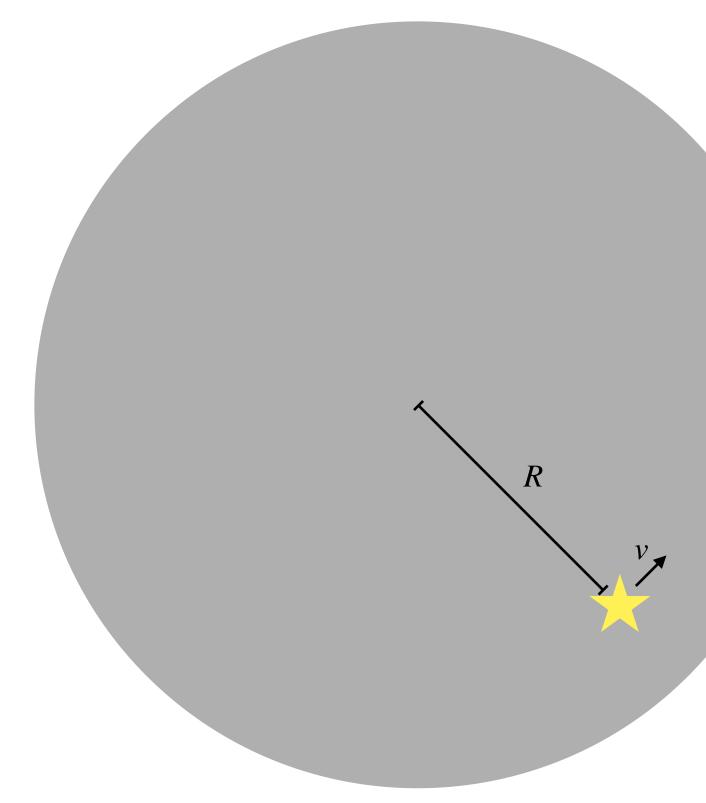
$$M \sim \rho R^3$$
, so $a \sim GM/R^2 \propto G\rho R$

• Consider typical star in galaxy of size R, moving at velocity $v \sim \sigma_v$



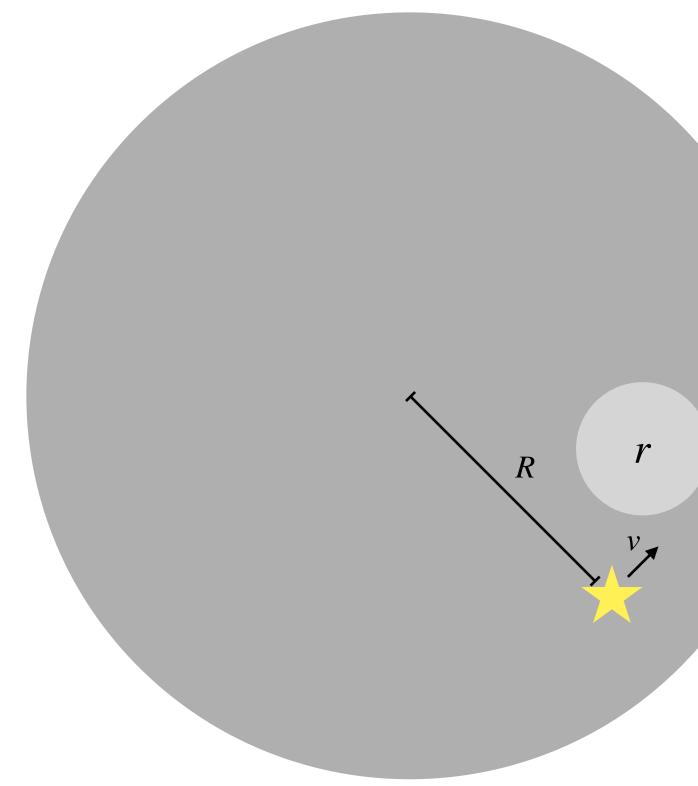


- Consider typical star in galaxy of size R, moving at velocity $v \sim \sigma_v$
- Enclosed mass is $M \sim 3\sigma_v^2 R/G$





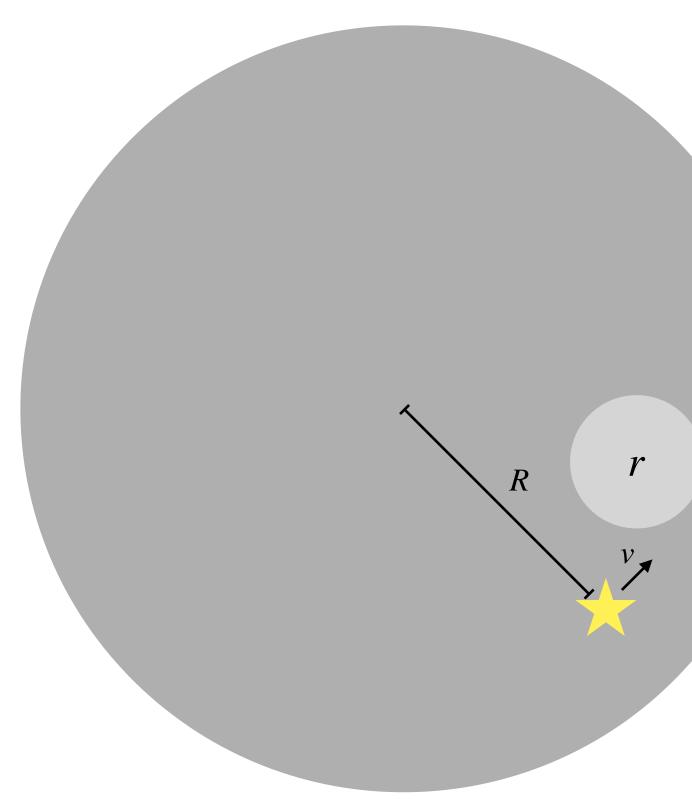
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- FDM fluctuation of size *r*





- Consider typical star in galaxy of size R, moving at velocity $v \sim \sigma_v$
- Enclosed mass is $M \sim 3\sigma_v^2 R/G$
- FDM fluctuation of size r
 - $\delta M \sim (r/R)^3 M$, $\delta \Phi \sim G \delta M/r \approx 3\sigma_v^2 (r/R)^2$

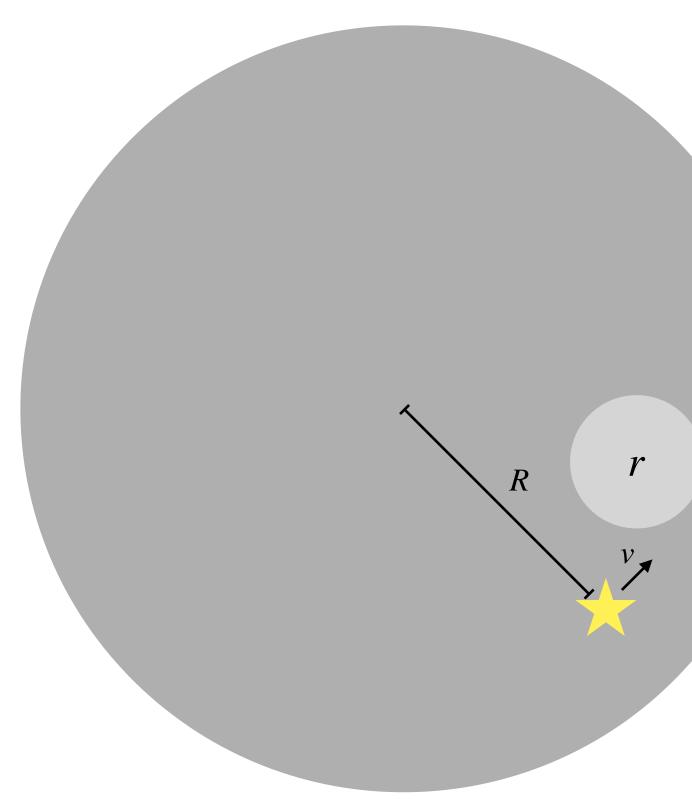






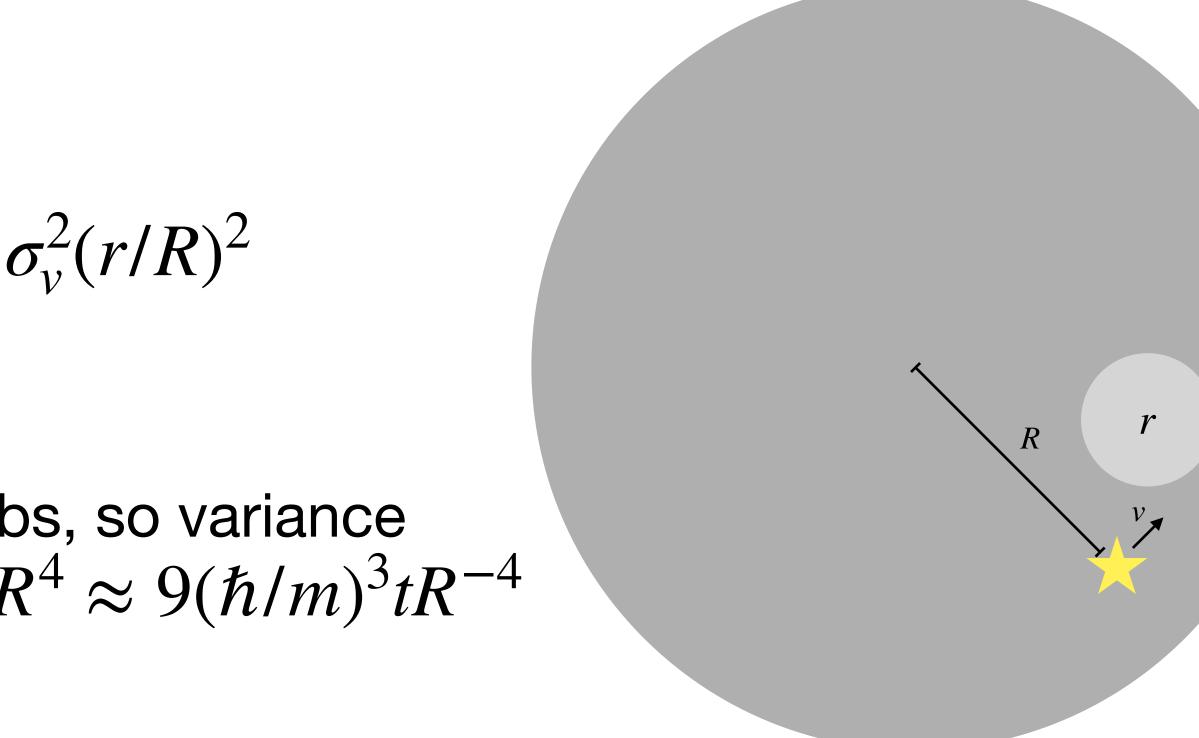
- Consider typical star in galaxy of size R, moving at velocity $v \sim \sigma_v$
- Enclosed mass is $M \sim 3\sigma_v^2 R/G$
- FDM fluctuation of size r
 - $\delta M \sim (r/R)^3 M$, $\delta \Phi \sim G \delta M/r \approx 3\sigma_v^2 (r/R)^2$
 - $\delta v \sim \delta \Phi / v \approx 3\sigma_v (r/R)^2$







- Consider typical star in galaxy of size R, moving at velocity $v \sim \sigma_v$
- Enclosed mass is $M \sim 3\sigma_v^2 R/G$
- FDM fluctuation of size r
 - $\delta M \sim (r/R)^3 M$, $\delta \Phi \sim G \delta M/r \approx 3\sigma_v^2 (r/R)^2$
 - $\delta v \sim \delta \Phi / v \approx 3\sigma_v (r/R)^2$
- In time t, star encounters $N \sim vt/r$ blobs, so variance increases by $\delta \sigma_v^2 \approx N(\delta v)^2 \approx 9\sigma_v^3 t r^3 / R^4 \approx 9(\hbar/m)^3 t R^{-4}$





- Consider typical star in galaxy of size R, moving at velocity $v \sim \sigma_v$
- Enclosed mass is $M \sim 3\sigma_v^2 R/G$
- FDM fluctuation of size r
 - $\delta M \sim (r/R)^3 M$, $\delta \Phi \sim G \delta M/r \approx 30$
 - $\delta v \sim \delta \Phi / v \approx 3\sigma_v (r/R)^2$
- In time t, star encounters $N \sim vt/r$ blobs, so variance increases by $\delta \sigma_v^2 \approx N(\delta v)^2 \approx 9\sigma_v^3 tr^3/R^4 \approx 9(\hbar/m)^3 tR^{-4}$
- So we can solve for mass *m* that makes $\delta \sigma_v^2 \approx \sigma_v^2$ in time *t*. Plugging in t = 10 Gyr, R = 50 pc, $\sigma_v = 3$ km/s gives $m \sim 10^{-19} \, \mathrm{eV}$

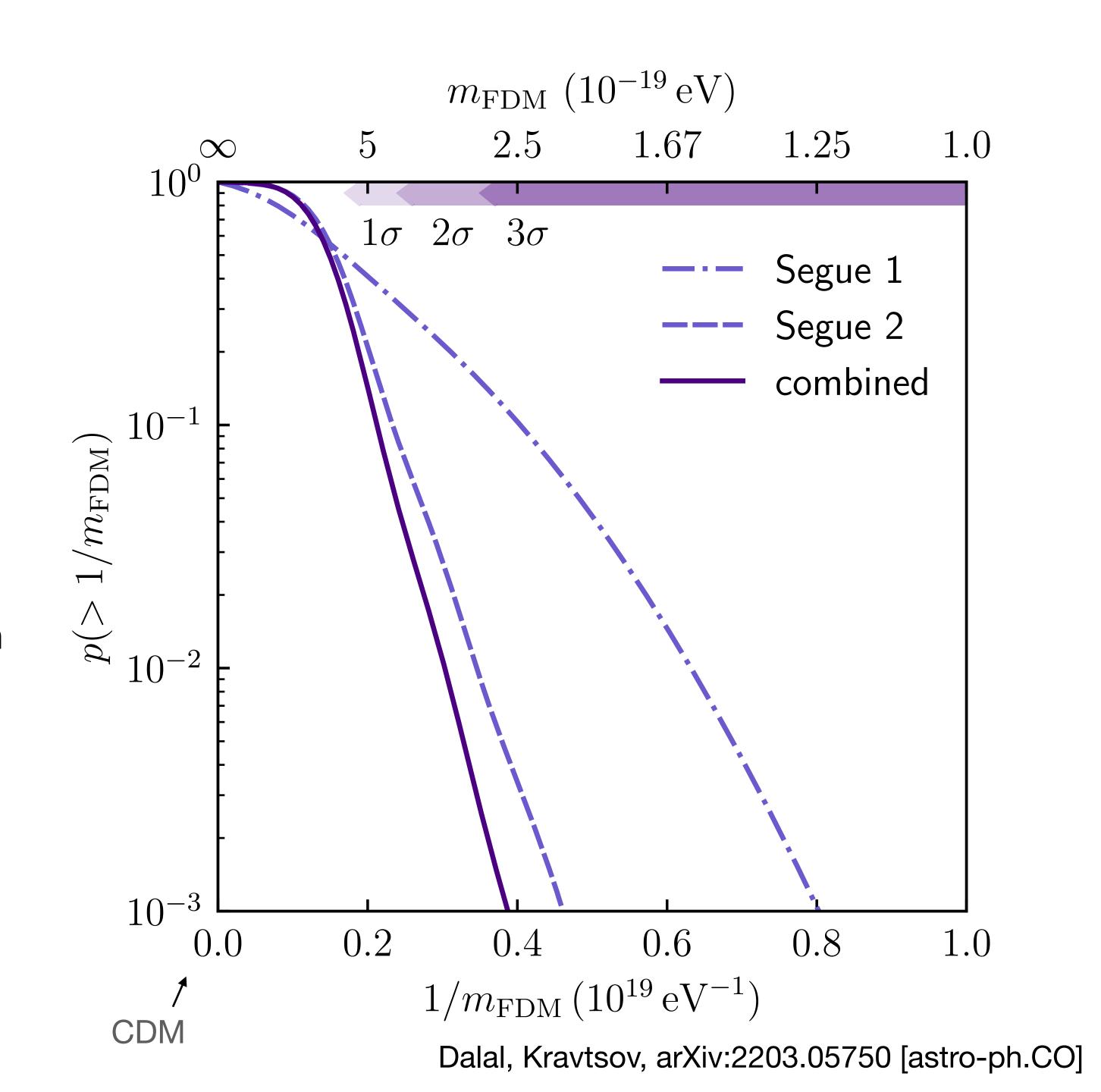
$$\sigma_v^2(r/R)^2$$



Results

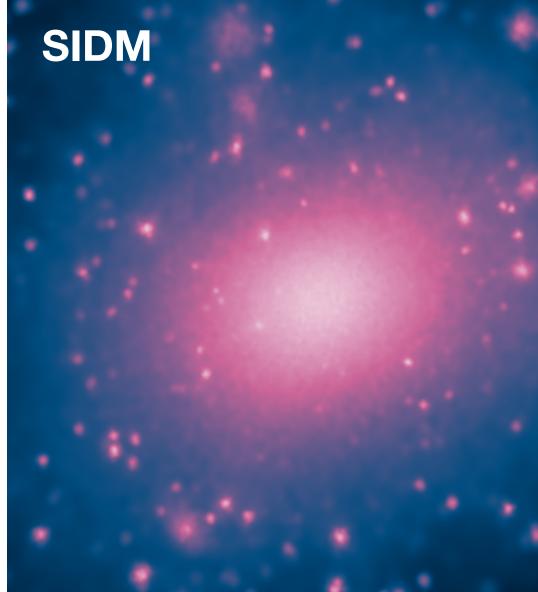
- Find $m_{\rm FDM} > 3 \times 10^{-19} \,{\rm eV}$ at >99% confidence, using Segue 1 & Segue 2. Previous bounds from Ly- α forests are $m \gtrsim 10^{-21} \,{\rm eV}$
- Essentially, rules out "fuzzy" regime:
 - linear power spectrum identical to LCDM out to $k \sim 200 \, {\rm Mpc}^{-1}$
 - halo mass function identical to LCDM down to $M\sim 2\cdot 10^5 M_\odot$

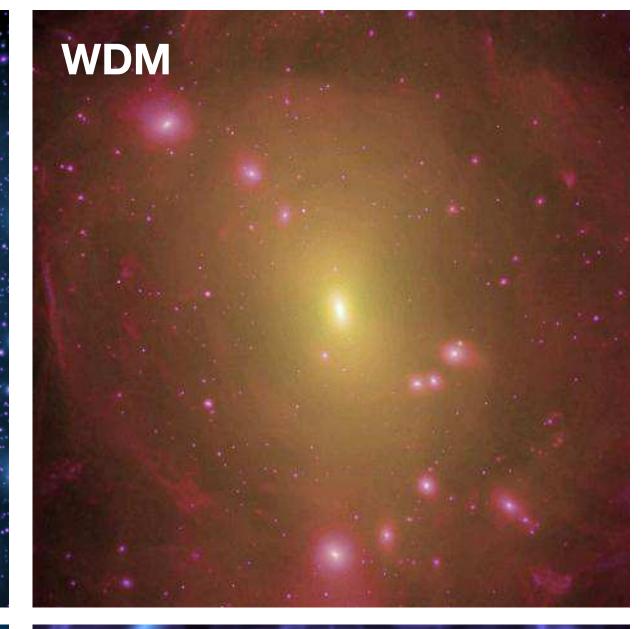
 Constraints of similar strength are obtained by a dynamical analysis of ultrafaint dwarf galaxies (Hayashi et al. 2021)

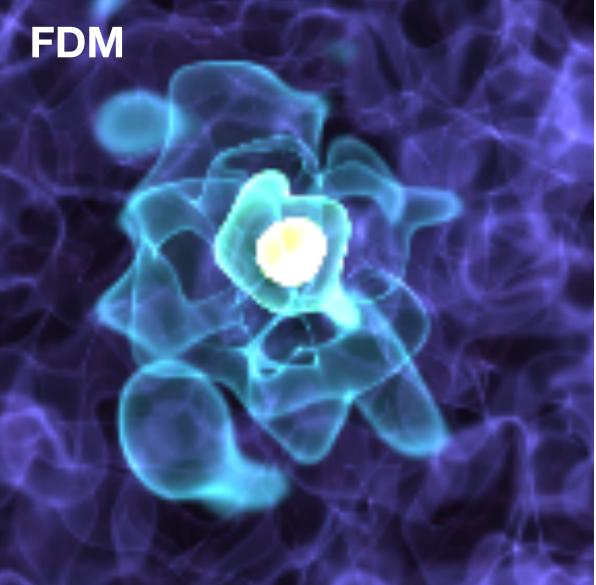


Conclusions









- FDM)
- Results worth highlighting:
 - Strong constraints on WDM and sterile neutrinos using satellite counts (and public codes SASHIMI)
 - Strong constraints on FDM using stellar velocity data in ultrafaint galaxies
- Exploration of CDM phase space, SIDM modeling, and numerical **simulations** is ongoing with exciting results
- More detailed talks by Horigome and Inoue
- Looking forward to collaborations with other groups for FY 2022!

Good progress is being made in all the four directions (CDM, WDM, SIDM,