Quantum Nucleation of Chiral Soliton Lattice

Based on : T. Higaki (Keio), KK, K. Nishimura (KEK), PRD 106 (2022) 096022, arXiv: 2207.00212 [hep-th] also see: M. Eto & M. Nitta, JHEP09 (2022) 077, arXiv:2207.00211 [hep-th] Kohei Kamada (RESCEU, U Tokyo) Workshop on Very Light Dark Matter 2023



29/03/2023, Marion Royal Kaikan, Chino



The message of this talk

1. In finite density (with "baryon" chemical potential), a layer of pion/axion domain wall is known to be stable once we apply a magnetic field, a. k. a. the Chiral Soliton Lattice.

2. We clarify the mechanism how to create such an object, through the quantum nucleation.



Chiral Soliton Lattice



 $\mathcal{L} = \frac{f^2}{2} (\partial_{\mu} \phi)^2 + f^2 m^2 (\cos \phi - 1)$





$$\mathcal{L} = \frac{f^2}{2} (\partial_\mu \phi)^2 + f^2 m^2 (\cos \phi)^2$$

allows a domain wall solution.





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Layer of domain walls are also allowed.





$$\frac{d}{dt} = \operatorname{sn}\left(\frac{m\xi}{k}, k\right)$$
$$(0 < k < 1)$$

k determines the distance between walls



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 $\cos\frac{\phi(\xi)}{2} = \operatorname{sn}\left(\frac{m\xi}{k}, k\right)$

k determines the distance between walls

Take into account the chiral anomaly

$\mathcal{L} = \frac{f^2}{2} (\partial_\mu \phi)^2 + f^2 m^2 (\cos \phi - 1) + c \alpha \phi F_{\mu\nu} \tilde{F}^{\mu\nu}$

in the vacuum; well-known Chern-Simons term

Not relevant for domain walls



Take into account the chiral anomaly

$$\mathcal{L} = \frac{f^2}{2} (\partial_\mu \phi)^2 + f^2 m^2 (\cos \phi)^2$$

In the medium with charged fermions:

Chiral magnetic effect (CME): $J_V =$

'80 Vilenkin, '06 Fukushima, Kharzeev, & Warringa

Chiral separation effect (CSE): $J_5 = J_R$

'80 Vilenkin, '04 Son & Zhitnitsky

At low energy, the chiral transformation turns to the shift of axions/pions

in the vacuum; well-known Chern-Simons term

 $(\phi - 1) + c\alpha\phi F_{\mu\nu}\tilde{F}^{\mu\nu}$

Not relevant for domain walls

$$J_R + J_L = \frac{e(\mu_R - \mu_L)}{2\pi^2} B$$
$$J_R - J_L = \frac{e(\mu_R + \mu_L)}{2\pi^2} B$$

$$-J_L = \frac{(1 + 1 + L)}{2\pi^2}I$$

>
$$\delta \mathcal{L} = \frac{1}{2} (\partial_{\mu} \phi) j_{5}^{\mu} \ni \frac{\mu_{B}}{4\pi^{2}} \nabla \phi \cdot B$$

 $(\mu_B = \mu_R + \mu_L; e = 1)$



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At low energy, the chiral transformation turns to the shift of axions/pions

also can be understood the WZW term with a spurious axial U(1) gauge field

in the medium; another "topological" term

 $(\phi - 1) + c\alpha \phi F_{\mu\nu} F^{\mu\nu}$

$$+\frac{\mu_B}{4\pi^2}\boldsymbol{B}\cdot\boldsymbol{\nabla}\phi$$

Relevant for domain walls...?

$$\boldsymbol{J}_R + \boldsymbol{J}_L = \frac{e(\mu_R - \mu_L)}{2\pi^2} \boldsymbol{B}$$

$$\boldsymbol{J}_R - \boldsymbol{J}_L = \frac{e(\mu_R + \mu_L)}{2\pi^2} \boldsymbol{B}$$

$$\delta \mathcal{L} = \frac{1}{2} (\partial_{\mu} \phi) j_{5}^{\mu} \ni \frac{\mu_{B}}{4\pi^{2}} \nabla \phi \cdot \boldsymbol{B}$$

'04 Son & Zhitnitsky



 $\cos\frac{\phi(\xi)}{2} =$

$$\operatorname{sn}\left(\frac{m\xi}{k},k\right)$$

then there are no effects from the topological term?



~ +b~r

Not really. Energy of the system changes. Total energy of the system of length *L* in the *z*-direction:

 $\mathcal{E} = \frac{Vm}{2kK(k)} \left[\frac{2E(k)}{k} + \left(k - \frac{1}{k}\right) K(k) - \frac{1}{k} \right] \right]$

$$\cos\frac{\phi(\xi)}{2} = \operatorname{sn}\left(\frac{m\xi}{k}, k\right)$$

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$$-\frac{\mu_B B}{2\pi}$$



thor

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Negative energy is induced from the spatial gradient of the field configuration.

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In the presence of non-zero chemical potential and external magnetic field

Domain wall solution is unchanged,

For sufficiently large magnetic field, $B > B_{\rm CSL} = \frac{16\pi m f^2}{\mu_B}$

Domain wall layer is energetically favored.

= Configuration at energy minimum k_{CSL} : "Chiral Soliton Lattice"

Negative energy is induced from the spatial gradient of the field configuration.





How it is interesting?

- Interest also for condensed matter system: known as Chiral Magnets.

- Forms in the neutron star? Helps to amplify magnetic field?

'13 Eto, Hashimoto, & Hatsuda





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Anyway we need to think how they form!

How Chiral Soliton Lattice forms

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- Classical move from the infinity would be unlikely.

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- Quantum tunneling of the layer of disk?

= disk wall + string loop on the edge

model it as the quantum mechanics of a disk with a Nambu-Goto-like action.

'91 Basu, Guth & Vllenkin, '20 Ai & Drewes

Effective action for disk

= disk wall + string loop on the edge

 $S_{\text{tot}} = S_{\text{wall}} + S_{\text{string}} = -\pi \int d\tau \left(2TR\sqrt{1 - \dot{R}^2} + \sigma R^2 \right)$

$$\begin{split} \sigma &= 8mf^2 - \frac{\mu B}{2\pi} \text{for one disk} \\ &< 0 \end{split} \qquad T \sim 2\pi \times 2f^2 \times \ln \frac{R_{\rm c}}{r_{\rm c}} \end{split}$$

Effective action for disk

$$S_{\rm tot} = S_{\rm wall} + S_{\rm string} = -\pi$$

R

Euclidean action

 $S_{\rm E}[R] = \pi \int d\tau_{\rm E} \left(2TR \sqrt{1 + \left(\frac{dR}{d\tau_{\rm E}}\right)^2} + \right)$

with bounce solution

$$\left(\frac{dR}{d\tau_{\rm E}}\right)^2 - 2V(R) = 0, \qquad V(R) = \frac{1}{2} \left[\frac{R^2}{\left(\epsilon - \frac{\sigma R^2}{2T}\right)^2}\right]$$

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$$\sigma = 8mf^2 - \frac{\mu B}{2\pi} \text{ for on}$$
$$< 0$$
$$T \sim 2\pi \times 2f^2 \times \ln \frac{R_c}{r_c}$$

$$+\sigma R^2$$

$$\left. 1 \right|, \quad \epsilon = \frac{E}{2\pi T} \to 0$$

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understood as the tunneling

Bounce action: $\mathcal{B} = 2 \times 2\pi T \int_0^{R_2} dR_N$

Coleman's formula: $P \simeq A e^{-B} = A e^{-B}$

'77 Coleman; Callan & Coleman

$$\sqrt{R^2 - (\sigma R^2 / 2T)^2} = \frac{16\pi T^3}{3\sigma^2}$$

$$e^{-16\pi T^3/3\sigma^2}$$
 $\mathcal{A} \simeq \left(\frac{\mathcal{B}}{2\pi}\right)^2 \frac{1}{R_2^4}$

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Exponential suppression factor (bounce action) is less the 1 for

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$B > B_c = \frac{16\pi m f^2}{m} \left(\frac{(\ln(R_c/r_c))^{3/2}}{\sqrt{3}} \frac{4\pi^2 f}{m} + 1 \right) (> B_{\rm CSL})$

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For CSL, we shall just replace σ

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$$\left(k-\frac{1}{k}\right)K(k)\right] - \frac{\mu B}{2\pi}$$
 smaller amplitude of tension
=> larger bounce action

Bounce actic

Coleman's fc

'77 Coleman

Exponential

For CSL, we

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 smaller amplitude of tension
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Whole picture?

 $B \bigstar$

single wall individually form

Whole picture?

 $B \bigstar$

expand, sometimes merge each other

Whole picture?

 $B \bigstar$

Concrete examples: QCD pion

It is difficult for CSL to form without exponential suppression

Concrete examples: ALP

- also difficult to setup the situation where $\mathcal{B} < 1$

 $B_{\mathcal{B}=1} \simeq \frac{64\pi^2 f^3}{\mu_B}$

after EWSB $\mu_B \ll B^{1/2} \lesssim T \implies$ for $f \gg T$ inevitably $\mathcal{B} \gg 1$

Concrete examples: ALP

- also difficult to setup the situation where $\mathcal{B} < 1$

 $B_{\mathcal{B}=1} \simeq \frac{64\pi^2 f^3}{\mu_B}$

Unfortunately, we cannot expect for the effect on cosmic birefringence in CMB.

after EWSB $\mu_B \ll B^{1/2} \lesssim T \implies$ for $f \gg T$ inevitably $\mathcal{B} \gg 1$ before EWSB, it can be $\mu_B \sim B^{1/2} \lesssim T$

with $f \sim T$ we might have $\mathcal{B} \sim 1$

Summary

- Chiral anomaly induces a nontrivial topological interaction between axion/pion and magnetic field at finite density (chemical potential). - It makes the domain wall layer energetically favorable,
- which is called as "Chiral Soliton Lattice".
- Its formation is described by the nucleation of a domain wall disk.
- The formation rate is relatively suppressed.
- Implication of light axions? Indeed, thin-wall approximation becomes bad.

