

# Sommerfeld Effect in Composite Dark Matter

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## Motivation

- ▶ Thermal freeze-out scenario is one of the most attractive scenarios for the production of Dark Matter (DM) in the early Universe.

$$\Omega_{\text{DM}} h^2 \sim 0.12 \frac{2 \times 10^{-26} \text{ cm}^3/\text{s}}{\langle \sigma v \rangle_{\text{ann}}}, \quad \langle \sigma v \rangle_{\text{ann}} \sim \frac{\alpha_{\text{ann}}^2}{m_{\text{DM}}^2}.$$

- ▶ Indirect detection sensitive to DM with  $\mathcal{O}(1-100)$  TeV mass [e.g. MAGIC Collaboration (2023), H.E.S.S. Collaboration (2018)]
- ▶  $\alpha_{\text{ann}} \sim \mathcal{O}(1)$  matches the relic abundance for  $m_{\text{DM}} \gtrsim \mathcal{O}(10)$  TeV.
- ▶ DM model with QCD-like gauge theory is a natural choice.  
→ **Composite DM Model**

## Composite DM Model

3-flavors of vector-like fermions  $\psi$  (dark quark),  $\bar{\psi}$  (anti-dark quark) and the  $SU(N_c)$  gauge interaction. [Bai and Hill (2010), Antipin, Redi, Strumia and Vigiani (2015)]

	$SU(N_c)$	$SU(3)_c$	$SU(2)_W$	$U(1)_Y$
$\psi$	$N_c$	<b>1</b>	<b>3</b>	0
$\bar{\psi}$	$\bar{N}_c$	<b>1</b>	<b>3</b>	0

Chiral symmetry breaking:  $SU(3)_L \times SU(3)_R \rightarrow SU(3)_V \supset SU(2)_W \rightarrow SU(2)_W$  triplet  $\chi$  ( $G$ -parity odd) and quintuplet  $\pi$  ( $G$ -parity even) dark pions

Chiral Lagrangian

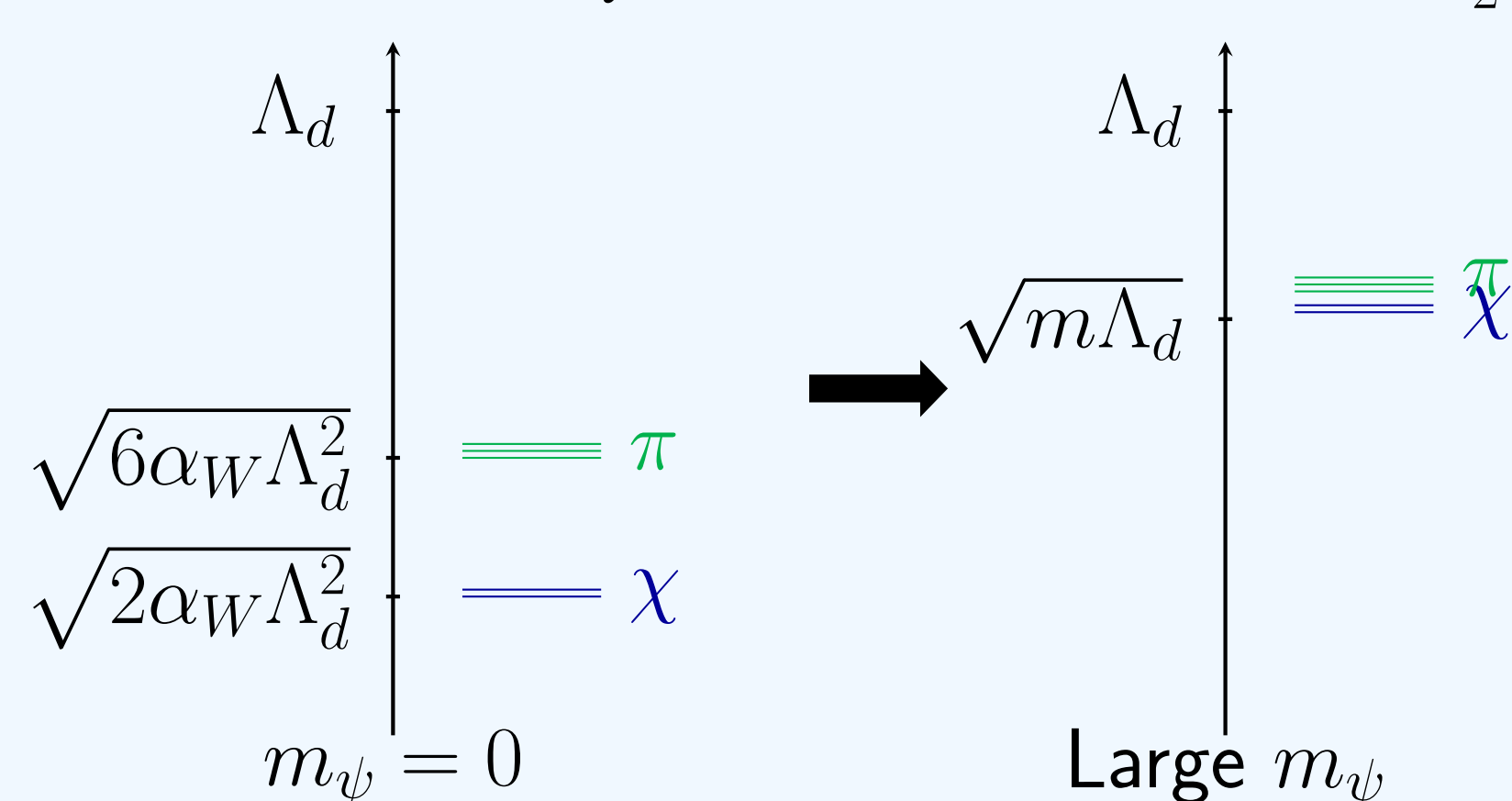
$$\mathcal{L} = \frac{f_d^2}{4} \text{Tr}[D_\mu U D^\mu U^\dagger] + v_d^3 \text{Tr}[MU + \text{h.c.}] + \mathcal{L}_{\text{WZW}}$$

- ▶  $\chi$  is stable because of  $G$ -parity :  $U \rightarrow U^T$
- ▶  $\pi$  decays to EW gauge bosons via WZW term.

## Previous Work [Abe, Sato and TY, JHEP 09(2024)]

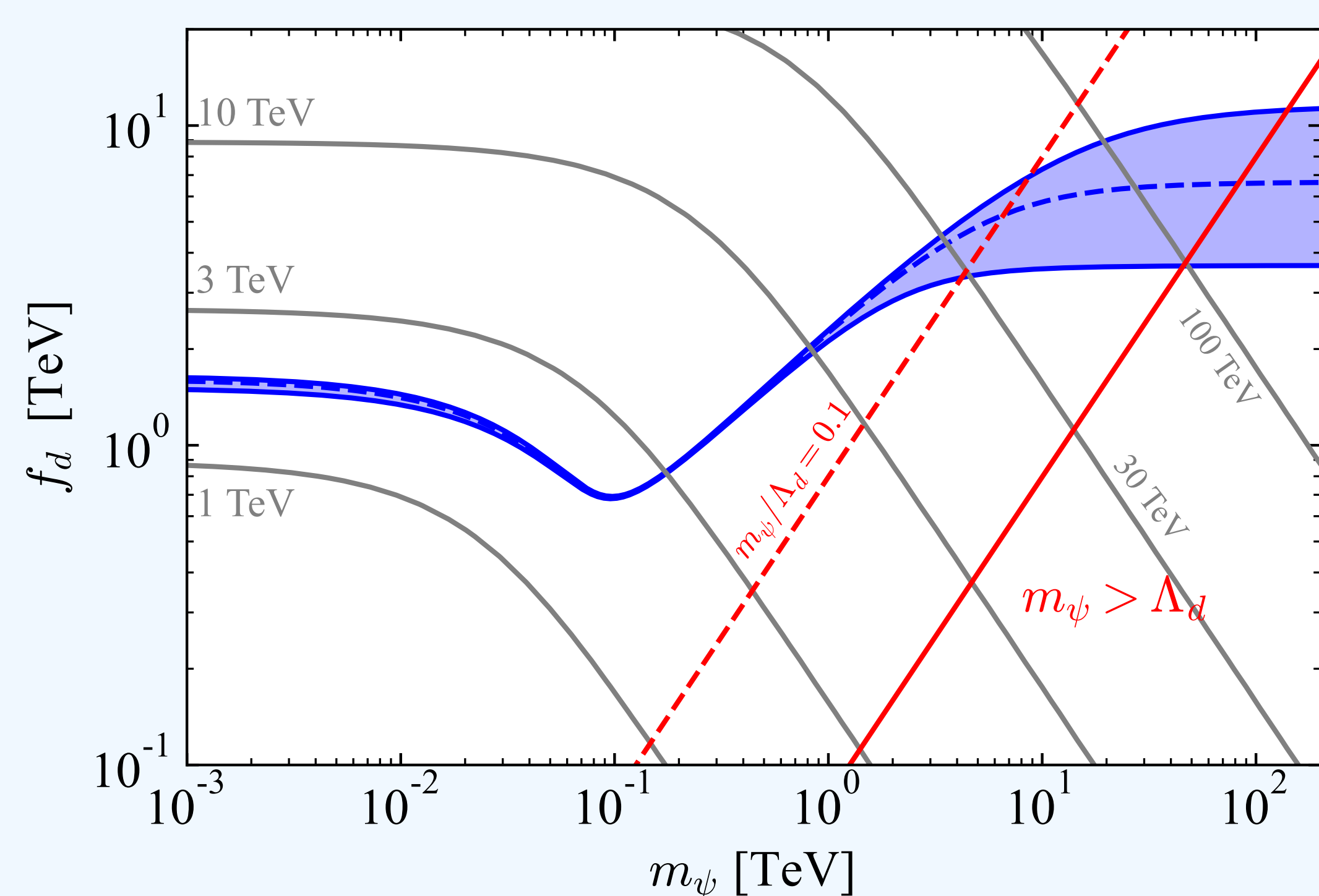
### Sources of dark pion mass

- ▶ Dark quark mass:  $m_\pi^2 \sim m_\psi \Lambda_d$
- ▶  $SU(2)_W$  radiative correction:  $\delta m_\pi^2 \sim C^2(R) \alpha_W \Lambda_d^2$
- ▶ EW symmetry breaking:  $\Delta \equiv m_Q - m_0 \sim Q^2 \alpha_W m_W \sin^2 \frac{\theta_W}{2} \sim \mathcal{O}(100)$  MeV



### DM annihilation (Leading Order)

- ▶  $m_\psi = 0$ :  $\chi\chi \rightarrow WW$   $\langle \sigma v \rangle_{WW} \propto \frac{\alpha_W^2}{m_\chi^2} \Rightarrow m_\chi \sim 1.8$  TeV.
  - ▶ Large  $m_\psi$ :  $\chi\chi \rightarrow \pi\pi$   $\langle \sigma v \rangle_{\pi\pi} \propto \frac{m_\chi^2}{f_d^4} \exp\left(-\frac{m_\pi - m_\chi}{T}\right) \Rightarrow m_\chi \sim \mathcal{O}(1-10)$  TeV
- Forbidden channels** [Abe, Sato and TY (2024)]

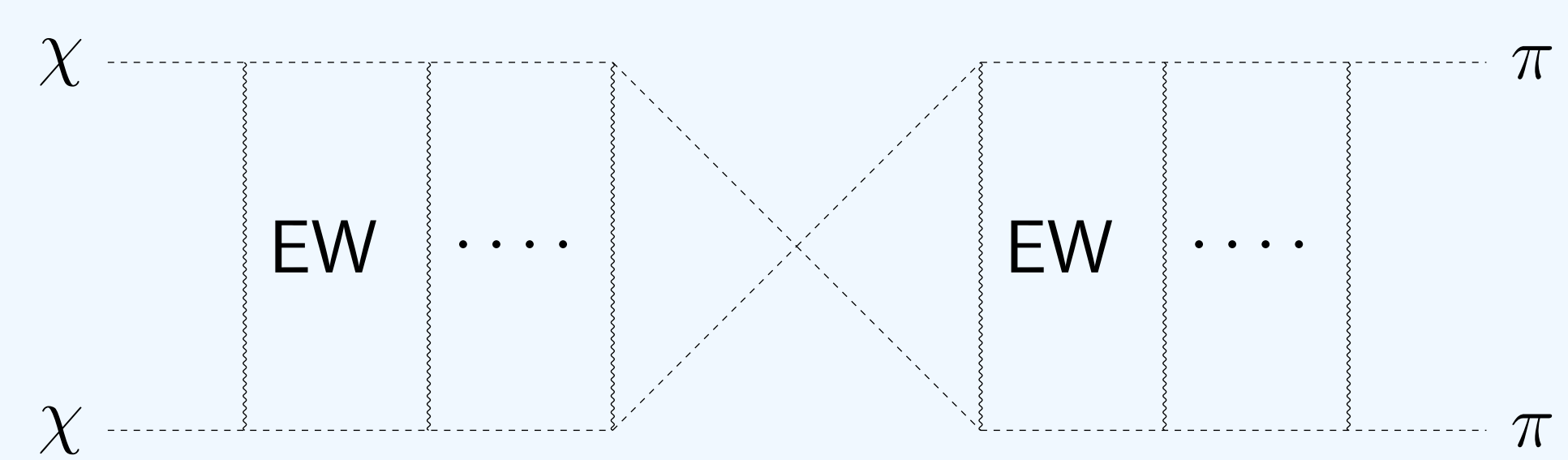


Gray contour indicates DM mass  $m_\chi$ . Chiral Lagrangian is valid for  $m_\psi \lesssim \Lambda_d$ .

**Forbidden channels can determine the DM abundance.**

## Sommerfeld Effect

Heavy EW interacting DM  $\rightarrow$  **Sommerfeld effect (SE)** must be taken into account. [Hisano, Matsumoto, Nojiri and Saito (2005)]



EW interactions also affect final two-body states in forbidden channels  
→ **SE in final states** [Cui and Luo (2021)]

## Toy Model

We consider  $\chi\chi \rightarrow \pi\pi$  via interaction terms w/o derivative:

$$\mathcal{L}_{\text{int}} = \frac{4m v_d^3}{3f_d^4} \left( \frac{1}{2}(\chi^0)^2 + \chi^+\chi^- \right) \left( \frac{1}{2}(\pi^0)^2 + \pi^+\pi^- + \pi^{++}\pi^{--} \right)$$

### Mixing of two-body states by EW interactions

- ▶  $\chi^0\chi^0 \leftrightarrow \chi^+\chi^-$
- ▶  $\pi^0\pi^0 \leftrightarrow \pi^+\pi^- \leftrightarrow \pi^{++}\pi^{--}$

SE factor  $\leftarrow$  solving Schroedinger eqns.

$$\left[ -\frac{1}{2\mu_\chi} \frac{d^2}{dr^2} + \mathbf{V}_\chi(r) - \frac{p_\chi^2}{2\mu_\chi} \right] \vec{\psi}_\chi(r) = 0,$$

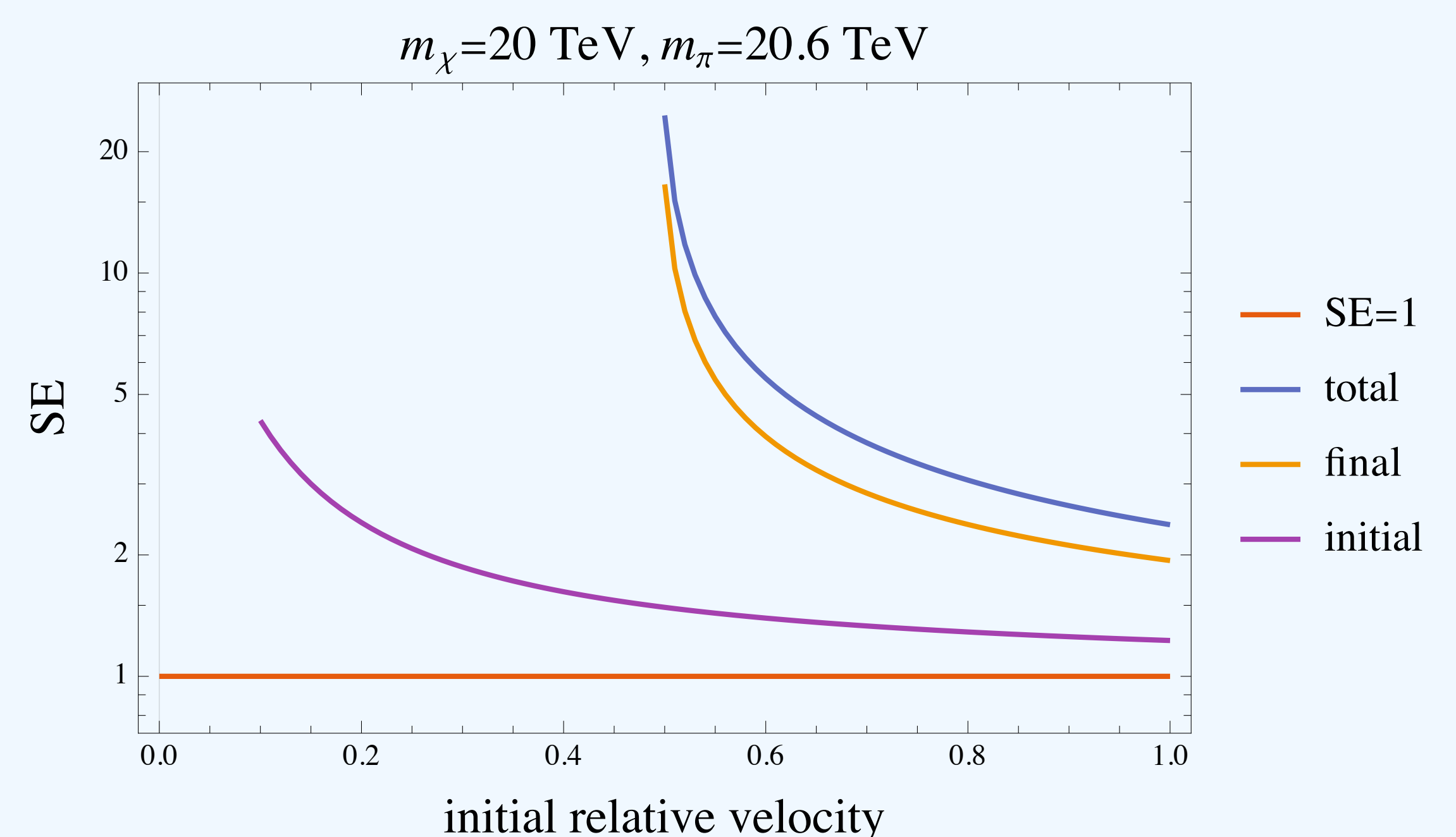
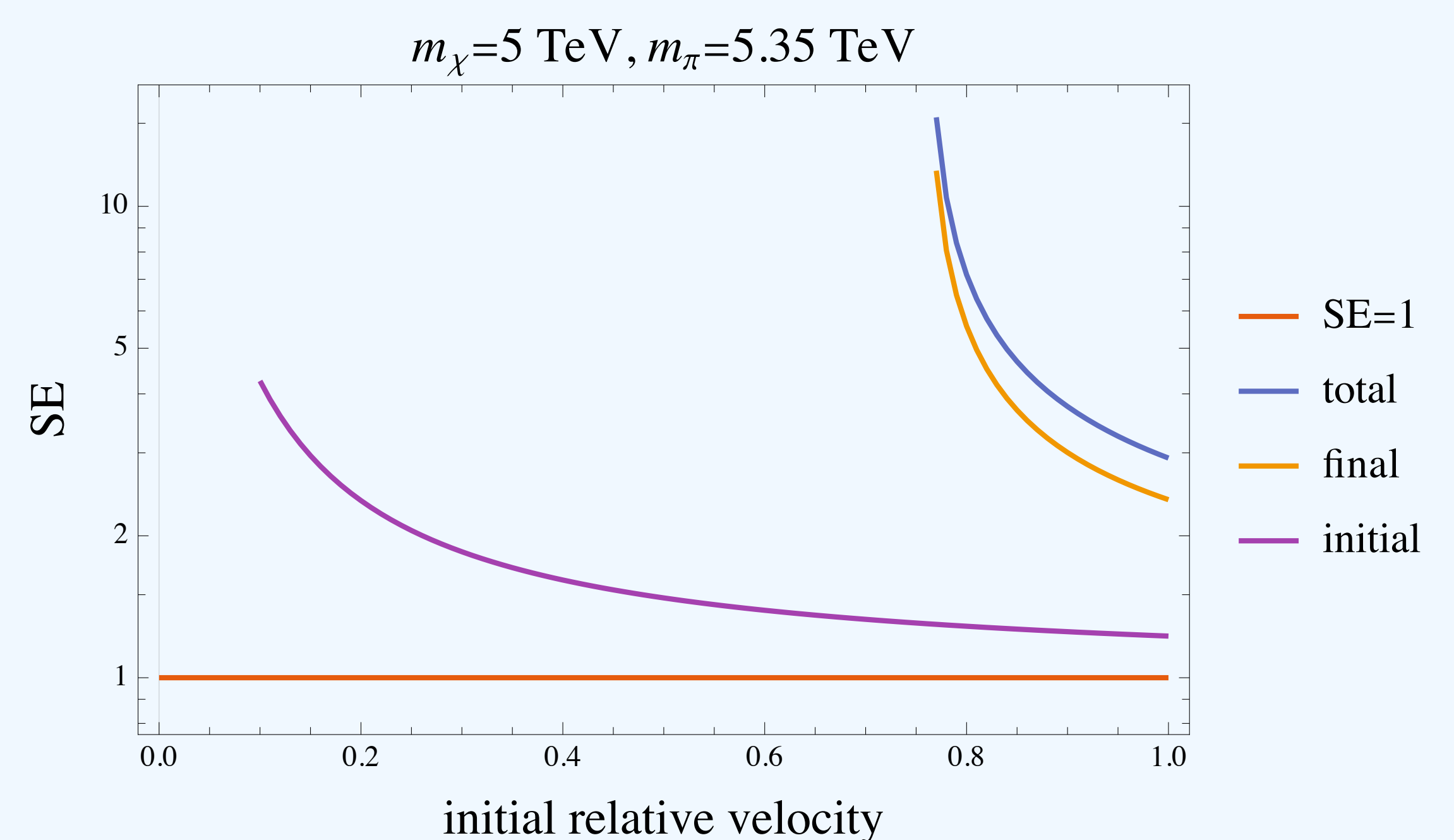
$$\mathbf{V}_\chi(r) \equiv \begin{pmatrix} 0 & -\sqrt{2}B \\ -\sqrt{2}B & -A + 2\Delta \end{pmatrix}, \quad \vec{\psi}_\chi(r) \equiv \begin{pmatrix} \psi_\chi^{00}(r) \\ \psi_\chi^{\pm\pm}(r) \end{pmatrix},$$

$$\left[ -\frac{1}{2\mu_\pi} \frac{d^2}{dr^2} + \mathbf{V}_\pi(r) - \frac{p_\pi^2}{2\mu_\pi} \right] \vec{\psi}_\pi(r) = 0,$$

$$\mathbf{V}_\pi(r) \equiv \begin{pmatrix} -4A + 8\Delta & -2B & 0 \\ -2B & -A + 2\Delta & -3\sqrt{2}B \\ 0 & -3\sqrt{2}B & 0 \end{pmatrix}, \quad \vec{\psi}_\pi(r) \equiv \begin{pmatrix} \psi_\pi^{\pm\pm}(r) \\ \psi_\pi^{\pm\mp}(r) \\ \psi_\pi^0(r) \end{pmatrix}$$

$$A \equiv \frac{\alpha}{r} + \frac{\alpha_W c_W^2}{r} e^{-m_Z r}, \quad B \equiv \frac{\alpha_W}{r} e^{-m_W r}$$

**Amplitude w/ SE factor** :  $\mathcal{M} \sim \vec{\psi}_\chi^\dagger(0) \mathcal{M}_{\text{LO}} \vec{\psi}_\pi(0)$



**Final state SE dominantly contributes to the total SE factor.**

## Future Work

- ▶ Adding interaction terms w/ derivative.
- ▶ Comparing w/ the current constraints from indirect detection experiments.