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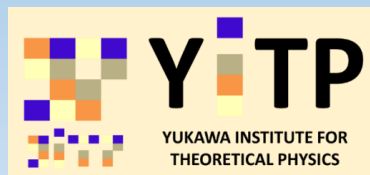
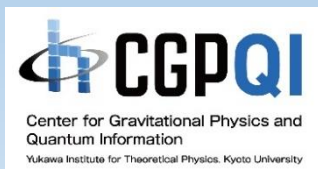
Hitoshi-fest: from particle physics,
cosmology to instrumentation
Dec.16-20, 2024

Emergence of Gravitational Spacetime from Quantum Information

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My wonderful days at IPMU [2008-2012]



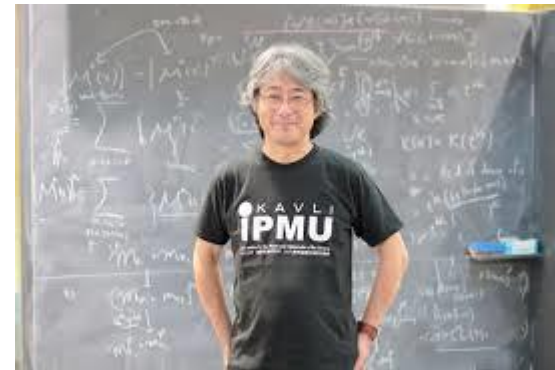
◆ I thank very much

- Great colleagues
- A lot of research time
- Tea time discussions

◆ I learned a lot from opportunities of

- Supervising post-docs and students
- Organizing international workshops
- Watching Murayama-san's effort to make IPMU internationally prominent

The valuable experiences in IPMU have been very helpful for me after I move to YITP, Kyoto.

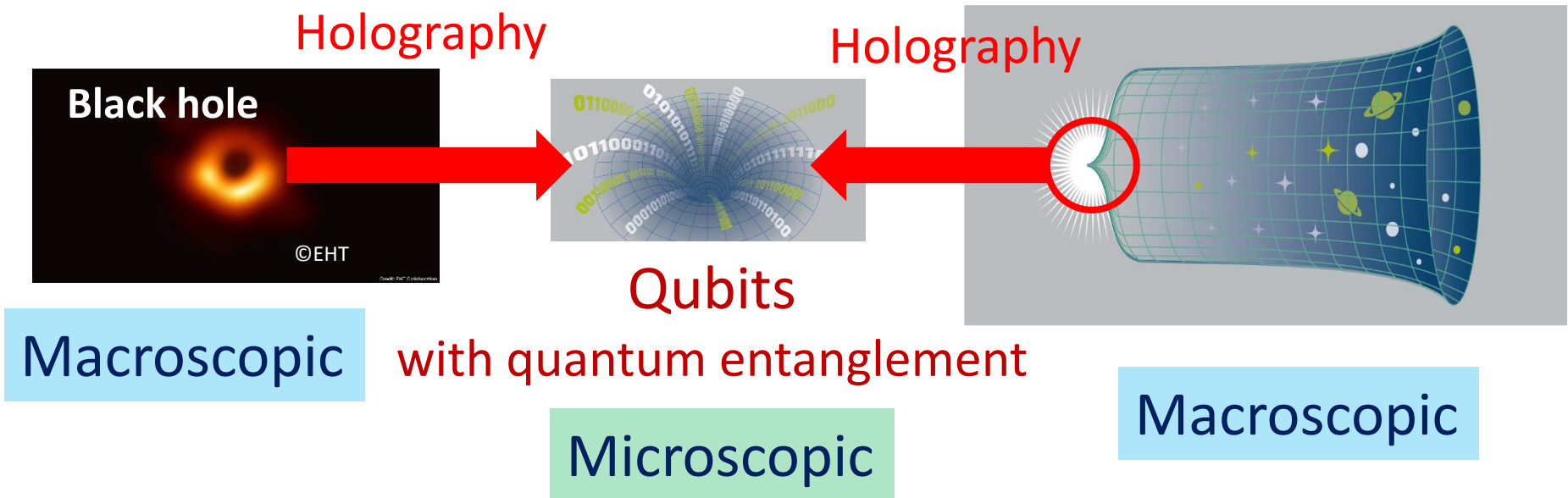


I am extremely grateful to Murayama-san for founding this amazing institute and for providing us with this ideal environment for researchers !

① Introduction

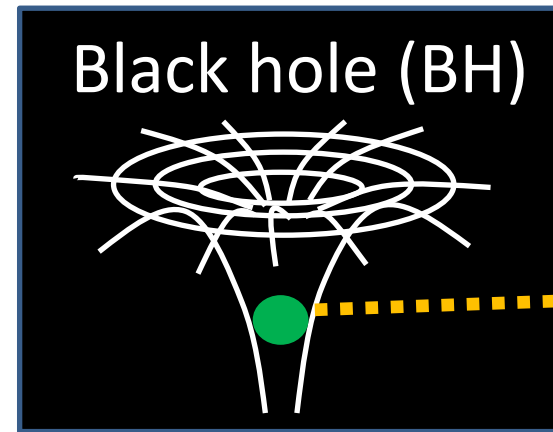
Holography is a very promising approach to quantum gravity, towards understanding of the creation of the Universe.

This is because holography is much like *a microscope* in thought experiments of quantum gravity.



Macroscopic Side: Black hole Entropy

A huge amount
of quantum information
is hidden inside black hole !



$$S_{BH} = \frac{k_B c^3}{\hbar} \times \frac{A_{BH}}{4G_N}$$



BH thermodynamics !

[Bekenstein, Hawking, 1972–1976]

A_{BH} = Surface Area of Black hole \Rightarrow Geometry

G_N = Newton constant \Rightarrow Gravity

\hbar = Planck constant \Rightarrow Quantum Mechanics

k_B = Boltzmann const. \Rightarrow Stat. Mech. , Quantum Info.

Quantum
Gravity!

Microscopic Side: Quantum Entanglement (QE)

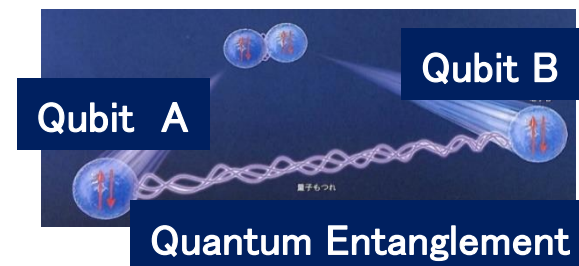
Quantum Entanglement \rightarrow Quantum correlations

$$|\Psi\rangle_{AB} \neq |\Psi_1\rangle_A \otimes |\Psi_2\rangle_B .$$

between A and B.

e.g. Bell state: $|\Psi_{Bell}\rangle = \frac{1}{\sqrt{2}} \left[|\uparrow\rangle_A \otimes |\downarrow\rangle_B + |\downarrow\rangle_A \otimes |\uparrow\rangle_B \right]$

\rightarrow **Minimal Unit of Entanglement**
 \leftrightarrow **Planck length**



The best (or only) measure of quantum entanglement for pure states is known to be **entanglement entropy (EE)**.

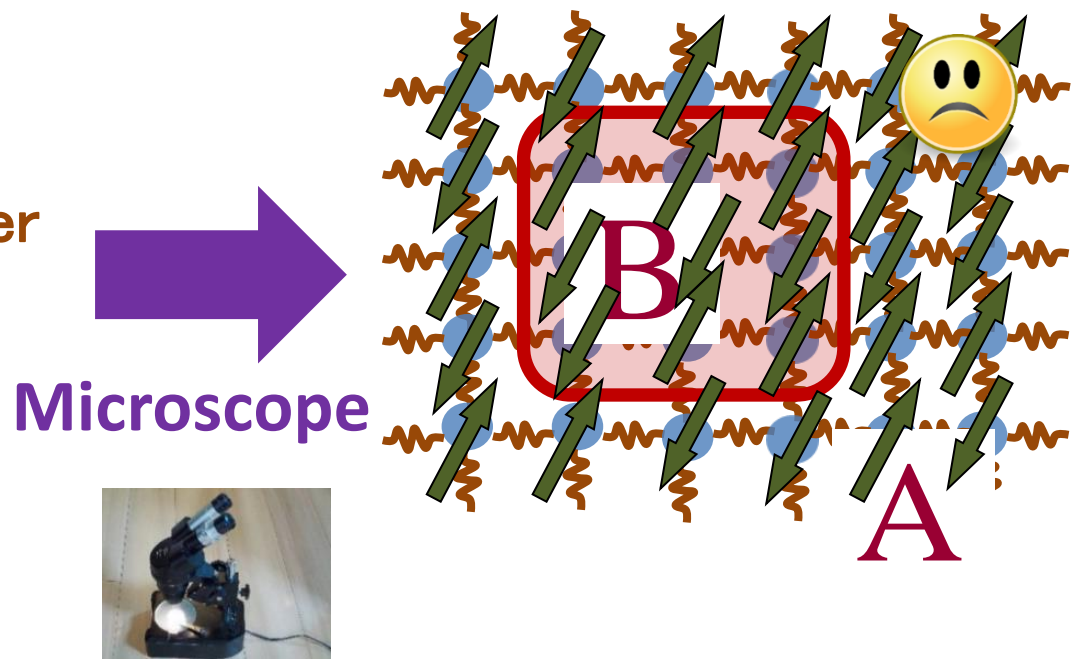
$$\mathbf{EE = S_A = \# \text{ of Bell states between A and B}}$$

From BHs to Qubits

Blackhole Spacetime

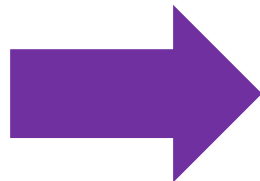


Quantum Many-body System



BH entropy SBH

Area law



Entanglement entropy SA

Area law

What is a useful microscope in quantum gravity ?

② Holography and Quantum Entanglement

Holography as microscope in quantum gravity

**BH Entropy
Formula**

$$S_{BH} = \frac{A_{BH}}{4G_N}$$



**Degrees of freedom
in Gravity \propto Area**

[' t Hooft 1993, Susskind 1994]

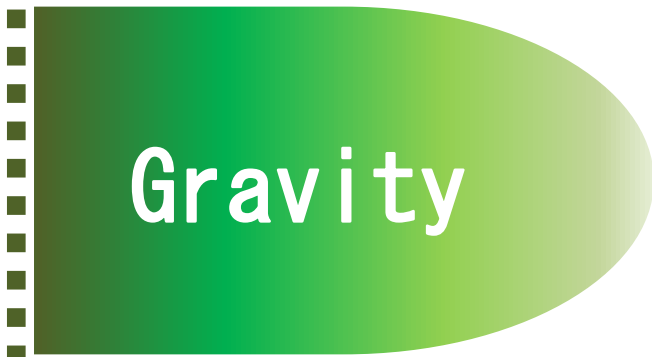
Holography

Gravity on M

=

Quantum Matter on ∂M

Boundary of M



Gravity

=

Matter



BH entropy (\propto Area) = Thermal Entropy of Matter (\propto Volume)

Gauge/Gravity Duality (AdS/CFT) – [Maldacena 1997]

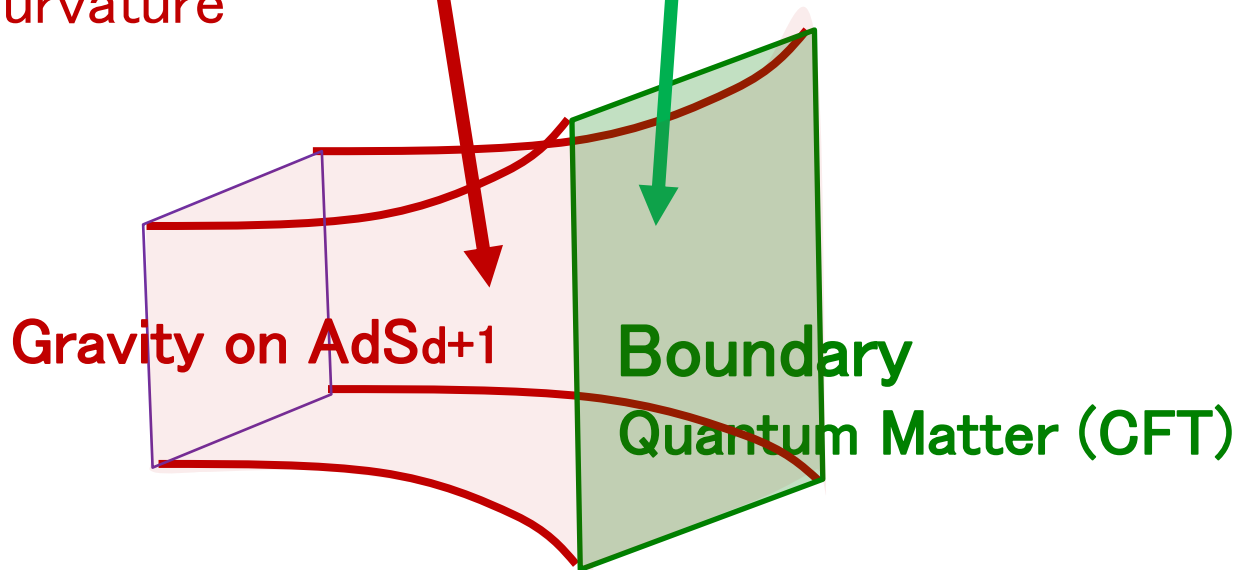
(Quantum) Gravity on
d+1 dim. Anti-de Sitter Space

=

d dim. Gauge theory
(or Conformal field theory)

Anti-de Sitter space (AdS)
→ Universe with
negative curvature
 $\Lambda < 0$

Conformal Field Theory (CFT)
→ quantum field theory
with only massless fields



Mining Qubits from Gravity

Holographic Entanglement Entropy (HEE)

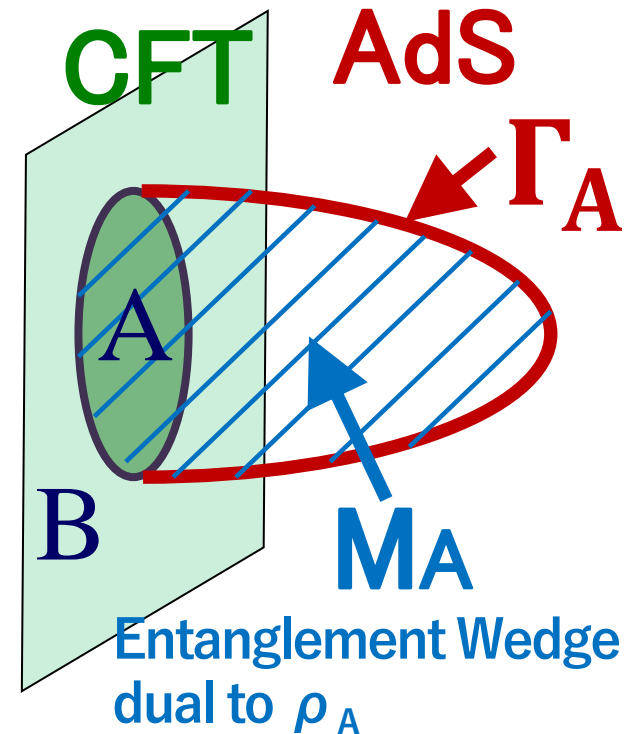
[Ryu-TT 2006, Hubeny-Rangamani-TT 2007]

$$S_A = \text{Min Ext}_{\Gamma_A} \left[\frac{\text{Area}(\Gamma_A)}{4G_N} \right]$$

This formula can compute entanglement entropy in time-dependent backgrounds.



Real time (Lorentzian) evolution



$$\rho_A = \text{Tr}_B \left[|\Psi_{tot}\rangle\langle\Psi_{tot}| \right]$$

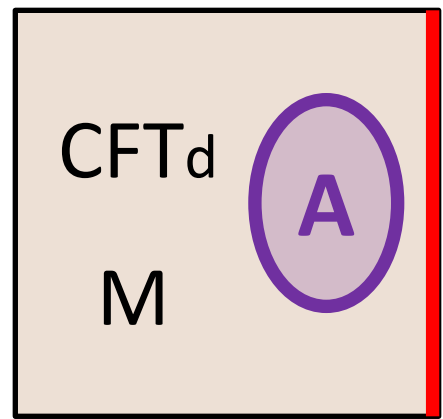
$$S_A = -\text{Tr}_A \rho_A \log \rho_A .$$

Extension 1: HEE in AdS/BCFT [TT 2011, Fujita-Tonni-TT 2011]

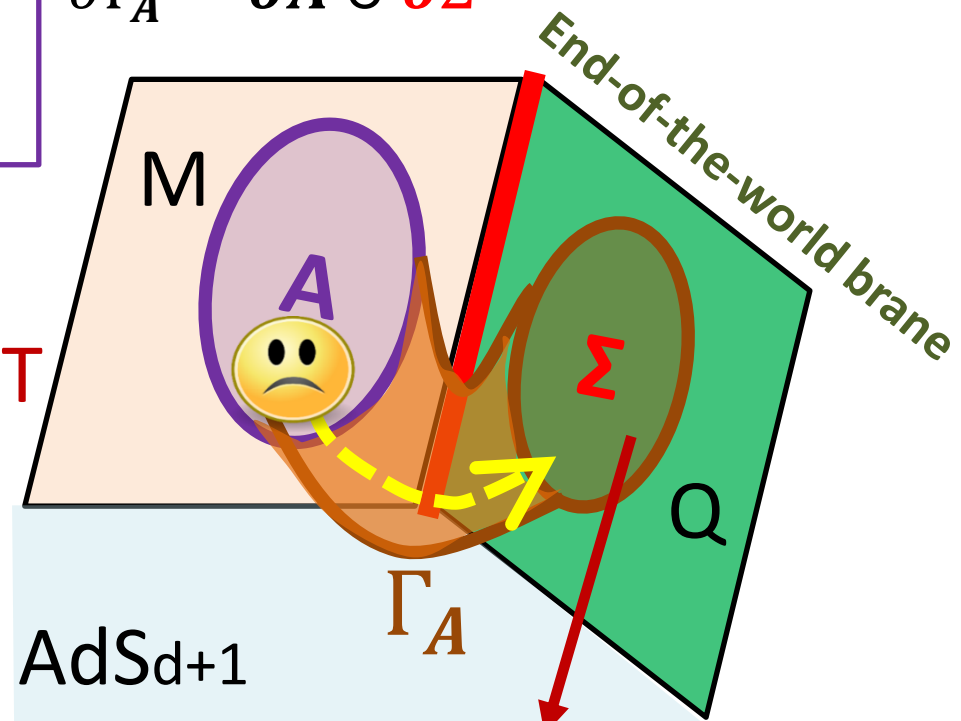
Q. What is holography and HEE for CFT on a space with boundaries ?

$$S_A = \text{Min Ext}_{\Gamma_A, \Sigma} \left[\frac{\text{Area}(\Gamma_A)}{4G_N} \right]$$

$$\partial\Gamma_A = \partial A \cup \partial\Sigma$$



AdS/BCFT =



Extremal Surfaces can end on Q !

The region **Σ** is now known as an **Island** ! ➡ A key to resolve BH information paradox
 [Penington 2019, Almheiri et.al. 2019]
 😞 can see BH interior !

Extension 2: HEE for non-Hermitian density matrices

Q. What is HEE in a *Euclidean time-dependent AdS* geometry ?

Holographic Pseudo Entropy

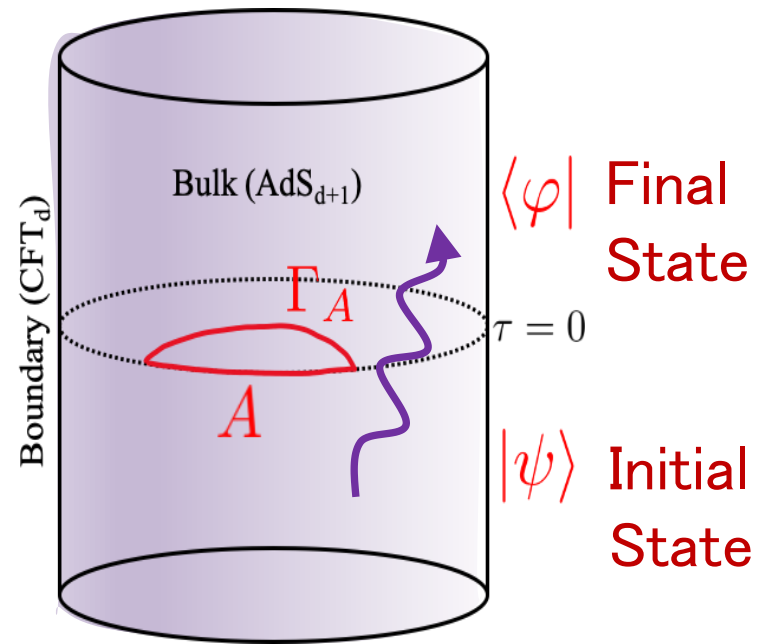
[Nakata-Taki-Tamaoka-Wei-TT, 2020]

$$S(\mathcal{T}_A^{\psi|\varphi}) = \min_{\Gamma_A} \frac{\text{Area}(\Gamma_A)}{4G_N}$$

Transition matrix

$$H_{tot} = H_A \otimes H_B \cdot \tau^{\psi|\varphi} = \frac{|\psi\rangle\langle\varphi|}{\langle\varphi|\psi\rangle}$$

$$\tau_A^{\psi|\varphi} = \text{Tr}_B \left[\tau^{\psi|\varphi} \right]$$

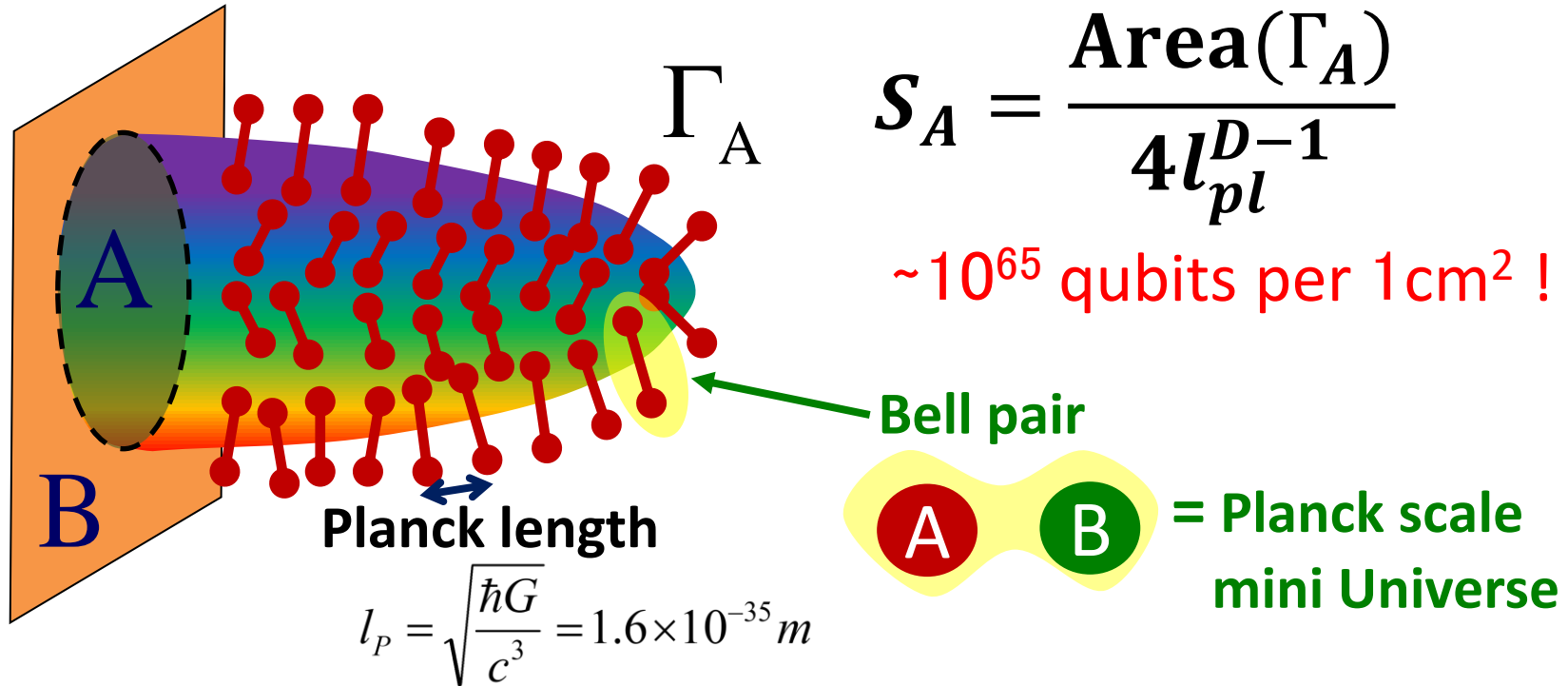


Pseudo entropy:

$$S \left(\tau_A^{\psi|\varphi} \right) = -\text{Tr} \left[\tau_A^{\psi|\varphi} \log \tau_A^{\psi|\varphi} \right].$$

In general, complex valued.

HEE/HPE suggests that in spacetimes with $\Lambda < 0$, there is one qubit of entanglement for each Planck length area !



$\Lambda < 0$ spacetime may emerge from entangled Qubits !

Q. How about de Sitter space $\Lambda > 0$?

③ De Sitter (dS) Holography and Holographic Entropy

An extension of AdS/CFT to de Sitter space (dS)

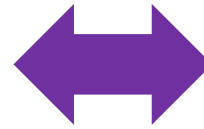
is known as the dS/CFT correspondence. [Strominger 2001, Maldacena 2002]

A Sketch of dS/CFT

Equivalent
 dS_{d+1}/CFT_d

**d+1 dim. Lorentzian
de-Sitter spacetime**

**Euclidean d. dim CFT
on S^d**



Bulk space

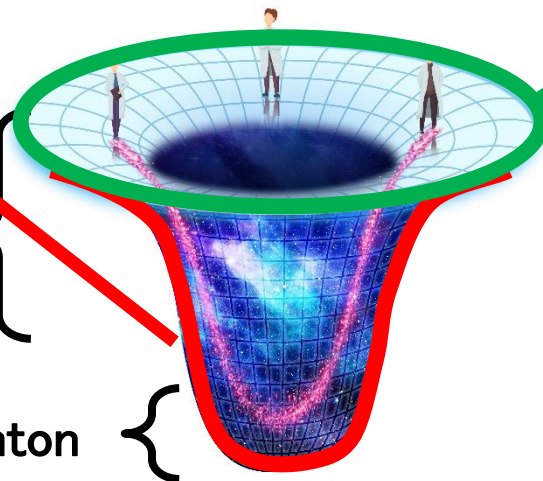
Lorentzian de Sitter space

$$ds^2 = L_{ds}^2(-dt^2 + \text{Cosh}^2 t d\Omega^2)$$

Euclidean instanton

Space-like bdy

Time



Why dS/CFT is much more difficult than AdS/CFT ?

[1] Dual Euclidean CFTs should be exotic and non-unitary !

A “standard” Euclidean CFTs is dual to gravity on hyperbolic space.
e.g. dS3/CFT2 → *Imaginary valued* central charge $c \approx i \frac{3L_{dS}}{2G_N}$!

[2] Time should emerge from Euclidean CFT !

From a usual Euclidean CFT, a space-like direction will emerge as RG scale.

How does a *time-like direction emerge* from a Euclidean CFT ?

[3] Entanglement entropy becomes complex valued !

Extremal surfaces in dS which end on its boundary are *time-like* !

Non-unitary CFT dual of 3 dim. dS

[Hikida–Nishioka–Taki–TT, 2021, 2022]

Large c limit of $SU(2)_k$ WZW model (a 2dim. CFT)
 = **Einstein Gravity** on 3 dim. de Sitter (radius L_{ds})

Level $k \approx -2 + \frac{4iG_N}{L_{ds}}$ **Central charge** $c = \frac{3k}{k+2} \approx i \frac{3L_{ds}}{2G_N}$

$$Z[S^3, R_j] = |S_j^0|^2 \approx e^{\frac{\pi L_{ds}}{2G_N} \sqrt{1-8G_N E}}$$

CFT partition function De Sitter Entropy

This non-unitary CFT is equivalent to the Liouville CFT

at $b^{-2} \approx \pm \frac{i}{4G_N}$ $I_{CFT}[\phi] = \int d^2x \left[\frac{1}{4\pi} (\partial_a \phi \partial_a \phi) + \underline{\mu e^{2b\phi}} \right]$.

complex !

[→Reproduced by Verlinde–Zhang 2024 via the Double Scaled SYK]



Holographic Entanglement Entropy in dS3/CFT2 ?

[Doi-Harper-Mollabashi-Taki-TT 2022, 2023]

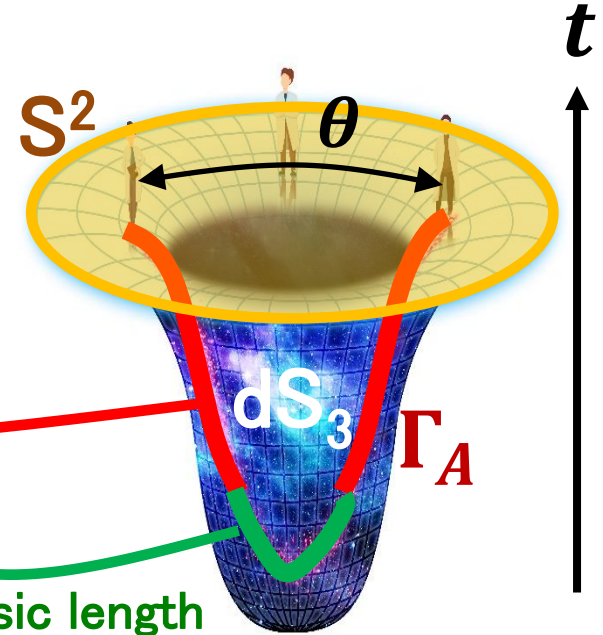
In dS3/CFT2, the geodesic Γ_A becomes time-like and we find:

$$S_A = \frac{L(\Gamma_A)}{4G_N} = i \frac{C_{ds}}{3} \log \left(\frac{2}{\epsilon} \sin \frac{\theta}{2} \right) + \underbrace{\frac{C_{ds}}{6} \pi}_{S_{dS}/2}$$

Time-like geodesics length
→ imaginary part !

Agree !

Space-like geodesic length
→ Real part



CFT calculation

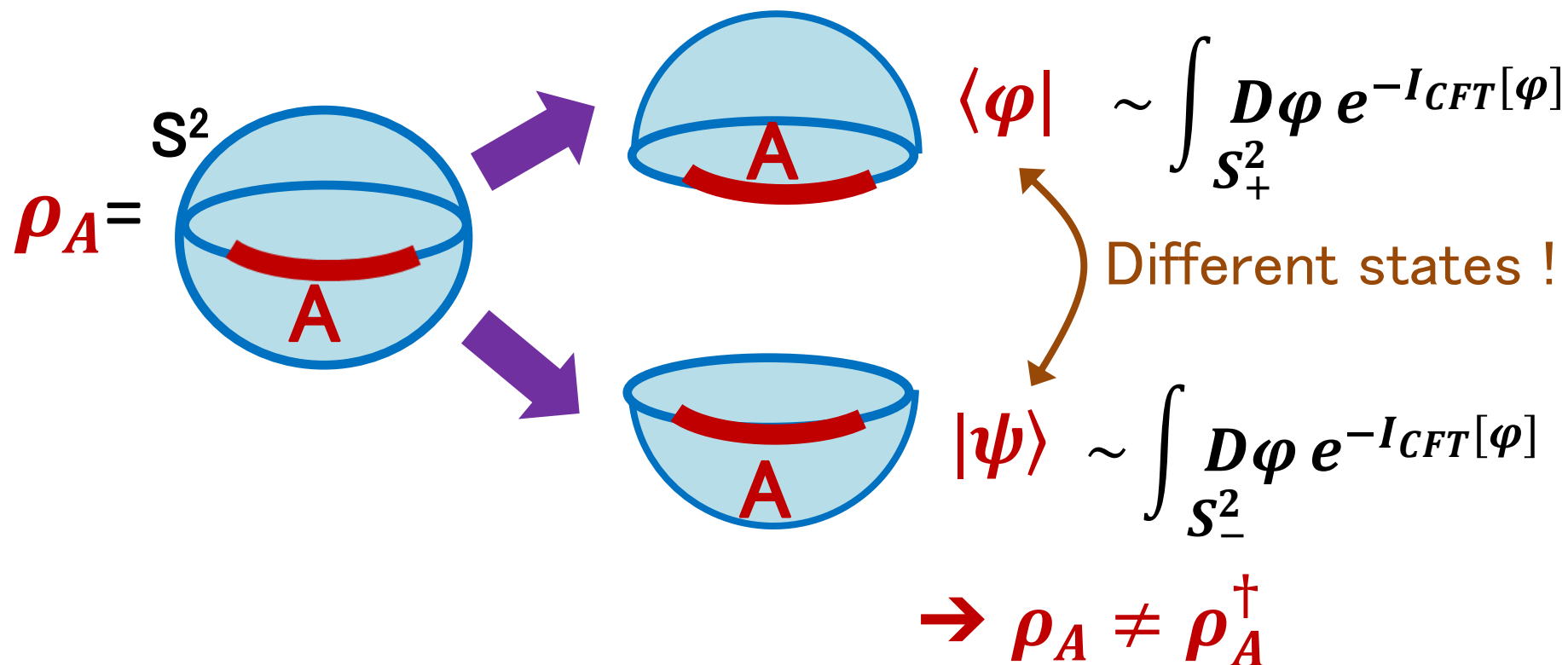
$$S_A = \frac{C_{CFT}}{6} \log \left[\frac{\sin^2 \frac{\theta}{2}}{\tilde{\epsilon}^2} \right], \text{ by setting}$$

Complex valued entropy !
(should not be EE !)

$$C_{CFT} = i C_{dS} \text{ and } \tilde{\epsilon} = i \epsilon = i e^{-t_\infty}.$$

We argue it is more properly considered as pseudo entropy (PE).

This is because the reduced density matrix ρ_A is not Hermitian !



Note: the emergent time coordinate = imaginary part of PE.

④ Probing De Sitter Space from CFT

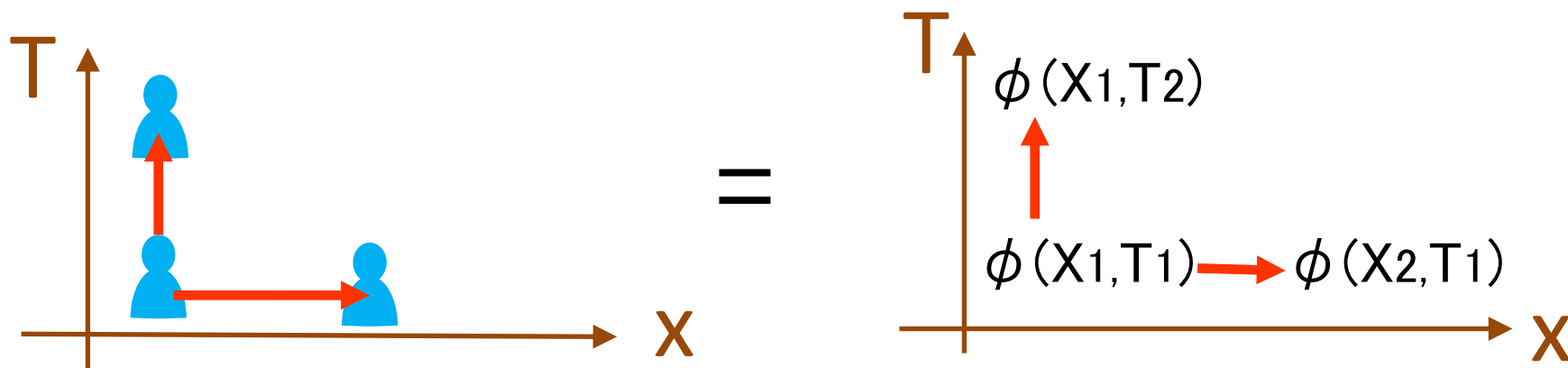
Consider an observer in 2d CFT.

➡ How does the observer feel that he or she lives in AdS or dS ?



To probe the spacetime, we introduce a local excitation.

“How many directions can the observer move ?”

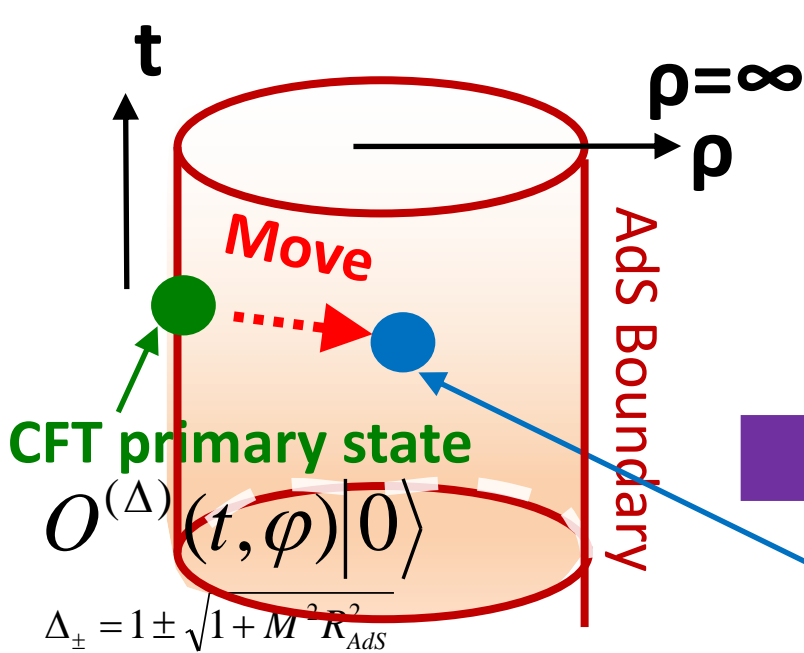


In flat spacetime, $|\Psi_{\text{Flat}}(t, x)\rangle = e^{-iHt} e^{iPx} \hat{\phi}(0)|0\rangle$

Warming up: Probing AdS from CFT

AdS3 metric

$$ds^2 = R_{AdS}^2 (-\text{Cosh}^2 \rho dt^2 + d\rho^2 + \text{Sinh}^2 \rho d\varphi^2)$$



Geometric Symmetry
= $SL(2, R)_L \times SL(2, R)_R$

Conformal
symmetry
of 2d CFT

$$L_0, L_{-1}, L_1 \quad \tilde{L}_0, \tilde{L}_{-1}, \tilde{L}_1$$

$$[L_n, L_m] = (n - m)L_{n+m}$$

$|\Psi_{AdS}^{(\Delta)}(\rho, t, \varphi)\rangle$ CFT dual of
bulk local excitation

The CFT dual of a localized excitation in AdS is given by:

$$|\Psi_{AdS}^{(\Delta)}(\rho, t, \varphi)\rangle = e^{-i(L_0 + \tilde{L}_0)t} e^{i\varphi(L_0 - \tilde{L}_0)} e^{-\frac{\rho}{2}(L_{-1} + \tilde{L}_{-1} - L_1 - \tilde{L}_1)} e^{\frac{\pi i}{2}(L_0 + \tilde{L}_0)} |I^{(\alpha)}\rangle.$$

[Miyaji-Numasawa-Shiba-Watanabe-TT 2015, equivalent to HLKK]

$SL(2, R)$ Ishibashi state

Main problem: Probing dS from CFT

[Doi-Ogawa-Shinmyo
-Suzuki-TT 2024]

First we note the “formal” relation between AdS and dS:

AdS metric

$$ds^2 = R_{AdS}^2 (-\text{Cosh}^2 \rho dt^2 + d\rho^2 + \text{Sinh}^2 \rho d\varphi^2)$$

$$R_{AdS}^2 = -R_{dS}^2$$

$$t = i\tau$$

$$\rho = i\theta$$

$$(L_0)^\dagger = L_0, \quad (L_{\pm 1})^\dagger = L_{\mp 1}$$

dS metric

$$ds^2 = R_{dS}^2 (-\text{Cos}^2 \theta d\tau^2 + d\theta^2 + \text{Sin}^2 \theta d\varphi^2)$$

Unusual conjugation

$$(L_0)^\dagger = -\tilde{L}_0, \quad (L_{\pm 1})^\dagger = \tilde{L}_{\pm 1}$$

$$H^\dagger = -H$$

Hamiltonian is anti-Hermitian ! \rightarrow Non-unitary CFT!

A naïve analytical continuation from AdS leads to the following quantum state for a localized excitation in dS:

$$\left| \tilde{\Psi}_{dS}^{(\Delta_+)}(\tau, \theta, \varphi) \right\rangle = e^{(L_0 + \tilde{L}_0)\tau} e^{i\varphi(L_0 - \tilde{L}_0)} e^{i\frac{\theta}{2}(L_1 + \tilde{L}_1 - L_{-1} - \tilde{L}_{-1})} e^{\frac{\pi i}{2}(L_0 + \tilde{L}_0)} \left| I^{(\alpha)} \right\rangle.$$



Non-unitary evolution
→ emergent time

SL(2,R) Crosscap State

SL(2,R) Ishibashi State

$$(L_n - \tilde{L}_{-n}) \left| I^{(\alpha)} \right\rangle = 0$$

Primary operator in CFT

$$O^{(\Delta_+)}(t, \varphi)$$

$$\Delta_{\pm} = 1 \pm i\sqrt{M^2 R_{dS}^2 - 1}$$

However, this state *leads to the confusing result* (due to the *unusual conjugation*):

$$\left\langle \tilde{\Psi}_{dS}^{(\Delta_+)}(\tau, \theta, \varphi) \left| \tilde{\Psi}_{dS}^{(\Delta_+)}(\tau', \theta', \varphi') \right. \right\rangle = 0.$$

Resolution

The correct answer is found by requiring the CPT invariant state:

$$\begin{aligned} \left| \Psi_E^{(\Delta_+)}(\tau, \theta, \varphi) \right\rangle &= \frac{1}{\sqrt{i}} \left| \tilde{\Psi}_{dS}^{(\Delta_+)}(\tau, \theta, \varphi) \right\rangle + CPT \cdot \left(\frac{1}{\sqrt{i}} \left| \tilde{\Psi}_{dS}^{(\Delta_+)}(\tau, \theta, \varphi) \right\rangle \right) \\ &= \frac{1}{\sqrt{i}} \left| \tilde{\Psi}_{dS}^{(\Delta_+)}(\tau, \theta, \varphi) \right\rangle + \sqrt{i} \left| \tilde{\Psi}_{dS}^{(\Delta_-)}(\tau, \pi - \theta, \varphi + \pi) \right\rangle. \end{aligned}$$

Antipodal map

[cf. Global sym. in gravity: Harlow–Ooguri 2018, Gauging CPT in QG: Harlow–Numasawa 2023]

Indeed, this reproduces the correct dS Green function at Euclidean vacuum

$$\left\langle \Psi_E^{(\Delta_+)}(x) \left| \Psi_E^{(\Delta_+)}(x') \right. \right\rangle = G_{dS}^E(x, x') = \frac{\sinh \mu(\pi - D_{dS}(x, x'))}{4\pi \sinh \pi\mu \cdot \sin D_{dS}(x, x')}.$$

Summary

- In dS/CFT, the Hamiltonian is anti Hermitian.
→ Emergence of time
- In dS/CFT, we need to gauge the CPT symmetry.

⑤ Conclusions

- ◆ Holography has successfully predicted that $\Lambda < 0$ (AdS) spacetimes emerge from quantum entanglement in CFTs.
- ◆ Holography for $\Lambda > 0$ (dS) spacetimes is still in the development stage due to many exotic properties of dual CFTs.
- ◆ Complex valued entanglement entropy in dS/CFT can be interpreted as pseudo entropy. ➔ Emergence of time
- ◆ An analysis of CFT dual of bulk local excitation again reveals non-Hermitian nature of the dual CFT.

Future problems

- ▶ CFT interpretations of Einstein equation in dS
- ▶ QI understandings of pseudo entropy
- ▶ CFT Models of creation of the Universe

:

Happy birthday, Murayama-san !



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