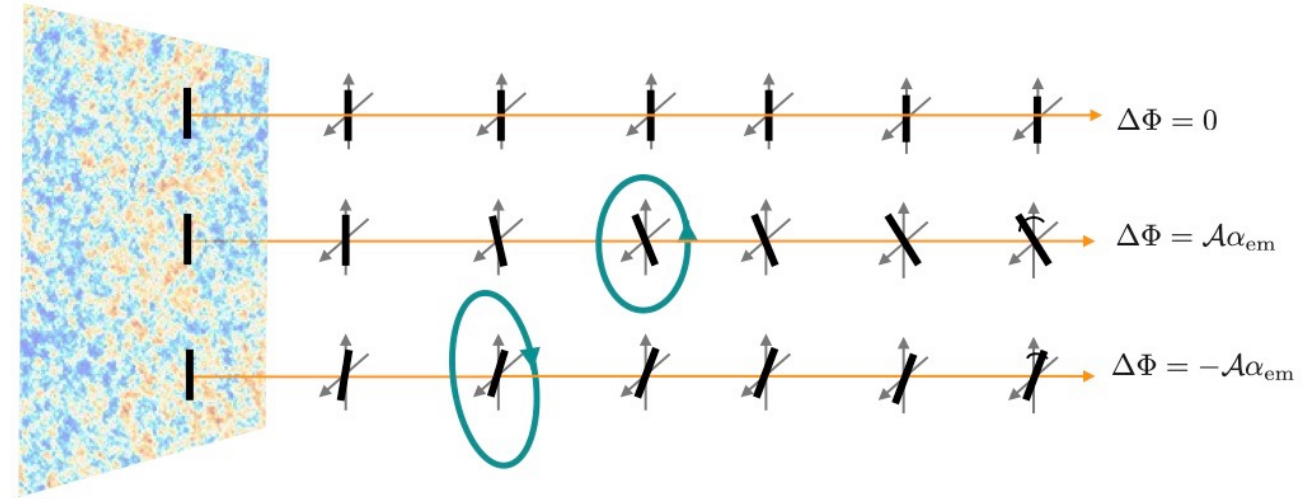


CMB Birefringence from Axion Strings



Extreme Mass Dark
Matter Workshop:

yeV zeV

ZeV YeV

Superlight → Superheavy

Andrew J. Long
Rice University
@ YITP
March 8, 2024



Summary

- If a [hyper-light axion-like particle](#) exists in Nature, the associated cosmological [network of axion strings](#) can leave an imprint on [CMB polarization](#) through birefringence
- We use existing [measurements of anisotropic birefringence](#) (Planck, SPT, ...) to place constraints on this scenario. Next-generation telescopes (CMB-S4) will probe $O(1)$ electromagnetic anomaly coefficients and thereby probe the axion's UV embedding
- We find that it is difficult (but not impossible!) to reconcile the [detection of isotropic birefringence](#) with strong limits on anisotropic birefringence coming from axion strings
- We argue that measurements of anisotropic birefringence could not only reveal the presence of a hyper-light ALP in Nature, but also lead to a [measurement of its mass](#)
- Our ongoing work (very early stages) seeks to use machine learning techniques (spherical CNN) to detect the subtle signal of axion strings in CMB polarization data

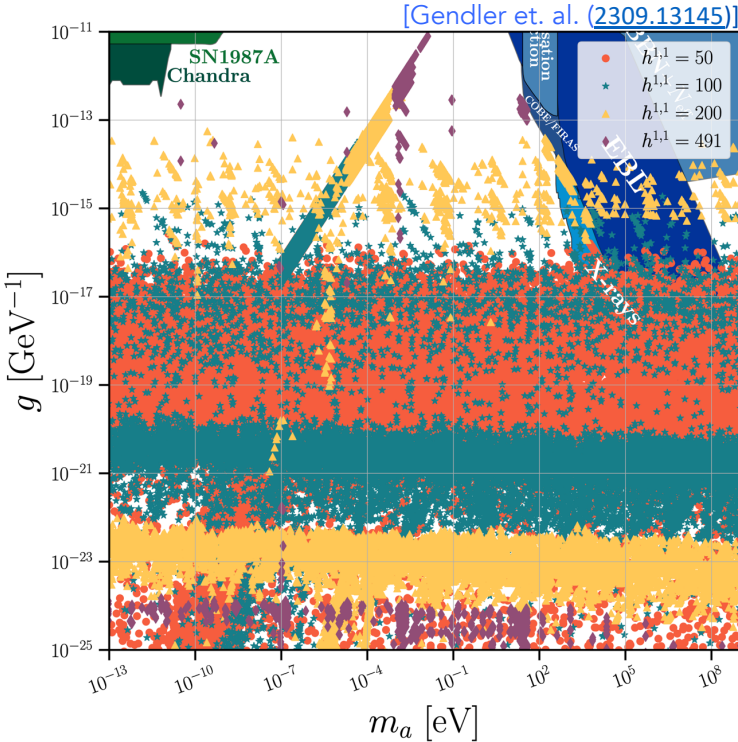
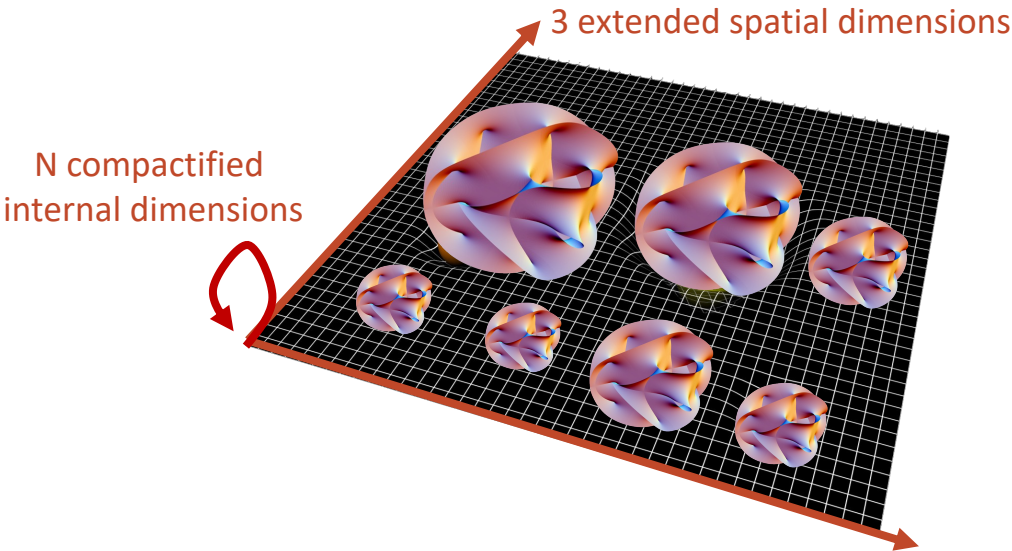
axion-like particles
& cosmic axion strings

Theory landscape: axion-like particles

axion-like particles

$$\mathcal{L} \supset \frac{1}{2}(\partial a)^2 - \frac{1}{2}m_a^2 a^2 - \frac{1}{4}g_{a\gamma\gamma} a F_{\mu\nu} \tilde{F}^{\mu\nu}$$

ALPs from extra dimensions
(such as string theory)



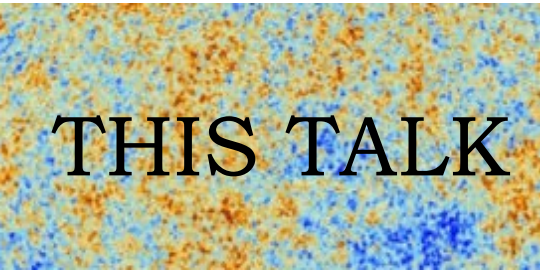
[Gendler et. al. (2309.13145)]

Theory landscape: axion-like particles

axion-like particles

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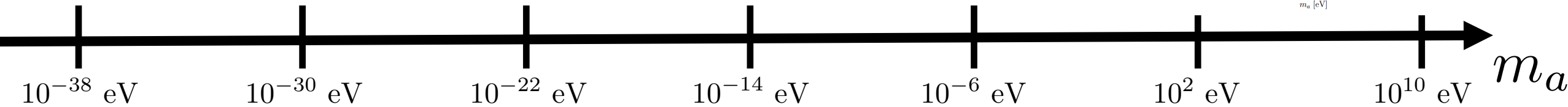
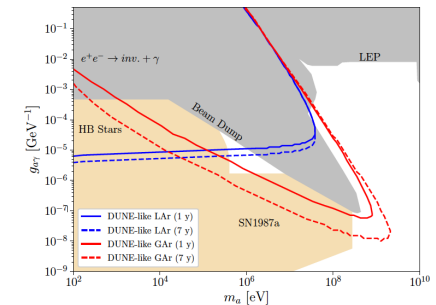
hyper-light axion-like particles
(testable with cosmology)



ultra-light axion-like particles
(dark matter candidate)



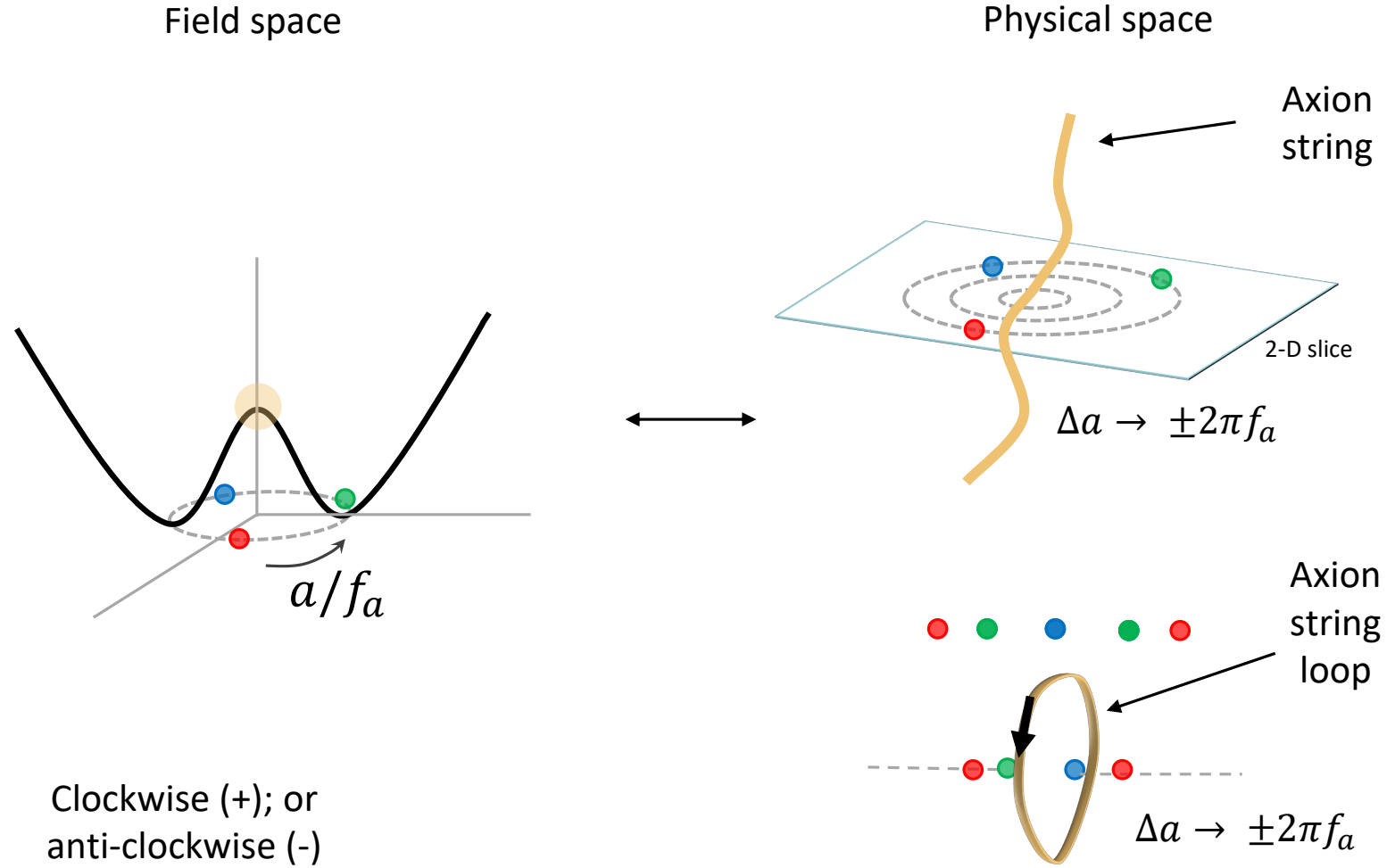
heavy axion-like particles
(testable in the lab)



ALPs form axion strings

[Kibble (1976)]
[Vilenkin & Vachaspati (1987)]

string formation:
early-universe phase transition



string thickness = microscopic

string length = cosmological

assume: $T_{RH} > f_a$

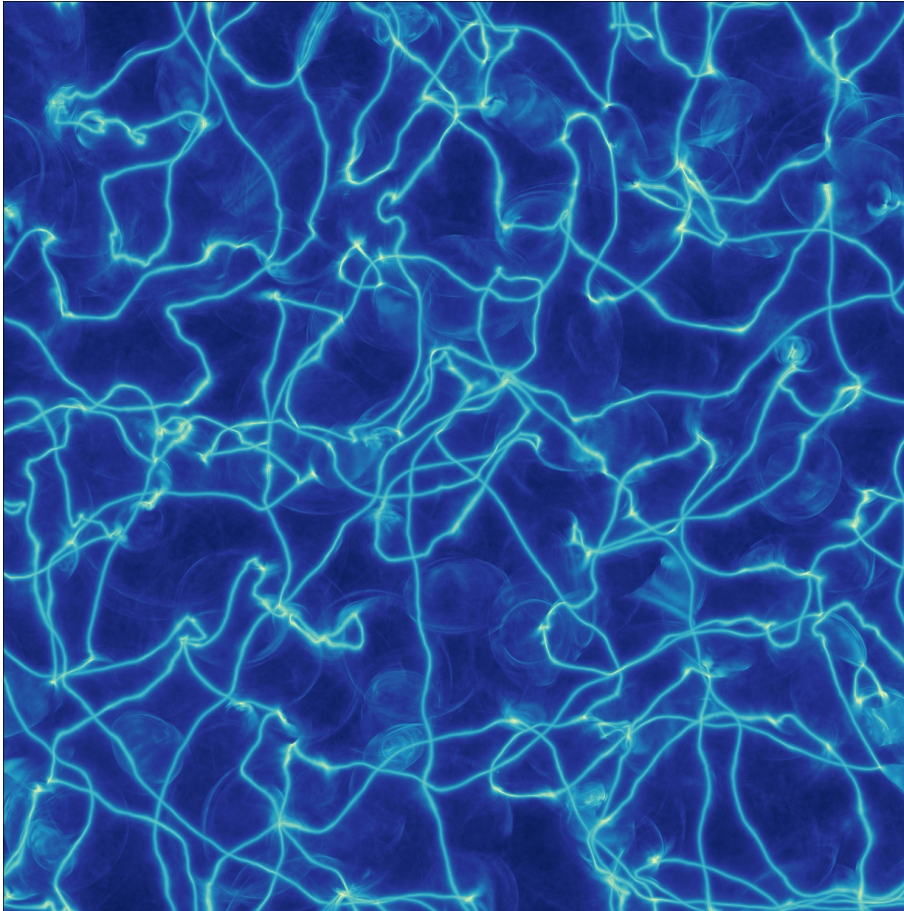
Clockwise (+); or
anti-clockwise (-)

image credit: Mudit Jain (2021)

A cosmic string network

[Saikawa et. al. (2024)]

string network simulation:



see talk by :
Ken'ichi
Saikawa

- string network is in scaling
- new loops are formed from reconnection
- loops emit axions and collapse
- typical string length tracks Hubble
- average energy density tracks Hubble
- today: $O(10)$ strings per Hubble volume

How can we detect axion strings in the Universe today?

[image credit: Ken'ichi Saikawa]

birefringence
from axion strings

How could we detect an axion string?

[Harvey & Naculich (1989)], [Carroll, Field, Jackiw (1990,91)], [Harari, Sikivie (1992)]
 [Fedderke, Graham, Rajendran (2019)], [Agrawal, Hook, Huang (2019)]
 [Yin, Dai, Ferraro (2021) & (2023)]

assume interaction
 with electromagnetism:
 standard Chern-Simons coupling

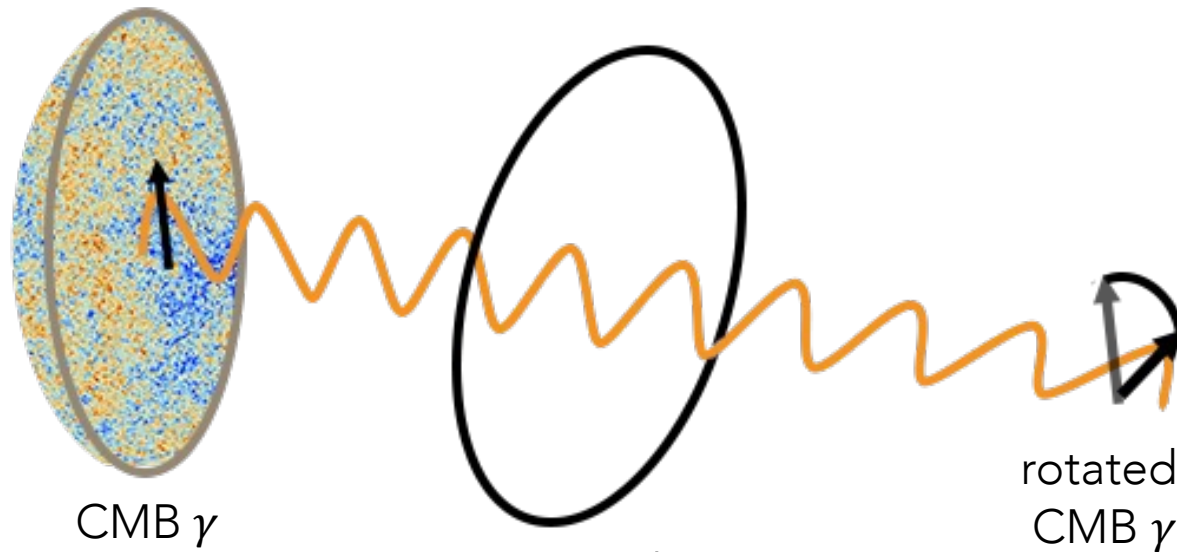
$$\mathcal{L}_{\text{int}} = -\frac{1}{4} g_{a\gamma\gamma} a F \tilde{F}$$

$$g_{a\gamma\gamma} = -\mathcal{A} \frac{\alpha_{\text{em}}}{\pi f_a}$$

$$\mathcal{A} = \sum Q_{\text{PQ}} Q_{\text{em}}^2 \sim \#/9$$

axion-induced birefringence:
 an electromagnetic wave
 traveling through a varying axion field
 has its plane of polarization rotated

$$\alpha = \frac{1}{2} g_{a\gamma\gamma} \int_C dX^\mu \partial_\mu a(X)$$



rotation angle

$$\alpha = g_{a\gamma\gamma} \pi f_a$$

$$\equiv -\mathcal{A} \alpha_{\text{em}}$$

$$\approx -0.42^\circ \mathcal{A}$$

axion string loop

$$\Delta a = 2\pi f_a$$

* birefringence can be measured through E-B cross correlation

The loop-crossing model

Assumptions

- All loops are circles
- Randomize loop orientation
- Randomize loop location in space
- All loops same radius at any time
- Loop radius evolves tracking Hubble

$$R(t) = \zeta_0 / H(t)$$

- Number of loops tracks Hubble

$$\rho(t) = \xi_0 \mu(t) H(t)^2$$

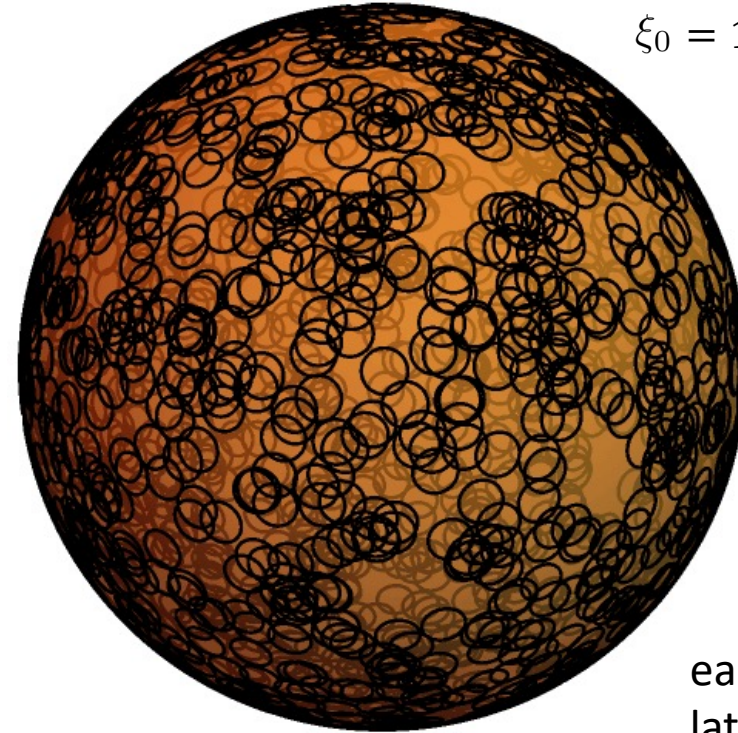
Model Parameters

$$\{m_a, \mathcal{A}, \zeta_0, \xi_0\}$$

loop-crossing model

$$\zeta_0 = 1.0$$

$$\xi_0 = 1.0$$



early time -> small loops
late time -> large loops

Expected birefringence signal

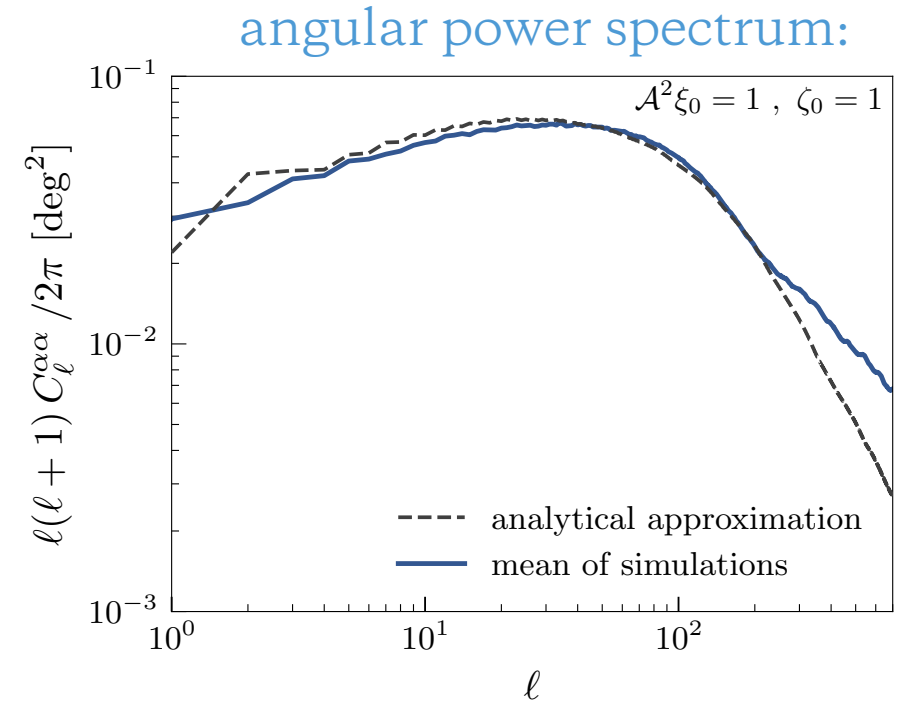
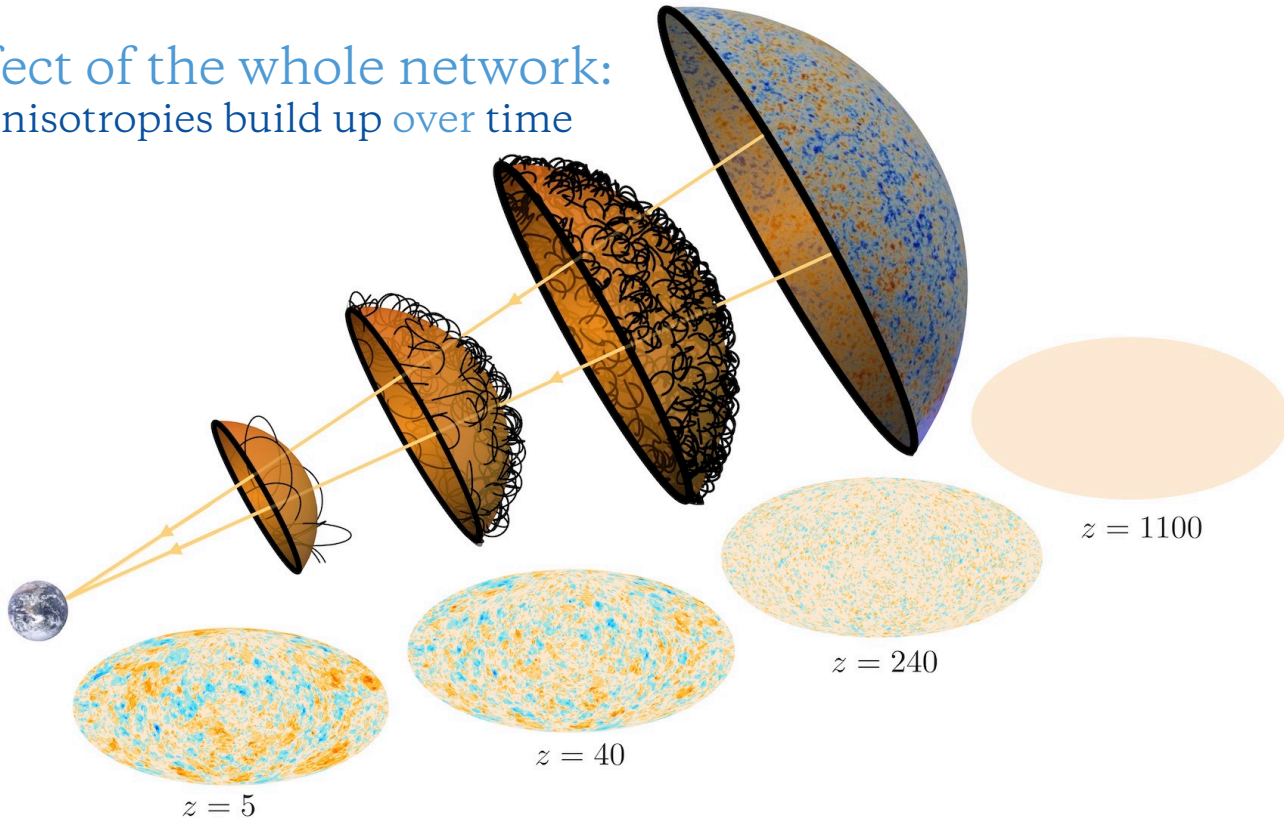
[Jain, AL, Amin, arXiv:2103.10962]

[Jain, Hagimoto, AL, Amin, arXiv:2208.08391]

loop-crossing model:
a network of circular loops

\mathcal{A} = dimensionless axion-photon coupling
 ξ_0 = dimensionless loop density (Hubble units)
 ζ_0 = dimensionless loop length (Hubble units)

effect of the whole network:
anisotropies build up over time



approx. scale invariant up to $l \sim 100$

degeneracy: $\langle \alpha \alpha \rangle \sim \mathcal{A}^2 \xi_0$

* need $m_a \lesssim 3H_{\text{cmb}} \approx 10^{-28}$ eV for the network to survive until after recombination

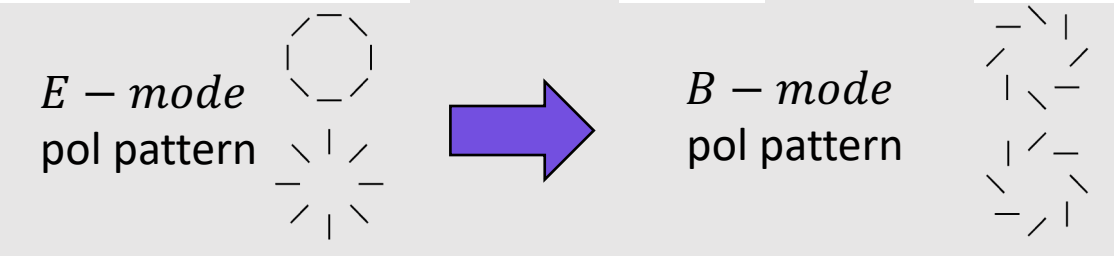
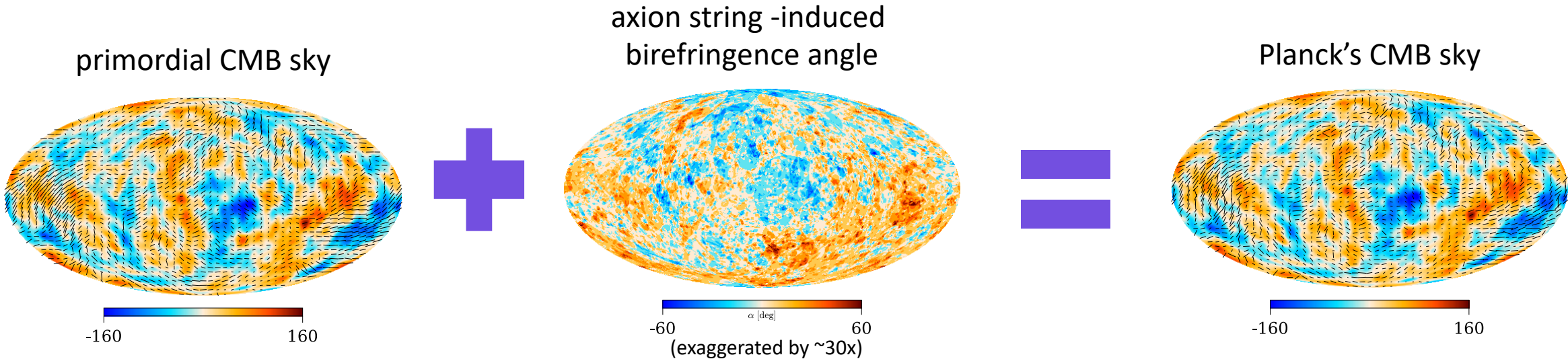
Effect on CMB polarization

How does birefringence affect the CMB's temperature and polarization?

$$T(\hat{n}) \rightarrow T(\hat{n})$$

$$[Q \pm iU](\hat{n}) \rightarrow [(Q \pm iU)e^{\pm 2i\Delta\Phi}](\hat{n})$$

see talk by :
Fumihito Naokawa



Signal of axion string-induced cosmological birefringence

$$\begin{cases} \langle TB \rangle \neq 0 \\ \langle EB \rangle \neq 0 \end{cases}$$

$$C_\ell^{EB} \sim \sin(4\Delta\Phi) (C_\ell^{EE} - C_\ell^{BB})$$

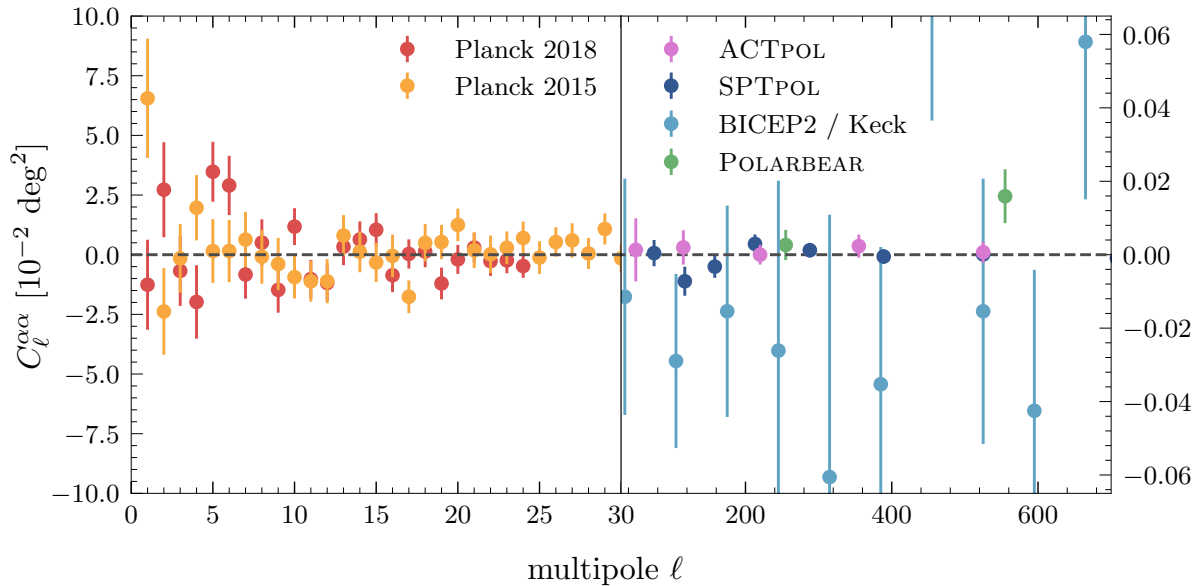
Constraints from anisotropic birefringence

[Jain, AL, Amin, arXiv:2103:10962]

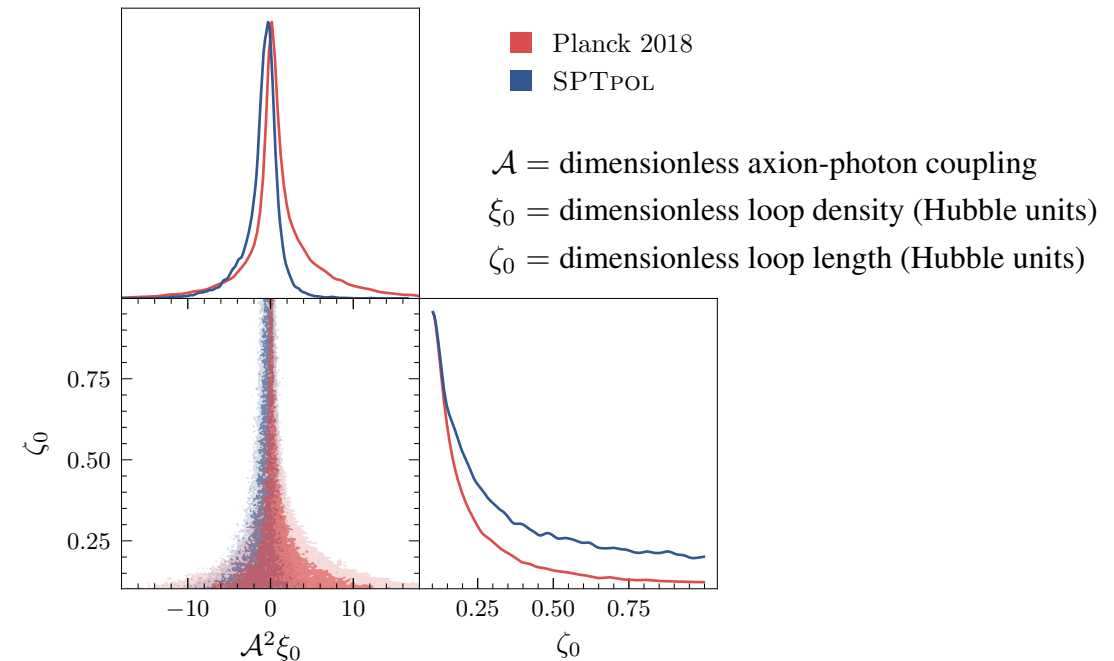
[Jain, Hagimoto, AL, Amin, arXiv:2208.08391]

see also: Yin, Dai, & Ferraro (2111.12741)

measurements of CMB polarization:
no evidence for anisotropic birefringence



a constraint on axion strings networks
& their coupling to electromagnetism:



constraints:

SPTPOL: $\mathcal{A}^2 \xi_0 < 3.7$ at 95% CL

already valuable!

if $\mathcal{A} = \frac{1}{3}$ and $\xi_0 = 30$ then $\mathcal{A}^2 \xi_0 \approx 3.3$

what about
isotropic
birefringence

Are strings responsible for isotropic birefringence?

[Jain, Hagimoto, AL, Amin, arXiv:2208.08391]

see talk by :
Eiichiro
Komatsu

reported detection of isotropic birefringence:
same rotation angle across the whole sky
(using *Planck* & *WMAP* data)

$$\alpha_{00} = -1.21^{\circ} {}^{+0.33^{\circ}}_{-0.32^{\circ}} \text{ (68\% CL)}$$

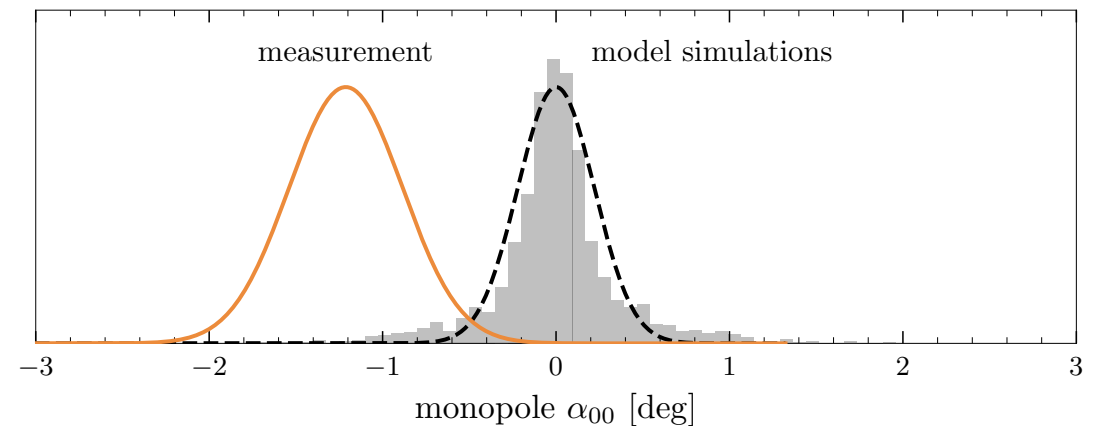
[Minami & Komatsu (2020)]

[Diego-Palazuelos et. al. (2022)]

[Eskilt (2022)], [Eskilt & Komatsu (2022)]

[Eskilt et. al. (2023)]

our conclusion: the isotropic signal is in tension
with limits on anisotropic BF if they both arise
from axion-string induced birefringence



note that: $\beta = -\alpha_{00}/\sqrt{4\pi} \approx 0.34^{\circ}$

Are strings responsible for isotropic birefringence?

[Ferreira, Gasparotto, Hiramatsu, Obata, & Pujolas (2023)]

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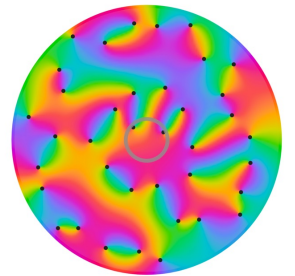
loopholes allowing large iso-BF

(1) environmental effects
a nearby loop in our Hubble volume
would dominate the isotropic signal

(2) Hubble-scale gradients
the massless axion field is expected to be
inhomogeneous on the Hubble scale

(3) late-forming network
if the string network is not present just after
recombination, the small-scale BF is suppressed

**please ask:
Ippei Obata**



note that: $\beta = -\alpha_{00}/\sqrt{4\pi} \approx 0.34^{\circ}$

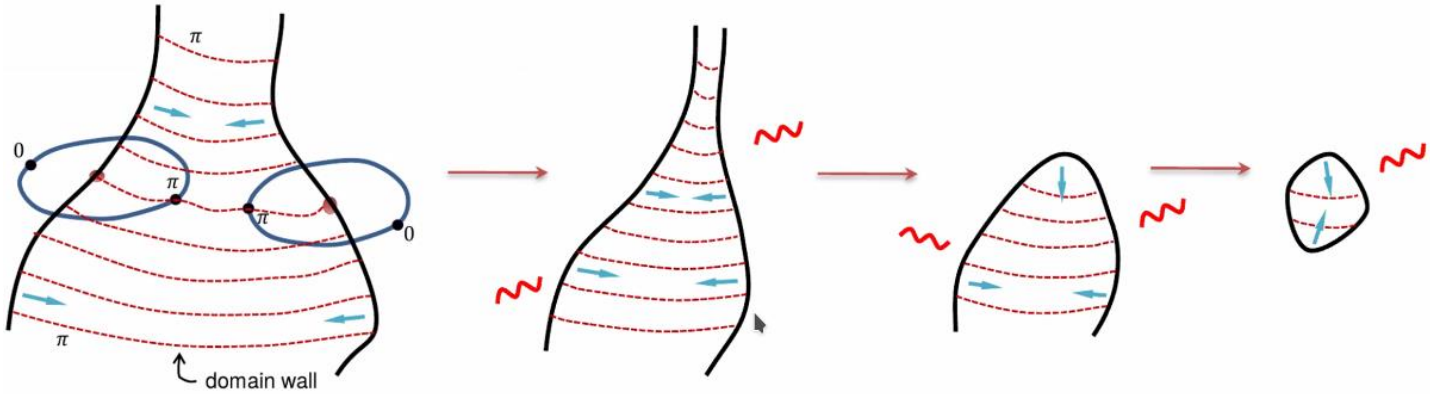
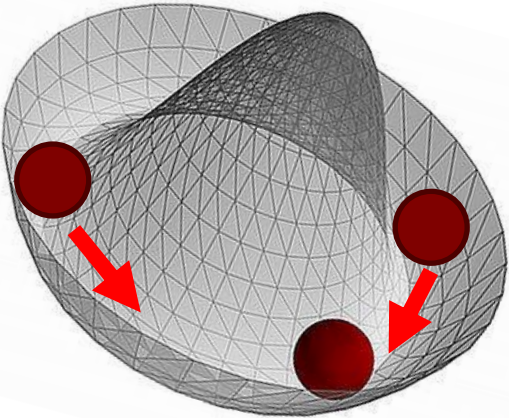
effect of
varying ALP mass

Collapse of the string-wall network

[Jain, Hagimoto, AL, Amin, arXiv:2208.08391]

Axion strings become connected together by domain walls

... the string-wall network collapses (for $N_{dw} = 1$)



let's consider:

$$\begin{cases} m_a \lesssim 3H_{\text{CMB}} \simeq 3 \times 10^{-29} \text{ eV} & \text{(string network survives until after recombination)} \\ m_a \gtrsim 3H_0 \simeq 5 \times 10^{-33} \text{ eV} & \text{(string network collapses before today)} \end{cases}$$

after the network collapses at redshift z_c the accumulation of birefringence is shut off

Impact on birefringence

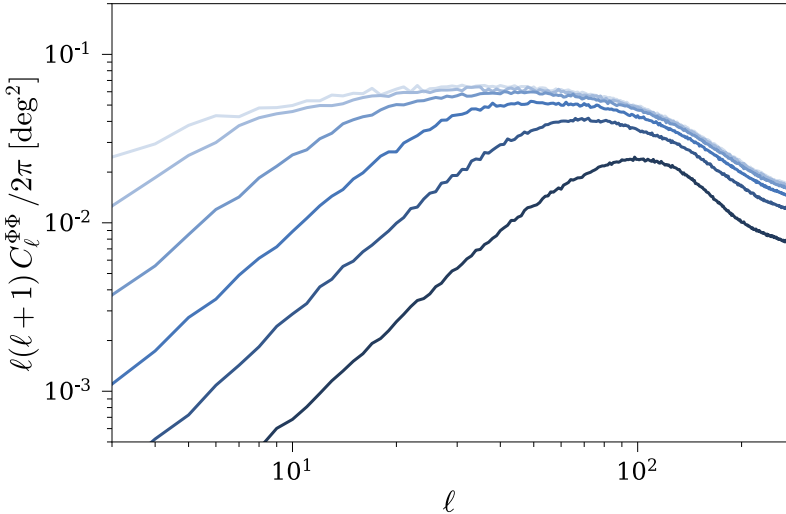
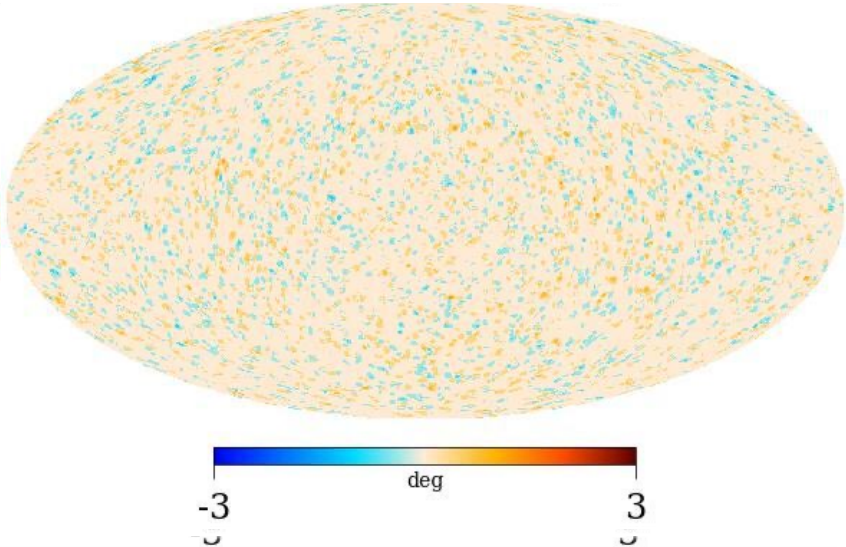
(assuming $N_{\text{DW}} = 1$)

raise the ALP mass
(network collapses earlier)

[Jain, Hagimoto, AL, Amin, arXiv:2208.08391]

see also: [Ferreira, Gasparotto, Hiramatsu, Obata, & Pujolas (2023)]

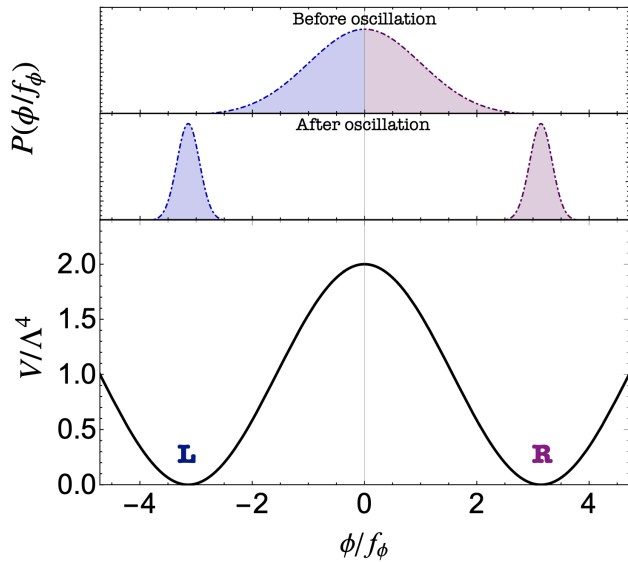
$$m_a = 2 \times 10^{-29} \text{ eV} \quad (z_c = 404)$$



strong scale dependence → possible to measure m_a

Complementary studies: stable axion domain walls

domain walls without strings
 expected if $H_{\text{inf}} \sim f_a$

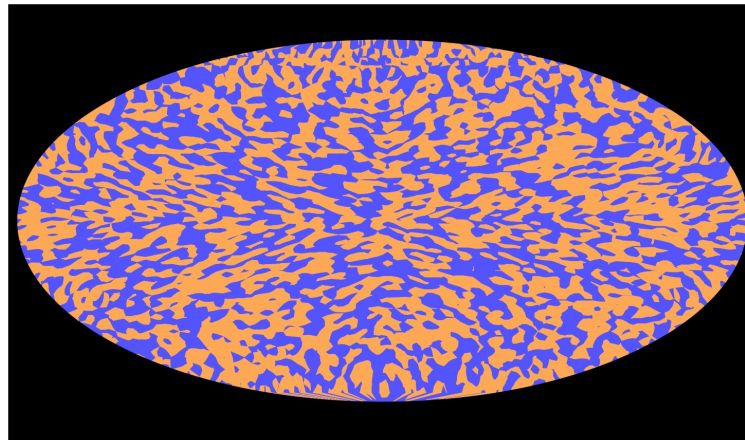


possible to evade DW problem
 imposes bound on mass & decay constant

$$\sigma_{\text{DW}} \simeq 8f_\phi^2 m_\phi \lesssim (1 \text{ MeV})^3,$$

$$f_\phi \lesssim 4 \times 10^9 \text{ GeV} \sqrt{\frac{10^{-20} \text{ eV}}{m_\phi}}.$$

birefringence signal
 independent of propagation



$$\Delta\Phi = 0 \quad \text{if LS}\gamma \text{ is from the vacuum } R$$

$$\Delta\Phi = c_\gamma \alpha \quad \text{if LS}\gamma \text{ is from the vacuum } L.$$

**possible to accommodate detection of
 isotropic BF and evade limits on anisotropic BF
 (no random-walk enhancement)**

[Takahashi & Yin (2020)]
 [Nakagawa, Takahashi, & Yamada (2021)]
 [Kitajima, Kozai, Takahashi, & Yin (2022)]
 [Gonzalez, Kitajima, Takahashi, & Yin (2022)]

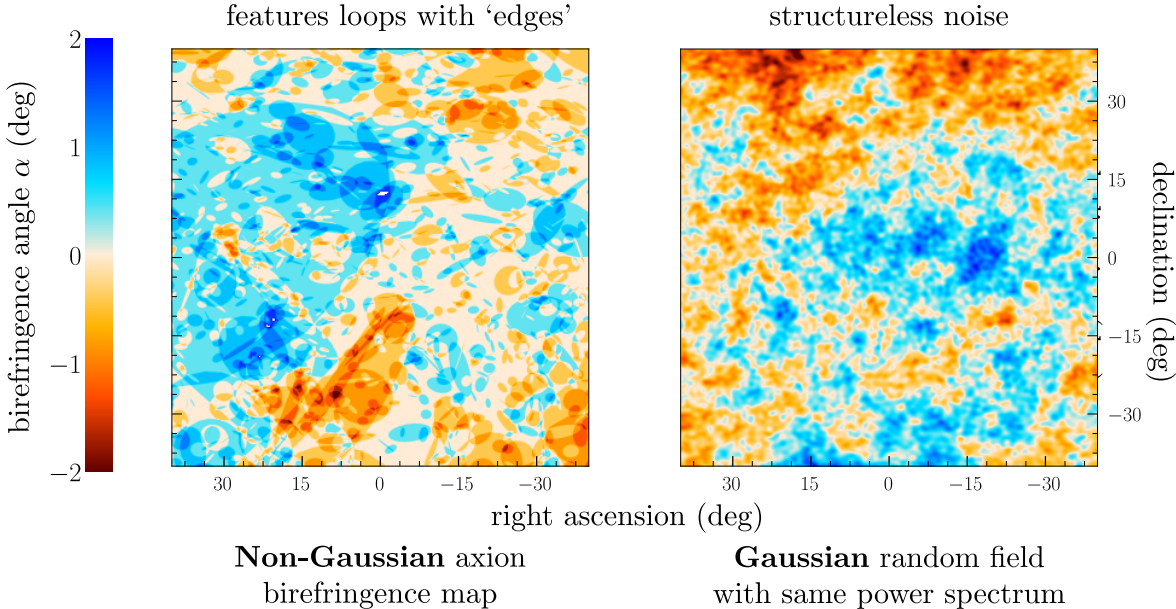
please ask:
Fuminobu Takahashi
Diego Gonzalez
Naoya Kitajima
Shota Nakagawa
Masaki Yamada
Wen Yin

signatures of
non-Gaussianity

Birefringence non-Gaussianity

[Hagimoto & AL, arXiv:2306.07351]
see also: Yin, Dai, Ferraro (2305.02318)

axion-string induced birefringence:
loop-like features are visibly non-Gaussian

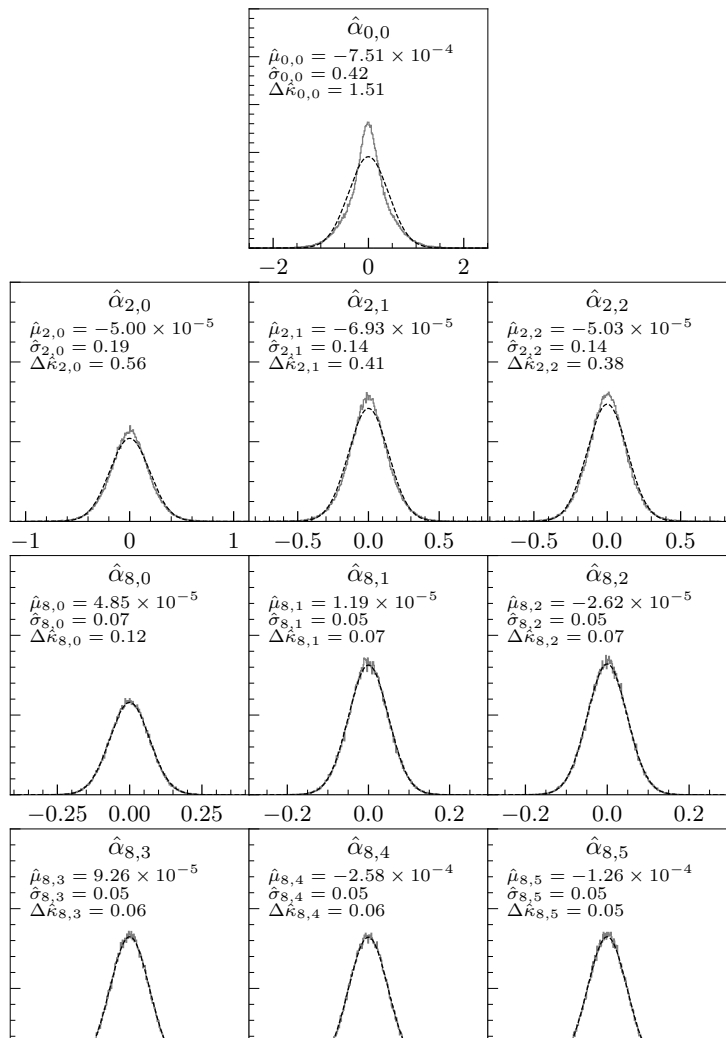


How to best quantify the non-Gaussian birefringence and develop tests to extract these features from the data?

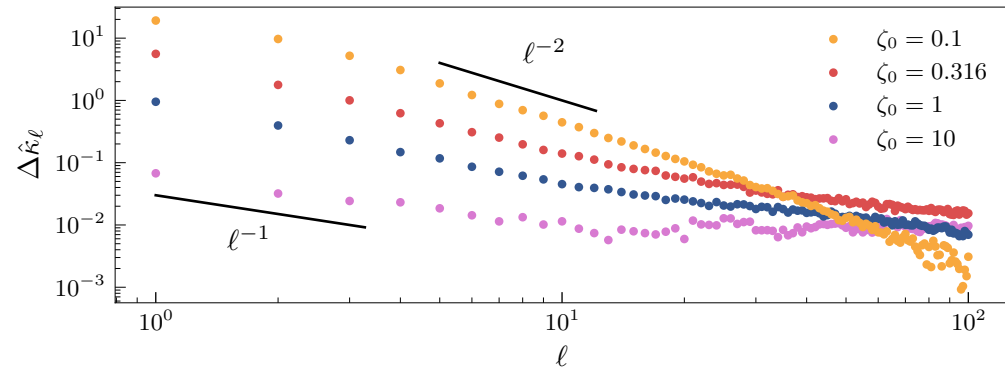
Birefringence non-Gaussianity

[Hagimoto & AL, arXiv:2306.07351]

excess kurtosis
can be O(1) at low multipoles

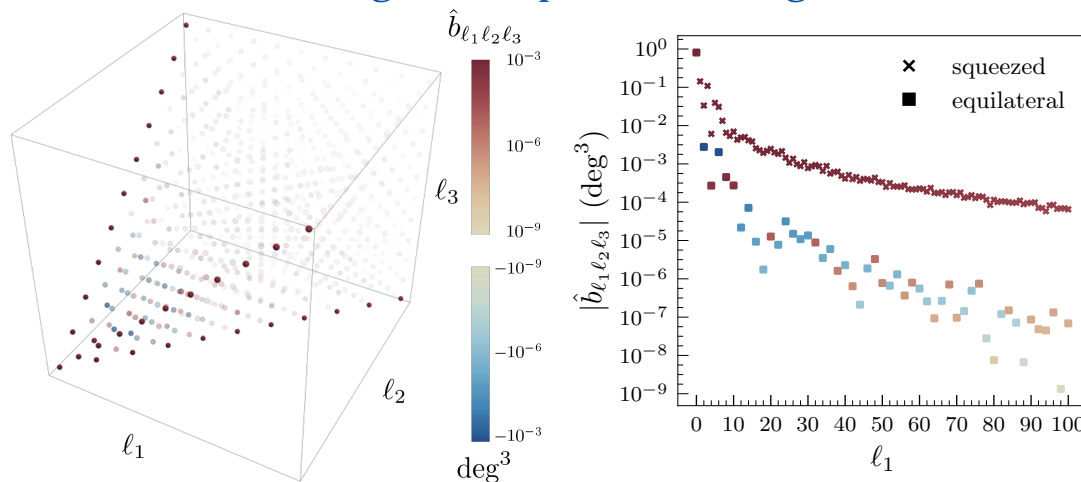


scaling with multipole index
matches well an analytic toy model



smaller NG on smaller scales (b/c of central limit theorem)

bispectrum
largest in squeezed triangle form



NG information breaks the $A^2 \xi_0$ degeneracy

a measurement of non-G will help to tell apart axion strings & other sources of birefringence

Probing non-Gaussianity with scattering transform

Yin, Dai, Ferraro (2023)

std. method
power spectrum

signal: $I_0(\mathbf{x})$
 plane wave: $\phi_{\mathbf{k}}(\mathbf{x})$
 $P_{\mathbf{k}}(\mathbf{x}) = \langle |I_0 * \phi_{\mathbf{k}}|^2 \rangle(\mathbf{x})$

new method
scattering transform

wavelet: $\psi^{j,l}(\mathbf{x})$

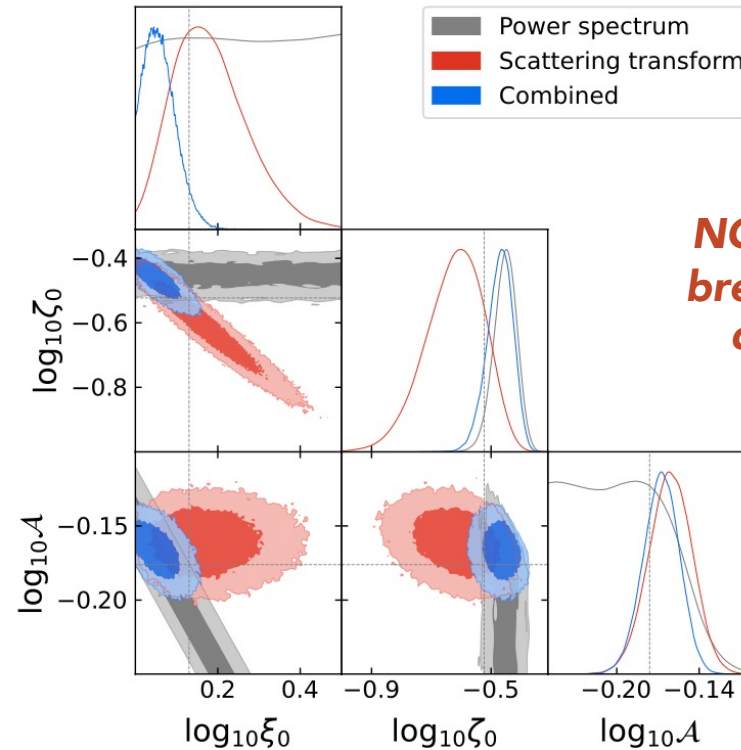
$$I_1^{j,l}(\mathbf{x}) = \langle |I_0 * \psi^{j,l}|^2 \rangle(\mathbf{x})$$

$$I_2^{j_1,l_1,j_2,l_2}(\mathbf{x}) = \langle |I_1^{j_1,l_1} * \psi^{j_2,l_2}|^2 \rangle(\mathbf{x})$$

$$s_1^j = \langle I_1^{j,l} \rangle_{\mathbf{x},l}$$

$$s_2^{j_1,j_2} = \langle I_2^{j_1,l_1,j_2,l_2} \rangle_{\mathbf{x},l_1,l_2}$$

comparison
pow-spec vs. scatt-transform



**NG information
breaks the $A^2 \xi_0$
degeneracy**

(b) $A^2 \xi_0 = 0.6$, $A = 2/3$, $\zeta_0 = 0.3$

Machine learning for axion strings

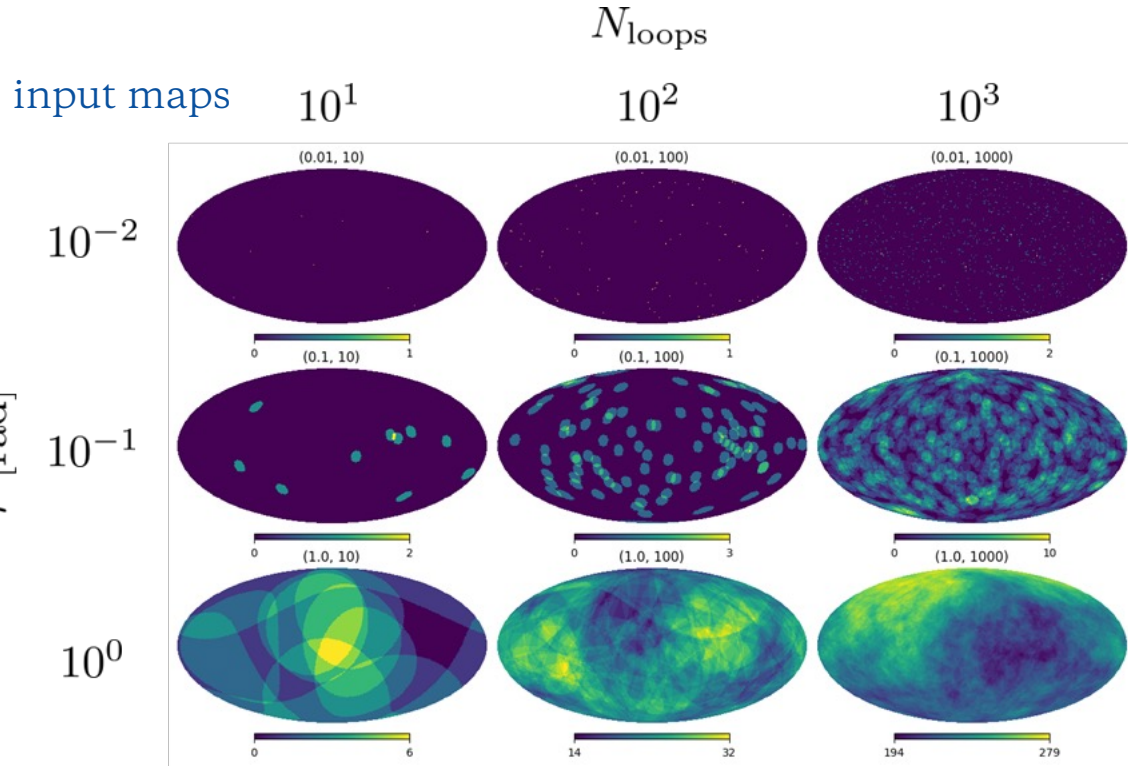
--- early stages ---

package: DeepSphere (Python)
architecture: 3 conv+pool layers

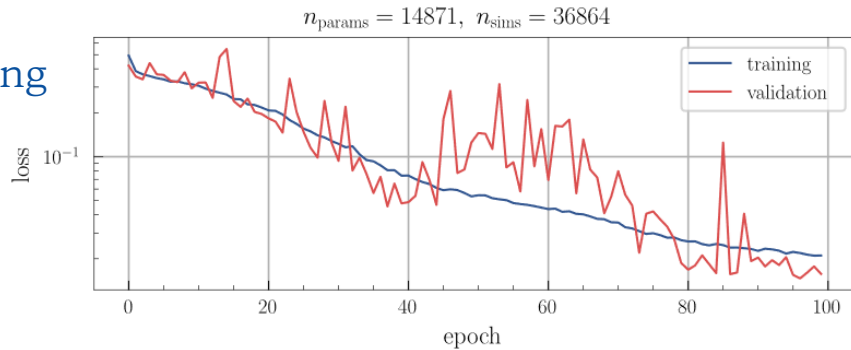
goal: to train an AI to identify features of axion strings in CMB polarization maps



Ray Hagimoto
(Rice U grad)

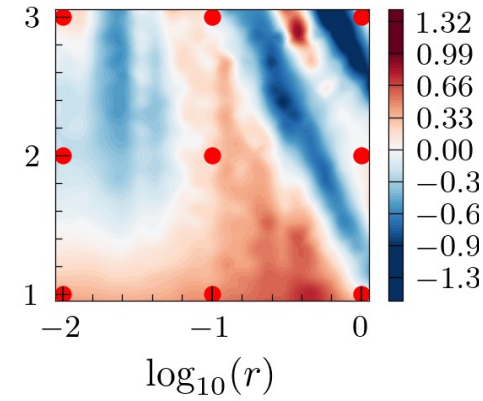


training

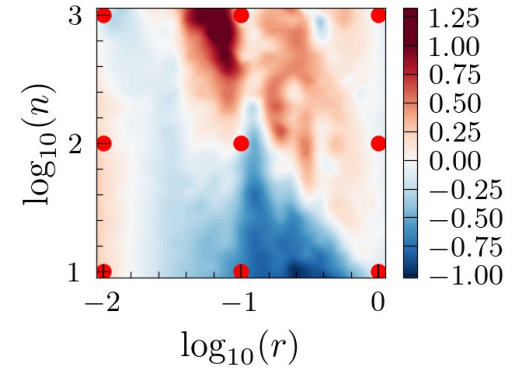


how well is it working? ... not bad!

average error in $\log_{10} n$



average error in $\log_{10} r$



summary
& conclusion

Summary

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- We use existing **measurements of anisotropic birefringence** (Planck, SPT, ...) to place constraints on this scenario. Next-generation telescopes (CMB-S4) will probe $O(1)$ electromagnetic anomaly coefficients and thereby probe the axion's UV embedding
- We find that it is difficult (but not impossible!) to reconcile the **detection of isotropic birefringence** with strong limits on anisotropic birefringence coming from axion strings
- We argue that measurements of anisotropic birefringence could not only reveal the presence of a hyper-light ALP in Nature, but also lead to a **measurement of its mass**
- Our ongoing work (very early stages) seeks to use machine learning techniques (spherical CNN) to detect the subtle signal of axion strings in CMB polarization data

backup slides



Projected sensitivity

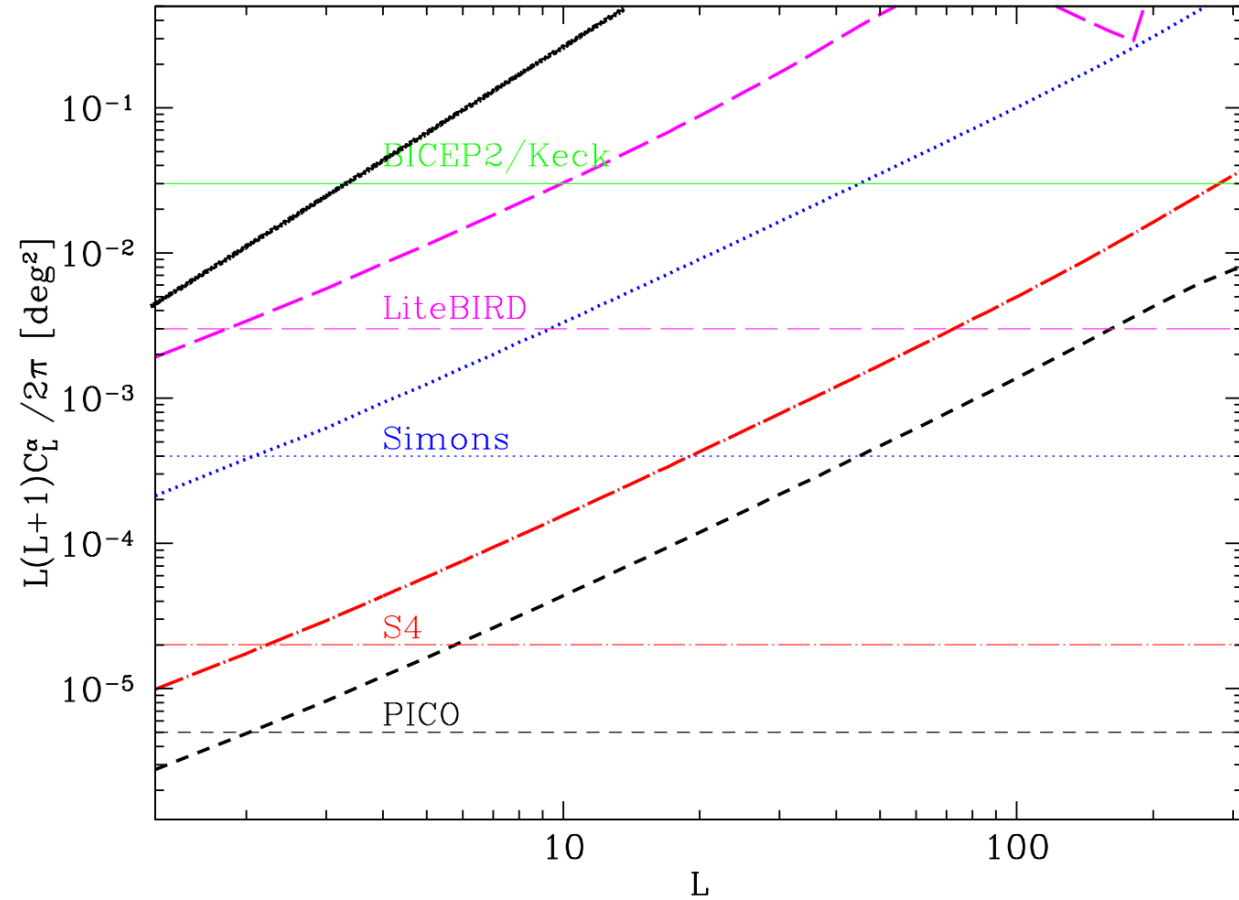
future telescopes
probes of isotropic + aniso. birefringence

Current			LiteBIRD			SO			CMB-S4-like			PICO		
α	A_α	$\sqrt{C_2^\alpha}$	α	A_α	$\sqrt{C_2^\alpha}$	α	A_α	$\sqrt{C_2^\alpha}$	α	A_α	$\sqrt{C_2^\alpha}$	α	A_α	$\sqrt{C_2^\alpha}$
/	10^{-2}deg^2	$/$	/	10^{-3}deg^2	$/$	/	10^{-4}deg^2	$/$	/	10^{-5}deg^2	$/$	/	10^{-5}deg^2	$/$
-	-	-	1.3	2.7	0.9	0.56	3	0.29	0.1	1.4	0.065	0.05	0.4	0.035
-	-	-	1.5	3.3	1.0	0.66	4	0.35	0.11	2.0	0.08	0.06	0.5	0.04
-	-	-	1.4	3.5	1.0	0.64	5.0	0.4	0.13	2.5	0.09	0.08	1.2	0.06
30	2	3	1.6	4.0	1.1	0.71	5.5	0.4	0.15	3.3	0.1	0.09	1.4	0.065

BLE II. Current and forecasted 68% CL bounds on the uniform and the anisotropic CPR parameters.

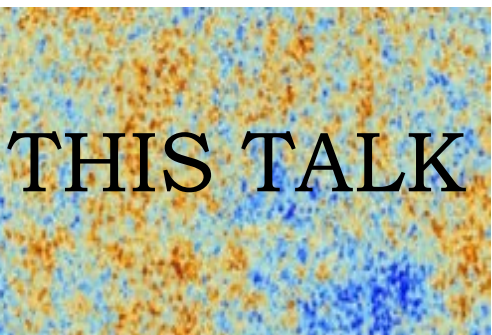
$$A_\alpha = L(L + 1)C_L^\alpha / 2\pi$$

diagonal = allows multipoles to vary independently
horizontal = restricts to a scale invariant spectrum

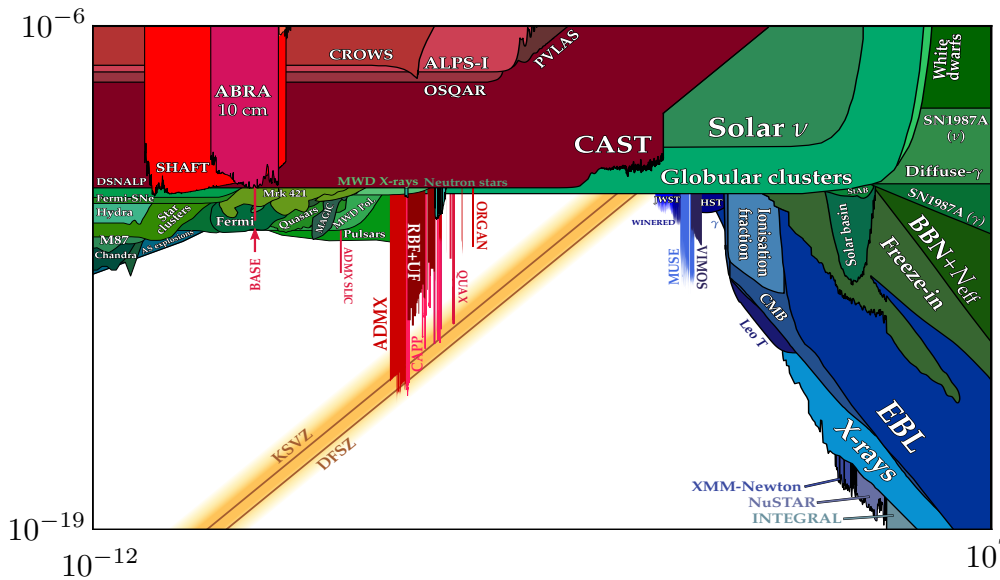
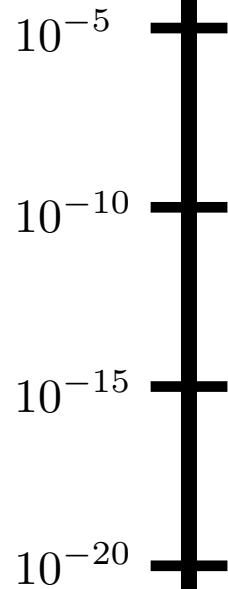


Theory landscape: axion-like particles

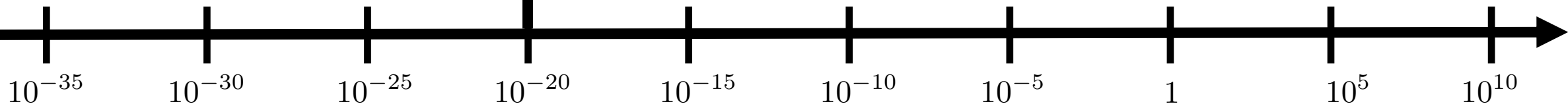
$$\mathcal{L} \supset \frac{1}{2}(\partial a)^2 - \frac{1}{2}m_a^2 a^2 - \frac{1}{4}g_{a\gamma\gamma} a F_{\mu\nu} \tilde{F}^{\mu\nu}$$



$g_{a\gamma\gamma} [\text{GeV}^{-1}]$

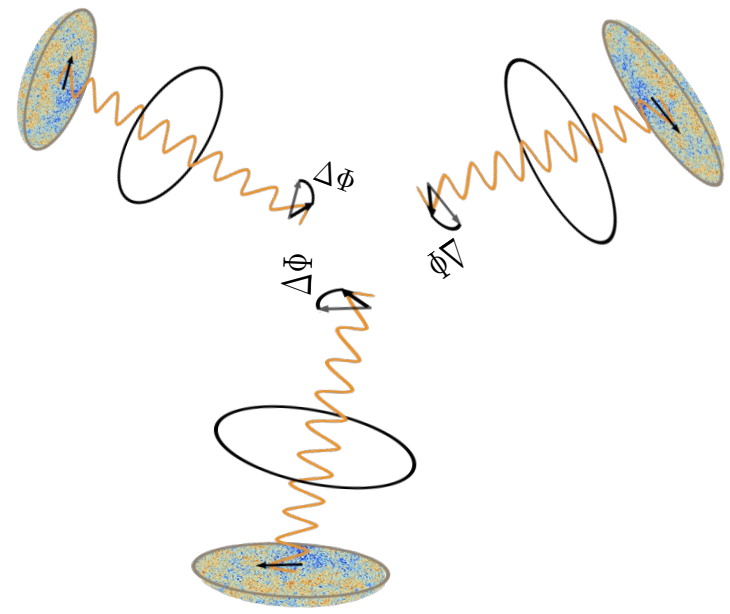


$m_a [\text{eV}]$



Net birefringence from the whole string network

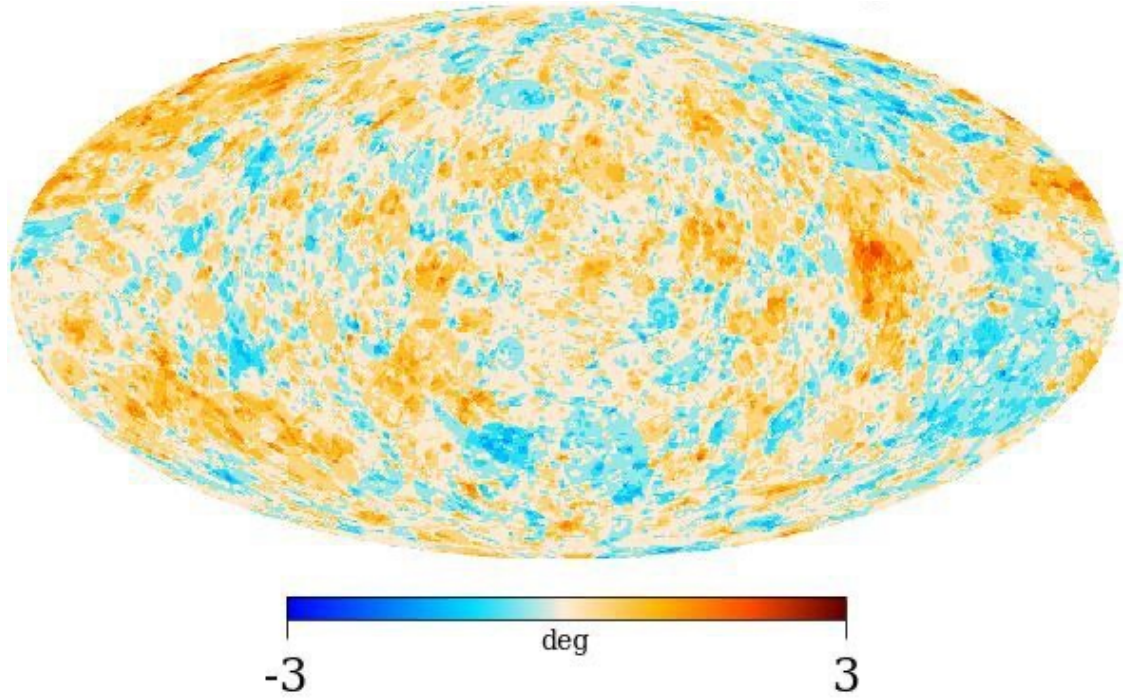
[Ray Hagimoto (2022)]



parameters:

$$\begin{cases} m_a = 0 \\ \mathcal{A} = 1 \\ \zeta_0 = 1 \\ \xi_0 = 1 \end{cases}$$

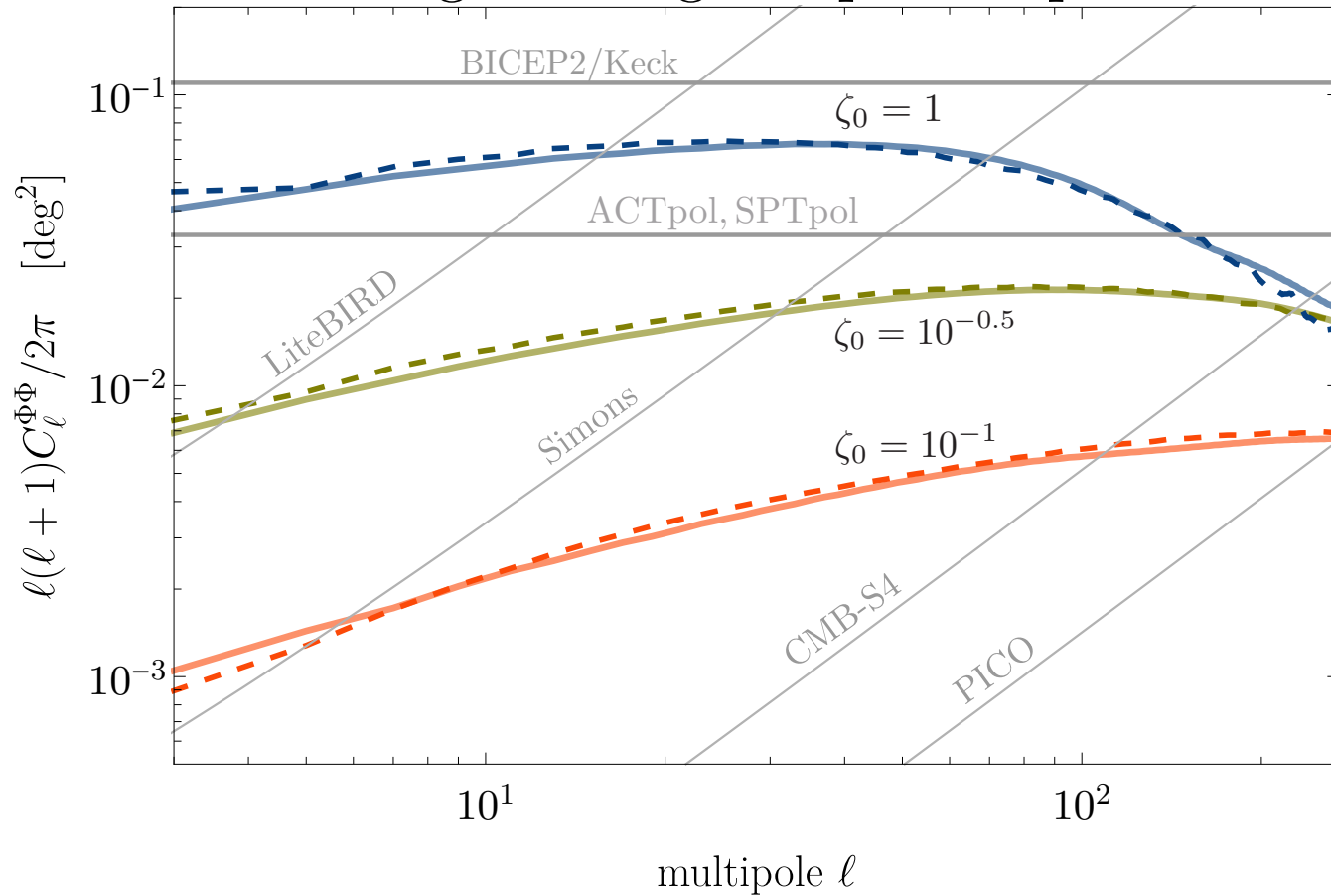
a map of $\Delta\Phi$ over the sky



Axion-string induced birefringence signal

[Jain, AL, Amin (2103.01962)]

birefringence angular power spectrum



assumes: $m_a = 0$ and $\xi_0 \mathcal{A}^2 = 1$

Key features

- Power spectrum is almost scale invar.
- Characteristic scale (l @ the peak) set by loop size at LSS
- Smaller loops (ζ_0) => weaker signal
- Trivial dependence on loop density (ξ_0) and anomaly coefficient (A) ... power scales with $\xi_0 A^2$

Testability

- Current telescopes (SPT/ACT) are already sensitive enough to test large loops ($\zeta_0=1$)
- Future surveys will be very powerful