

WIMP search at future muon collider

Atsuya Niki (U. Tokyo)

based on JHEP 02 (2024) 214 (arXiv: 2310.07162)
collaboration with H. Fukuda, T. Moroi and S-F. Wei

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Outline

- Minimal dark matter candidate as WIMP
- Muon collider as a future high energy collider
 - indirect search
 - direct search
- Result
- Summary

Weakly Interacting Massive Particle (WIMP)

- We consider new SU(2) multiplet (n -plet). Cirelli, Fornengo, Strumia (2005)
- Well-known **dark matter candidate** (WIMP miracle) and many models predict its existence, like MSSM. Jungman, Kamionkowski, Griest (1995)

Our target is where the WIMPs can explain the dark matter relic abundance through the freeze-out mechanism (**thermal mass target**).

We consider the femionic WIMP candidate:

- **Higgsino** ($n = 2, Y = 1/2$) : ~ 1 TeV
 - **Wino** ($n = 3, Y = 0$) : ~ 2.7 TeV
 - **Quintuplet minimal dark matter** ($n = 5, Y = 0$) : ~ 14 TeV
- Y : hypercharge

Hisano, Matsumoto, Nagai, Saito, Senami (2007)
Cirelli, Strumia, Tamburini (2007)
Mitridate, Redi, Smirnov, Strumia (2017)

Collider search for WIMP

WIMP is one of the main target at **high energy collider experiment**.

- **LHC**

Charged state of **Wino** and **Quintuplet** is long-lived, $c\tau \sim \mathcal{O}(\text{cm})$, and the decay product is soft due to small mass splitting. → **Disappearing Track**

Chen, Drees, Gunion (1996)

Ostdiek (2015)

ATLAS collaboration (2022)

Higgsino is not so long-lived due to large mass splitting. Even in this case, **displaced soft pion** can be used to discriminate the background.

Fukuda, Nagata, Oide, Otono, Shirai (2020)

- **ILC**

Lepton collider is a **clean environment**, where the kinematics can be fully reconstructed. This provides a great sensitivity, but the center of mass energy is not so large ($\sqrt{s} = 250\text{-}1000$ GeV).

ILC international development team (2021)

Muon collider

For WIMP search, **muon collider** has many **advantages**.

- high energy: $\sqrt{s} = \mathcal{O}(1-10 \text{ TeV})$
- large luminosity (circular collider): more than ab^{-1}
- low hadronic BG

But muon collider suffer from **Beam-Induced Background (BIB)**.

It may become a main background of disappearing track signal...

Capdevilla, Meloni, Simoniello, Zurita (2021)

Then, WIMP search **without** using track information is important to show the capability of muon collider.

□ mono-X search at $\mu^+\mu^-$ collider

Han+ (2021), Bottaro+ (2022)

□ muon polarization, cut optimization, etc.

Fukuda, Moroi, AN, Wei, in progress

$\mu^+ \mu^+$ collider

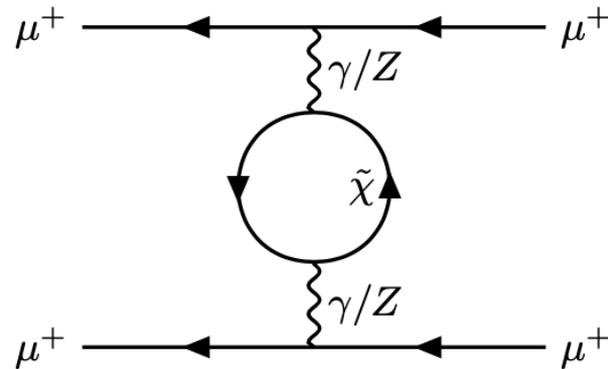
Recently, μ^+ cooling technique is developed, and **$\mu^+ \mu^+$ collider seems a realistic option as a future collider experiment.**

Hamada, Kitano, Matsuda, Takaura, Yoshida (2022)

In this work, we consider WIMP search by **two different ways**.

- 1. Indirect search:** angular distribution of $\mu^+ \mu^+$ elastic scattering
- 2. Direct search:** WIMP production and mono- μ channel

Indirect search



WIMP affects the $\mu^+ \mu^+$ elastic scattering through the gauge boson propagator:

$$\mathcal{L} = \mathcal{L}_{\text{SM}} + \frac{C_{WW}}{4} W_{\mu\nu}^a \Pi(-D^2/m^2) W^{a\mu\nu} + \frac{C_{BB}}{4} B_{\mu\nu} \Pi(-\partial^2/m^2) B^{\mu\nu} + \dots,$$

where

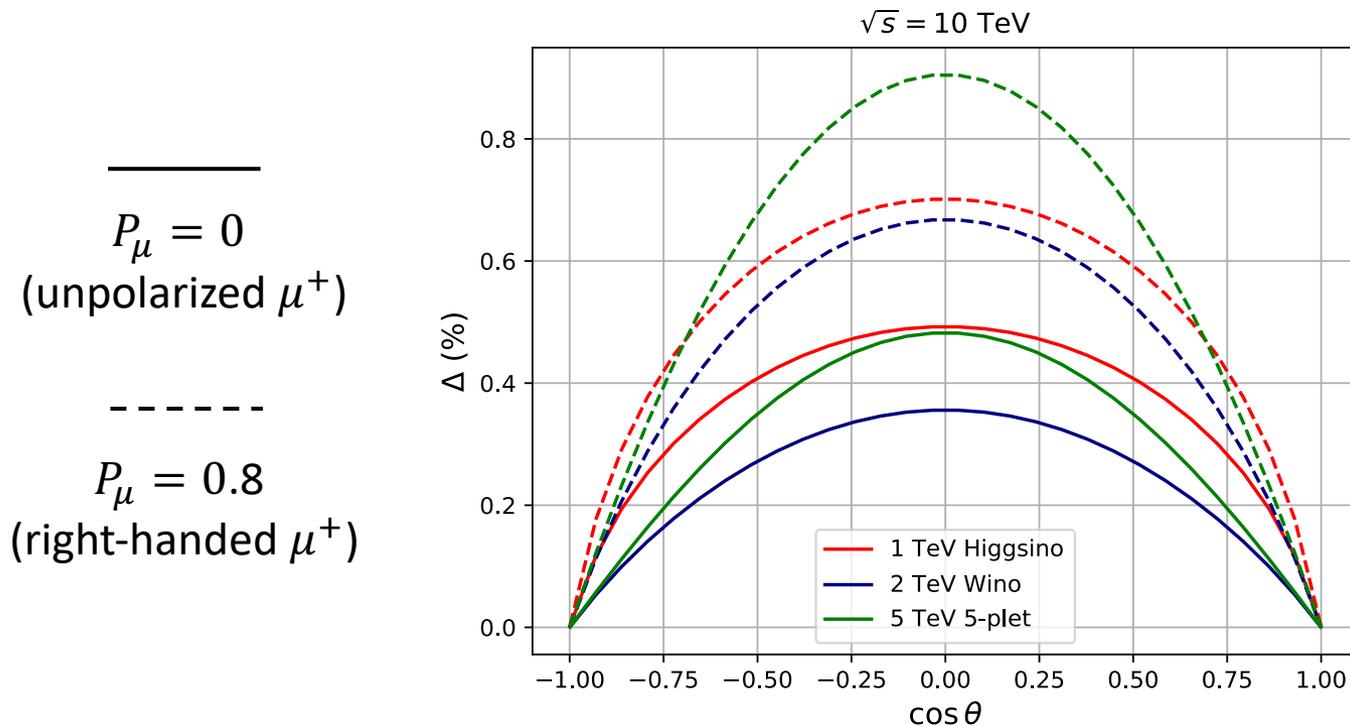
$$\Pi(x) = \frac{1}{16\pi^2} \int_0^1 dy y(1-y) \log \left(\frac{m^2 - xy(1-y)m^2}{\mu^2} \right),$$

From the t,u-dependence of the self-energy Π ,
the angular distribution of μ is distorted.

Indirect search

(WIMP contribution to cross section) / (SM contribution)

$$\Delta(\theta) \equiv \frac{d(\sigma_{\text{BSM}} - \sigma_{\text{SM}}) / d \cos \theta}{d\sigma_{\text{SM}} / d \cos \theta}$$



Indirect search

Statistical method : **shape analysis**

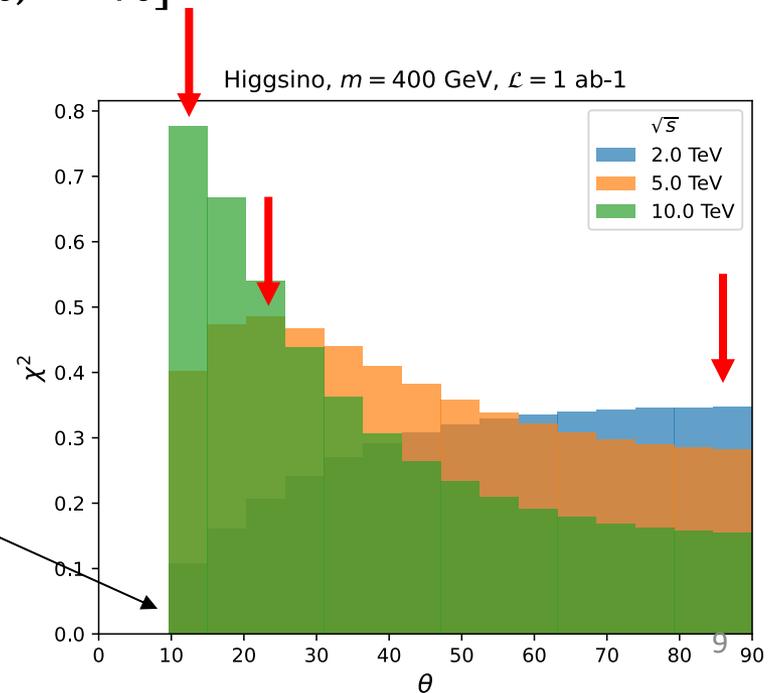
$$\chi^2 = \sum_{i \in \text{bin}} \frac{\left(N_i^{(SM+WIMP)} - N_i^{(SM)} \right)^2}{N_i^{(SM)} + (N_i^{sys})^2},$$

Bin: 15 intervals of the scattering angle, which satisfy $0 < \eta < 2.5$.

Systematic error: $N^{sys} = \epsilon N^{SM}$, $\epsilon \in [0\%, 0.3\%]$

The most contributing bin is where $t \sim m^2$, and this makes **the peak structure** of each bin contribution to χ^2 .

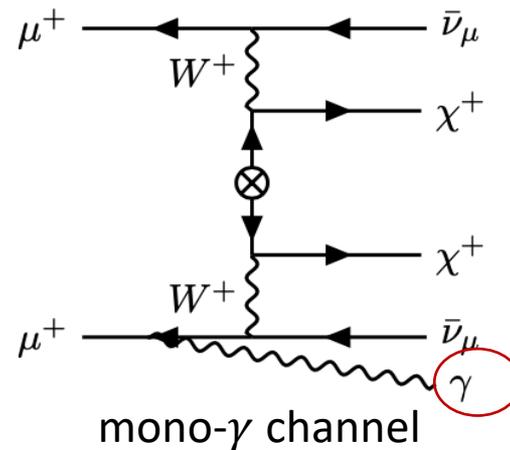
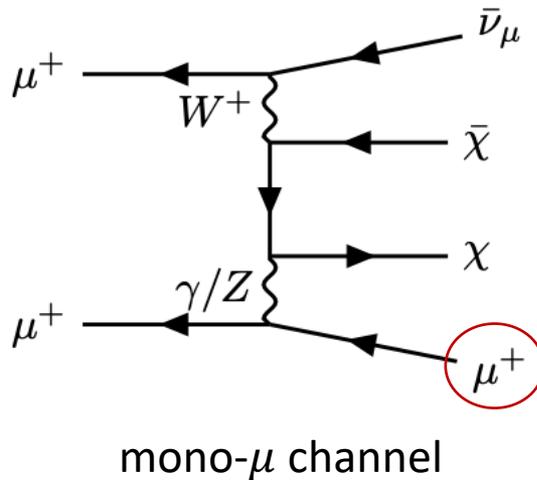
$\eta = 2.5$



Direct search

At $\mu^+\mu^-$ collider, Drell-Yan production process is dominant, not at $\mu^+\mu^+$ collider.

Then **VBF process** is the dominant production process at $\mu^+\mu^+$ collider.



Signal channel is mono- X ($X=\mu, \gamma, W, \dots$).

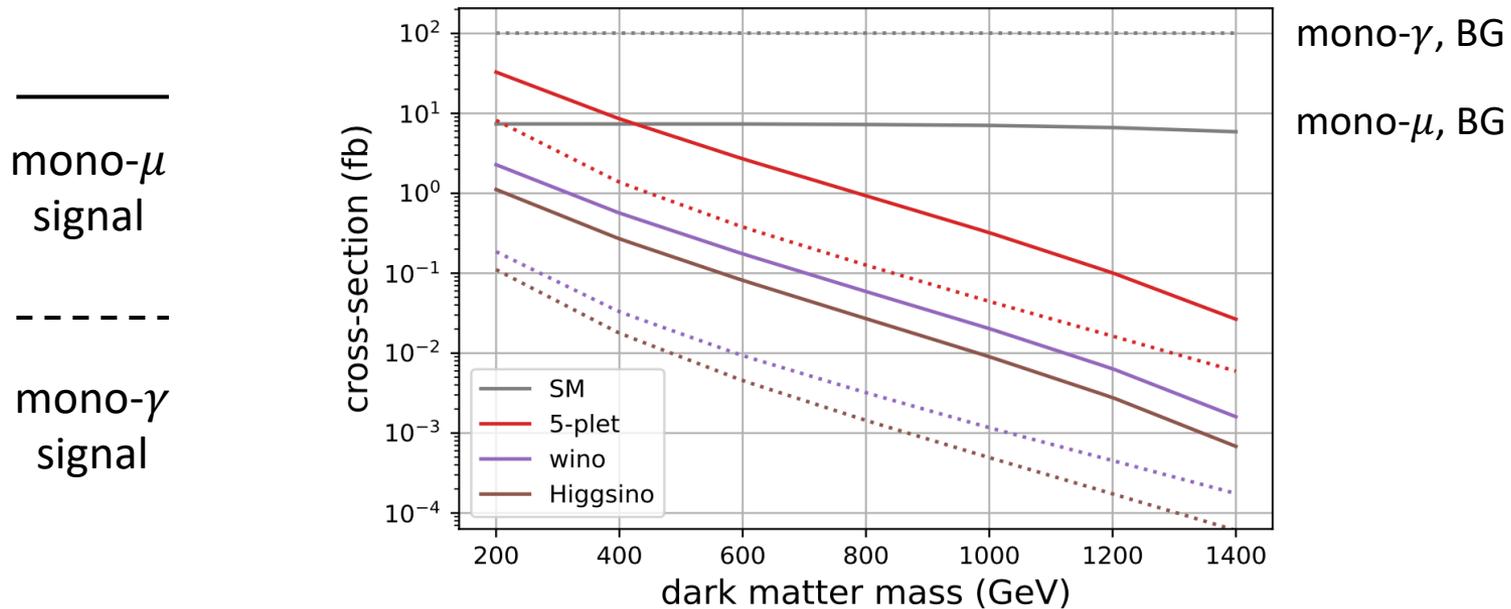
Direct search

In the case of $\mu^+\mu^-$ collider, **mono- μ channel** is most sensitive because background can be discriminated **by the kinematical cut**.

Han, Liu, Wang, Wang (2021)

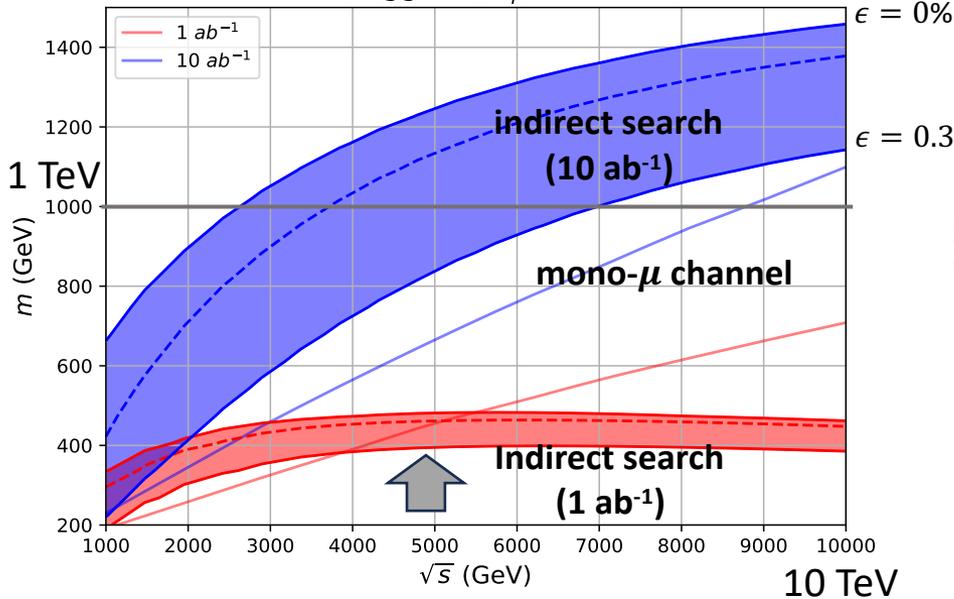
This is also true at $\mu^+\mu^+$ collider.

We show **only mono- μ channel** in our result.

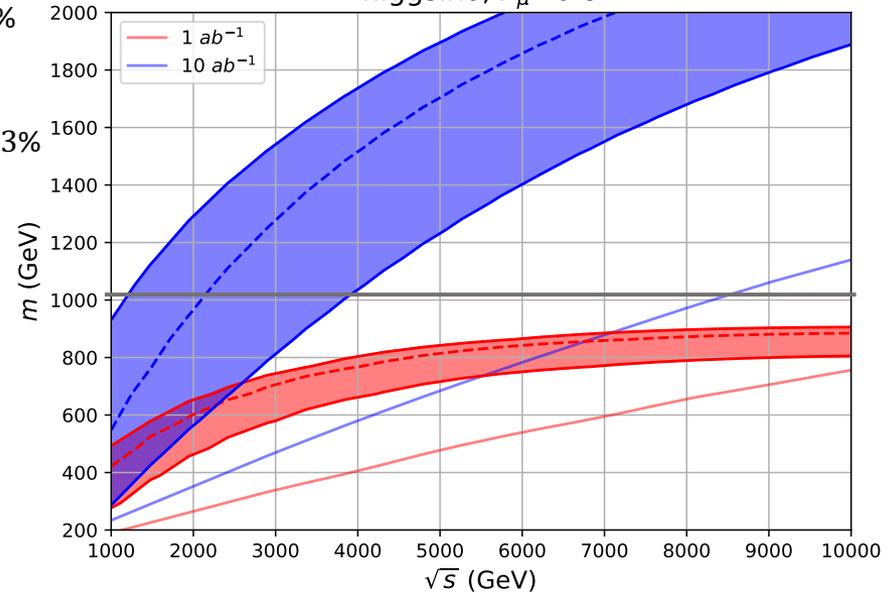


Result: Higgsino

unpolarized
higgsino, $P_\mu=0.0$



(right-handed) polarized
higgsino, $P_\mu=0.8$



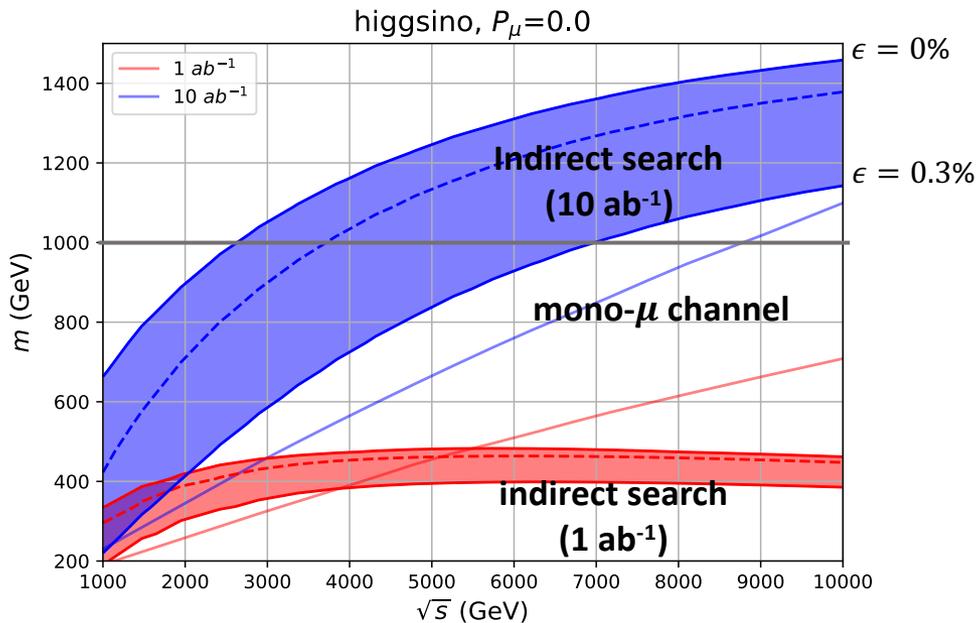
1 ab^{-1} : indirect search is more sensitive than direct search with $\sqrt{s} \lesssim 5$ TeV.

10 ab^{-1} : indirect search is more sensitive than direct search.

With 10 ab^{-1} luminosity, both searches can probe **the thermal target of higgsino.**

With **polarized muon**, sensitivity of the indirect search is much enhanced, due to the increase of the effective luminosity and SN ratio.

Result: Higgsino



\mathcal{L} : luminosity

Significance

Indirect search: $\chi^2 \propto \mathcal{L}/m^3$

Direct search: $\frac{N_{signal}}{\sqrt{N^{BG}}} \propto \sqrt{\mathcal{L}}/m^4$

fixed significance
e.g. 95% CL

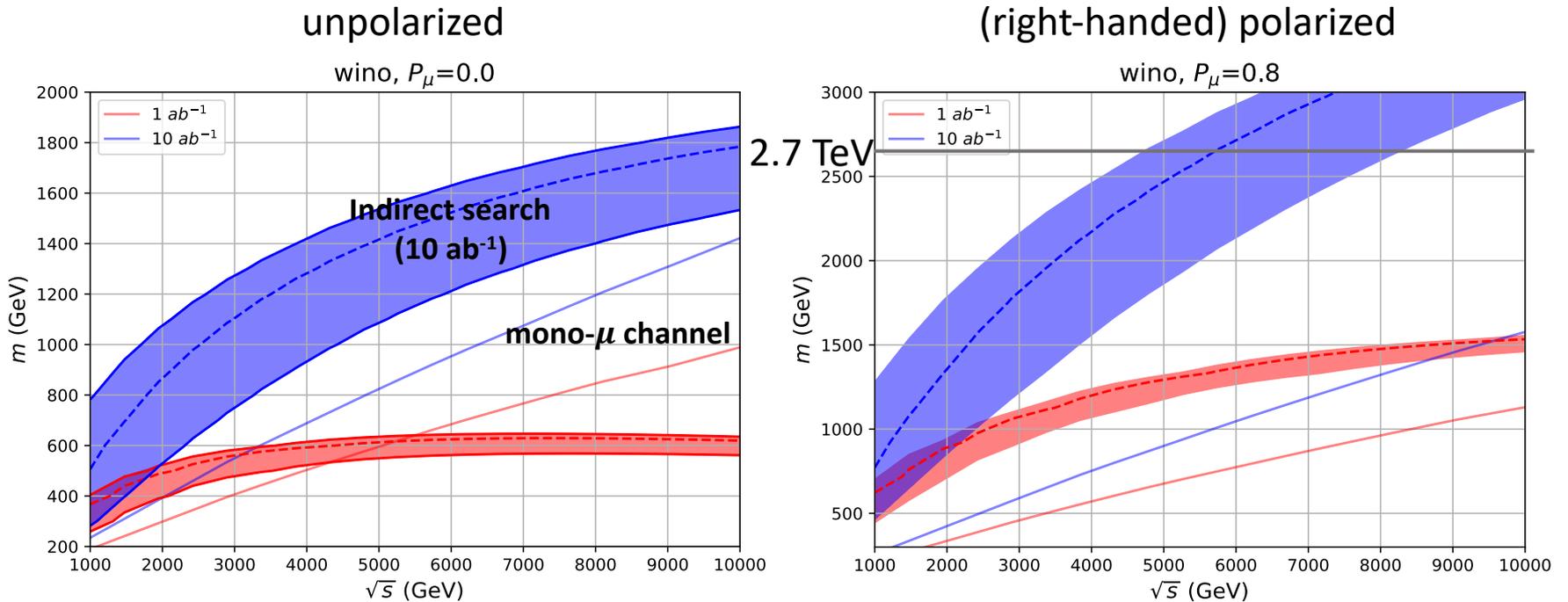


Indirect search: $m \propto \mathcal{L}^{1/3}$
Direct search: $m \propto \mathcal{L}^{1/8}$

Indirect search has **larger** sensitivity than direct search (mono- μ channel) with **large luminosity**.

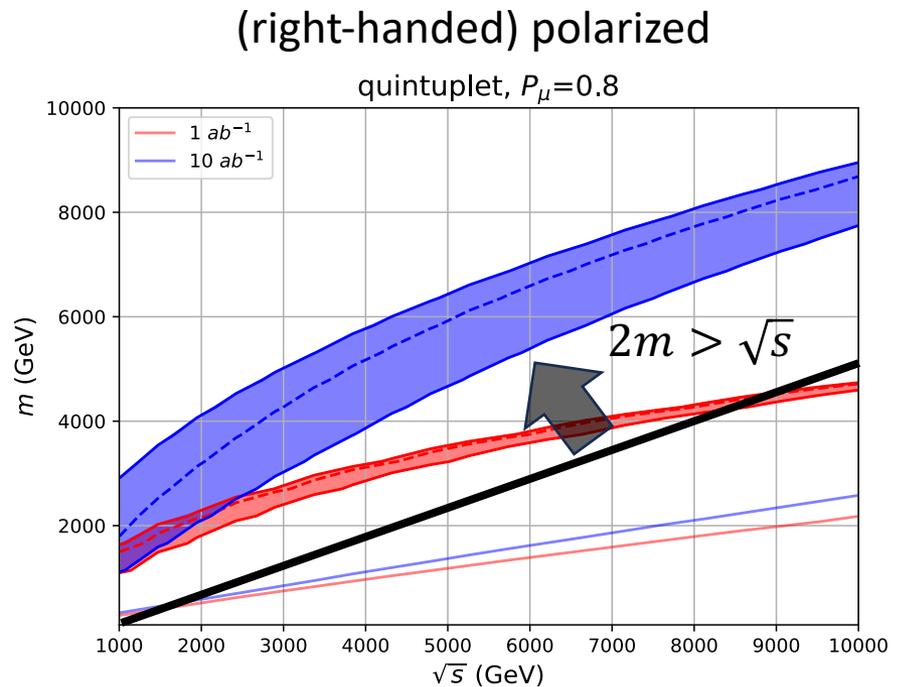
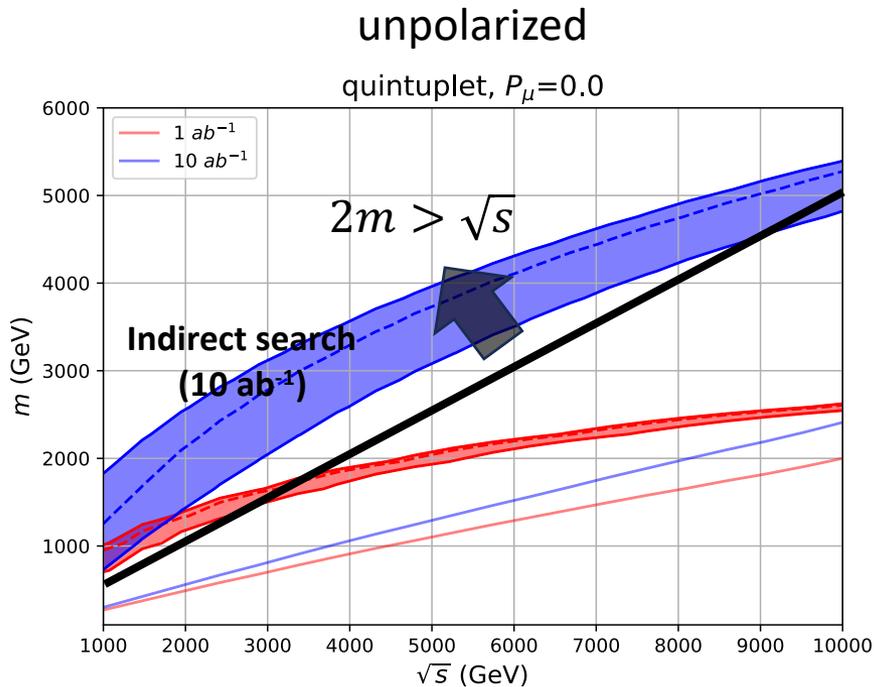
* When $\sqrt{s} \gg m^2$, χ^2 mass-dependence is different from what is discussed above due to the forward angular cutoff. See backup or our paper for more detail.

Result: Wino



With **polarized muon beam**, the indirect search is sensitive to **the thermal target** ($m = 2.7 \text{ TeV}$) with $\sqrt{s} \sim 5 \text{ TeV}$, 10 ab^{-1} collider.

Result: Quintuplet



The indirect search is sensitive to the parameter regions, $2m > \sqrt{s}$, where the direct search cannot search due to the kinematics.

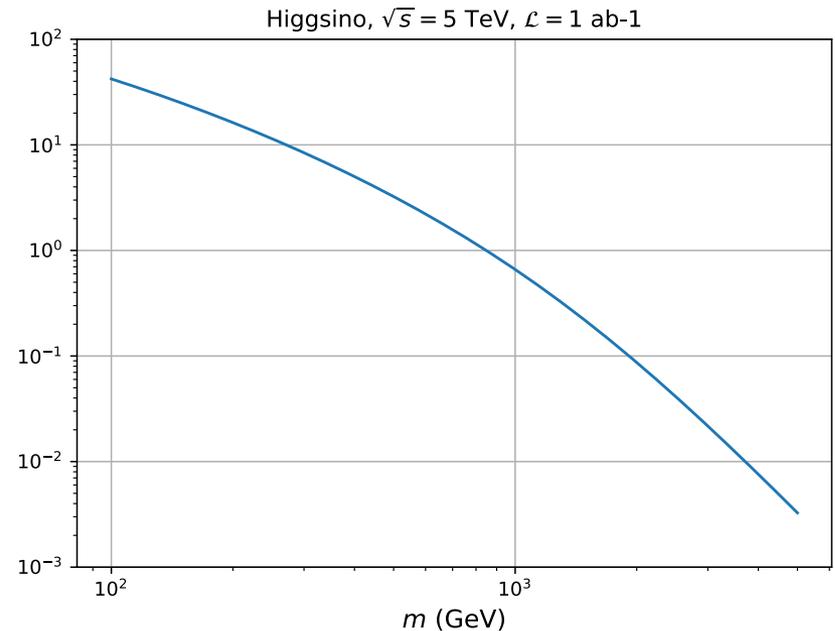
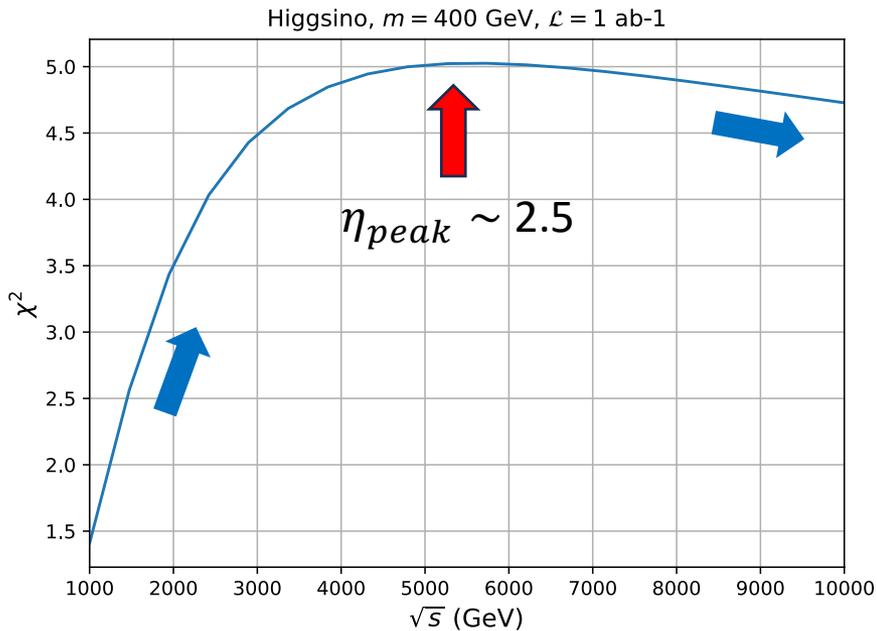
Summary

We estimate the sensitivity of the **indirect** and **direct** search for WIMP at $\mu^+\mu^+$ collider.

- Quantum correction from WIMP modifies the angular distribution of $\mu^+\mu^+$ elastic scattering (indirect search).
- For the direct search, mono- μ channel has larger sensitivity than mono- γ channel.
- Indirect search has **an advantage** over the direct search with sufficient luminosity due to the difference of mass dependence.
- Initial state muon polarization enhances the sensitivity of the indirect search.
- With 10 ab^{-1} and polarized beam, the thermal target for Higgsino (Wino) can be probed when $\sqrt{s} \sim 2 (6) \text{ TeV}$.

Backup

Indirect search



$d\chi^2/d\theta$ has a peak around $t \sim m^2$.

If this peak is inside the observed range, $\chi^2 \propto \sqrt{s} \frac{\mathcal{L}}{m^3}$.

If the peak is outside, χ^2 slightly decreases with large \sqrt{s} .

Direct production search

Kinematical cut

- $E_\mu > 0.23\sqrt{s}$
- $|\eta| < 2.5$
- $(p_{\mu,1}^{in} + p_{\mu,2}^{in} - p_\mu^{out})^2 > 4m^2$

BG process for mono- μ channel

