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Cosmological 21cm line observations to test scenarios of super-Eddington accretion on to seed BHs of high-z SMBHs

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理論センタ

Introduction

• We do not know origins of Super-Massive Black Holes

 $10^{15} M_{\odot}$ to be PBH Dark matter $10^{9} M_{\odot}$ observed at z=7.642 (PBH origins are excluded)

 We need seed (primordial) BHs before z >> 7 which had evolved to the SMBHs through the (super-) Eddington accretion

 M_{sBH} = 10² M_{\odot} at z~30 \rightarrow accretions \rightarrow 10⁹ M_{\odot} at z=7

 $\Omega_{\rm sBH}/\Omega_{\rm CDM} \sim 10^{-10} \left(\frac{n_{\rm seed,0}}{10^{-3} {\rm Mpc}^{-3}}\right) \left(\frac{M_{\rm BH,ini}}{10^2 M_{\odot}}\right) \left(\frac{M_{\rm SMBH}}{10^9 M_{\odot}}\right) \left(\frac{M_{\rm gal}}{10^{12} M_{\odot}}\right)^{-1}$

 By 21cm data by EDGES, we can obtain upper bounds on emissions from accretions on to seed BHs and exclude the seed masses at z=17

 $M_{\rm BH,ini} \gtrsim 10^2 M_{\odot}$ for $n_{\rm seed}(z=0) = 10^{-3} Mpc^{-3}$

Origin of seed BHs with $M_{sBH} = 10^2 - 10^6 M_{\odot}$ at z>>7

Astrophysical

Collapses of massive stars/gas clouds Mergers of massive stars/ black holes

• Primordial

Formation temp. $T_{form} \sim O(10) \text{ MeV} - O(0.1) \text{ MeV}$ Curvature pertrurbation $P_{\zeta} \sim 10^{-1.5}$





proton-electron's spin-spin interaction



Spin temperature Ts

 Defined by the ratio of the occupation numbers in two states

$$\frac{n_{upper}}{n_{lower}} = \frac{g_{upper}}{g_{lower}} Exp \left[-\frac{\Delta E}{T_s} \right]$$

 $\Delta E = 2\pi v_{21} = 5.8 \times 10^{-6} eV$

 g_i = degree of reedom for a level "i"

Cosmological 21cm emission line emitted at the reionization epoch



21cm absorption by EDGES

Judd D. Bowman, et al., Nature 555 (2018) 67 Steven R. Furlanetto et al., arXiv:1903.06212



 $T_{21cm} = -500^{+200}_{-500} \,\mathrm{mK}$ (99% CL) We can constrain any new heating mechanism at least such as accretions on to BHs or annihilating DMs at z~17

Impacts of new small-scale N-body simulations on dark matter annihilations

constrained from cosmological 21cm line observations

Nagisa Hiroshima, Kazunori Kohri, Toyokazu Sekiguchi, Ryuichi Takahashi, arXiv:2103.14810 [astro-ph.CO]



The Eddington limit in accretions



$$L_{\rm E} \sim \frac{G_{\rm Newton} m_{\rm proton} M_{\rm BH}}{\sigma_{\rm Thomson}} \simeq 1.3 \times 10^{38} \, {\rm erg \ sec^{-1}} \left(\frac{M_{\rm BH}}{M_{\odot}}\right)$$

The Super-Eddington accretion

Accretion rate in unit of the Eddington accretion

$$\dot{M}_{\rm crit} \equiv \eta_{\rm eff}^{-1} L_E \simeq 1.4 \times 10^{18} \text{ g sec}^{-1} \left(\frac{\eta_{\rm eff}^{-1}}{10}\right) \left(\frac{M_{\rm BH}}{M_{\odot}}\right)$$
$$\dot{m} = \frac{\dot{M}}{\dot{M}_{\rm crit}}$$

Mass evolutions in the Eddington accretion

$$M_{\rm BH}(t) \sim M_{\rm BH,ini} \exp\left(10\dot{m} \frac{t - t_{\rm ini}}{\tau_E}\right)$$

$$\tau_E \equiv \frac{M_{\rm BH}c^2}{L_E} = \frac{\sigma_T c}{4\pi\mu G m_p} \simeq 0.45 {\rm Gyr}.$$

Luminosity of accretion disks





Optically thick?

Xuelei Chen, Marc Kamionkowski, arXiv:astro-ph/0310473



Ionization fraction x_e and the gas temperature T_m

Ionization fraction

$$\frac{dx_e}{dt} = -C \left[\alpha_{\rm H}(T_m) x_e^2 n_H - \beta_{\rm H}(T_\gamma) (1 - x_e) e^{-E_\alpha/T_\gamma} \right] + \frac{dE_{\rm inj}}{dV dt} \frac{1}{n_{\rm H}} \left[\frac{f_{\rm ion}(t)}{E_0} + \frac{(1 - C) f_{\rm exc}(t)}{E_\alpha} \right],$$

$$C = \frac{\Lambda n_{\rm H} (1 - x_e) + \frac{1}{2\pi^2} E_{\alpha}^3 H(t)}{\Lambda n_{\rm H} (1 - x_e) + \frac{1}{2\pi^2} E_{\alpha}^3 H(t) + \beta_H n_H (1 - x_e)},$$

• Gas temperature

$$\frac{dT_m}{dt} = -2H(t)T_m + \Gamma_C(T_\gamma - T_m) + \frac{dE_{\text{inj}}}{dVdt} \frac{1}{n_{\text{H}}} \frac{2f_{\text{heat}}(z)}{3(1 + x_e + f_{\text{He}})}$$
$$T_{21\text{cm}}(z) = \frac{T_s(z) - T_\gamma(z)}{1 + z} \tau_{21\text{cm}}(z) \qquad \Gamma_C = \frac{8\sigma_T a_r T_\gamma^4}{3m_e} \frac{x_e}{1 + f_{\text{He}} + x_e}$$

Time evolutions of temperature

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Upper bounds on accretion rates on seed BHs at z=17 evolved to SMBHs until z=7

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Conclusion

- By the EDGES data, we can obtain upper bounds on accretion on to seed BHs, which evolved to high-z SMBHs
- We exclude the seed BHs with their masses

$$\begin{split} \mathrm{M}_{\mathrm{BH,ini}} \gtrsim 10^2 M_{\odot} & \text{for } n_{\mathrm{seed}}(z=0) = 10^{-3} \mathrm{Mpc}^{-3} \\ & \text{Number counts of SMBHs at z=0 (the strongest assumption)} \\ \mathrm{M}_{\mathrm{BH,ini}} \gtrsim 10^6 M_{\odot} & \text{for } n_{\mathrm{seed}}(z=0) = 10^{-7} \mathrm{Mpc}^{-3} \end{split}$$

Observations of SMBHs at high-redshift at z=6 (conservative)