

Continuing Hitoshi's search for axions in the universe

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Kavli IPMU

Hitoshifest

18/Dec/2024

Congratulations, Hitoshi!



Thank you Hitoshi!



KAVLI
IPMU

I am very grateful for you to have founded this amazing institute, the perfect environment for researchers!

Congratulations, Hitoshi!



Hitoshi's papers one of the reasons I work on dark matter and why I came to



What is *dark matter*?

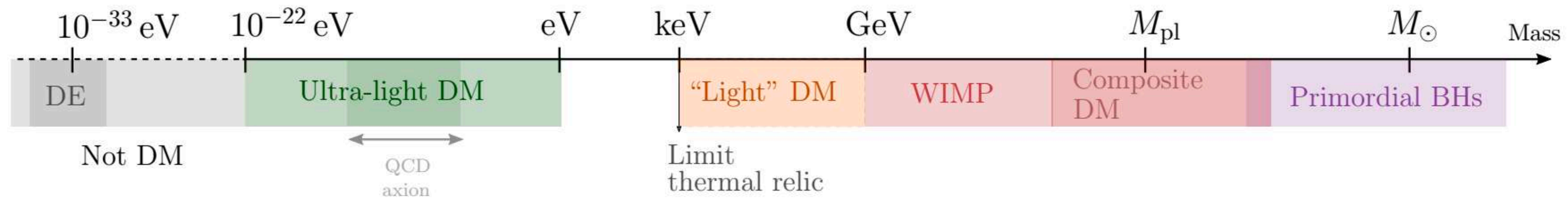
- What is the nature of DM?

State of the “art”



Mass scale of DM

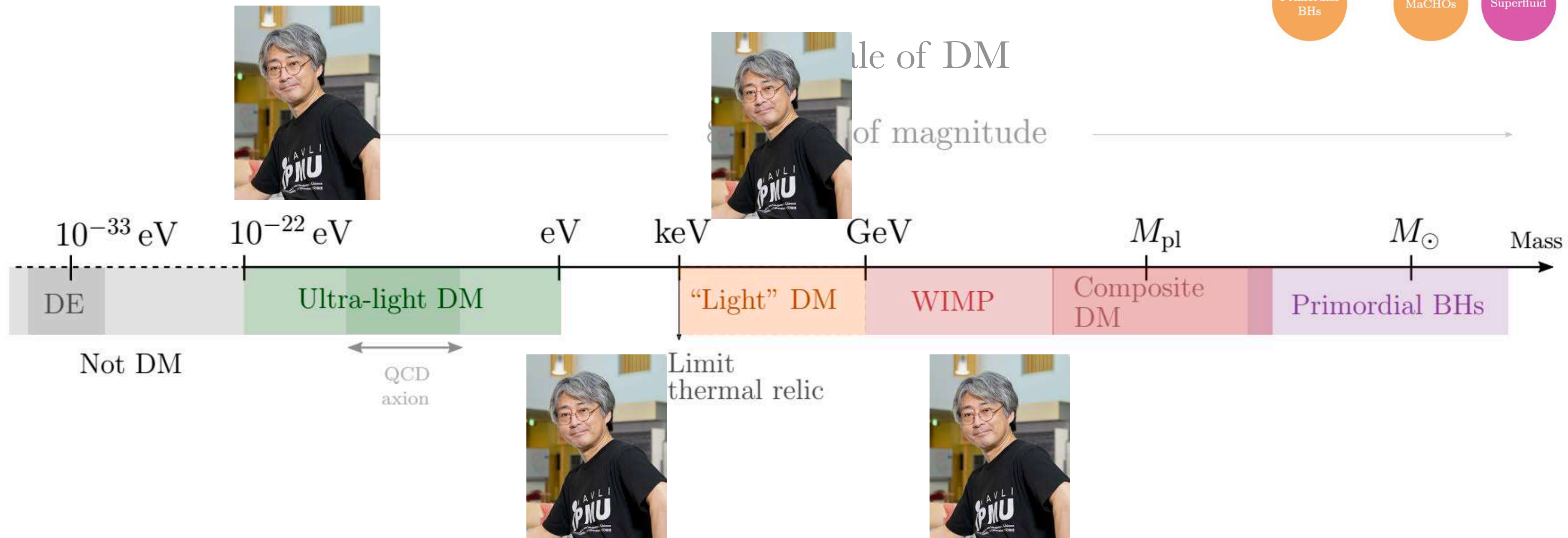
80 orders of magnitude



What is *dark matter*?

- What is the nature of DM?

State of the “art”



Dark matter

@ **KAVLI**
IPMU



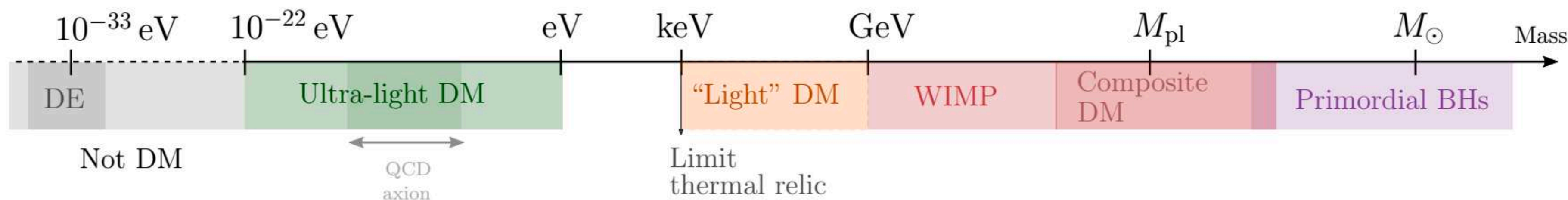
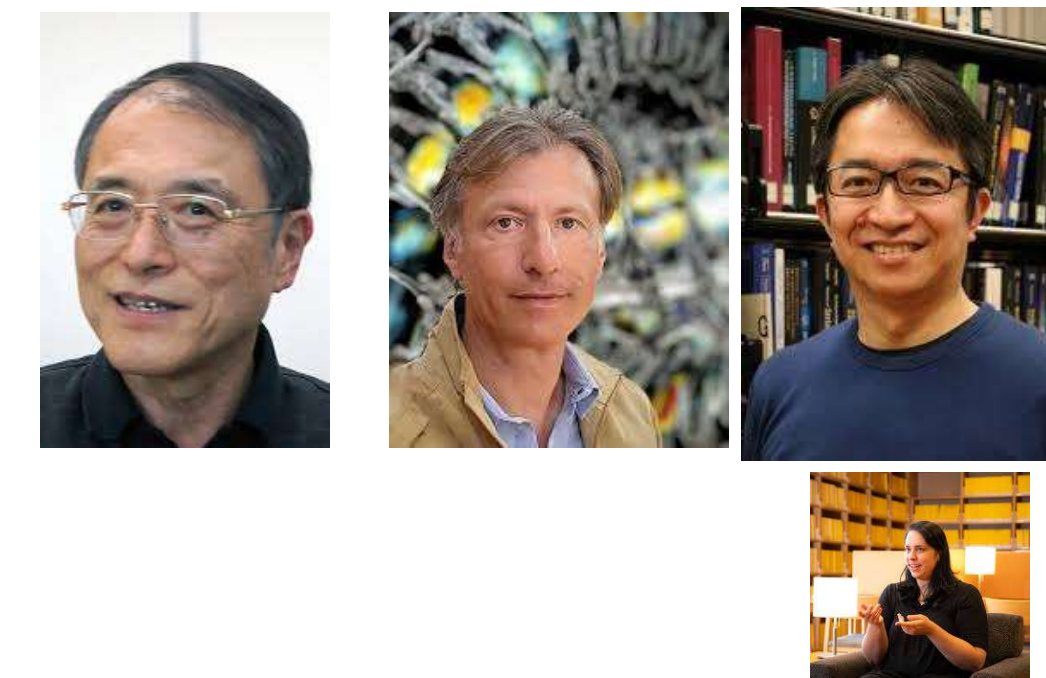
The search for *dark matter* @

DM science is very strong at IPMU.
We work on this entire range of DM candidates...
one of **Hitoshi legacy!**

Heavy DM



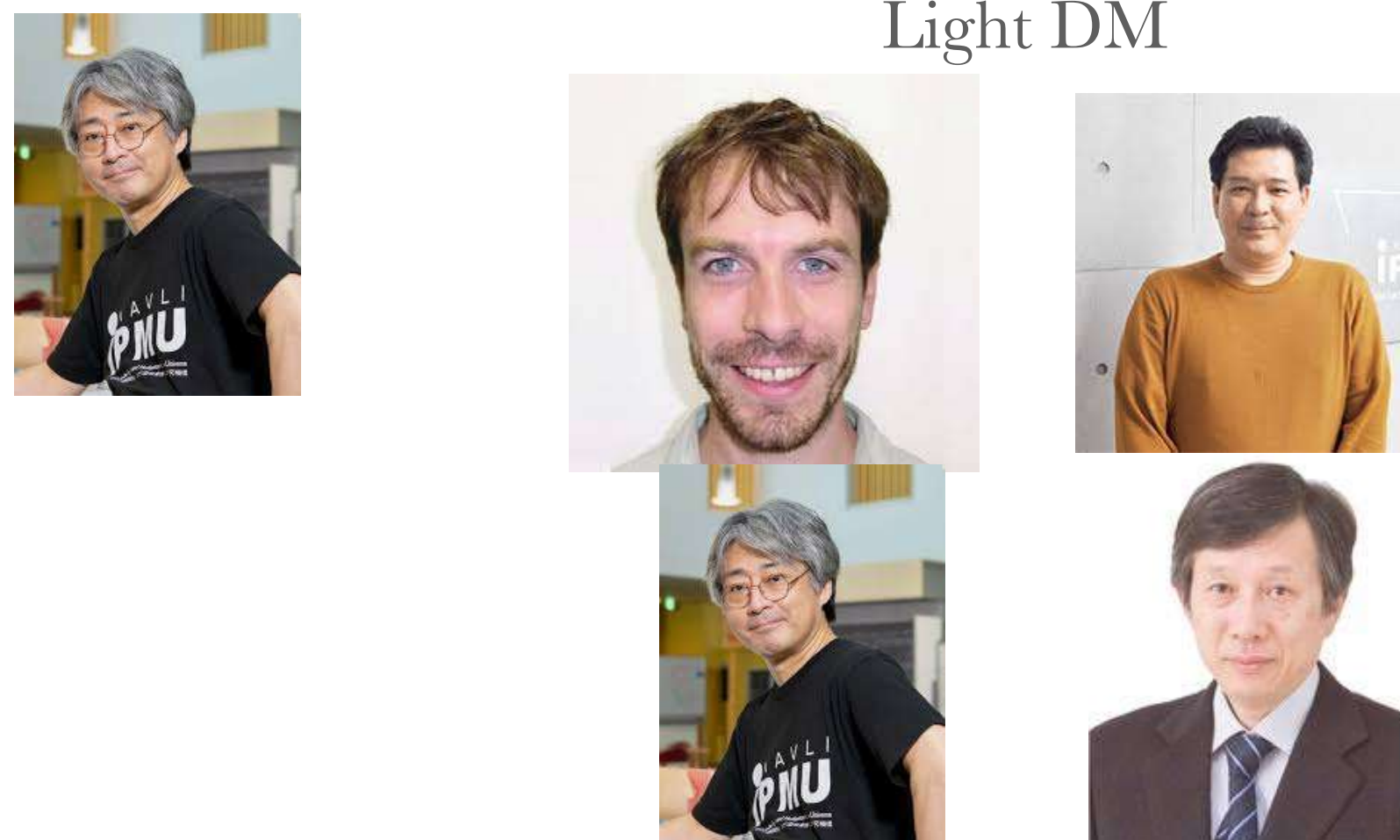
Primordial BHs



Ultra-light DM



Light DM



COSI

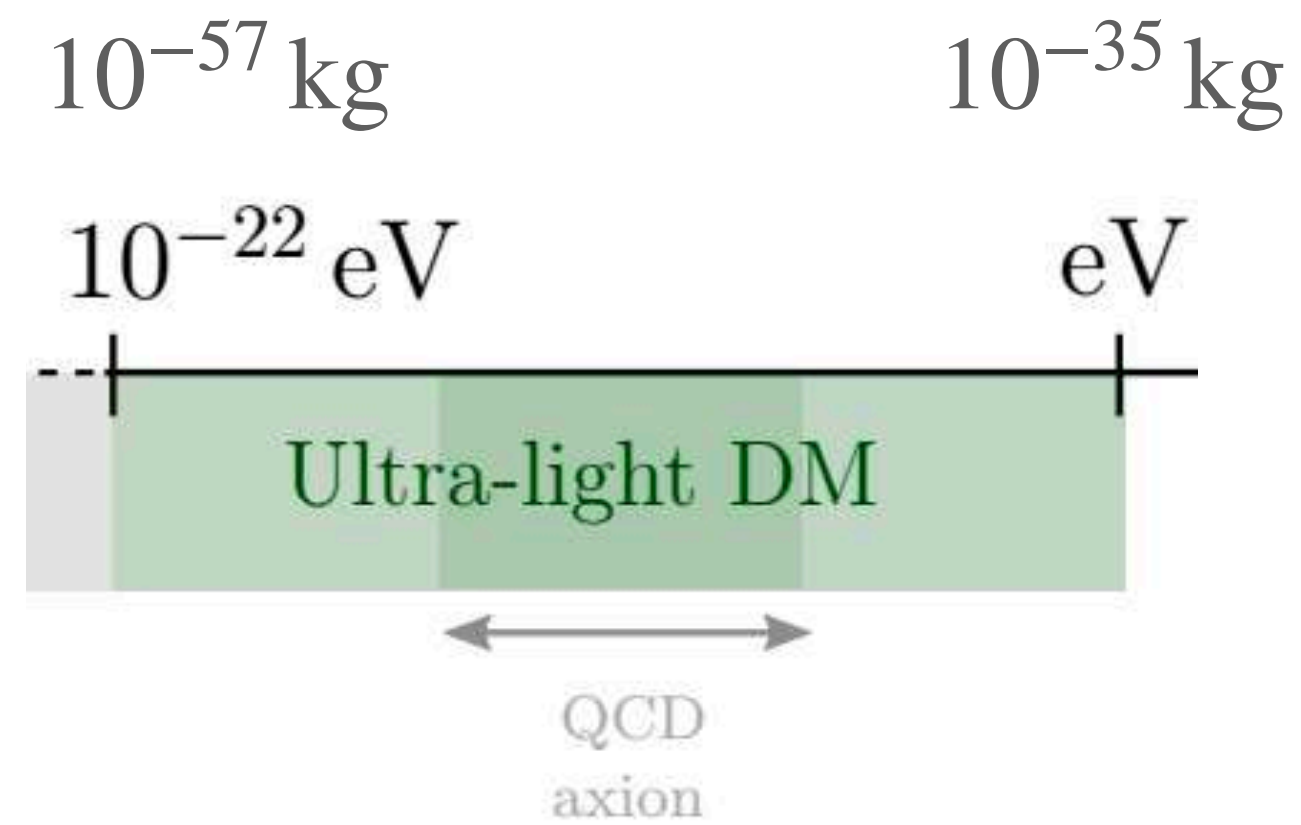
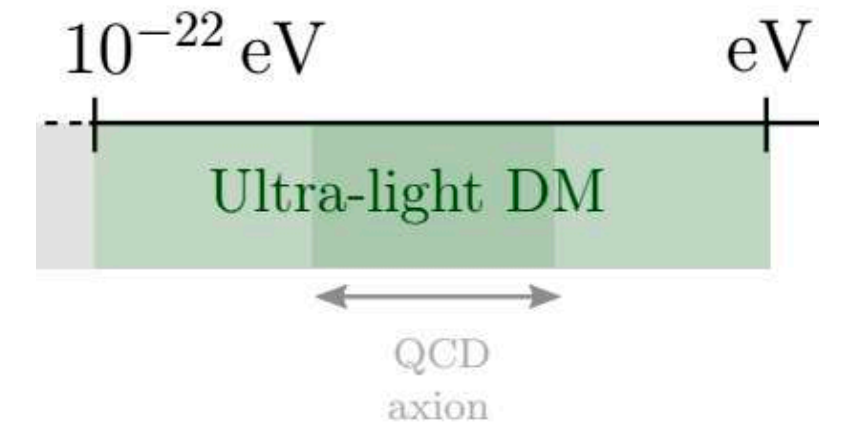
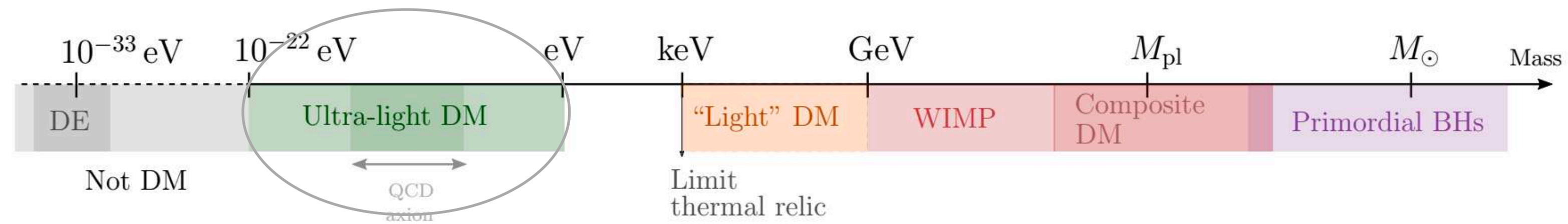
DM at high z



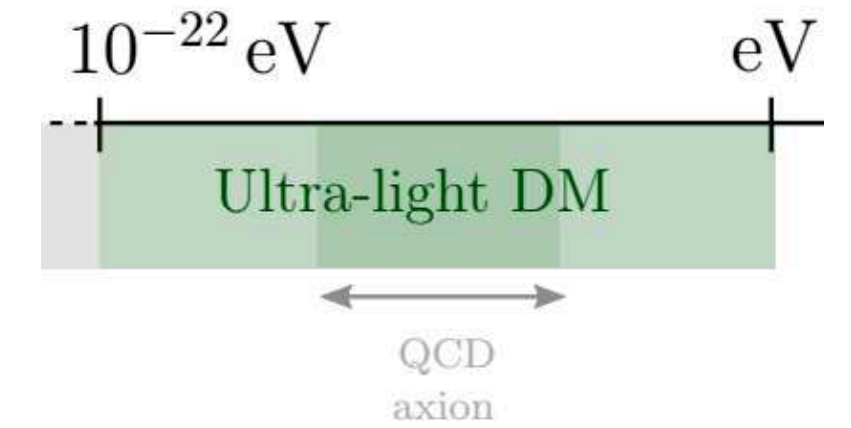
Grav. probe DM



Axion/ALP Dark Matter



Axion/ALP Dark Matter



Axion detection via superfluid ^3He ferromagnetic phase and quantum measurement techniques

So Chigusa (LBL, Berkeley and UC, Berkeley), Dan Kondo (Tokyo U., IPMU), Hitoshi Murayama (LBL, Berkeley and UC, Berkeley and Tokyo U., IPMU)

Search for a Dark-Matter-Induced Cosmic Axion Background with ADMX

ADMX Collaboration · T. Nitta (Washington U., Seattle) [Show All\(47\)](#)

Mar 10, 2023

7 pages

Cosmic axion background

Jeff A. Dror (UC, Santa Cruz and UC, Santa Cruz, Inst. Part. Phys. and UC, Berkeley and LBNL, Berkeley), Hitoshi Murayama (UC, Berkeley and LBNL, Berkeley and Tokyo U., IPMU), Nicholas L. Rodd (UC, Berkeley and LBNL, Berkeley)

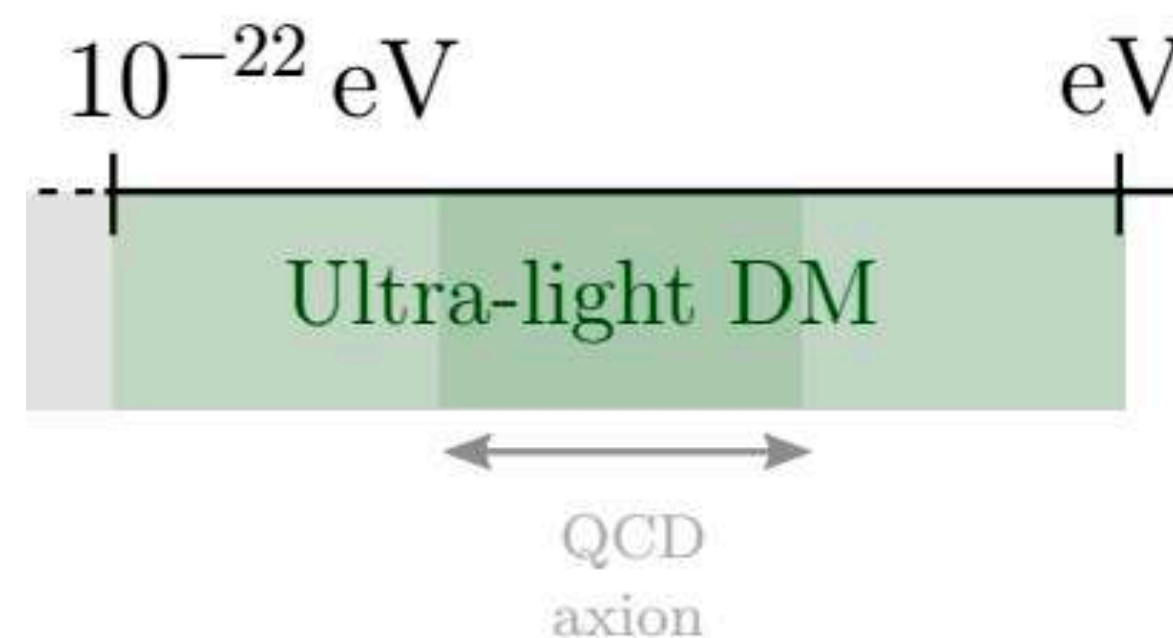
Jan 22, 2021

Axion detection via superfluid ^3He ferromagnetic phase and quantum measurement techniques

So Chigusa (LBL, Berkeley and UC, Berkeley), Dan Kondo (Tokyo U., IPMU), Hitoshi Murayama (LBL, Berkeley and UC, Berkeley and Tokyo U., IPMU), Risshin Okabe (Tokyo U., IPMU), Hiroyuki

Sep 17, 2023

41 pages



Axion strings are superconducting

Hajime Fukuda (LBL, Berkeley and UC, Berkeley), Aneesh V. Manohar (UC, San Diego), Hitoshi Murayama (LBL, Berkeley and UC, Berkeley and Tokyo U., IPMU), Ofri Telem (LBL, Berkeley and UC, Berkeley)

Oct 6, 2020

A keV String Axion from High Scale Supersymmetry

Brian Henning (UC, Berkeley and LBL, Berkeley), John Kehayias (Tokyo U., IPMU), Hitoshi Murayama (UC, Berkeley and LBL, Berkeley and Tokyo U., IPMU), David Pinner (UC, Berkeley and LBL, Berkeley), Tsutomu T. Yanagida (Tokyo U., IPMU)

Aug 1, 2014

11 pages

Quark mass uncertainties revive Kim-Shifman-Vainshtein-Zakharov axion dark matter

Matthew R. Buckley (UC, Berkeley and LBL, Berkeley), Hitoshi Murayama (UC, Berkeley and LBL, Berkeley)

May 2007



How to detect dark matter



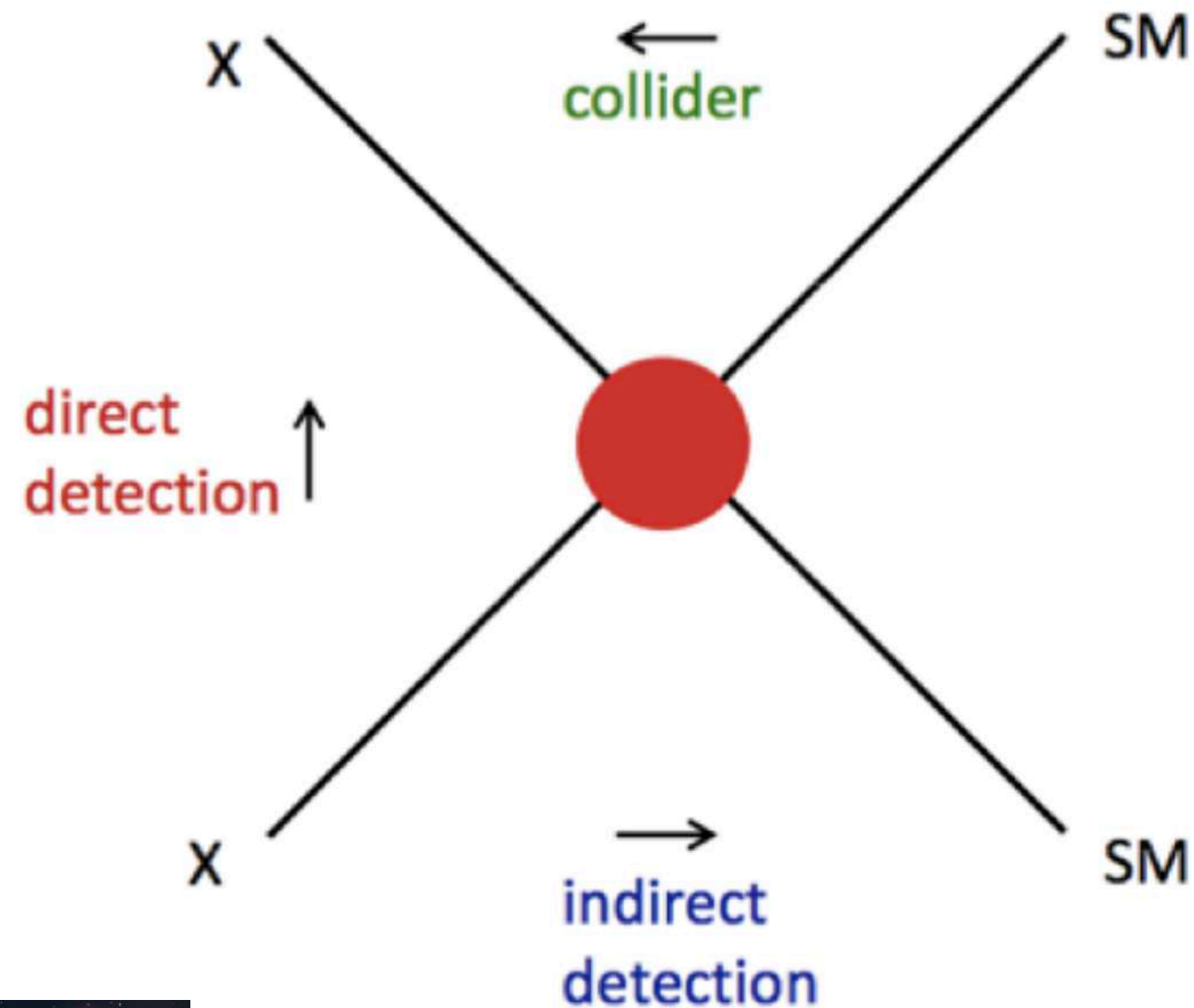
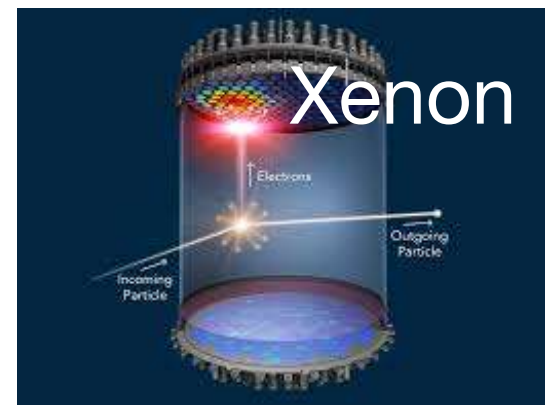
All of these experiments have involvement or leadership

Interaction with the SM

Talks by Lindley Winslow, Kai Martens and Risshin Okabe



Muon colliders



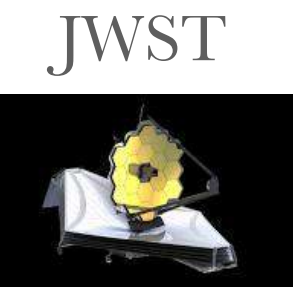
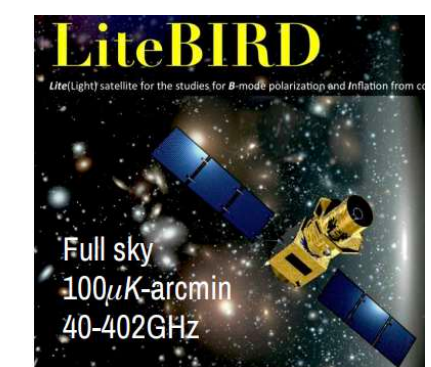
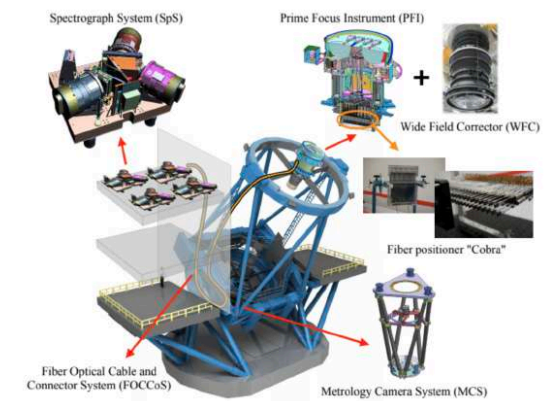
+

Gravitationally Cosmological and astrophysical searches

Hyper Suprime-Cam (HSC)



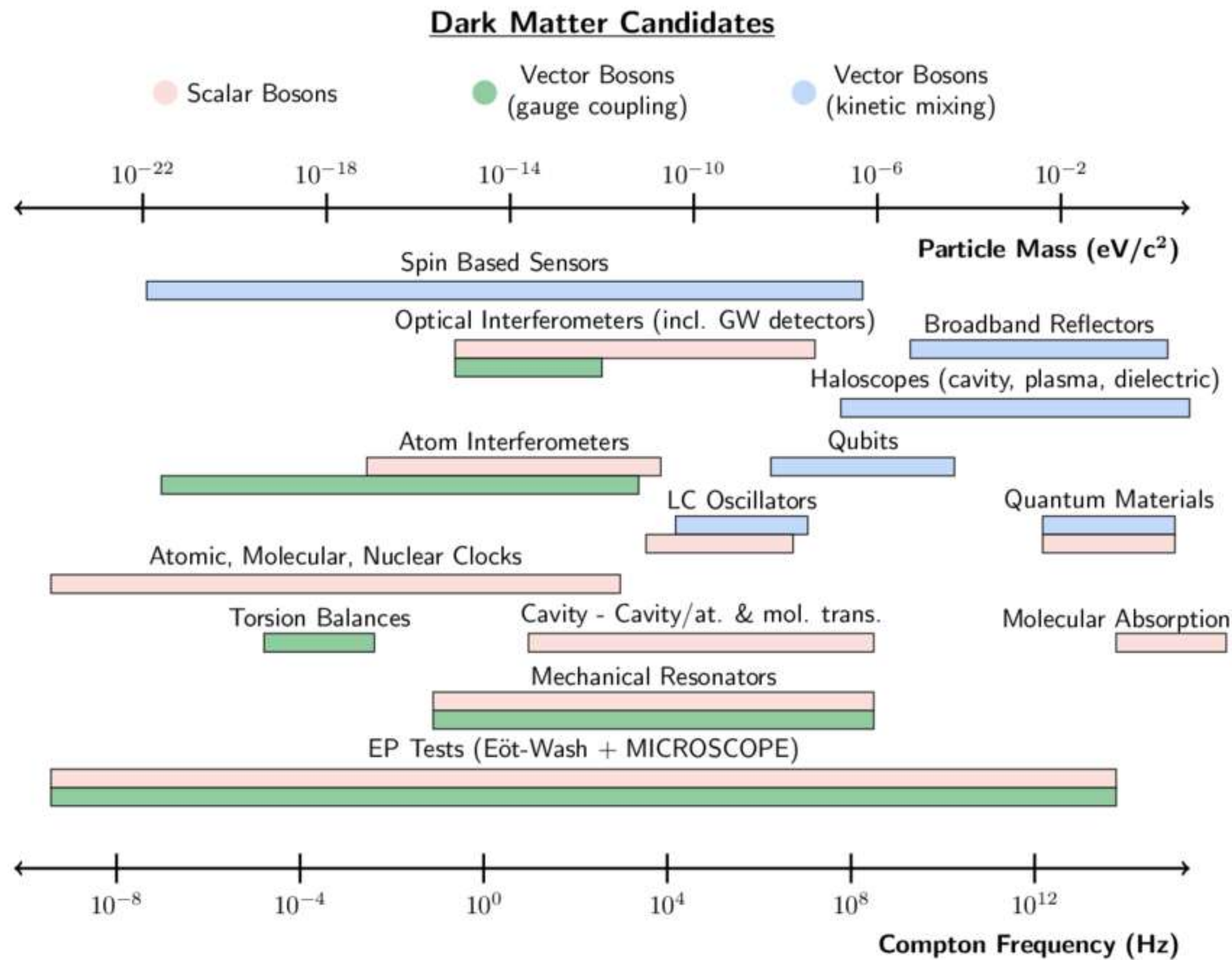
Prime Focus Spectrograph (PFS)



* not a complete list

How to detect *axion/ALP* dark matter

Interaction with the SM



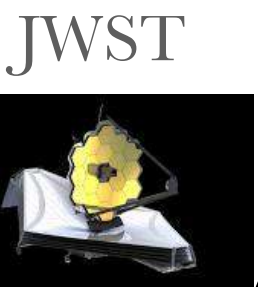
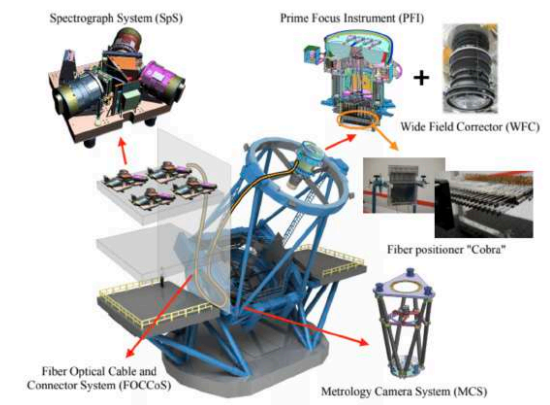
+

Gravitationally Cosmological and astrophysical searches

Hyper Suprime-Cam (HSC)



Prime Focus Spectrograph (PFS)

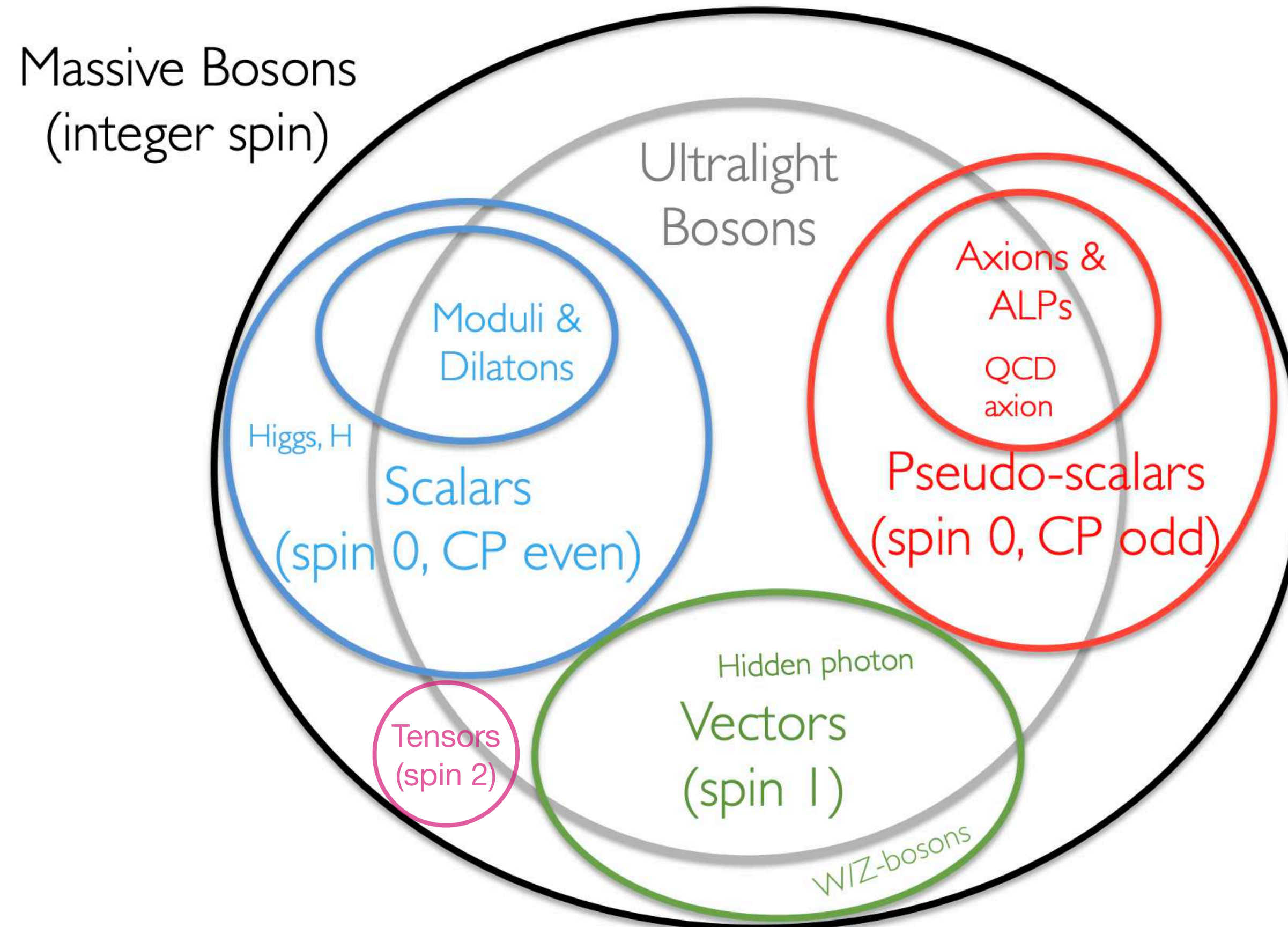
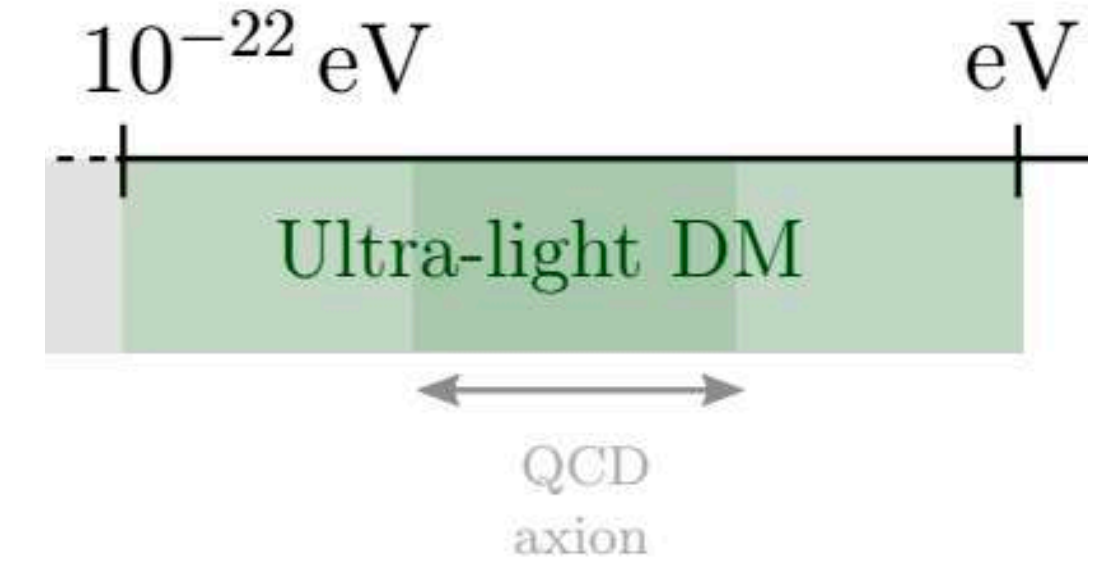


Talks by *Lindley Winslow* and
Risshin Okabe

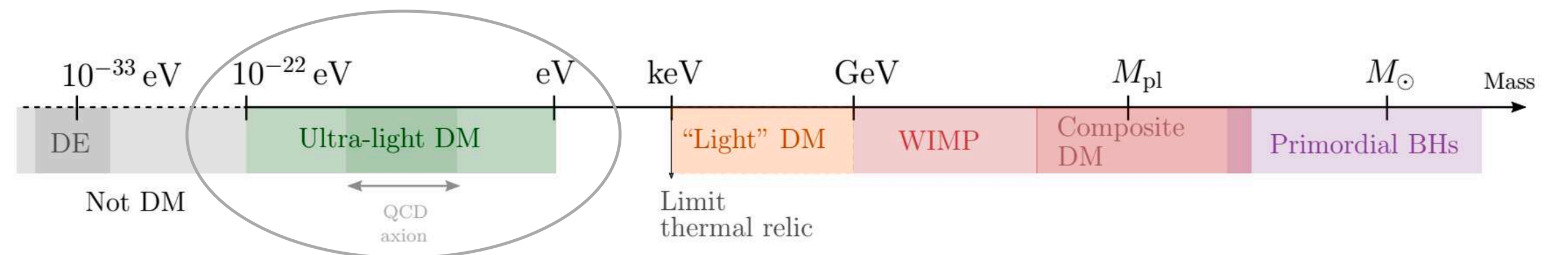
* not a complete list

Ultra-light *Dark Matter*

Since we are interested in the gravitational signatures...



Ultra-light dark matter



“Ultra-light dark matter”, **E.Ferreira**, 2020. The Astronomy and Astrophysics Review.



Andrew Eberhardt



Qiuyue Liang



Ippei Obata



Shun'ichi Horigome



Dongdong Zhang



Dan Kondo

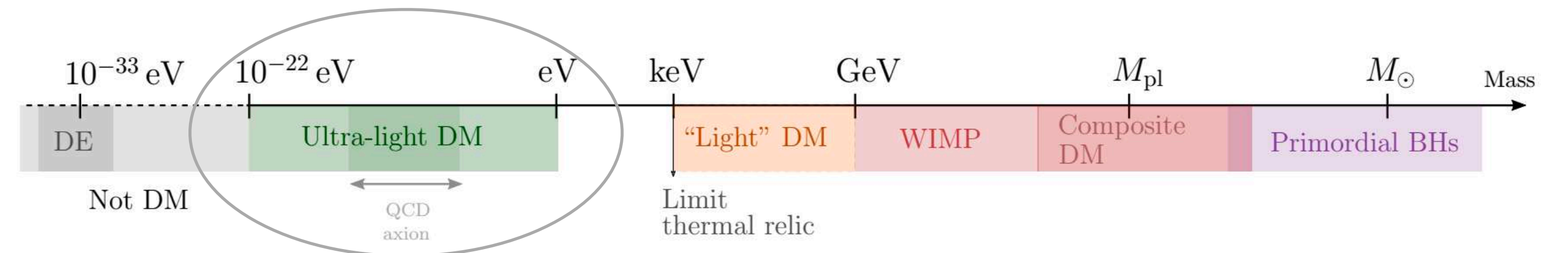


Margot Imbach



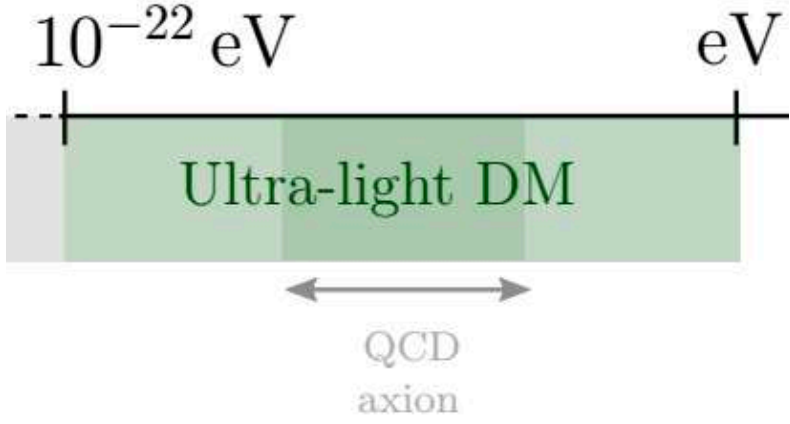
Fernanda Matos

Ultra-light dark matter



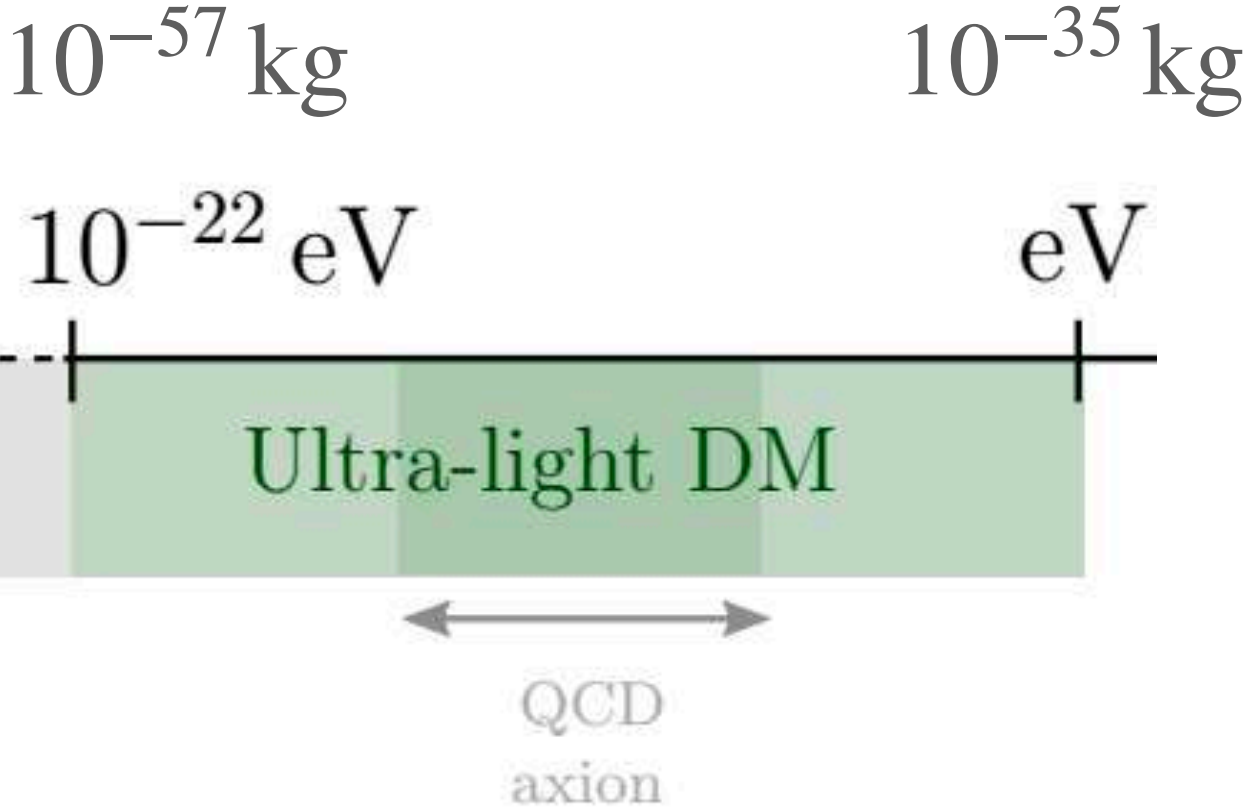
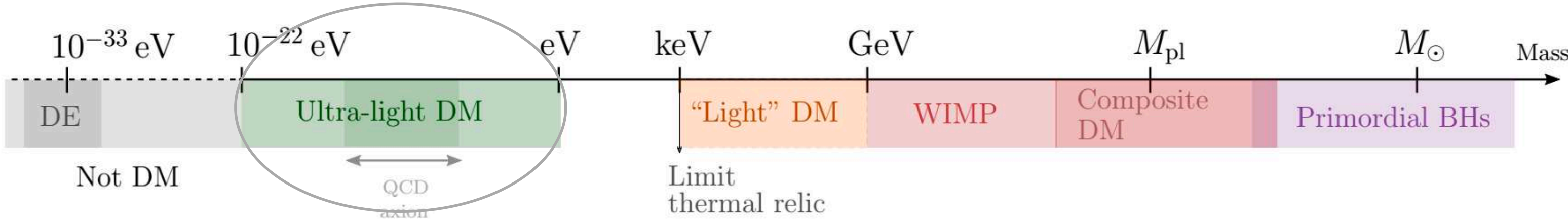
"Ultra-light dark matter", E.Ferreira, 2020. The Astronomy and Astrophysics Review.

Ultra-light *Dark Matter*



Ultra-light candidate, cold \longrightarrow Large $\lambda_{\text{dB}} \sim 1/mv$

Lightest possible candidate for DM



\longrightarrow

Bosons

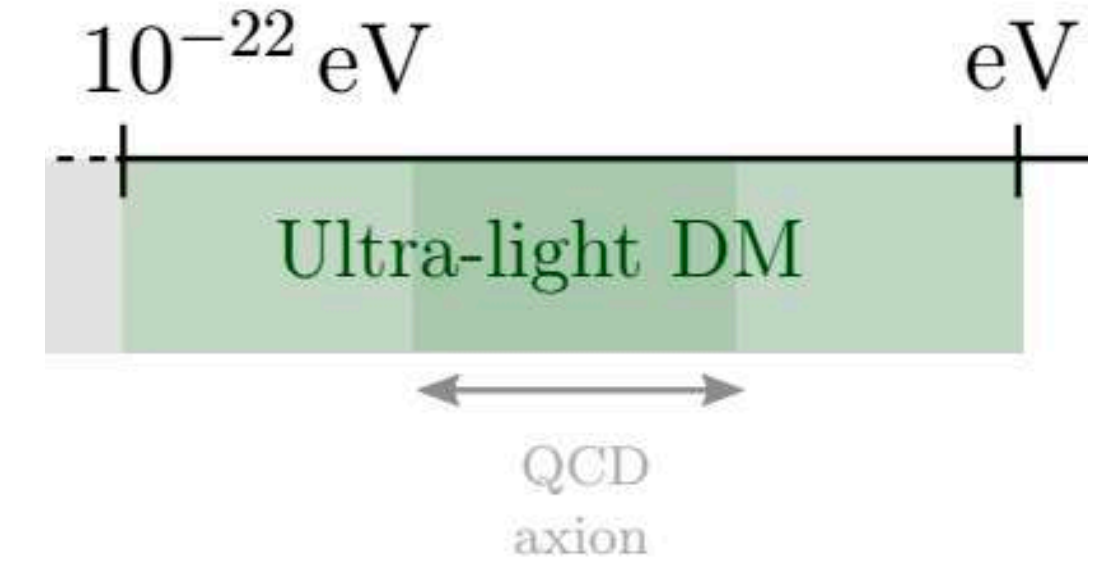
Non-thermally produced

Motivation: *particle physics*

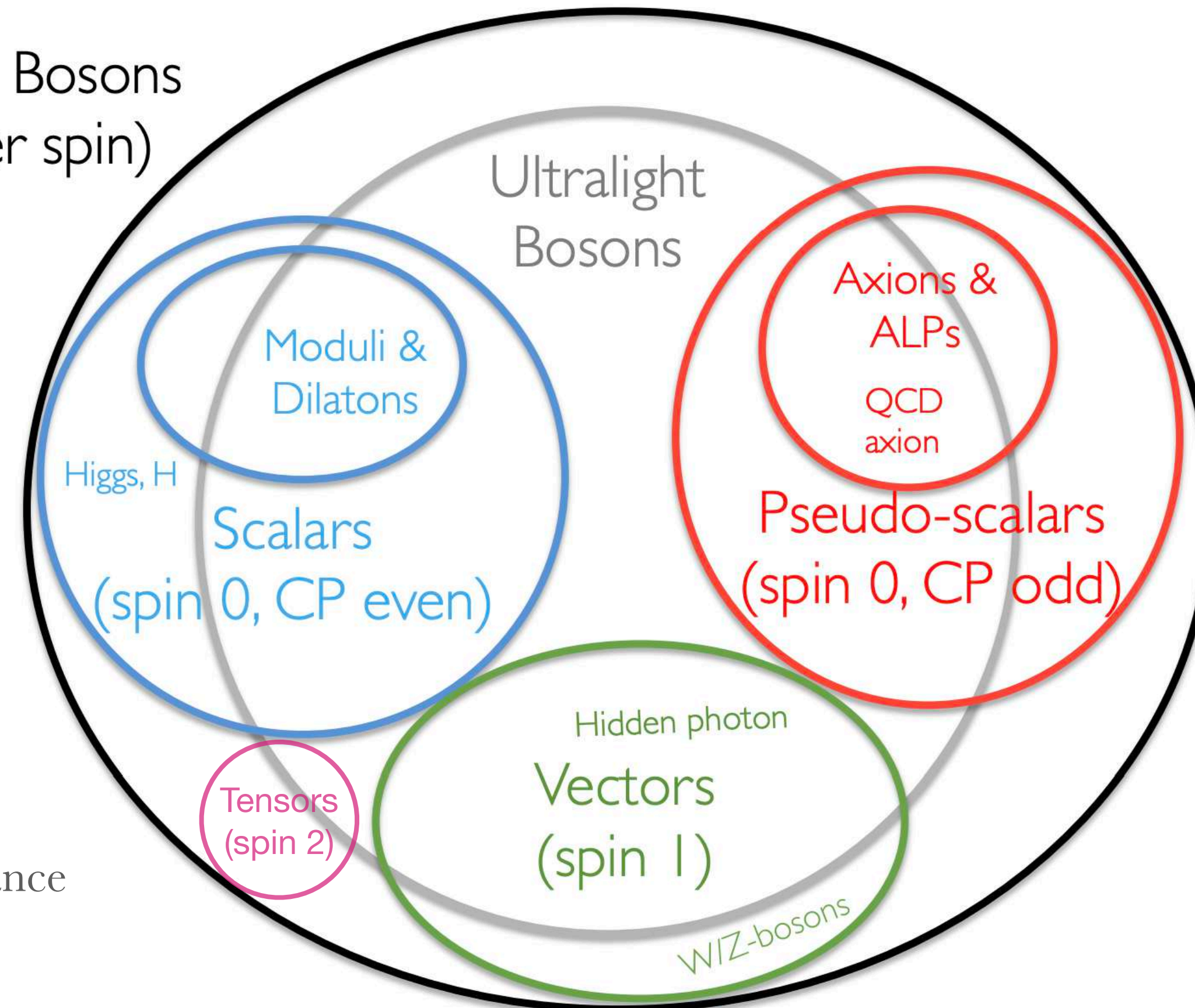
ULDM candidates

Natural candidate for a light scalar field is a pseudo-Nambu Goldstone boson

Many extensions of the Standard Model predict additional massive bosons



Massive Bosons
(integer spin)



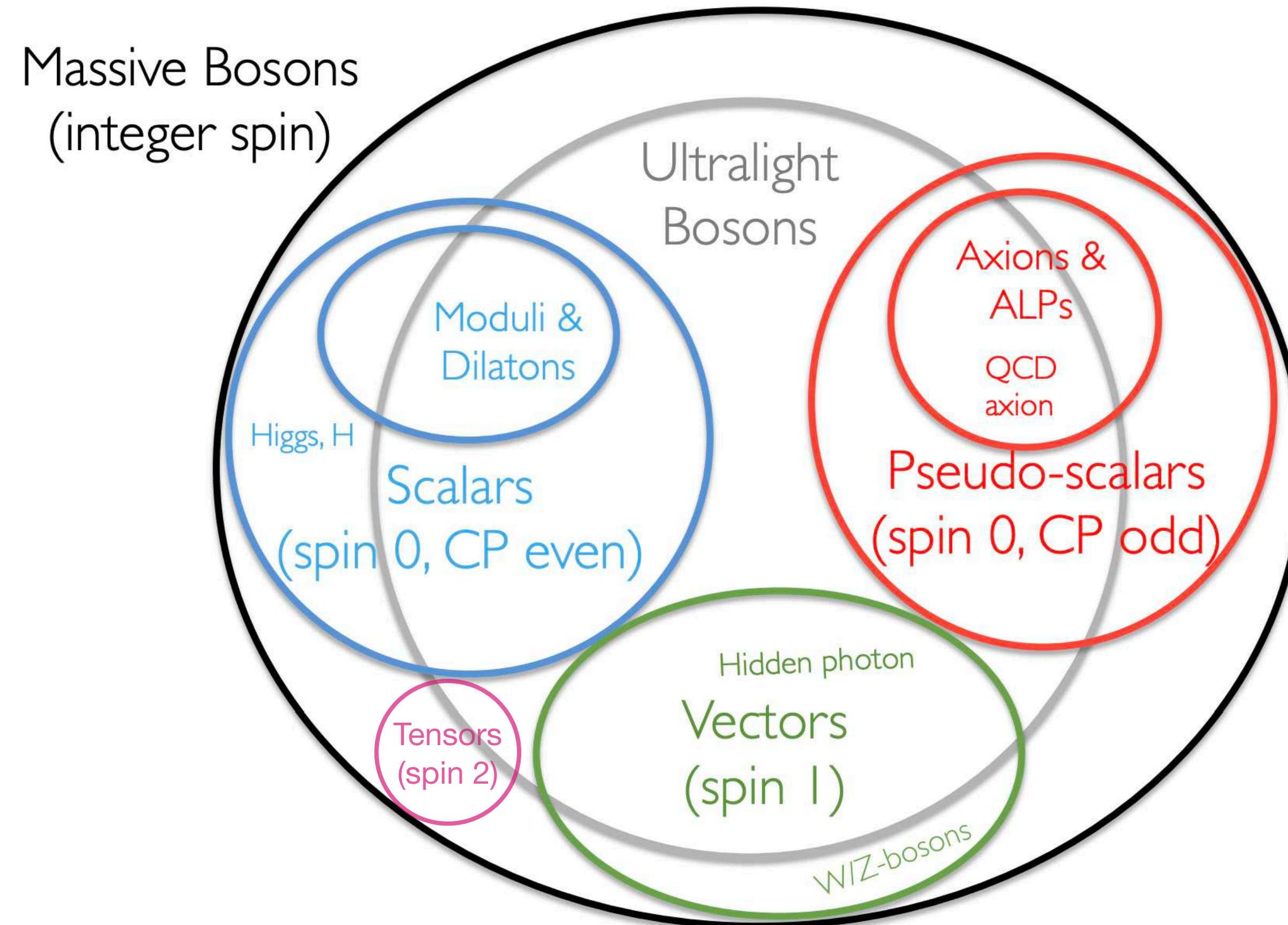
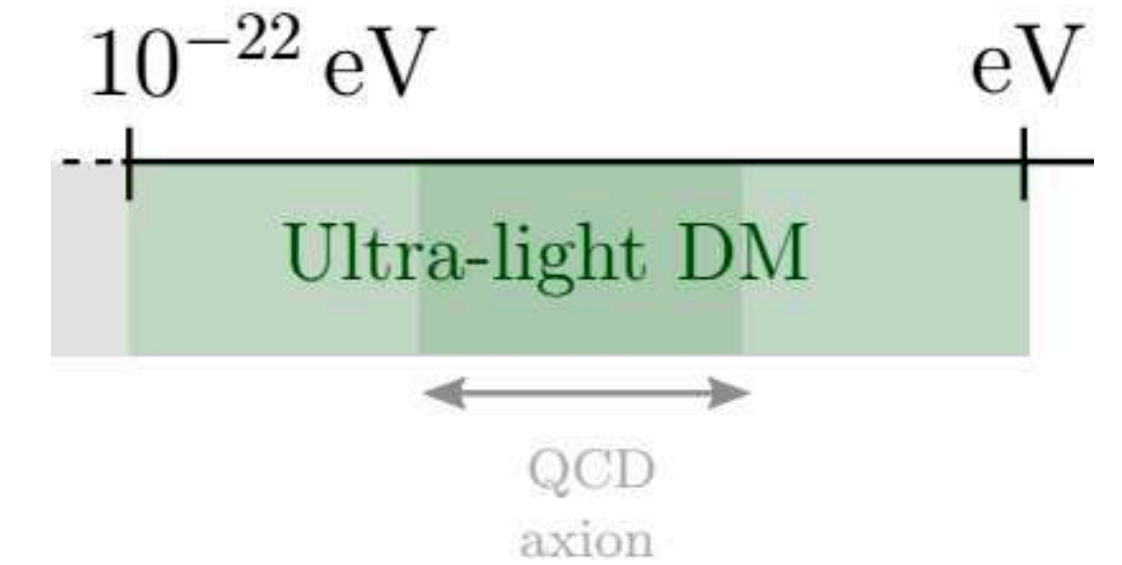
- Formation mechanism: needs to have a relic abundance that gives the correct DM abundance

A lot of research in this at IPMU! (e.g. Kaloian, Ippei)

Motivation: *particle physics*

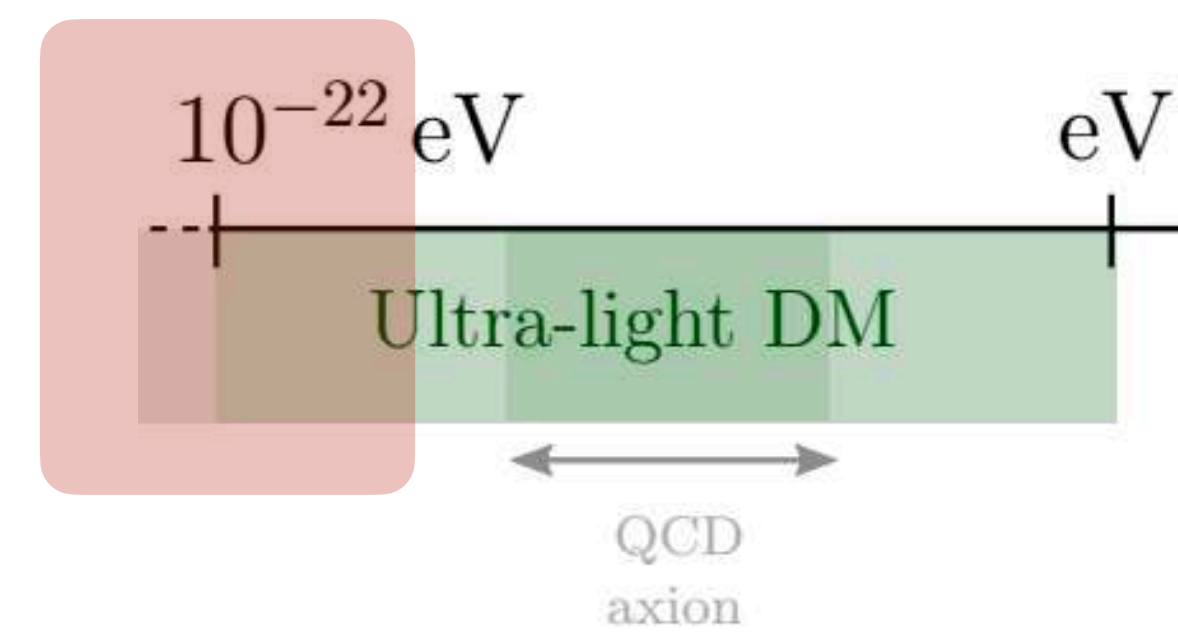
ULDM candidates

Many extensions of the Standard Model predict additional massive bosons

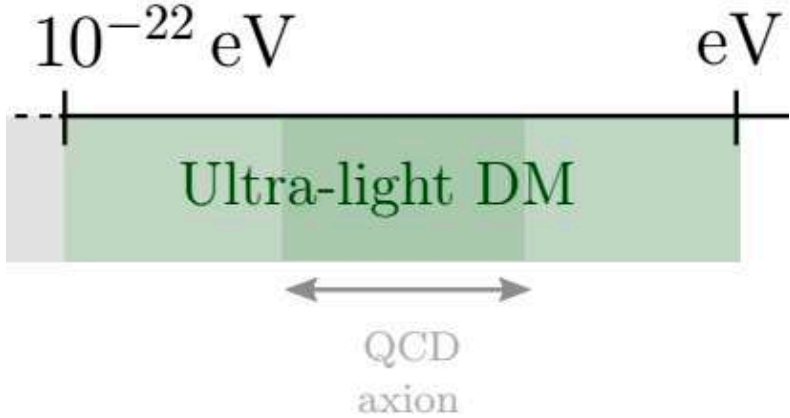


For most of this talk:
Gravitational signatures!

Cosmological signatures



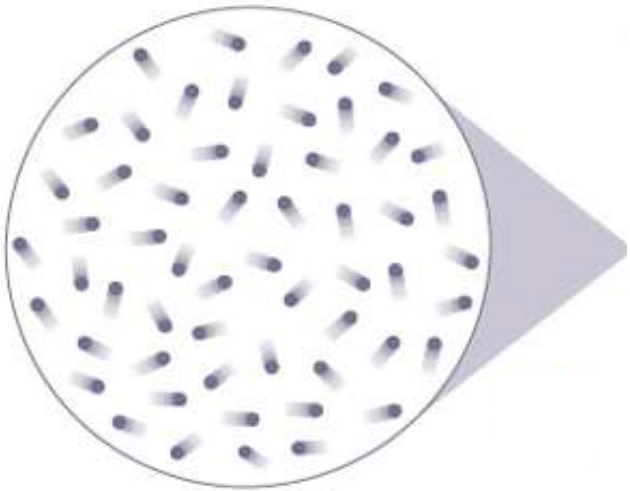
Ultra-light *Dark Matter*



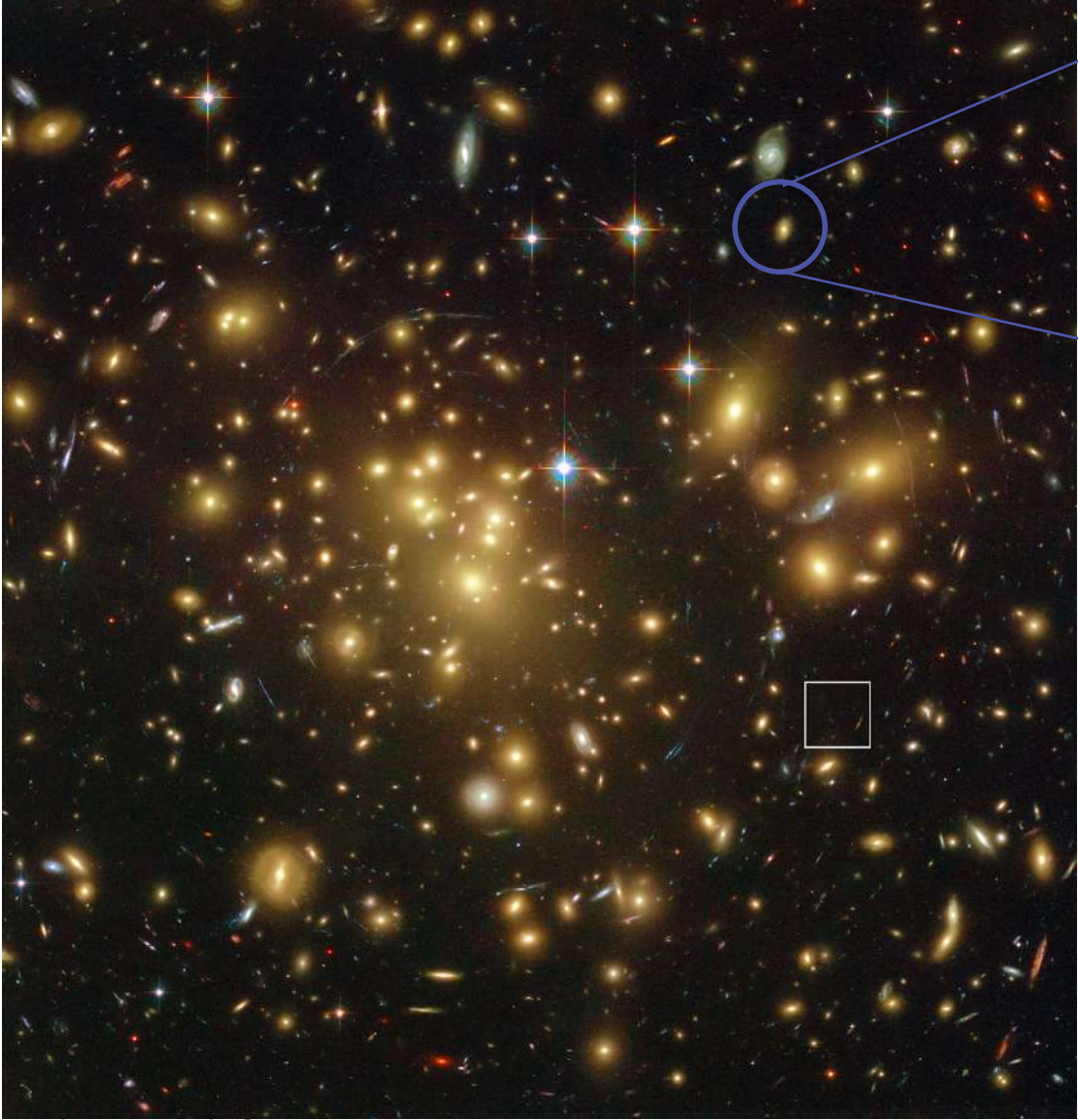
Ultra-light candidate \longrightarrow Large $\lambda_{dB} \sim 1/mv$

Lightest possible candidate for DM

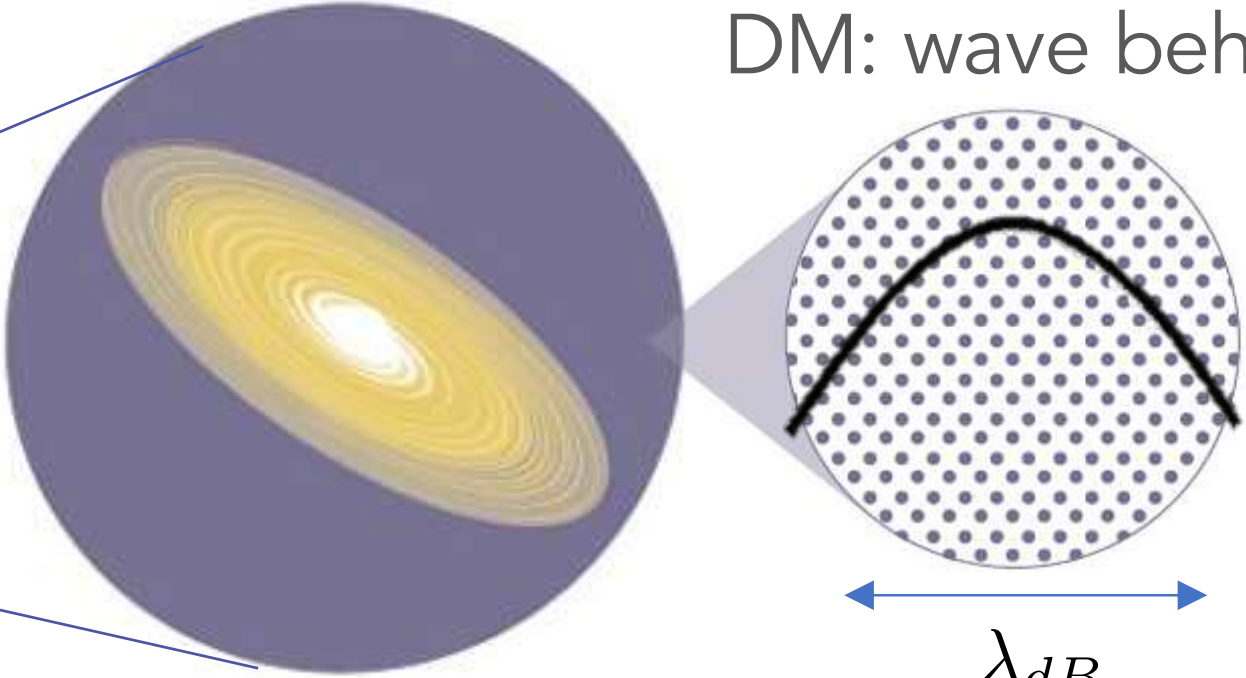
Large scales:
DM behaves like standard particle DM (**CDM**).



DM: particles
 $d \gg \lambda_{dB}$



Adapted from Quanta



Galaxy halo

DM: wave behaviour

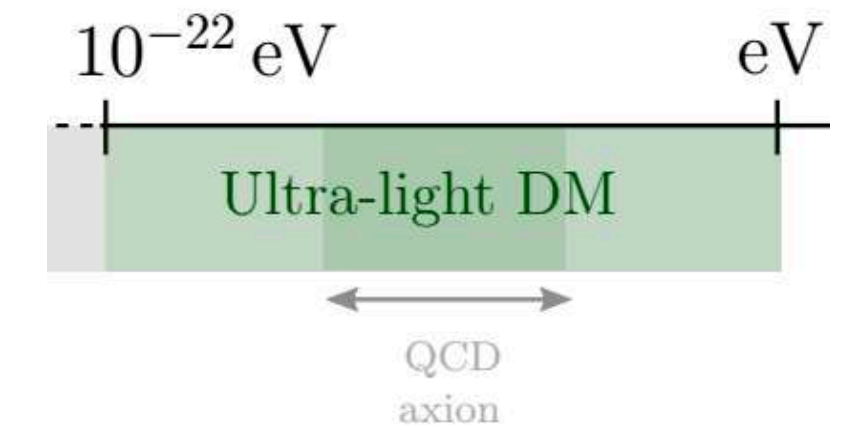
λ_{dB}
 $d \ll \lambda_{dB}$

Small scales:
DM behaves like a **wave**

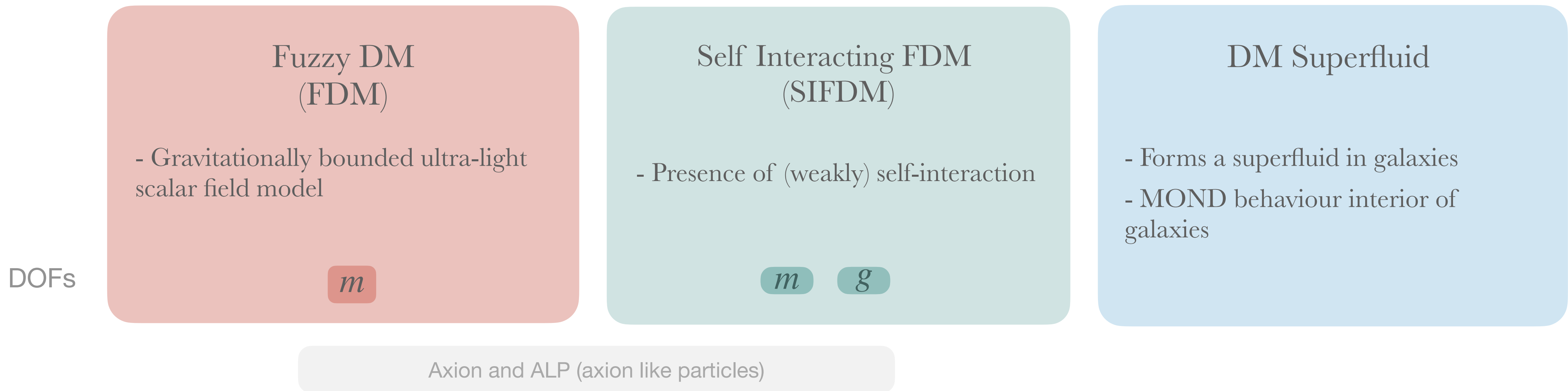
$$10^{-25} \text{ eV} \lesssim m \lesssim \text{eV}$$

$$\lambda_{dB}^{ULDM} \sim \text{pc} - \text{kpc}$$

Ultra-light Dark Matter -classes



3 classes:



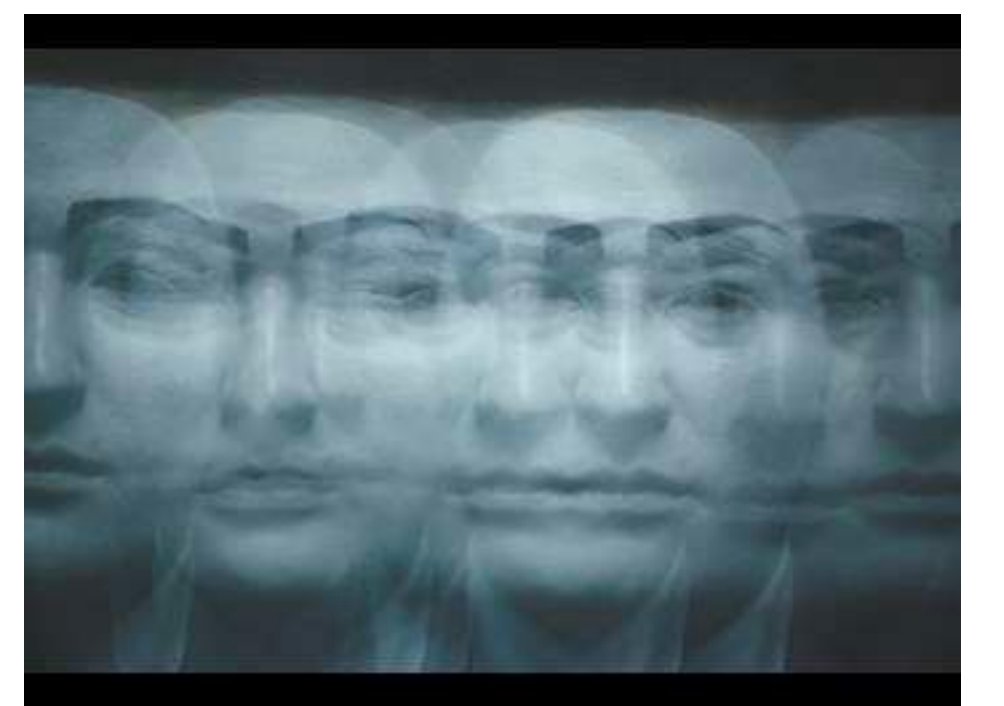
$$i\dot{\psi} = \left(-\frac{1}{2m} \nabla^2 + \frac{g}{8m^2} |\psi|^2 - m\Phi \right) \psi$$

$$\mathcal{L} = P(X)$$

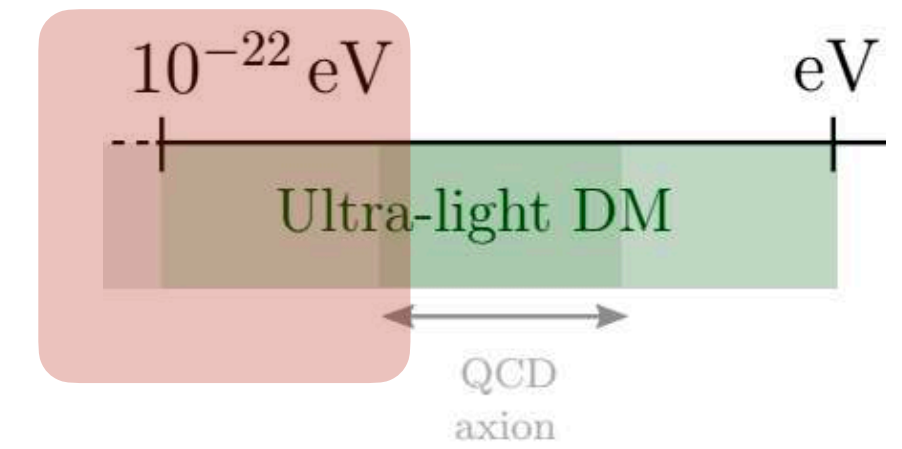
→ Connection with condensed matter and particle physics!

“Ultra-light dark matter”, **E.Ferreira**, 2020. *The Astronomy and Astrophysics Review*.

Fuzzy dark matter



Fuzzy Dark Matter



Fuzzy DM (FDM)

- Gravitationally bounded ultra-light scalar field model

m

Wave DM Ultra-light axions

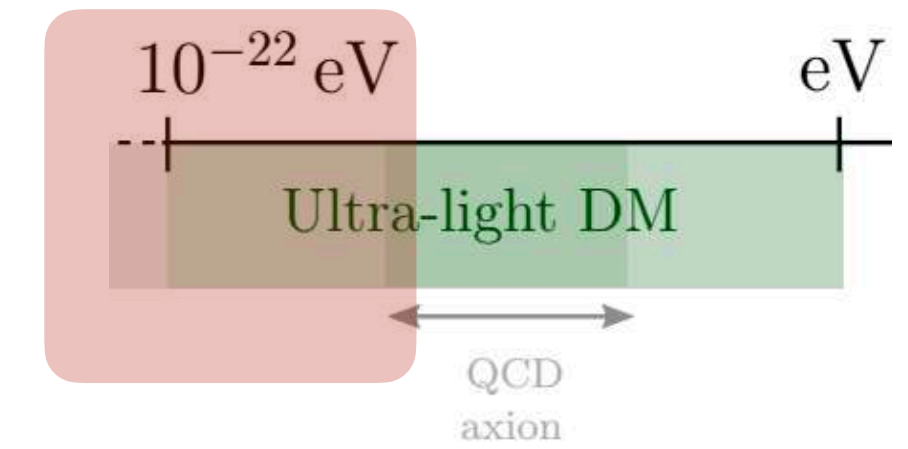
Focus *more* on spin 0 particles here!

$$10^{-22} \text{ eV} < m < 10^{-18} \text{ eV}$$

Hu W, Barkana R, Gruzinov A (2000 a,b)

(Reviews: *EF (2021), J. Niemeyer (2019), L. Hui (2021)*)

Structure formation - *non-relativistic regime*



Evolution on small scales: take non-relativistic regime of the theory, relevant for structure formation.

Schrödinger-Poisson system : describe the FDM and the SIFDM

$$\left\{ \begin{array}{l} i\dot{\psi} = \left(-\frac{1}{2m} \nabla^2 + \frac{g}{8m^2} |\psi|^2 - m\Phi \right) \psi \\ \nabla^2 \Phi = 4\pi G(m|\psi|^2 - \bar{\rho}) \end{array} \right.$$

Schrödinger equation
(Gross-Pitaevskii)

Poisson equation

$g = 0 \rightarrow$ FDM
 $g \neq 0 \rightarrow$ SIFDM

Fundamentally different than
CDM/WDM/SIDM!

Madelung equations ($\psi \equiv \sqrt{\rho/m} e^{i\theta}$ and $\mathbf{v} \equiv \nabla\theta/m$)

$$\dot{\rho} + \nabla \cdot (\rho \mathbf{v}) = 0$$

$$\dot{\mathbf{v}} + (\mathbf{v} \cdot \nabla) \mathbf{v} = -\frac{1}{m} \left(V_{grav} - P_{int} - \frac{1}{2m} \frac{\nabla^2 \sqrt{\rho}}{\sqrt{\rho}} \right)$$

$$P_{int} = K \rho^{(j+1)/j} = \frac{g}{2m^2} \rho^2$$

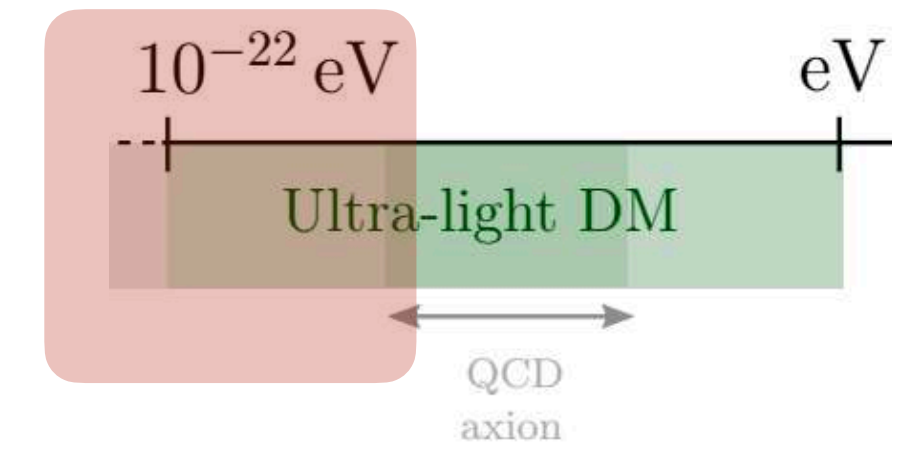
Quantum pressure

Finite Jeans length -
Suppresses
structure formation
on small scales

FLUID
DESCRIPTION

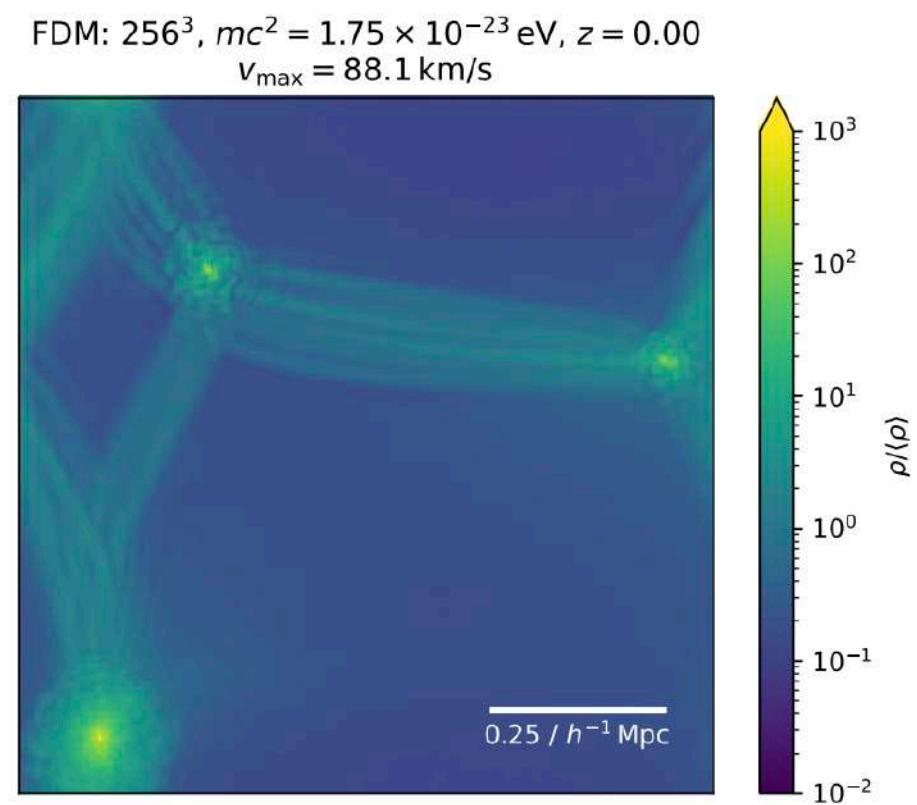
Phenomenology

RICH PHENOMENOLOGY ON SMALL SCALES

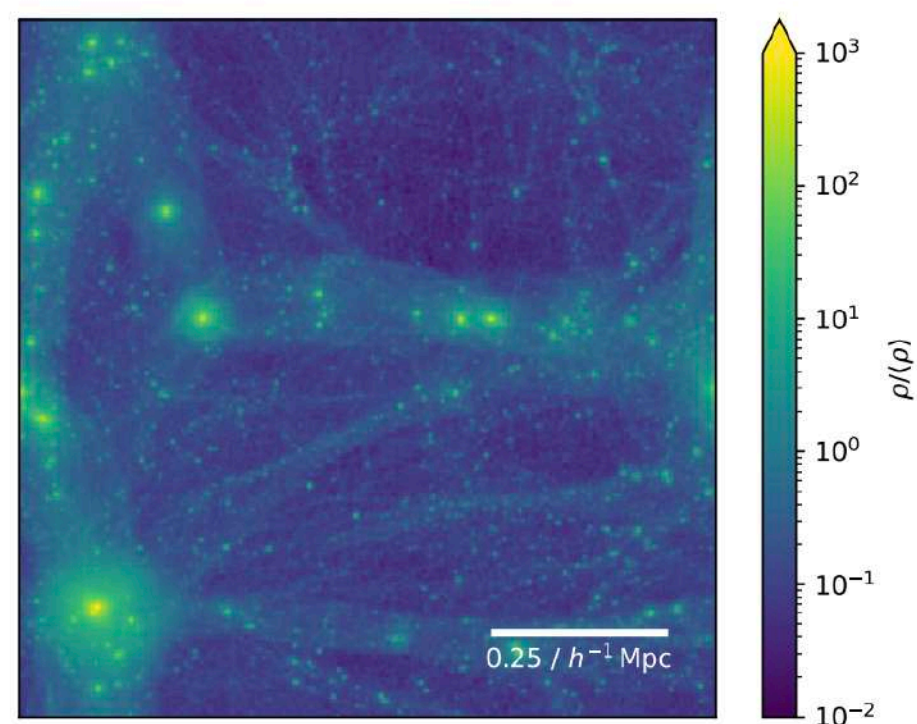


* Focus only in gravitational signatures

Suppression of small structures

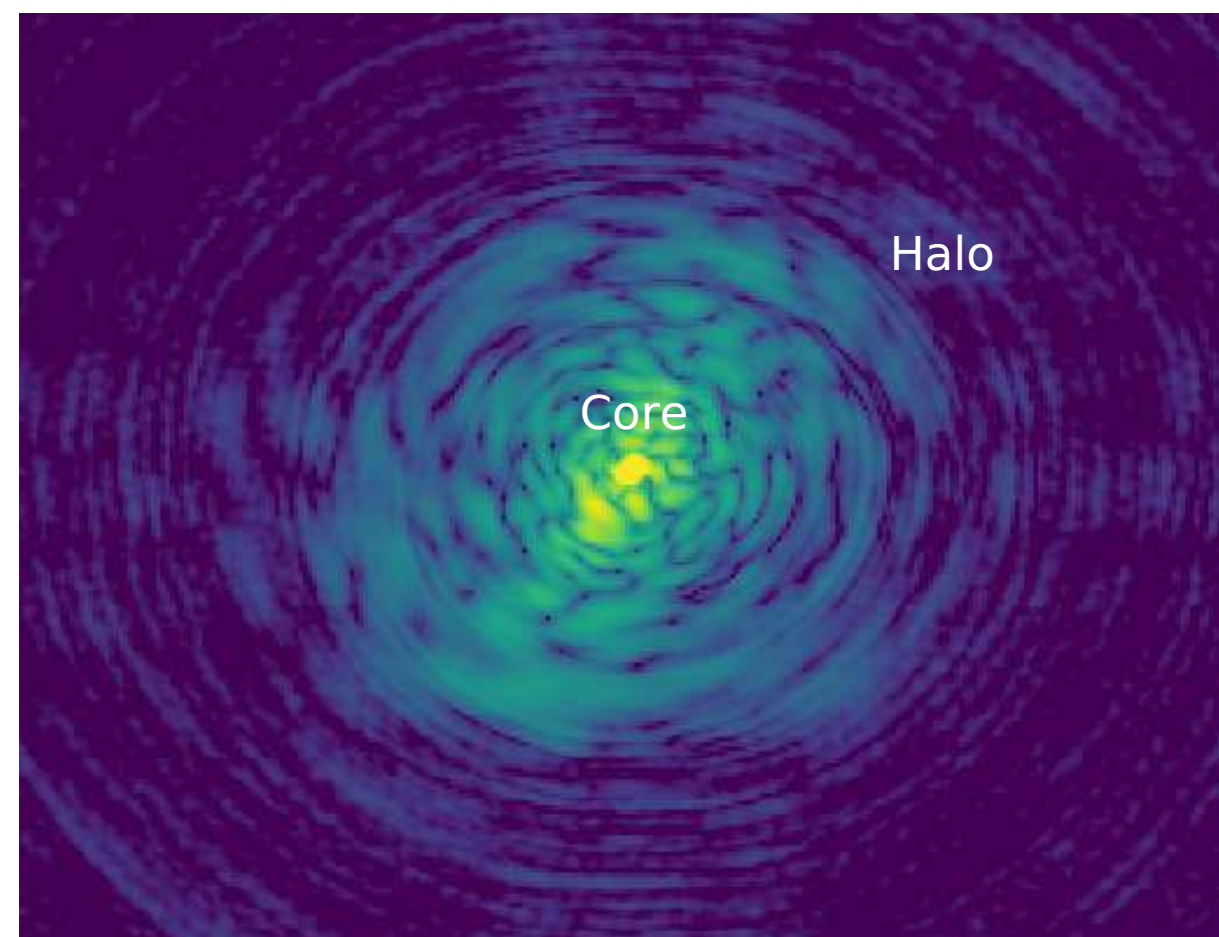


CDM: 256^3 , $z = 0.00$

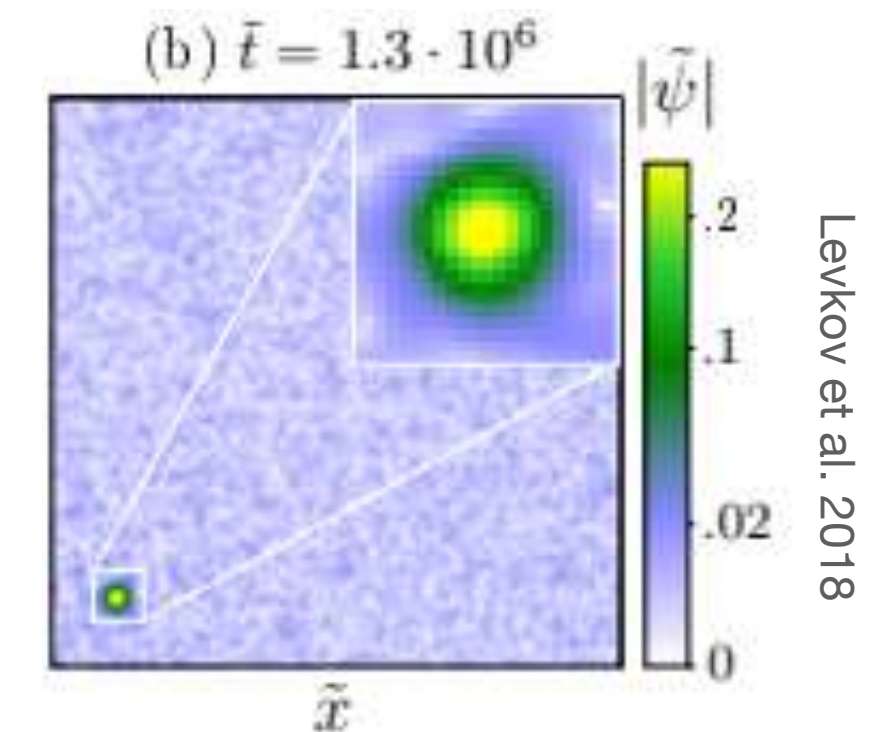


S. May et al. 2021

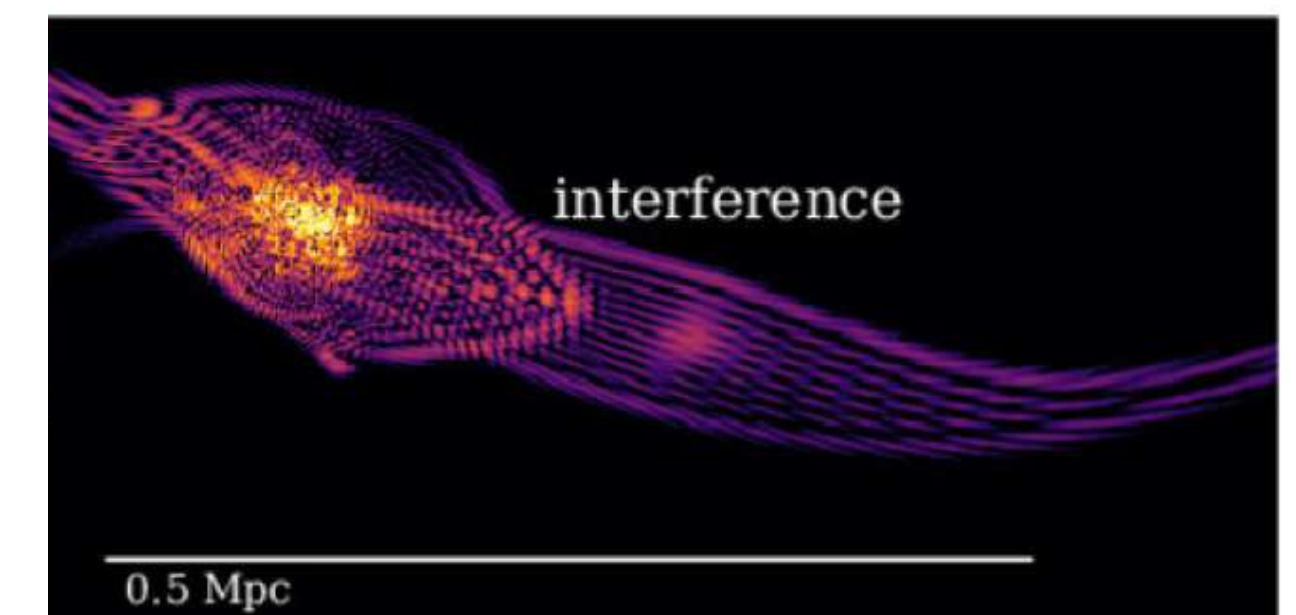
Formation of a solitonic core



Dynamical effects

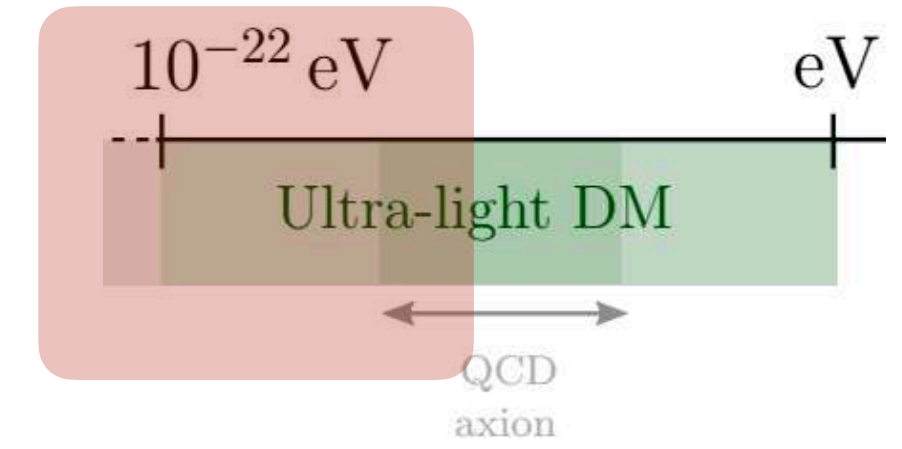


Wave interference



Mocz et al. 2017

Observational implications and constraints

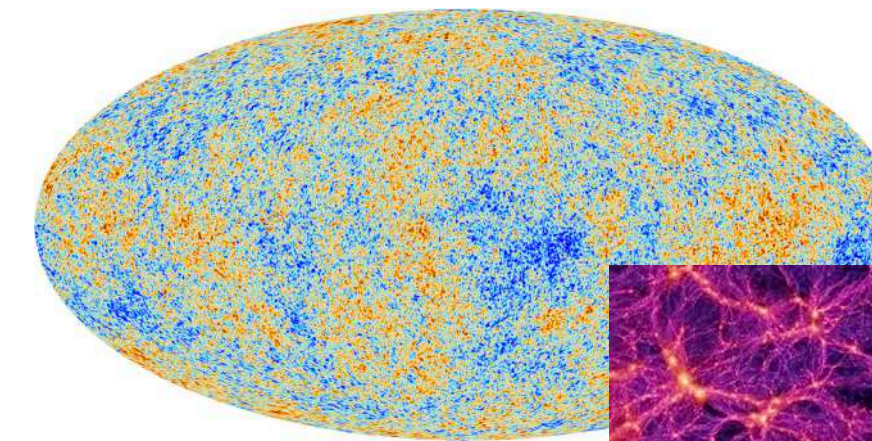


Galaxies

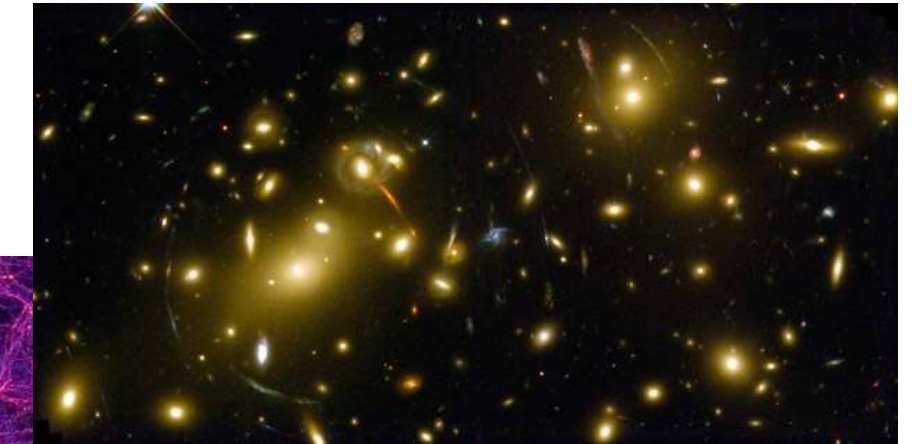


NASA and ESA

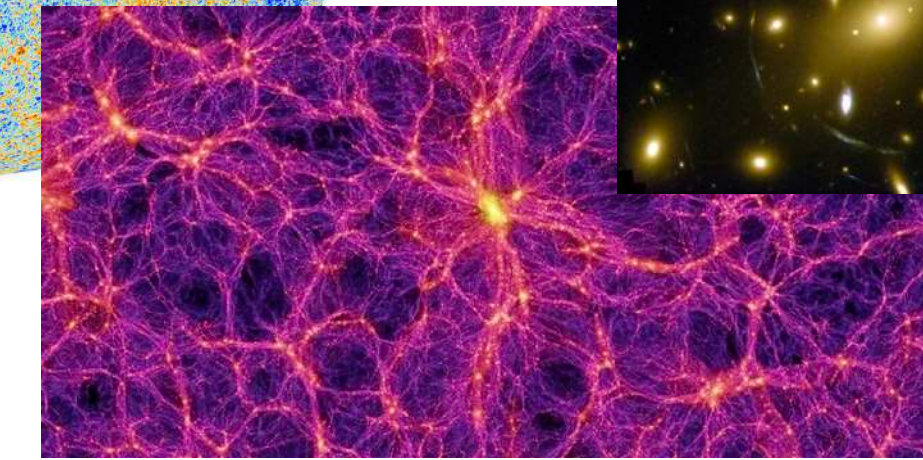
CMB+LSS



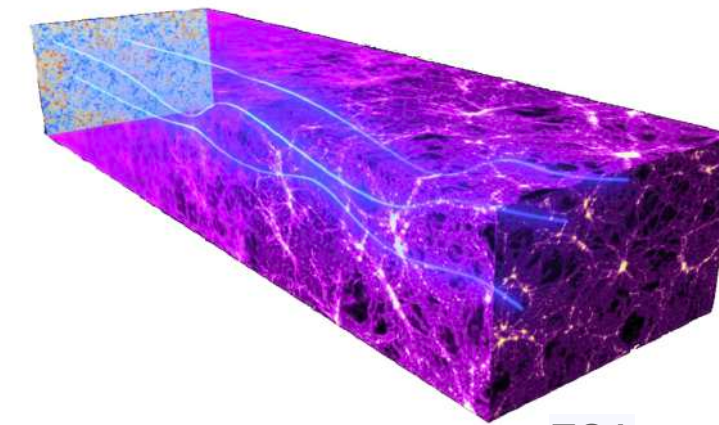
ESA and the Planck Collaboration



NASA and ESA

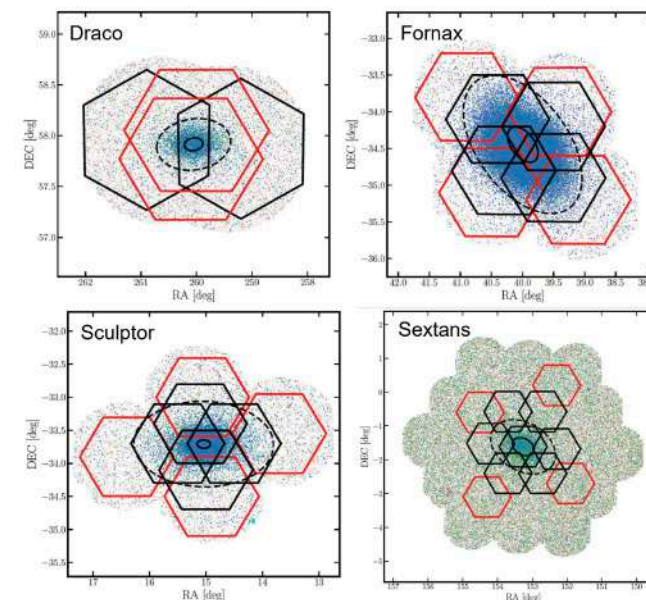


Springel & others / Virgo Consortium

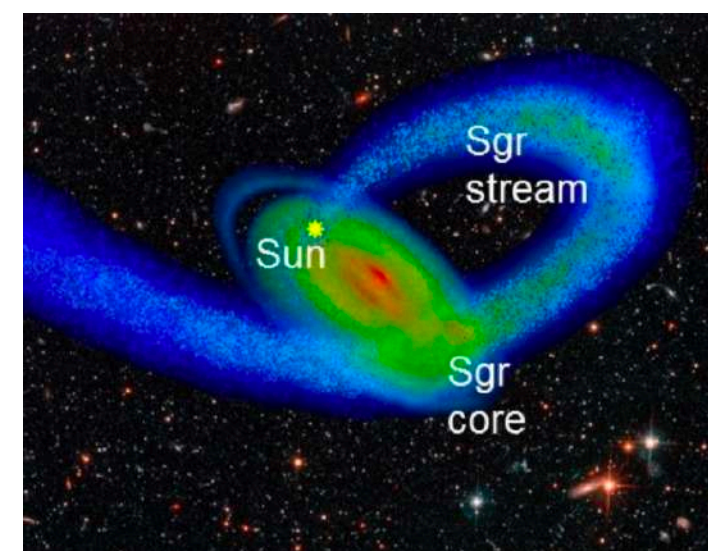


ESA

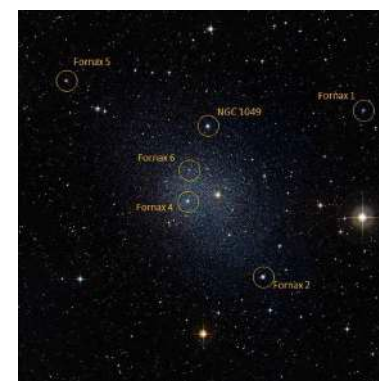
Dwarfs



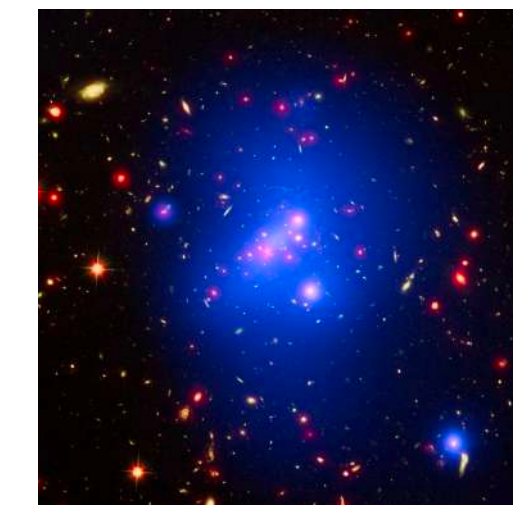
Stellar stream



Globular clusters

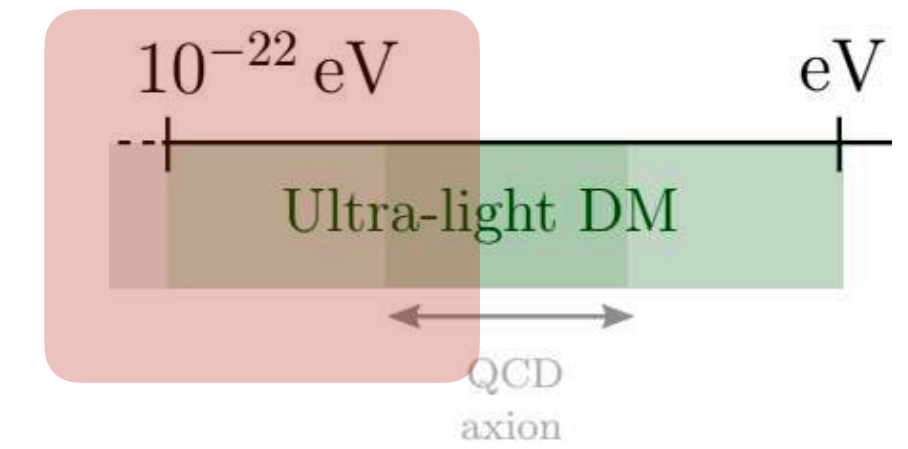


Clusters

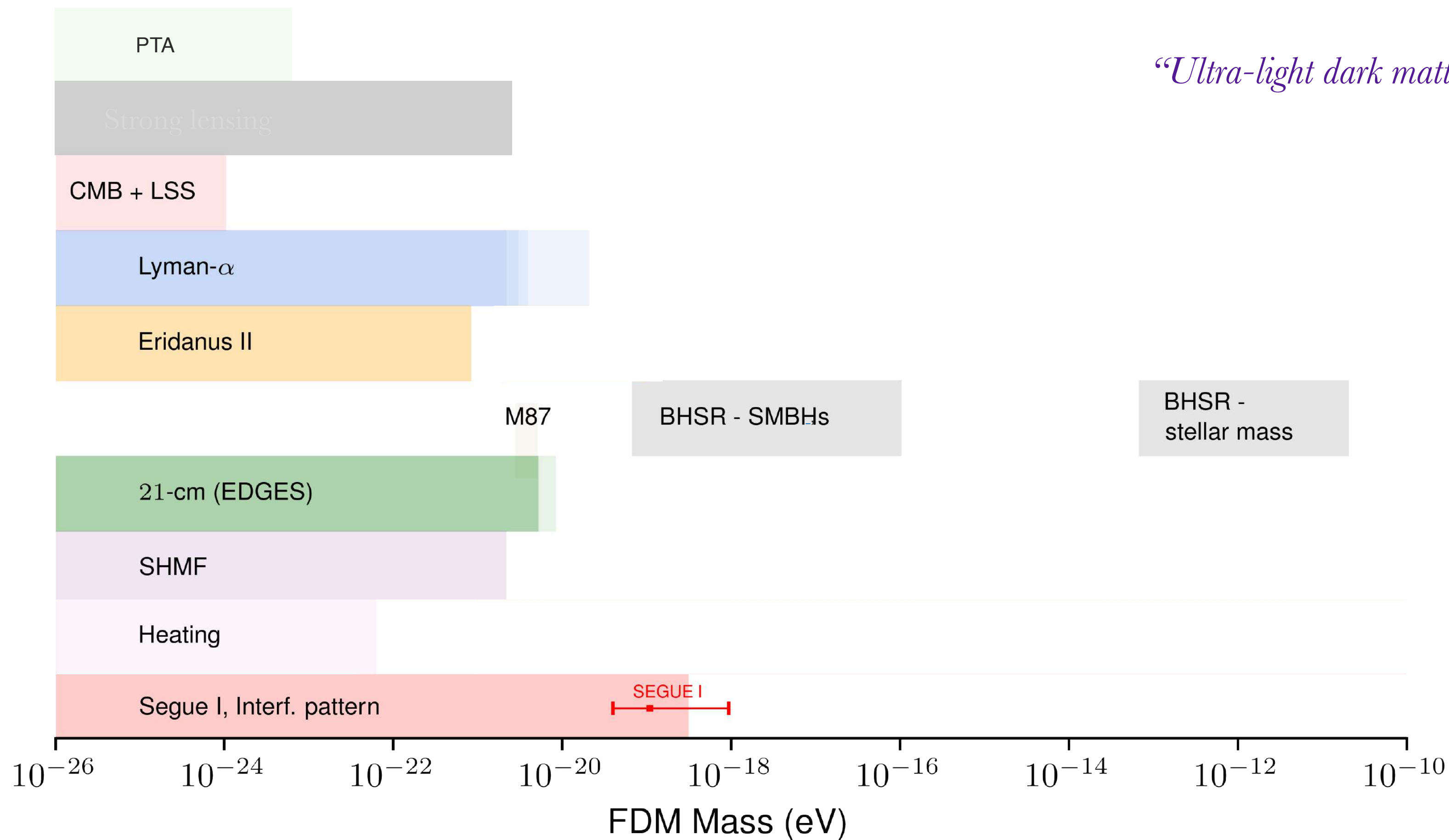


Observational implications and constraints

Fuzzy Dark Matter - bounds on the mass



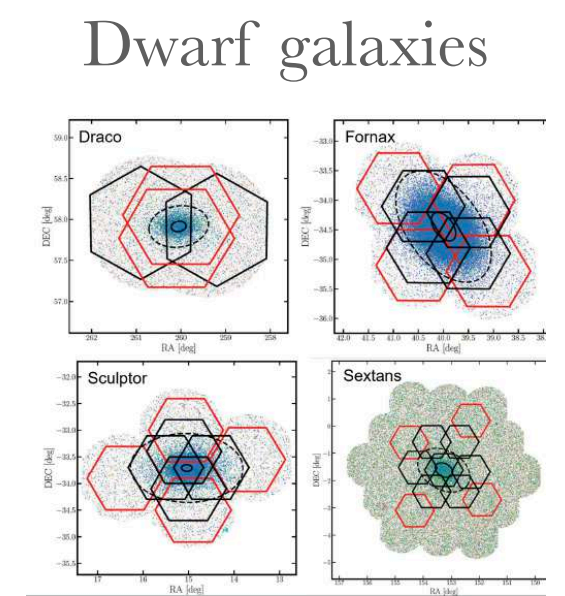
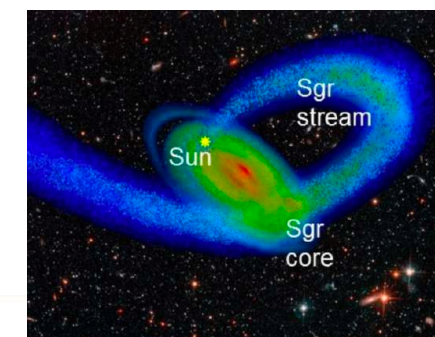
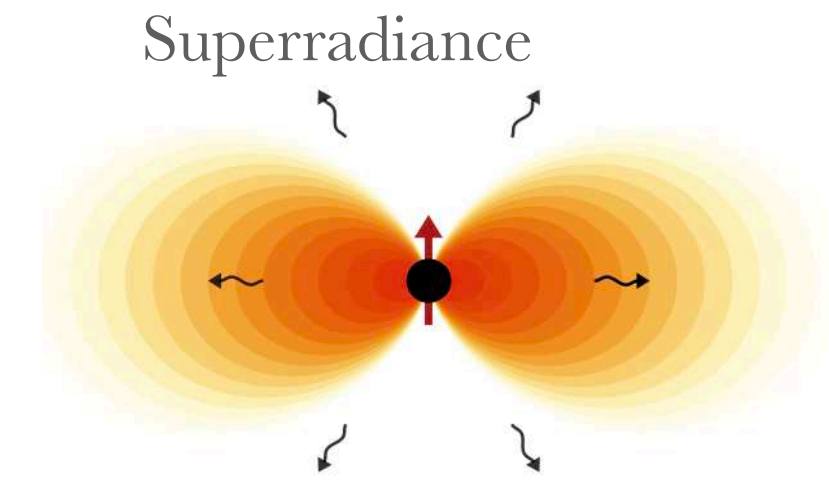
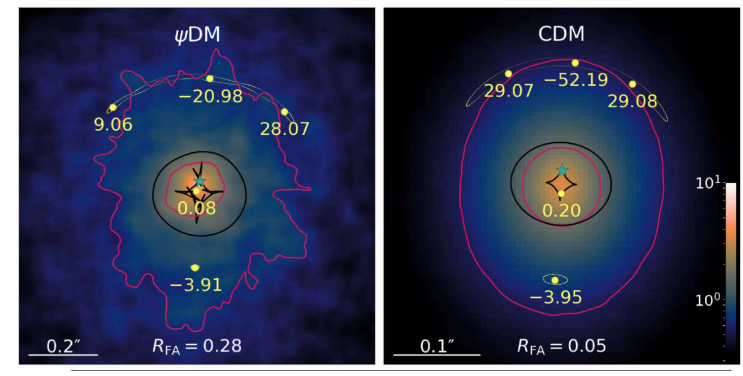
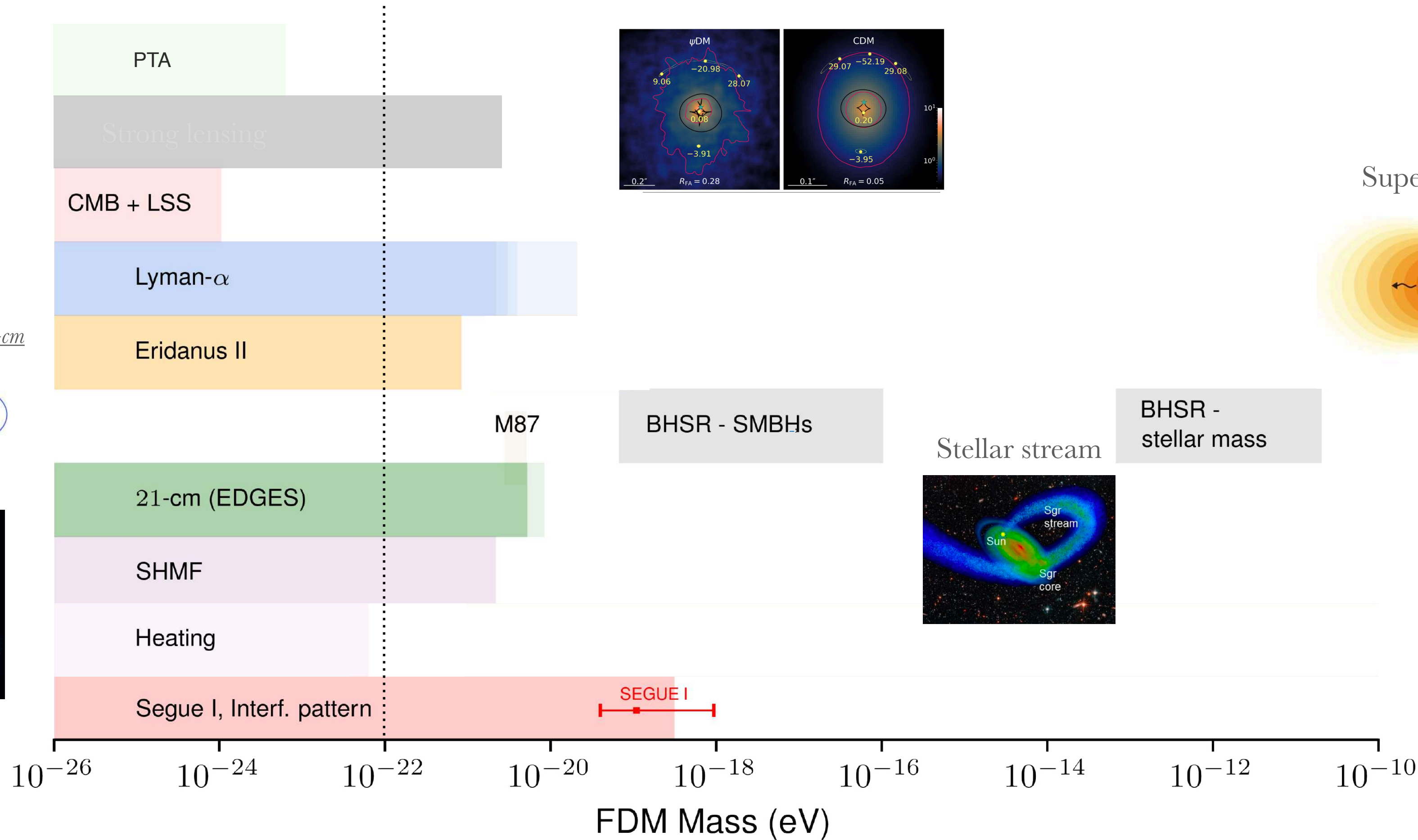
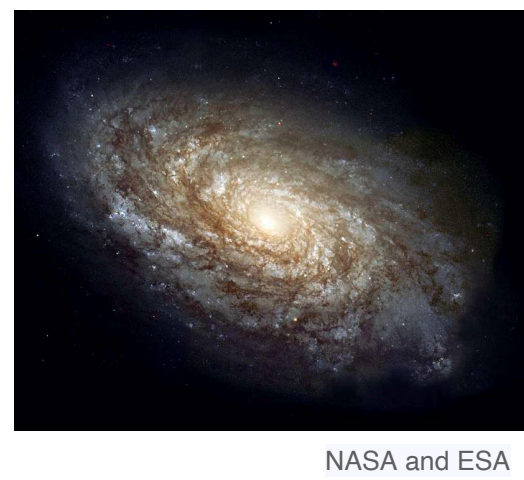
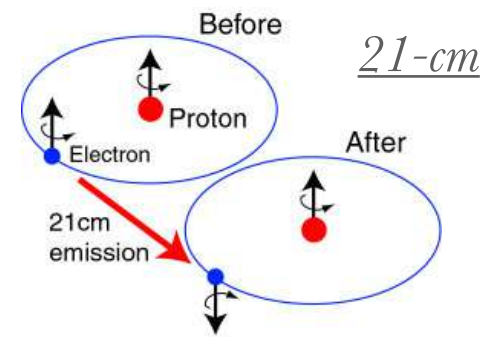
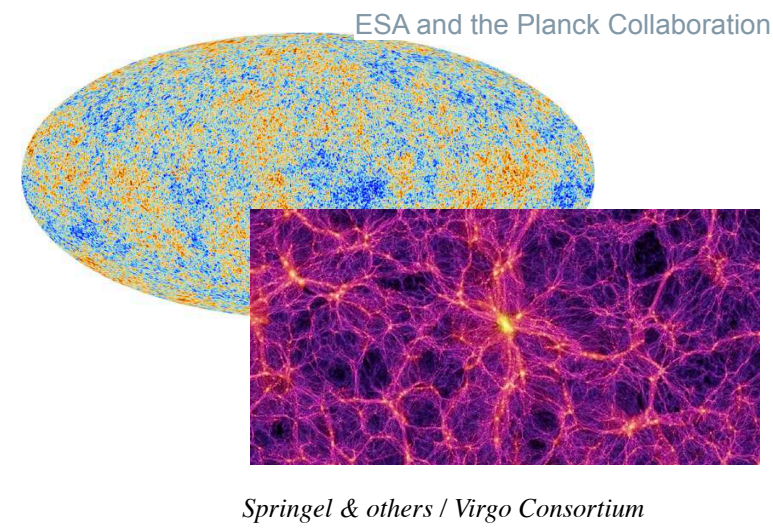
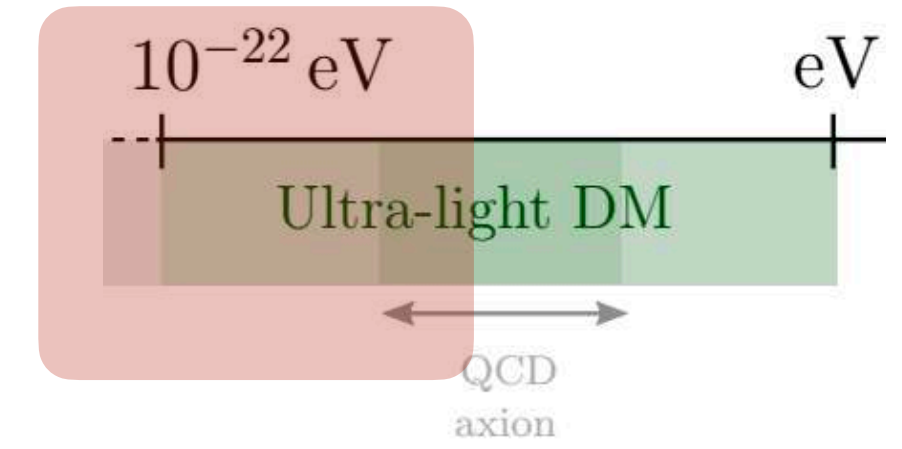
“Ultra-light dark matter”, **E.F.**, 2020. *The Astronomy and Astrophysics Review.*



Bounds consider FDM is *all* DM

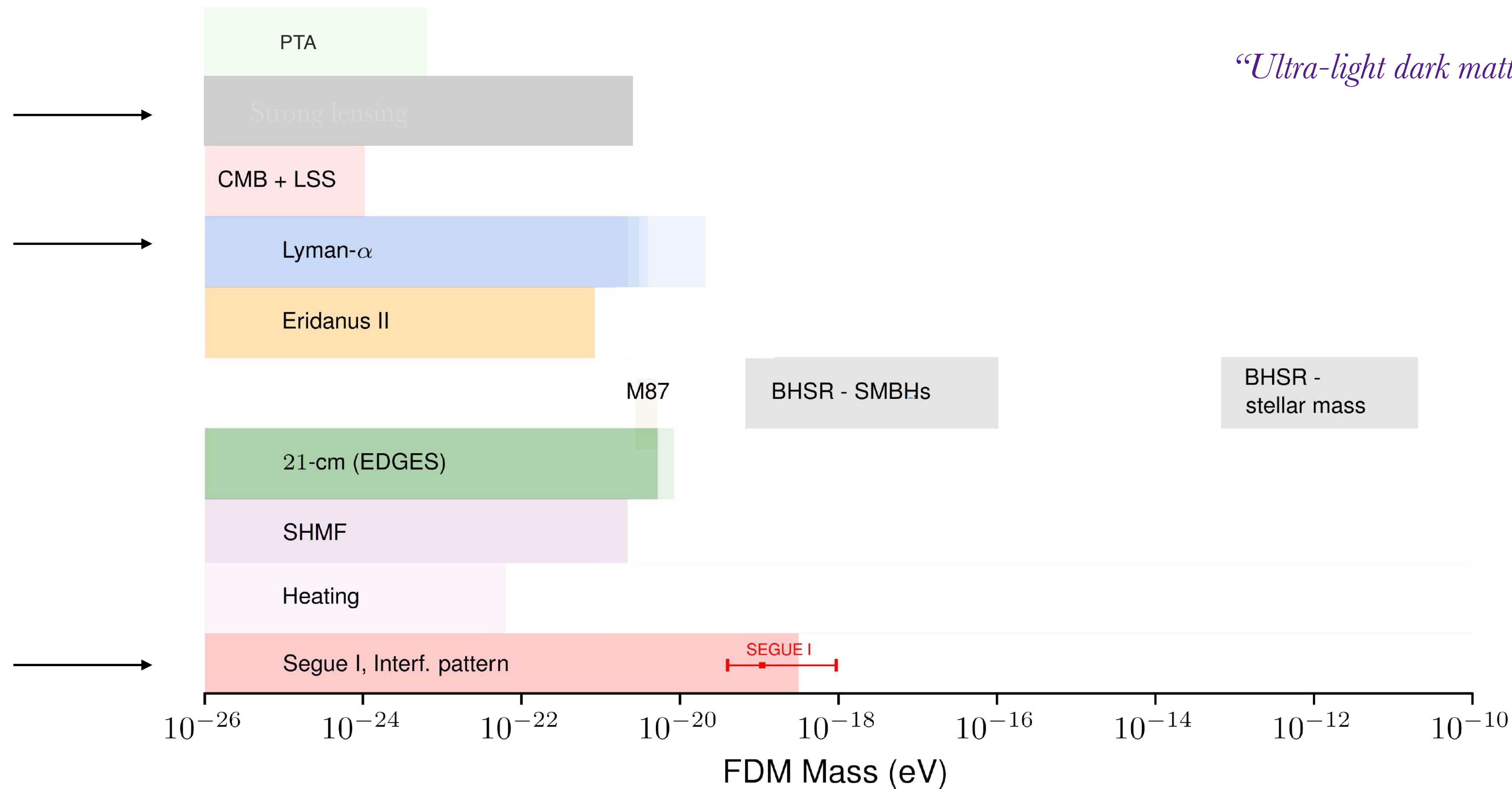
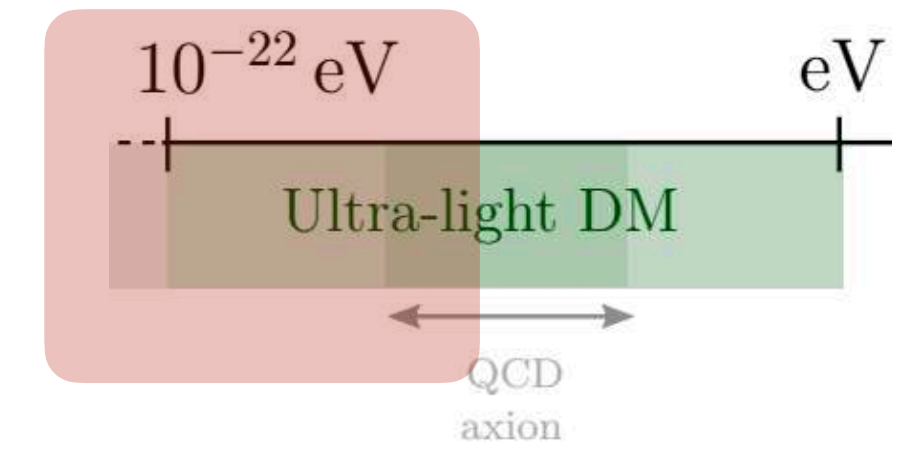
Current status

Fuzzy Dark Matter - bounds on the mass



Observational implications and constraints

Fuzzy Dark Matter - bounds on the mass

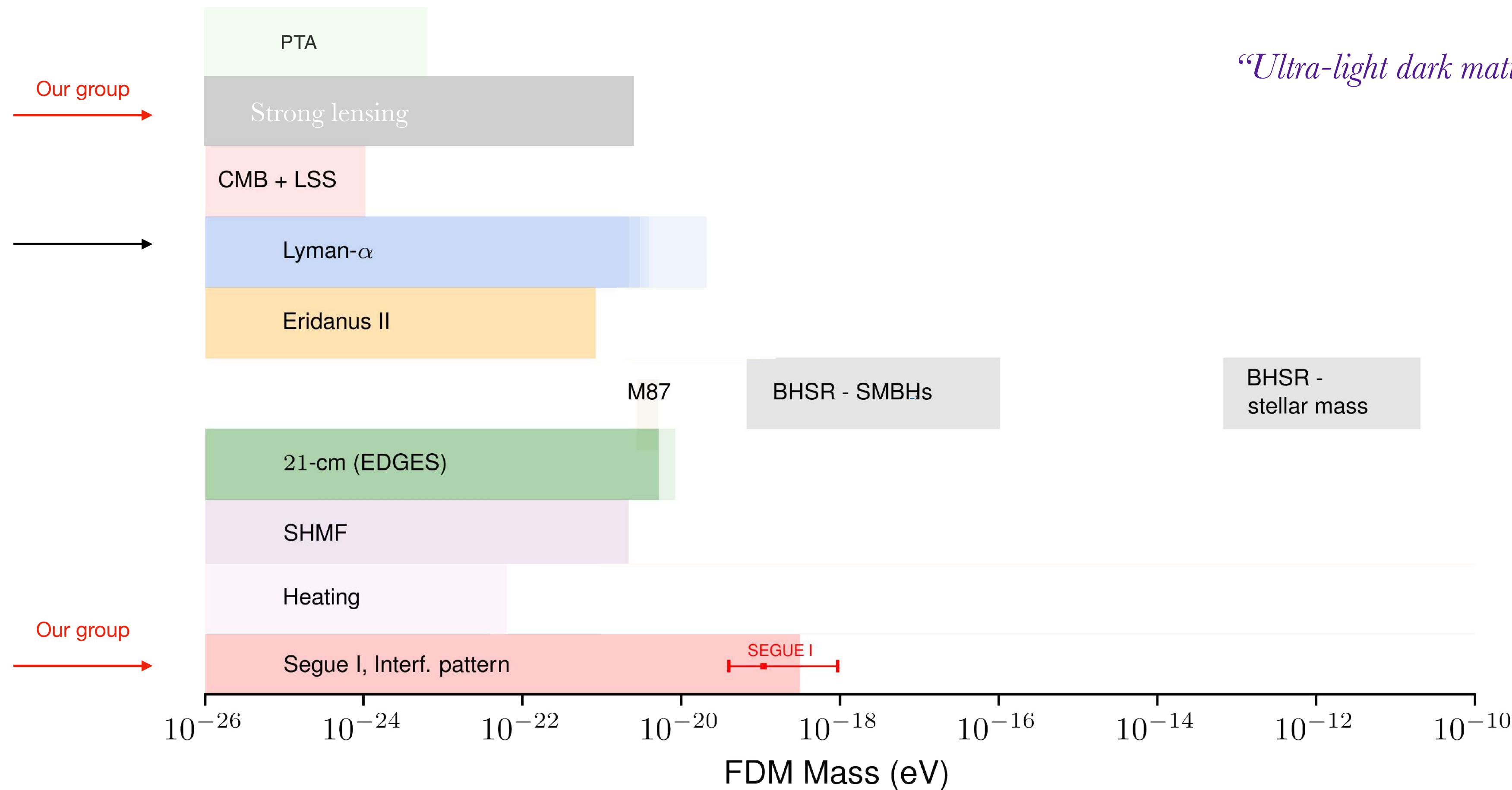
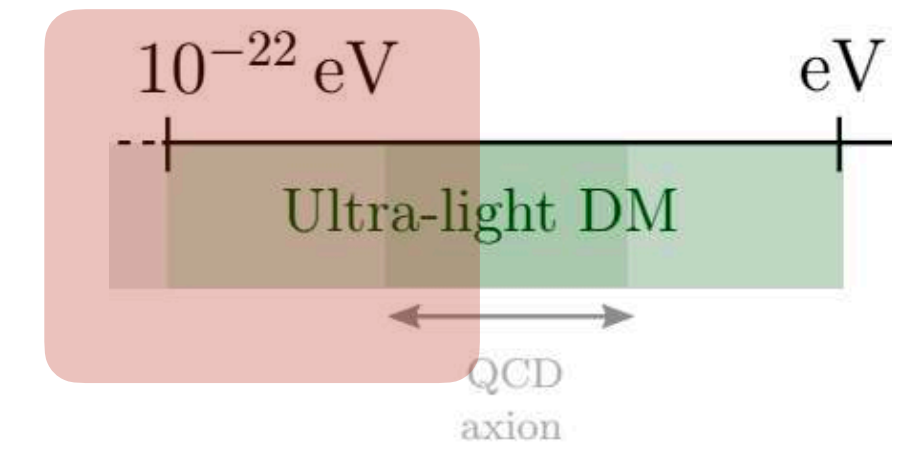


“Ultra-light dark matter”, **E.F.**, 2020. The Astronomy and Astrophysics Review.

Bounds consider FDM is *all* DM

Observational implications and constraints

Fuzzy Dark Matter - bounds on the mass



“Ultra-light dark matter”, **E.F.**, 2020. The Astronomy and Astrophysics Review.

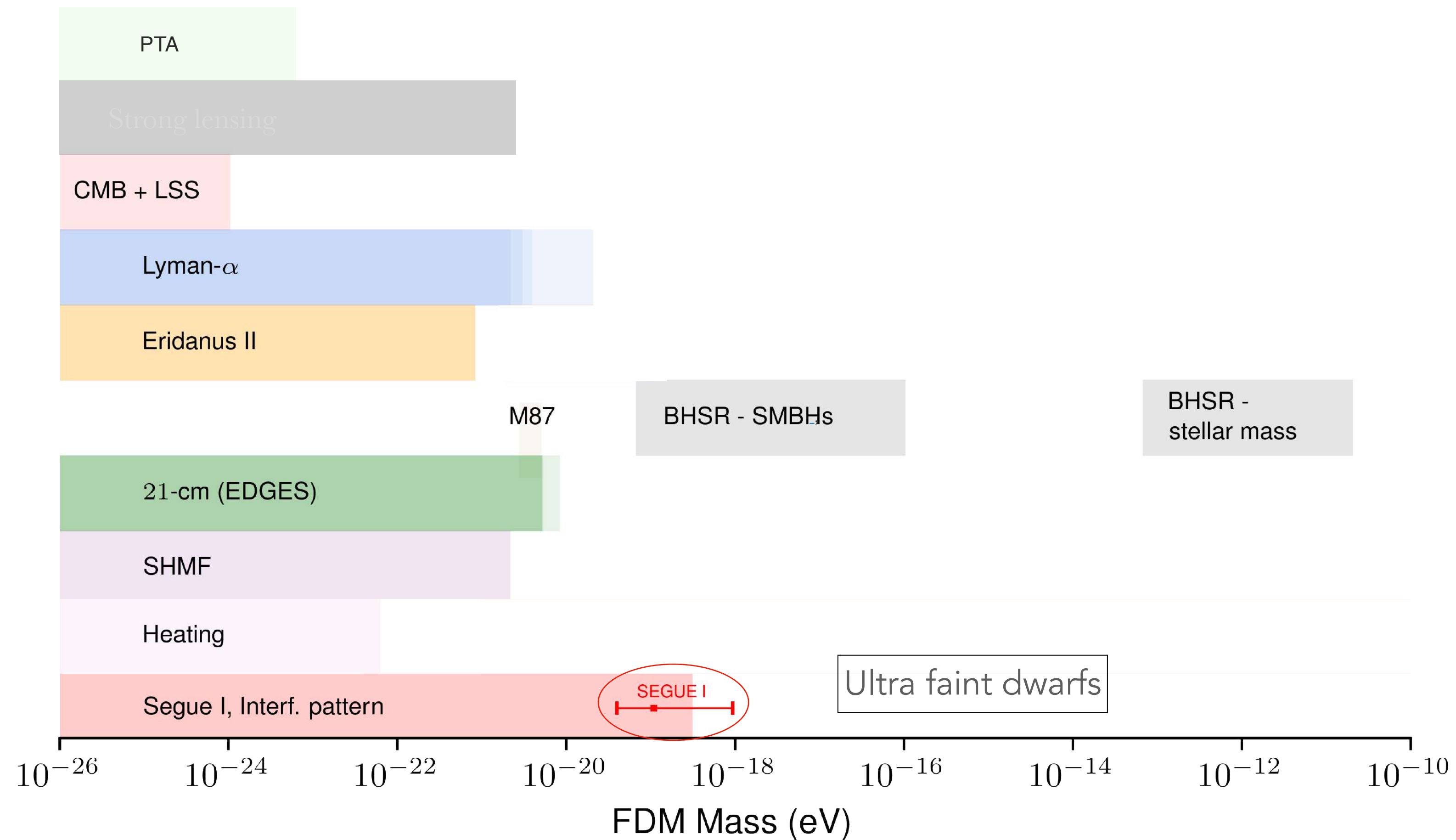
Bounds consider FDM is *all* DM

Our group has some of the strongest bounds on these models!

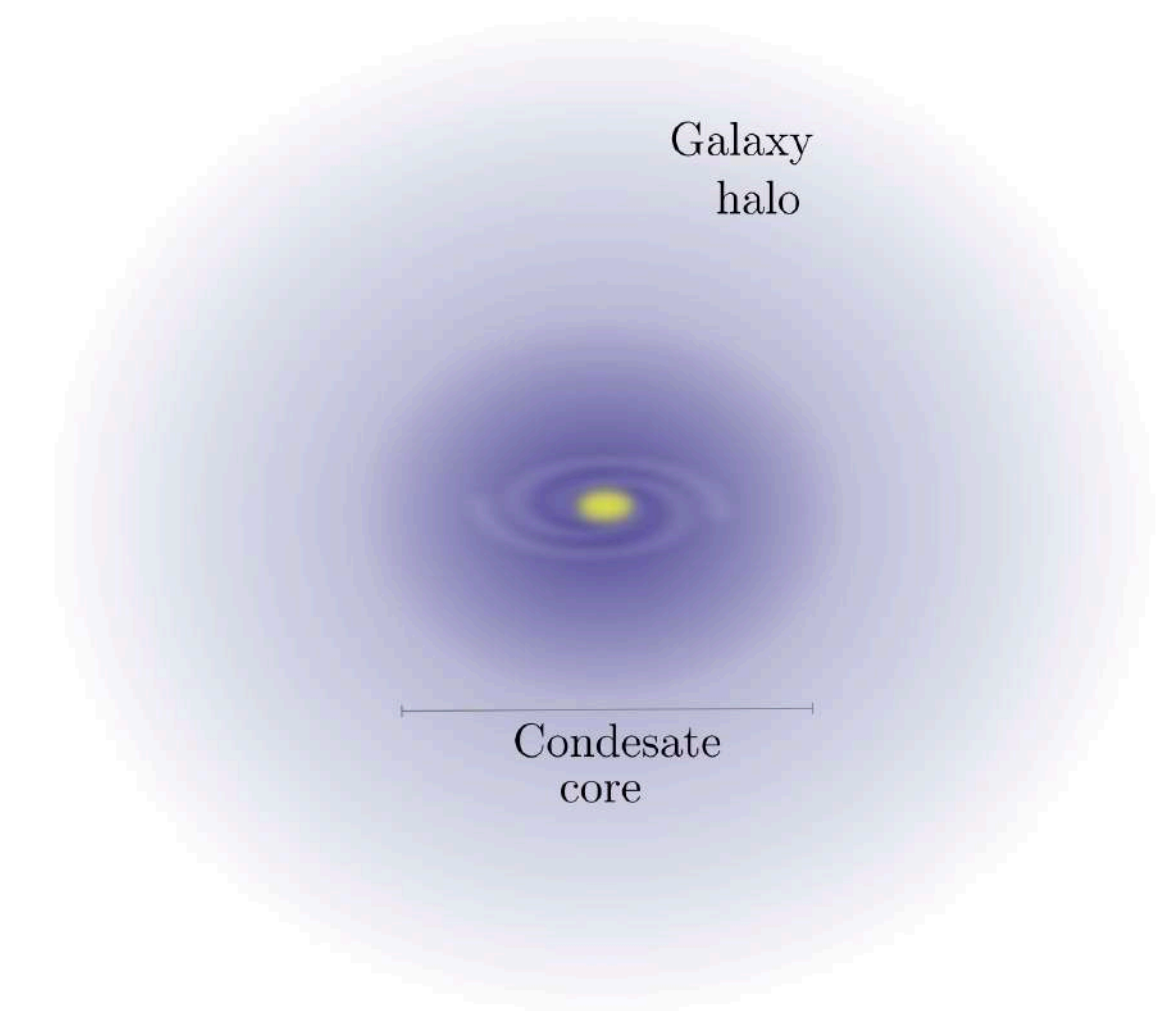
Observational implications and constraints

Fuzzy Dark Matter - bounds on the mass

“Narrowing the mass range of Fuzzy Dark Matter with Ultra-faint Dwarfs”, J. Chan, E.F., K. Hayashi, 2021.

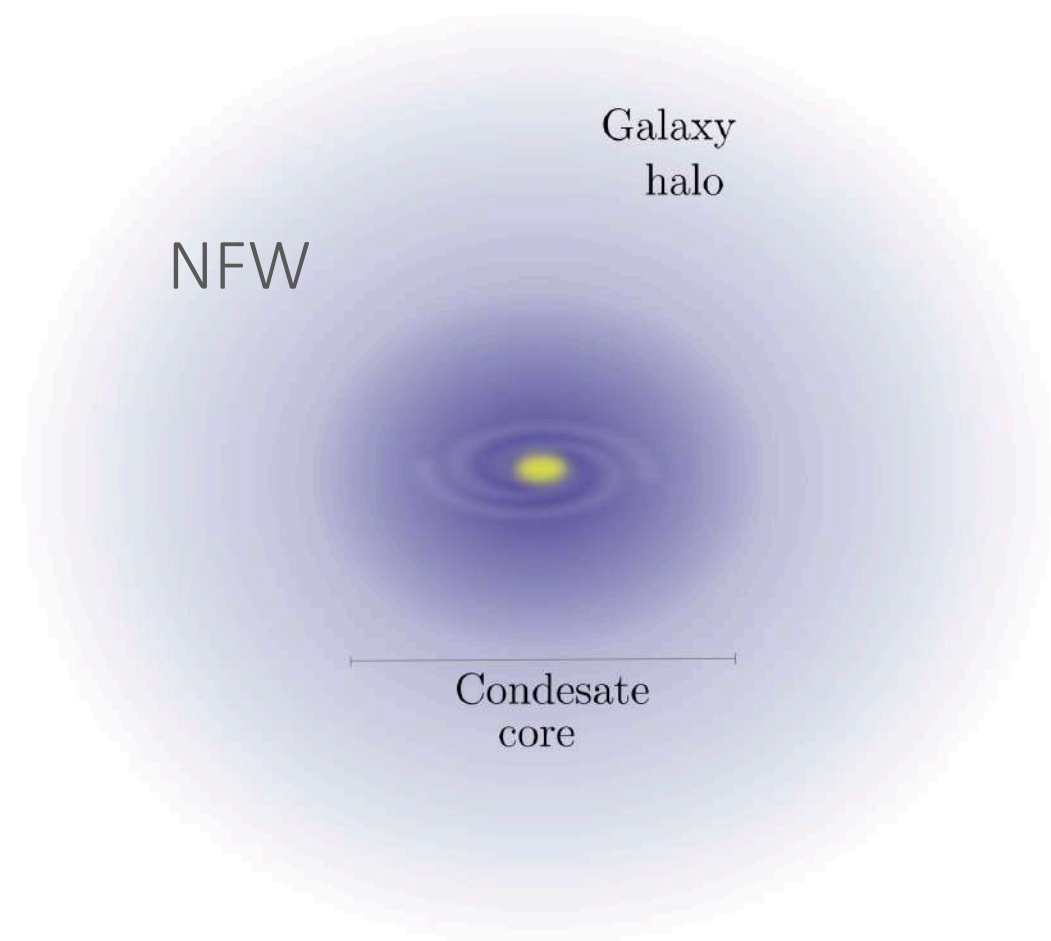


Presence of a core

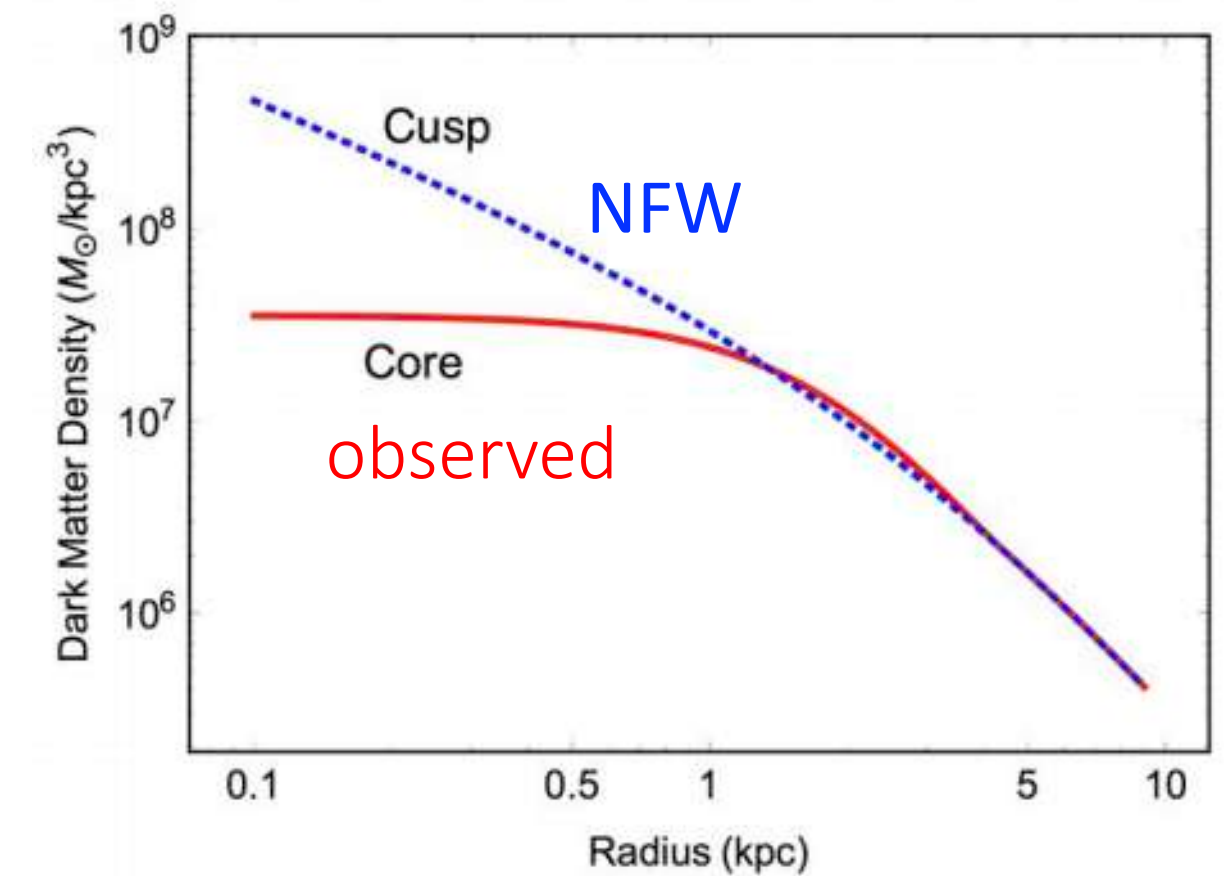
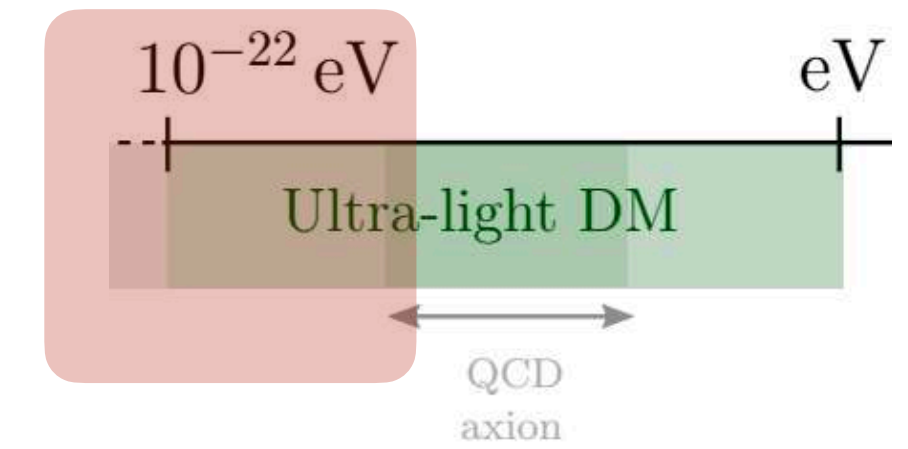


Phenomenology

Formation of **cores**



$$\rho(r) \simeq \begin{cases} \rho_c & \text{for } r \leq r_c \\ \rho_{\text{NFW}} & \text{for } r \geq r_c \end{cases}$$



FDM

From simulations Schive et al. 2014, fitting function:

Updated by our group (*Chan, EF et al 2021*)

$$\rho_c \simeq \frac{1.9 \times 10^{-2}}{[1 + 0.091 (r/R_{1/2,c})^2]^8} \left(\frac{m}{10^{-22} \text{ eV}} \right)^{-2} \left(\frac{r_c}{\text{kpc}} \right)^{-4} M_\odot \text{ pc}^{-3},$$

$$r_c \simeq 0.16 \left(\frac{m}{10^{-22} \text{ eV}} \right)^{-1} \left(\frac{M}{10^{12} M_\odot} \right)^{-1/3} \text{ kpc}.$$

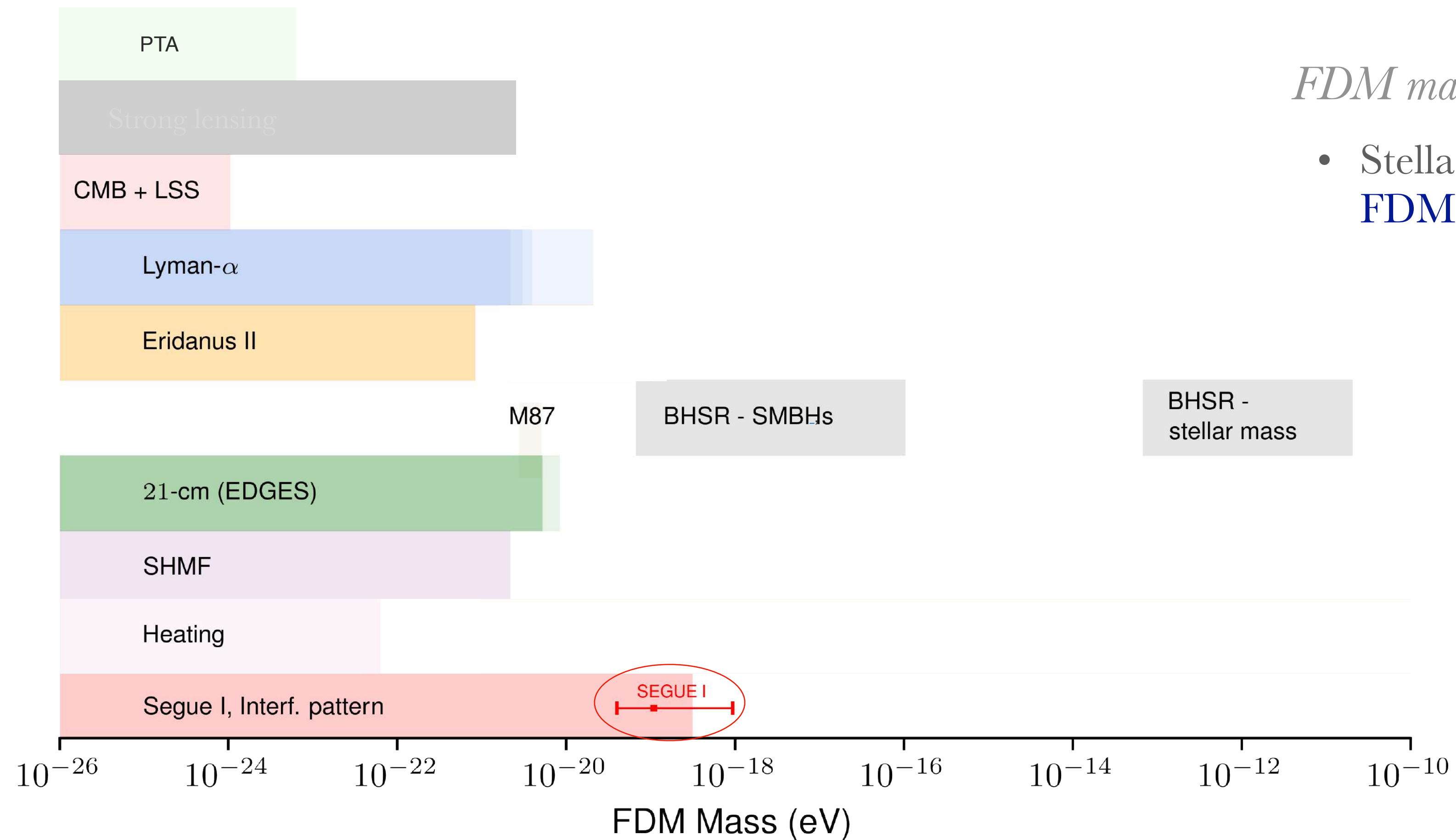
Relations used to compare with **observations**

Ultra-light Dark Matter

Fuzzy Dark Matter - bounds on the mass

Ultra faint dwarfs

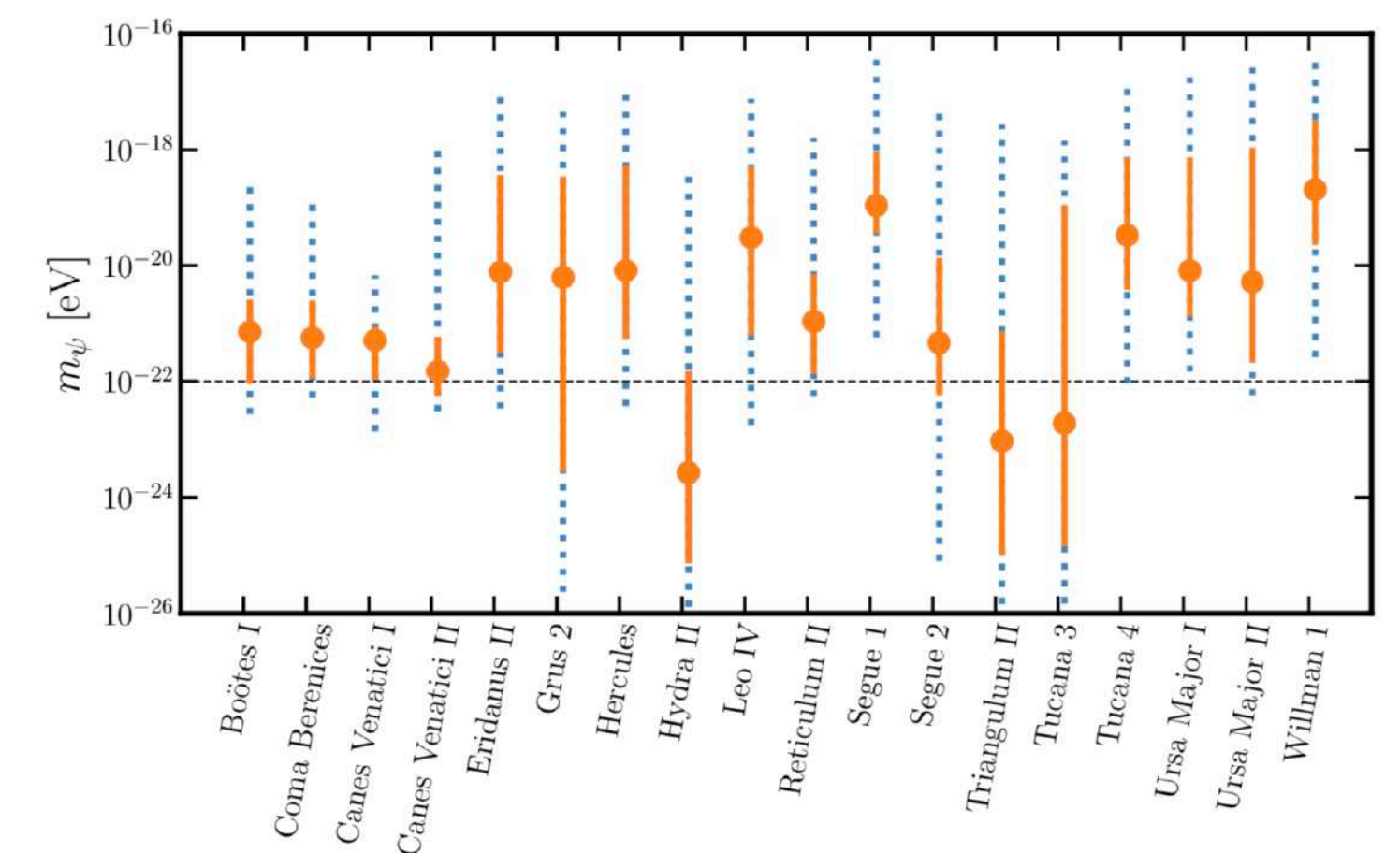
Hayashi, E.F, Chan, 2021.



FDM mass from Ultra-faint dwarfs

- Stellar kinematic data from 18 UFDs to fit the FDM profile from simulations

$$m_{\text{FDM}}^{(\text{Seg1})} = 1.1^{+8.3}_{-0.7} \times 10^{-19} \text{ eV}$$



Preference for higher mass

Ultra-light Dark Matter

Fuzzy Dark Matter - bounds on the mass

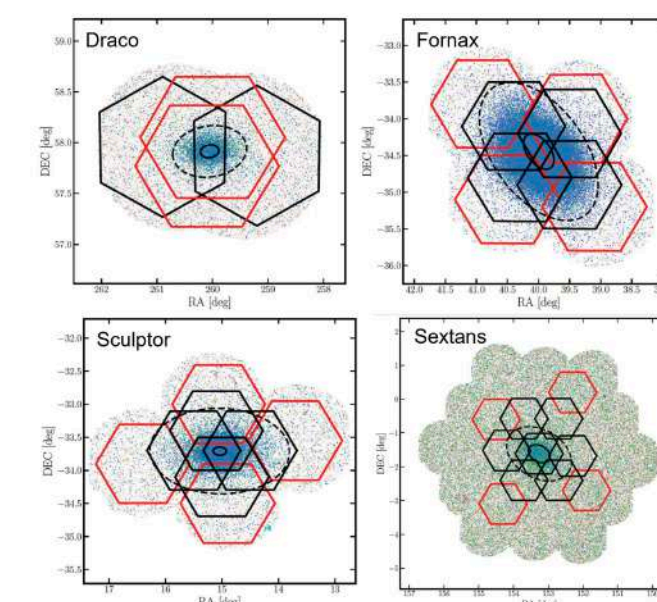
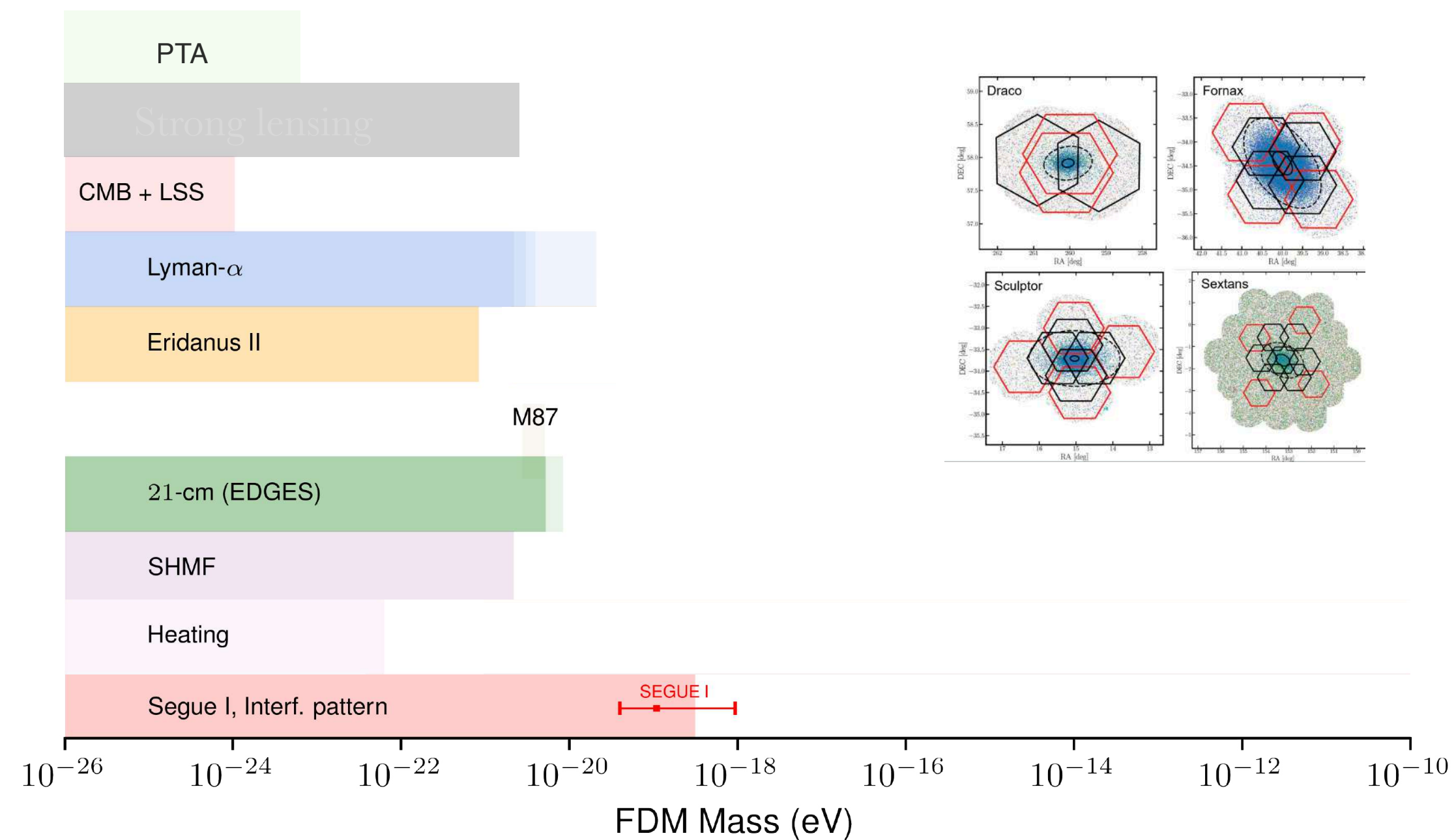
(Work in progress)



Shun'ichi
Horigome



Shin'ichiro
Ando



FDM mass from Ultra-faint dwarfs

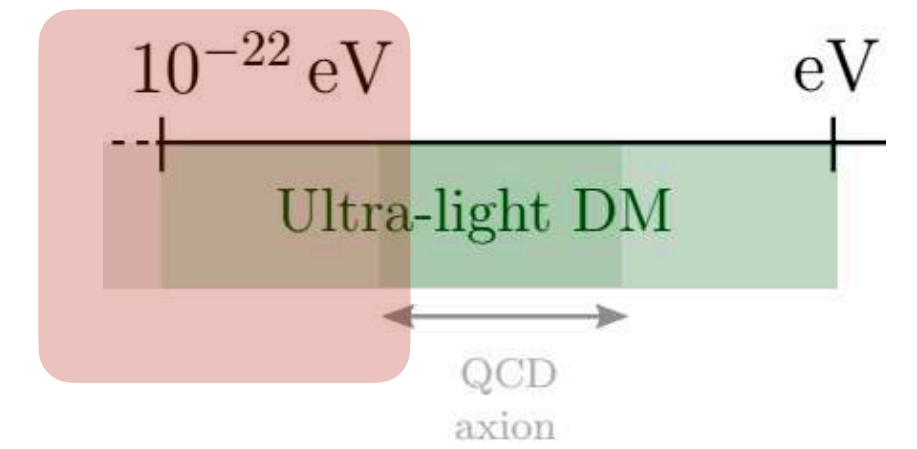
- Stellar kinematic data from 18 UFDs to fit the **FDM** profile from simulations
- Repeat the previous analysis with:
 - Improved modelling of FDM based on previous work and simulation from **our** group
 - Using environmental effects to put informed priors for Bayesian analysis the dwarfs quantities - SASHIMI



Result: more conservative and improved bounds
Lessons for **PFS** analysis!

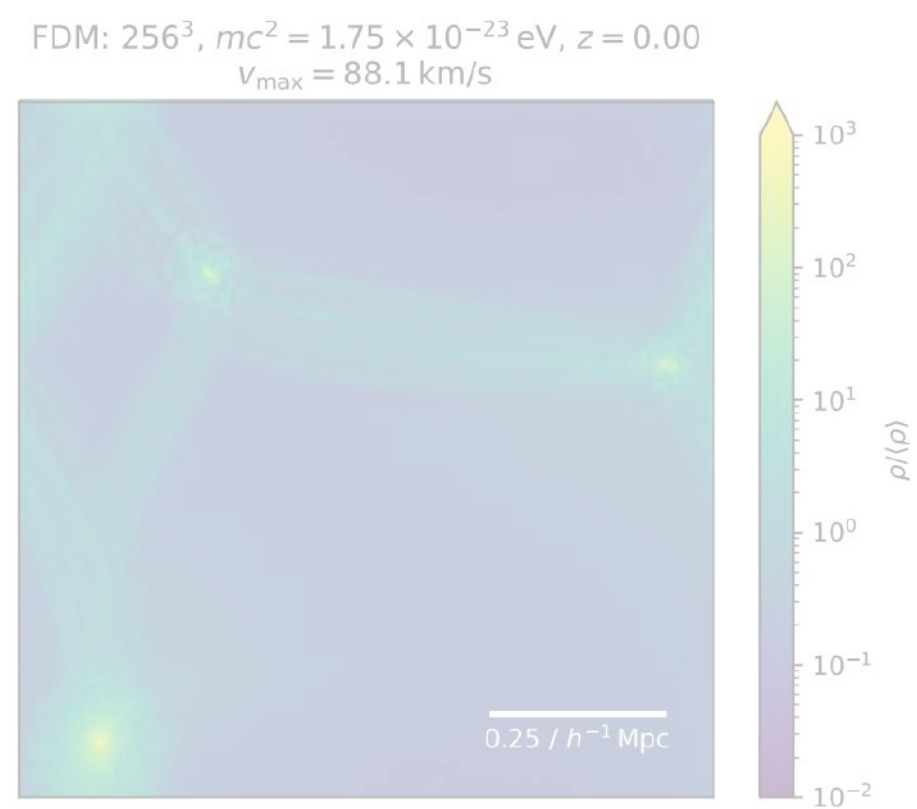
Preparation for PFS!

Phenomenology

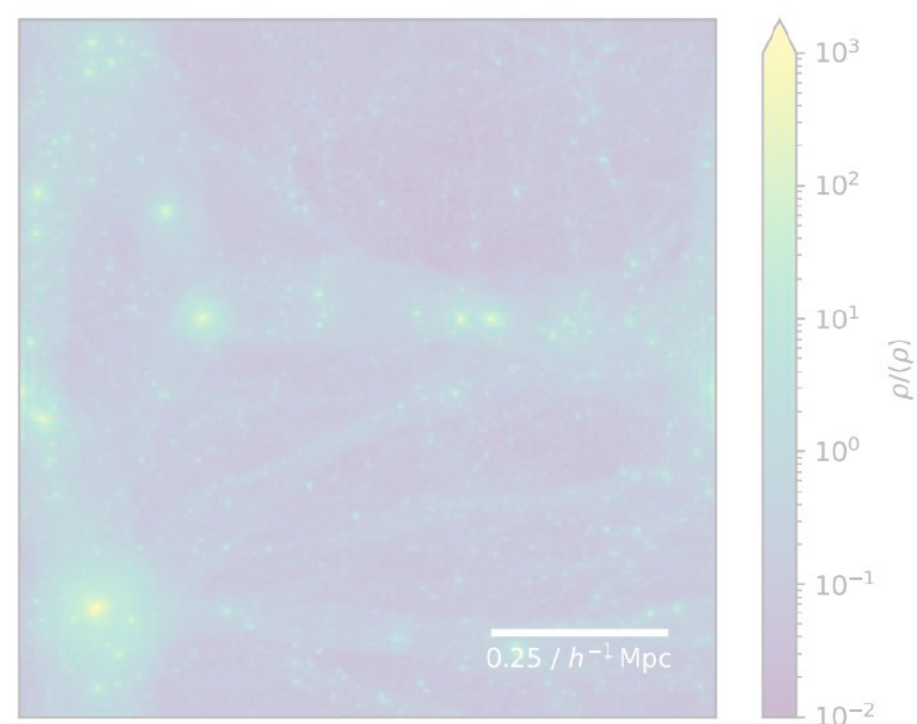


RICH PHENOMENOLOGY ON SMALL SCALES

Suppression of small structures

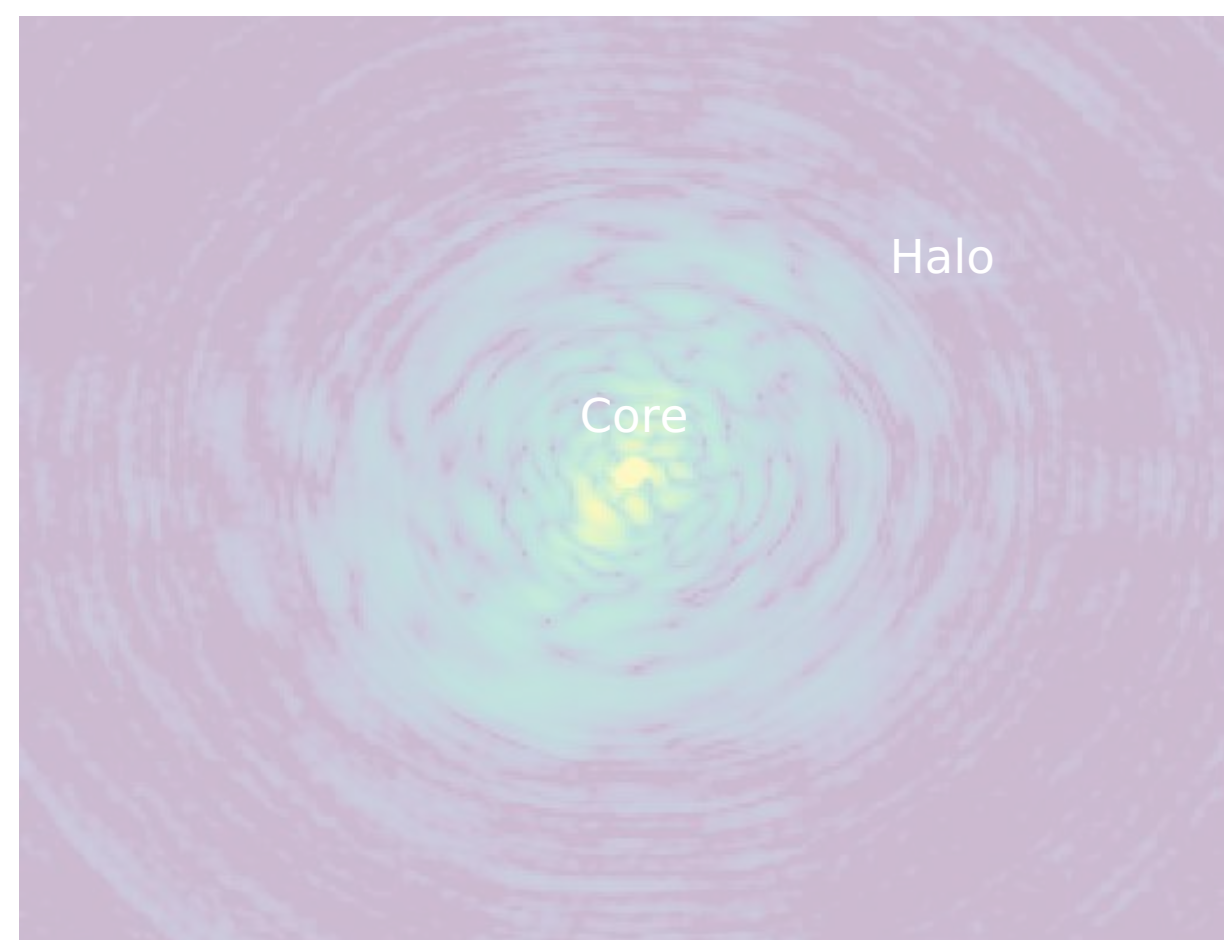


CDM: 256^3 , $z = 0.00$

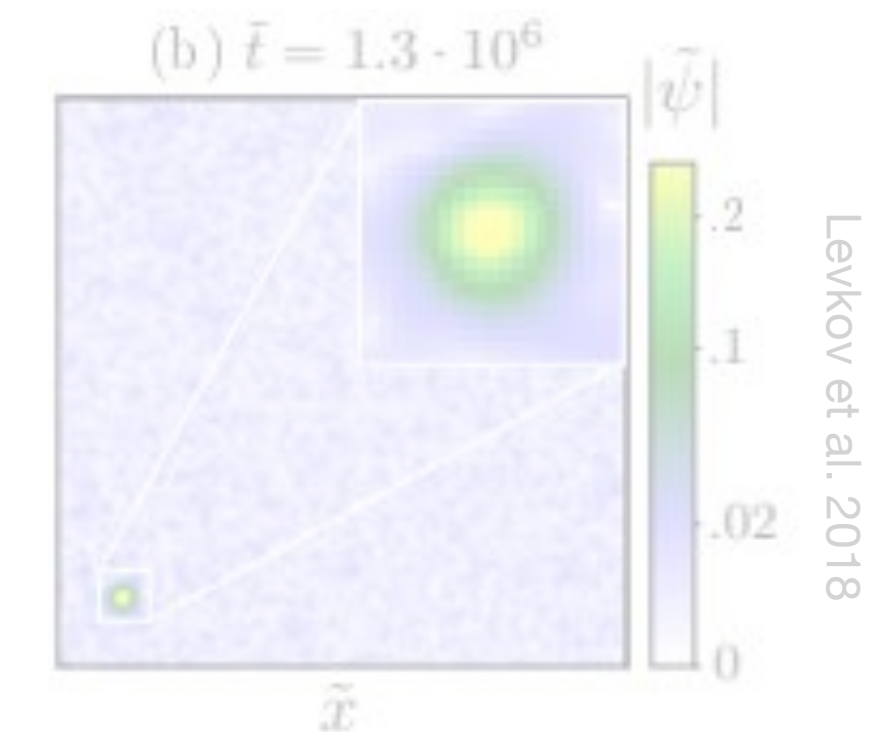


S. May et al. 2021

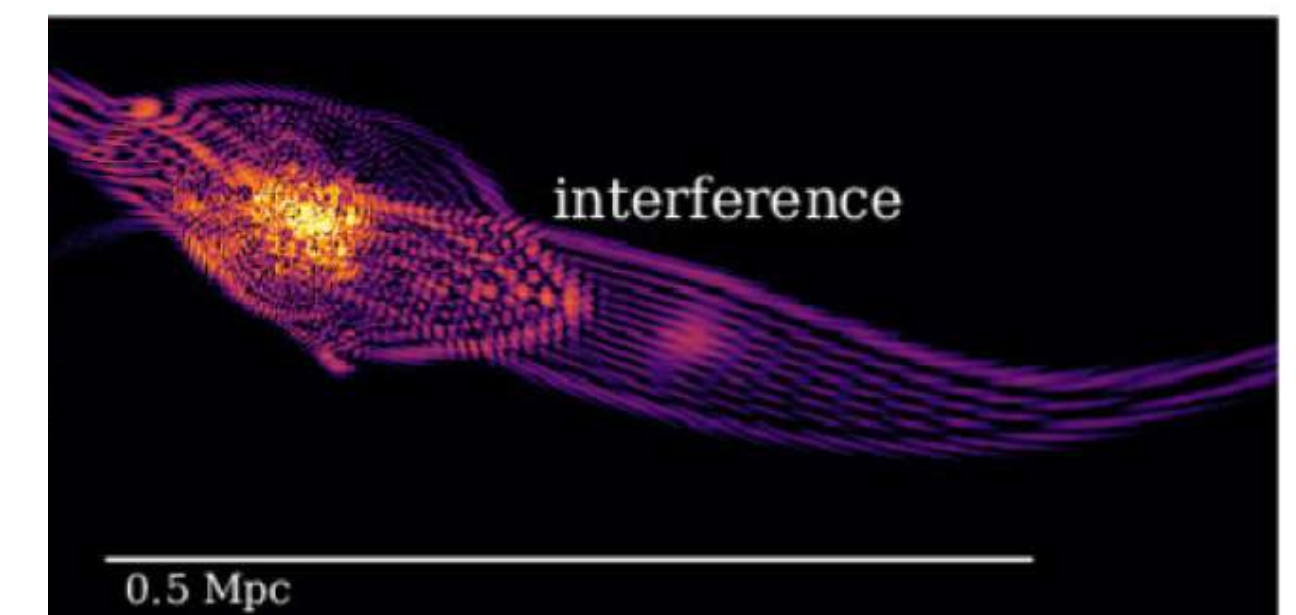
Formation of a solitonic core



Dynamical effects



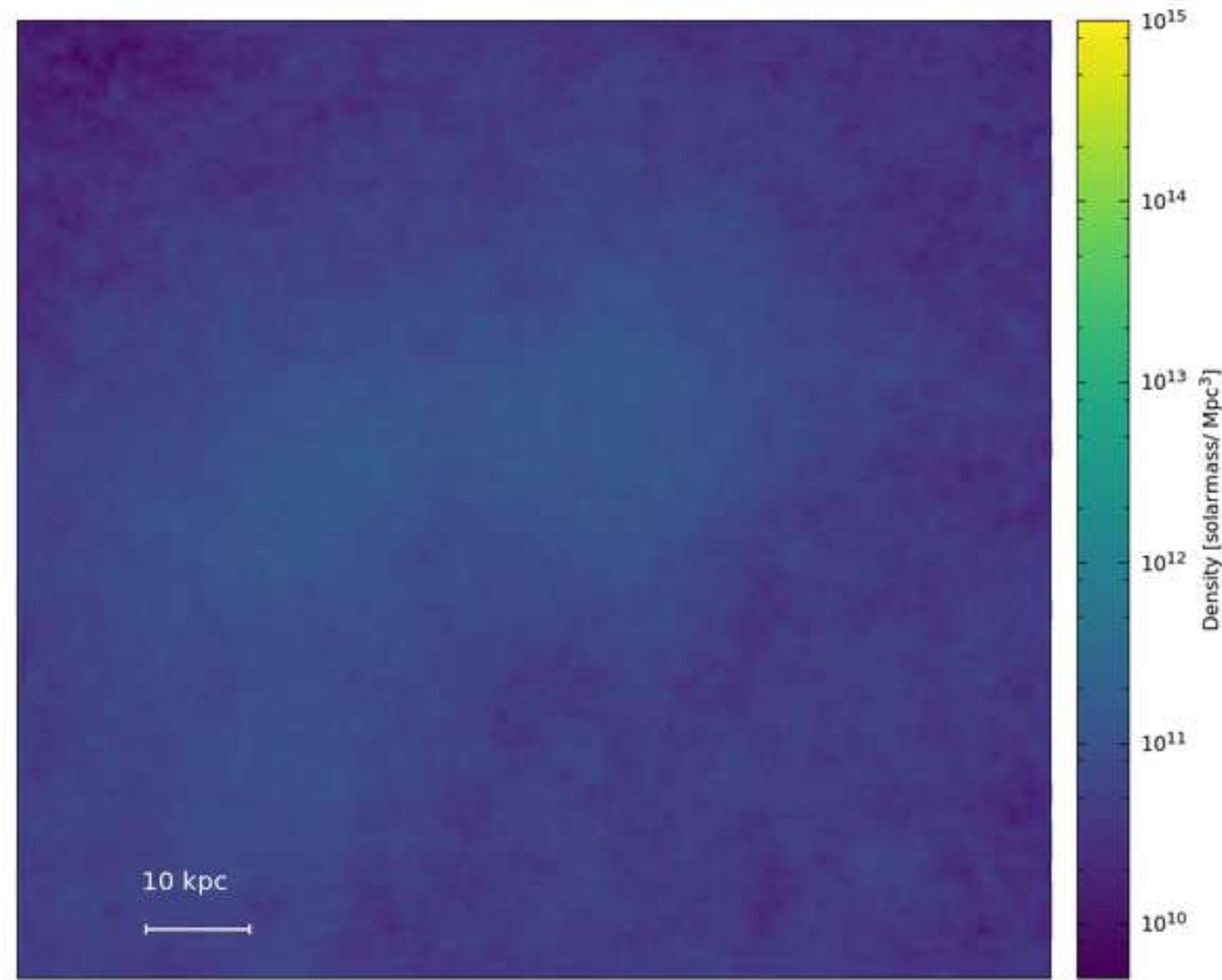
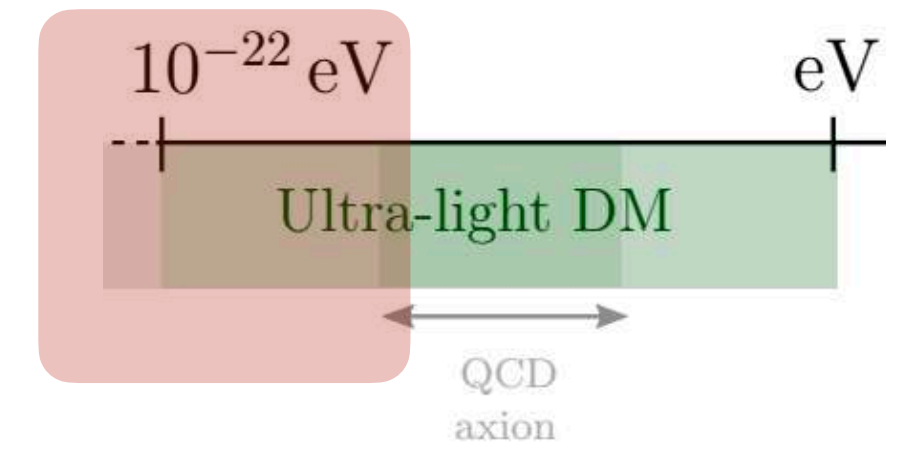
Wave interference



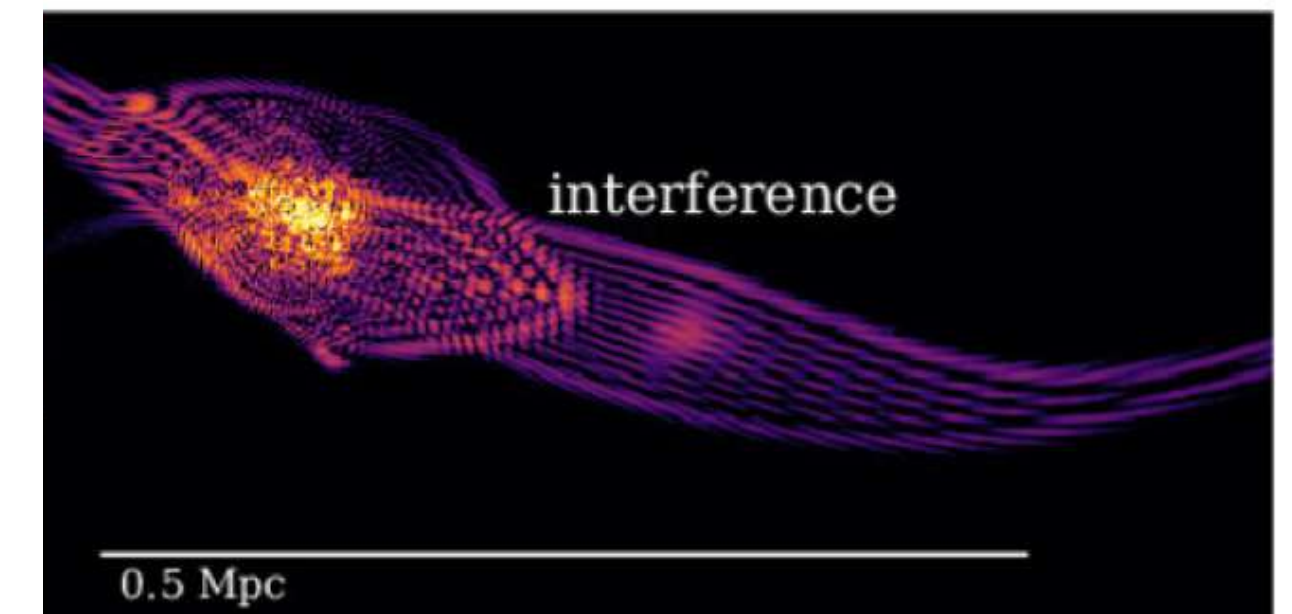
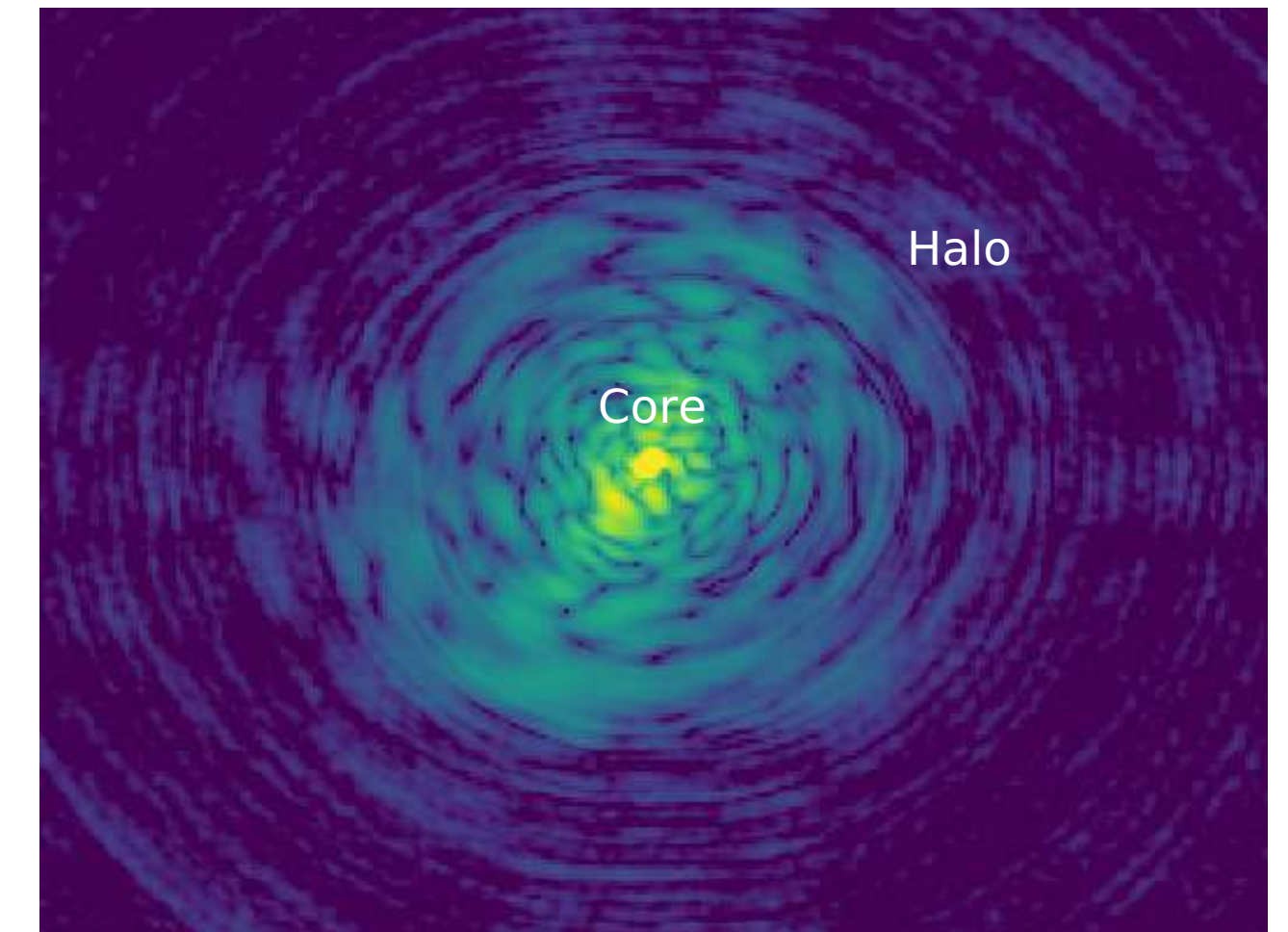
Mocz et al. 2017

Phenomenology

Wave interference: granules and vortices



Simulation by Jowett Chan



Mocz et al. 2017

Order one fluctuations in density \longrightarrow

Constructive interference: granules

Destructive interference

$\sim \lambda_{dB}$

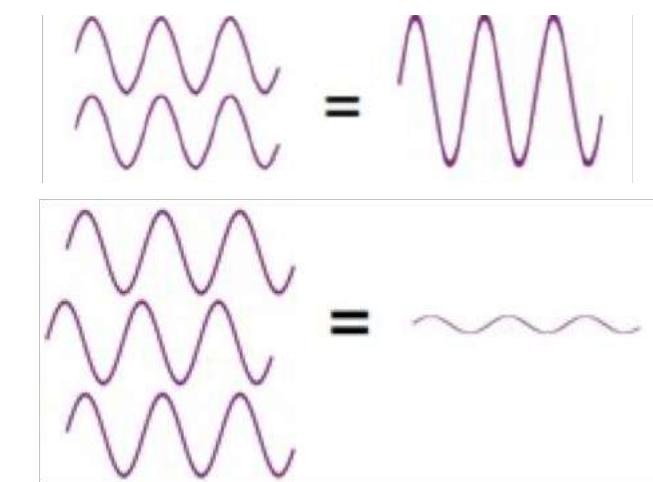
Hard to observe!

Vector, higher spin or multicomponent *FDM*

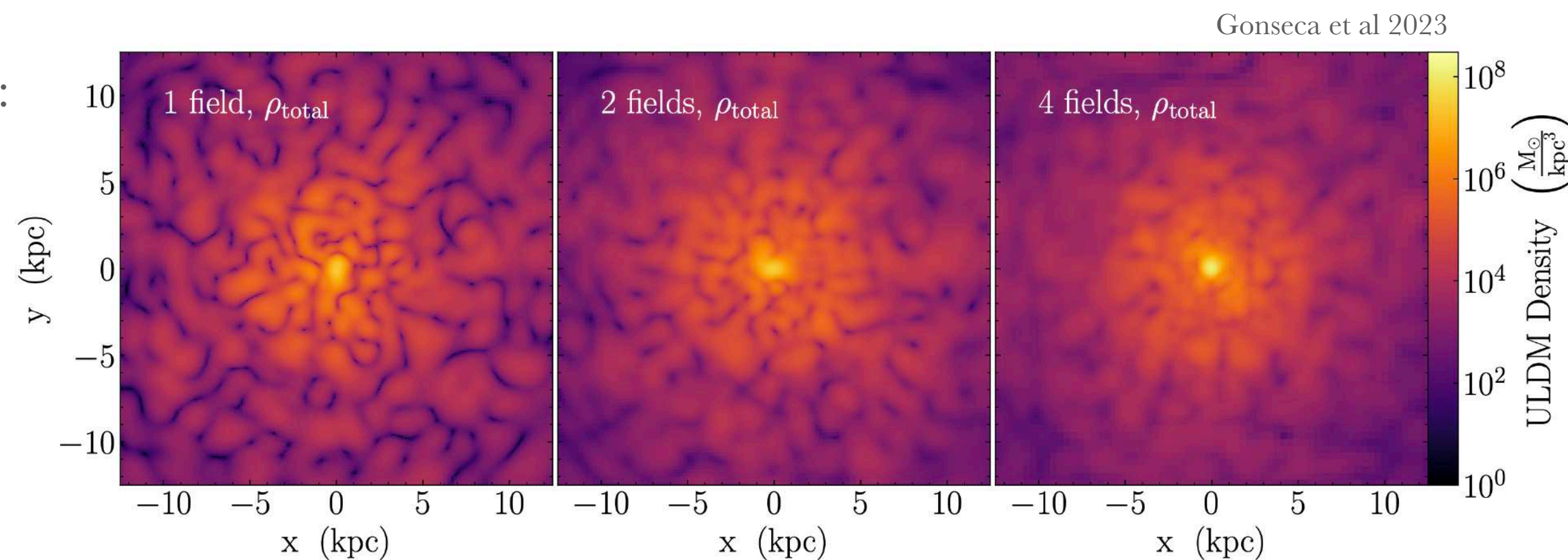
ULDM or ULA are a coherent wave - same frequency and constant phase difference

Multiple coherent waves

Interference patterns



For ULDM:



Multiple FDM or VFDM (or higher spin s FDM) *attenuates* the granule amplitude by

$$\frac{[\delta\rho/\rho]_{\text{nfdm},s}}{[\delta\rho/\rho]_{\text{fdm}}} \propto \frac{1}{\sqrt{(2s+1)}} = \frac{1}{\sqrt{N}}$$

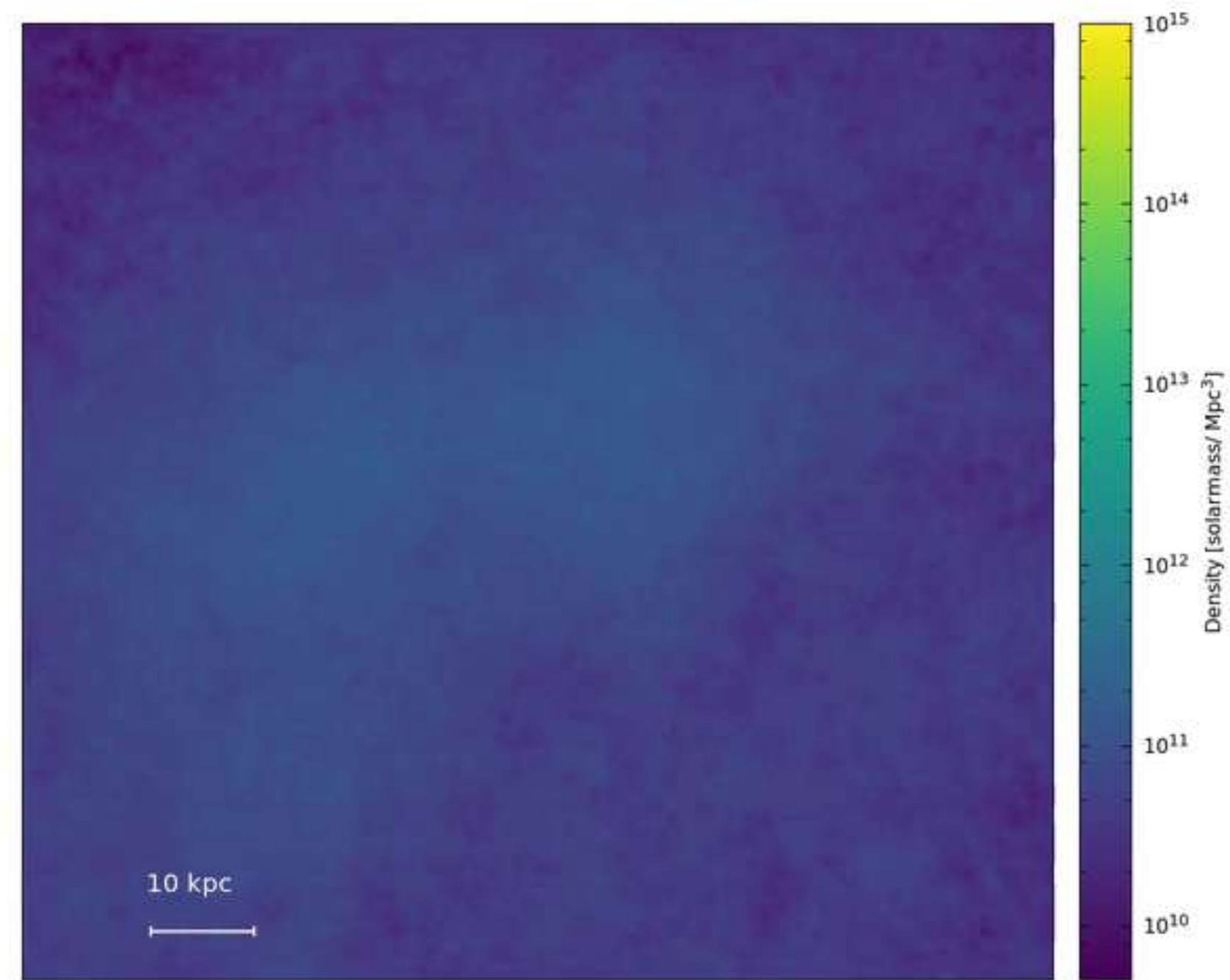
(Amin et al 2022)

Vector (and higher-spin) FDM Amin et al 2022

(Vector FDM = 3 x same mass FDM (spin 0))

Multicomponent FDM Gonseca et al 2023

Interference pattern

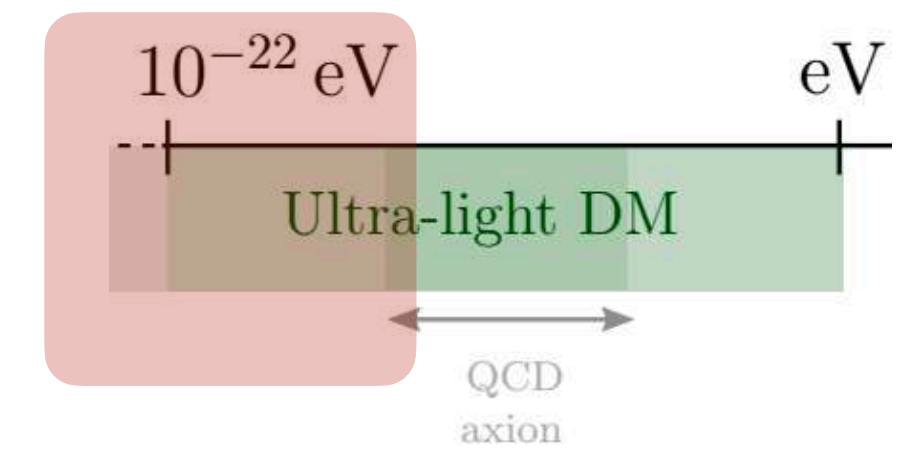


$\mathcal{O}(1)$ fluctuations in density $\rightarrow \sim \lambda_{\text{dB}}$

Best probe for ULDM

PROBES:

- Strong lensing
- Stellar streams
- Heating

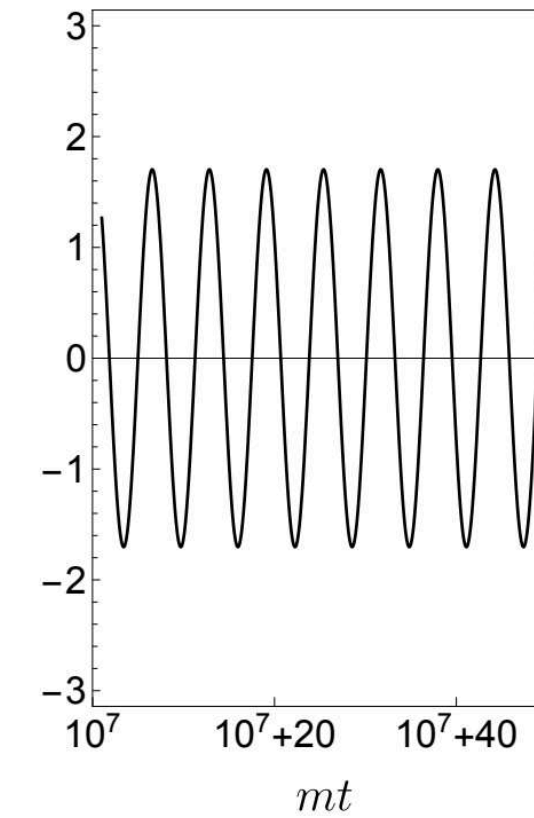


Modeling a *granular halo*

Coherent wave oscillation of ULDM

$$\phi(t, \vec{x}) = m^{-1} \sqrt{2\rho(t, \vec{x})} \cos[mt + \theta]$$

Fixed Constant
freq. phase



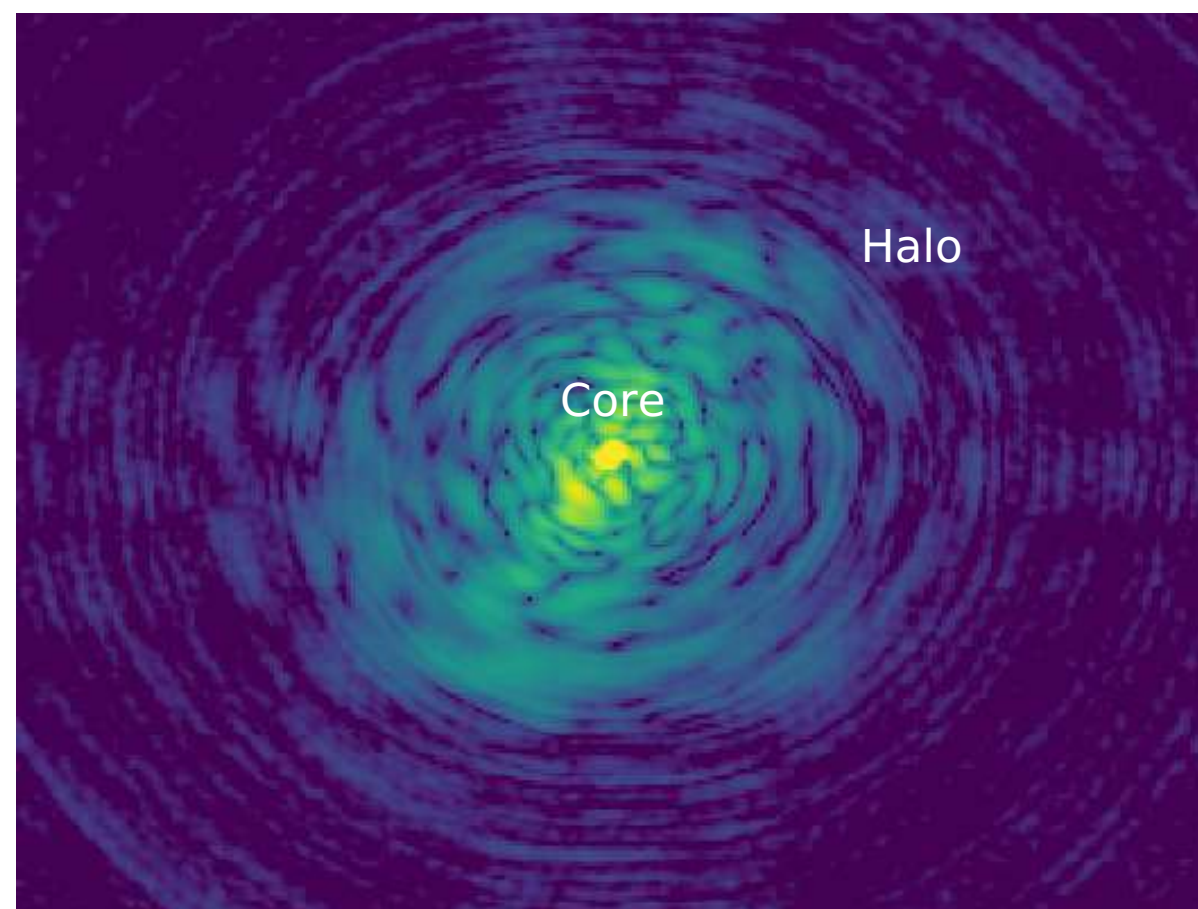
Modeling a *granular halo*

Coherent wave oscillation of ULDM

$$\phi(t, \vec{x}) = m^{-1} \sqrt{2\rho(t, \vec{x})} \cos[mt + \theta]$$

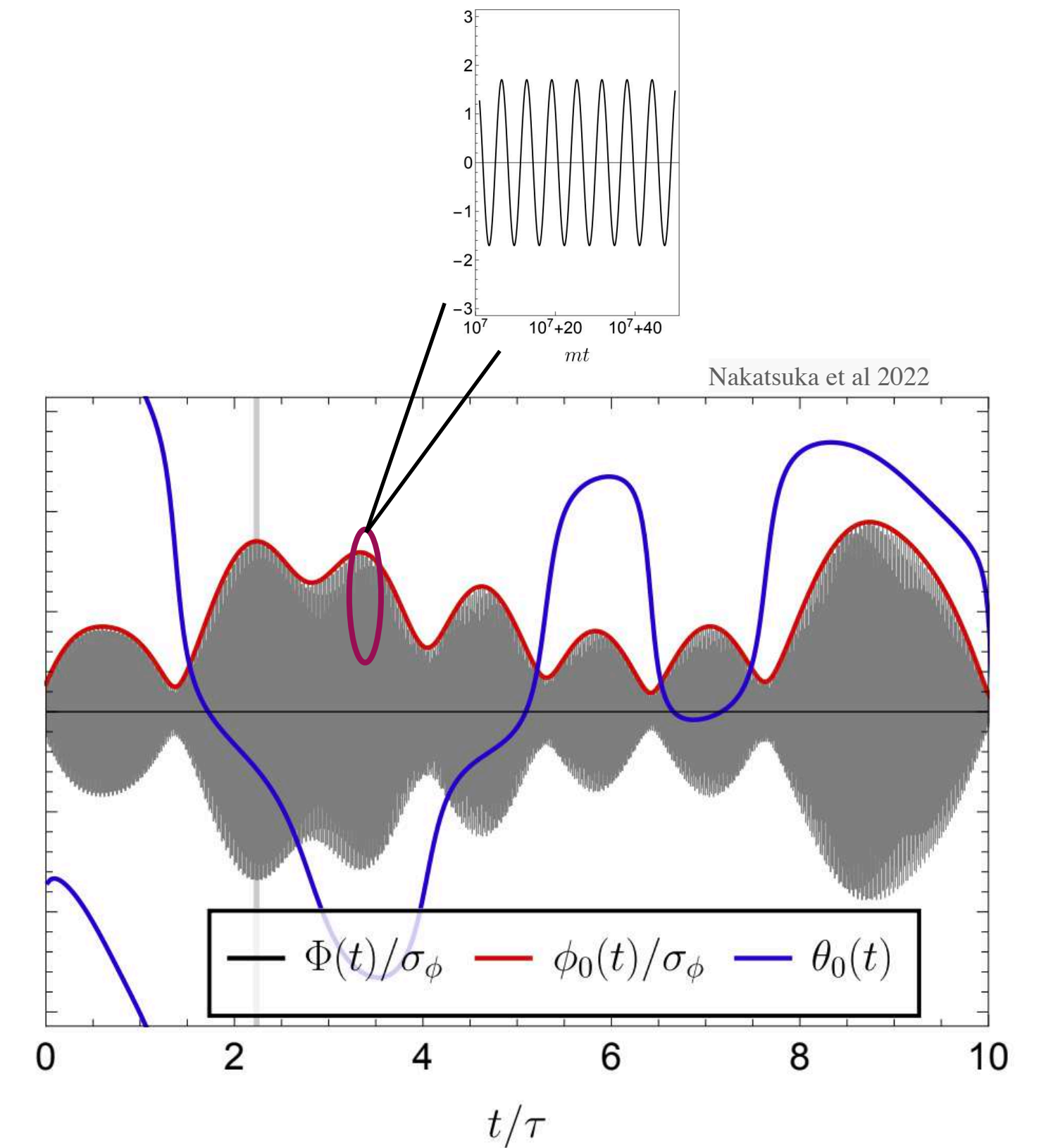
Fixed
freq. Constant
phase

But, the halo in these models is like this:



Superposition of plane waves

$$\left| \begin{matrix} \text{wavy line} \\ \text{wavy line} \end{matrix} \right| = \left| \text{smooth wave} \right|$$



Modeling a *granular halo*

Full SP simulations can describe perfectly this interference pattern (while fluid ones *cannot* describe it)

OR

We can adopt simpler descriptions of the galactic halo to describe this effect.

1) A simple model of a galactic halo, consider a **superposition of plane waves**:

$$\psi(t, \vec{x}) = \sum_{\vec{k}} A_{\vec{k}} e^{iB_{\vec{k}}} e^{i\vec{k} \cdot \vec{x} - i\omega_k t}$$

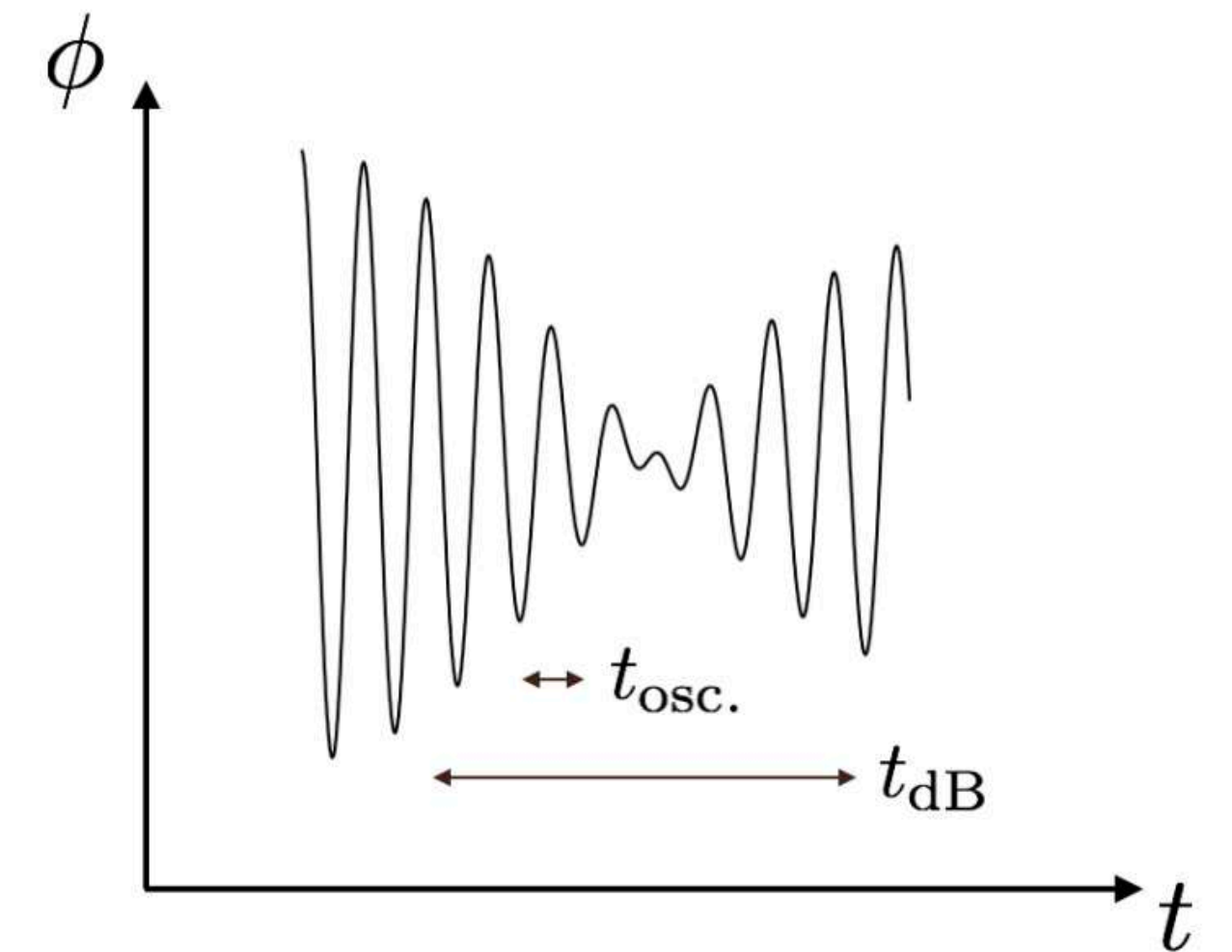
Randomly distributed

Wave interference produces de-Broglie-scale, order unity density fluctuations which vary on time scale of t_{dB}

This collection of plane waves can also be represented like this:

$$\phi(t, \vec{x}) = A(\vec{x}) \cos(mt + \alpha(\vec{x}))$$

describes the interference patterns



$$t_{osc.} = 2\pi/m$$

$$t_{dB} = 2\pi/(mv^2)$$

$$= 1.9 \times 10^6 \text{ yr} \left(\frac{10^{-22} \text{ eV}}{m} \right) \left(\frac{250 \text{ km/s}}{v} \right)^2$$

Modeling a *granular halo*

Full SP simulations can describe perfectly this interference pattern (while fluid ones *cannot* describe it)

OR

We can adopt simpler descriptions of the galactic halo to describe this effect.

1) A simple model of a galactic halo, consider a **superposition of plane waves**:

2) A more realistic model would **superimpose eigenstates** of a desired gravitational potential (Lin et al. 2018, Li et al. 2021)

Perform an eigenmode decomposition of the halo wavefunction, where the eigenmodes are for a fixed gravitational potential
→ ω_k is the energy of each eigenmode (labeled abstractly by k), with $e^{i\vec{k}\cdot\vec{x}}$ replaced by the corresponding eigenfunction.

$$\psi(r, \theta, \phi, t) = \sum_{n,l,m} A_{nlm} F_{nlm}(r, \theta, \phi) e^{-iE_{nl}t/\hbar}$$

Energy eigenvalue

$$F_{nlm}(r, \theta, \phi) = R_{nl}(r) Y_l^m(\theta, \phi)$$

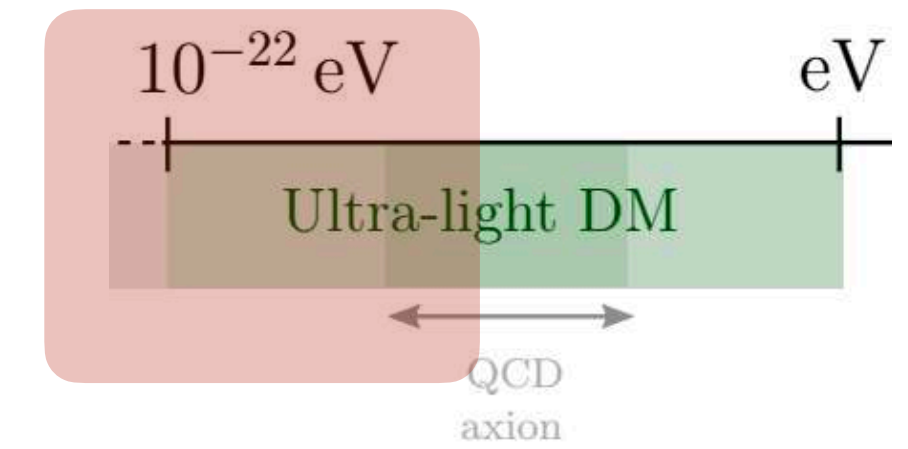
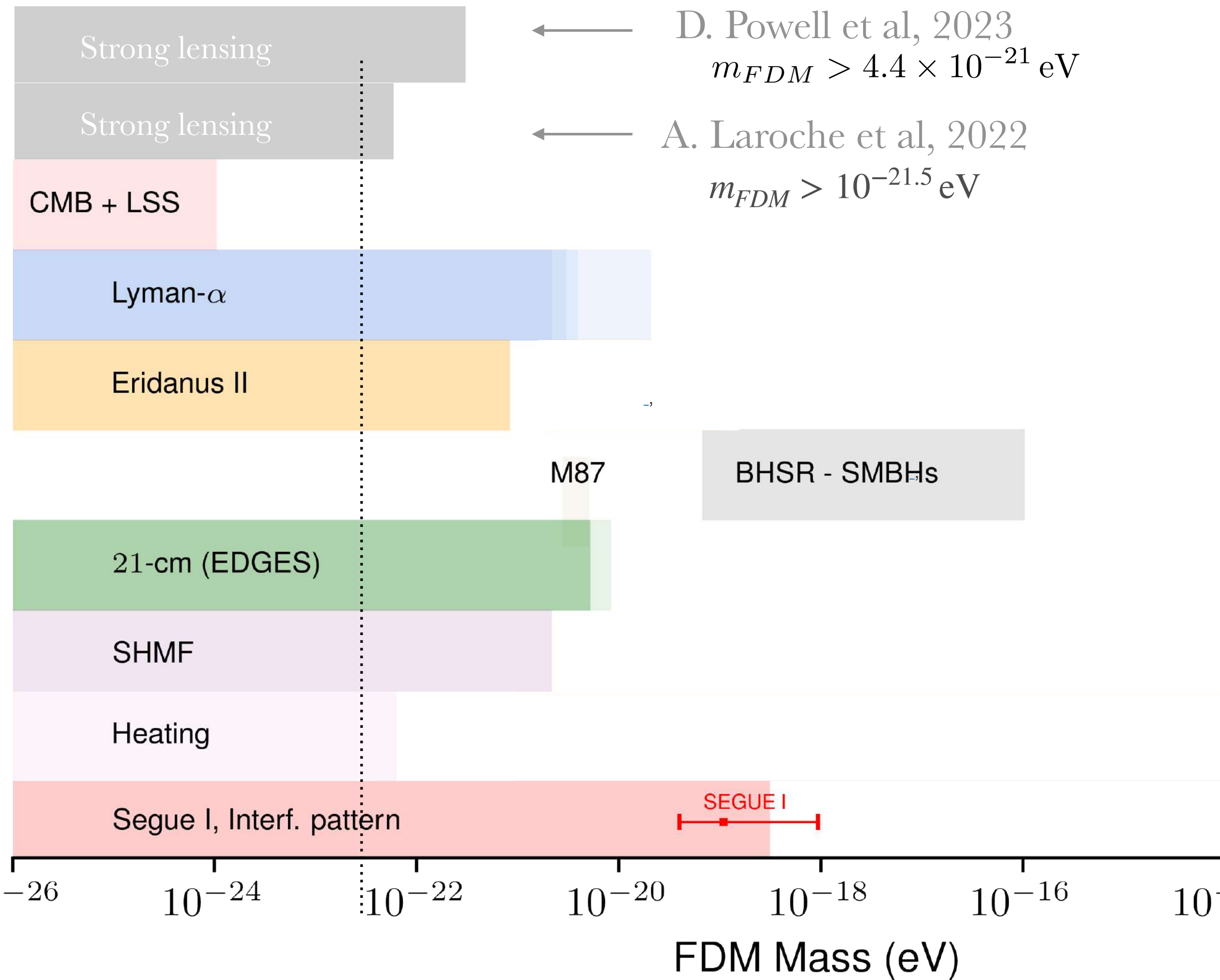
Radial eigenfunction **Spherical harmonics**

→ energy eigenmodes of the gravitational potential of the virialized halo

Interference patterns - granules

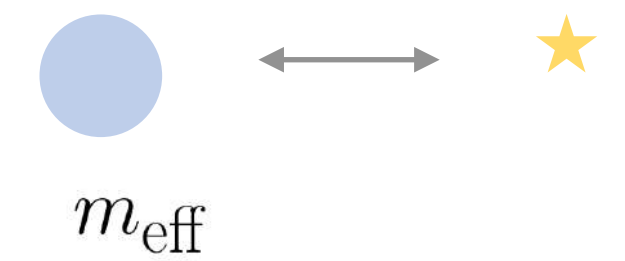
Strong lensing

Constraints 2.0

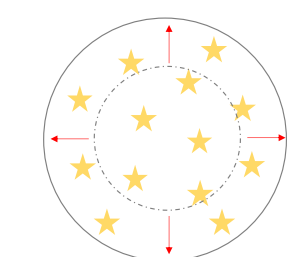


Heating

FDM granule

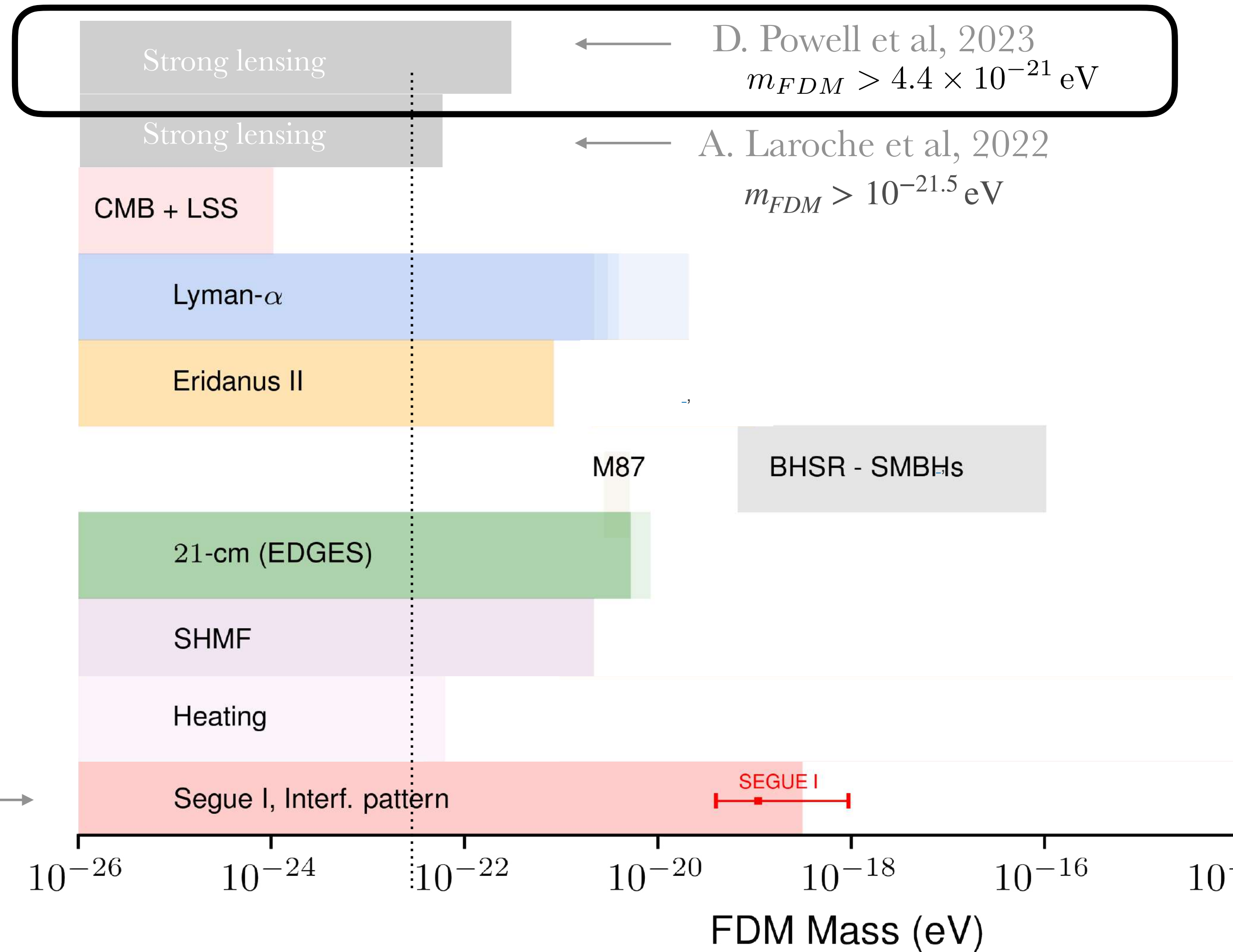
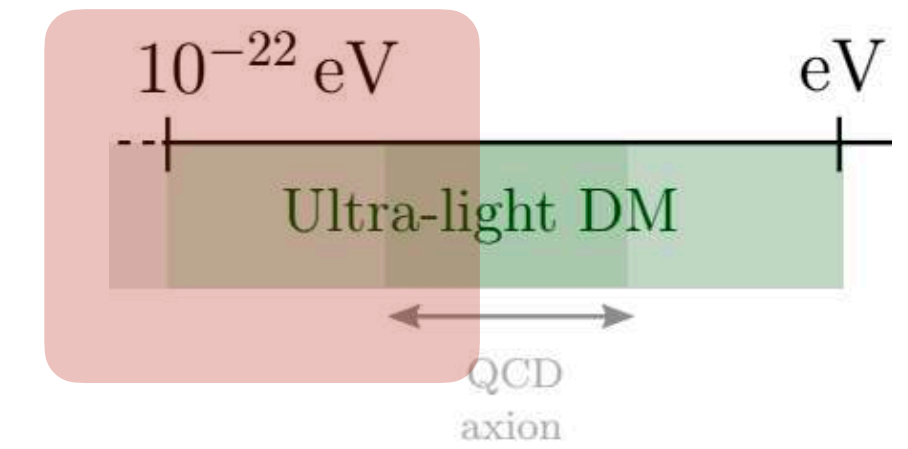


System (star) gains energy



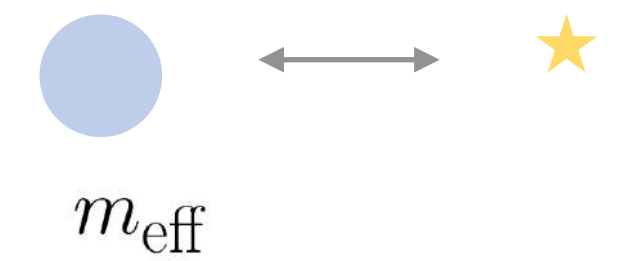
Interference patterns - granules

Strong lensing

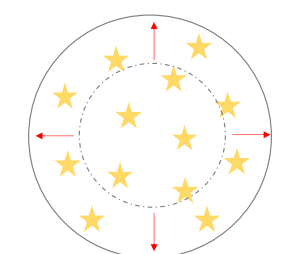


Heating

FDM granule

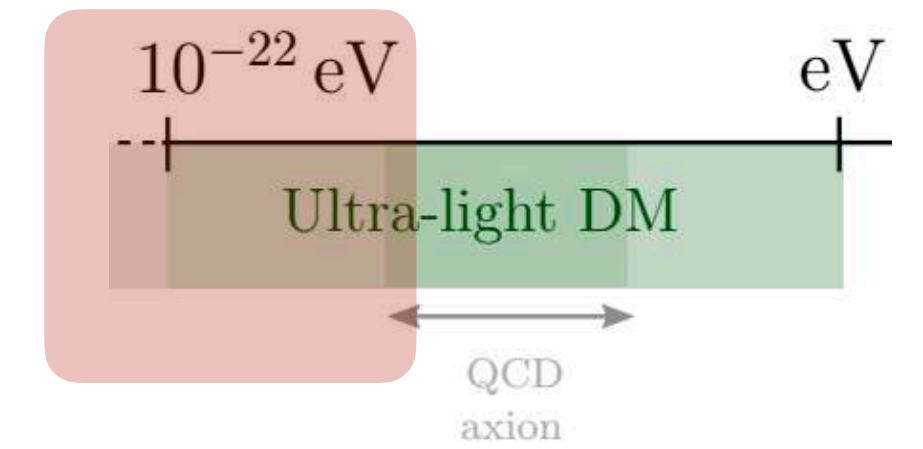
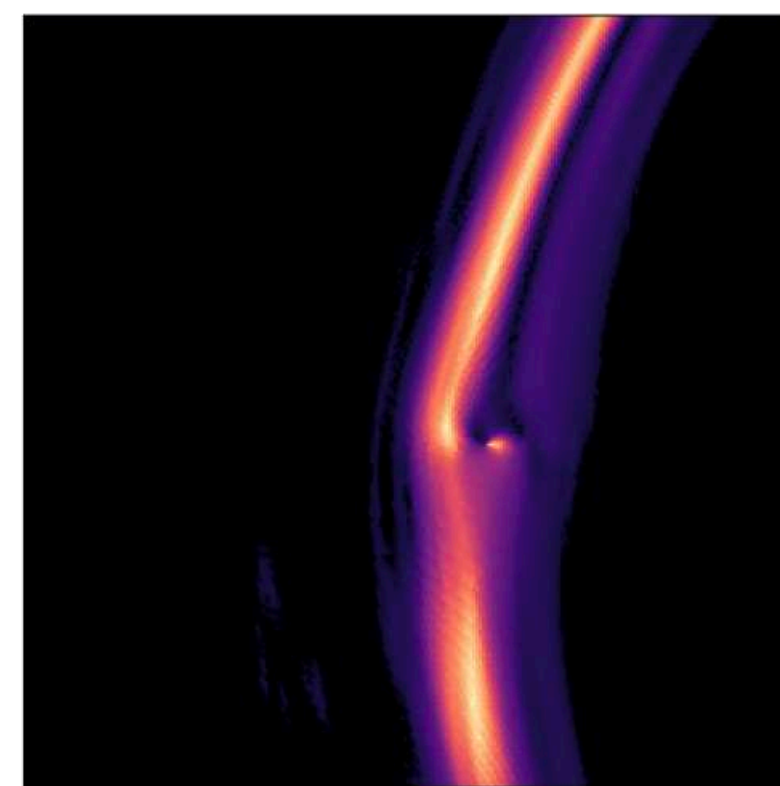
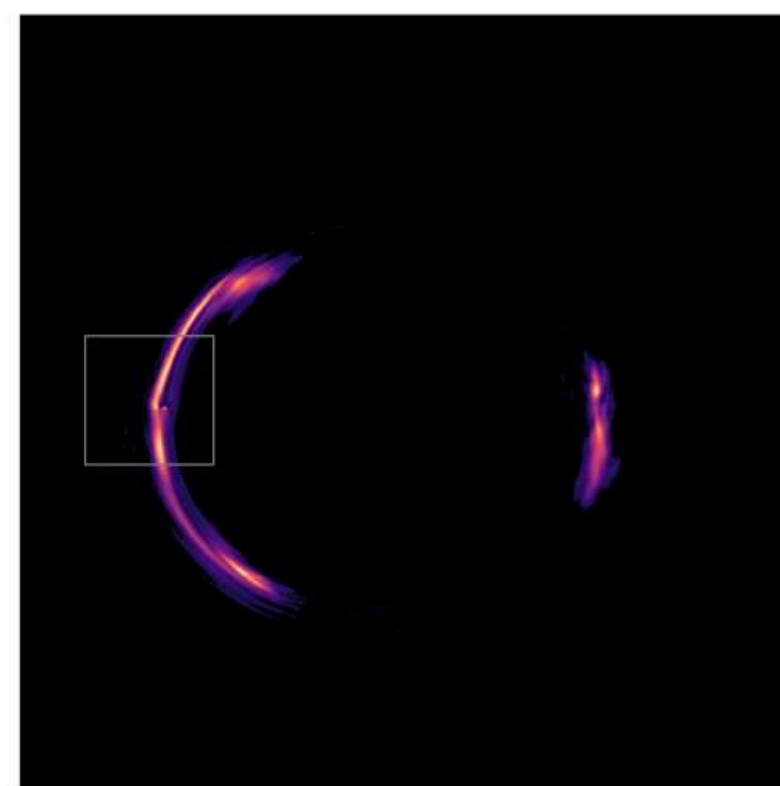
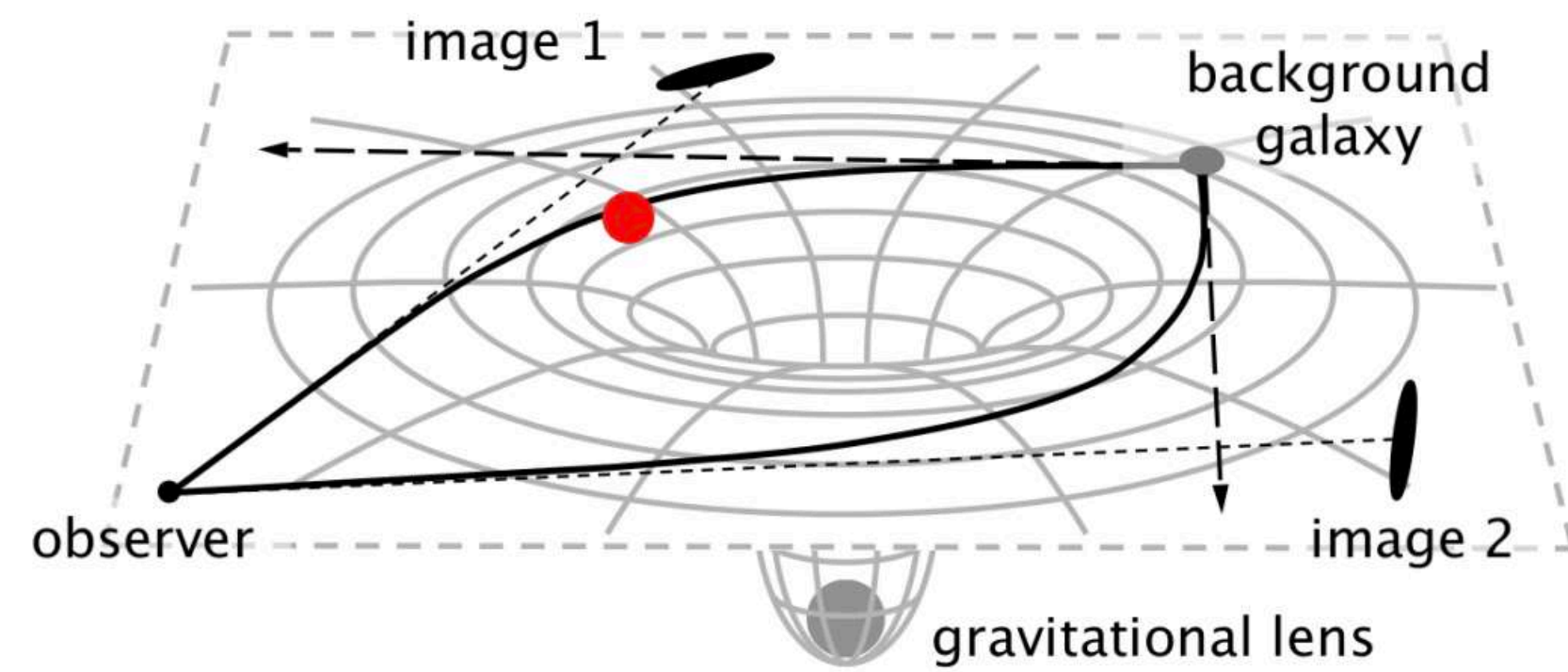


System (star) gains energy



Strong *lensing*

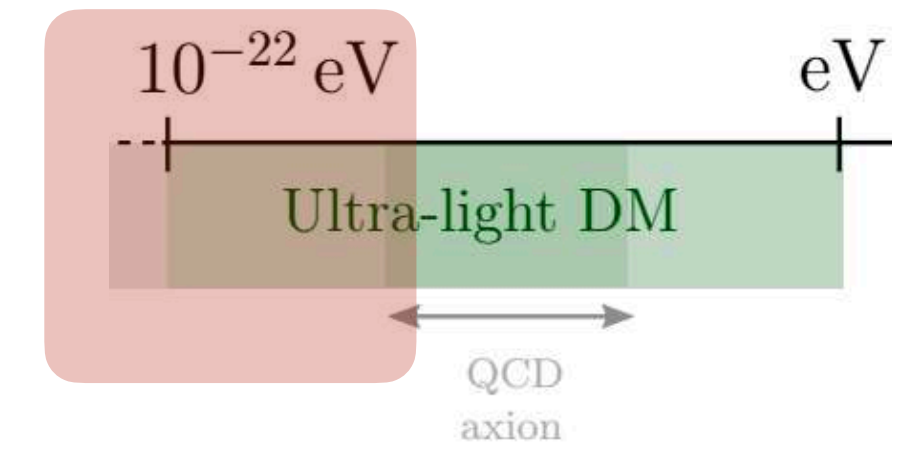
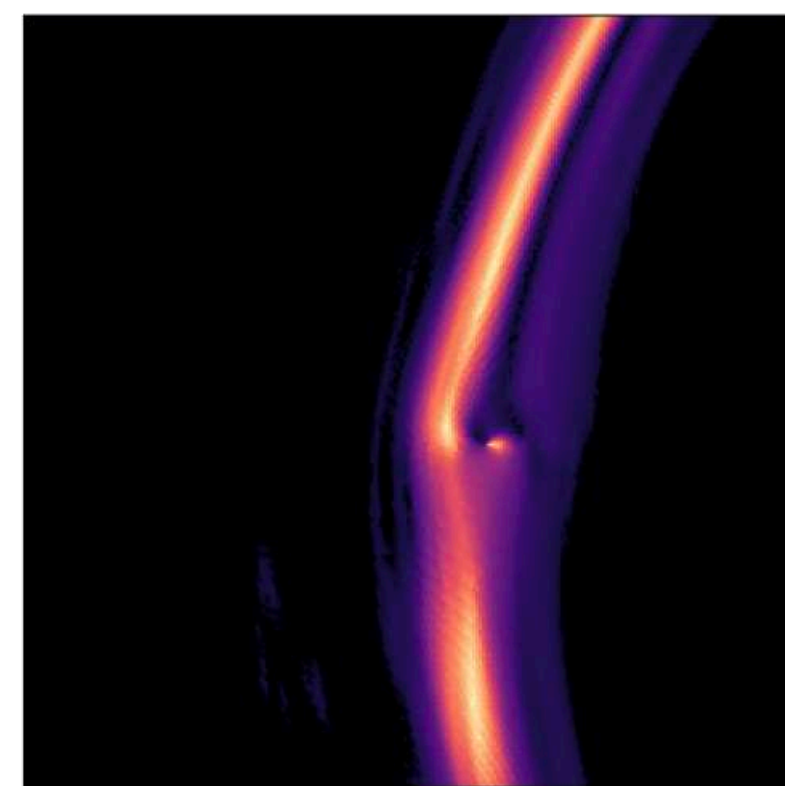
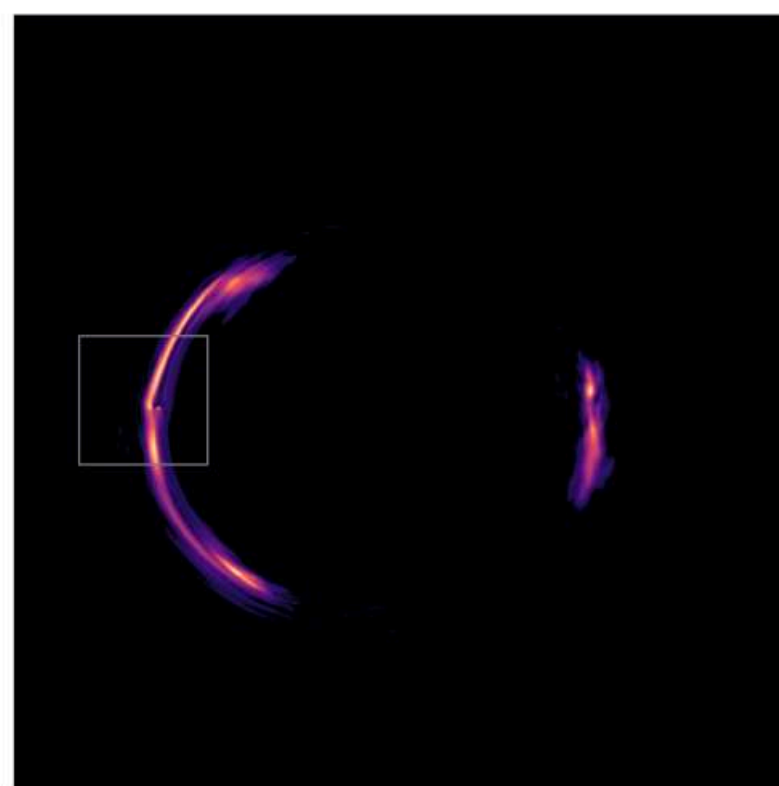
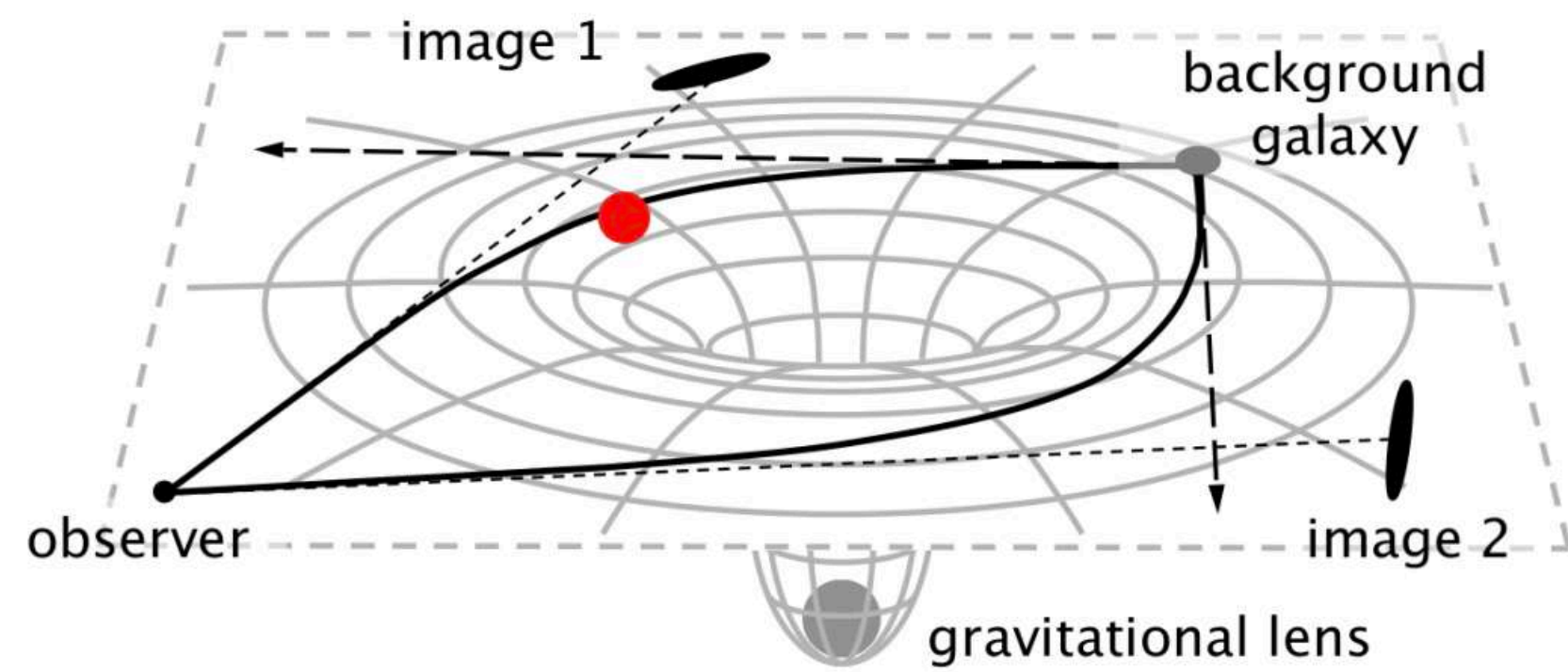
Low mass perturber with lensing



- Strong lensing: powerful probe of substructure
- Sensitivity is limited by angular **angular resolution**
- Roughly speaking, the resolution must be better than the scale radius of the perturber

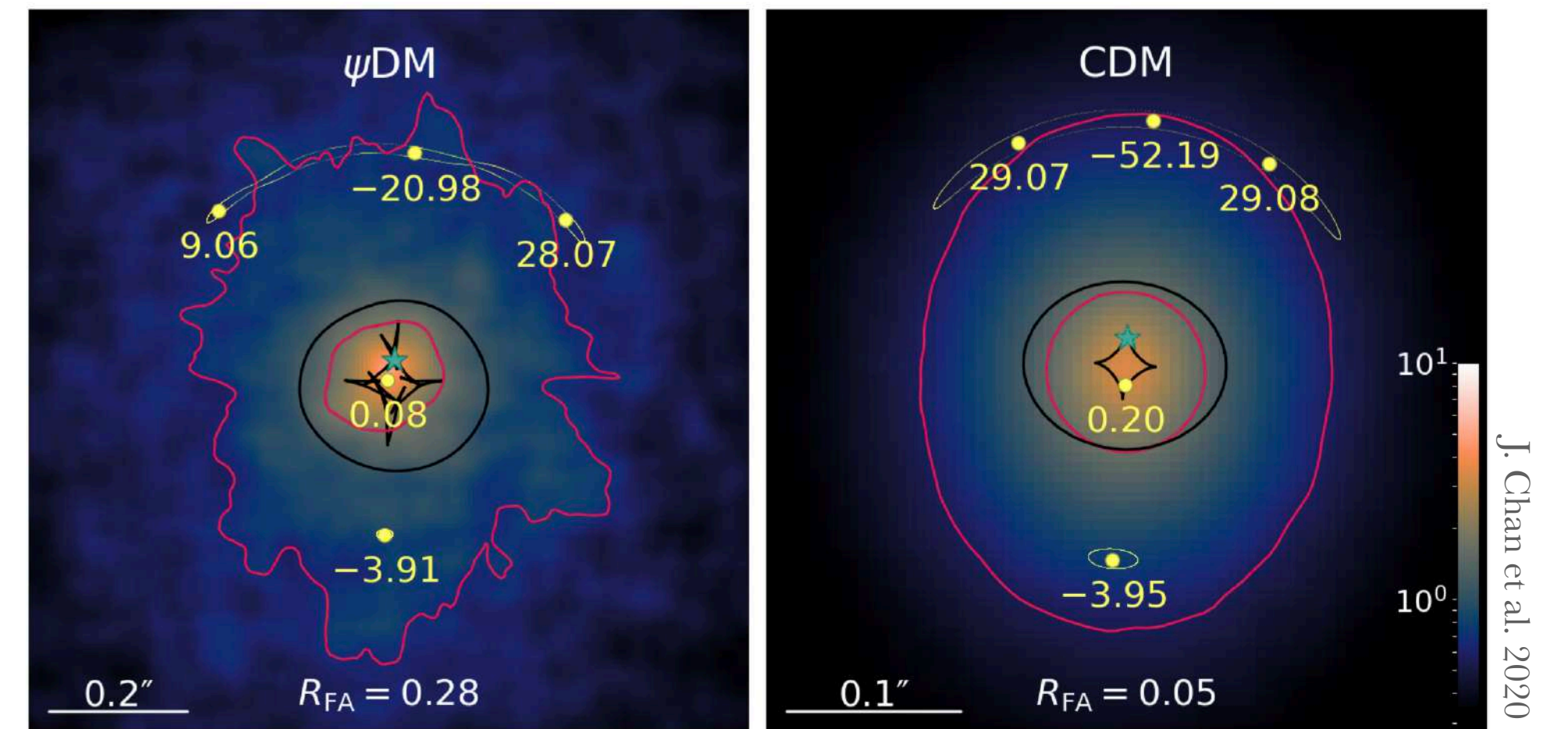
Strong *lensing*

Low mass perturber with lensing



Presence of granules

Surface densities overlaid with sources and quad images for fuzzy and smooth lenses



Fuzzy lens: fluctuating tangential critical curve; flux ratio anomalies also sizable.

Previous works:

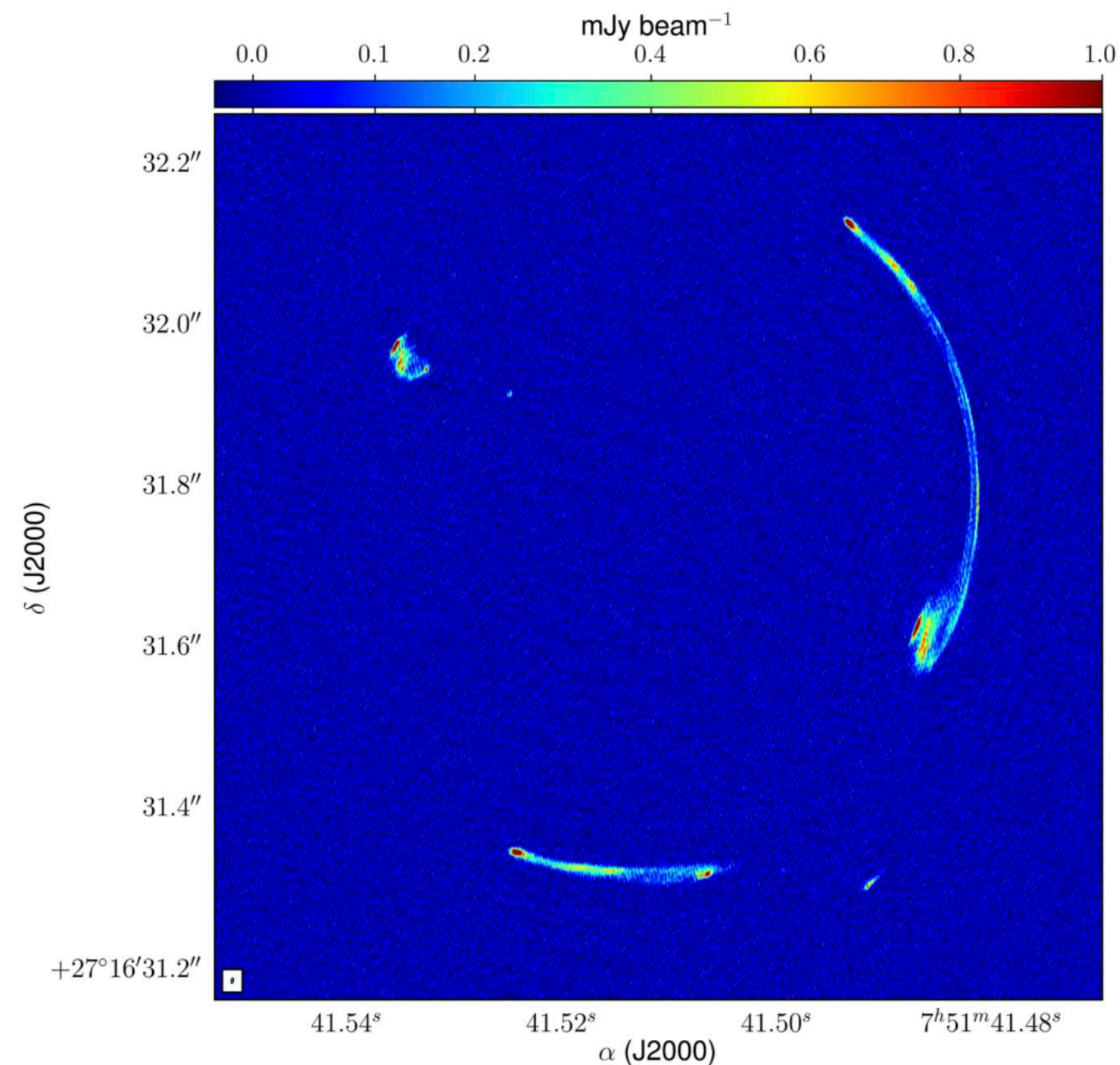
- J. Chan, H. Schive, S.g. Wong, T. Chiueh, T. Broadhurst, 2020
- A. Laroche, Daniel Gilman, X. Li, J. Bovy, X. Du, 2022

Strong *lensing*

A lensed radio jet at milli-arcsecond resolution II: Constraints on fuzzy dark matter from an extended gravitational arc

D. Powell, S. Vegetti, J.P. McKean, S. White, EF, S. May, C. Spingola

MG J0751+2716



- Lensed radio jet, observed with global VLBI
- First image of a lensed radio jet!
- Source structure allows us to “image” the lens surface density
- Extended lensed radio arcs and the milli-arcsecond resolution provide direct sensitivity to the presence of **FDM granules** in the halo of the lens galaxy

Bayesian approach to jointly inferring the lens mass model and source surface brightness distribution

Data taken at 1.6 GHz using global very long baseline interferometry (VLBI) with an angular resolution, measured as the full width at half maximum (FWHM) of the main lobe of the dirty beam response, of $5.5 \times 1.8 \text{ mas}^2$

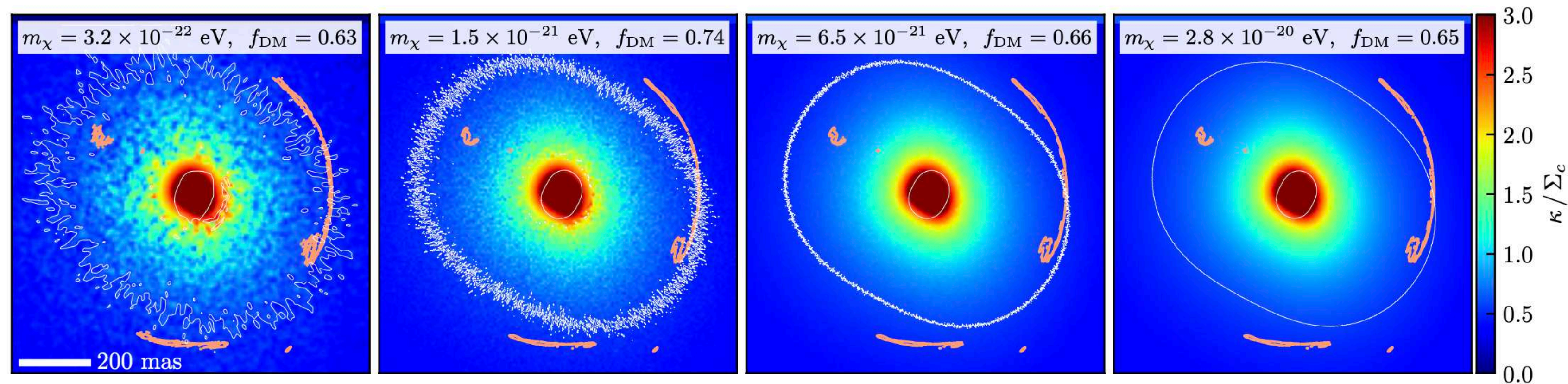
(Suyu et al. 2006; Vegetti & Koopmans 2009; Hezaveh et al. 2016; Rizzo et al. 2018)

Strong *lensing*

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D. Powell, S. Vegetti, J.P. McKean, S. White, EF, S. May, C. Spingola

Example convergence maps with corresponding MAP surface mass density maps (κ , in units of the critical density Σ_c) reconstruction for 4 random realizations of MG J0751+2716 in an FDM cosmology - the model lensed images in orange contours



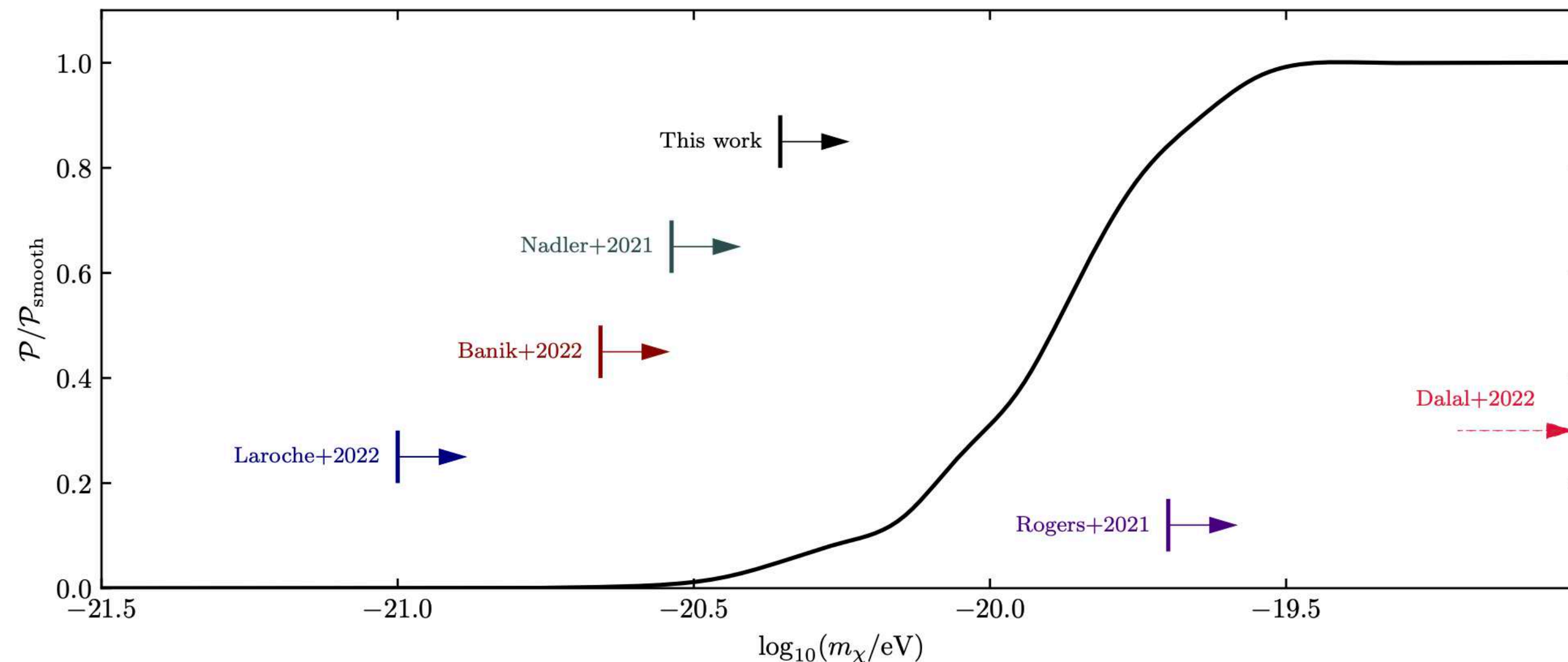
The lensing effect of the FDM granules is apparent: The critical curves wiggle back and forth across the lensed arcs, which would require the presence of multiple images of the same region of the source along the arc.

Strong lensing

A lensed radio jet at milli-arcsecond resolution II: Constraints on fuzzy dark matter from an extended gravitational arc

D. Powell, S. Vegetti, J.P. McKean, S. White, EF, S. May, C. Spingola

Results quoted in terms of posterior odds ratio (POR) between FDM with a particle mass m_{fdm} and the smooth model, $\mathcal{P}/\mathcal{P}_{smooth}$



Fuzzy dark matter
(Single spin-0 particle)

$$m_{fdm} > 4.4 \times 10^{-21} \text{ eV}$$

Vector fuzzy dark matter
(spin-1 particle)
OR 3 same mass FDM

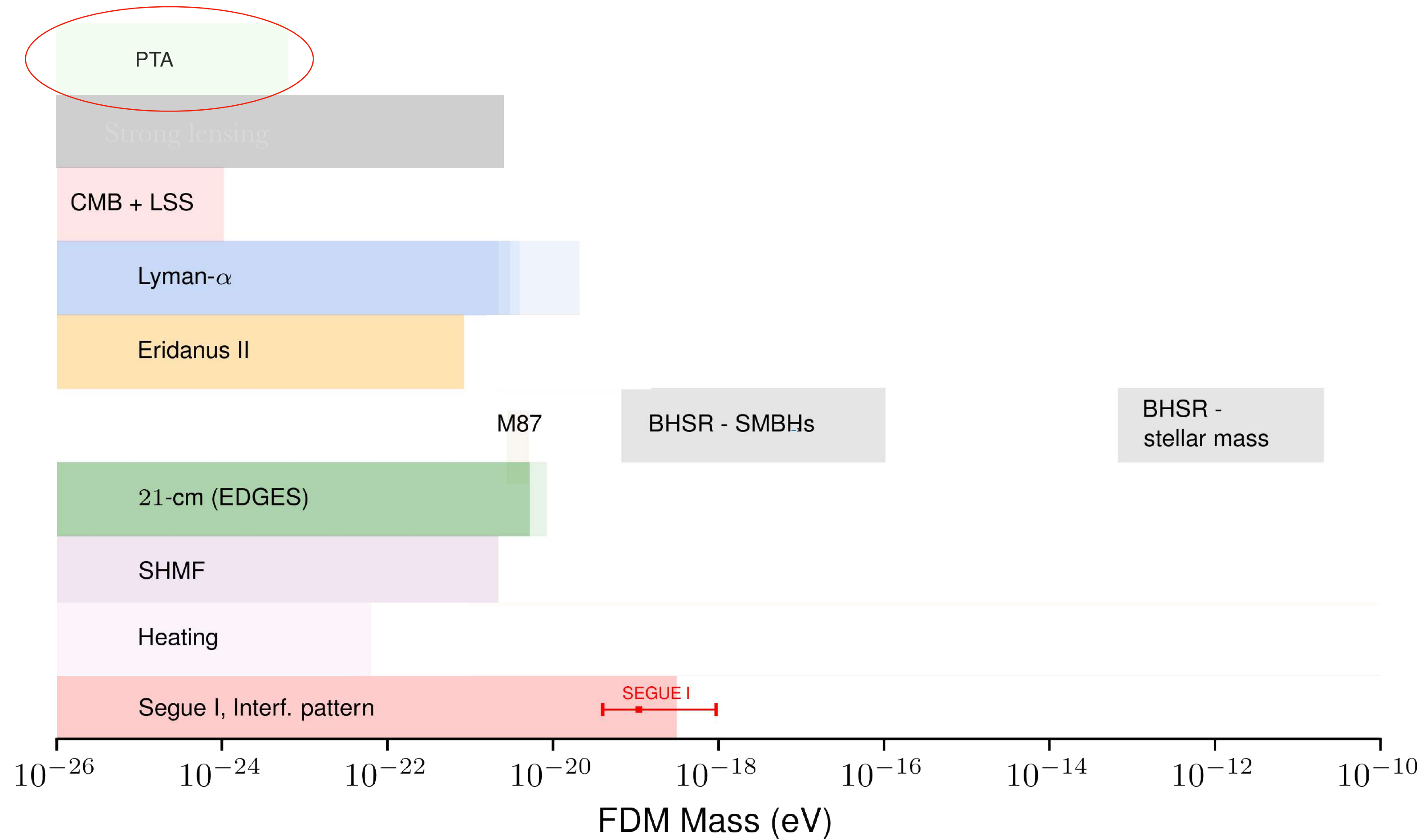
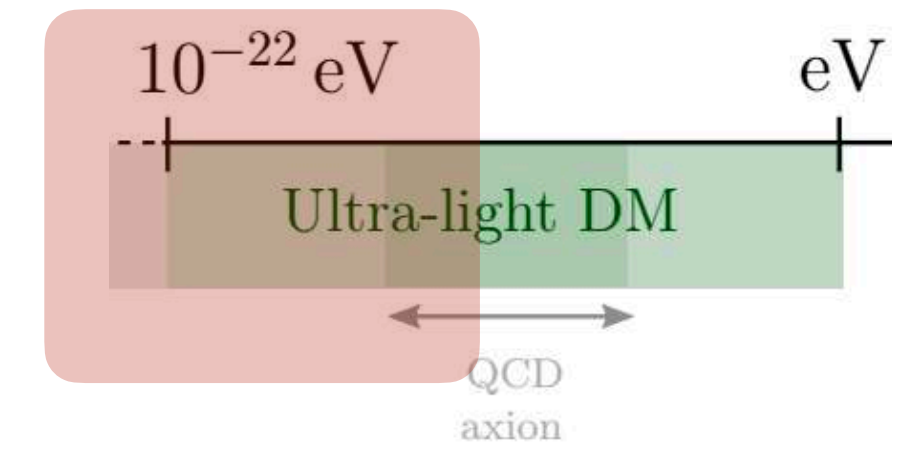
$$m_{vdm} > 1.4 \times 10^{-21} \text{ eV}$$

Spin-2 FDM

$$m_{spin-2} > 8.8 \times 10^{-22} \text{ eV}$$

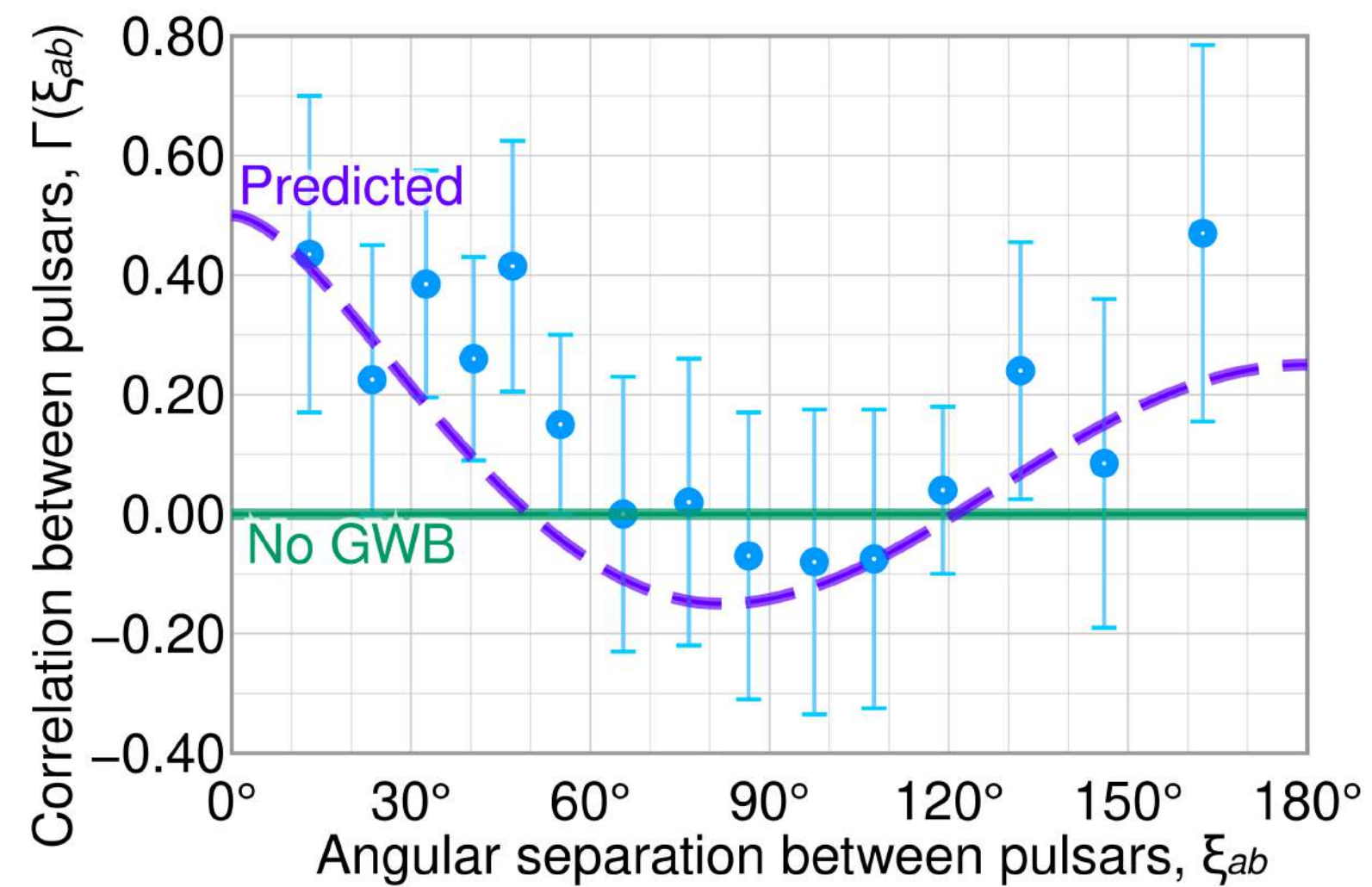
Current status

Fuzzy Dark Matter - bounds on the mass

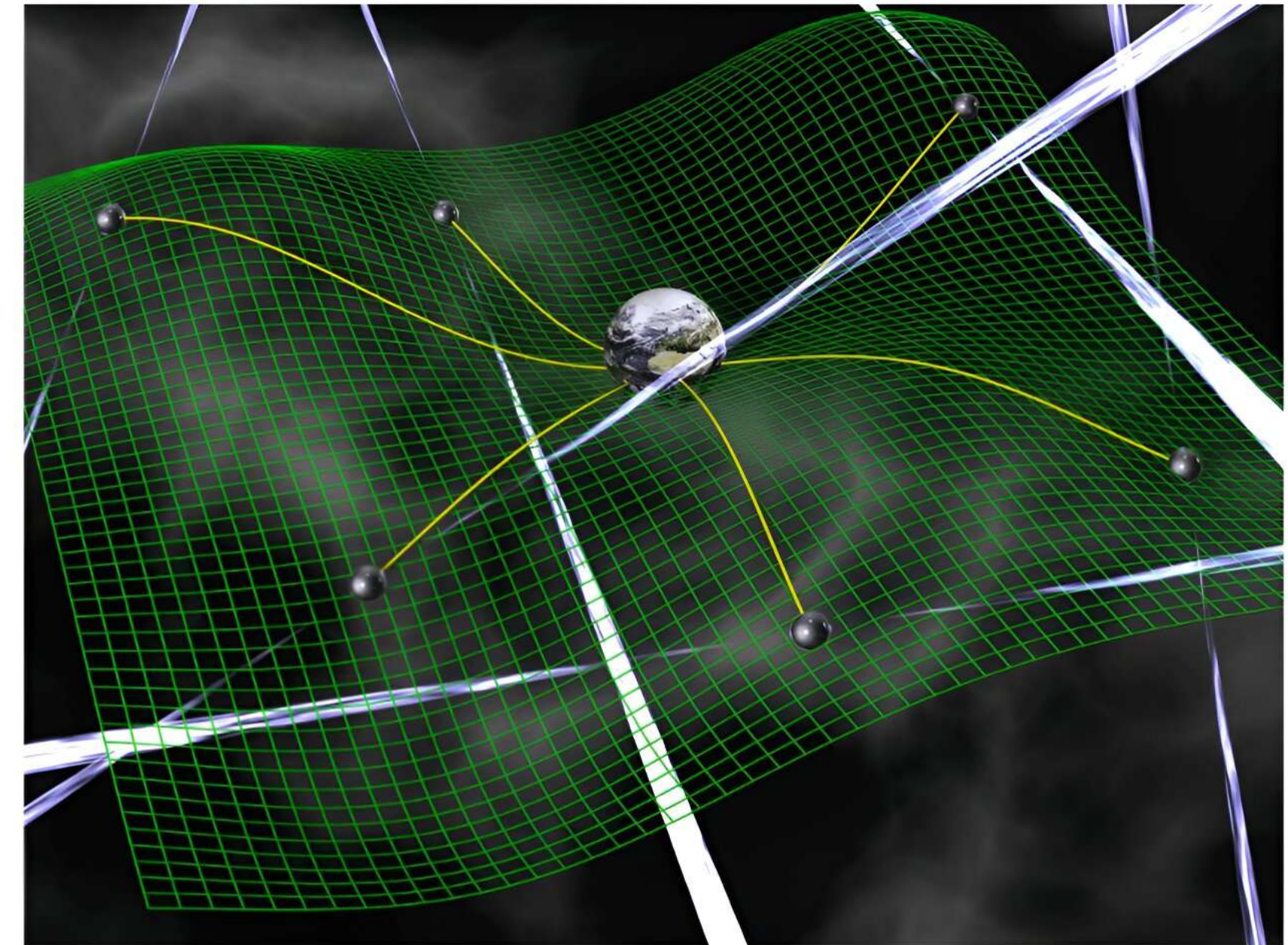


Pulsar Timing array

NANOGrav, EPTA, PPTA, and InPTA announced that they found evidence for a gravitational wave background



EPTA - 3σ
CPTA - 4.6σ

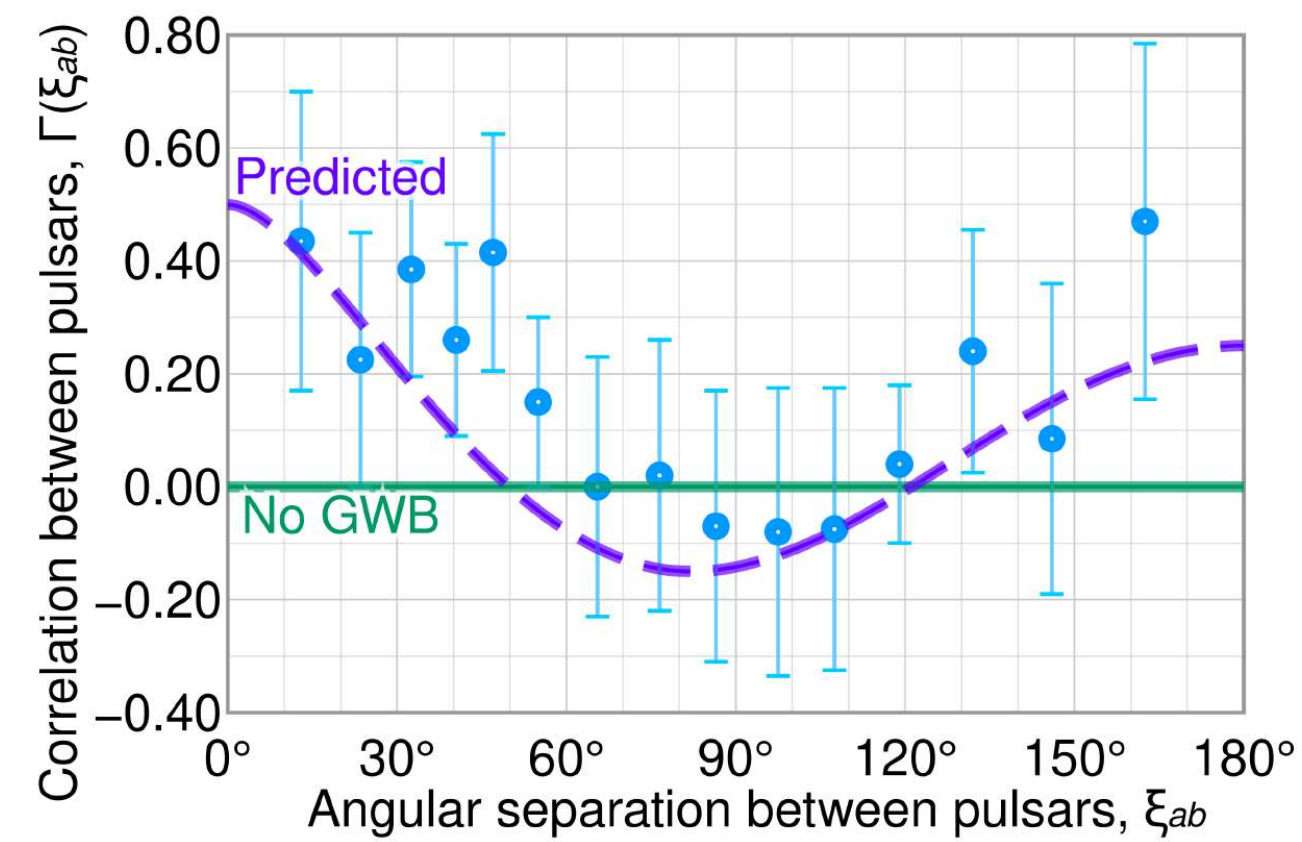


Artistic representation of the Pulsar Timing Array concept. Credit: David Champion / MPIfR

Hellings-Downs curve refers to the wave-like shape predicted to appear in a plot of timing residual correlations versus the angle of separation between pairs of pulsars

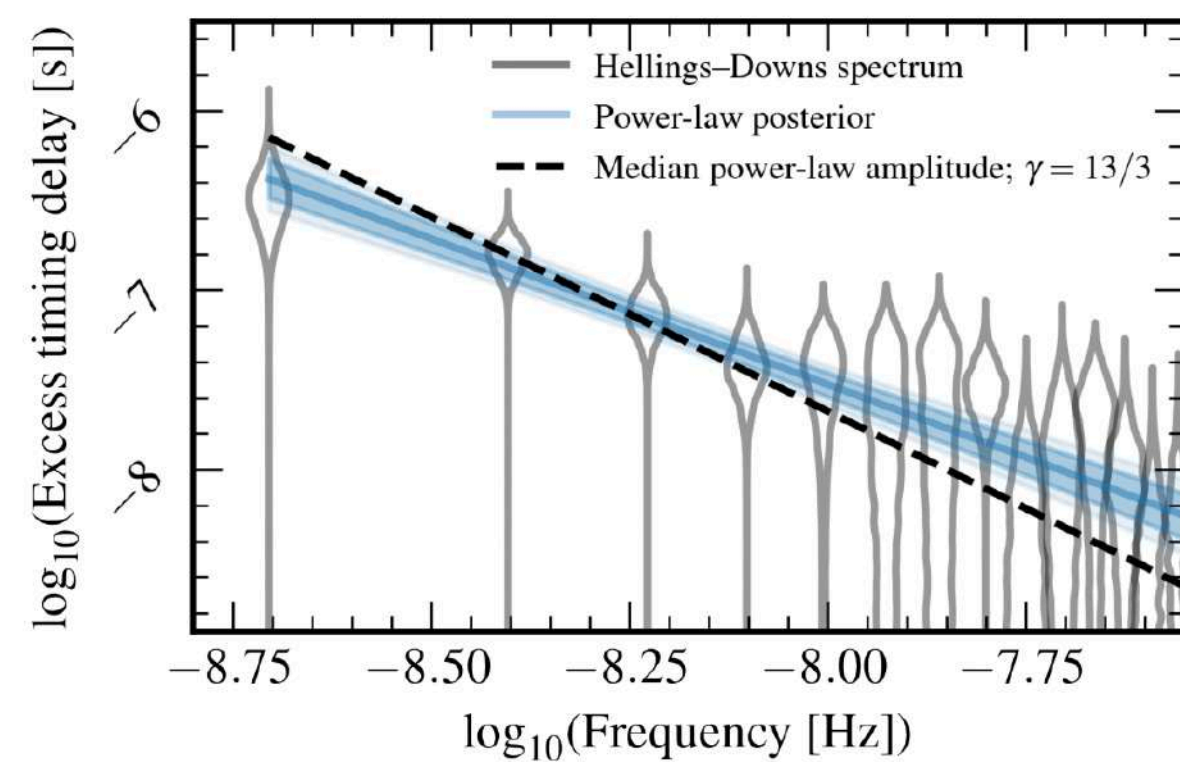
Pulsar Timing array

NANOGrav, EPTA, PPTA, and InPTA announced that they found evidence for a gravitational wave background



Possibility of being a **Gravitational Wave Background** signal caused by:

- Ensemble of binary supermassive black-holes
- New physics



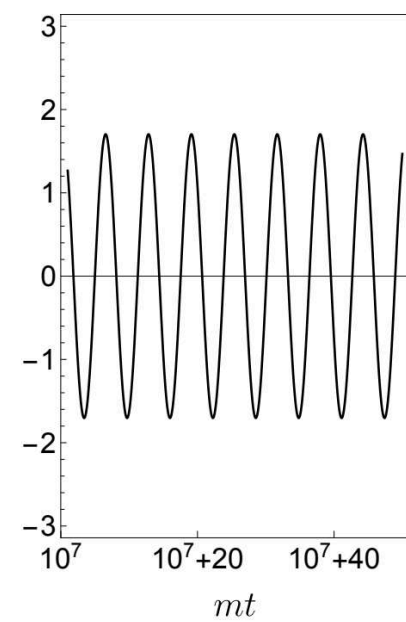
Pulsar Timing *array*

The presence of ULDM induces an oscillating gravitational potential that affects the light travel time of radio pulses emitted by pulsars. PTAs can be used to test the presence of ULDM particles in the MW \rightarrow how ULDM affects the SGWB in PTA system

Ultra-light DM: Coherent wave oscillation of ULDM

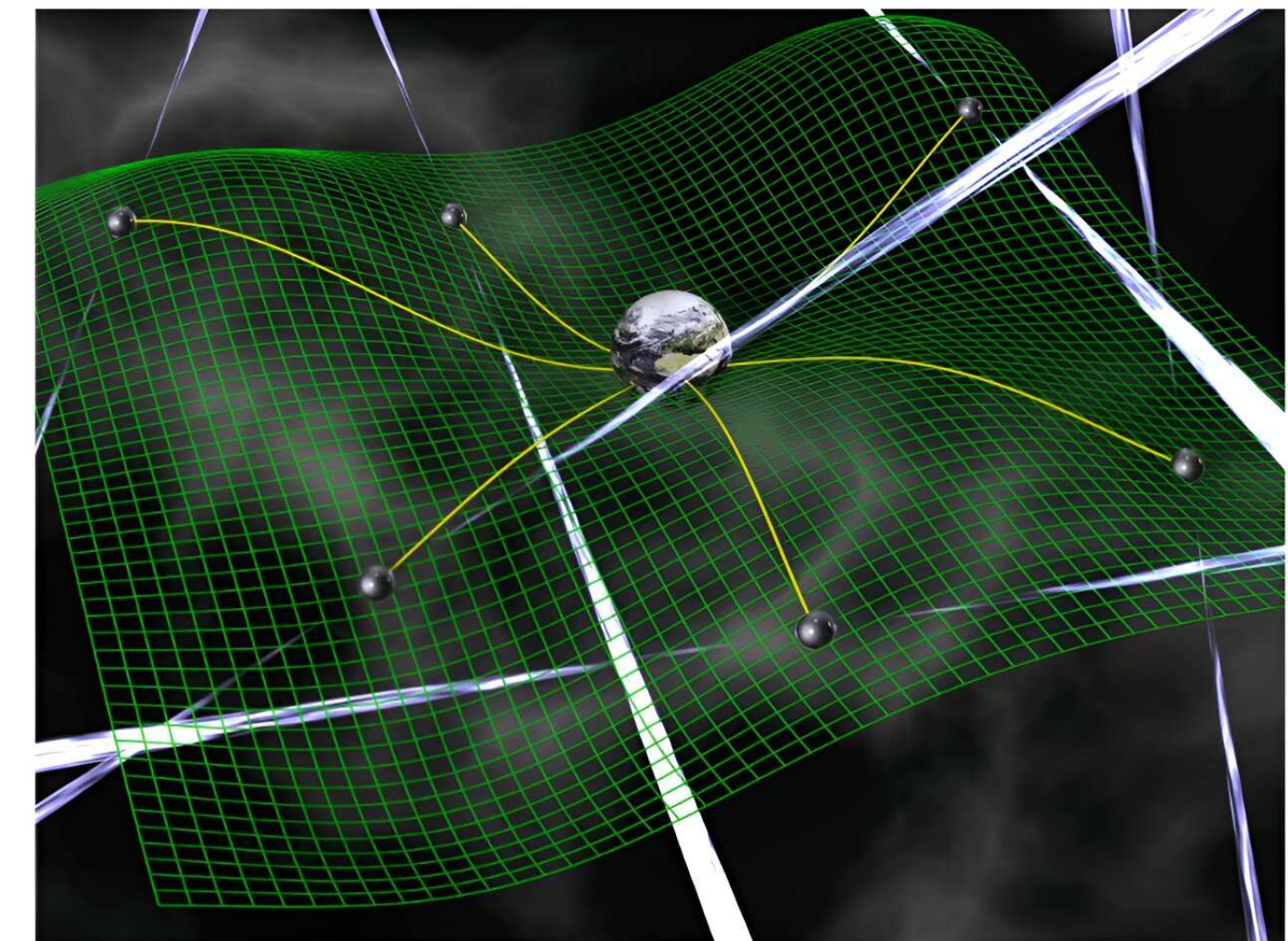
"Second Data Release from the European Pulsar Timing Array: Challenging the Ultralight Dark Matter Paradigm"

Clemente Smarra *et al.* (European Pulsar Timing Array), Phys. Rev. Lett. **131**, 171001



$$\phi(t, \vec{x}) = m^{-1} \sqrt{2\rho(t, \vec{x})} \cos[mt + \theta(t, \vec{x})]$$

Sources an oscillating gravitational potential



Artistic representation of the Pulsar Timing Array concept. Credit: David Champion / MPIfR

Pulsar Timing *array*

"Second Data Release from the European Pulsar Timing Array: Challenging the Ultralight Dark Matter Paradigm"
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The presence of ULDM induces an oscillating gravitational potential that affects the light travel time of radio pulses emitted by pulsars. PTAs can be used to test the presence of ULDM particles in the MW \rightarrow how ULDM affects the SGWB in PTA system

$$\phi(t, \vec{x}) = m^{-1} \sqrt{2\rho(t, \vec{x})} \cos[mt + \theta(t, \vec{x})]$$

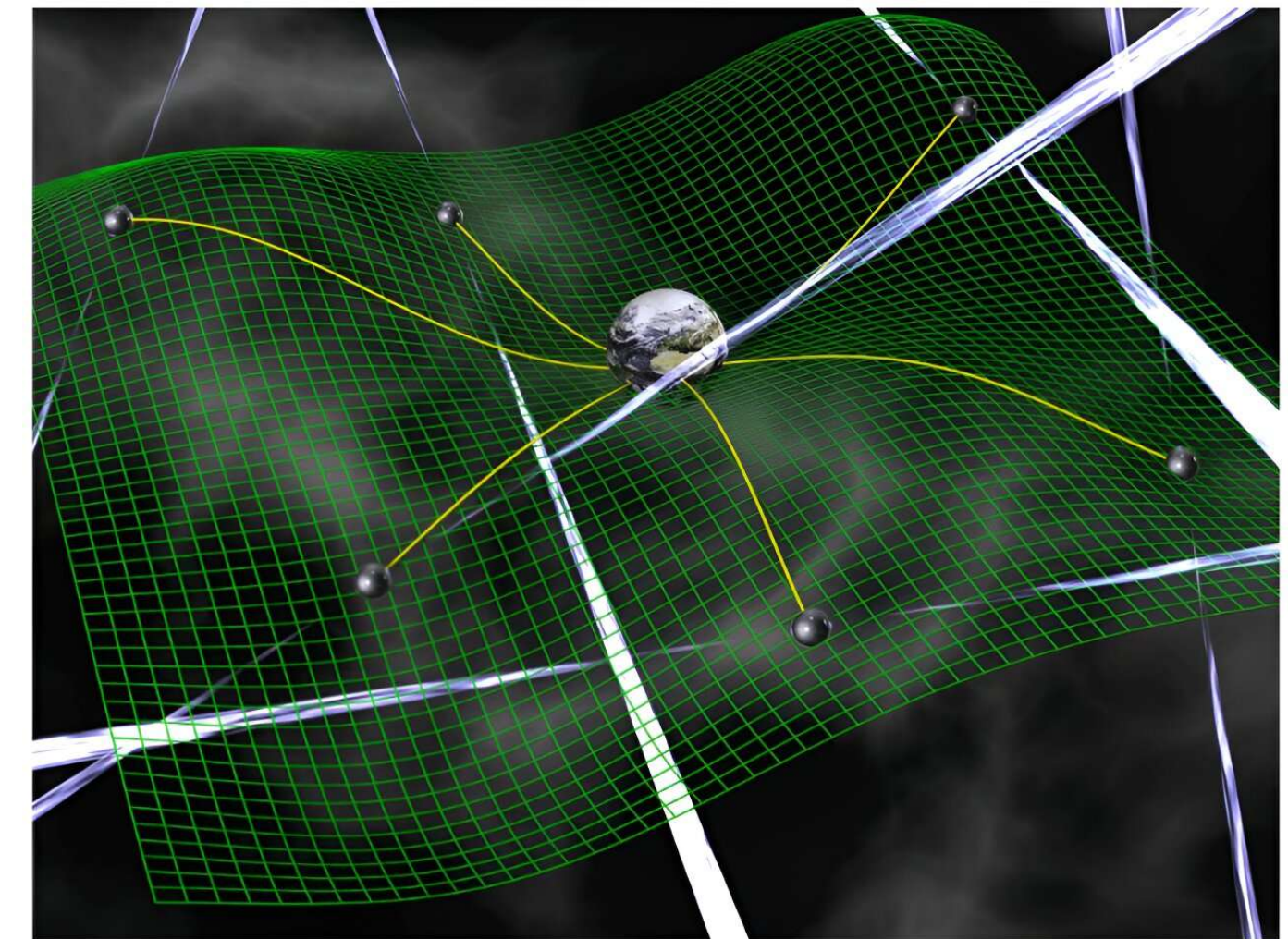
Sources an oscillating gravitational potential. Solving the 00 and ij components of Einstein's equations at first order:

1) 00: constant that obeys Poisson eq. $\nabla^2 \Phi = 4\pi G(m|\psi|^2 - \bar{\rho})$

2) Trace part if ij: oscillating part obeying: $-\ddot{\Phi} \sim 4\pi GP$

with $P = (\dot{\phi}^2 - m^2\phi^2)/2$ oscillating with frequency $2m$

\Rightarrow Φ **oscillates** with frequency $2m$ (and amplitude $\pi G\rho/m^2$)



Artistic representation of the Pulsar Timing Array concept. Credit: David Champion / MPIfR

In the Milky Way: the **constant** part of Φ is of the order 10^{-6} ; the **oscillating** part is $\sim 10^{-12}$

Pulsar Timing *array*

"Second Data Release from the European Pulsar Timing Array: Challenging the Ultralight Dark Matter Paradigm"

Clemente Smarra *et al.* (European Pulsar Timing Array), Phys. Rev. Lett. **131**, 171001

Coherent wave oscillation of ULDM: Φ *oscillates* with frequency $2m$ (and amplitude $\pi G\rho/m^2$)

$$\Phi(t, \vec{x}) \simeq \Phi_0(\vec{x}) + \Phi_c(\vec{x}) \cos(\omega t + 2\alpha(\vec{x})) + \Phi_s(\vec{x}) \sin(\omega t + 2\alpha(\vec{x}))$$

From Poisson eq.: $\Phi_0 \sim G\rho_{dm}/k^2$

Oscillating part: $\Phi_c(\vec{x}) = (1/2)\pi G A(\vec{x})^2 = \pi \frac{G\rho_d m(\vec{x})}{m^2}$
(smaller Φ_0 by a factor of $k^2/m^2 = v^2$)

This oscillation induces a **time-dependent frequency shift and a time delay for any propagating signal**. This is a displacement in the time on arrival (TOAs) of radio pulses emitted by pulsars

$$\delta t_{dm}(t) = - \int_0^t \frac{\Omega(t) - \Omega_0}{\Omega_0} dt'$$

$$\simeq \frac{2\Phi_c}{\omega} \sin\left(\frac{\omega D}{2} + \alpha(\vec{x}) - \alpha(\vec{x}_p)\right)$$

$$\left\{ \begin{array}{l} \Omega(t) \text{ pulse arrival frequency at the detector at the moment } t \\ \Omega_0 \text{ frequency in the absence of this oscillation - coincides with the pulse emission frequency at the pulsar} \end{array} \right.$$

This expression depends on the distance to the pulsar and the scalar field phase α at the locations of the pulsar and the detector.

Pulsar Timing *array*

Comparing with GWs

ULDM

$$\delta t_{dm} \simeq \frac{2\Phi_c}{\omega} \sin\left(\frac{\omega D}{2} + \alpha(\vec{x} - \alpha(\vec{x}_p))\right)$$

Root mean square values of the time residuals, averaged over the distance to the pulsar:

$$\sqrt{\langle \delta t_{dm}^2 \rangle} = \sqrt{2} \frac{\Phi_c}{\omega}$$

ULDM has same effect on the pulsar timing measurements as GWB with characteristic strain:

$$h_c = 2\sqrt{3}\Phi_c = 2 \times 10^{-15} \left(\frac{\rho_{dm}}{0.3\text{GeV/cm}}\right) \left(\frac{10^{-23}\text{eV}}{m}\right)^2$$

At frequency

$$f \equiv 2\pi\omega = 5 \times 10^{-9}\text{Hz} \left(\frac{m}{10^{-23}\text{eV}}\right)$$

GWs

For a single monochromatic gravitation wave with frequency ω and characteristic strain h_c the amplitude of the timing residual

$$\delta t_{gw} = \frac{h_c}{\omega} \sin\left(\frac{\omega D(1 - \cos(\theta))}{2}\right) (1 + \cos\theta) \sin(2\psi)$$

Polarization angle
of GW

Direction to the
source

RMS: averaged over D ($\omega D \ll 1$), θ and ψ :

$$\sqrt{\langle \delta t_{gw}^2 \rangle} = \frac{1}{\sqrt{3}} \frac{h_c}{\omega}$$

Pulsar Timing *array*

Comparing with GWs

ULDM

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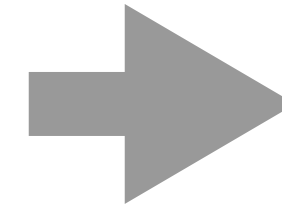
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At frequency

$$f \equiv 2\pi\omega = 5 \times 10^{-9}\text{Hz} \left(\frac{m}{10^{-23}\text{eV}}\right)$$

EPTA: for the ULDM coherent oscillation, the PTA frequency coverage corresponds to

$$m \sim 10^{-24} - 10^{-22}\text{eV}$$



Pulsar Timing *array*

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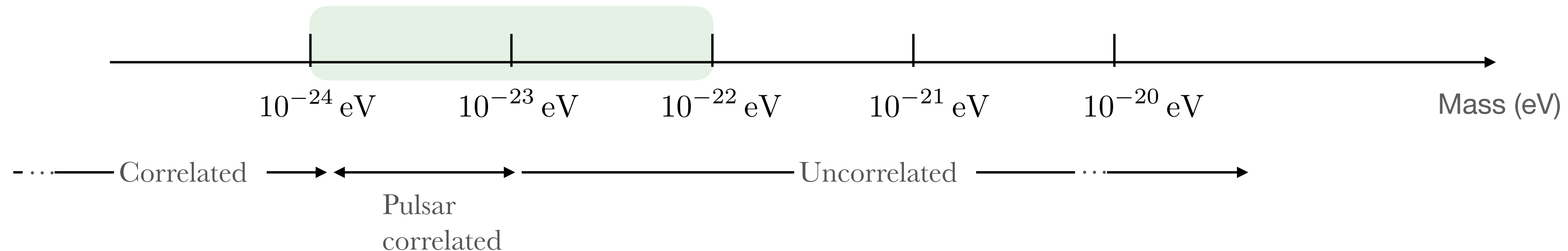
$$\delta t_{dm} = \frac{\Phi(\vec{x})}{2m} [\hat{\phi}_E^2 \sin(2m + \alpha_E) - \hat{\phi}_P^2 \sin(2m + \alpha_P)]$$

$$\Phi(\vec{x}) \sim 6.52 \times 10^{-18} \left(\frac{10^{-22} \text{ eV}}{m} \right)^2 \left(\frac{\rho_\phi}{0.4 \text{ GeV/cm}^3} \right)$$

$$\alpha_P \equiv 2\alpha(\vec{x}_p) - 2md_p/c;$$

$$\alpha_E \equiv 2\alpha(\vec{x}_E)$$

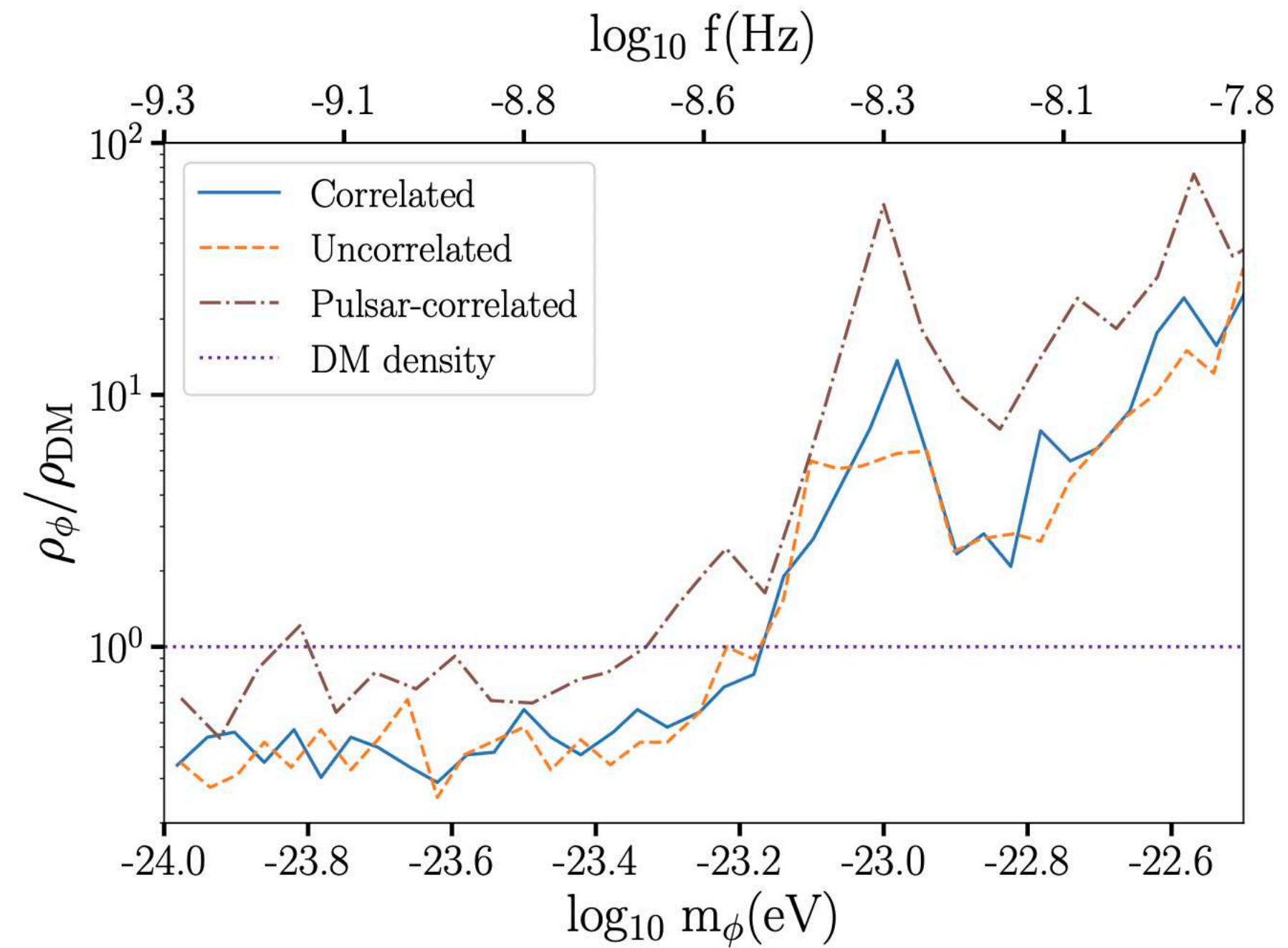
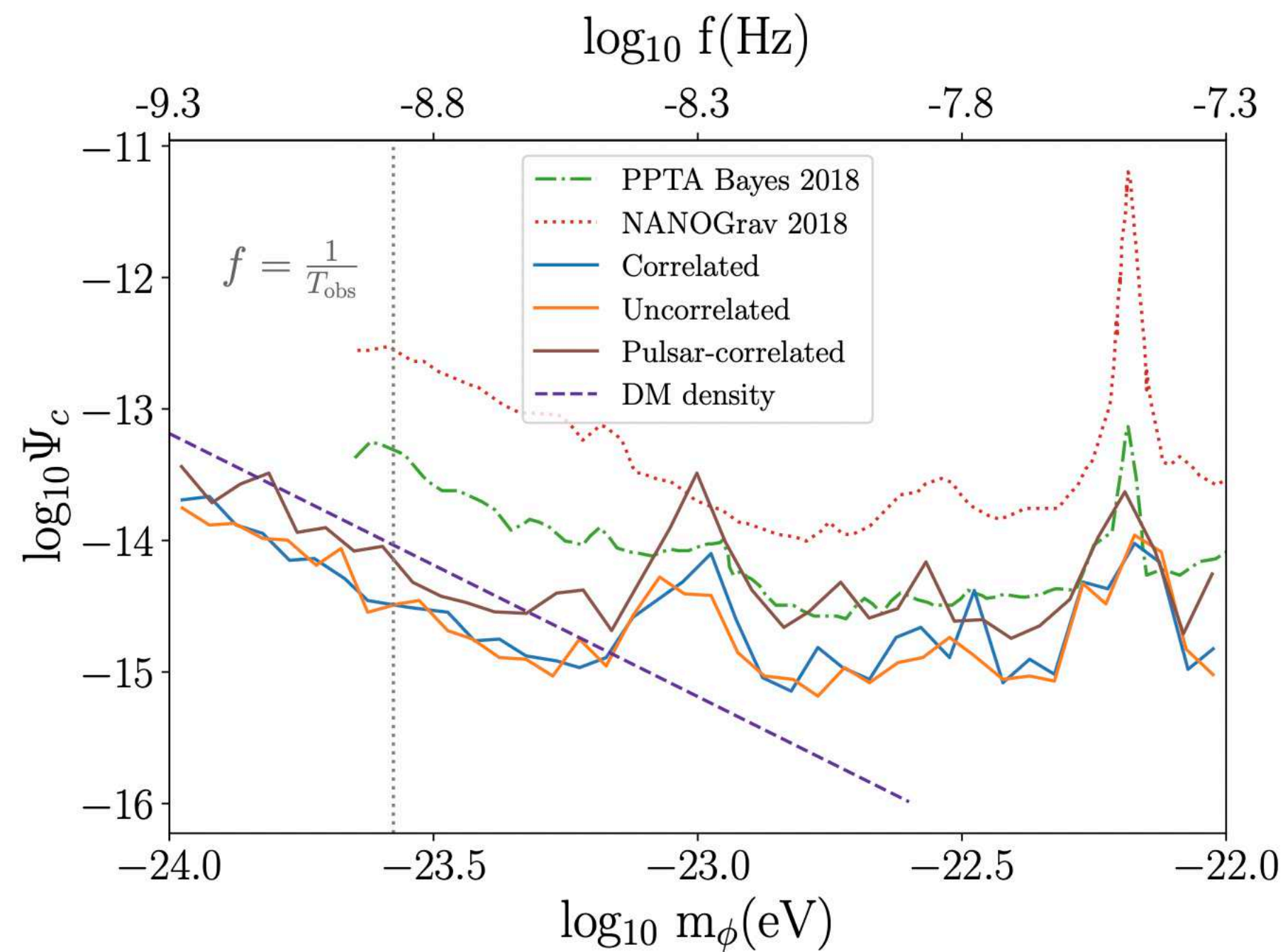
$\hat{\phi}(\vec{x})$ accounts for the interference pattern in the proximity of \vec{x}



Pulsar Timing array

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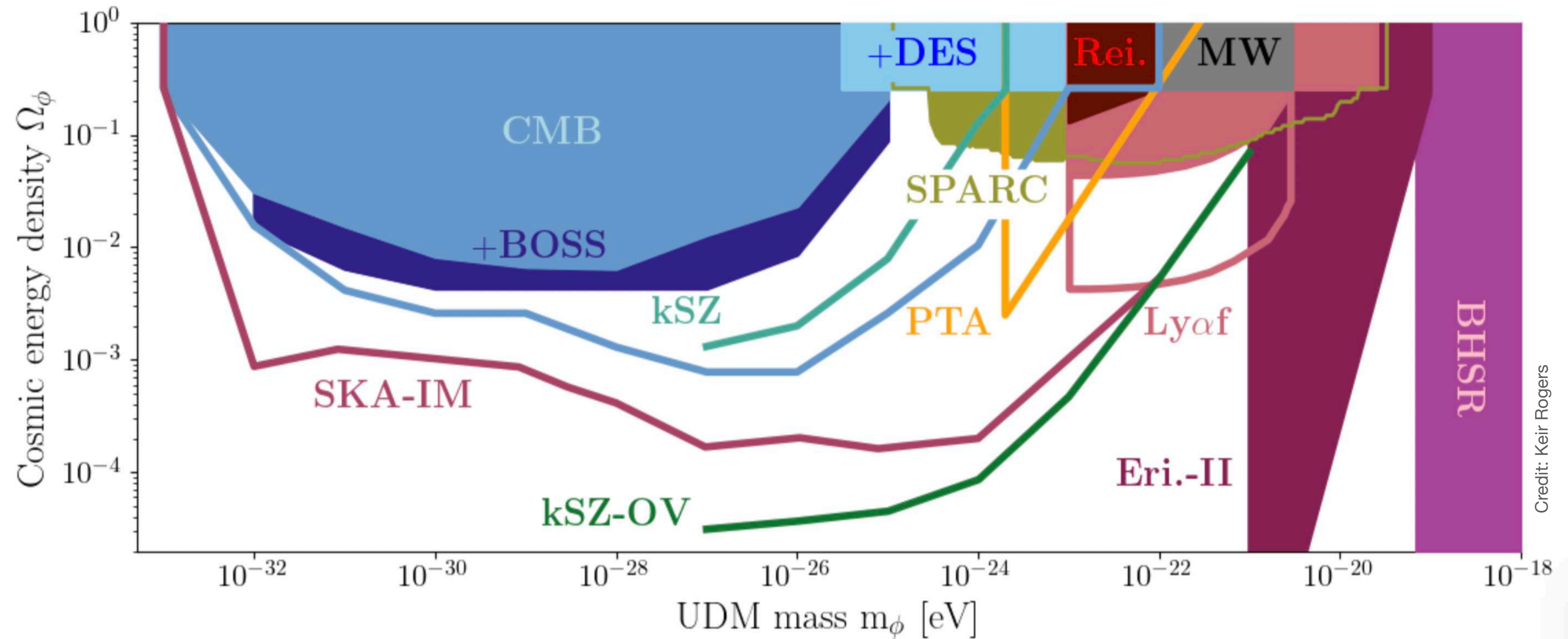
$-24.0 < \log_{10} (m_\phi/\text{eV}) < -23.7$ \longrightarrow $f_\phi \lesssim 30 - 40 \%$

$-23.7 < \log_{10} (m_\phi/\text{eV}) < -23.3$ \longrightarrow up to $f_\phi \sim 70 \%$

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Credit: Keir Rogers

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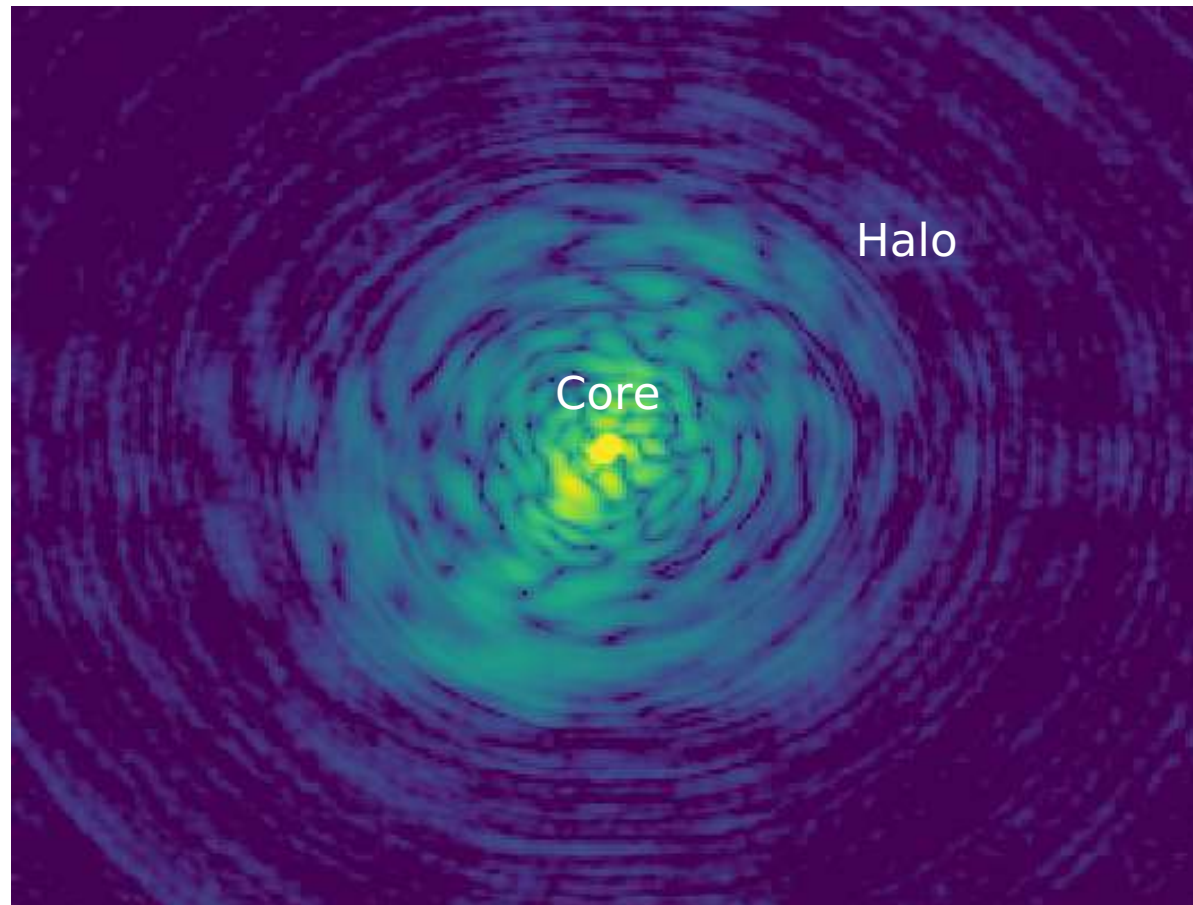
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de Broglie scale *time delays*

(Work in progress)

In collaboration with Andrew Eberhardt
and Qiuyue Liang

Our analysis: Considering the interference pattern

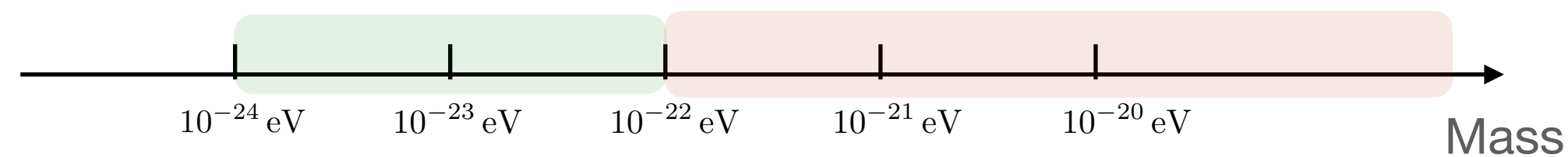


~~$$\phi(t, \vec{x}) = m^{-1} \sqrt{2\rho(t, \vec{x})} \cos[mt + \theta(t, \vec{x})]$$~~

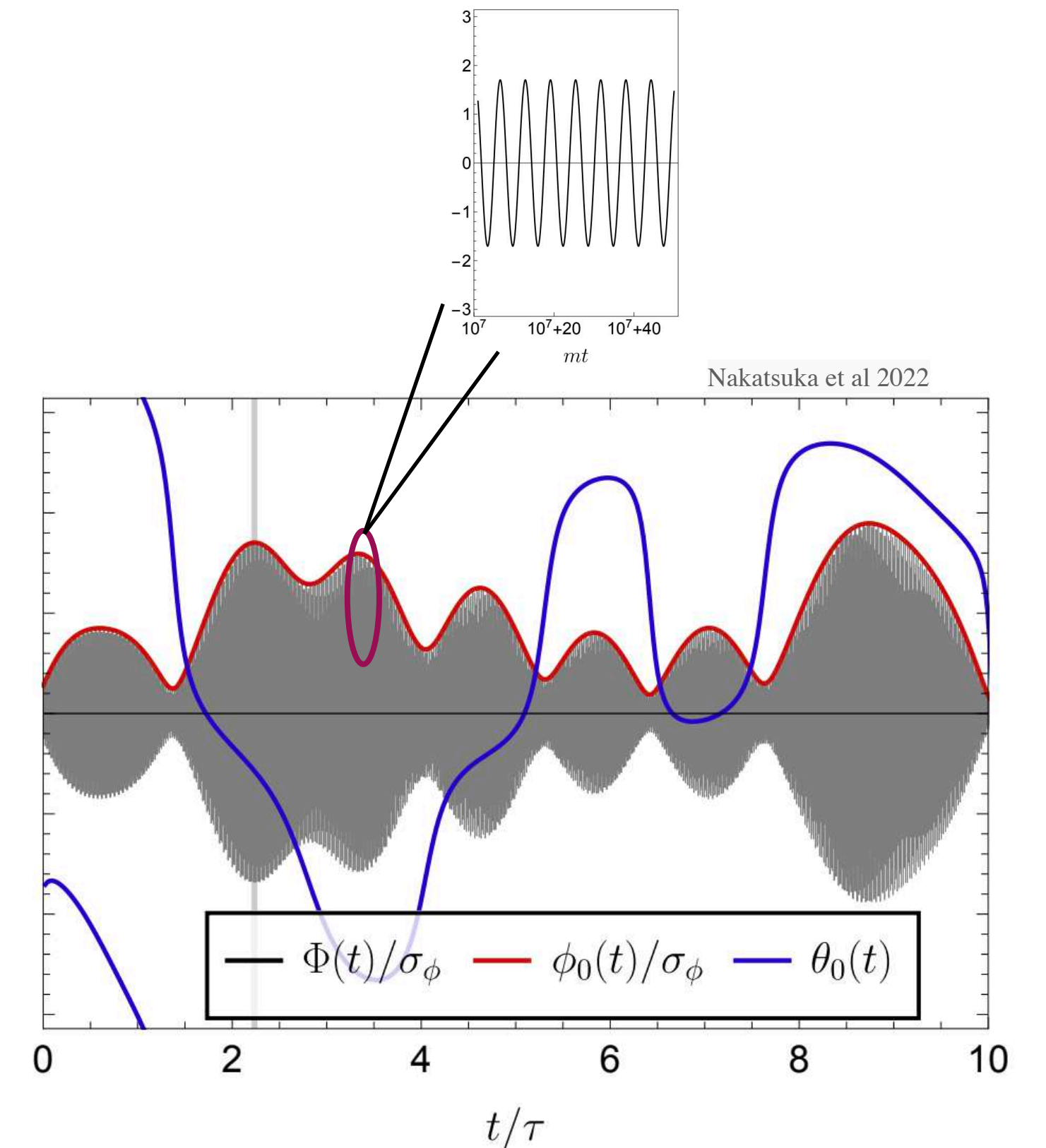
Random phase halo model: as a simple model of a galactic halo, consider a **superposition of plane waves**:

$$\phi(t, \vec{x}) = \sum_i^N \frac{\phi(0)}{\sqrt{N}} \cos \left(mt + \frac{m}{2} v_i^2 t - m\vec{v}_i \cdot \vec{x} + \theta_i \right)$$

Our work



de Broglie scale time delays



Pulsar Timing *array*

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In collaboration with Andrew Eberhardt
and Qiuyue Liang

Our analysis: Considering the interference pattern

$$\tau_{coh} = \frac{2\pi}{\Delta E \phi} \sim 10^6 T_{osc}$$

$$m t \ll \tau_{coh}$$

$$m t = t_{propagating} \sim \tau_{coh}$$

$$m t = t_{obs} \sim \tau_{coh}$$

Correlated

$$\Phi(t) \sim \cos(mt)$$

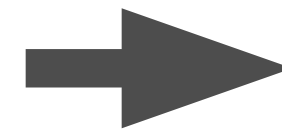
+

Large granules:

Ex.: $m \sim 10^{-22}$ eV, granules of 1/10 of kpc

Angular correlation of PTA signal within a granular size would be enhanced due to the superposition.

$$t_{propagating} \sim 3000 \text{ yrs}$$



$$3000 \text{ yrs} \sim 10^6 / m$$

$$m \sim 4.38 \times 10^{-20} \text{ eV}$$

Pulse would go through many granules during the propagation.

$$t_{obs} \sim 10 \text{ yrs}$$



$$30 \text{ yrs} \sim 10^6 / m$$

$$m \sim 1.3 \times 10^{-17} \text{ eV}$$

*de Broglie scale **time delays***

*de Broglie scale time delays in pulsar networks for
ultralight dark matter*

Andrew Eberhardt, Qiuyue Liang, EF, 2411.18051

$$m \gtrsim 10^{-17} \text{ eV}$$

granules fluctuate on timescales similar to observational timescales

de Broglie scale *time delays*

*de Broglie scale time delays in pulsar networks for
ultralight dark matter*

Andrew Eberhardt, Qiuyue Liang, EF, 2411.18051

Time delays:

$$\delta t_{dm}(t) = - \int_0^t \frac{\Omega(t') - \Omega_0}{\Omega_0} dt'$$

$$\Delta\Omega = \underbrace{\Delta\Omega_{sh}^{osc} + \Delta\Omega_{sh}}_{\text{Shapiro time delay}} + \underbrace{\Delta\Omega_{gr}^{osc} + \Delta\Omega_{gr}}_{\text{Gravitational redshift}}$$

Gravitational redshift

$$\frac{\Delta\Omega_{gr}}{\Omega_0} = \Phi_P - \Phi_E$$



The frequency change comes from the time variation of this path

$$\frac{\Delta\Omega_{gr}}{\Omega_0} = \frac{\partial_t \ell}{1 - \partial_t \ell}$$

Shapiro time delay

The travel time as an integral along the path from the pulsar to Earth

$$\ell = \int dl' (1 + 2\Phi(l'))$$

$$\longrightarrow \delta t_{sh} = \int_p^e \frac{dl}{c} 2\Phi(l)$$

de Broglie scale *time delays*

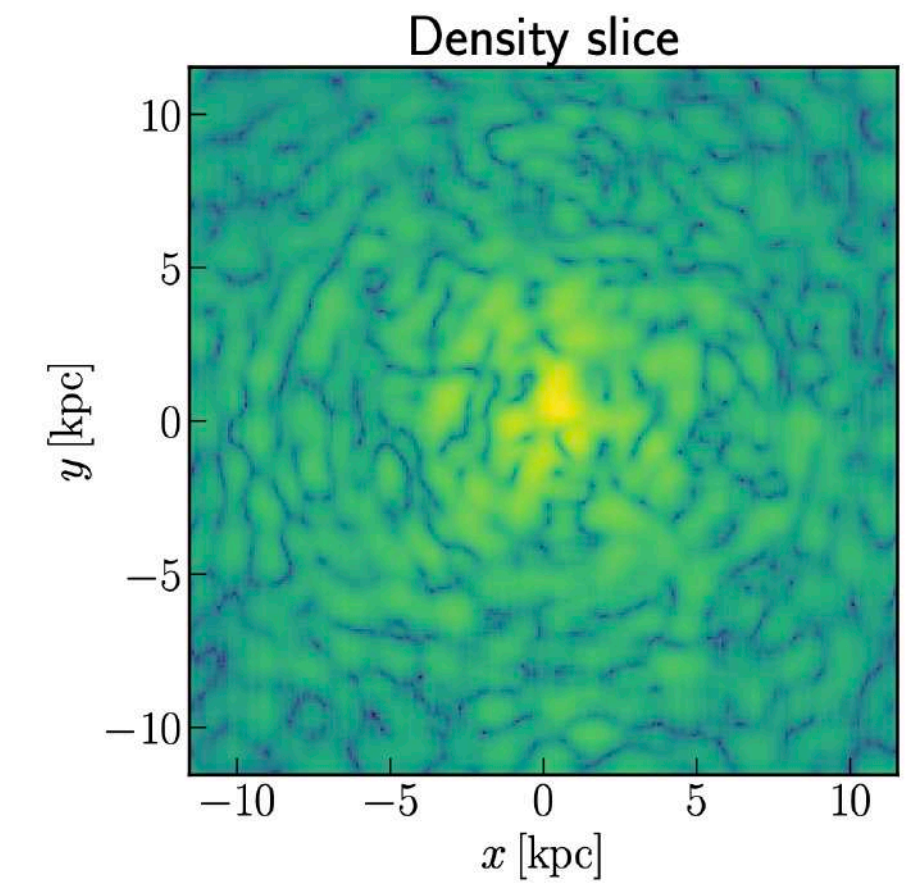
*de Broglie scale time delays in pulsar networks for
ultralight dark matter*

Andrew Eberhardt, Qiuyue Liang, EF, 2411.18051

$m \gtrsim 10^{-17}$ eV granules fluctuate on timescales similar to observational timescales

We simulate arrays of mock pulsars in a fluctuating granular density field

$$\delta t_{\text{rms}}^{\text{sh}} \sim 5 \times 10^{-2} \left(\frac{10^{-17} \text{ eV}}{m} \right)^{7/3} \left(\frac{200 \text{ km/s}}{\sigma} \right)^{7/3} \left(\frac{\rho}{10^7 \text{ M}_{\odot}/\text{kpc}^3} \right) \left(\frac{D}{\text{kpc}} \right)^{2/3} \text{ ns},$$
$$\delta t_{\text{rms}}^{\text{z}} \sim 4 \times 10^{-4} \left(\frac{10^{-17} \text{ eV}}{m} \right)^{3/2} \left(\frac{200 \text{ km/s}}{\sigma} \right)^4 \left(\frac{\rho}{10^7 \text{ M}_{\odot}/\text{kpc}^3} \right) \left(\frac{T}{30 \text{ yrs}} \right)^{3/2} \text{ ns},$$



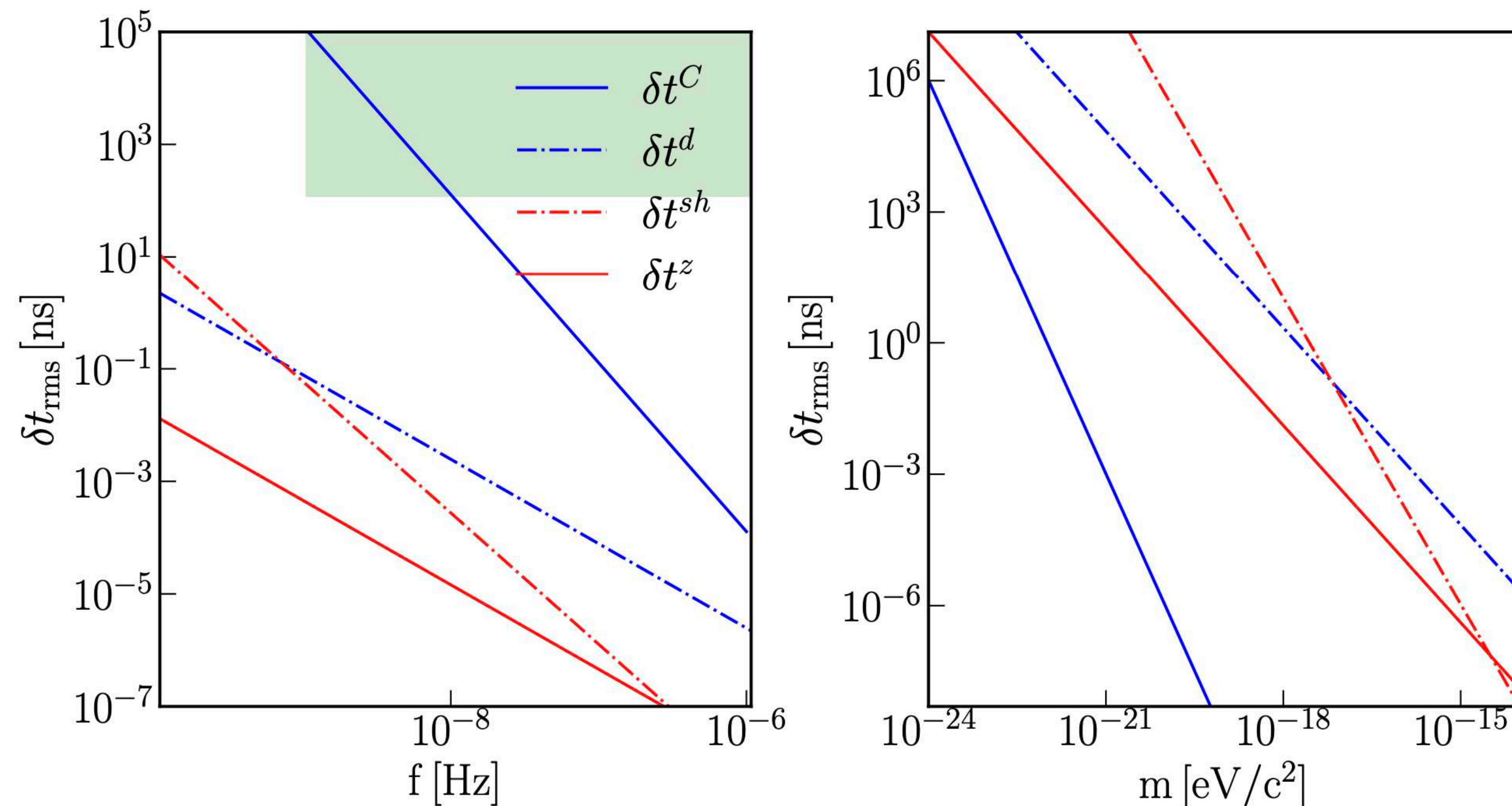
de Broglie scale *time delays*

*de Broglie scale time delays in pulsar networks for
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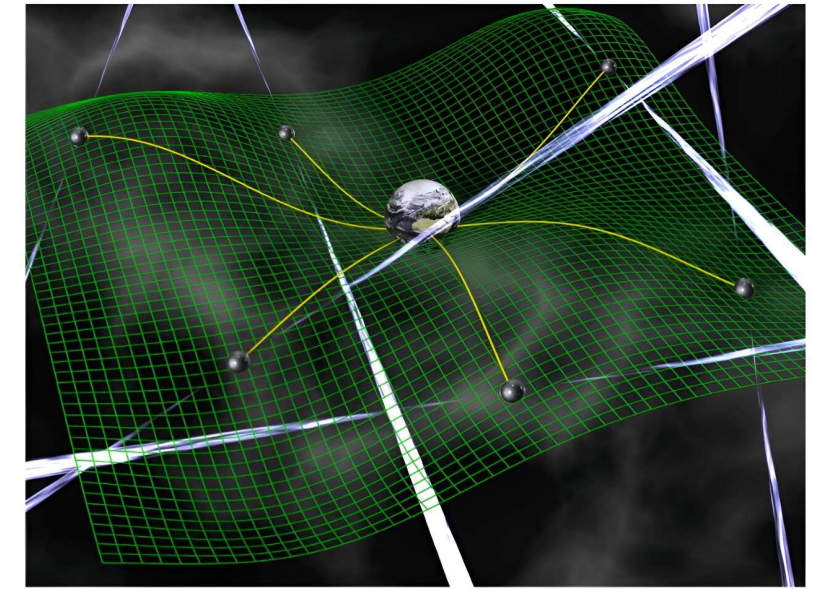
We simulate arrays of mock pulsars in a fluctuating granular density field



- Not the dominant effect for pulsar timing experiments
- dB scale effects on observable timescales could potentially provide sensitivity to mass scales higher than those probed thus far by small-scale structure
- Current largest limitation appears to be the longest de Broglie time to which experiments are sensitive
- Experiments with longer runtimes or somehow sensitive to **longer time scales** would substantially improve the signal from the effects studied here

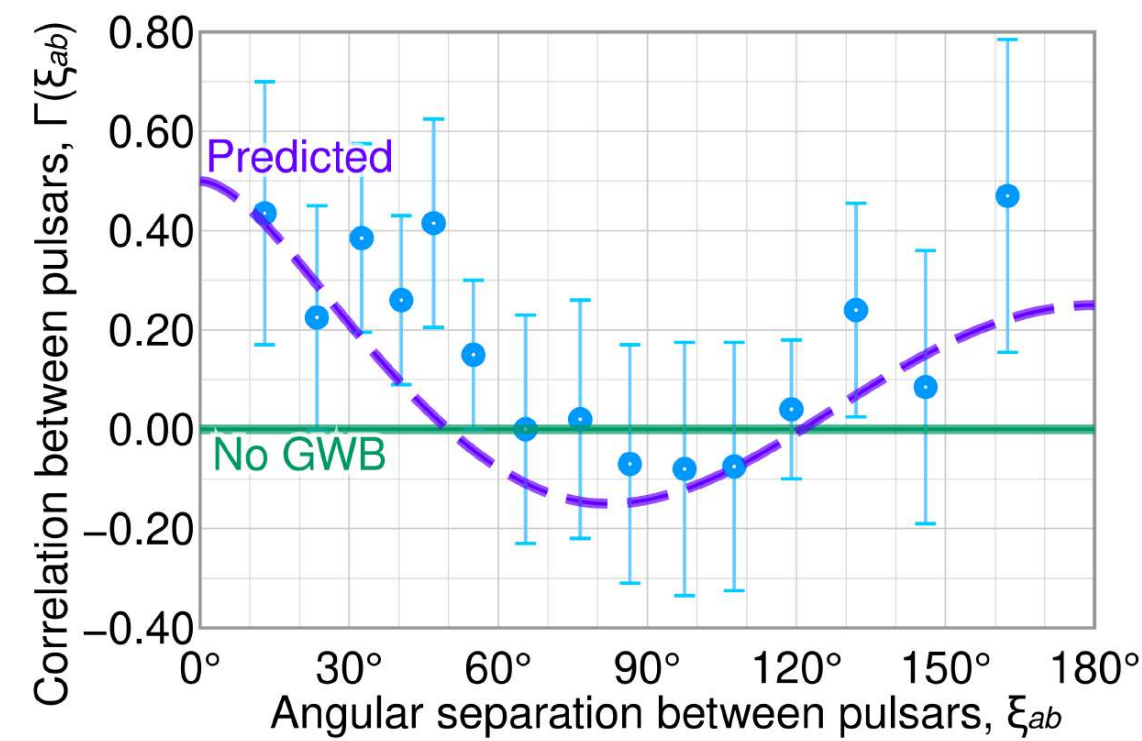
The left edge of the shaded regions in green corresponds to when the relevant timescale (dB time for dB scale effects or Compton time for Compton scale effects) is 30 yrs

Pulsar Timing array



Artistic representation of the Pulsar Timing Array concept. Credit: David Champion / MPhIR

NANOGrav, EPTA, PPTA, and InPTA announced that they found evidence for a gravitational wave background



Possibility of being a **Gravitational Wave Background** signal caused by:

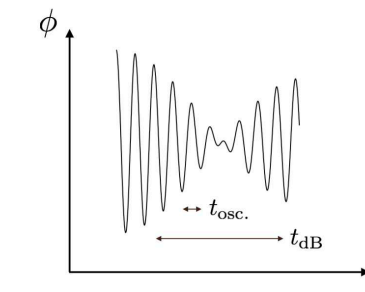
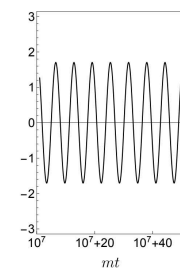
- Ensemble of binary supermassive black-holes
- New physics

The presence of ULDM induces an oscillating gravitational potential that affects the light travel time of radio pulses emitted by pulsars. PTAs can be used to test the presence of ULDM particles in the MW \rightarrow how ULDM affects the SGWB in PTA system

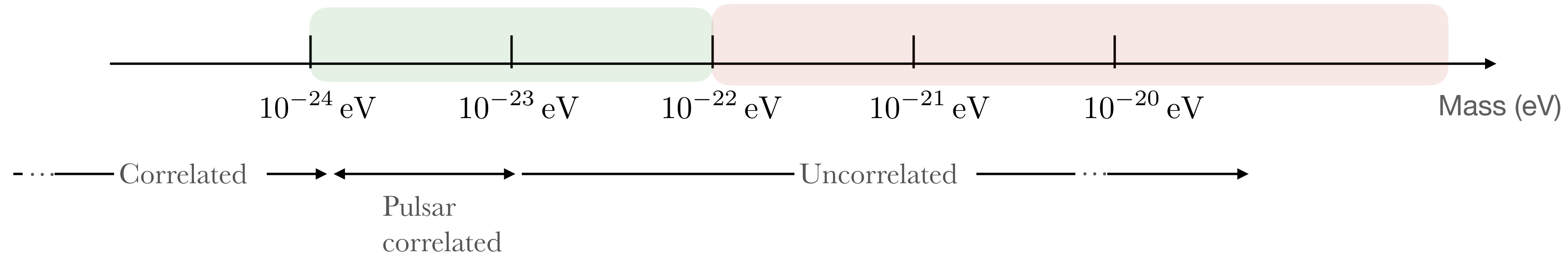
EPTA, Smarra et al 2024

$$-24.0 < \log_{10}(m_\phi/\text{eV}) < -23.7 \longrightarrow f_\phi \lesssim 30 - 40\%$$

$$-23.7 < \log_{10}(m_\phi/\text{eV}) < -23.3 \longrightarrow \text{up to } f_\phi \sim 70\%$$



Our work



Pulsar Timing *array*

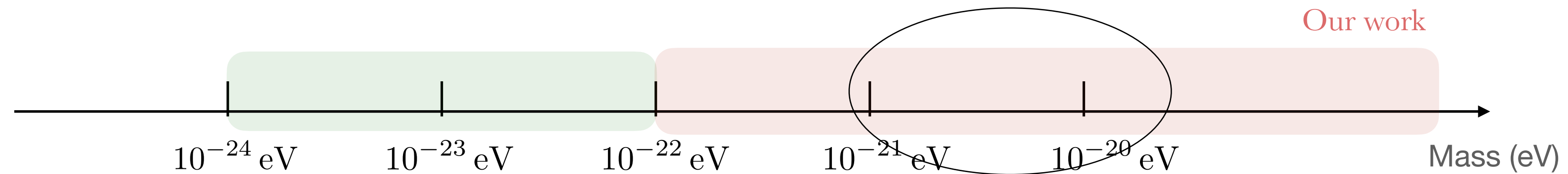
(Work in progress)

In collaboration with Andrew Eberhardt
and Qiuyue Liang

The presence of ULDM induces an oscillating gravitational potential that affects the light travel time of radio pulses emitted by pulsars.

PTAs can be used to test the presence of ULDM particles in the MW \rightarrow how ULDM affects the SGWB in PTA system

Considering the interference pattern



de Broglie scale time delays

Pulse would go through many granules
during the propagation.

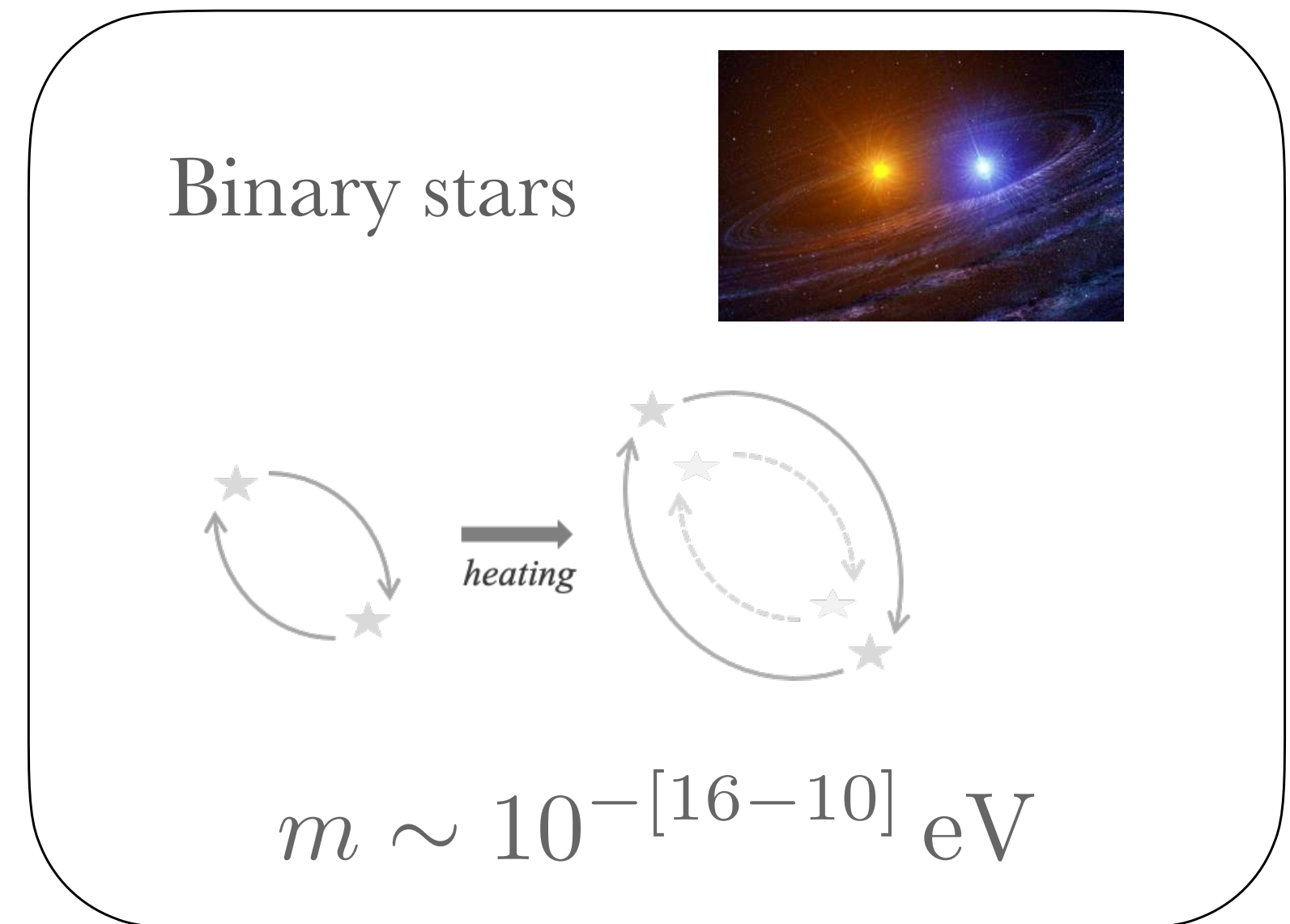
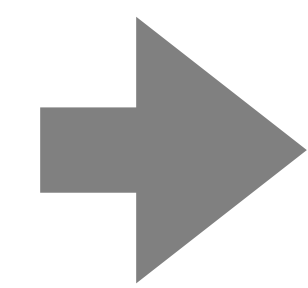
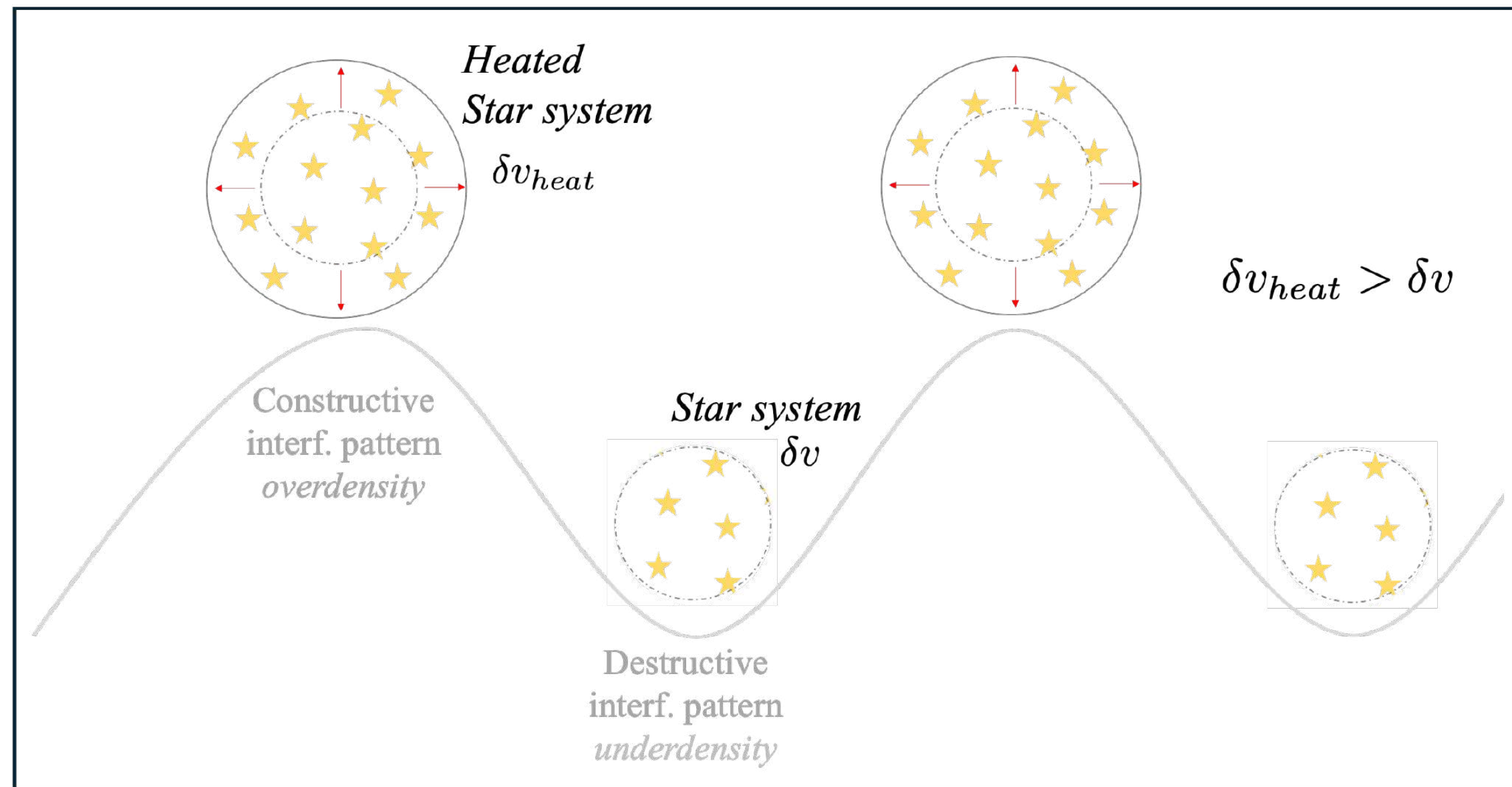
Angular correlation

Binary stars - as a probe of ULDM

(Work in progress)

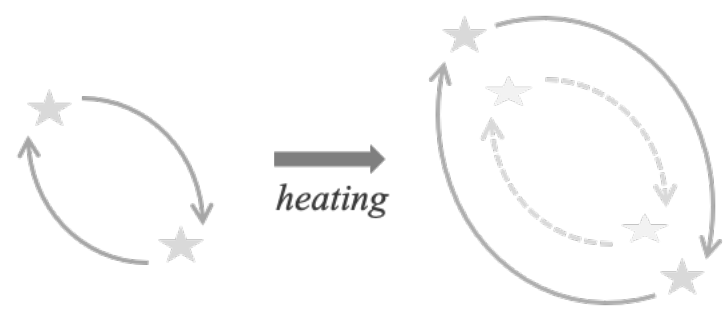
In collaboration with Andrew Eberhardt,
Margot Imbach and Naoki Yoshida

Stellar heating



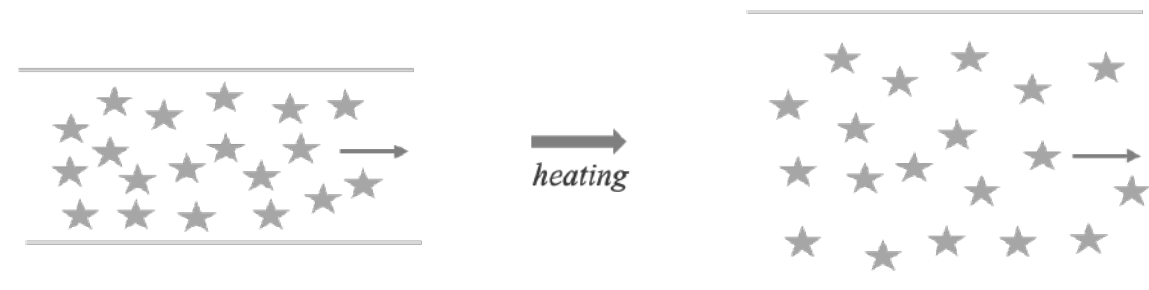
Preparation for PFS!

Binary stars



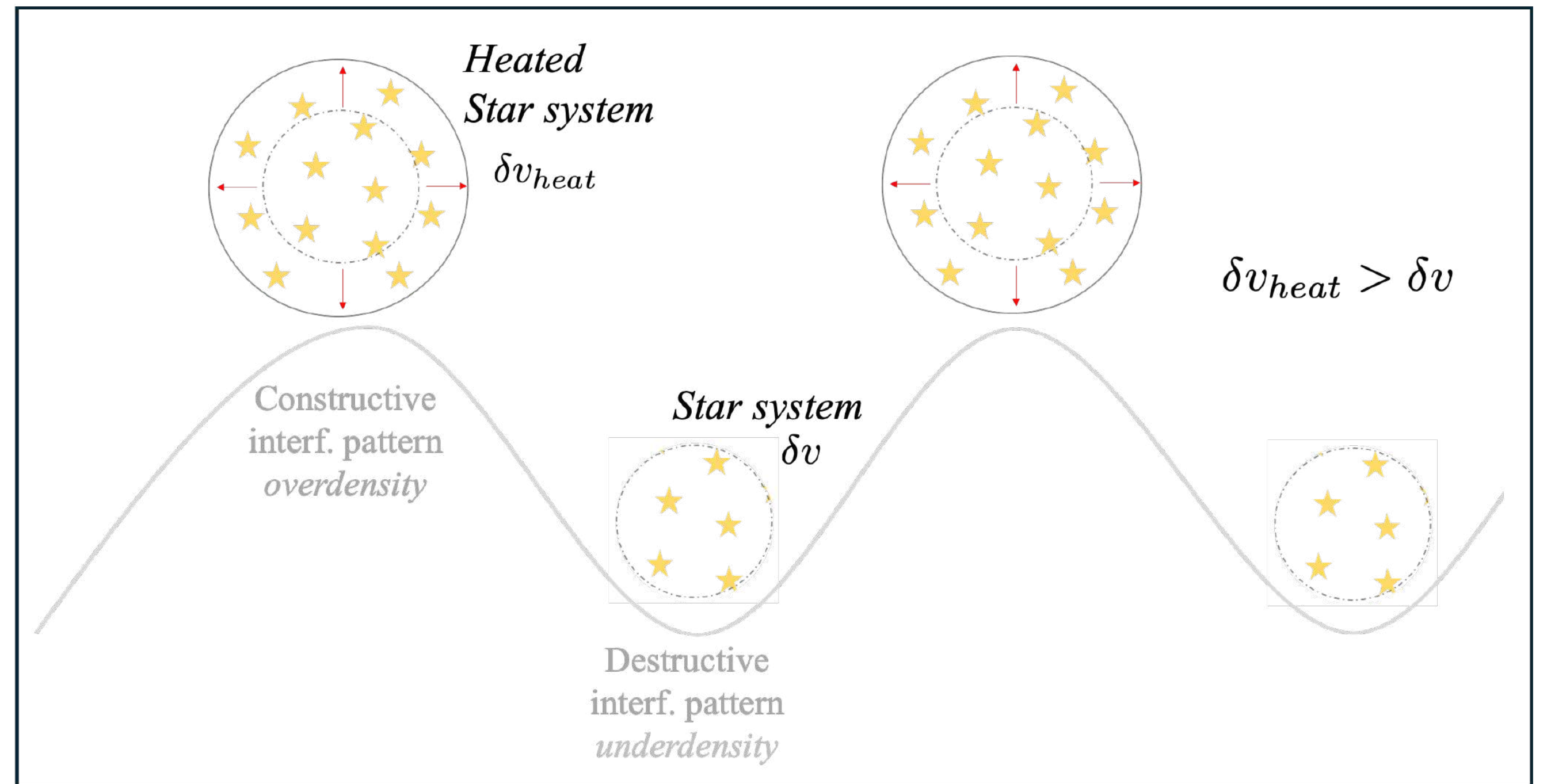
$$m \sim 10^{-[16-10]} \text{ eV}$$

Stellar streams



$$m \sim 10^{-[23-18]} \text{ eV}$$

Stellar heating

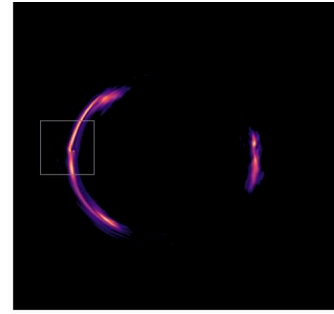


(Work in progress)

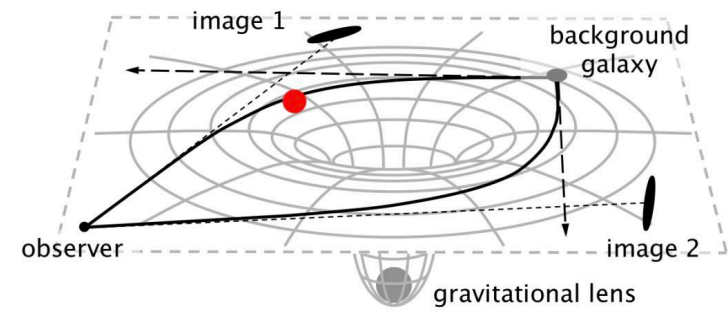
In collaboration with Andrew Eberhardt,
and Fabian Schmidt

Preparation for PFS!

*The search for **dark matter (ULDM)** is a
multi-probe/multi-scale endeavour...*



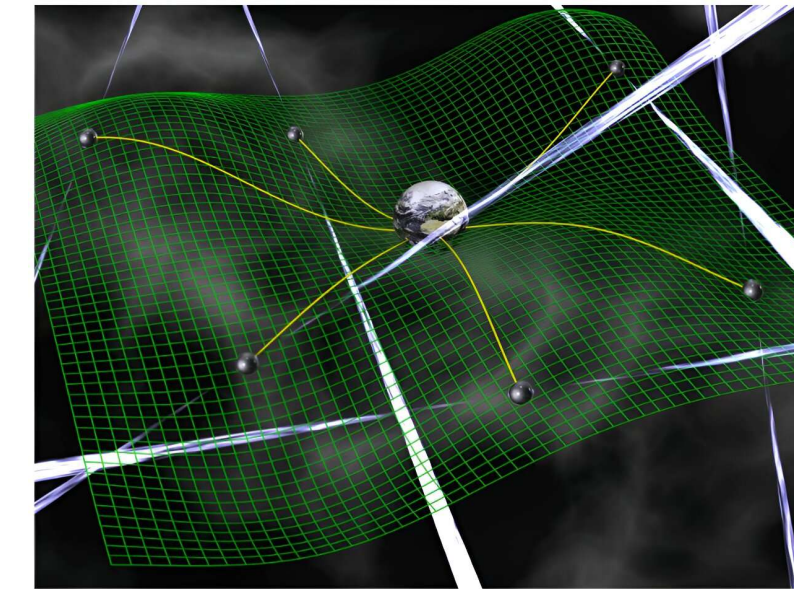
Strong lensing



$$m \sim 10^{-[22-17]} \text{ eV}$$

(Work in progress)

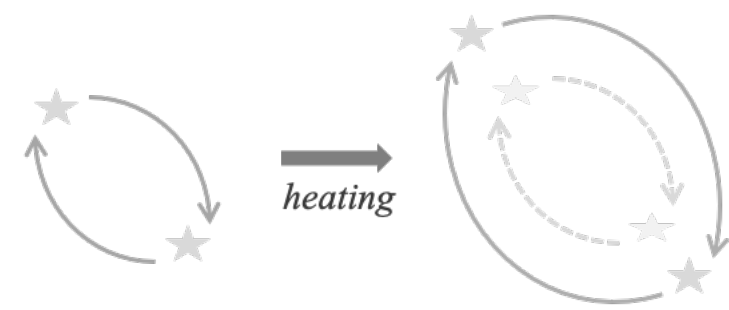
Pulsar Timing Array



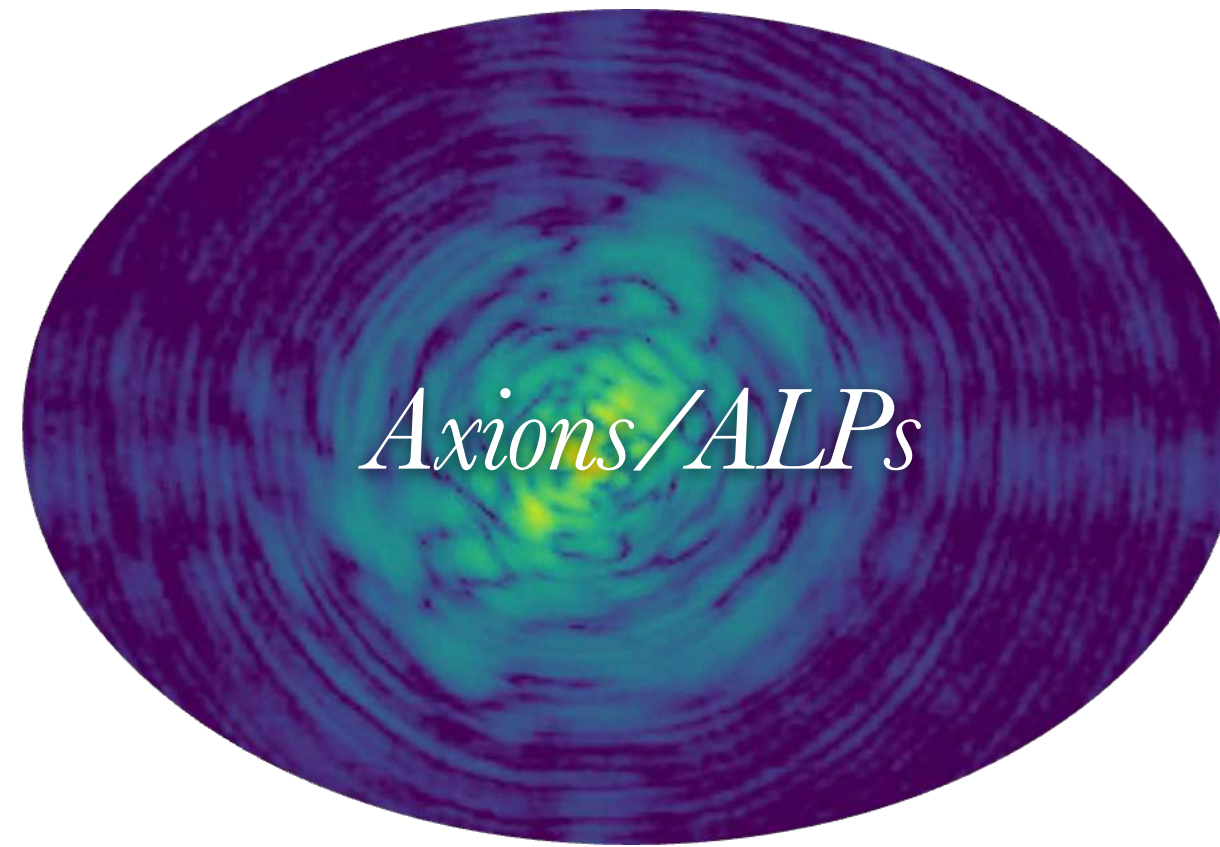
Artistic representation of the Pulsar Timing Array concept. Credit: David Champion / MPIfR

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Binary stars



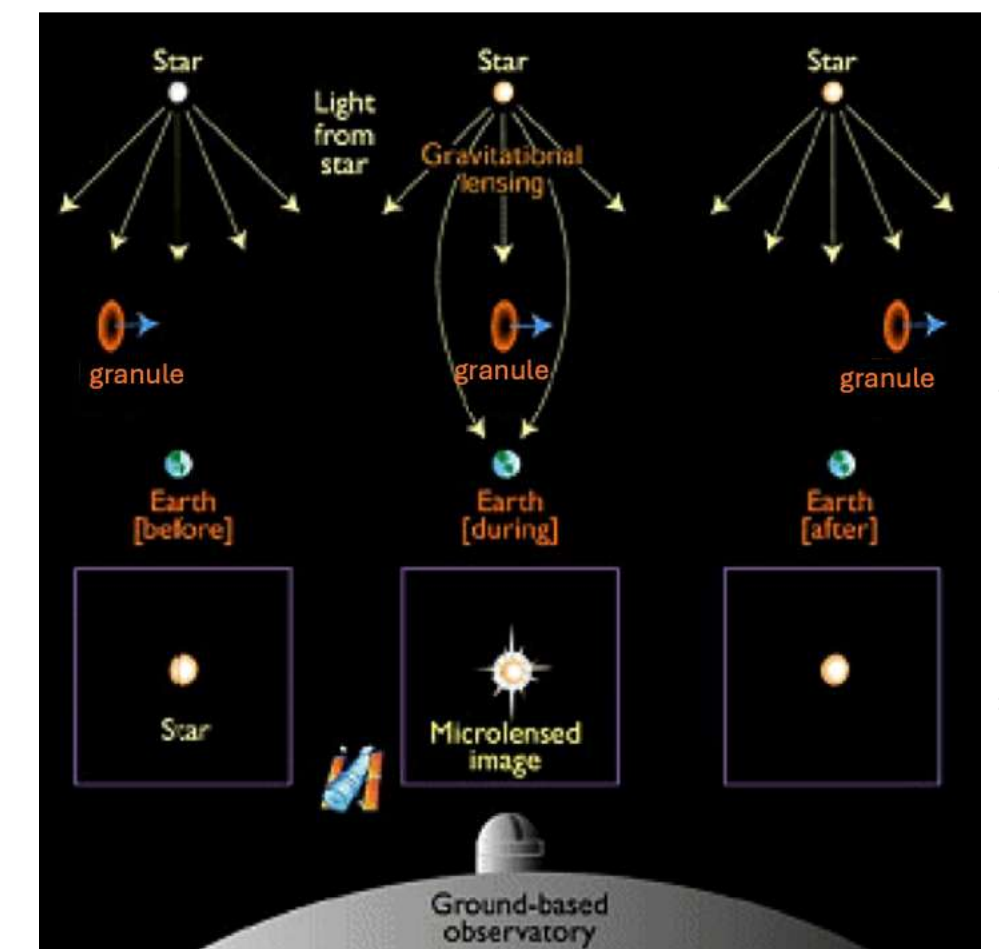
$$m \sim 10^{-[16-10]} \text{ eV}$$



Axions/ALPs

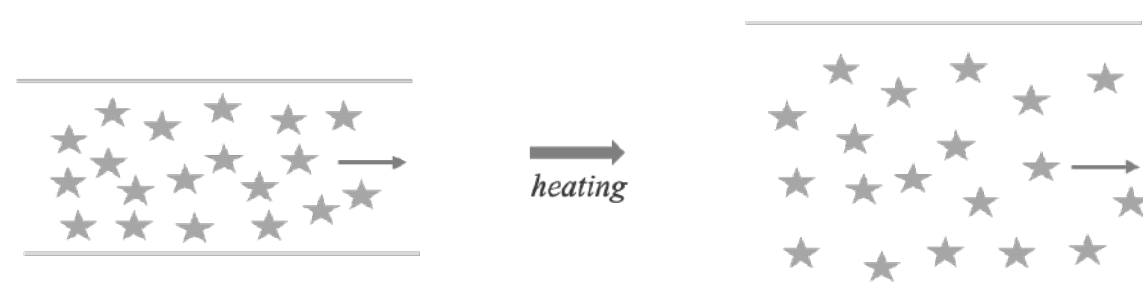
Mass, spin (# particles), fraction self interaction axion-photon coupling

Microlensing



$$m \sim 10^{-[23-18]} \text{ eV}$$

Stellar streams



$$m \sim 10^{-[23-18]} \text{ eV}$$

Solar system



$$m \sim 10^{-[16-10]} \text{ eV}$$

axion-photon coupling

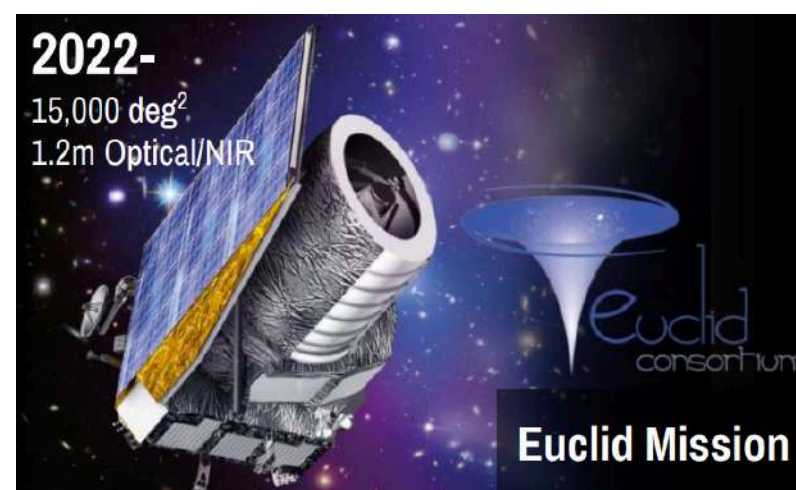
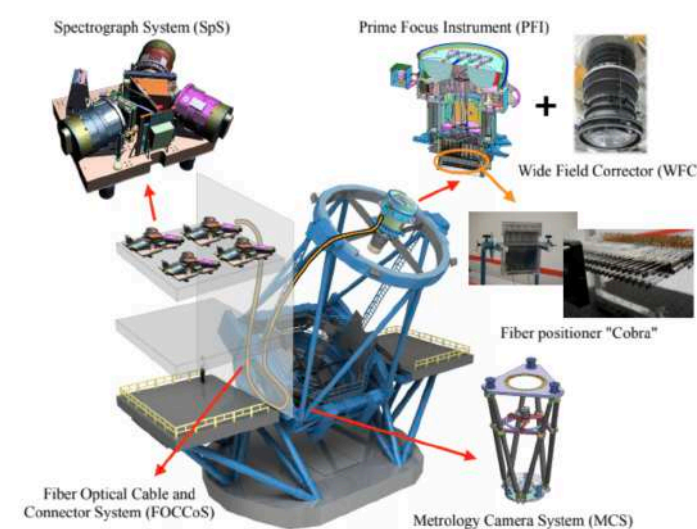
Improving these bounds

Observations

Photometric and spectroscopic surveys



Prime Focus Spectrograph (PFS)



21cm



CMB



CMB-S4
Next Generation CMB Experiment

GWs

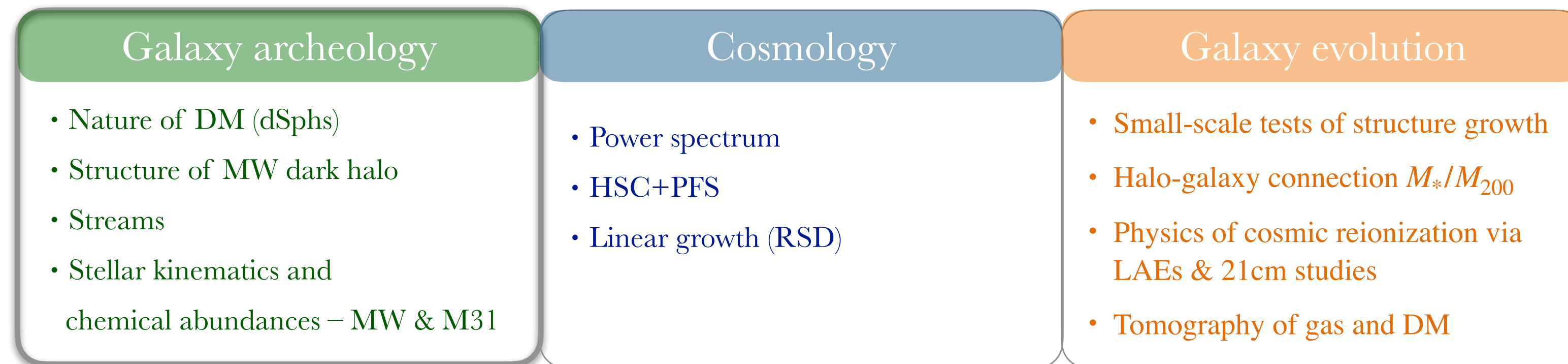
+ direct detection experiments

PFS (Prime Focus Spectrograph)

PFS is going to be exquisite to measure the properties of DM

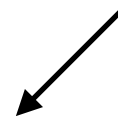
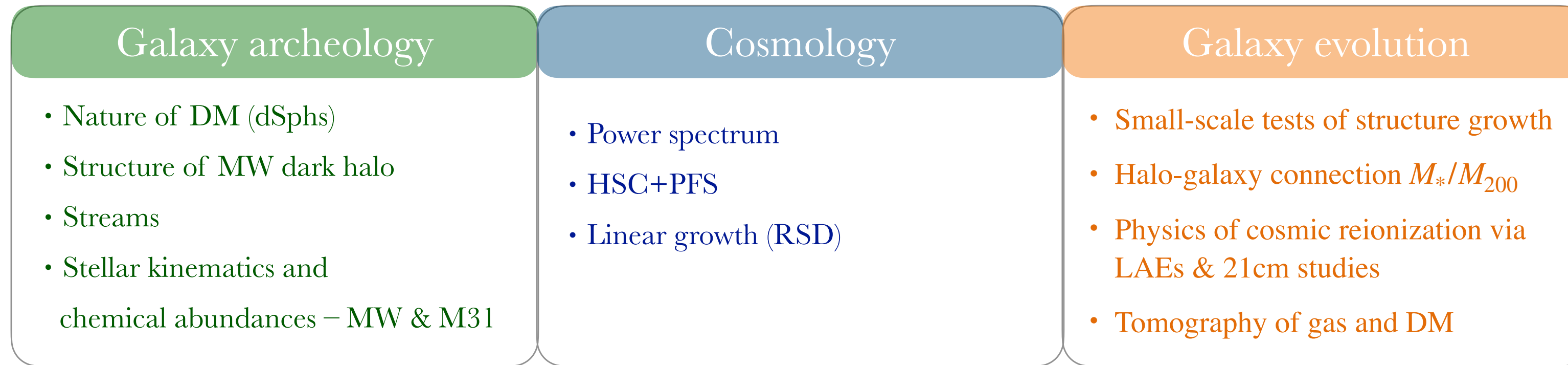
PFS: spectroscopy part of *SuMIRe project*

DM with PFS → synergy between science goals



Wide & deep survey of MW dwarf galaxies w. Subaru/PFS

DM with PFS



- Science with dwarf galaxies

Core:

- Presence of a core or not (slope)
- Size of the core
- Profile

FDM

SIDM

ULA

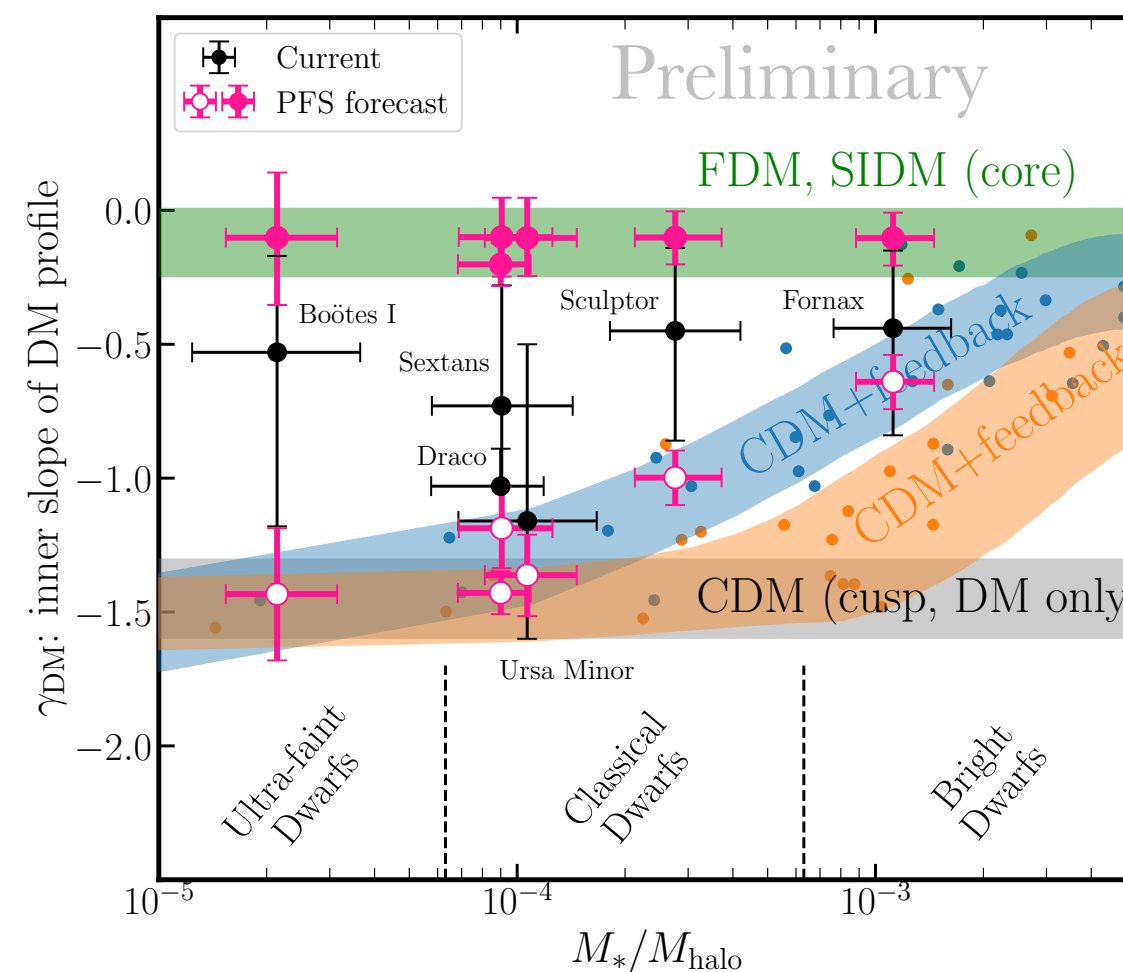
- Inner density
- Transition radius
- Abundance data to understand the role of baryons in each system

- Beyond the core

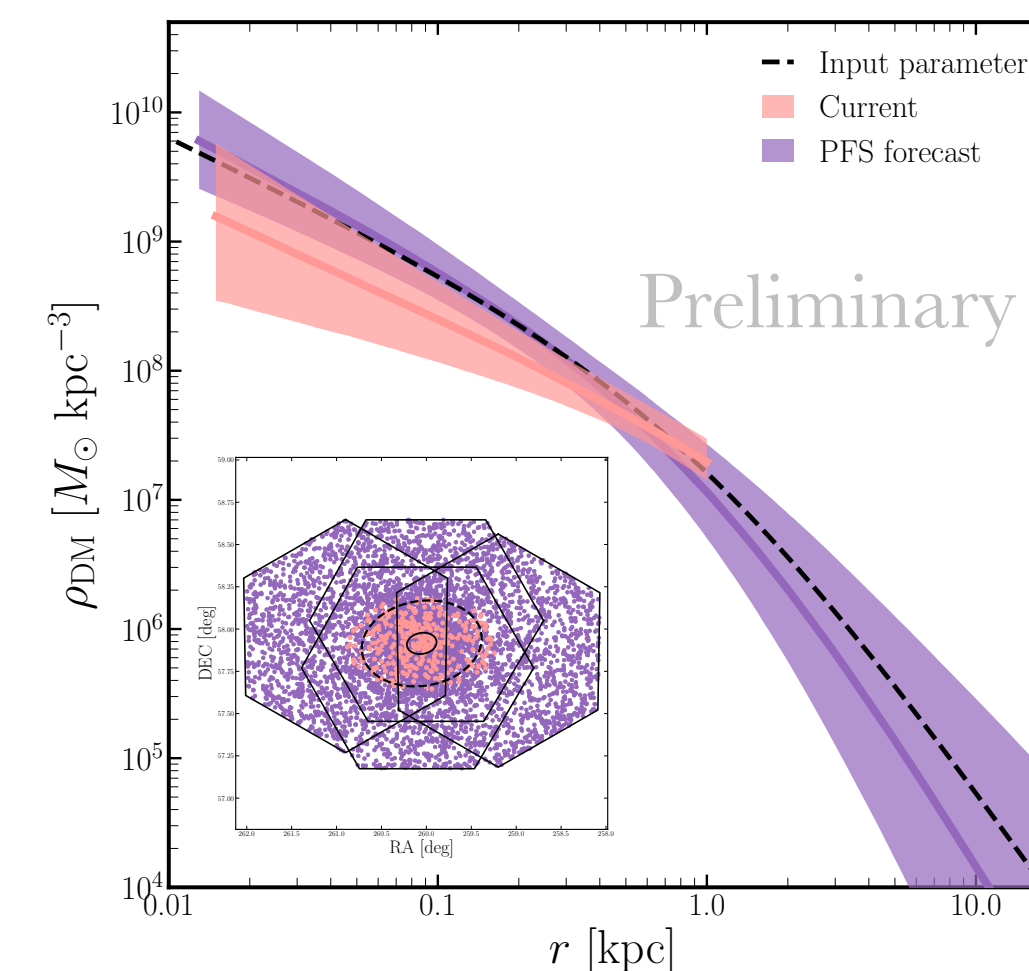
- Granules: heating of stars (dwarfs)
 - Angular momentum
- Stellar streams
- Binary stars

FDM

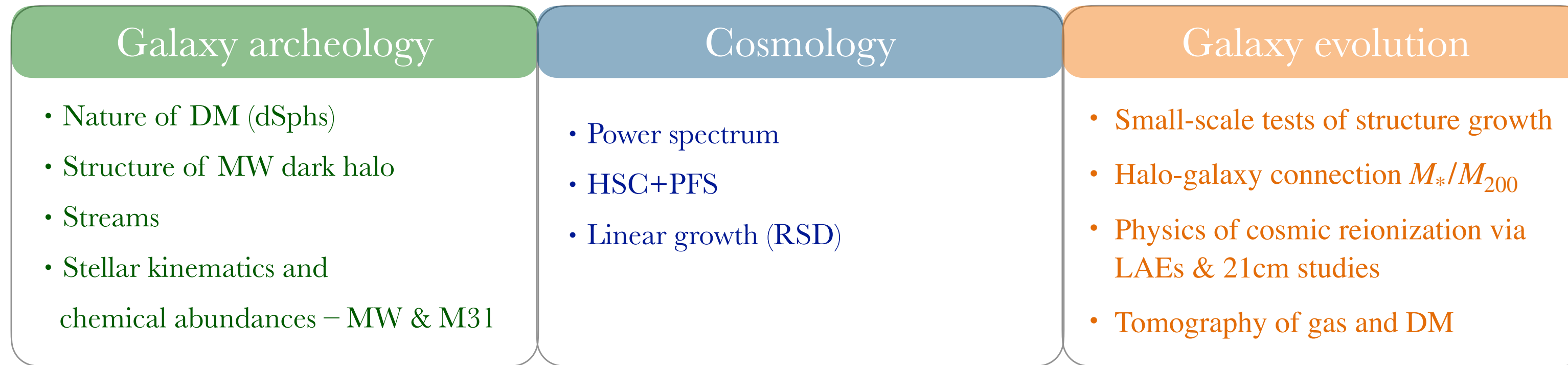
...



Figures by Kohei Hayashi



DM with PFS



- Science with dwarf galaxies

Core:

- Presence of a core or not (slope)
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 - Transition radius

FDM

SIDM

ULA

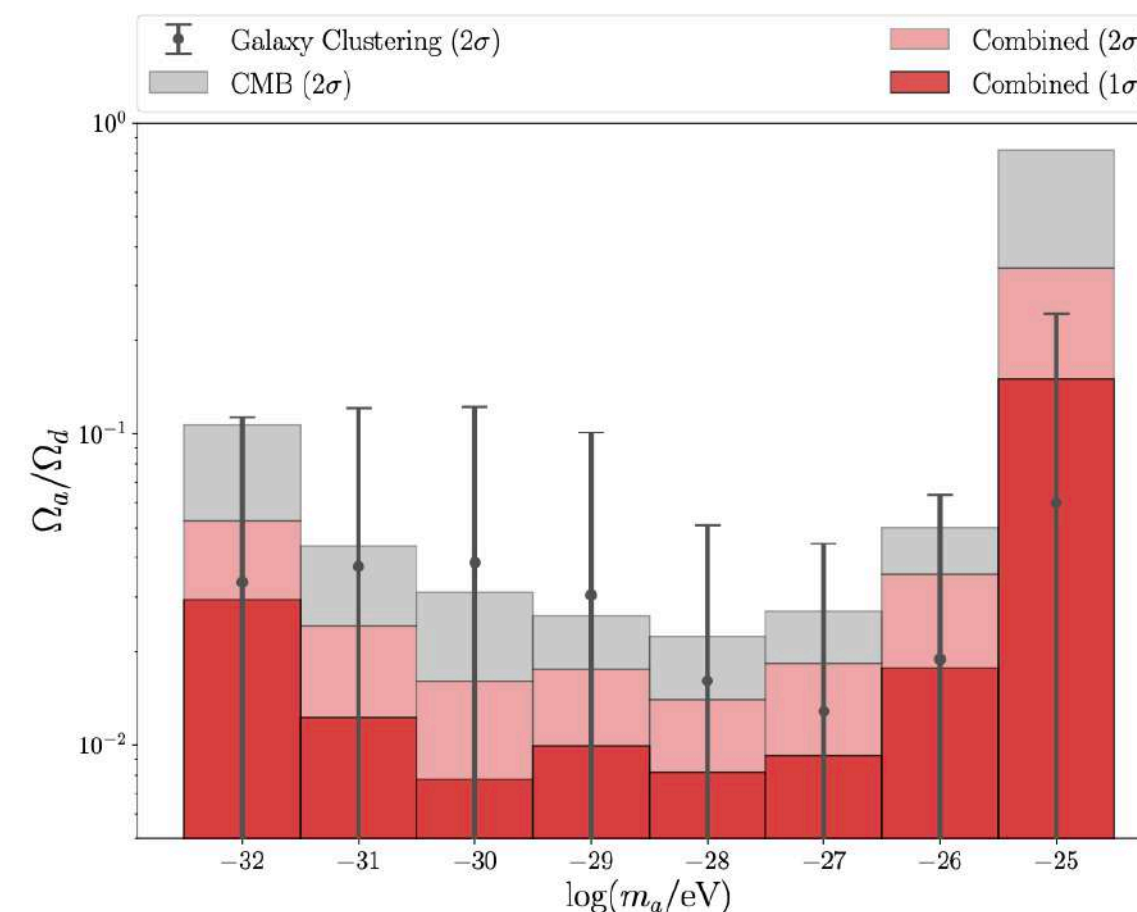
- Abundance data to understand the role of baryons in each system

- Beyond the core

- Granules: heating of stars (dwarfs)
 - Angular momentum
- Stellar streams
- Binary stars

...

Fraction of axions in the dark sector:
 $10^{-32} \text{ eV} < m < 10^{-25} \text{ eV}$



Lague et al 2021

The small-scale Ly- α forest power spectrum

ULA

Halo mass function

FDM

WDM

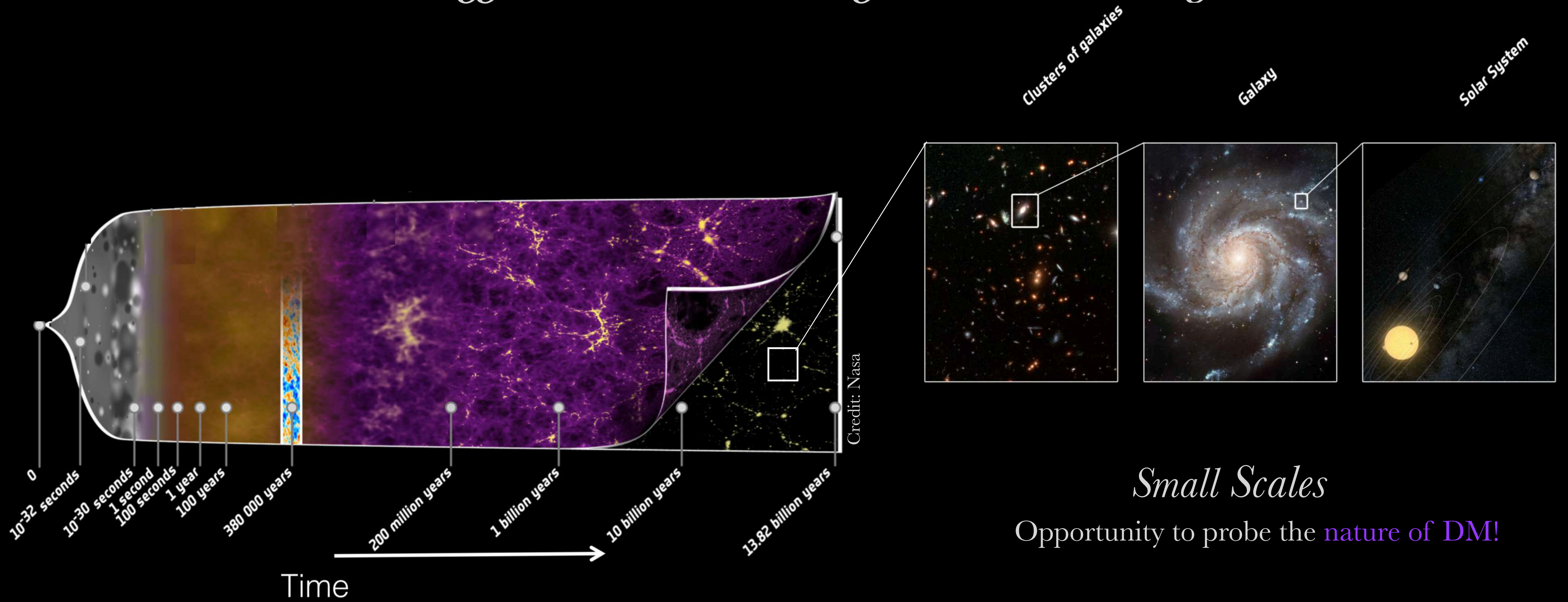
SIDM

Constraints on the optical depth:

Constraint the ULDM mass

Kinematic Sunyaev–Zel’dovich effect: sensitive to the duration of the reionization

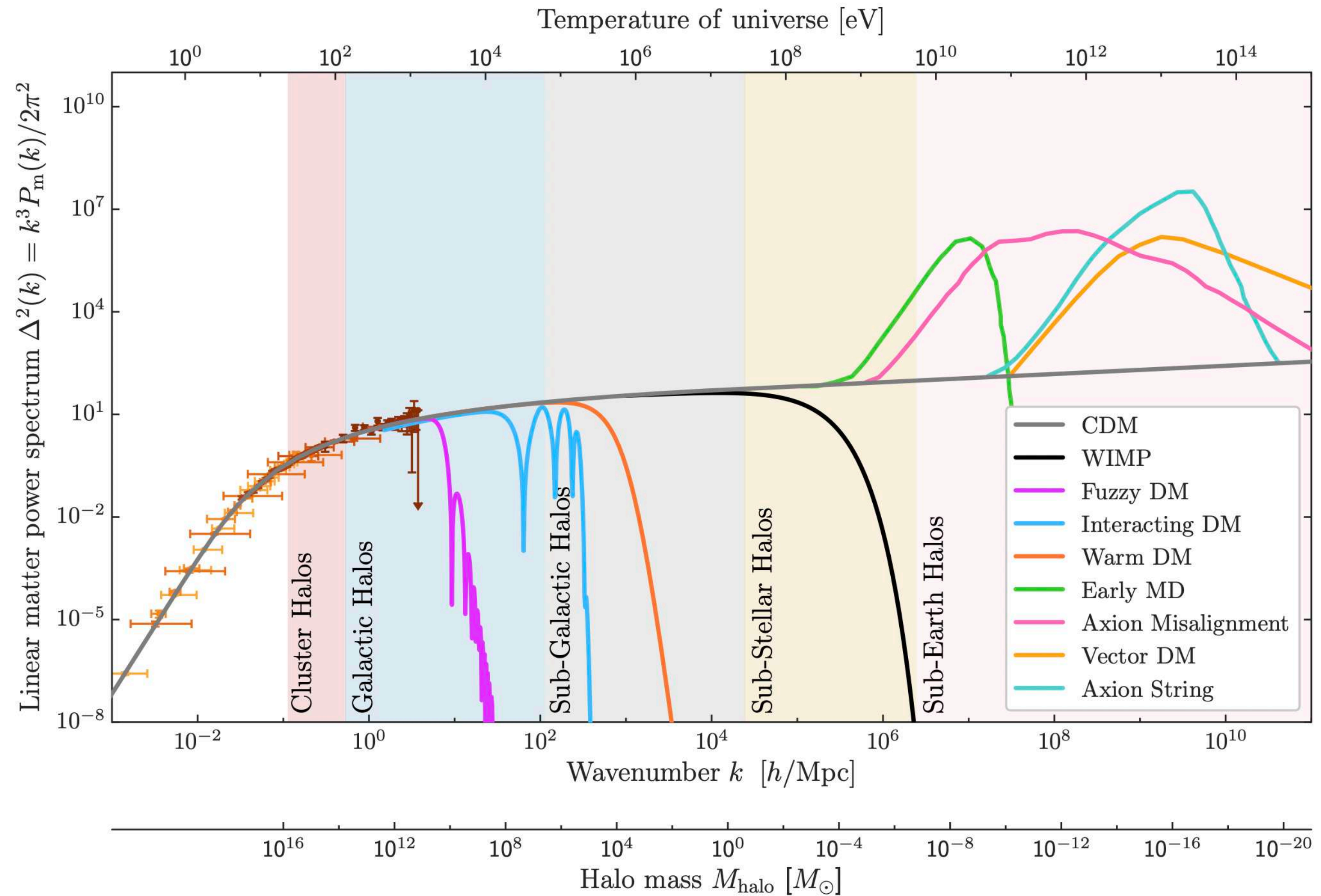
Small scales can offer some *hints* of the nature of DM



Small Scales
Opportunity to probe the *nature of DM!*

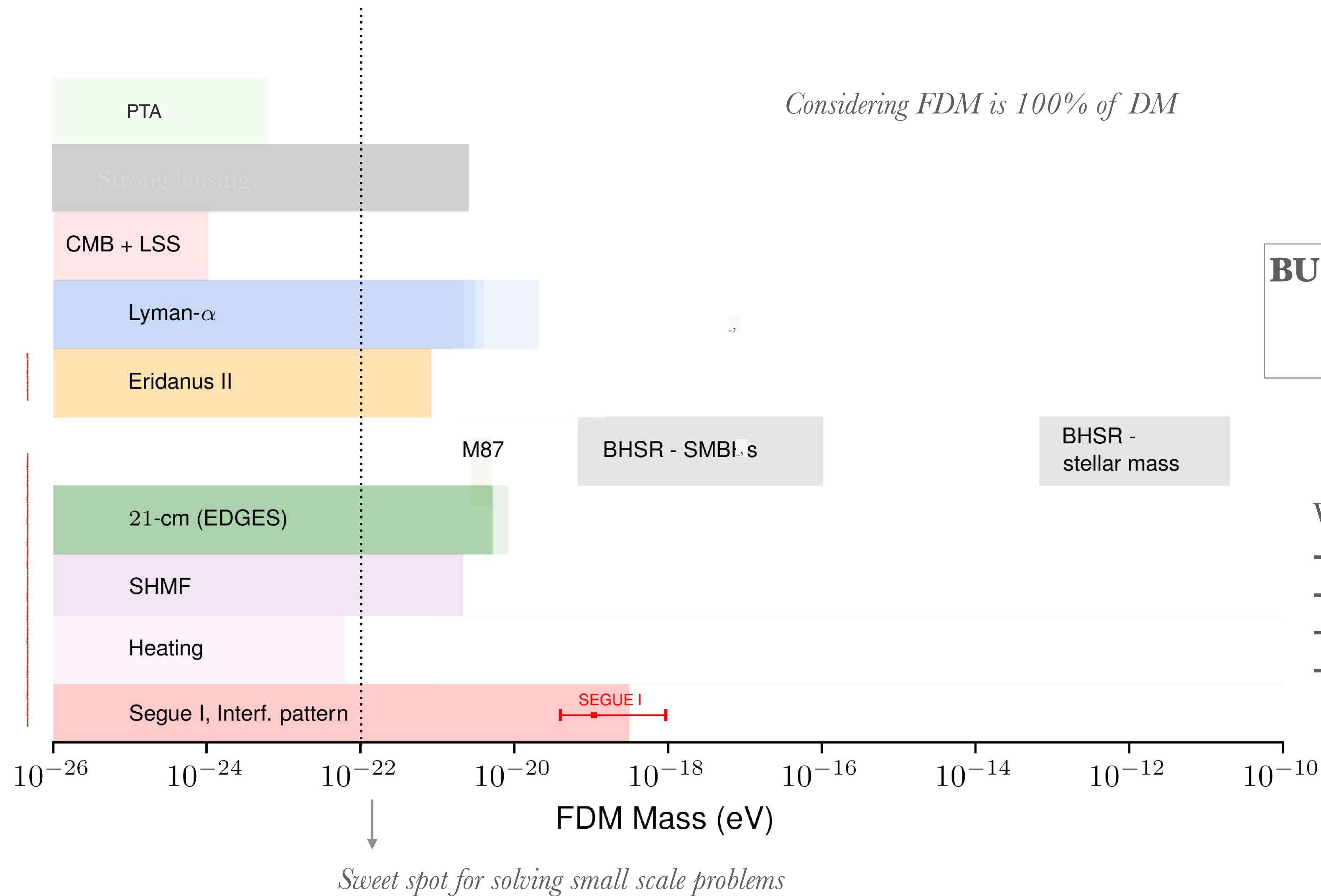
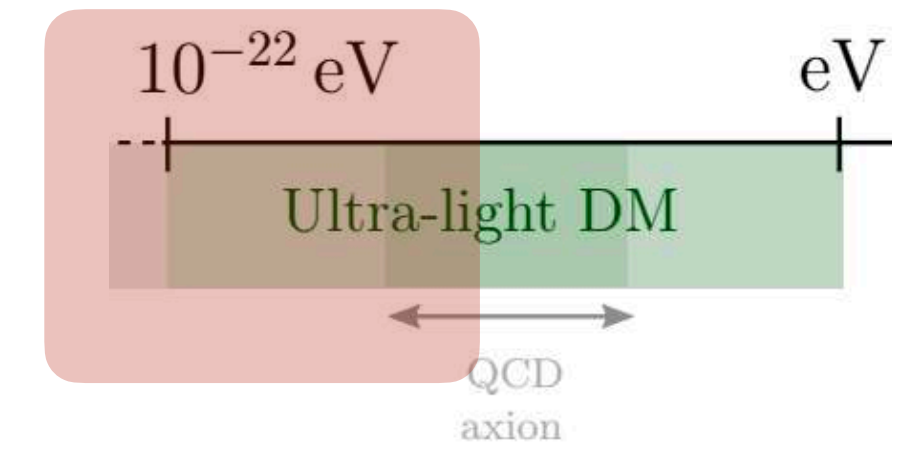


Small scales can offer some *hints* of the nature of DM



Current status

Fuzzy Dark Matter - bounds on the mass



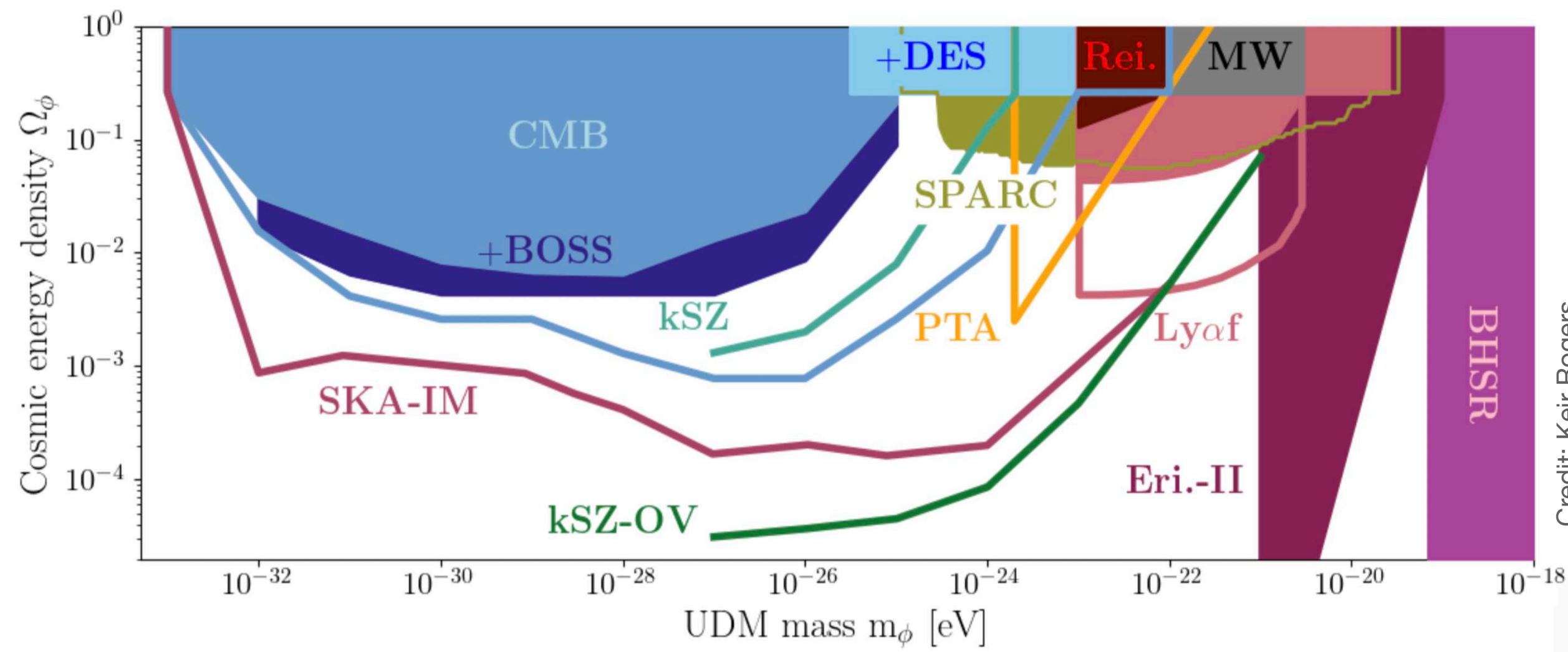
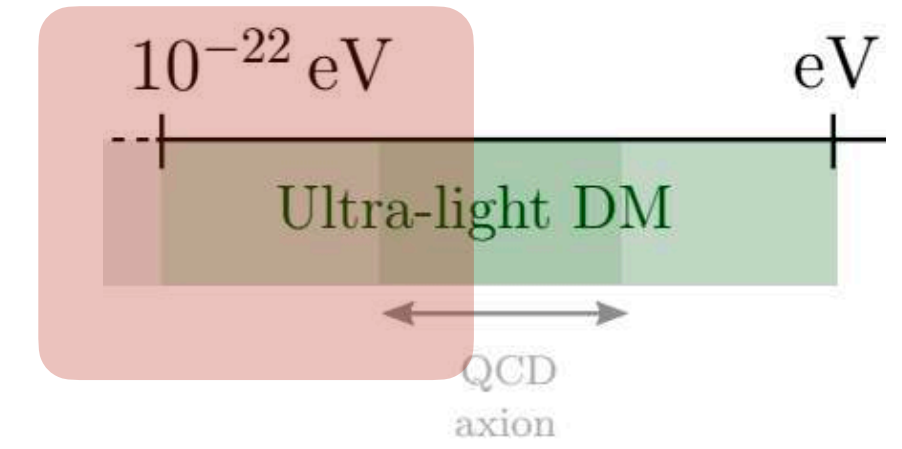
BUT: - systematic effects!!
 - dynamics of FDM not fully understood.

What if:
 - DM not 100%
 - Mixed CDM+ULDM
 - More than one field - axiverse
 - ...

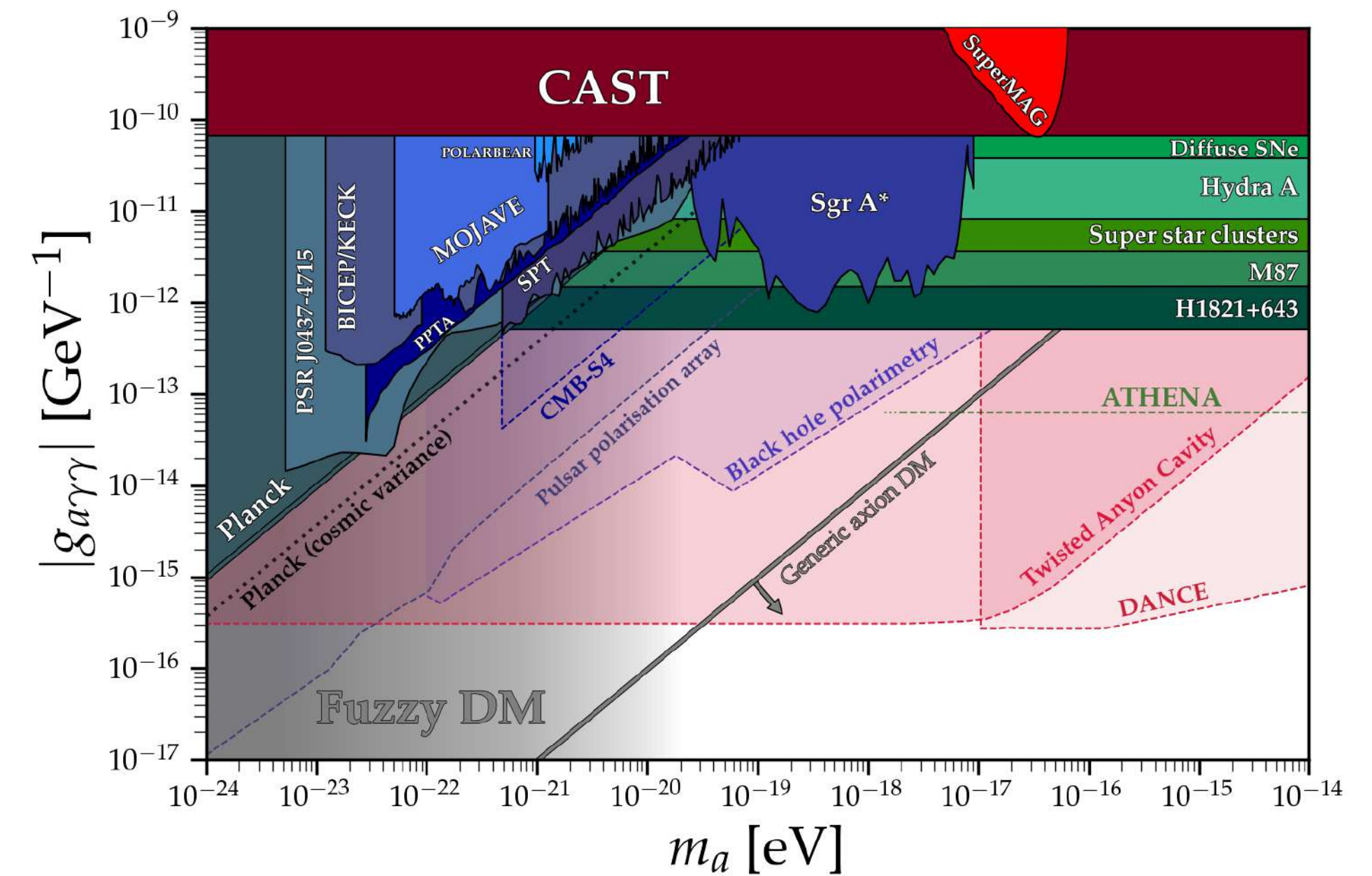
These bounds would change!

Current status

Fuzzy Dark Matter - bounds on the mass



Credit: Keir Rogers



*The search for **dark matter (ULDM)** is a
multi-probe/multi-scale endeavour...*

*exciting times for **axion dark matter***

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Congratulations, Hitoshi!

