# **Returning CP-observables to the frames they belong**

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> Based on arXiv: 2308.00027









- Conventional LHC analysis involves comparing measured data with MC events simulated under NP hypothesis.
  - Reconstructed LHC events present a convoluted version of the true underlying physics.
  - Forward simulation chain can be highly resource intensive.

Invert simulation chain  $\rightarrow$  apply on measured data  $\rightarrow$  reconstruct parton-level

 $\rightarrow$  compare new physics hypotheses at the parton-level.

### • Bin-by-bin unfolding:

• Correct the information in each bin using correction factor  $\mathscr{C}_i$  computed from MC data.



- Matrix inversion:
  - Build response matrix  $R \rightarrow$  each cell  $\{i, j\}$  represents the fraction of events which have a true value in bin
    - i but get reconstructed in bin j.
- Iterative unfolding
  - Augment R<sub>ii</sub> by correction factors computed by comparing generated with true parton level data.

Bin-dependent unfolding.

# Unfolding

rton truth  $\mathscr{C}_i = N_{\text{truth},i}/N_{\text{reco},i}$ Unfolded distribution:  $x_{p,i} = x_{d,i} \times \mathscr{C}_i$ **Detector level** 

Dimensionality limitations.



# **ML-Unfolding**



### Possible with machine learning based generative models.

• Generative Adversarial Networks (GAN)

Variational Auto Encoders (VAE)

[Bellagente, Butter, Kasieczka, Plehn, Winterhalder (2020)] [Bellagente, Butter, Kasieczka, Plehn, Rousselot, Winterhalder, Ardizzone, Kothe (2020)] [Andreassen, Komiske, Metodiev, Nachman, Thaler (2020)] [Komiske, McCormack, Nachman (2021)]

Normalizing Flows (NF)



























$$D(x_p) \to 1, \quad D(x_G) \to 0$$







$$D(x_p) \to 1, \quad D(x_G) \to 0$$

[Bellagente, Butter, Kasieczka, Plehn, Winterhalder(2019)] [Butter, Plehn, Winterhalder(2019)]



 $L_{\rm D} = \left\langle -\log D(x) \right\rangle_{x \sim P_p} + \left\langle -\log(1 - D(x)) \right\rangle_{x \sim P_G}$ 







- Discriminator works to distinguish generated data  $\{x_G\}$  from truth data  $\{x_p\}$ .  $[D(x_P) \rightarrow 1, D(x_G) \rightarrow 0]$
- Generator works to fool the discriminator such that  $D(x_G) \rightarrow 1$ .

[Bellagente, Butter, Kasieczka, Plehn, Winterhalder(2019)] [Butter, Plehn, Winterhalder(2019)]



 $L_{\rm D} = \left\langle -\log D(x) \right\rangle_{x \sim P_p} + \left\langle -\log(1 - D(x)) \right\rangle_{x \sim P_G}$  $L_{\rm G} = \left\langle -log D(x) \right\rangle_{x \sim P_{\rm G}}$  $D(x_p) \rightarrow 1, \quad D(x_G) \rightarrow 0$ 





- Targeted loss terms required for sharp kinematic features.
- Fails if training and test data to not statistically similar.



[Bellagente, Butter, Kasieczka, Plehn, Winterhalder(2019)]





Figure taken from Bellagente, Butter, Kasieczka, Plehn, Winterhalder(2019)





[Image adopted from Bellagente, Butter, Kasieczka, Plehn, Winterhalder (2019)]









 $L_{\rm D}^{FC} = \langle -logD(x, y) \rangle_{r}$ 

 $L_{\rm G} = \left\langle -\log D(x, y) \right\rangle_{x \sim P_G, y \sim P_d}$ 

- Event-by-event matching  $\rightarrow$  exploit the pairing information between parton and detector level.
- Trained network can be applied to statistically different regions of phase space.

Test data

 $2\ell + 2j$  exclusive @ detector level

 $30 \text{ GeV} < p_{T,j_1} < 60 \text{ GeV}, 30 \text{ GeV} < p_{T,j_2} < 50 \text{ GeV},$ (  $\sim 14\%$  of events)

### Fails with harsher cuts !

$$x \sim P_p, y \sim P_d + \langle -log(1 - D(x, y)) \rangle_{x \sim P_G, y \sim P_d}$$

 $\times 10^{-2}$ Truth 4.0Eq.(7)..... Eq.(8)-----3.0 $\frac{1}{\sigma} \frac{\mathrm{d}\sigma}{\mathrm{d}p_{T,j_1}}$  $1.0^{-1}$ 0.0 25100 125 50150750 175 $p_{T,j_1}$  [GeV]

> [Image adopted from Bellagente, Butter, Kasieczka, Plehn, Winterhalder (2019)]





→ Regress detector response function starting from a generator-level jet

[Touranakou, Chernyavskaya, Duarte, Gunopulos, Kansal, Orzari, Pierini, Tomei, Vlimant (2022)] [Otten, Caron, Swart, Beekveld, Hendriks, Leeuwen, Podareanu, Austri, Verheyen (2019)]



**Reconstruction Loss** 



### $\eta KL(p(z \mid x) \mid | q(Z))$ KL divergence term

[lanazi, Sato, Ambrozewicz, Blin, Melnitchouk, Battaglieri, Liu, Li (2021)]



→ Regress detector response function starting from a generator-level jet

[Touranakou, Chernyavskaya, Duarte, Gunopulos, Kansal, Orzari, Pierini, Tomei, Vlimant (2022)]



![](_page_14_Picture_4.jpeg)

![](_page_14_Picture_6.jpeg)

**Process:**  $pp \rightarrow WW \rightarrow (W \rightarrow jj)(W \rightarrow jj)$ 

![](_page_15_Figure_2.jpeg)

Good agreement between reco and predicted distributions, but jet substructure quantities not well reproduced.

Figures taken from Touranakou, Chernyavskaya, Duarte, Gunopulos, Kansal, Orzari, Pierini, Tomei, Vlimant (2022)

![](_page_15_Picture_7.jpeg)

![](_page_16_Figure_2.jpeg)

- essential features.
- The decoder maps z to the parton level p' = D(z) = E(l(d)).

![](_page_16_Figure_6.jpeg)

![](_page_16_Picture_7.jpeg)

• The Encoder maps the input detector data d to a more tractable latent space z = E(d) while preserving the

[Otten, Caron, Swart, Beekveld, Hendriks, Leeuwen, Podareanu, Austri, Verheyen (2019)]

![](_page_16_Picture_11.jpeg)

**Exact likelihood estimation** 

Invertibility :

- NF:capable of bi-directional mapping w/o information loss.
- > VAEs: not strictly invertible due to stochasticity of the latent space.
- FCGANs: focus on generation (sharper data), and invertibility is not strictly defined.

Flexibility:

- Image NF can model intricate distributions without making strict assumptions.
- the distributions.

•GANs focus on generating data that matches the target distribution  $\rightarrow$  no explicit latent mapping and less statistical robustness.

 $\odot$  VAEs assume a Gaussian latent space  $\rightarrow$  may not always capture the complexity of

- Building blocks  $\rightarrow$  Invertible coupling layers. [Dinh, Krueger, Bengio (2016), Dinh, Sohl-Dickstein, Bengio (2016)]

![](_page_18_Figure_4.jpeg)

# **Normalizing flows**

# • Series of bijective layers that transform complex (Y) to simple probability distributions (Z).

### • Learns both directions of the mapping in parallel $\rightarrow$ bijectivity encoded in the same network.

![](_page_18_Figure_9.jpeg)

Inversion

[Image adapted from Nguyen, Ardizzone, Kothe (2019) and talk by A. Butter at Pheno-2022]

![](_page_19_Figure_0.jpeg)

• In the coupling layers, the coupling functions  $s_2$ and  $t_2$  take  $x_2$  as input, and scale/translate  $x_1$ .

Forward pass:  

$$y_1 = x_1 \odot e^{S_2(x_2)} + t_2(x_2)$$
  
 $y_2 = x_2 \odot e^{S_1(y_1)} + t_1(y_1)$ 

• Fully invertible coupling layer  $\rightarrow [x_1, x_2]$ can be reconstructed given  $[y_1, y_2]$ 

> Inverse transformations:  $x_1 = (y_1 - t_2(x_2)) \odot e^{-s_2(x_2)}$  $x_2 = (y_2 - t_1(y_1)) \odot e^{-s_2(y_1)}$

![](_page_20_Figure_0.jpeg)

# For a coupling block transformation $f(x) \sim y$ tractable Jacobian $J_f(x)$ :

→ rule of change of variables  $p_Y(x_d) = p_Z(x_p) \times |det(J_f(x_p))|^{-1}$ 

$$: \frac{\partial f(x)}{\partial x} = \begin{bmatrix} e^{S_2(x_2)} & \text{finite} \\ 0 & e^{S_1(y_1)} \end{bmatrix}$$

 $\rightarrow$  ensures bijective transformations and exact likelihood estimation

### **Normalizing flows**

![](_page_21_Figure_1.jpeg)

- Coupling layers stacked together  $\rightarrow$  Invertible Neural Network (INN).
- inverse mapping ambiguous  $\rightarrow$  Not an issue with INNs.

[Image adapted from Nguyen, Ardizzone, Kothe (2019) and talk by A. Butter at Pheno-2022]

• Typically, DNNs suffer an inherent information loss in the forward direction, making the

![](_page_21_Picture_8.jpeg)

![](_page_22_Figure_0.jpeg)

![](_page_23_Figure_0.jpeg)

![](_page_24_Figure_0.jpeg)

![](_page_25_Figure_0.jpeg)

Loss: 
$$L_{MSE}(x_p) + L_{MSE}(x_d) + L_{MMD}$$

### **Naive INN unfolding**

[Bellagente, Butter, Kasieczka, Plehn, Rousselot, Winterhalder, Ardizzone, Kothe (2020)]

**Process:**  $pp \rightarrow ZW \rightarrow$ 

![](_page_26_Picture_3.jpeg)

![](_page_26_Picture_4.jpeg)

$$\rightarrow (Z \rightarrow \ell^+ \ell^-)(W \rightarrow jj)$$

```
Forward simulation: \bar{g} \rightarrow
                                          (X_{\text{detector}})
           \leftarrow Unfolding: g
```

![](_page_26_Picture_7.jpeg)

![](_page_27_Figure_0.jpeg)

- Differences between generated and parton truth deviate in the soft  $p_{T,j_2}$  region and tails.
- Typically inefficient in the inversion of features not included in event parametrization.
- Dimensionality limitations.

**Process:**  $pp \to ZW \to (Z \to \ell^+ \ell^-)(W \to jj)$ 

[Figure taken from Bellagente, Butter, Kasieczka, Plehn, Rousselot, Winterhalder, Ardizzone, Kothe (2020)]

![](_page_27_Figure_6.jpeg)

![](_page_27_Picture_7.jpeg)

![](_page_28_Figure_0.jpeg)

- Allows mapping between unequal degrees of freedom at the parton and detector level.
- MMD terms included for each observable and gaussian input  $\rightarrow$  improves unfolding in the low and high  $p_T$  regions.

**Process:**  $pp \to ZW \to (Z \to \ell^+ \ell^-)(W \to jj)$ 

![](_page_28_Figure_5.jpeg)

### **Noise-extended INN: Limitations and Challenges**

- Inclusive detector level information requires using large number of random variables.
- Calibration of weights associated to different loss terms.
- Combination of several loss terms pose training challenges.

![](_page_29_Figure_4.jpeg)

### **Conditional INN**

![](_page_30_Figure_1.jpeg)

**QCD** jet radiation Particle decays

• Generate probability distributions at the parton-level, given detector-level events  $x_{detector}$ 

> [Image adapted from Nguyen, Ardizzone, Kothe (2019) and talk by A. Butter at Pheno-2022]

### Target phase space for unfolding can be chosen flexibly to include:

### **Unfolding semileptonic** *tth* **events**

➡ Parton-level:  $1\ell + 2b + 2\gamma + \nu + 2j$ ➡ Detector-level:  $1\ell + 2b + 2\gamma + MET + \leq 6$  jets inclusive

Challenges:

\* Can the unfolding model correctly reconstruct the two hard jets at the parton level from a variable number of jets at the detector level?

+ How well can the dedicated BSM observables be reconstructed?

How model-dependent is the training?

 $pp \rightarrow t\bar{t}h \rightarrow (t \rightarrow \ell\nu b)(\bar{t} \rightarrow jj\bar{b})(h \rightarrow \gamma\gamma)$ Acceptance cuts  $|\eta_b| < 4, |\eta_i| < 5, |\eta_\ell| < 4, |\eta_\gamma| < 4$  $p_{T,b} > 25 \text{ GeV}, \quad p_{T,i} > 25 \text{ GeV}, \quad p_{T,\ell} > 15 \text{ GeV}, \quad p_{T,\gamma} > 15 \text{ GeV}$ 

![](_page_31_Picture_7.jpeg)

![](_page_31_Picture_8.jpeg)

- state particles  $\rightarrow$  may include redundant d.o.f.
- Reconstruction of sharp kinematic features like mass peaks can be challenging: Can be improved by adding targeted maximum mean discrepancy loss: [Butter, Plehn, Winterhalder (2019)] **Marget Affects only the target distributions** [Bellagente, Butter, Kasieczka, Plehn, Rousselot, **Marge model dependence** Winterhalder, Ardizzone, Kothe (2020)] Complications in training and performance limitations.

### <u>Alternative approach:</u>

 $\rightarrow$  directly learn invariant mass features and important observable with appropriate phasespace parametrization.

 $\rightarrow$  may provide direct access to the most important BSM observables.

• Event information at the parton level can be parametrised through the 4-momentum of the final

### **Conditional INN**

- We use the Bayesian version of cINN [Butter, Heimel, Hummerich, Krebs, Plehn, Rousselot, Vent (2021)]
  - Stable network predictions
  - Allows the estimation of training-related uncertainties.
- Degrees of freedom: Parton-level:  $(t \rightarrow \ell \nu b)(\bar{t} \rightarrow j\bar{j}\bar{b})h$ 22 d.o.f.

A natural parametrization involving top mass:  $\left\{ m_t, p_{T,t}, \eta_t, \phi_t, m_W, \eta_W^t, \phi_W^t, \eta_{\ell,u}^W, \phi_{\ell,u}^W \right\}$ 

 Alternatively, redefine the parton level parametrization including the important CP observables

![](_page_33_Figure_7.jpeg)

### 46 d.o.f. Detector-level: $1\ell + 2b + 2\gamma + MET + \leq 6$ jets inclusive

 $\vec{p}_{t\bar{t}}, m_{t_{\ell}}, |\vec{p}_{t_{\ell}}^{\mathsf{CS}}|, \theta_{t_{\ell}}^{\mathsf{CS}}, \phi_{t_{\ell}}^{\mathsf{CS}}, m_{t_{h}},$   $\operatorname{sign}(\Delta \phi_{\ell \nu}^{t\bar{t}}) m_{W_{\ell}} |\vec{p}_{\ell}^{t\bar{t}}|, \theta_{\ell}^{t\bar{t}}, \phi_{\ell}^{t\bar{t}}, |\vec{p}_{\nu}^{t\bar{t}}|$ sign( $\Delta \phi_{du}^{t\bar{t}}$ )  $m_{W_h}$ ,  $|\vec{p}_d^{t\bar{t}}|$ ,  $\theta_d^{t\bar{t}}$ ,  $\Delta \phi_{\ell d}^{t\bar{t}}$ ,  $|\vec{p}_u^{t\bar{t}}|$ 

![](_page_33_Picture_12.jpeg)

![](_page_33_Picture_13.jpeg)

### **CP** measurement in Higgs-top interactions

- New sources of CPV interactions can explain the matter-antimatter asymmetry in the universe.
- One such scenario: CP violation in the Higgs sector.

### • CPV in *hVV* interactions is extensively tested at the LHC.

[See for instance: G. Aad et al. (1506.05669), G. Aad et al. (1602.04516), A. M. Sirunyan et al. (1707.00541), A. M. Sirunyan et al. (1903.06973), A. M. Sirunyan et al. (1901.00174), G. Aad et al. (2002.05315), Bernreuther, Gonzalez, Wiebusch (2010), Englert, Goncalves, Mawatari, Plehn (2012), Djouadi, Godbole, Mellado, Mohan (2013), Anderson, Bolognesi, Caola, Gao et al. (2013)]

### • CPV in *hff* couplings manifest at tree-level: $\rightarrow$ desirable choice: $ht\bar{t}$

![](_page_34_Figure_8.jpeg)

# **Direct probes at the LHC**

$$\mathscr{L} = -\frac{m_t}{v}\kappa_t$$

- $pp \rightarrow h$  (+ jets): indirect constraints.
- $pp \rightarrow t\bar{t}h$  stands out as the viable direct probe:
  - ◆ Small rate at the LHC and complex topology.
    ◆ Silver Lining: Observation at 5.2σ by ATLAS [2004.04545] and 6.6σ by CMS [2003.10866]
- Current limits:  $|\alpha| < 43^{\circ}$  (ATLAS) and  $|\alpha| < 55^{\circ}$  (CMS), at 95 % CL.

### Improved statistics @ HL-LHC paves the pathway for precision studies.

 $h\bar{t}(\cos\alpha + i\gamma_5\sin\alpha)t$ 

SM:  $(\kappa_t, \alpha) = (1, 0)$ 

[Duca, Kilgore, Oleari, Schmidt, Zeppenfeld (2001), Klamke, Zeppenfeld (2007), Grojean et al. (2013), Dolan, Harris, Jankowiak, Spannowsky (2014)]

![](_page_35_Figure_10.jpeg)

### Importance matrix at the **non-linear level**

![](_page_36_Figure_2.jpeg)

### Sensitive to non-linear new physics effects.

![](_page_36_Picture_4.jpeg)

-even observables nation					0.58
				0.4	0.69
			0.4	0.48	0.69
		0.61	0.69	0.69	0.65
	0.43	0.69	0.53	0.53	0.62
0.51	0.65	0.62	0.59	0.59	0.65
0.59	0.65	0.69	0.46	0.46	0.73
0.62	0.49	0.73	0.45	0.45	0.72
0.52	0.44	0.61	0.4	0.4	0.58
$b_4$	$\Delta \Phi_{t\bar{t}}$	$\Delta \eta_{t\bar{t}}$	$m_{th}$	$m_{\bar{t}h}$	$\theta^{*}$

[RKB, Goncalves, Kling (2021)]

### **CP-odd observables**

- Short lifetime for  $t (10^{-25} s) \rightarrow$  Spin correlations can be traced back from their decay products.
- CP-odd observables constructed from antisymmetric tensor products  $\epsilon(p_t, p_{\bar{t}}, p_i, p_j) \sim \epsilon_{\mu\nu\rho\sigma} p_t^{\mu} p_{\bar{t}}^{\nu} p_i^{\rho} p_j^{\sigma}$ :

$$\Delta \phi_{ij}^{t\bar{t}} = \mathbf{sgn} \left[ \vec{p}_t \cdot (\vec{p}_i \times \vec{p}_j) \right] at$$

![](_page_37_Figure_4.jpeg)

[RKB, Goncalves, Kling (2021)]

![](_page_37_Figure_6.jpeg)

 $\leftarrow \text{ Spin correlations scale with the spin analysing power } \beta_i.$ [Mileo, Kiers, Szynkman, Crane, Gegner (2016); Goncalves, Kong, Kim (2018)]; RKB, Goncalves, Kling (2021)]

$\frac{1}{-1} (1 \pm \beta P \cos \xi)$	Fisher Info = $\mathbb{E}$	$\int \partial \log p(\mathbf{x}   \kappa_{\mathrm{t}}, \alpha)$	$\partial \log p(\mathbf{x} \mid \kappa_t, \sigma)$
$2^{(1+p_i r_t \cos \varsigma_i)}$		$d\alpha$	$d\alpha$

### • Kinematic reconstruction efficiency is limited at the detector level

Use Machine learning techniques to maximize the extraction of NP information from CP observables.

![](_page_37_Picture_11.jpeg)

![](_page_37_Picture_12.jpeg)

### Back to results from unfolding with cINN...

### Challenges:

★ Can the unfolding model correctly reconstruct the two hard jets at the parton level from a variable number of jets at the detector level?

+ How well can the dedicated BSM observables be reconstructed?

How model-dependent is the training?

# **Jet combinatorics**

![](_page_39_Figure_2.jpeg)

\* Unfolded distributions in good agreement with parton level truth despite added combinatorial ambiguity at the detector level.

Parton level truth and unfolded top invariant masses  $m_{t_{e}}$  and  $m_{t_{h}}$ 

![](_page_39_Figure_5.jpeg)

### Back to results from unfolding with cINN...

### Challenges:

\* Can the unfolding model correctly reconstruct the two hard jets at the parton level from a variable number of jets at the detector level?

+ How well can the dedicated BSM observables be reconstructed?

How model-dependent is the training?

# **Reconstruction of dedicated observables**

![](_page_41_Figure_2.jpeg)

★ Unfolded distributions in close agreement with truth: ✓ Close agreement even for observables not included in event parametrization. ✓ Full phase space reconstruction.

\* Potential differences from the truth are covered by the uncertainty estimates of the Bayesian network.

### Parton level truth and unfolded SM for $\theta_{CS}$ , $\Delta \phi_{t_{e}t_{h}}$ and $b_{4}$ .

### Back to results from unfolding with cINN...

### Challenges:

\* Can the unfolding model correctly reconstruct the two hard jets at the parton level from a variable number of jets at the detector level?

+ How well can the dedicated BSM observables be reconstructed?

How model-dependent is the training?

# Model dependence

![](_page_43_Figure_2.jpeg)

 $b_4$  distributions.

Unfolding SM events using networks trained on events with different amounts of CP-violation.

 $\star$  Networks trained on  $\alpha = \pi/4$  and  $-\pi/4$  show only a slight bias towards broader  $\theta_{CS}$  and flatter

 $\star \sim 10 - 20\%$  bias  $\rightarrow$  much smaller than the changes at parton truth from varying  $\alpha$ .

# Model dependence

![](_page_44_Figure_3.jpeg)

 $\bigstar$ Again, the effect of bias is much smaller than the effect of  $\alpha$  on the data.

Unfolding events with CP-violation using a network trained on SM events.

![](_page_44_Figure_6.jpeg)

Unfold  $\alpha = + \pi/4, - \pi/4$  and SM dataset

![](_page_44_Figure_8.jpeg)

![](_page_44_Figure_9.jpeg)

- Generative unfolding makes it possible to invert high-dimensional distributions and full phasespace reconstruction.
- The trained cINN behaves as an efficient kinematic reconstruction algorithm capable of tackling complex reconstruction challenges.
- The trained unfolding network was able to
  - extract various CP observables at the parton level with appropriate phase space parametrization.
  - resolve jet combinatorial ambiguity.
  - absolve any large model-dependence.
- While this study is clearly not the last word on this analysis technique, it presents a promising outlook for an experimental study, with a proper treatment of statistical limitations, continuum backgrounds, calibration, and iterative improvements of the unfolding network.

# Outlook

![](_page_45_Figure_9.jpeg)

![](_page_46_Picture_1.jpeg)