

Dark Matter direct detection using Qubits

(1) *Phys. Rev. Lett.* 131 (2023) 21, 211001

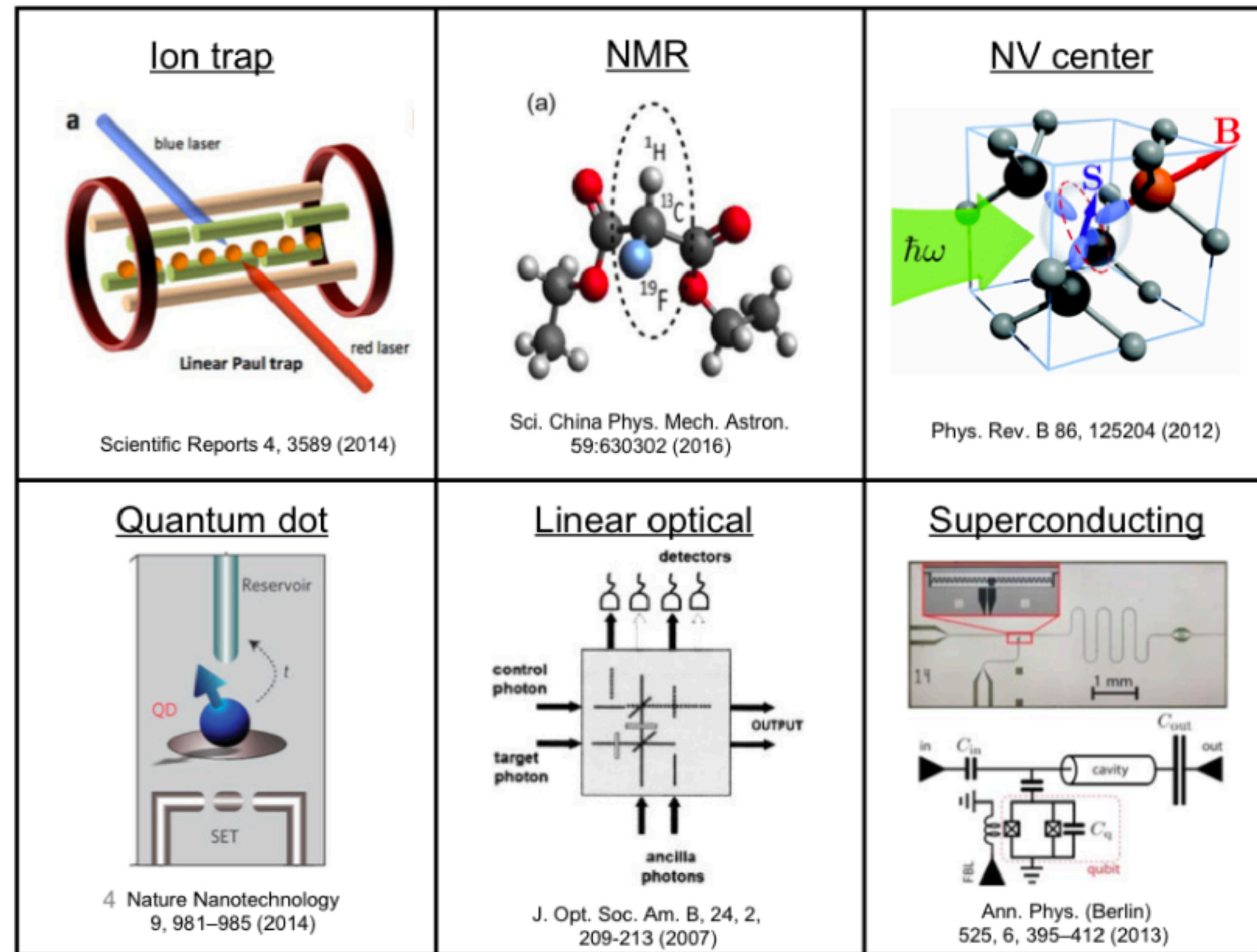
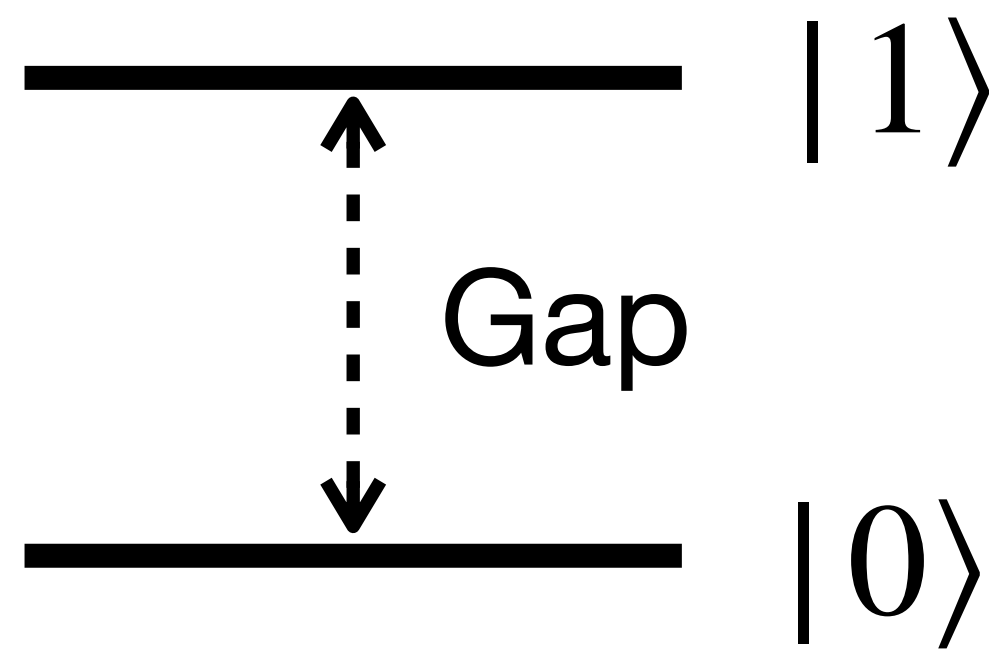
(2) *arXiv: 2311.10413*

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Background and Motivation

Qubit is two-level system developed for computation

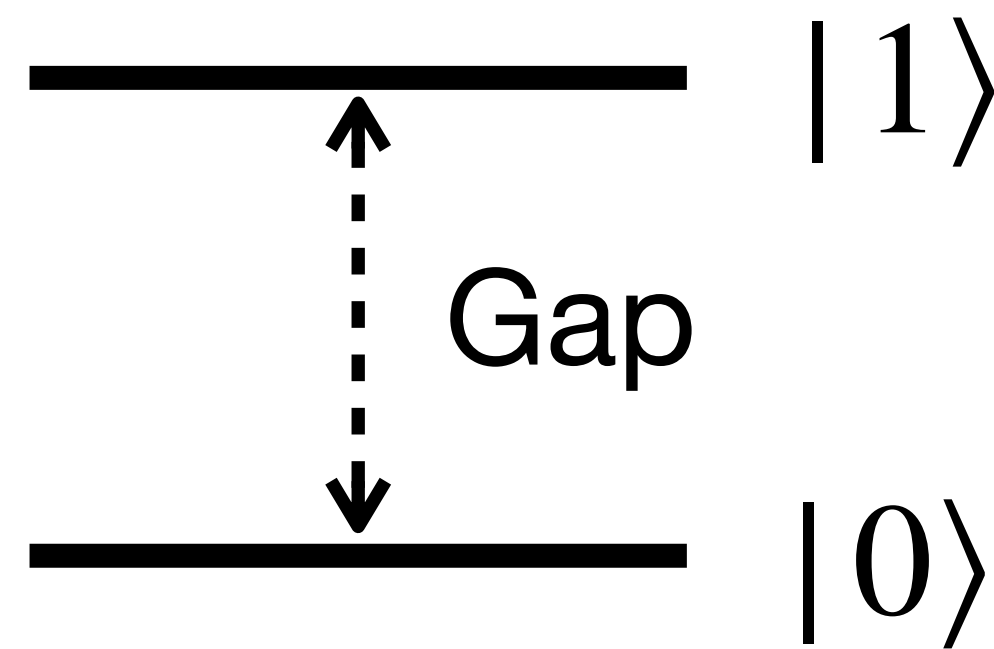


James Amundson, Elizabeth Sexton-Kennedy, EPJ Web of Conferences 214, 09010 (2019)

Background and Motivation

Qubit is two-level system developed for computation

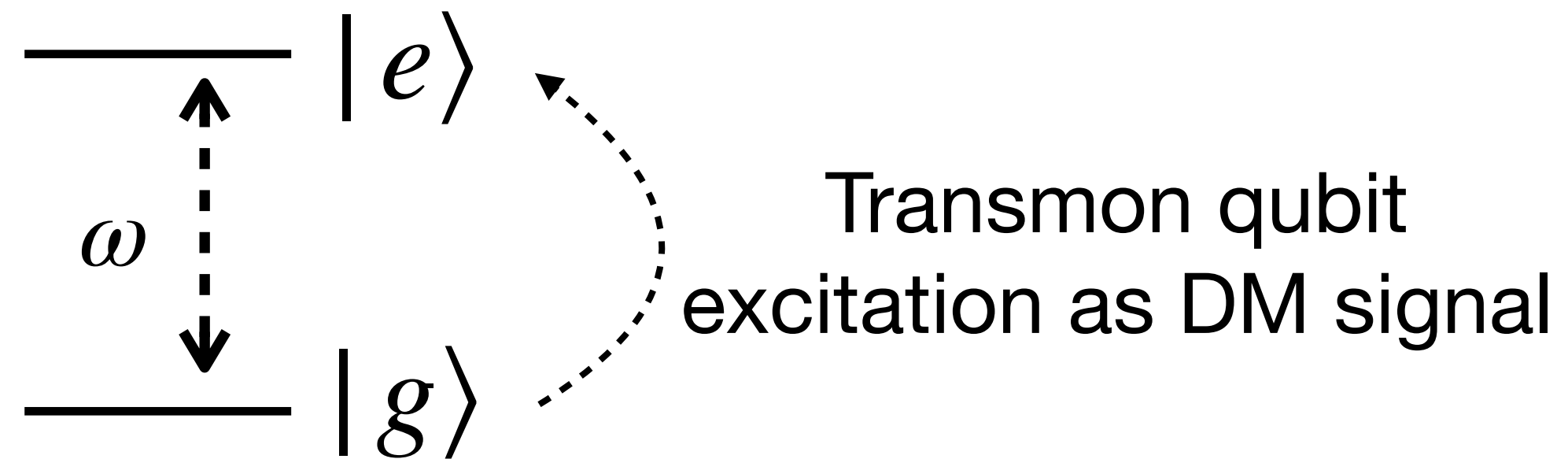
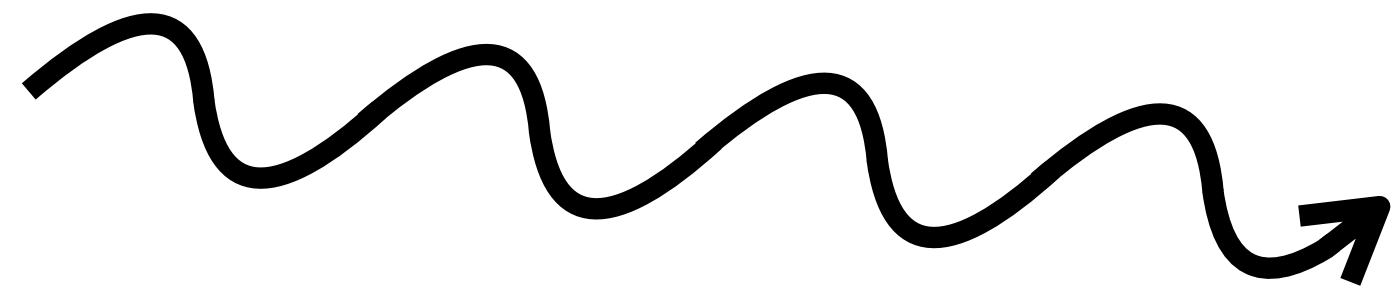
In the same time, it could be a good quantum sensor



- precise readout
- state controllability
- tunable energy gaps

Dark matter search using qubit

Hidden photon DM, mass m_X



Outline

- (1) Hidden photon DM search with transmon qubits

[Chen, Fukuda, Inada, Moroi, Nitta, TS, *Phys. Rev. Lett.* 131 (2023) 21, 211001]

- (2) Quantum computation to enhance detection

[Chen, Fukuda, Inada, Moroi, Nitta, TS, *arXiv: 2311.10413*]

Dark matter detection using transmon qubit excitation

[Chen, Fukuda, Inada, Moroi, Nitta, TS, *Phys. Rev. Lett.* 131 (2023) 21, 211001]

Hidden photon DM

In kinetic-mixing basis,

$$\mathcal{L} = -\frac{1}{4}F^{\mu\nu}F_{\mu\nu} - \underbrace{\frac{\epsilon}{2}F^{\mu\nu}X_{\mu\nu}}_{\text{Kinetic mixing}} - \frac{1}{4}X^{\mu\nu}X_{\mu\nu} + \frac{m_X^2}{2}X_\mu X^\mu + e\bar{\Psi}_e\gamma^\mu A_\mu\Psi_e$$

In mass-eigenstate basis ($A_\mu \rightarrow A_\mu + \epsilon X_\mu$),

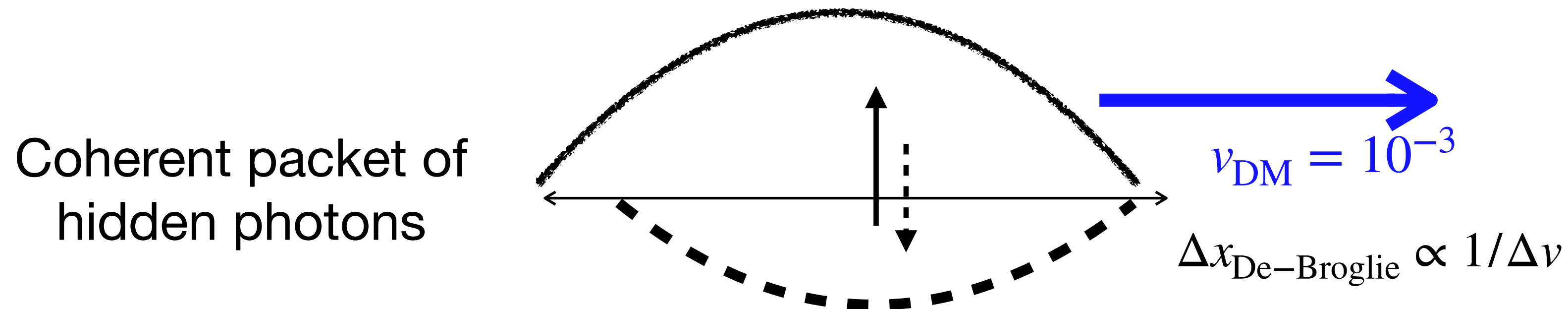
$$\mathcal{L} = -\frac{1}{4}F^{\mu\nu}F_{\mu\nu} - \frac{1}{4}X^{\mu\nu}X_{\mu\nu} + \frac{m_X^2}{2}X_\mu X^\mu + \underbrace{e\bar{\Psi}_e\gamma^\mu(A_\mu + \epsilon X_\mu)\Psi_e}$$

Direct interaction between DM and matters

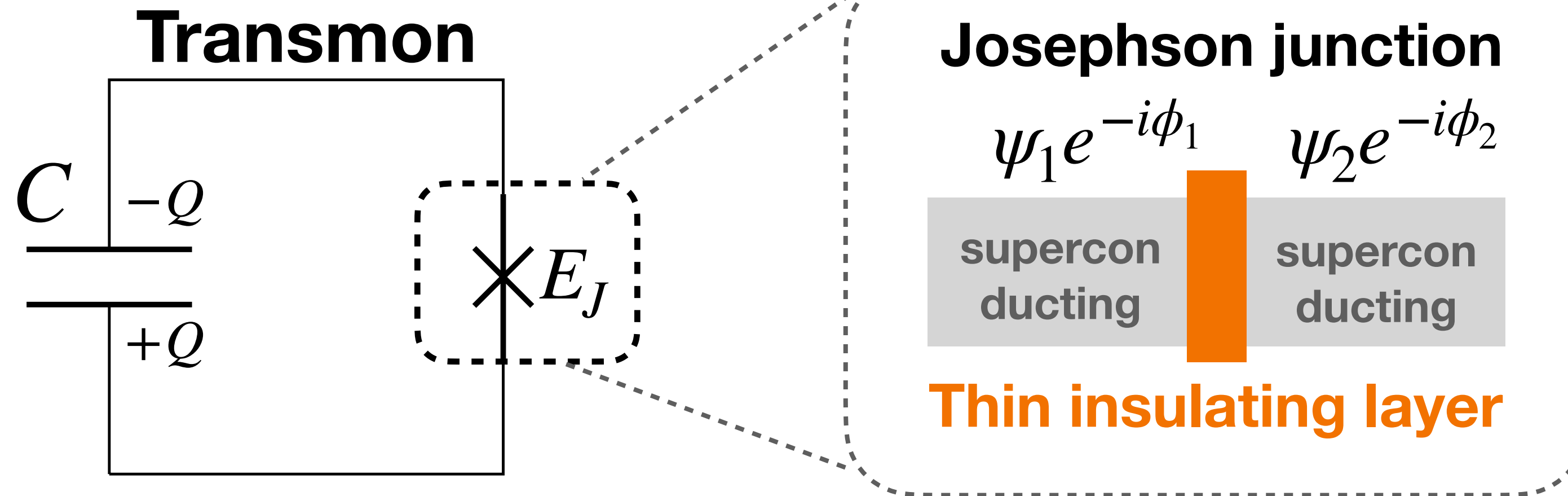
Hidden photon DM as effective E-field

<https://arxiv.org/abs/1711.10489>

- Mass $m_X \ll 1 \text{ eV} \Rightarrow$ [# particles within De-Broglie volume] $\gg 1$
- Hidden photon field $\vec{X} = \bar{X} \vec{n}_X \sin(m_X t + \alpha)$
- DM density $\rho_{\text{DM}} = \frac{1}{2} m_X^2 \bar{X}^2$ (local $\rho_{\text{DM}} \sim 0.45 \text{ GeV/cm}^3$)
- Effective E-field $\vec{E}^{(X)} := -\epsilon \partial_t \vec{X} = \bar{E}^{(X)} \vec{n}_X \cos(m_X t + \alpha)$; $\bar{E}^X = \epsilon \sqrt{2\rho_{\text{DM}}}$



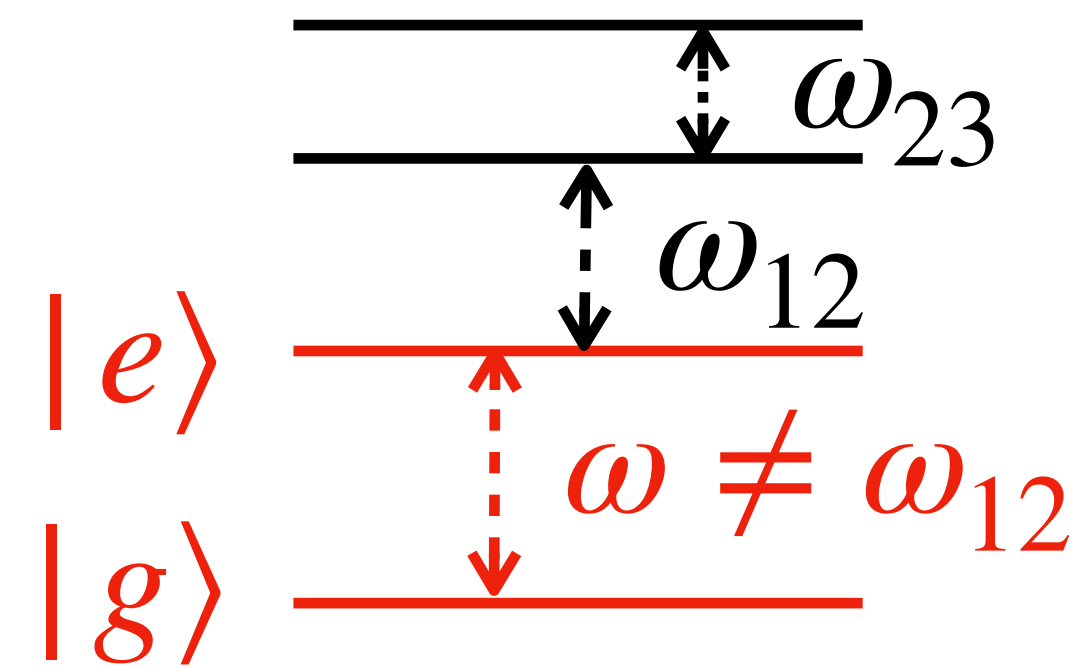
Transmon qubit as DM sensor



Hamiltonian

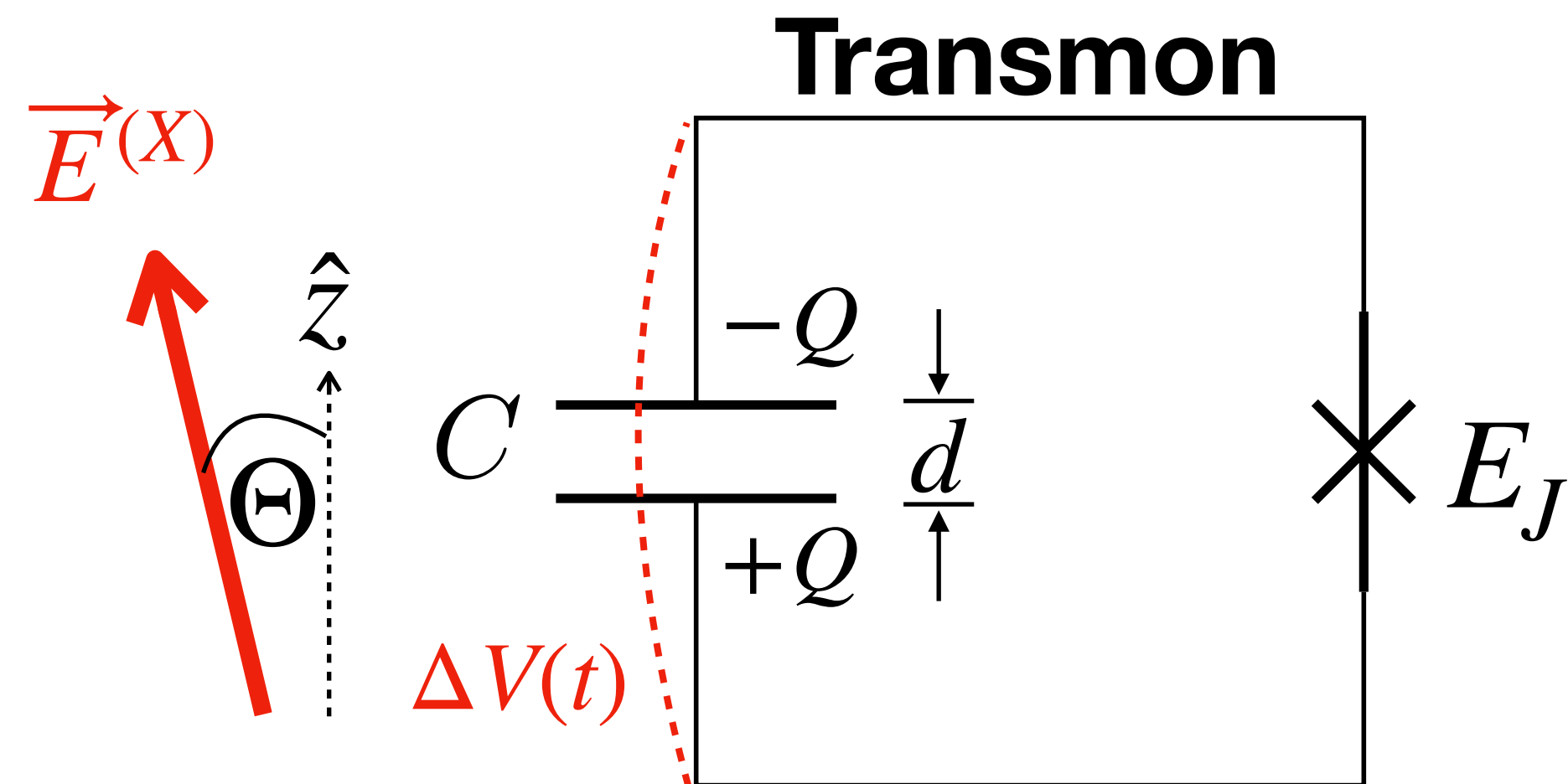
$$H_0 = Q^2/2C - E_J \cos \theta ; \theta \equiv \phi_1 - \phi_2$$

$$H_0 \simeq \omega |e\rangle\langle e| \text{ with } \omega = \sqrt{4e^2 E_J / C}$$



Gap not the same

Transmon qubit as hidden photon DM sensor



Induced voltage on capacitor:

$$\Delta V(t) = d \vec{E}^{(X)}(t) \cdot \hat{z}$$

Interaction hamiltonian:

$$H_{\text{int}} = Q \Delta V$$

$$H = \omega |e\rangle\langle e|$$

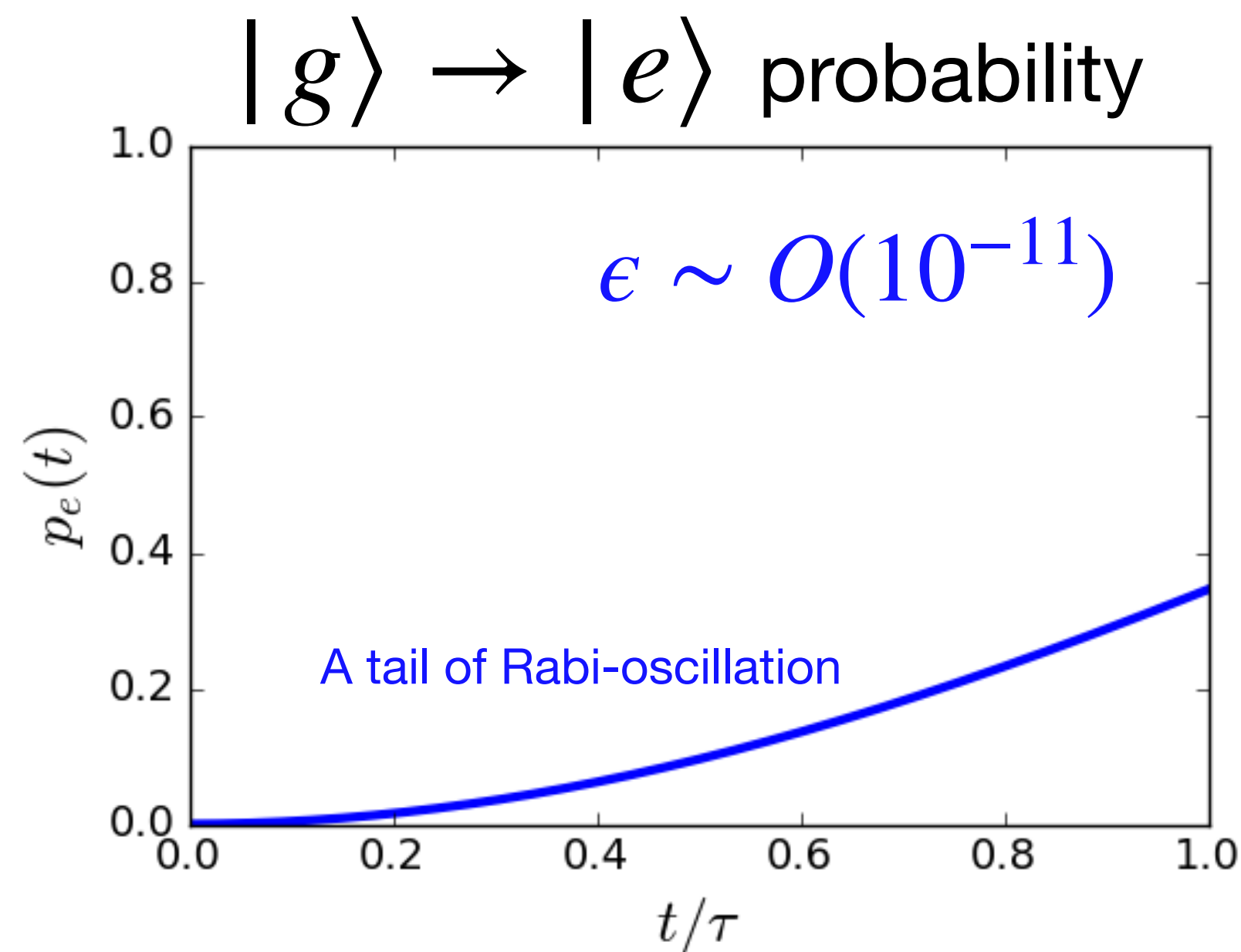
$$-2\eta \cos(m_X t + \alpha) (|g\rangle\langle e| + |e\rangle\langle g|)$$

For $m_X = \omega$, the Schrödinger Eq. is

$$\frac{d}{dt} \begin{pmatrix} \psi_g \\ \psi_e \end{pmatrix} = \begin{pmatrix} 0 & ie^{-i\alpha}\eta \\ ie^{i\alpha}\eta & 0 \end{pmatrix} \begin{pmatrix} \psi_g \\ \psi_e \end{pmatrix}$$

$$; \eta \equiv \sqrt{\omega C d \bar{E}^{(X)}} \cos \Theta / 2\sqrt{2}$$

Qubit direct excitation due to external field



$$p_e(\tau) \equiv |\langle g | e \rangle|^2 \simeq \delta^2 ; \delta \equiv \eta\tau$$

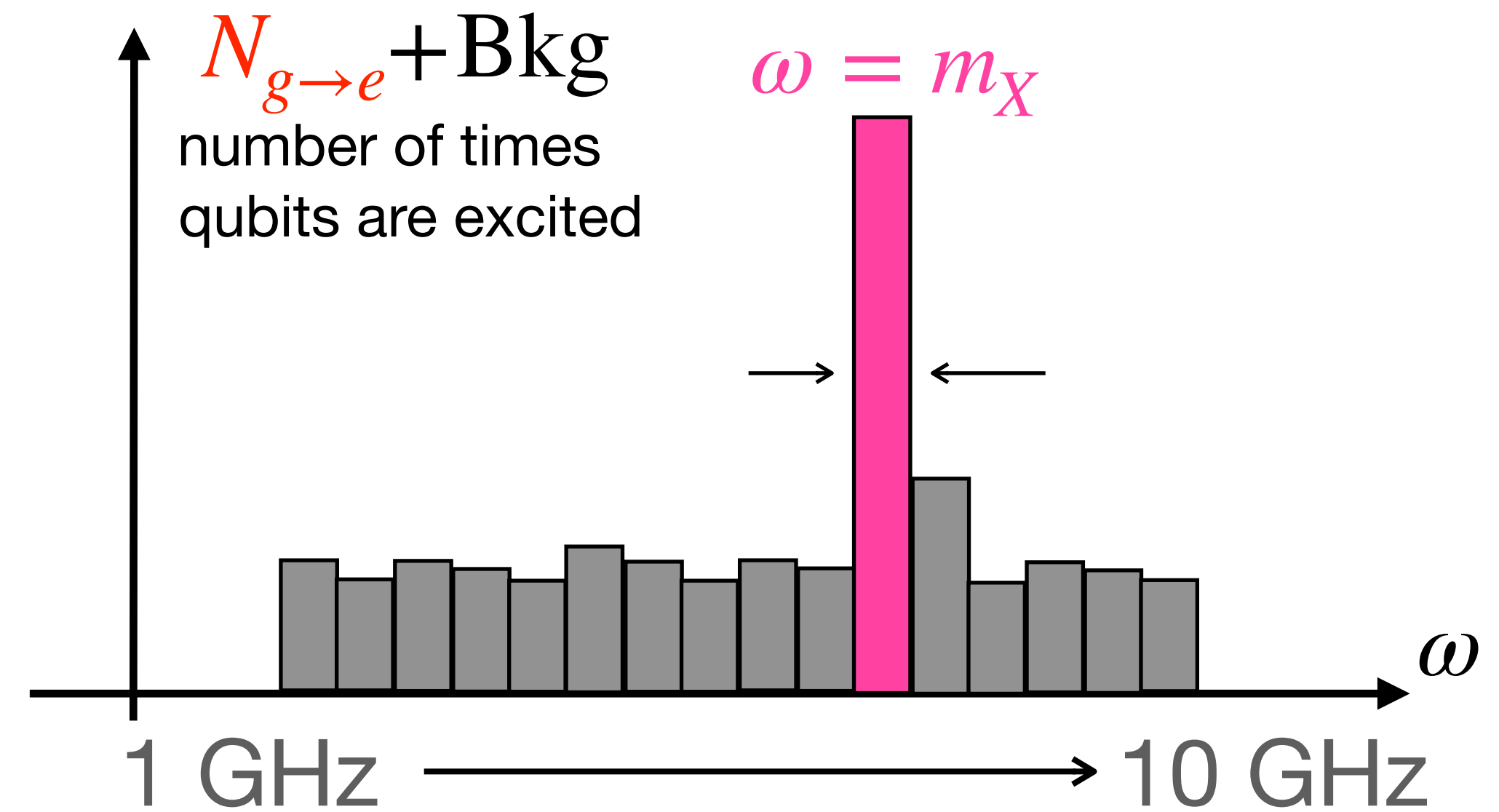
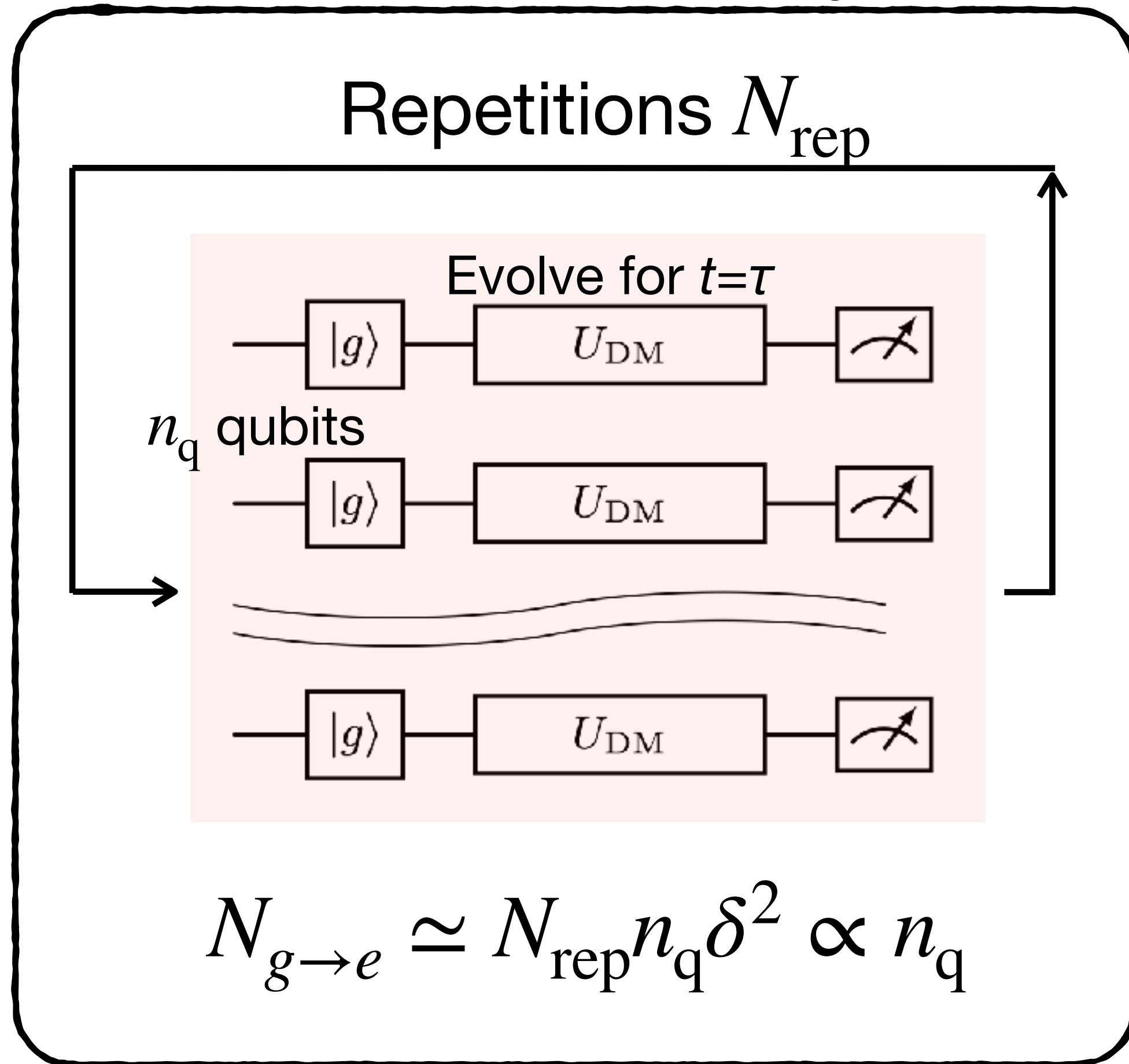
Gradual growth of probability

$$; \tau = \min\{\tau_{\text{DM}}, \tau_{\text{qubit}}\} \sim 100 \mu\text{s}$$

$$p_e(\tau) \simeq 0.12 \times \cos^2 \Theta \left(\frac{\epsilon}{10^{-11}} \right)^2 \left(\frac{f}{1 \text{ GHz}} \right) \\ \times \left(\frac{\tau}{100 \mu\text{s}} \right)^2 \left(\frac{C}{0.1 \text{ pF}} \right) \left(\frac{d}{100 \mu\text{m}} \right)^2 \left(\frac{\rho_{\text{DM}}}{0.45 \text{ GeV/cm}^3} \right)$$

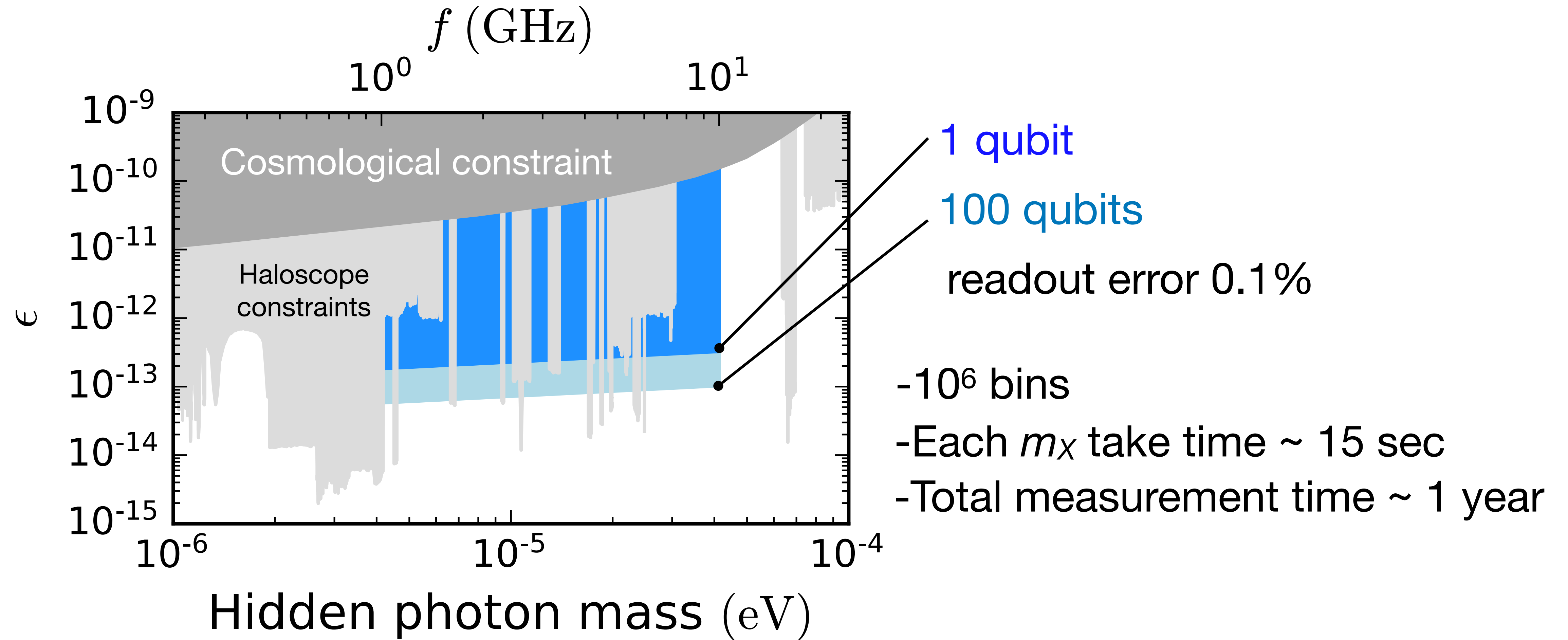
Process of measurement

For each frequency bin



Significance $\sigma = N_{g \rightarrow e} / \sqrt{\text{Bkg}}$
; $\text{Bkg} = p_{\text{noise}} n_q N_{\text{rep}}$

Sensitivity plot



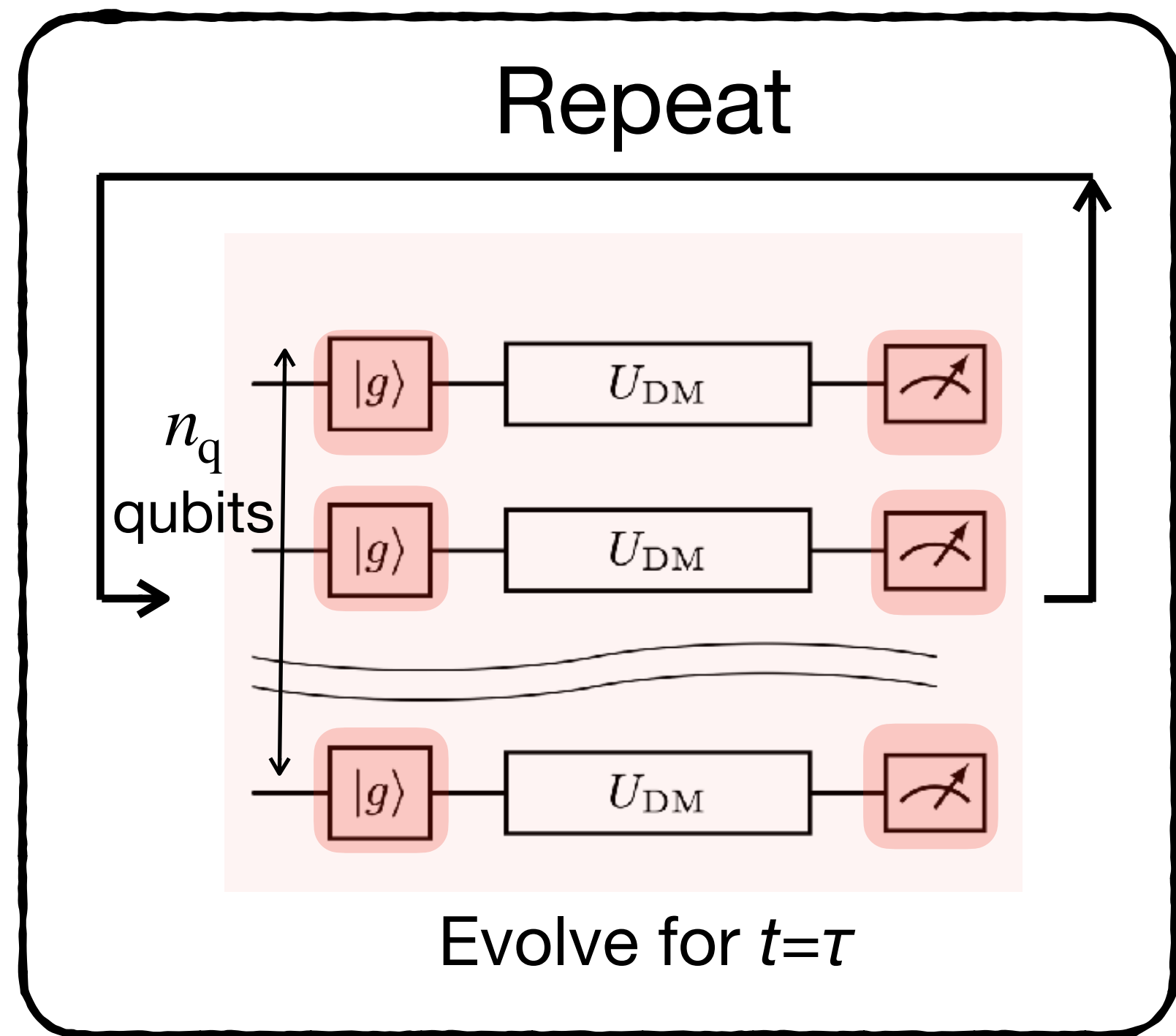
Enhancement using quantum computation

[Chen, Fukuda, Inada, Moroi, Nitta, TS, *arXiv: 2311.10413*]

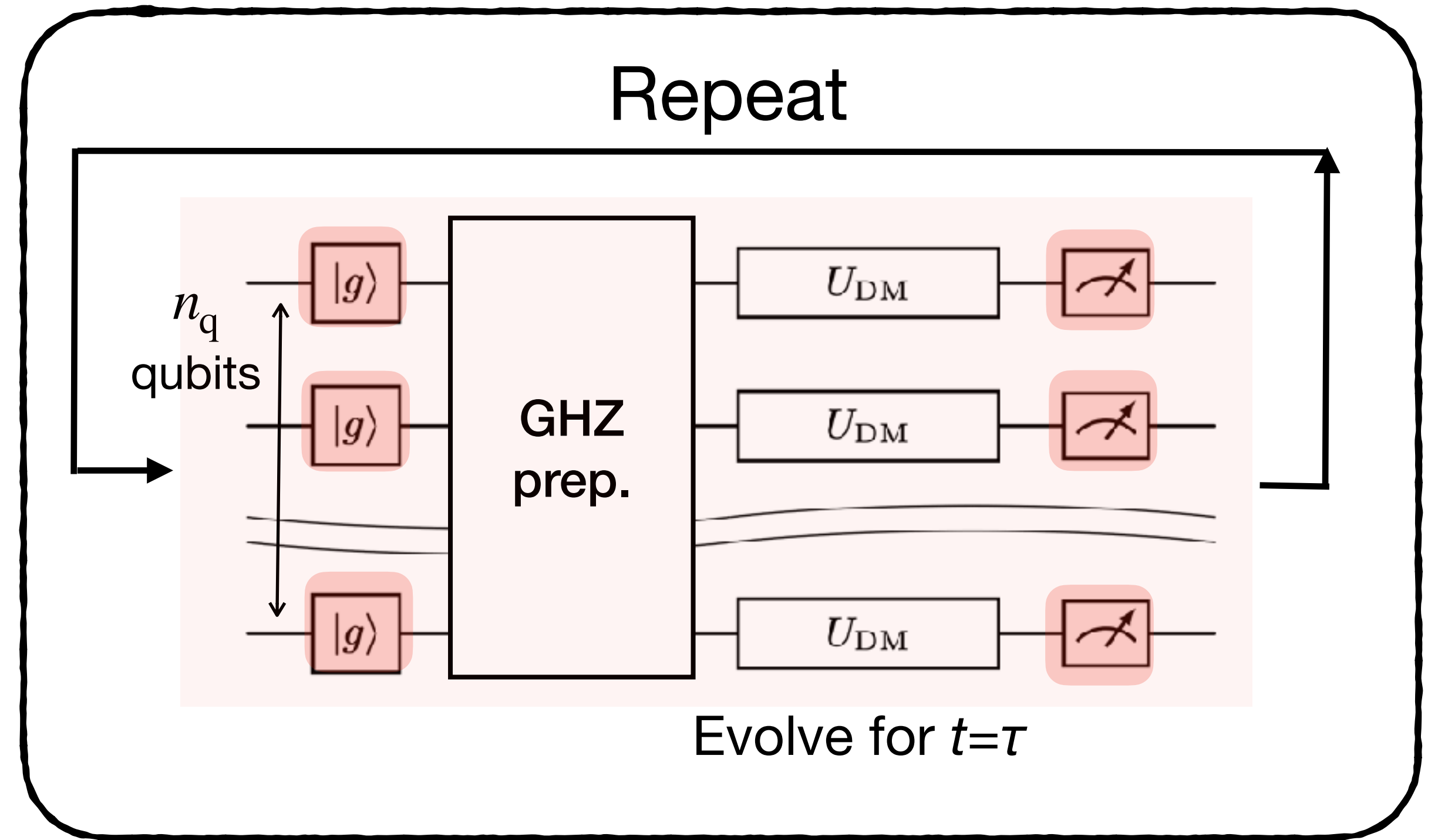
Quantum enhancement by quantum circuits

(2) arXiv: 2311.10413

Individual measurement



Using quantum circuit



The number that qubits are excited

$$N_{\text{signal}} \propto n_q$$

The number of DM signal*

$$N_{\text{signal}} \propto n_q^2$$

Accumulated phase with the effect of DM

For $m_X = \omega$, the resonant evolution is

$$\frac{d}{dt} \begin{pmatrix} \psi_g \\ \psi_e \end{pmatrix} = \begin{pmatrix} 0 & ie^{-i\alpha}\eta \\ ie^{i\alpha}\eta & 0 \end{pmatrix} \begin{pmatrix} \psi_g \\ \psi_e \end{pmatrix} \Rightarrow U_{\text{DM}} \simeq \begin{pmatrix} 1 & ie^{-i\alpha}\eta t \\ ie^{i\alpha}\eta t & 1 \end{pmatrix}$$

For $\alpha = 0$: eigenstate $U_{\text{DM}} |\pm\rangle = e^{\pm i\delta} |\pm\rangle$ with $\delta \equiv \eta\tau$,

For n_q qubit state

$$U_{\text{DM}}^{\alpha=0} |\pm\rangle^{\otimes n_q} = e^{\pm i n_q \delta} |\pm\rangle^{\otimes n_q}$$

Phase accumulated coherently !

; $|\pm\rangle \equiv (|g\rangle \pm |e\rangle)/\sqrt{2}$.

Quantum computation basics

• Hadamard: $\text{---} \boxed{H} \text{---} \Rightarrow H|g\rangle = |+\rangle, H|e\rangle = |-\rangle; |\pm\rangle \equiv (|g\rangle \pm |e\rangle)/\sqrt{2}.$

• X gate: $\text{---} \boxed{X} \text{---} \Rightarrow X = |g\rangle\langle e| + |e\rangle\langle g|$

• Z gate: $\text{---} \boxed{Z} \text{---} \Rightarrow Z = |g\rangle\langle g| - |e\rangle\langle e| = |+\rangle\langle -| + |-\rangle\langle +|$

• Controlled NOT:

$$\begin{array}{c} \bullet \\ | \\ \oplus \end{array} \Rightarrow |g\rangle\langle g| \otimes I + |e\rangle\langle e| \otimes X$$

• Controlled Z:

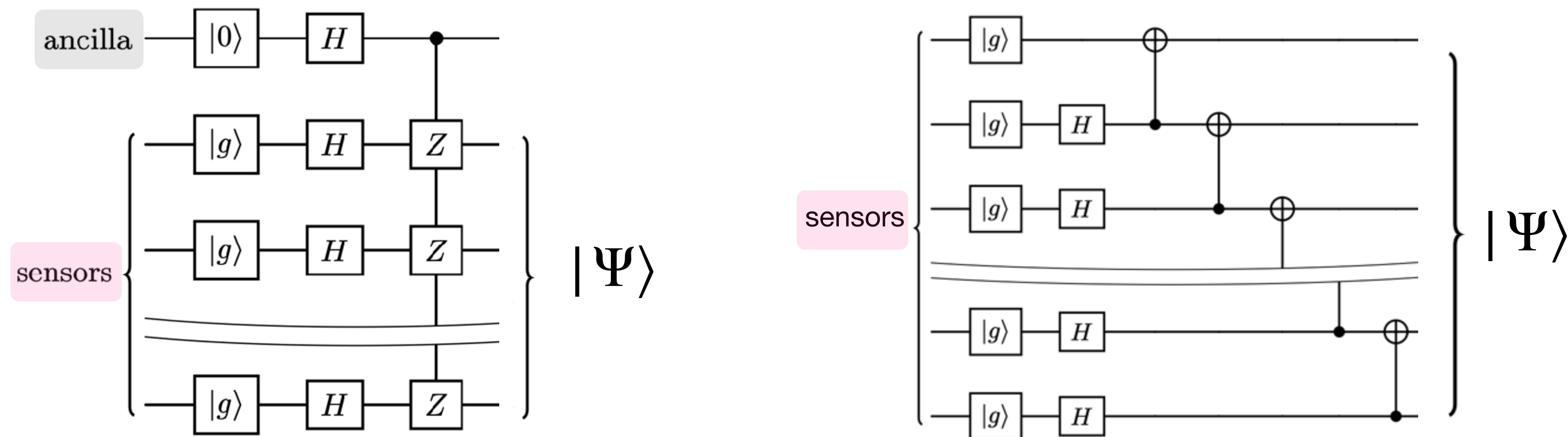
$$\begin{array}{c} \bullet \\ | \\ \boxed{Z} \end{array} \Rightarrow |g\rangle\langle g| \otimes I + |e\rangle\langle e| \otimes Z$$

GHZ state

GHZ state in Hadamard basis (highly entangled state)

$$|\Psi\rangle = |+\rangle^{\otimes n_q} + |-\rangle^{\otimes n_q} \Rightarrow U_{\text{DM}}^{\alpha=0} |\Psi\rangle = e^{in_q\delta} |+\rangle^{\otimes n_q} + e^{-in_q\delta} |-\rangle^{\otimes n_q}$$

Example circuit to prepare the GHZ state in Hadamard basis



DM signal = Observing odd excitations

$$\text{Prep. } |\Psi\rangle = |+\rangle^{\otimes n_q} + |-\rangle^{\otimes n_q}$$

$$U_{\text{DM}}^{\alpha=0} |\Psi\rangle = e^{in_q\delta} |+\rangle^{\otimes n_q} + e^{-in_q\delta} |-\rangle^{\otimes n_q}$$

$$= \cos(n_q\delta) (|+\rangle^{\otimes n_q} + |-\rangle^{\otimes n_q}) + i \sin(n_q\delta) (|+\rangle^{\otimes n_q} - |-\rangle^{\otimes n_q})$$

Contain only even $|e\rangle$ states

$$Z_1 Z_2 \dots Z_n = +1$$

Induced by effect of DM

Contain only odd $|e\rangle$ states

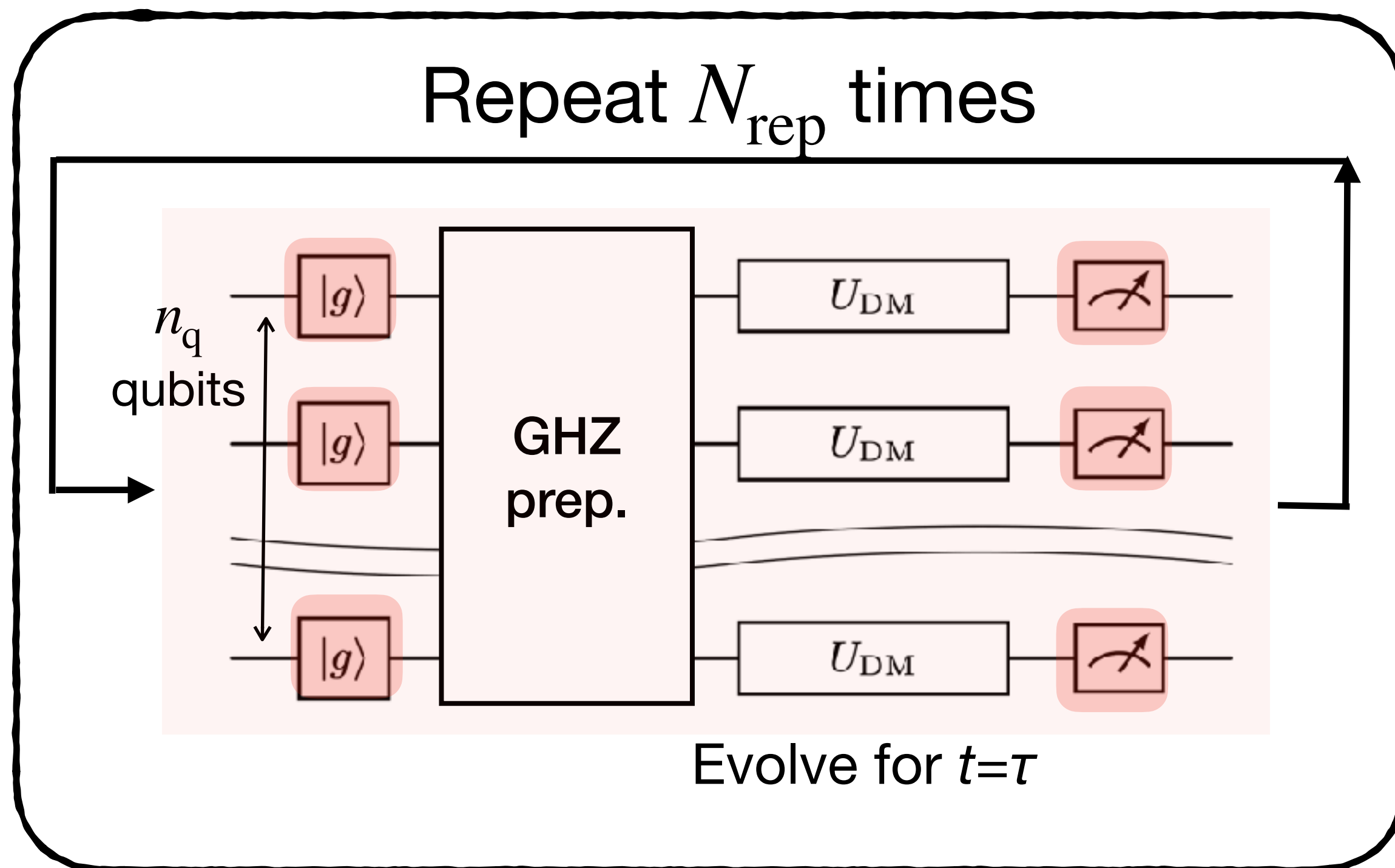
$$Z_1 Z_2 \dots Z_n = -1$$

DM signal = observing odd qubits excited; Prob. $p_{\text{odd}} = \sin^2(n_q\delta) \simeq (n_q\delta)^2$

For general α : $p_{\text{odd}} \simeq (n_q\delta)^2 \cos^2 \alpha$

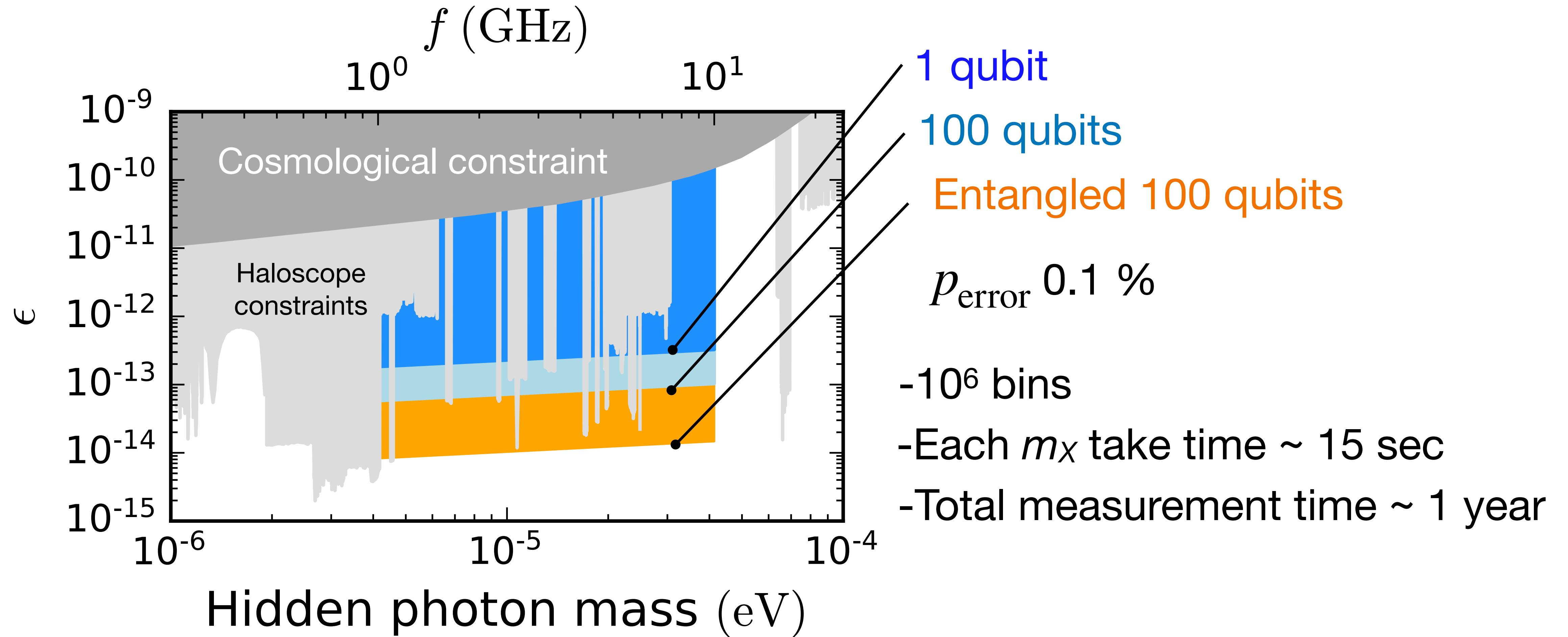
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Using quantum circuit



- Signal with N_{rep} repetition of experiment
- $$N_{\text{signal}} = \frac{1}{2} n_q^2 \delta^2 N_{\text{rep}} \propto n_q^2$$
- Noise with N_{rep} repetition of experiment
- $$\text{Bkg} = p_{\text{noise}} n_q N_{\text{rep}}$$
- Significance $\sigma = N_{\text{signal}} / \sqrt{\text{Bkg}}$ for claiming detection

Sensitivity plot (Hidden photon)



Summary

- Coherent wave-like dark matter can excite qubits, resulting in detectable signal
- Transmon has good sensitivity, reaching unexplored parameter regions of hidden photon dark matter
- Enhancement with the quantum circuit is possible
- Extension to axion detection

Sensitivity plot (Axions)

