

# A01: Light dark matter

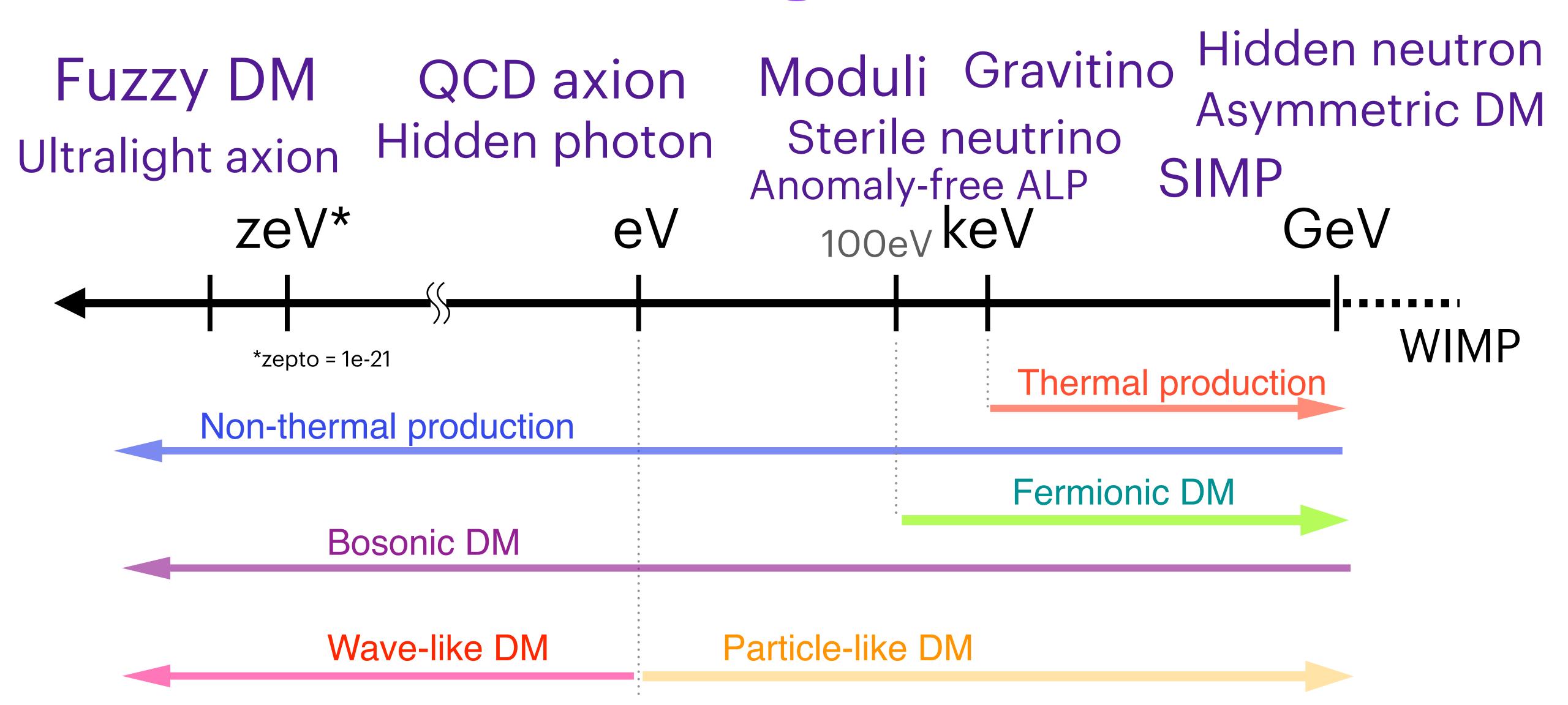
1.0e+00 1.0e+00 0.2 0.05 0.005 0.005 0.005 0.005 0.005

Fumi Takahashi (Tohoku)

Mar. 7. 2024 @

"What is dark matter? - Comprehensive study of the huge discovery space in dark matter"

# Mass scale of light dark matter



#### Members

Masahiro Kawasaki

Naoya Kitajima

Fuminobu Takahashi

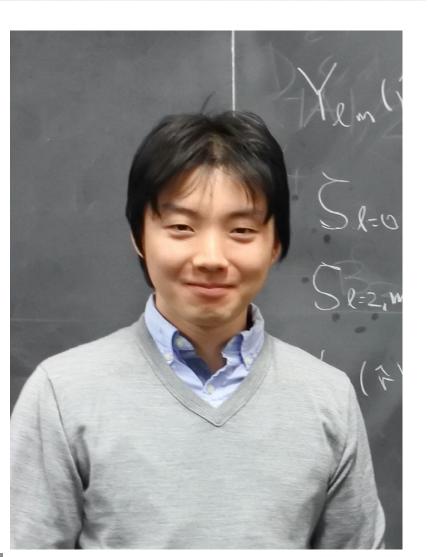
Masaki Yamada

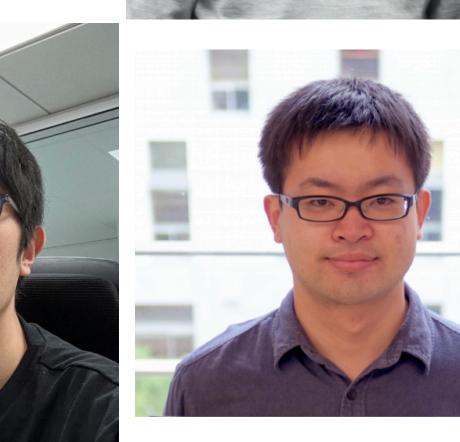
Wen Yin

Shota Nakagawa

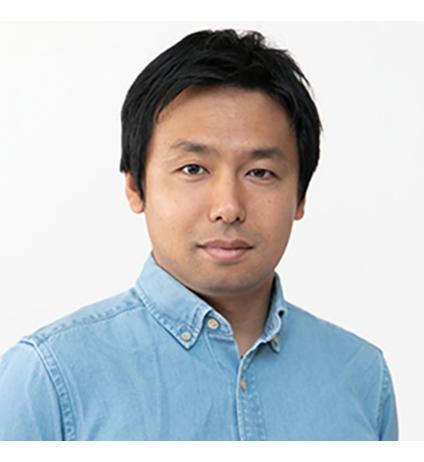
Kai Murai

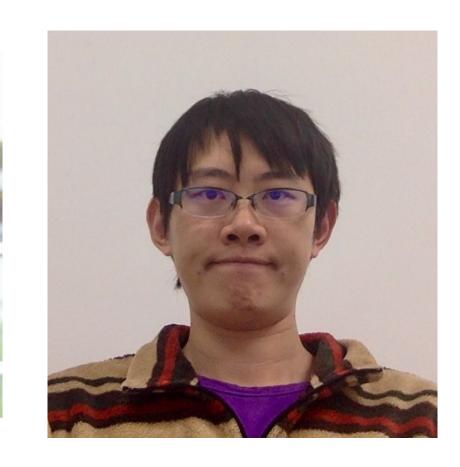












## Papers

28 papers appeared on arXiv since Apr. 2023. (104 papers since Oct. 2020)









WWW.PHDCOMICS.COM

# Highlights

• First Result for Dark Matter Search by WINERED See talk by Yin

Wen Yin, Taiki Bessho, Yuji Ikeda, Hitomi Kobayashi, Daisuke Taniguchi et al. 2402.07976

Misalignment production of vector boson dark matter from axion-SU(2) inflation

Tomohiro Fujita, Kai Murai, Kazunori Nakayama, Wen Yin, 2312.06889

See talk by Murai

Cosmic strings from pure Yang–Mills theory

collaboration with BO1

Masaki Yamada, Kazuya Yonekura 2204.13123, 2204.13125, 2307.06586

See talk by Kitajima

Dark photon dark matter production

Naoya Kitajima and Kazunori Nakayama 2303.04287, FT, Naoya Kitajima, 2303.05492

Axion dark matter from FOPT

Shota Nakagawa, FT, Masaki Yamada, and Wen Yin 2210.10022 Junseok Lee, Kai Miurai, FT, and Wen Yin 2402.09501

Primordial black holes from QCD axion bubbles

Kentaro Kasai, Masahiro Kawasaki, Naoya Kitajima, Kai Murai, Shunsuke Neda, FT 2305.13023, 2310.13333

Gravitational waves from domain walls

Naoya Kitajima, Junseok Lee, Kai Murai, FT, and Wen Yin 2306.17146 Naoya Kitajima, Junseok Lee, FT, and Wen Yin 2311.14590

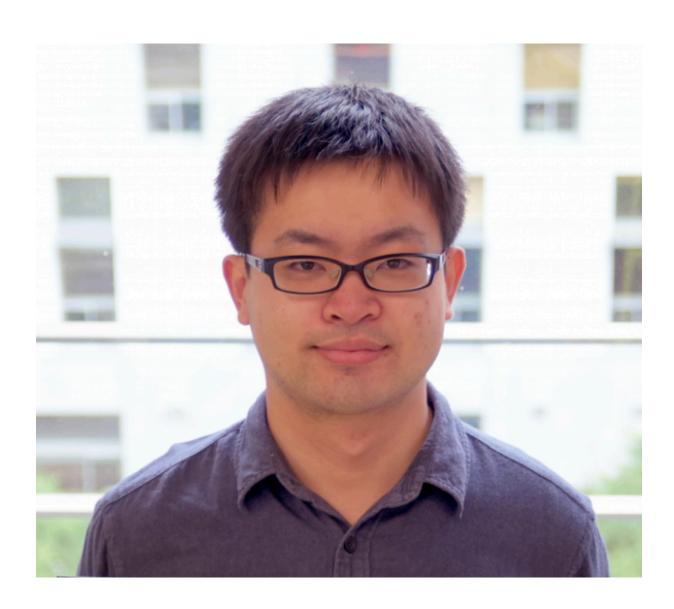
• QCD axion hybrid inflation Yuma Narita, FT and Yin, 2308.12154

#### Bubble misalignment mechanism

Junseok Lee, Kai Murai, FT and Wen Yin, 2402.09501



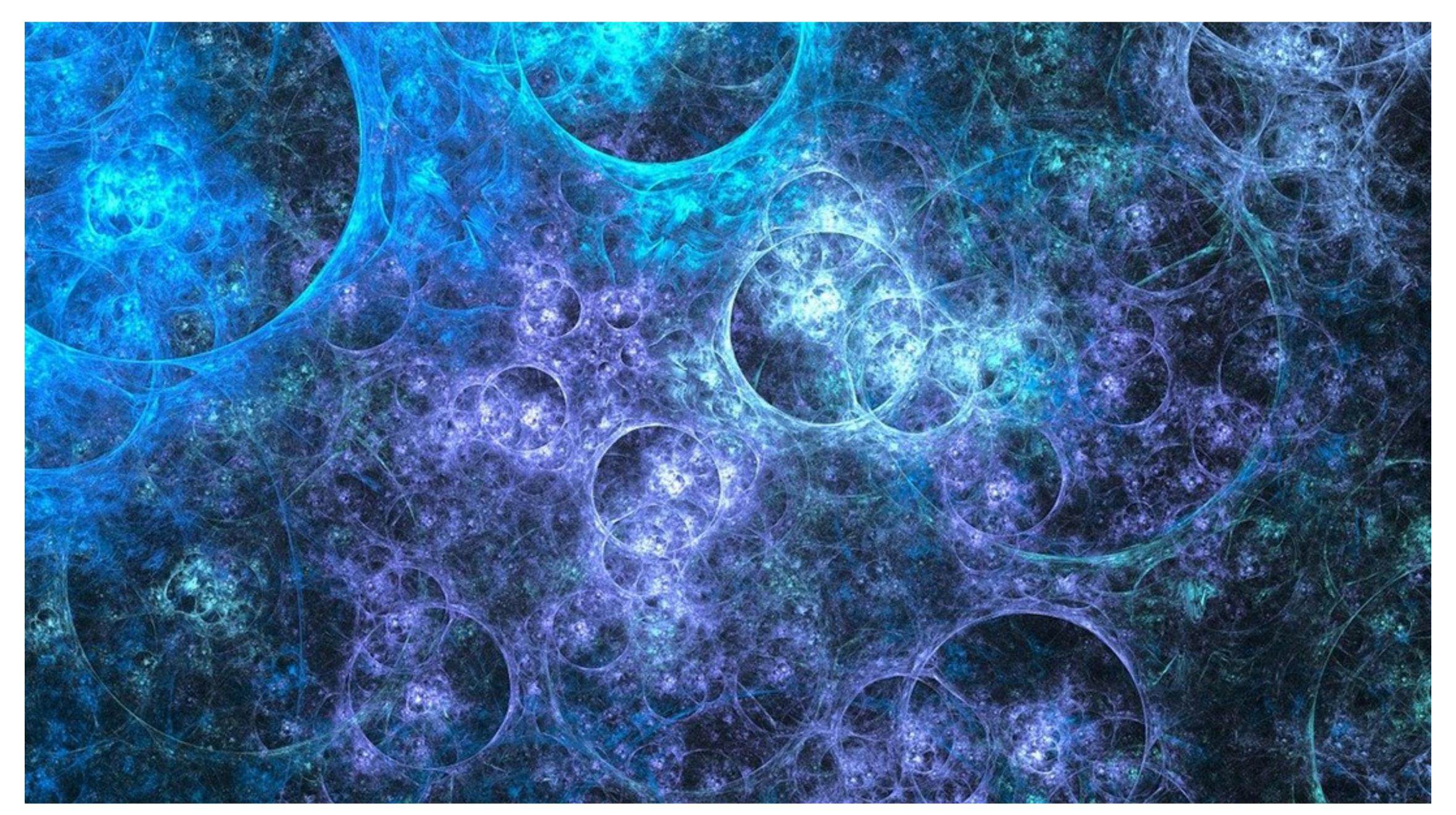
Junseok Lee



Kai Murai



Wen Yin



https://www.advancedsciencenews.com/footprints-of-phase-transitions-in-the-early-universe/

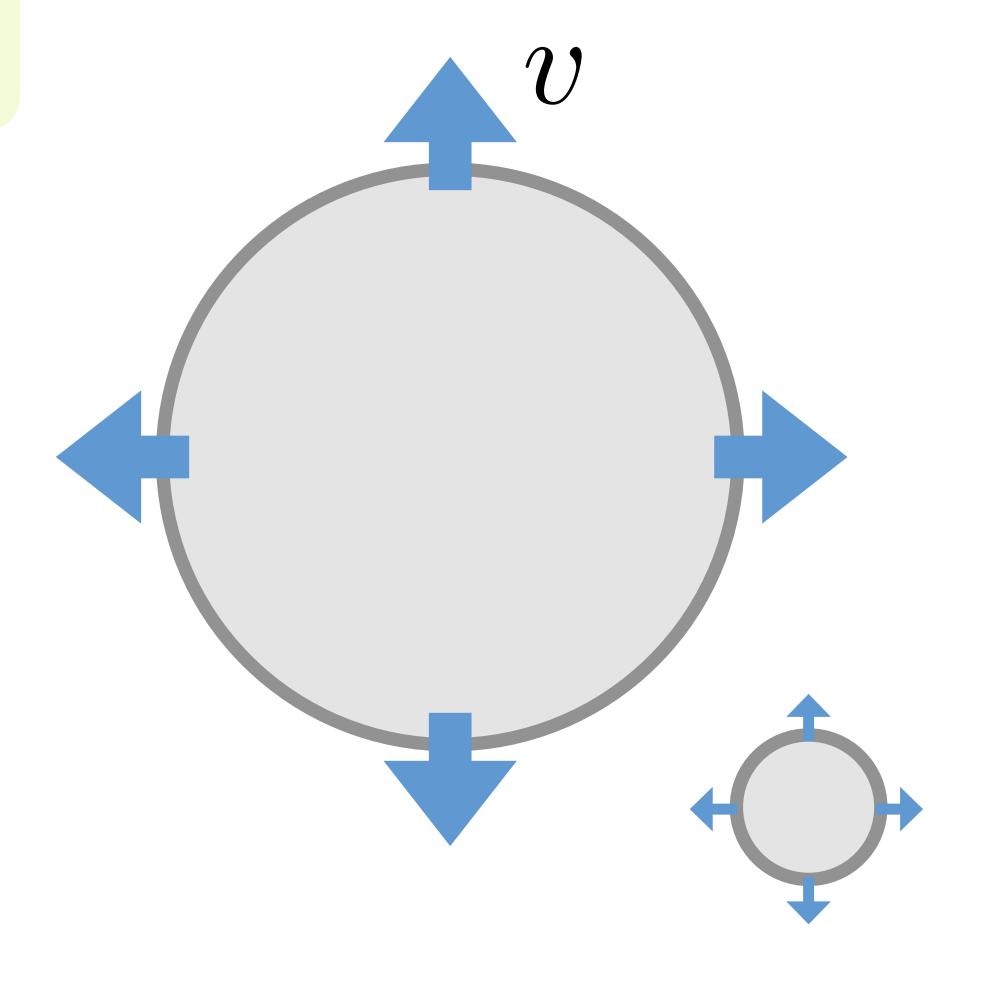
#### Bubble nucleation in FOPT

We introduce two parameters: v and  $\beta$ 

- Once bubbles nucleate, they expand with velocity  $\boldsymbol{v}$ .
- The bubble nucleation rate generically depends on temperature (or time);

$$\Gamma \approx \Gamma(t_0) \exp \left[\beta(t-t_0) + \cdots\right]$$

 $\beta^{-1}$  defines the duration of FOPT.



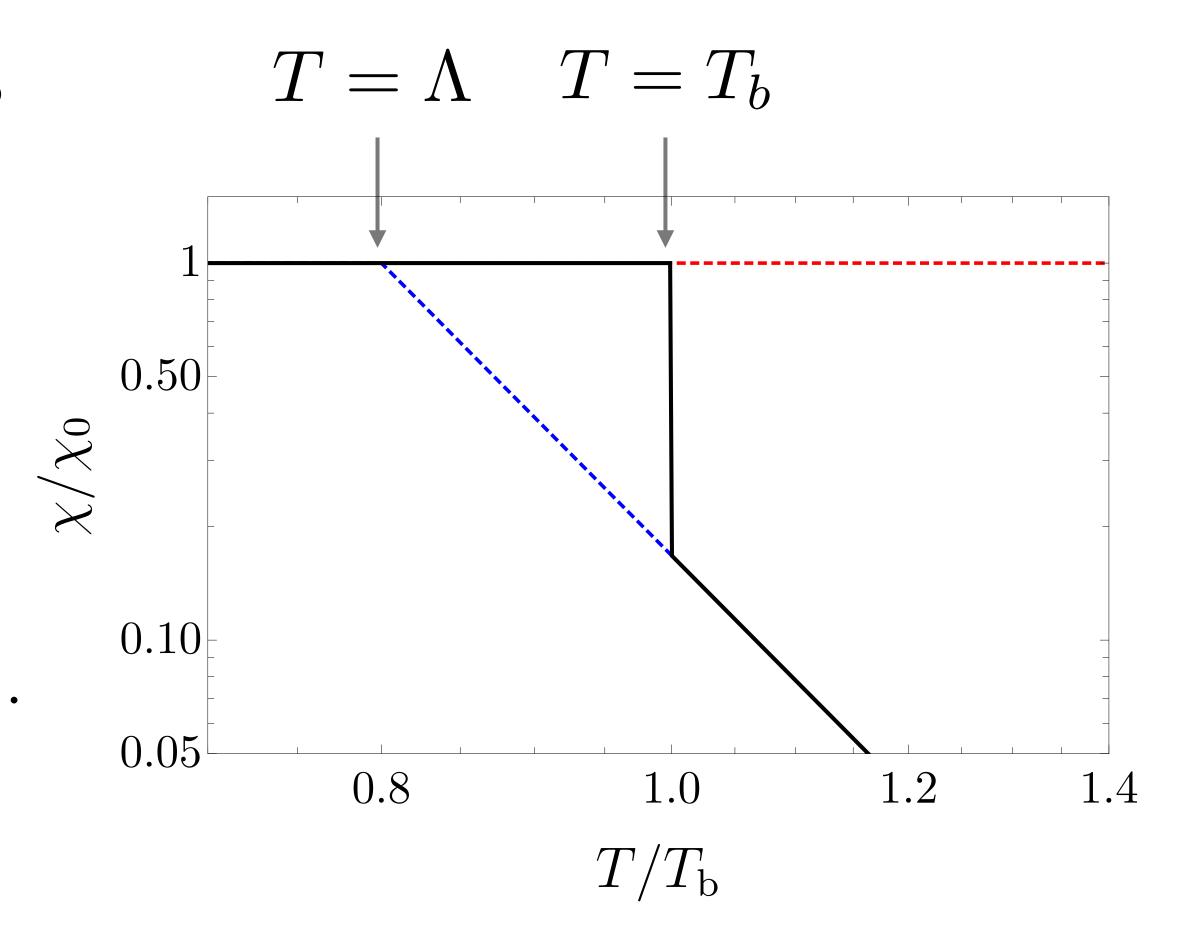
We assume  $\beta > H_{\rm b}$ .

The topological suceptibility changes discontinuously at the transition.

$$V(\phi) = \chi(T) \left[ 1 - \cos\left(\frac{\phi}{f_{\phi}}\right) \right]$$

with 
$$\chi(T)\simeq\left\{ egin{array}{ll} \chi_0 & (T<\Lambda) \\ \chi_0 \Big(rac{T}{\Lambda}\Big)^{-p} & (T\geq\Lambda) \end{array} 
ight. .$$

$$m_{\rm b} \equiv \frac{\sqrt{\chi(T_{\rm b})}}{f_{\phi}} < m_0 \equiv \frac{\sqrt{\chi_0}}{f_{\phi}}$$



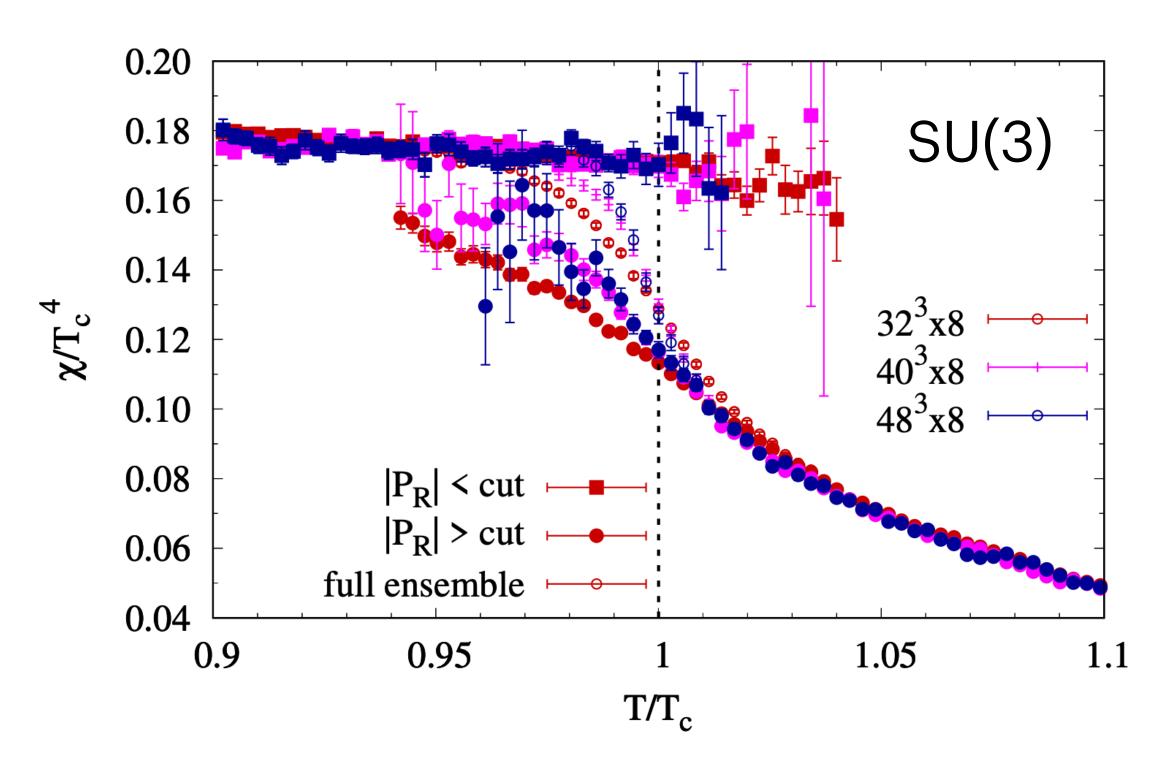
The axion mass changes from  $m_{\rm b}$  to  $m_0$  at  $T=T_{\rm h}$ .

The topological suceptibility changes discontinuously at the transition.

$$V(\phi) = \chi(T) \left[ 1 - \cos\left(\frac{\phi}{f_{\phi}}\right) \right]$$

with 
$$\chi(T)\simeq\left\{egin{array}{ll} \chi_0 & (T<\Lambda) \ \chi_0 \Big(rac{T}{\Lambda}\Big)^{-p} & (T\geq\Lambda) \end{array}
ight.$$

$$m_{\rm b} \equiv \frac{\sqrt{\chi(T_{\rm b})}}{f_{\phi}} < m_0 \equiv \frac{\sqrt{\chi_0}}{f_{\phi}}$$



S. Borsányi et al, 2212.08684

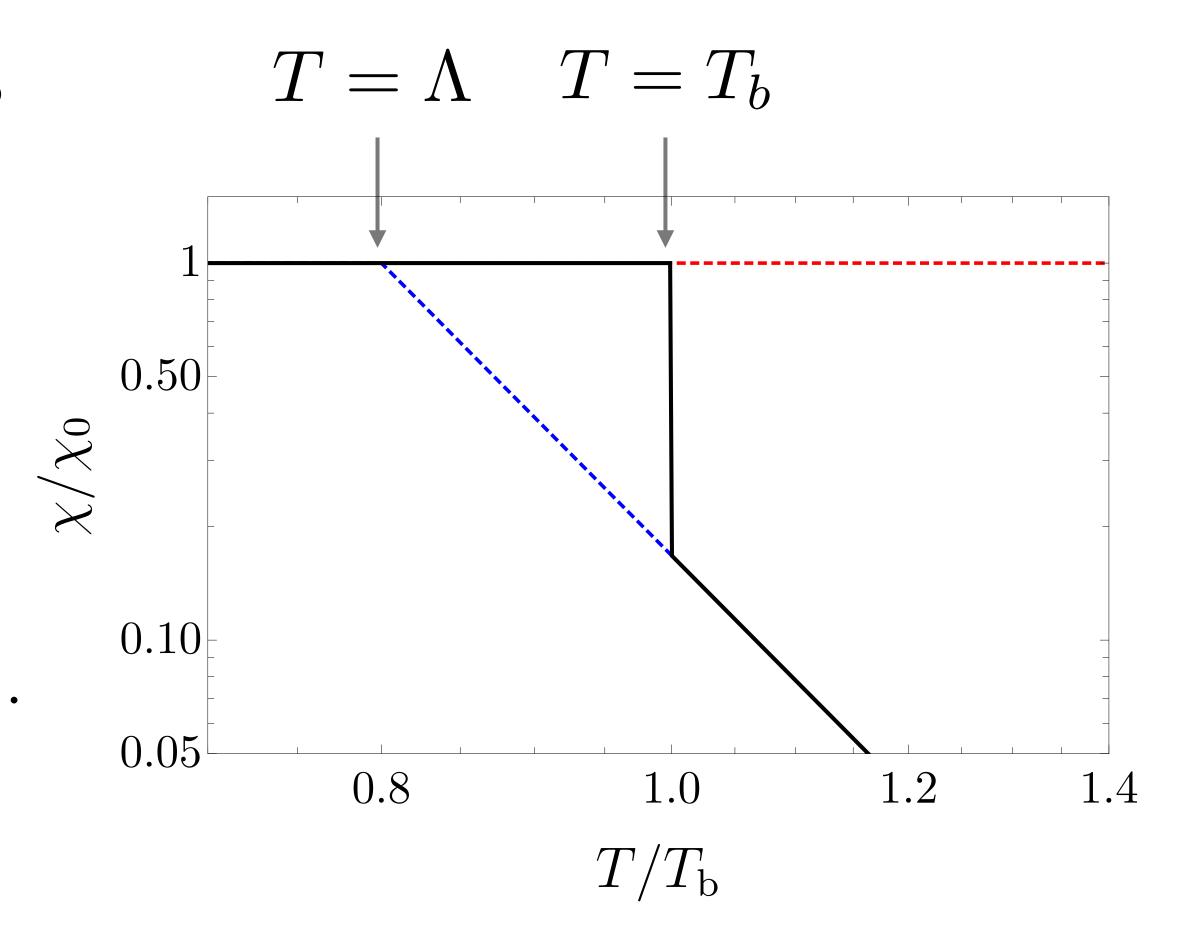
The axion mass changes from  $m_{\rm b}$  to  $m_0$  at  $T=T_{\rm h}$ .

The topological suceptibility changes discontinuously at the transition.

$$V(\phi) = \chi(T) \left[ 1 - \cos\left(\frac{\phi}{f_{\phi}}\right) \right]$$

with 
$$\chi(T)\simeq\left\{ egin{array}{ll} \chi_0 & (T<\Lambda) \\ \chi_0 \Big(rac{T}{\Lambda}\Big)^{-p} & (T\geq\Lambda) \end{array} 
ight. .$$

$$m_{\rm b} \equiv \frac{\sqrt{\chi(T_{\rm b})}}{f_{\phi}} < m_0 \equiv \frac{\sqrt{\chi_0}}{f_{\phi}}$$



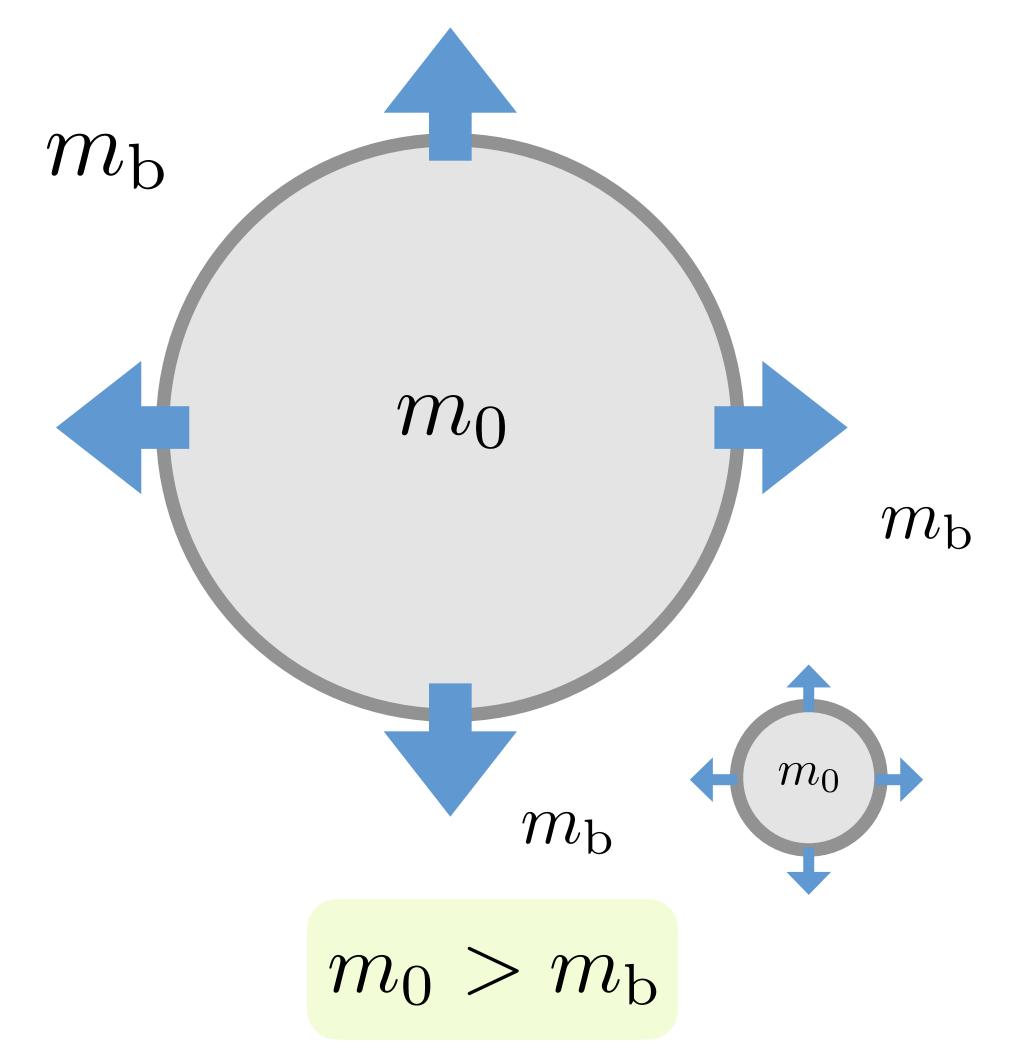
The axion mass changes from  $m_{\rm b}$  to  $m_0$  at  $T=T_{\rm h}$ .

The topological suceptibility changes discontinuously at the transition.

$$V(\phi) = \chi(T) \left[ 1 - \cos\left(\frac{\phi}{f_{\phi}}\right) \right]$$

with 
$$\chi(T)\simeq\left\{egin{array}{ll} \chi_0 & (T<\Lambda) \ \chi_0 \Big(rac{T}{\Lambda}\Big)^{-p} & (T\geq\Lambda) \end{array}
ight.$$

$$m_{\rm b} \equiv \frac{\sqrt{\chi(T_{\rm b})}}{f_{\phi}} < m_0 \equiv \frac{\sqrt{\chi_0}}{f_{\phi}}$$



The axion mass changes from  $m_{\rm b}$  to  $m_0$  at  $T=T_{\rm b}.$ 

\_ee, Murai, FT and Yin <u>2402.09501</u>

The most important criterion is whether the axion can move inside the bubble during the FOPT.

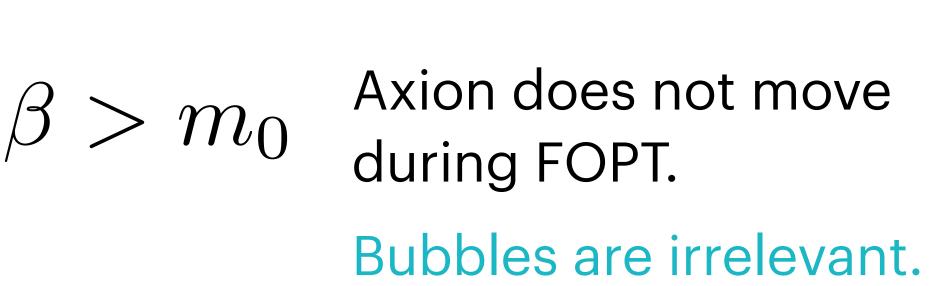
$$\beta > m_0$$
 Axion does not move during FOPT.

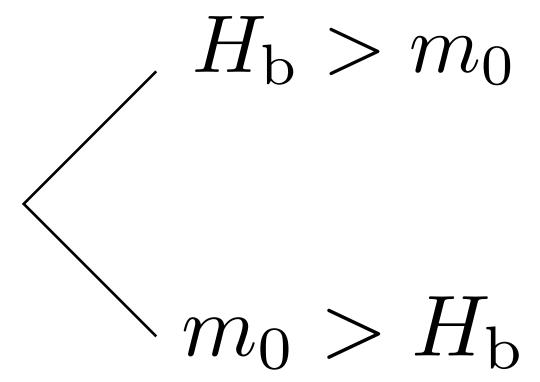
Bubbles are irrelevant.

\_ee, Murai, FT and Yin <u>2402.09501</u>

The most important criterion is whether the axion can move inside the bubble during the FOPT.

 $\beta > H_{\rm b}$ 





Axion begins to oscillate some time after FOPT.



Axion begins to oscillate right after FOPT.

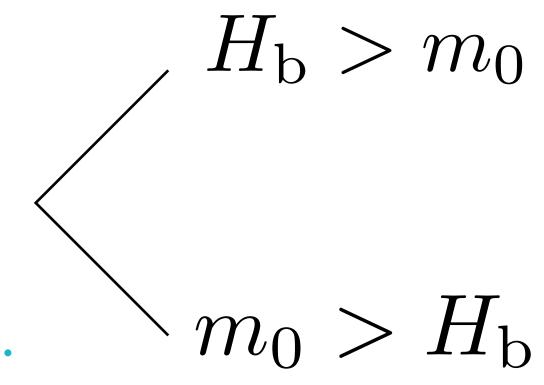


Lee, Murai, FT and Yin 2402.09501

The most important criterion is whether the axion can move inside the bubble during the FOPT.

$$\beta > H_{\rm b}$$

$$eta > m_0$$
 Axion does not move during FOPT. Bubbles are irrelevant.



Axion begins to oscillate some time after FOPT.



Constant axion mass.

Axion begins to oscillate right after FOPT.



Sudden change in axion mass.

$$m_0 > eta \quad {
m Axion inside bubbles} \ {
m moves during FOPT.}$$

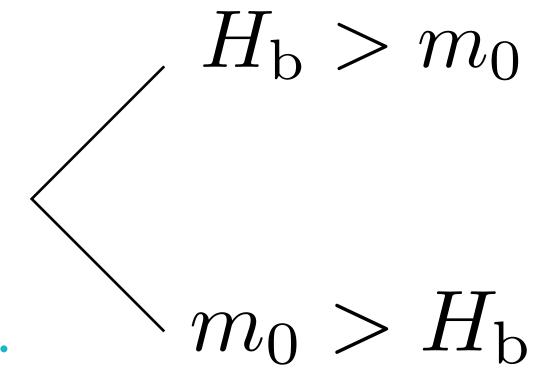
Bubbles significantly influence the axion dynamics!

Lee, Murai, FT and Yin 2402.09501

The most important criterion is whether the axion can move inside the bubble during the FOPT.

$$\beta > H_{\rm b}$$

$$eta > m_0$$
 Axion does not move during FOPT. Bubbles are irrelevant.



Axion begins to oscillate some time after FOPT.



Axion begins to oscillate right after FOPT.



Sudden change in axion mass.

Axions may or may not oscillate outside bubbles during FOPT.

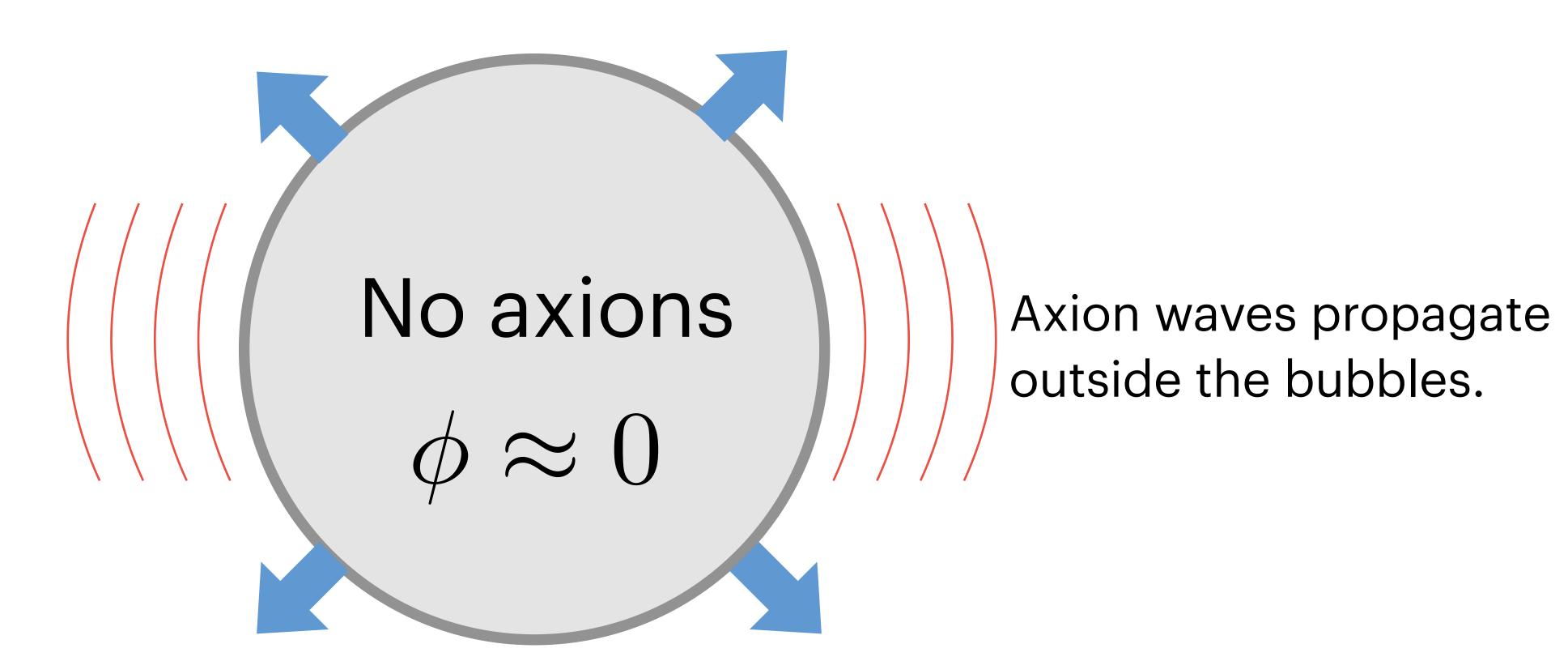
$$m_0>eta$$
 Axion inside bubbles  $m_{
m b}>eta$  moves during FOPT.  $m_{
m b}>eta$ 

Bubbles significantly influence the axion dynamics!

Lee, Murai, FT and Yin 2402.09501

Lee, Murai, FT and Yin 2402.09501

1) Axions are expelled from the inside of bubbles.

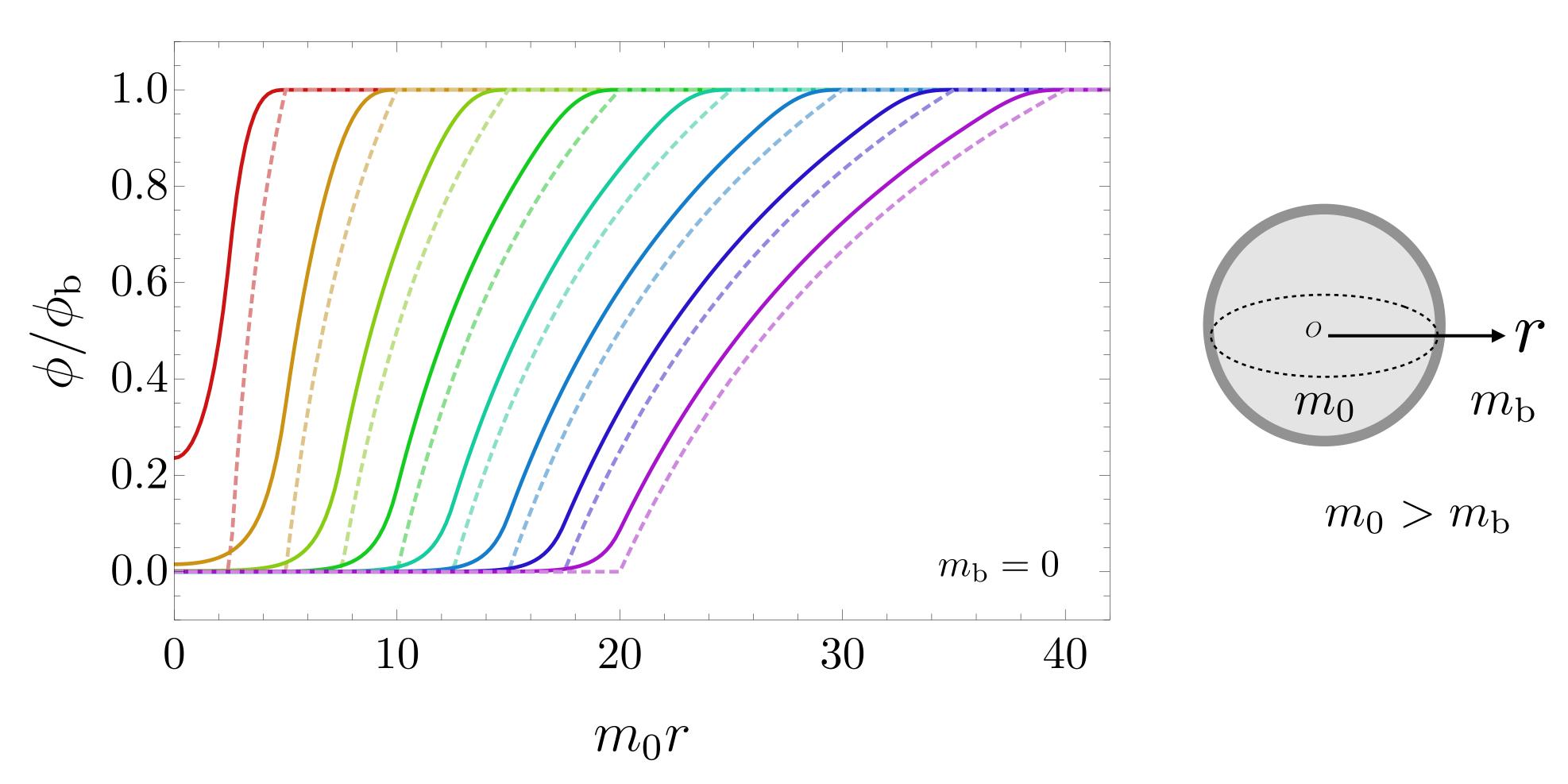


For  $m_{\rm b}>\beta$  , axion oscillates in the entire space during FOPT.

For  $m_{\rm b} < \beta$  , axion shock wave is induced near the wall.

Lee, Murai, FT and Yin 2402.09501

1) Axions are expelled from the inside of bubbles.

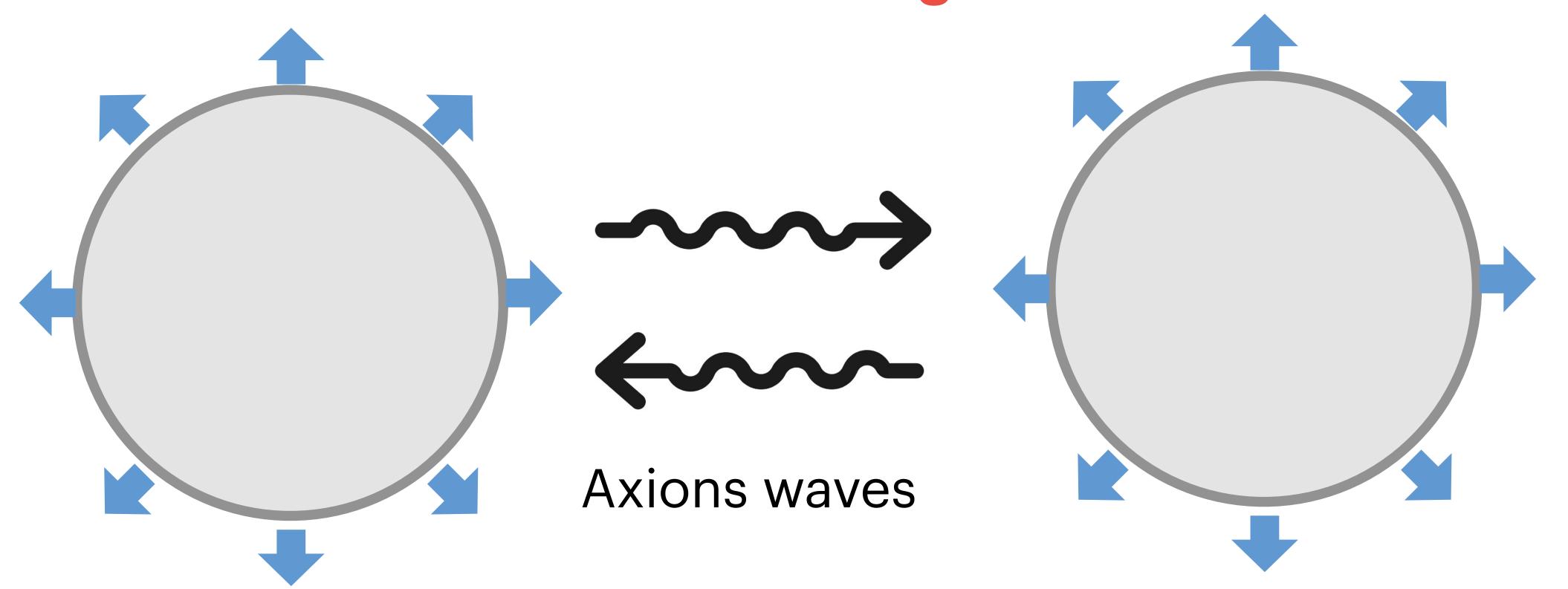


For  $m_{\rm b} < H_{\rm b}$  , axion shock wave is induced outside the wall.

Lee, Murai, FT and Yin 2402.09501

2 Axion waves accumulate outside the bubbles and repeatedly scatter off the bubble walls. They are accelerated with each scattering.

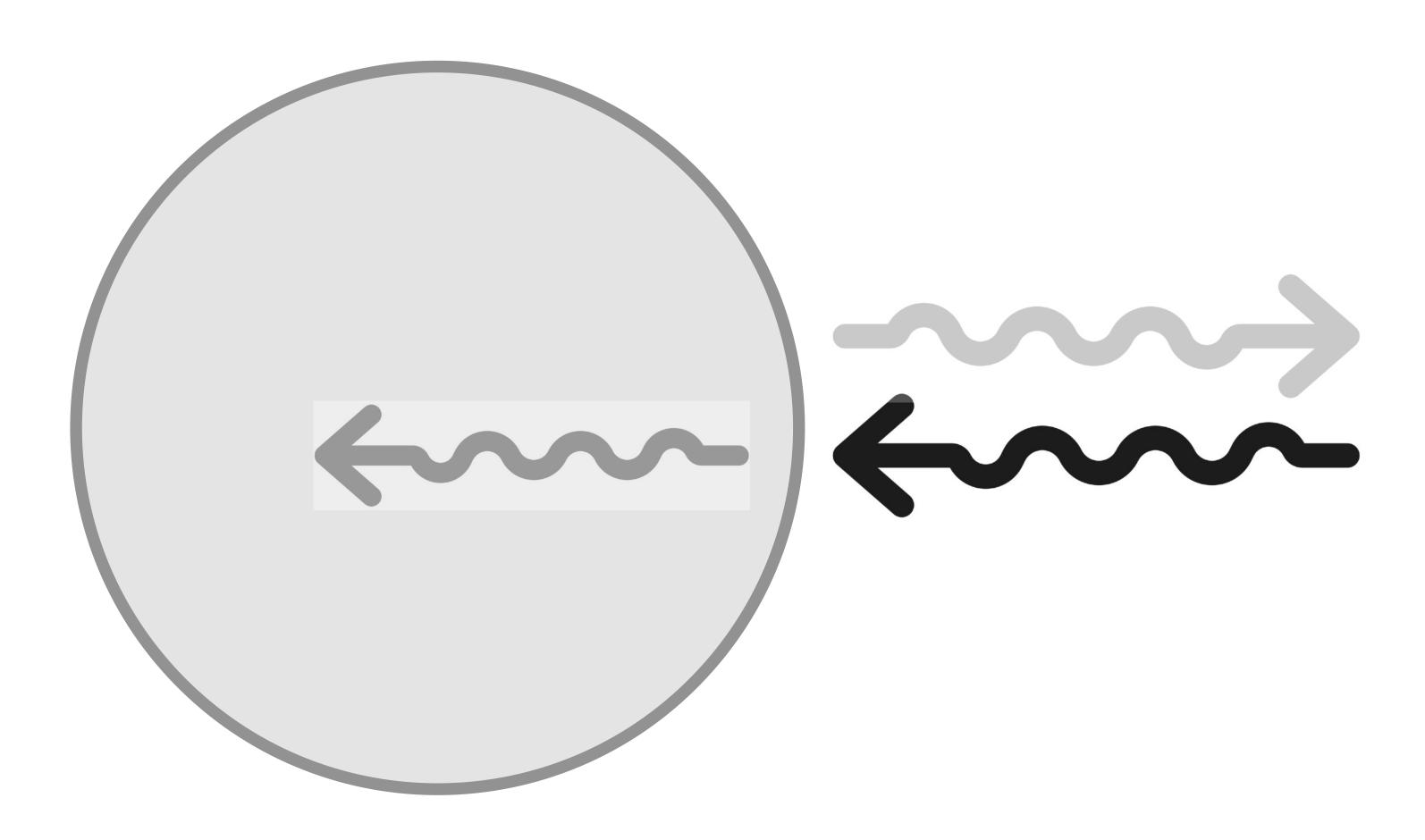
This is analogous to Fermi acceleration!



Lee, Murai, FT and Yin 2402.09501

3 Axion waves enter the bubbles when they gain enough energy.

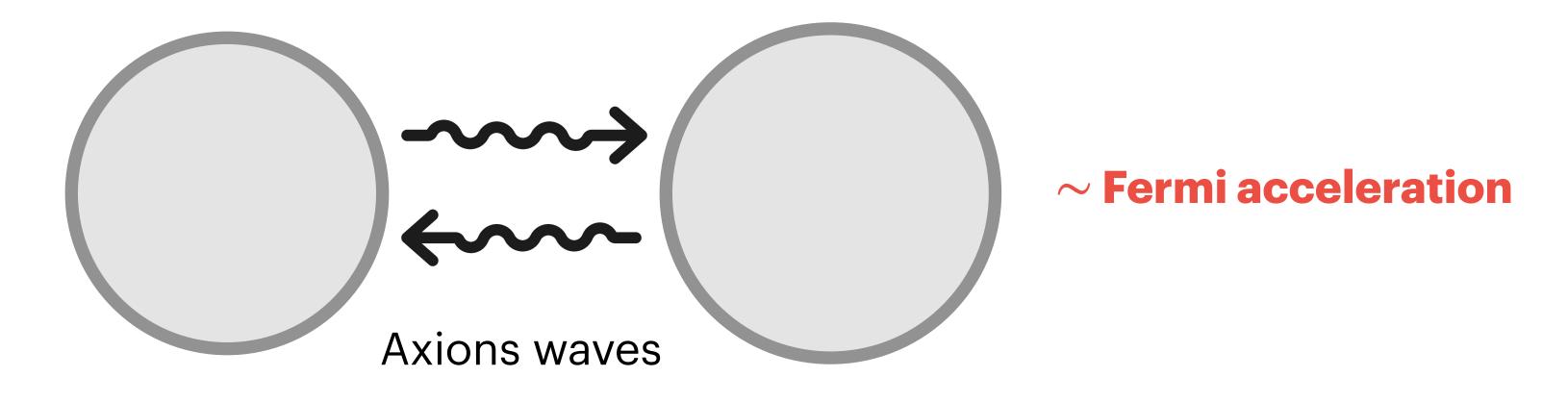
The axion number increases at the transmission.



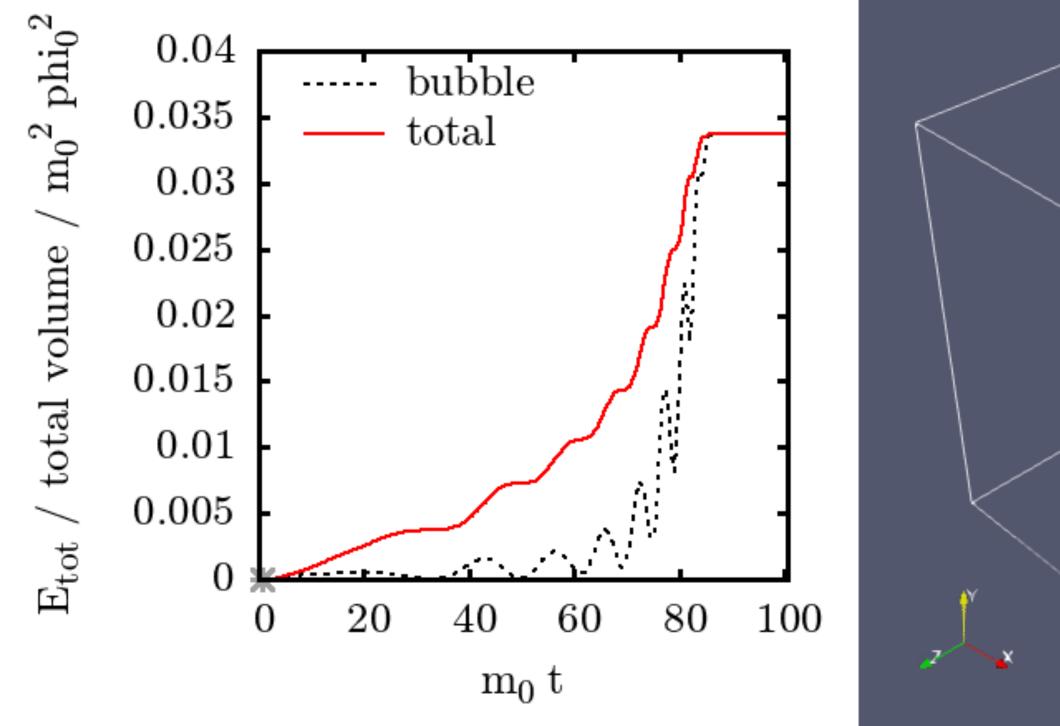
Lee, Murai, FT and Yin 2402.09501

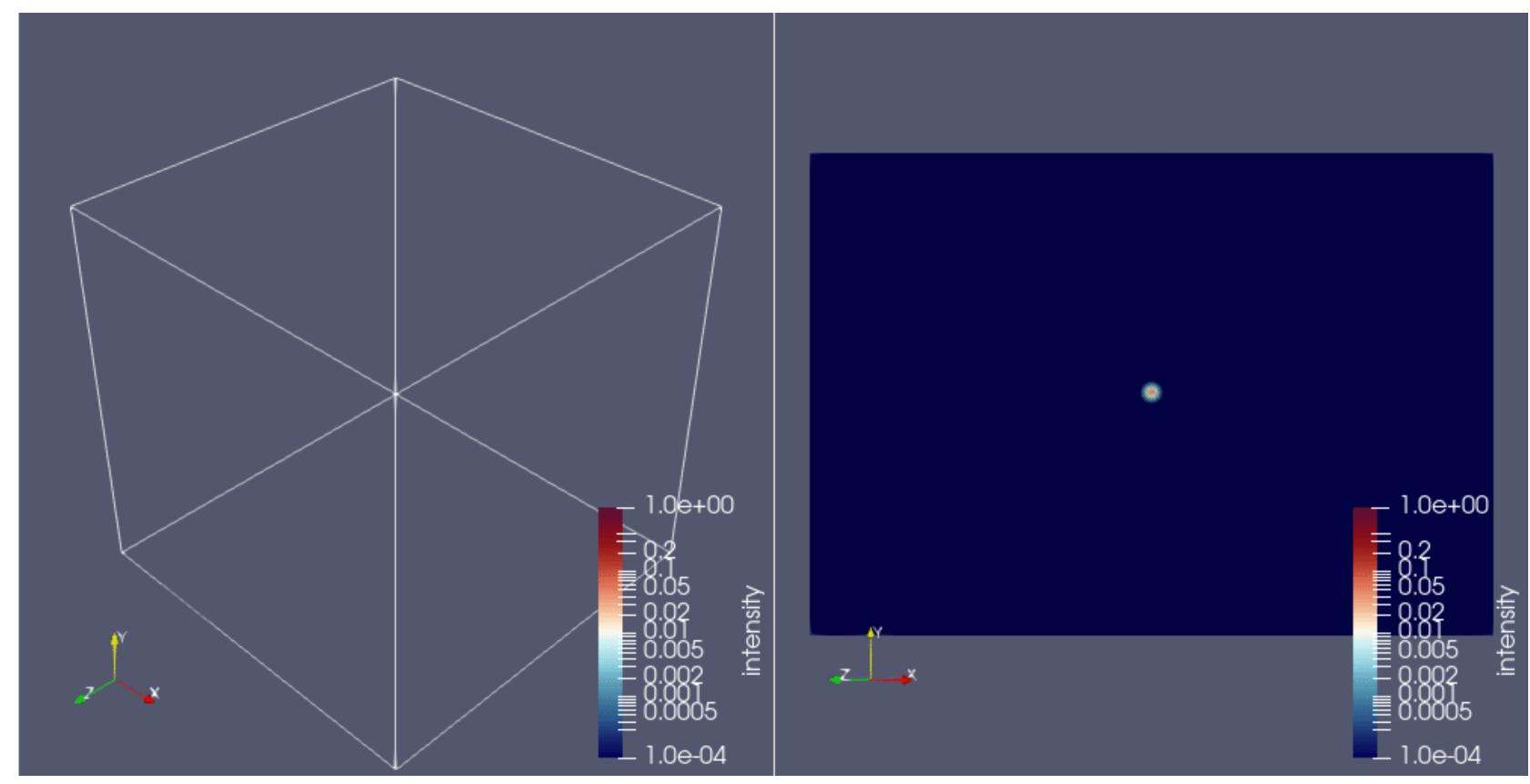
- Axions are expelled from the inside of bubbles.
- No axions

Axions waves are accumulated outside the bubbles and repeatedly scatter off bubble walls.



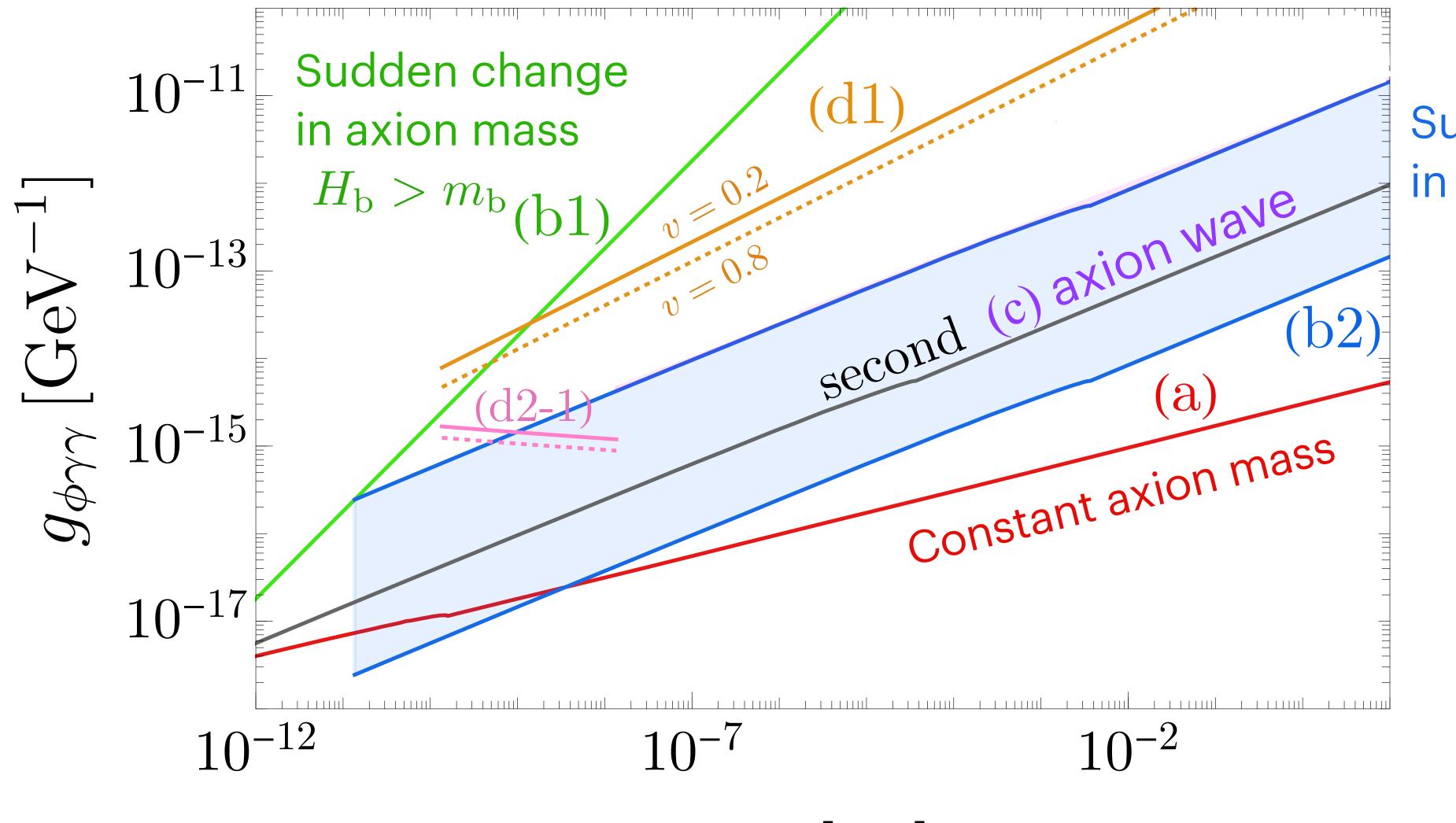
Axions waves enter the bubbles if they acquire enough energy.





# Viable region for axion dark matter

Axion shock wave



Sudden change in axion mass  $H_{
m b} < m_{
m b}$ 

$$g_{\phi\gamma\gamma} = lpha/(2\pi f_{\phi})$$
  
 $\phi_0 = f_{\phi}$   
 $p = 8$   
 $T_{
m b} = 10 \, {
m MeV}$ 

$$m_0 \, [\mathrm{eV}]$$

#### Summary

- We studied the axion evolution in the FOPT, taking account of the bubble dynamics; "Bubble misalignment mechanism"
- We find that axion is expelled from the interior of the bubbles and that Fermi acceleration occurs.
- The axion abundance can be significantly increased compared to the case of constant axion mass.
- Much to be done: analysis of realistic bubble nucleation, oscillon/I-ball formation, axion minicluster, production of dark photon dark matter, etc.