

# Characteristic Mass of Dark Matter Haloes for the Cusp-Core Transition

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## Abstract

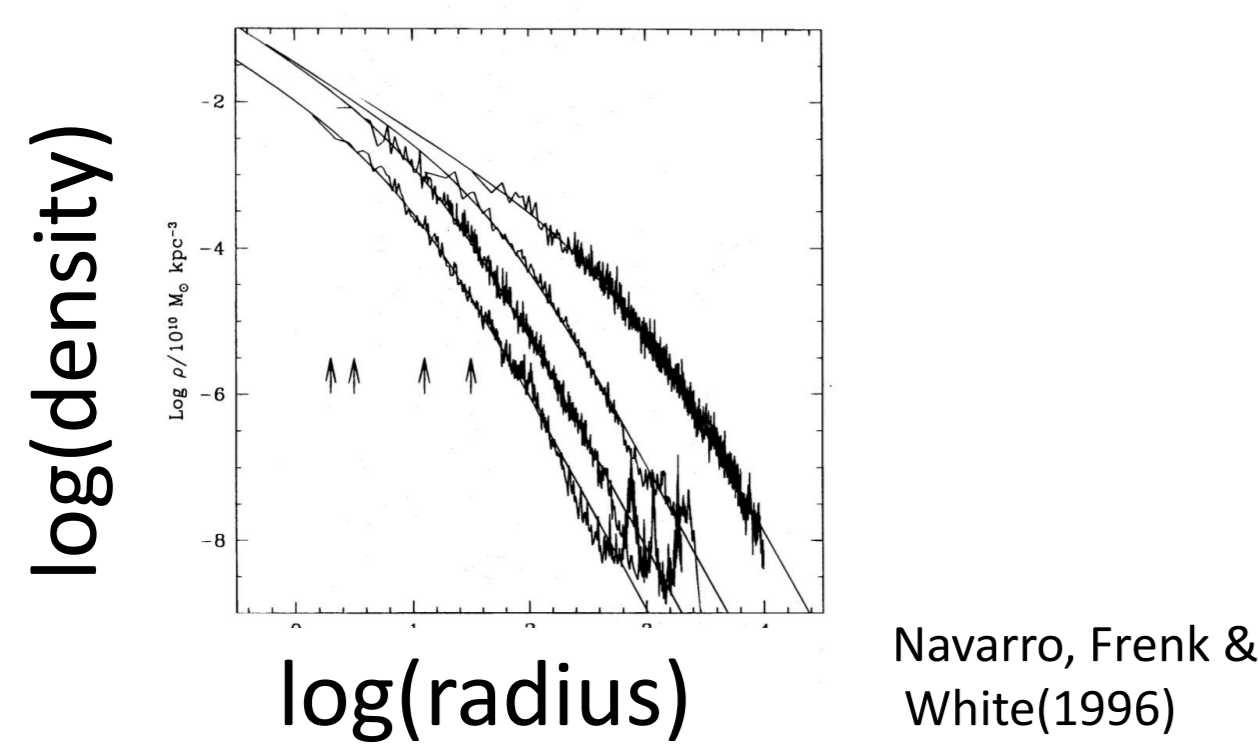
The dark matter haloes surrounding galaxies exhibit remarkable universal properties. Recently, Kaneda, Mori, and Otaki (2024) investigated the cusp-core problem of dark matter haloes by analysing their internal structures, derived from the ultra high resolution cosmological simulation by Ishiyama et al. (2021). The cusp-core problem refers to a discrepancy between theory and observation, wherein observations of dwarf galaxies often reveal flatter cores, despite simulations predicting a high density cusp at the centre. Their findings suggest that dark matter haloes associated with galaxies exceeding a mass of  $10^{11} M_{\odot}$  generally exhibit a cusp structure, whereas less massive galaxies typically display a core structure. In this study, we present a model of the cusp-core transition in cold dark matter haloes, identifying the critical halo mass at which this transition occurs and examining the physical conditions under which supernova feedback primarily drives this transition.

Background: Ishiyama et al. 2021

## Introduction

### Cusp-core problem

#### CDM simulations: cusp

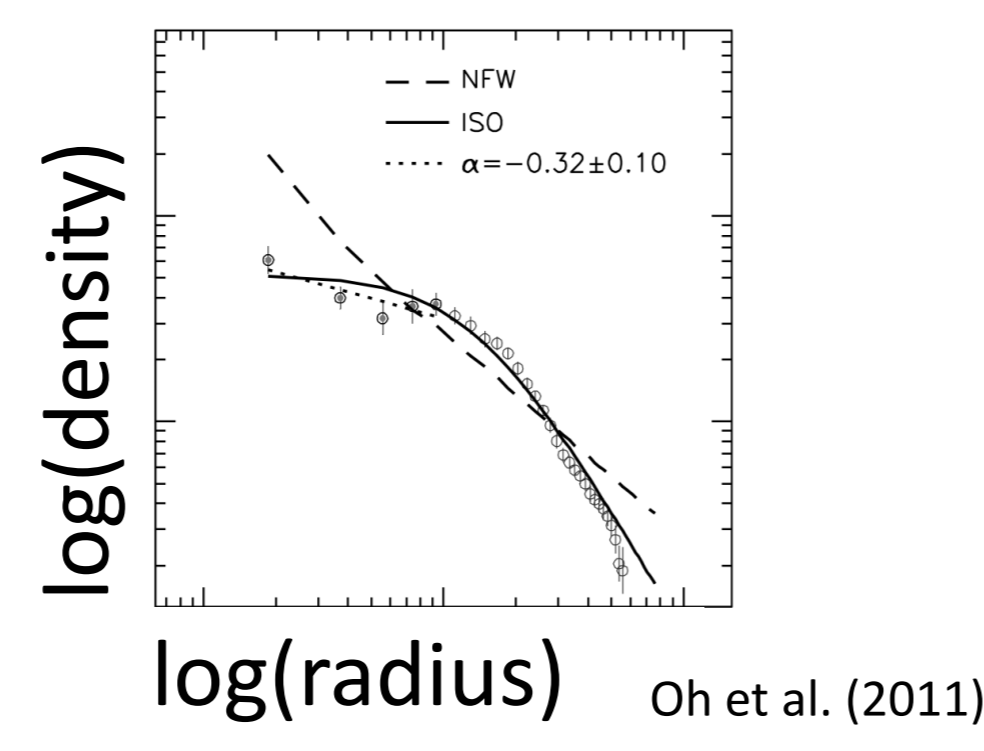


ex) NFW profile

$$\rho(r) = \frac{\rho_N r_N^3}{r(r+r_N)^2}$$

$r_N$ : scale radius,  $\rho_N$ : scale density

#### dwarf galaxy observation: core



ex) Burkert profile

$$\rho(r) = \frac{\rho_B r_B^3}{(r+r_B)(r^2+r_B^2)}$$

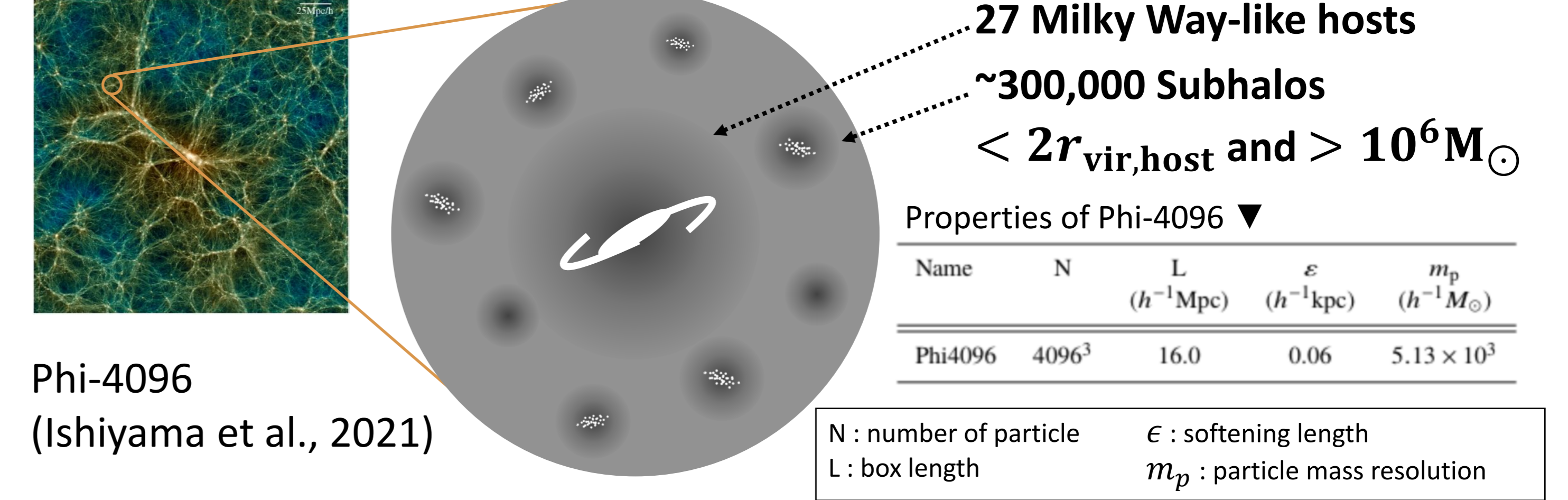
$r_B$ : core radius,  $\rho_B$ : central density

### Solutions

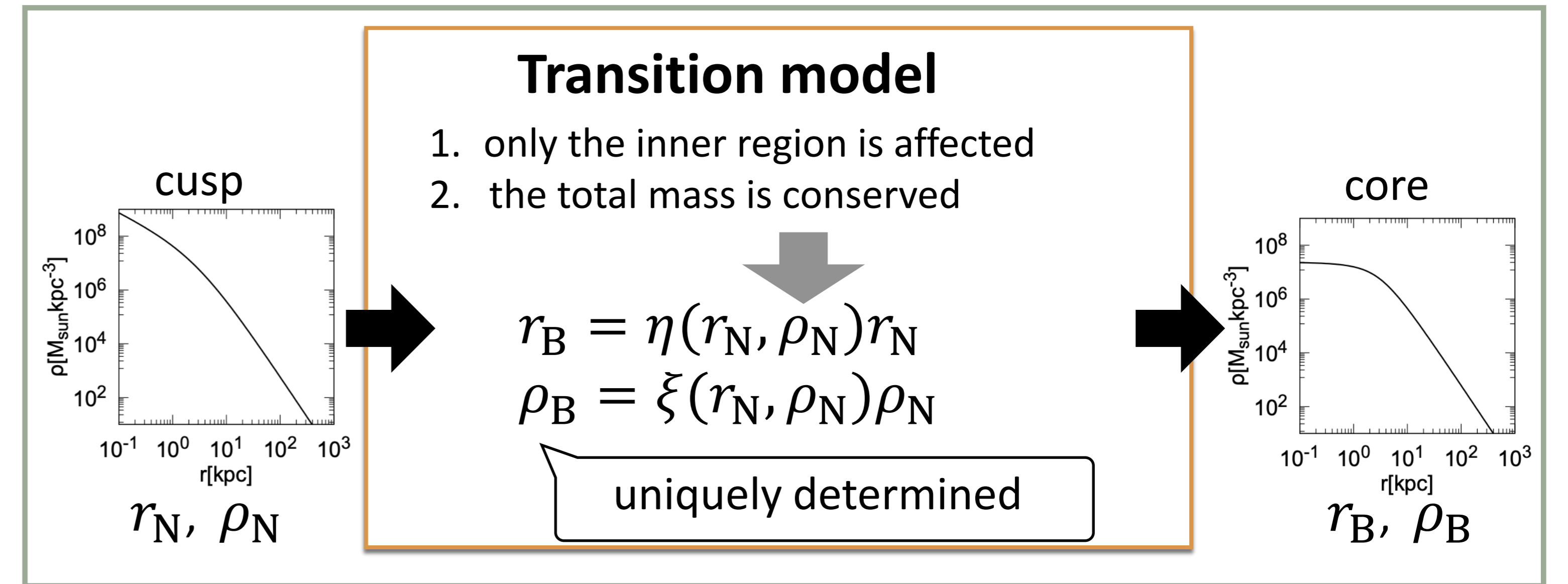
- Considering new candidates for dark matter particles that intrinsically produce cores
- **Taking the effect of baryonic feedback processes into account in the CDM simulations to alter cusps to cores**

## Methods

### Cosmological N-body simulation

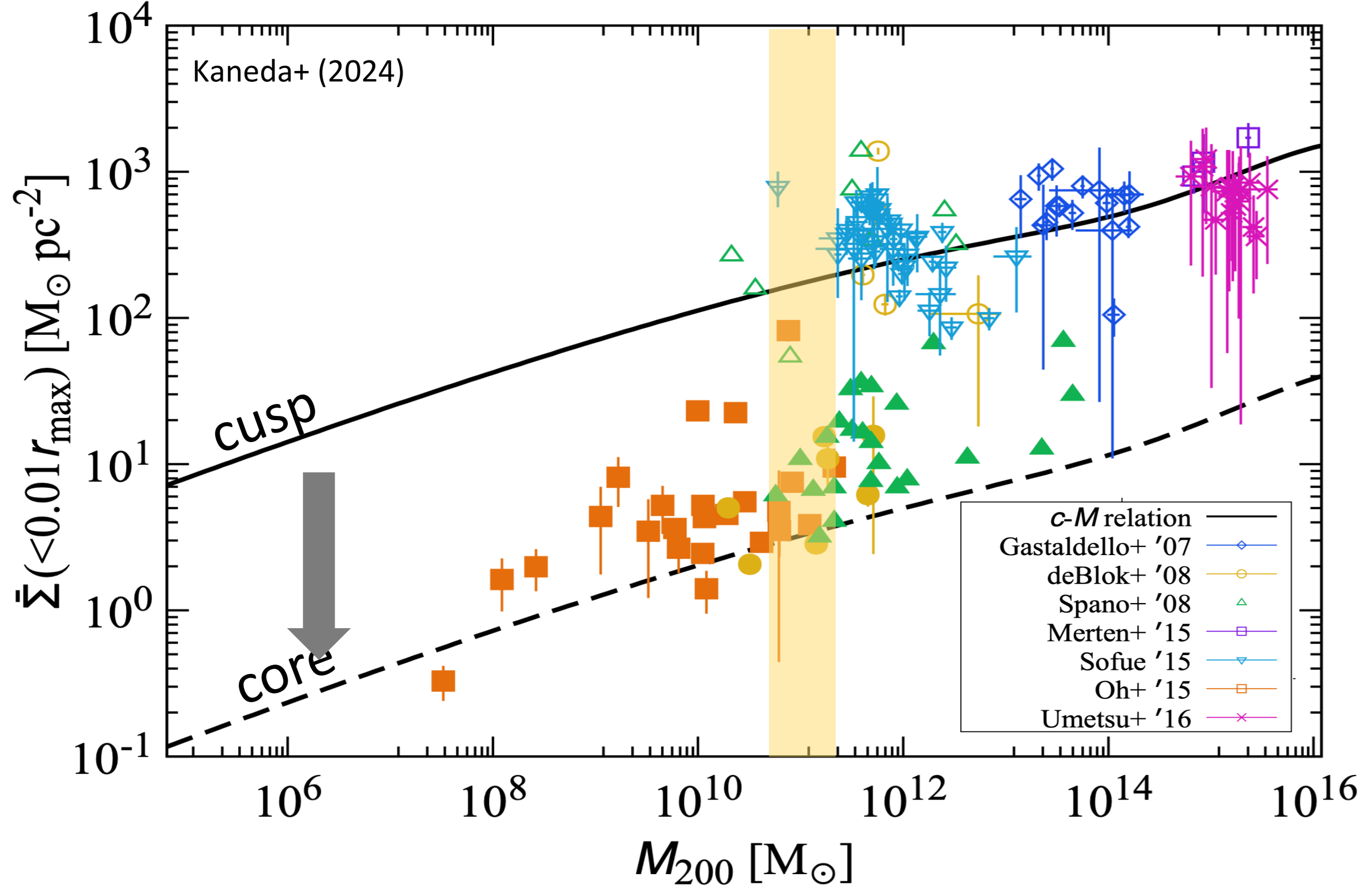


### Cusp-to-core transition model

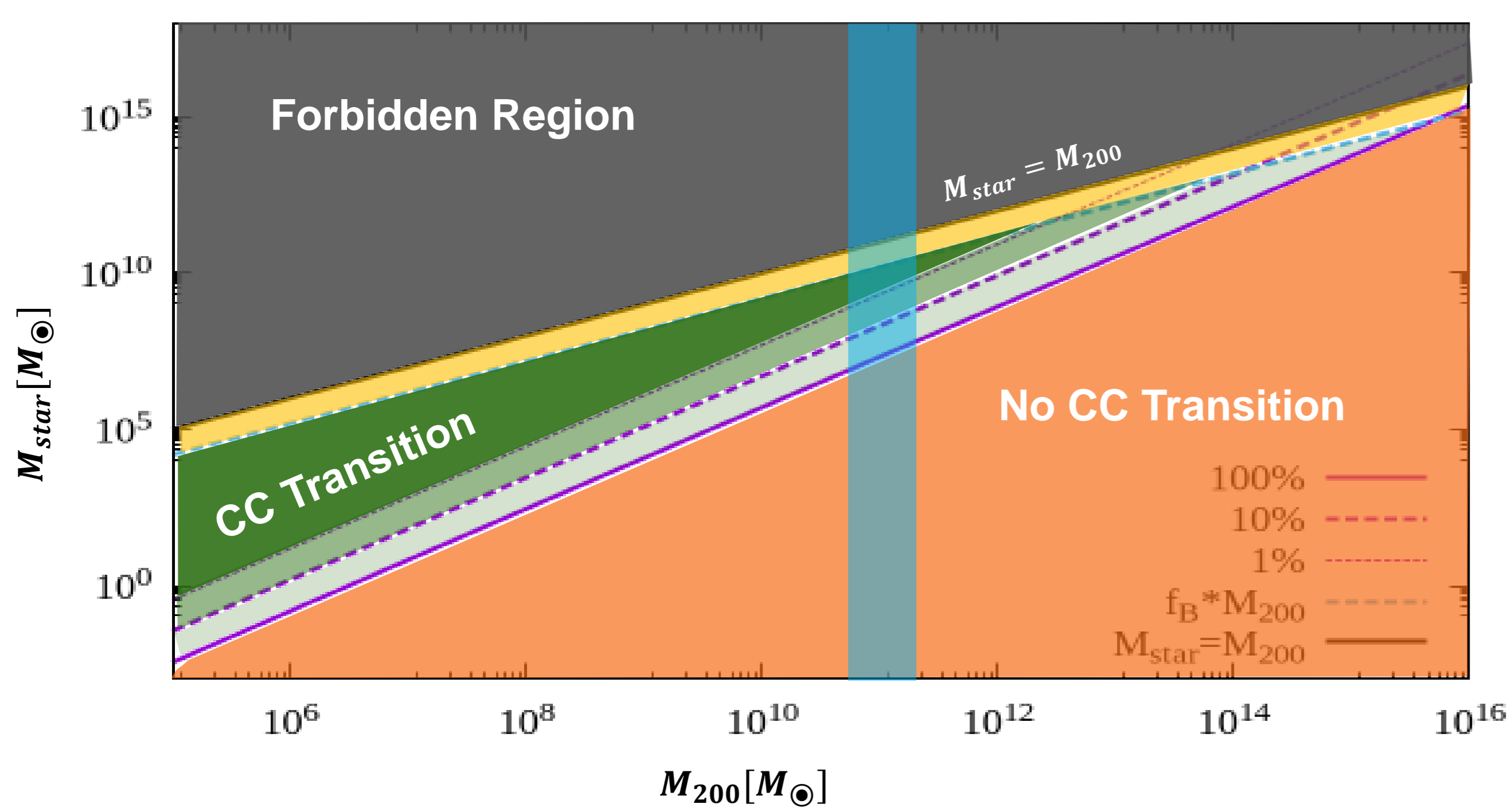


## Results

### Theoretical scaling relation with CC transition



### Critical mass and SN energy transport efficiency



### Physical conditions for the CC transition

$$M_* > M_{*,crit} = \frac{\Delta E}{\epsilon \epsilon_{51} f_{SN II}}$$

$\Delta E \equiv E_{BKT} - E_{NFW}$   
 $E_{BKT}, E_{NFW}$ : total energy  
 $\epsilon$ : energy conversion efficiency  
 $\epsilon_{51}$ : SN II energy,  $10^{51}$  erg  
 $f_{SN II}$ : Number of SN II per  $1 M_{\odot}$

### Chabrier's Initial Mass Function

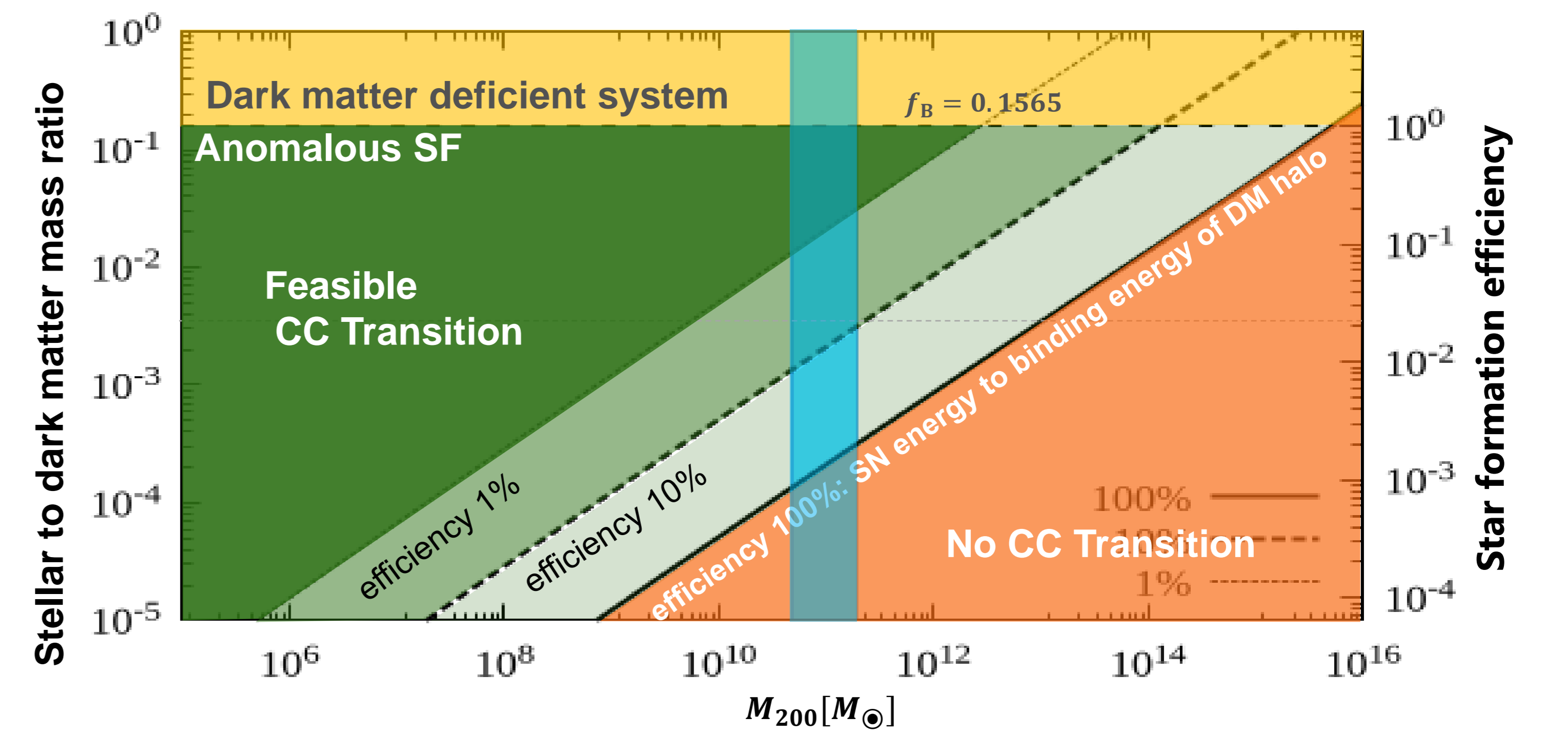
$$f_{SN II} = \int_{8 M_{\odot}}^{100 M_{\odot}} \phi(m) dm \quad \phi(m) \propto \begin{cases} \exp\left[-\frac{(\log m - \log m_c)^2}{2\sigma^2}\right] & \text{for } m/M_{\odot} \leq 1.0 \\ m^{-1.35 \pm 0.3} & \text{for } m/M_{\odot} > 1.0 \end{cases}$$

Chabrier+ (2005)

### Star Formation Efficiency

$$SFE = \frac{M_*}{f_B M_{200}} \quad f_B = \Omega_B / \Omega_m = 0.1565$$

$\Omega_B$ : baryon density parameter  
 $\Omega_m$ : mass density parameter



## Summary

- The  $c$ - $M$  relation is consistent with observations from dwarf galaxies to galaxy clusters.
- For cored haloes, a scaling relation based on the  $c$ - $M$  relation and our cusp-to-core transition model is proposed.
- This scaling relation aligns well with observations, and the transition from cusp to core is suggested to occur at  $M_{200} \sim 10^{11} M_{\odot}$ .
- This critical mass depends on SF efficiency and the efficiency of SN feedback transfer to the binding energy of the DM halo.