

Black Holes as a Window into Quantum Gravity

Yasunori Nomuraa

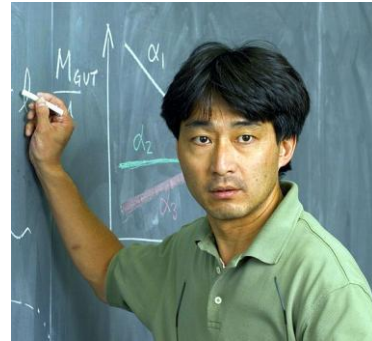
UC Berkeley; LBNL; RIKEN iTHEMS; Kavli IPMU



iTHEMS



My career with Hitoshi



~25 years, 6 papers



Compact Supersymmetry Hitoshi Murayama (UC, Berkeley and LBL, Berkeley and Tokyo U., IPMU), Yasunori Nomura (UC, Berkeley and LBL, Berkeley), Satoshi Shirai (UC, Berkeley and LBL, Berkeley), Kohsaku Tobioka (Unlisted, JP and Tokyo U., IPMU and Tokyo U.) (Jun, 2012) Published in: <i>Phys.RevD</i> 86 (2012) 115014 • e-Print: 1206.4993 [hep-ph]	#1
pdf DOI cite claim reference search 46 citations	
More visible effects of the hidden sector Hitoshi Murayama (UC, Berkeley and LBL, Berkeley), Yasunori Nomura (UC, Berkeley and LBL, Berkeley), David Poland (UC, Berkeley and LBL, Berkeley) (Sep, 2007) Published in: <i>Phys.RevD</i> 77 (2008) 015005 • e-Print: 0709.0775 [hep-ph]	#2
pdf DOI cite claim reference search 112 citations	
Simple Scheme for Gauge Mediation Hitoshi Murayama (LBL, Berkeley and UC, Berkeley), Yasunori Nomura (LBL, Berkeley and UC, Berkeley) (Jan, 2007) Published in: <i>Phys.RevD</i> 75 (2007) 095011 • e-Print: hep-ph/0701231 [hep-ph]	#3
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Gauge Mediation Simplified Hitoshi Murayama (LBL, Berkeley and UC, Berkeley), Yasunori Nomura (LBL, Berkeley and UC, Berkeley) (Dec, 2006) Published in: <i>Phys.Rev.Lett.</i> 98 (2007) 151803 • e-Print: hep-ph/0612186 [hep-ph]	#4
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Wilson lines and symmetry breaking on orbifolds Lawrence J. Hall (UC, Berkeley and LBL, Berkeley), Hitoshi Murayama (UC, Berkeley and LBL, Berkeley), Yasunori Nomura (UC, Berkeley and LBL, Berkeley) (Jul, 2001) Published in: <i>Nucl.Phys.B</i> 645 (2002) 85-104 • e-Print: hep-th/0107245 [hep-th]	#5
pdf DOI cite claim reference search 134 citations	
Viable ultraviolet insensitive supersymmetry breaking Nima Arkani-Hamed (UC, Berkeley and LBL, Berkeley), David E. Kaplan (Chicago U., EFI and Argonne), Hitoshi Murayama (UC, Berkeley), Yasunori Nomura (UC, Berkeley) (Dec, 2000) Published in: <i>JHEP</i> 02 (2001) 041 • e-Print: hep-ph/0012103 [hep-ph]	#6
pdf DOI cite claim reference search 107 citations	

Two pillars of modern physics

Quantum mechanics

$$i\hbar \frac{\partial}{\partial t} |\Psi\rangle = \mathcal{H} |\Psi\rangle$$

General relativity

$$R_{\mu\nu} - \frac{1}{2} g_{\mu\nu} R = \frac{8\pi G_N}{c^4} T_{\mu\nu}$$



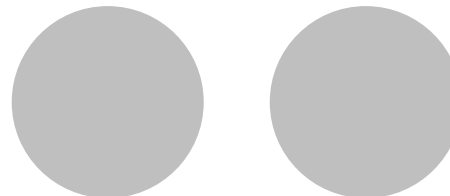
photos: Nobel Foundation archive

not get along well

No problem in “usual” circumstances (below, $c = \hbar = 1$)

p e^-
• •

$$G_N \frac{m_p m_e}{r^2} \sim 10^{-39} \frac{e^2}{4\pi\epsilon_0 r^2}$$



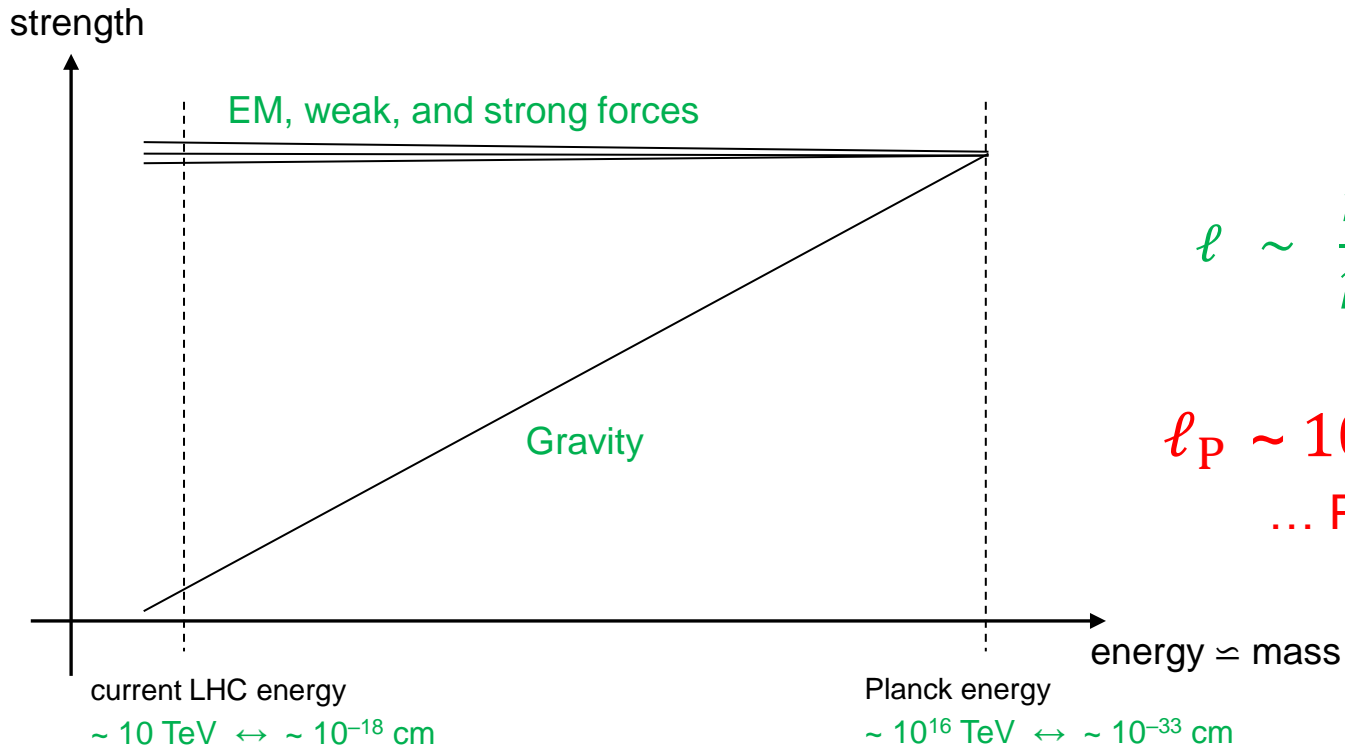
macroscopic body
($N \gg 1$ particles)

quantum interference $\sim \epsilon^N \ll 1$

→ A “patchwork” is enough.

Interesting things occur in “unusual” situations

cf. $v \gg v_{\text{usual}}$ in Newtonian mechanics



$$\ell \sim \frac{\hbar}{p} \sim \frac{\hbar}{E/c}$$

$$\ell_P \sim 10^{-33} \text{ cm}$$

... Planck length

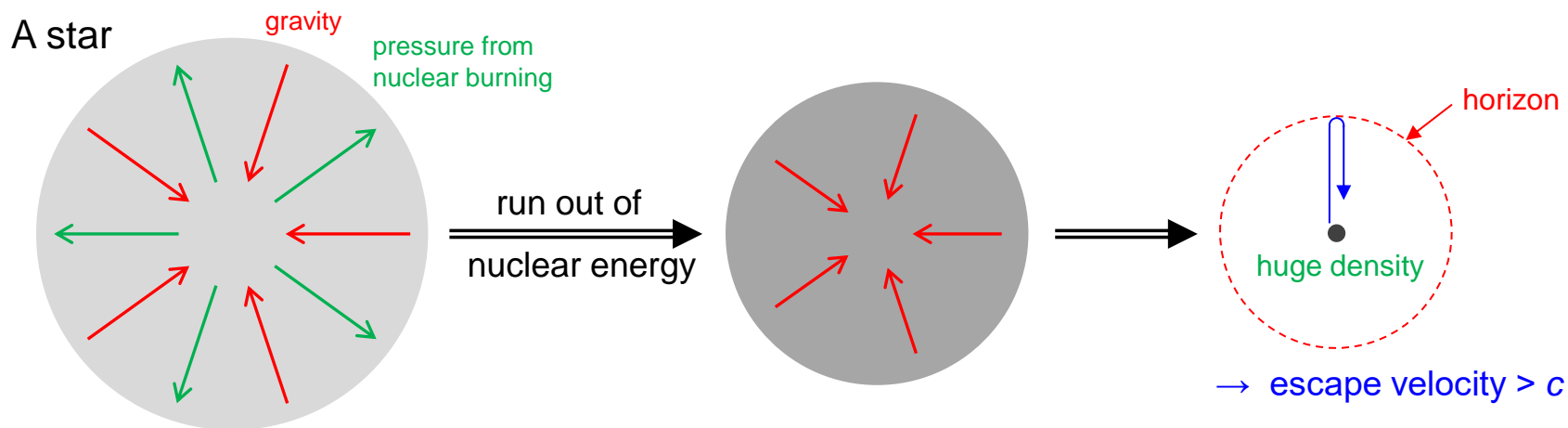
At $\sim \ell_P$, theoretical control of quantum field theory
 (point particles in continuous spacetime) is lost.

→ string theory

There is a problem that the current formulation
 of string theory cannot (directly) address. (today's theme)

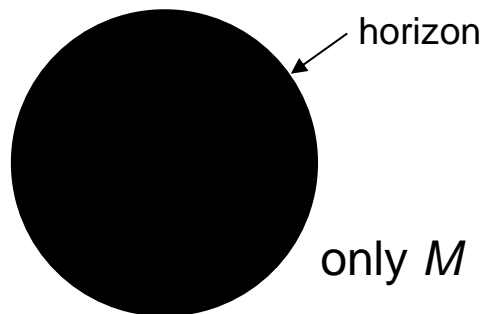
Unusual situations can occur at long distances:

Black Holes



The interior cannot be seen *from the outside*.

— “No hair” theorem



⇒ What are the quantum properties?

Black Hole Thermodynamics

A puzzle

Another pillar of modern physics

Statistical mechanics

$$S = k_B \ln W$$

(below, $k_B = 1$)

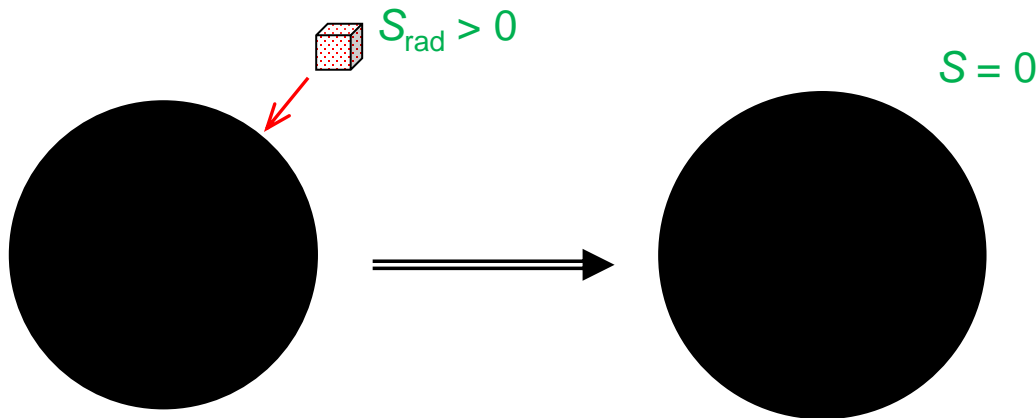


photo: Univ. Frankfurt

for any system
for all practical purposes

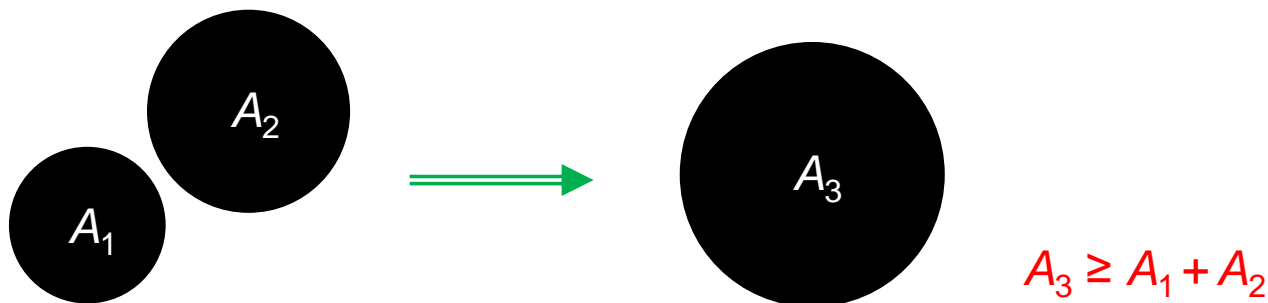
$$\Delta S \geq 0$$

What happens if matter falls into a black hole?



... $\Delta S < 0$!?

A peculiar property of BHs in general relativity



A proposal [Bekenstein, 1973]



photo: APS

The entropy of a BH is proportional to its horizon area.

$$S_{\text{BH}} = \frac{A}{4G_{\text{N}}}$$

Note: $G_{\text{N}} = \ell_{\text{P}}^2 \sim (10^{-33} \text{ cm})^2 \rightarrow$ huge entropy
 e.g. A solar mass BH has $S \sim 10^{78}$ while the sun has $\sim 10^{60}$.

Indeed, $\Delta \left(\frac{A}{4G_{\text{N}}} + S_{\text{matter}} \right) \geq 0$

Does this make sense?

$$\frac{A}{4G_{\text{N}}} = 4\pi G_{\text{N}} M^2 \rightarrow \frac{\partial S}{\partial E} = \frac{1}{T} \rightarrow \text{finite temperature}$$

Doesn't a BH only absorb stuff?

Black holes radiate [Hawking, 1974]



photo: NASA

The horizon is “smooth.”



Quantum mechanical effect

There must be radiation corresponding to $T_H \sim \frac{1}{8\pi M G_N}$.

Hawking temperature

Black holes radiate [Hawking, 1974]



photo: NASA

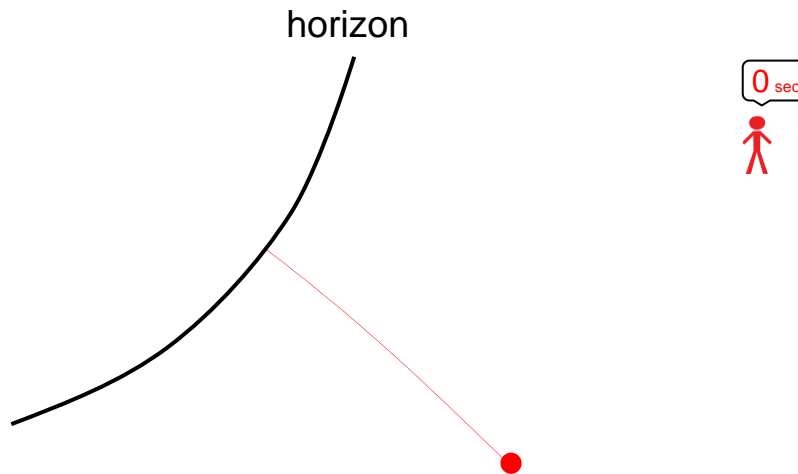
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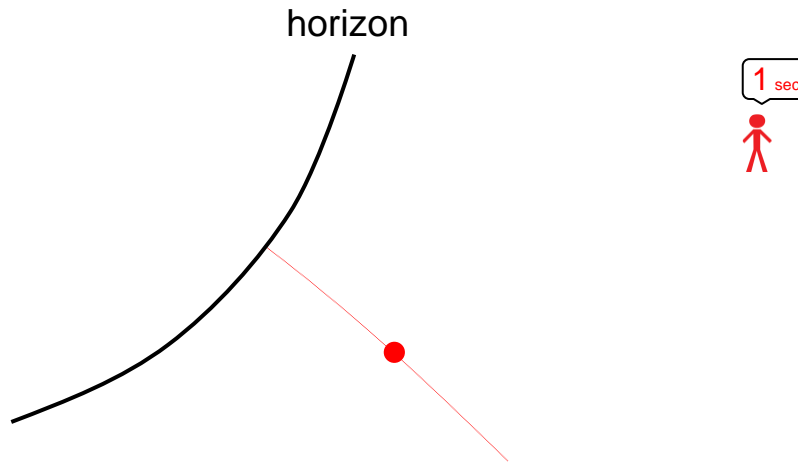
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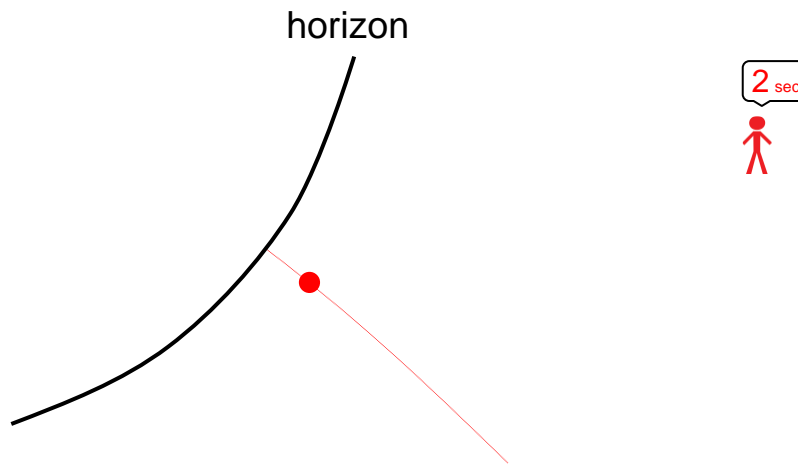
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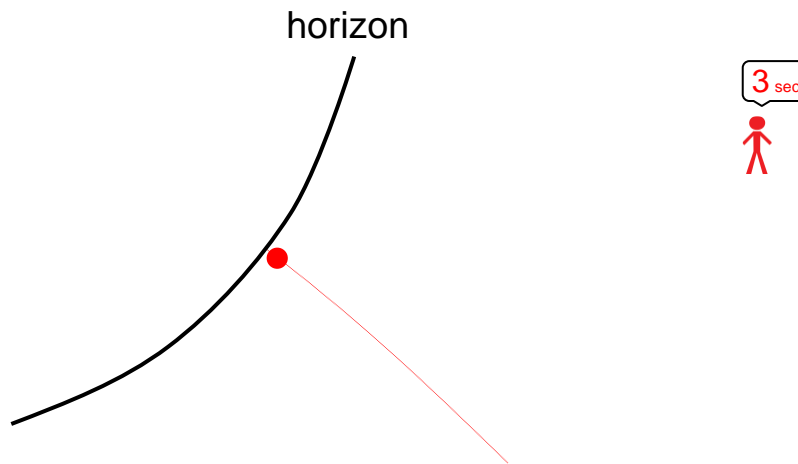
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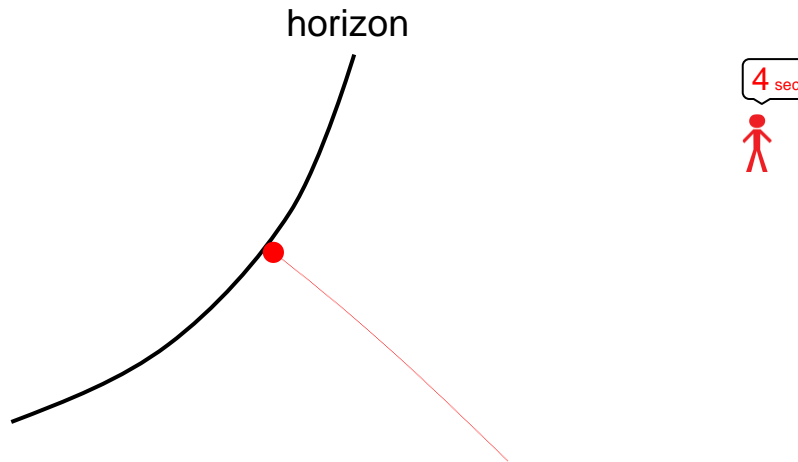
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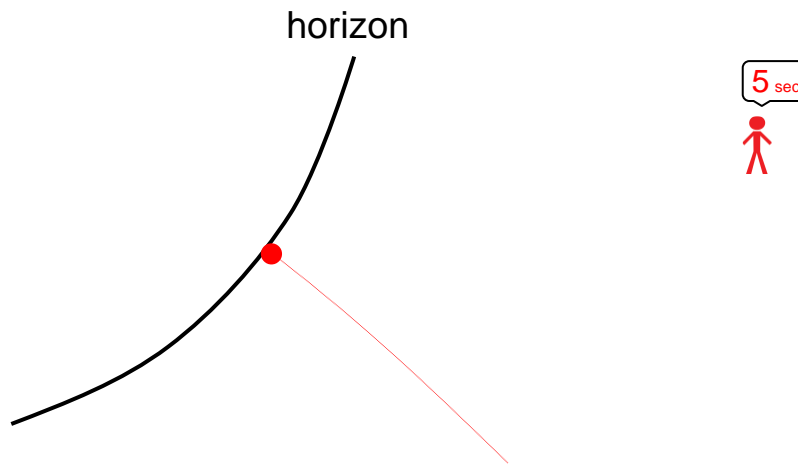
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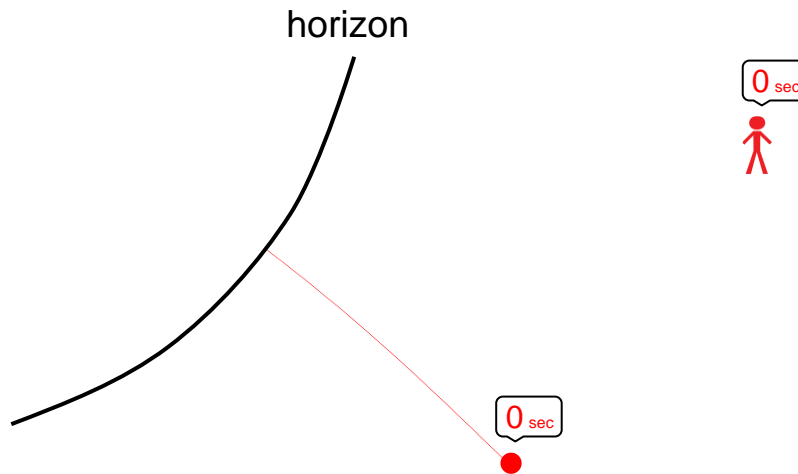
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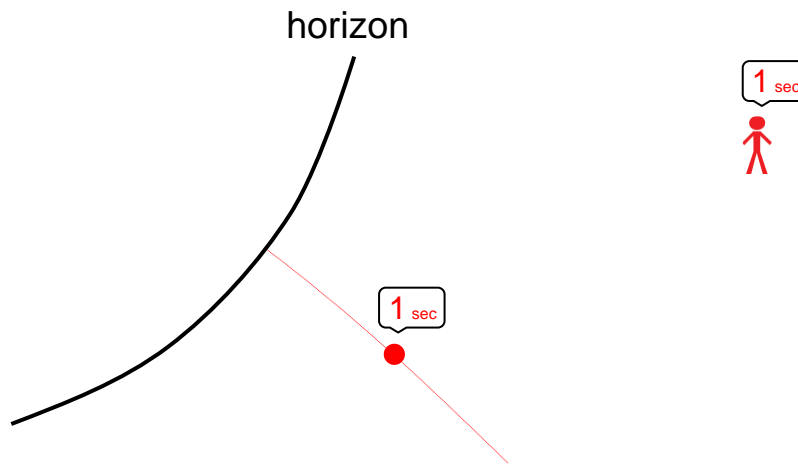
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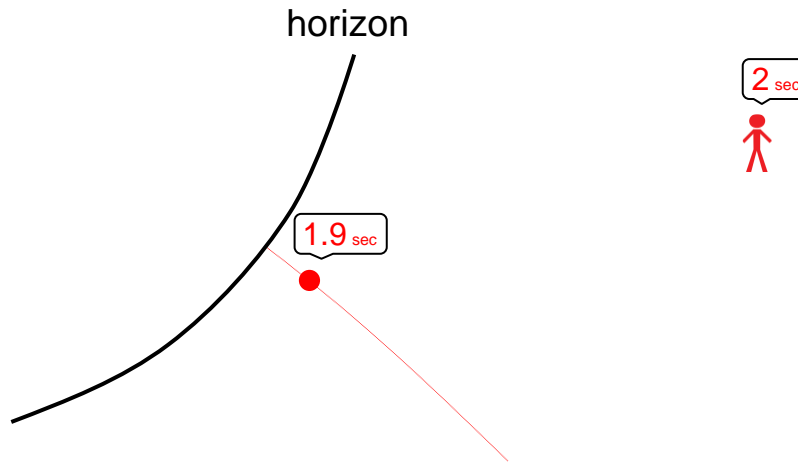
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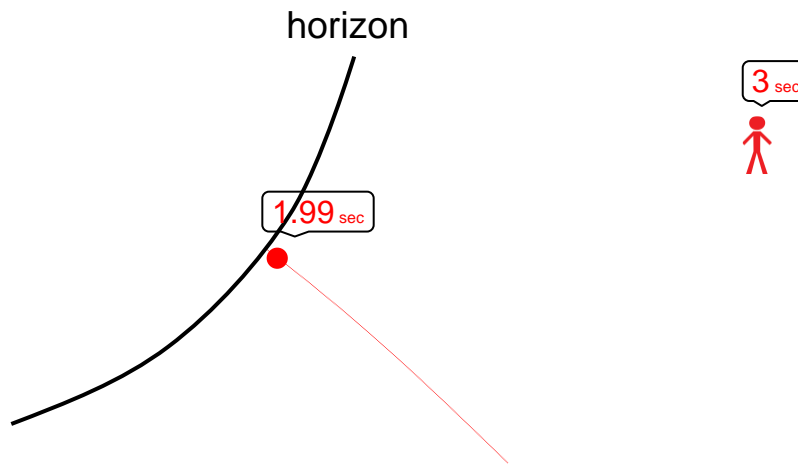
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photo: NASA

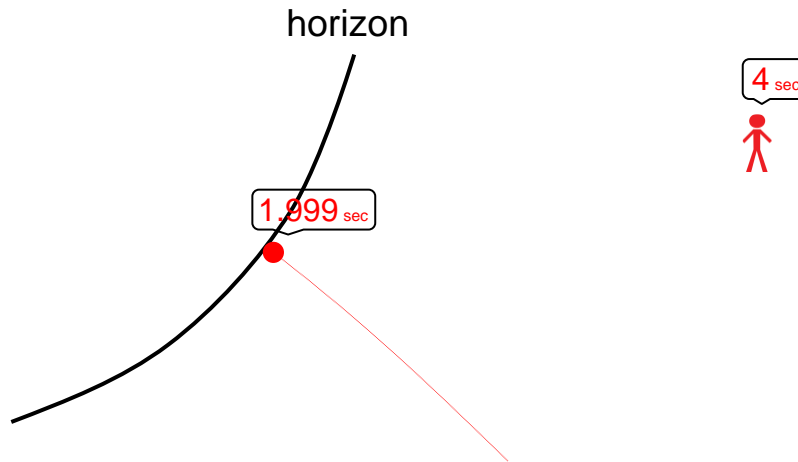
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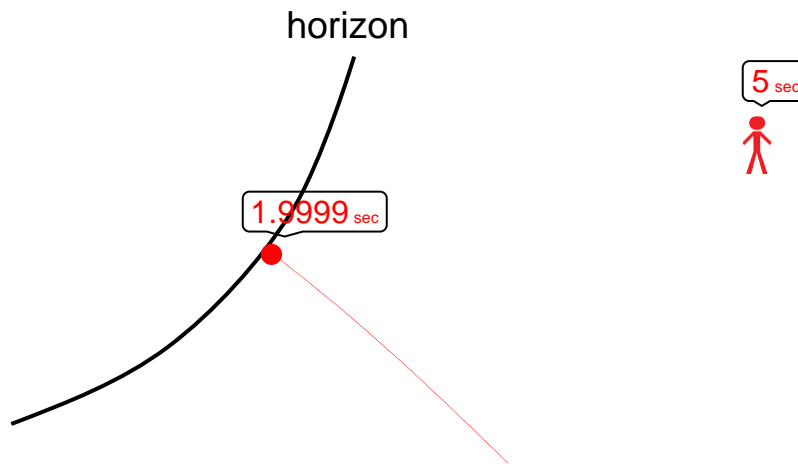
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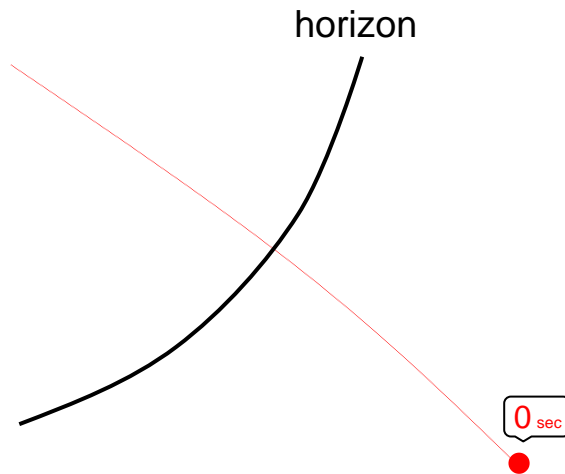
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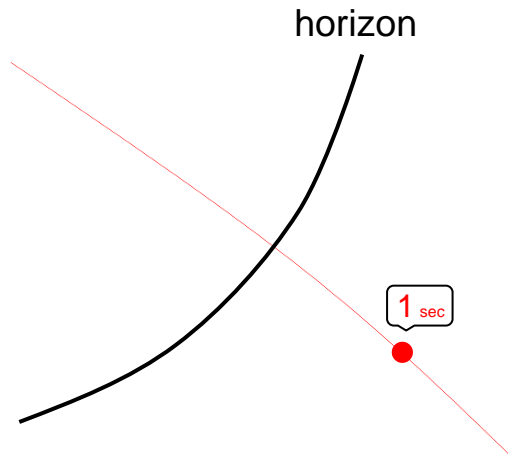
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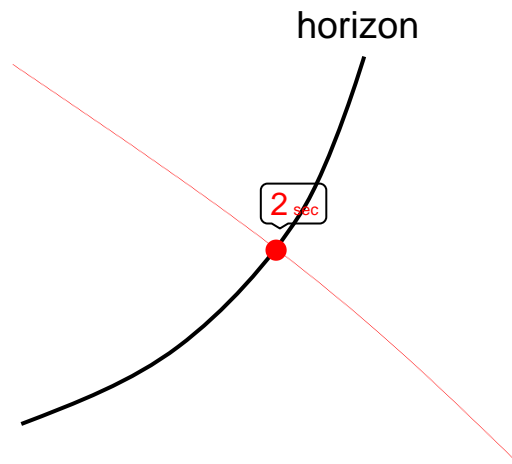
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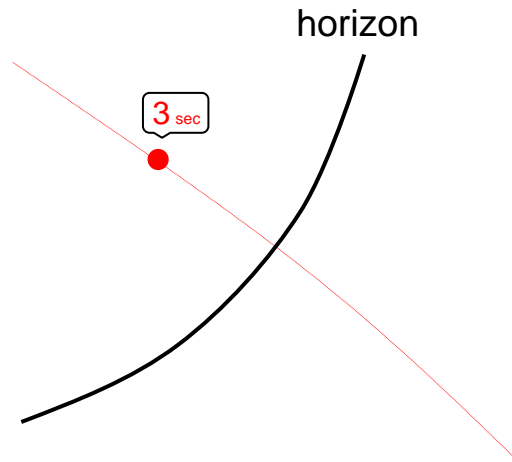
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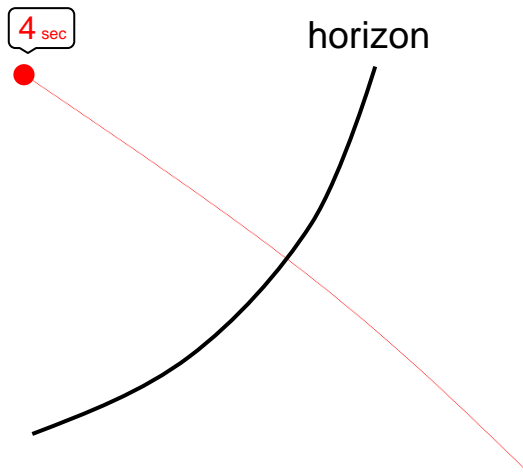
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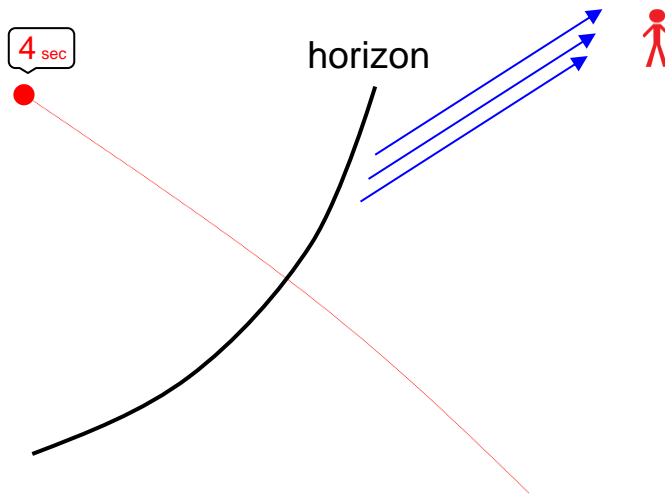
Quantum mechanical effect

There must be radiation corresponding to $T_H \sim \frac{1}{8\pi M G_N}$.

Hawking temperature

BHs are thermodynamic objects.

→ Spacetime is composed of microscopic d.o.f.s!



Black holes radiate [Hawking, 1974]



photo: NASA

The horizon is “smooth.”



Quantum mechanical effect

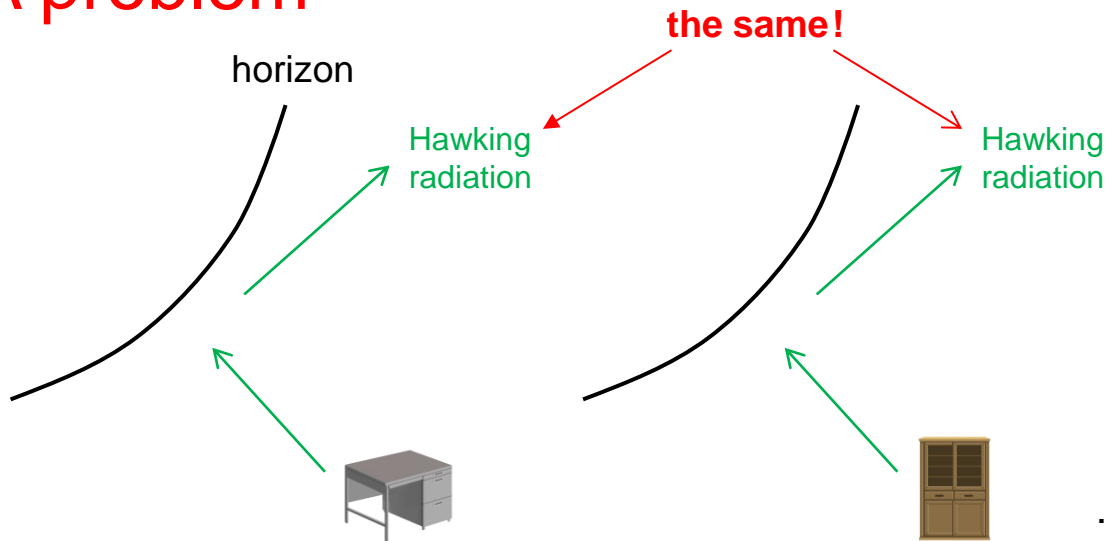
There must be radiation corresponding to $T_H \sim \frac{1}{8\pi M G_N}$.

Hawking temperature

BHs are thermodynamic objects.

→ Spacetime is composed of microscopic d.o.f.s!

A problem



The time evolution
is **not** one-to-one!
(not unitary)

... (the original form of)
BH information problem

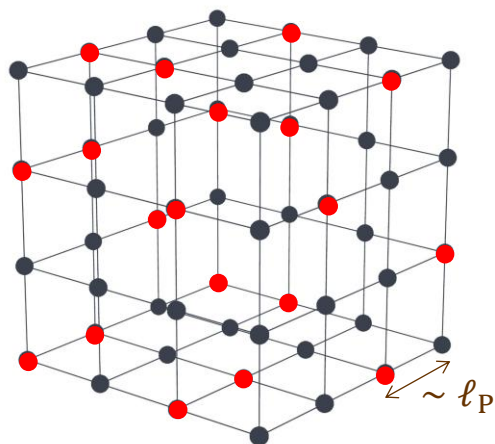
Holography

A clue comes from the BH physics itself

A BH is the highest entropy state of the region,

$$\text{and still } S \propto \frac{A}{\ell_P^2}$$

Strange!



$$S \sim \ln 2^{V/\ell_P^3} \propto \frac{V}{\ell_P^3} \gg \frac{A}{\ell_P^2}$$

$$(\ell_P \sim 10^{-33} \text{ cm})$$

The concept that spacetime exists down to $\sim \ell_P$ is an illusion!

→ suggests that there is a formulation of quantum gravity
in spacetime **one less dimension** than the naïve one.

Sounds crazy?

AdS/CFT correspondence [Maldacena, 1997]

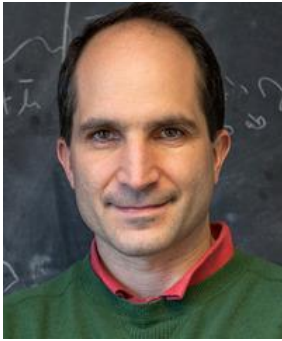
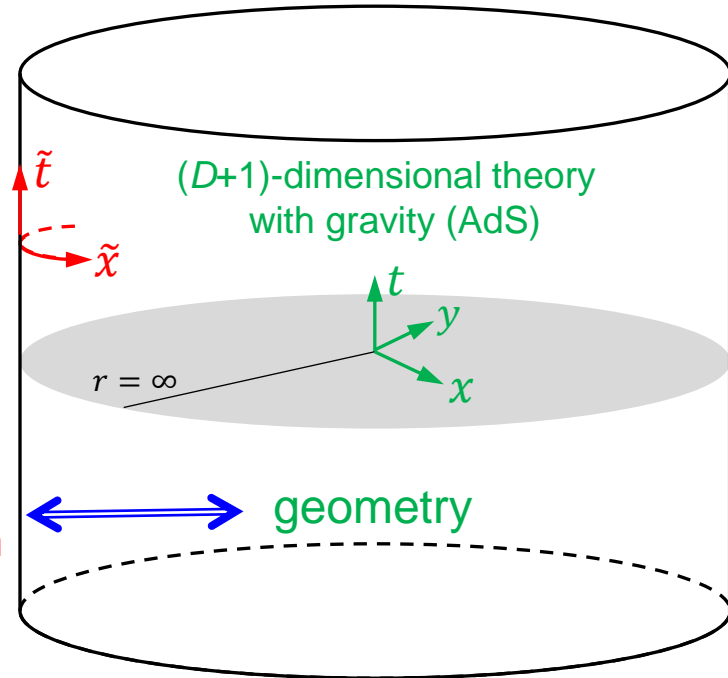


photo: IAS

D -dimensional theory
without gravity (CFT)

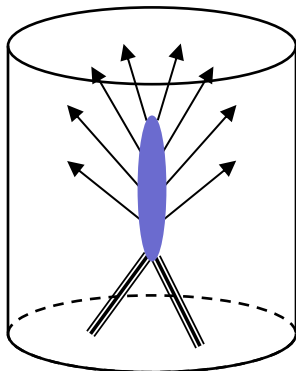


$(D+1)$ -dimensional theory
with gravity (AdS)

quantum
information

geometry

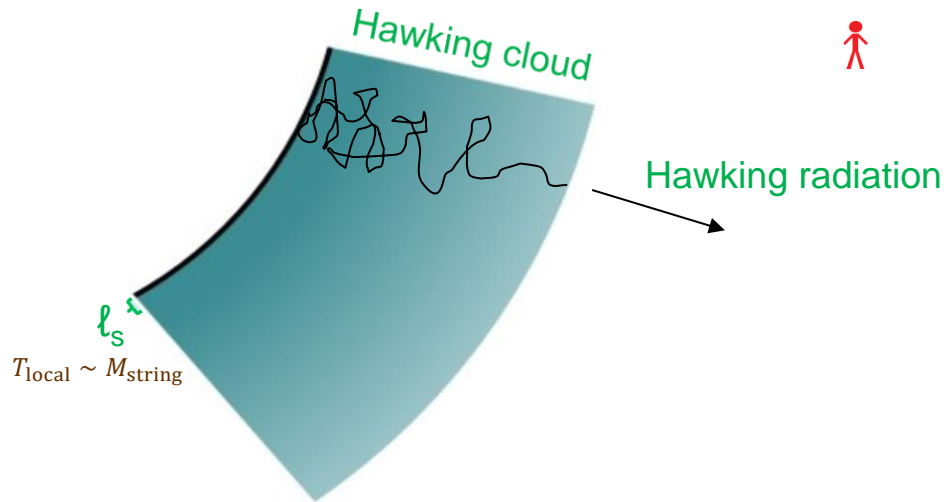
BH evolution **must be** unitary.



=

A process in non-gravitational
(unitary) theory

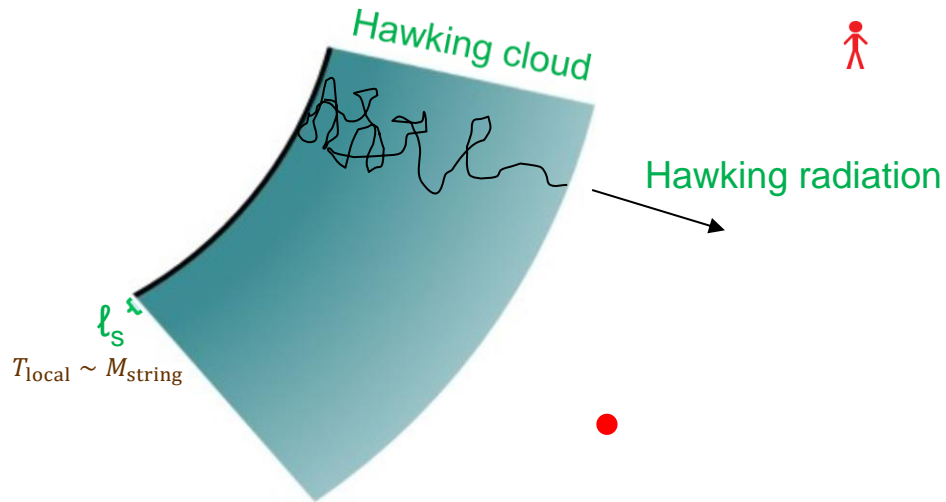
BH at the quantum level



The horizon behaves
as the surface of regular material.

... no issue with unitarity

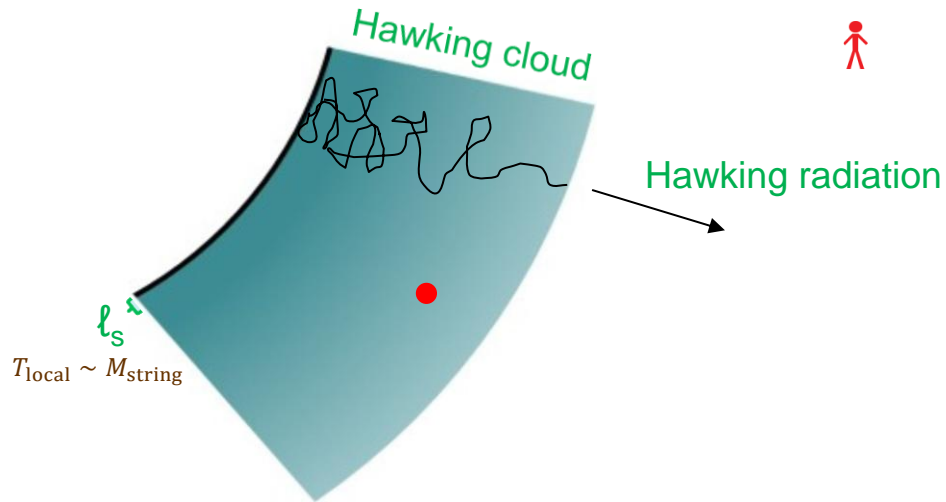
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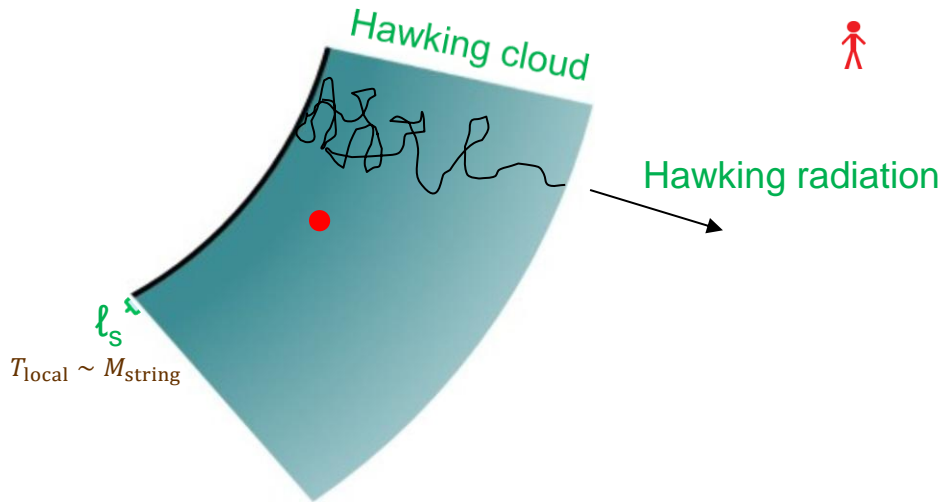
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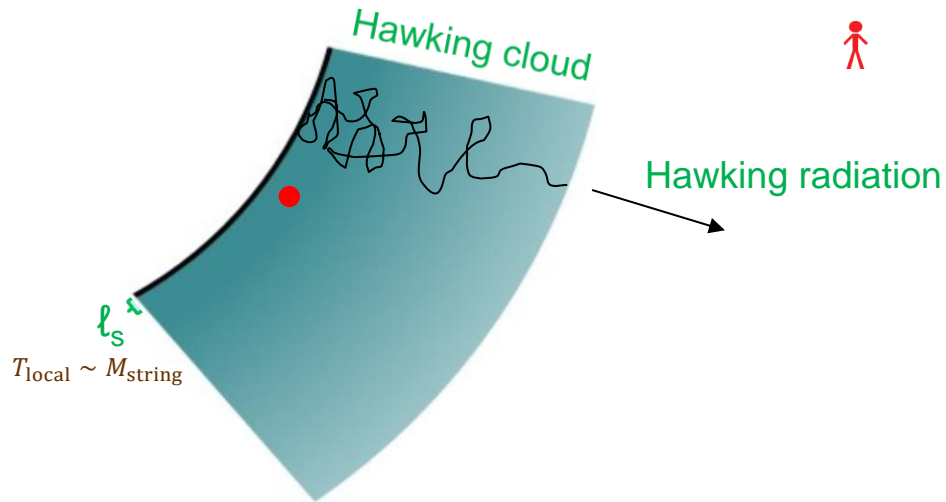
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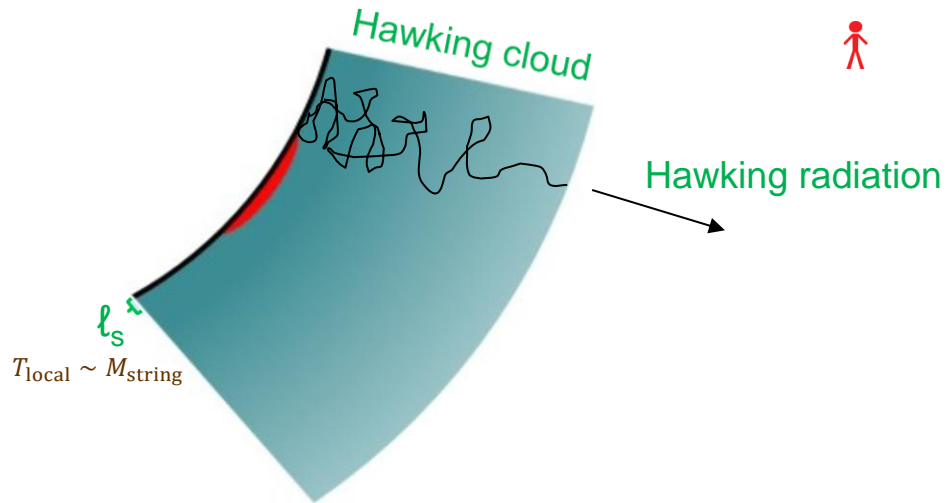
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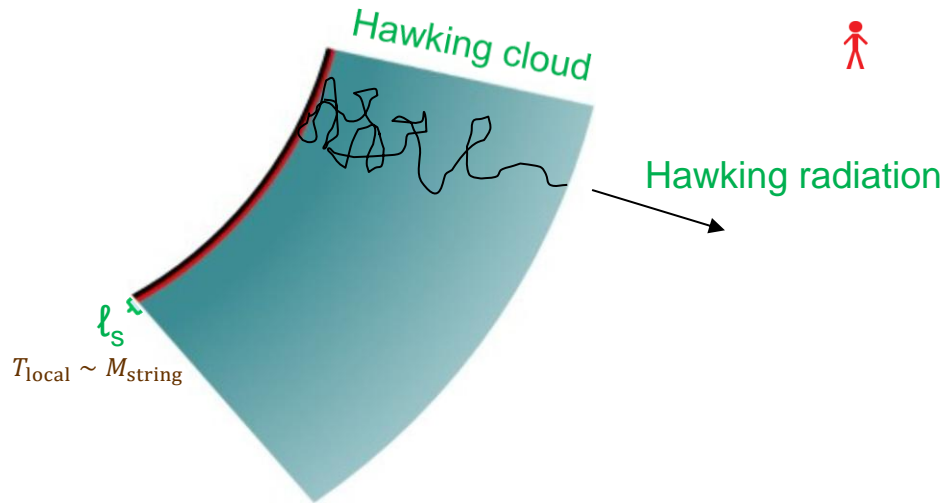
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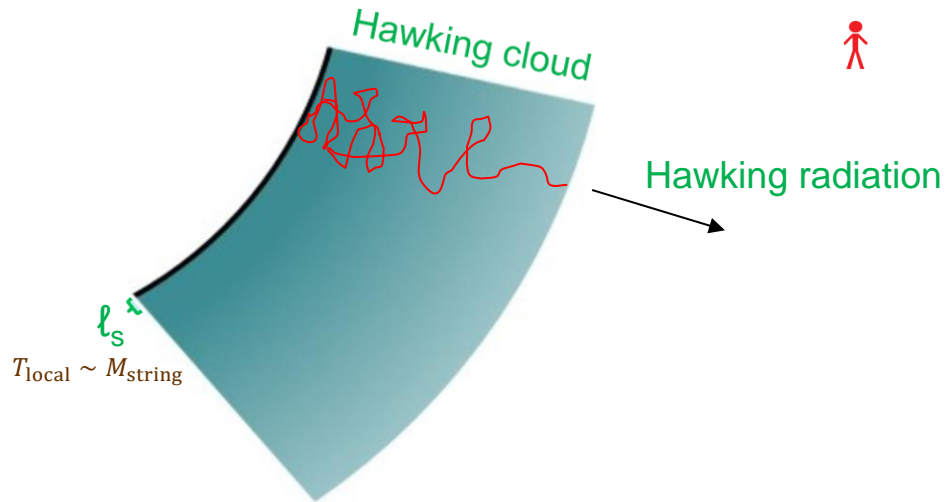
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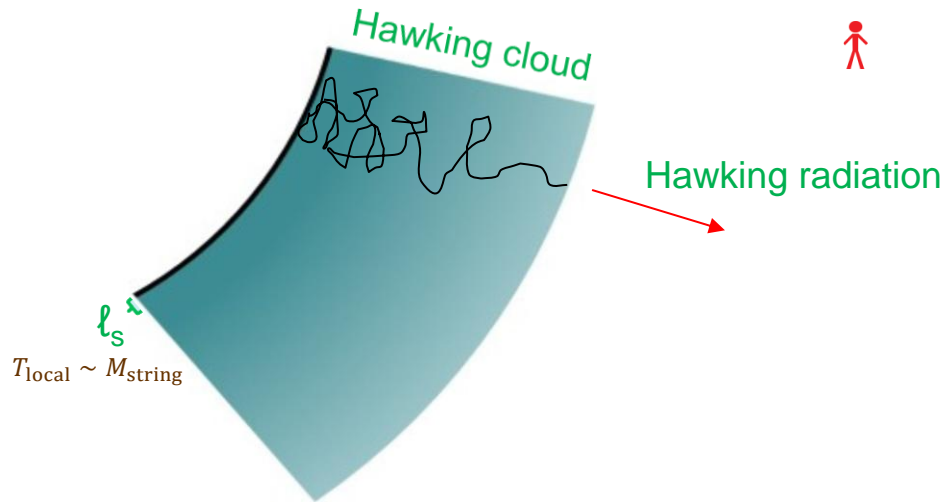
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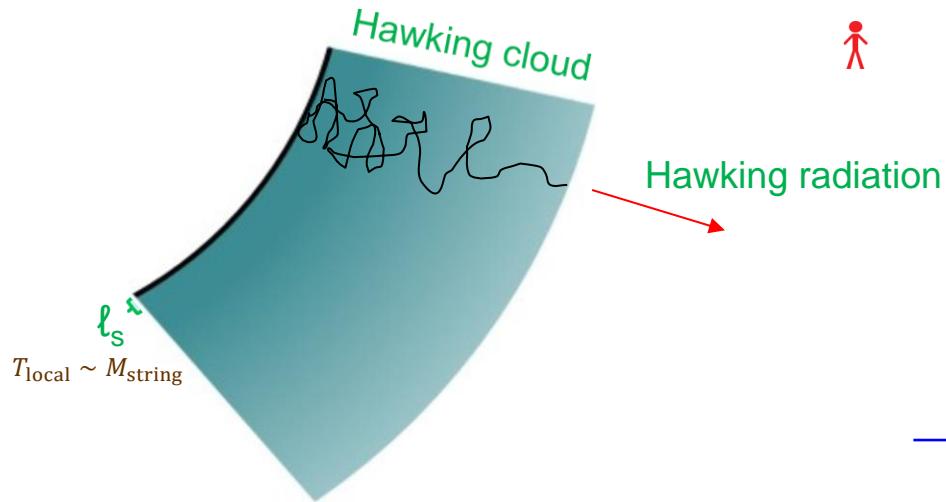
BH at the quantum level



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BH at the quantum level

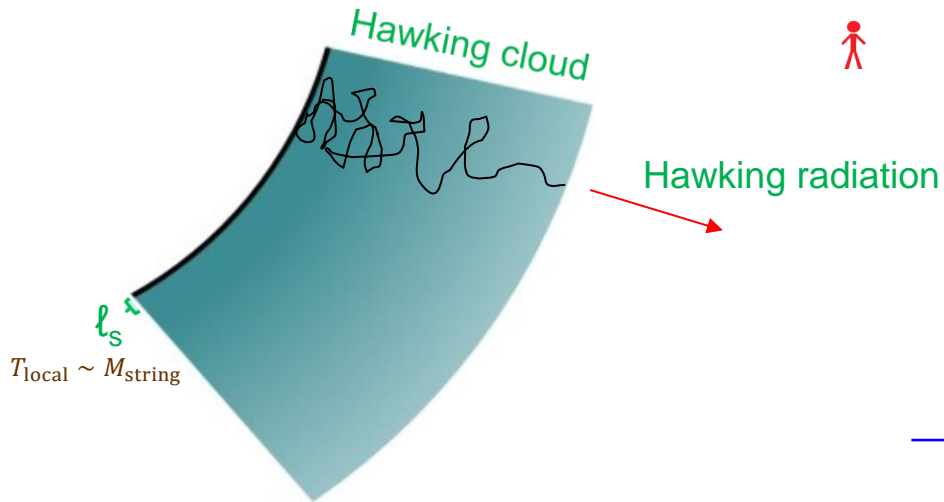


The horizon behaves
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→ What about the interior?

BH at the quantum level

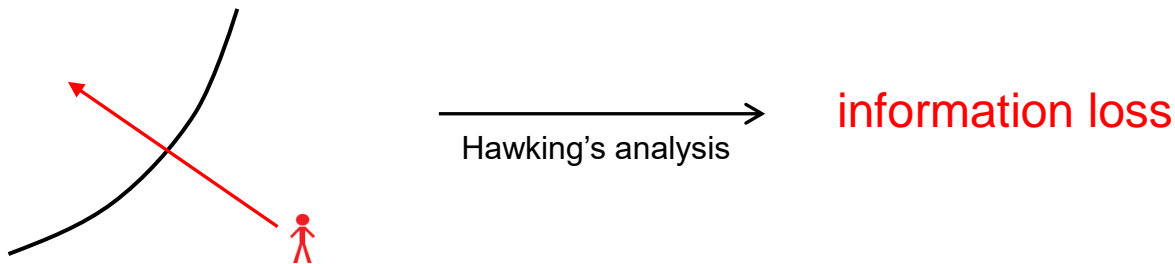


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→ What about the interior?

Alternatively

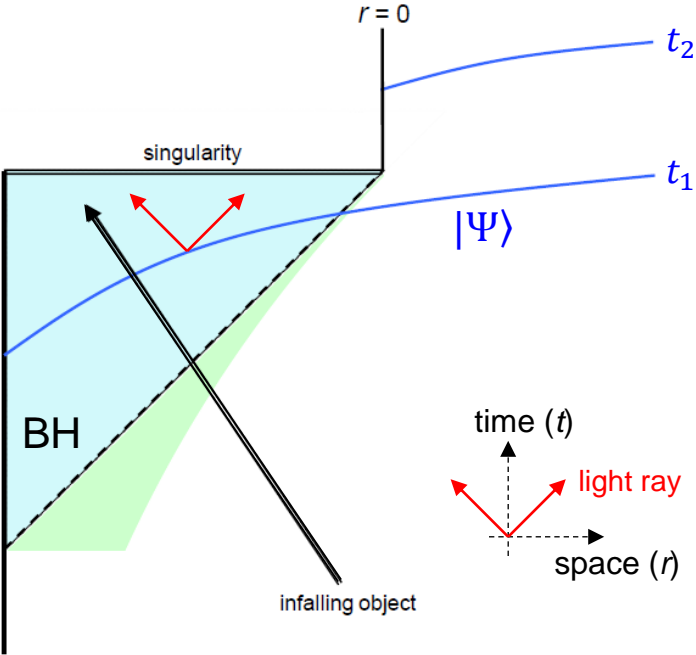


→ What was wrong with Hawking's analysis?

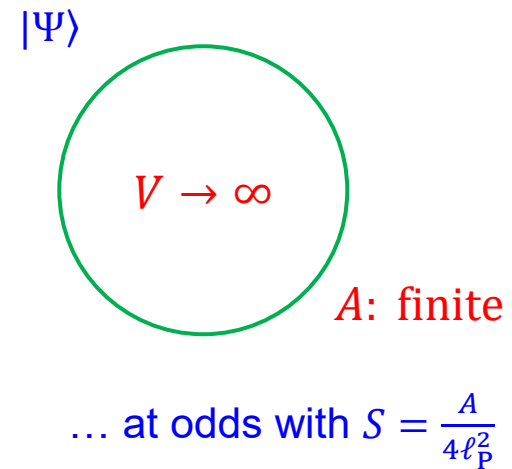
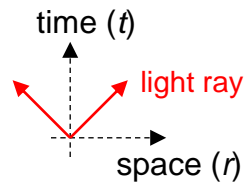
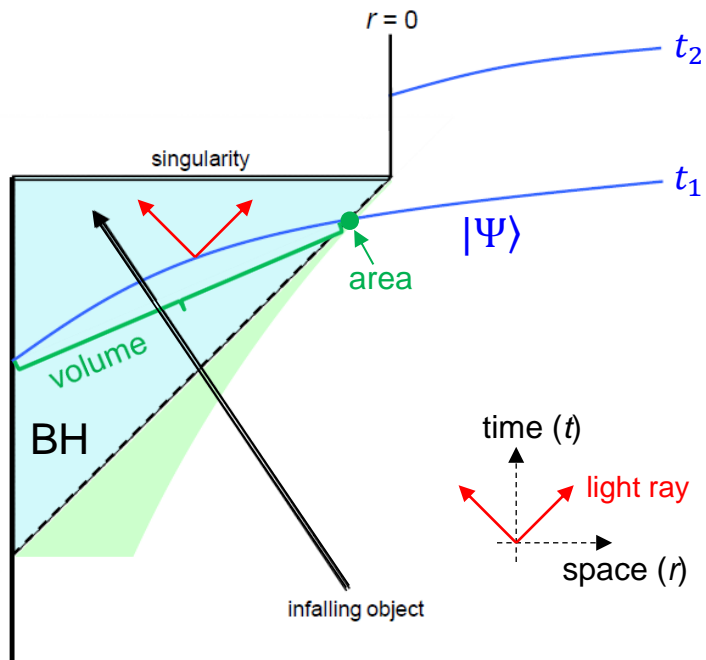
Recent Progress I

— replica wormholes —

Start with “global spacetime”

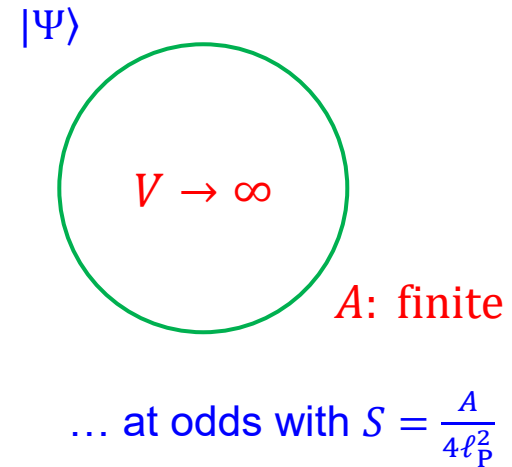
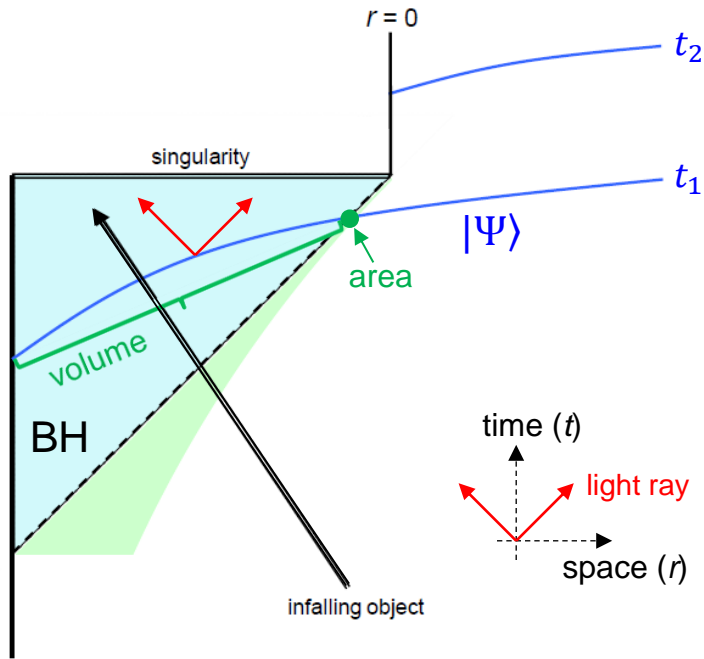


Start with “global spacetime”



Hugely redundant!

Start with “global spacetime”



Hugely redundant!

$$\langle \Psi_1 | \Psi_2 \rangle = 0 \quad \longrightarrow \quad \langle \Psi_1 | \Psi_2 \rangle \sim e^{-\frac{S}{2}}$$

semiclassical
(QFT in curved spacetime)

quantum gravity

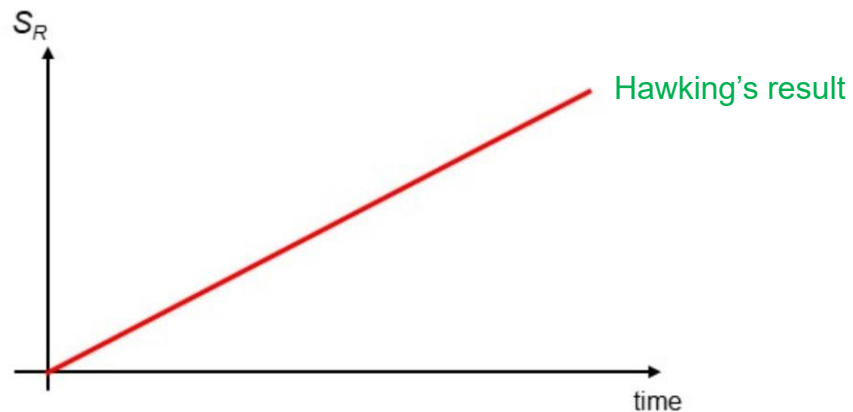
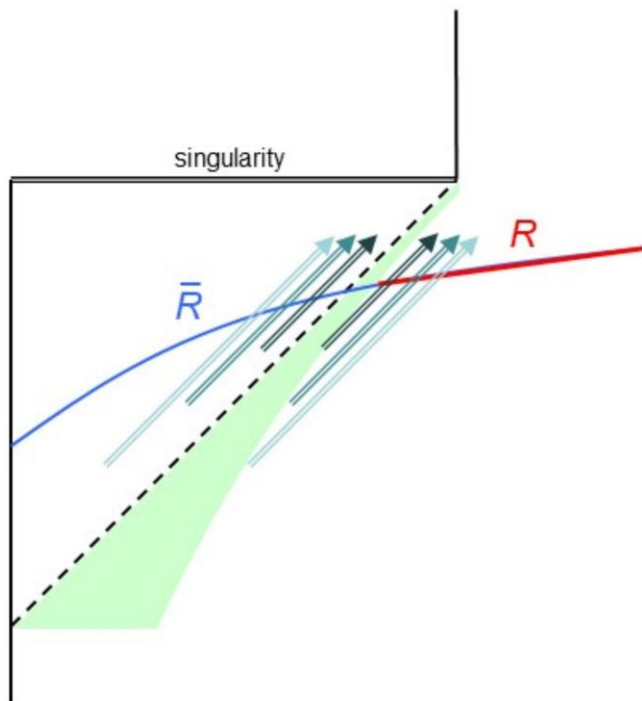
... only e^S independent states

$$|\Psi\rangle = \sum_{i=1}^{e^S} c_i |\psi_i\rangle \quad c_i \sim e^{-\frac{S}{2}}$$

$$\langle \Psi_1 | \Psi_2 \rangle = \sum_{i=1}^{e^S} c_{1,i}^* c_{2,i} \sim e^{\frac{S}{2}} e^{-S} \sim e^{-\frac{S}{2}}$$

$\rightarrow e^{e^S}$ approximately orthogonal states

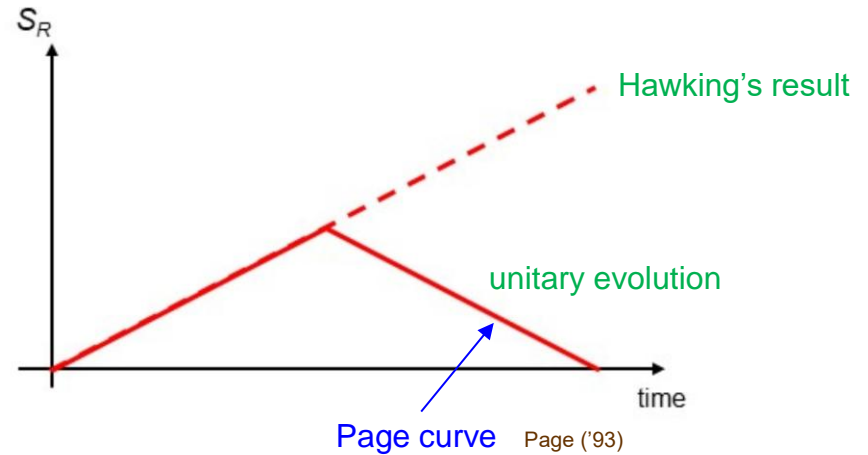
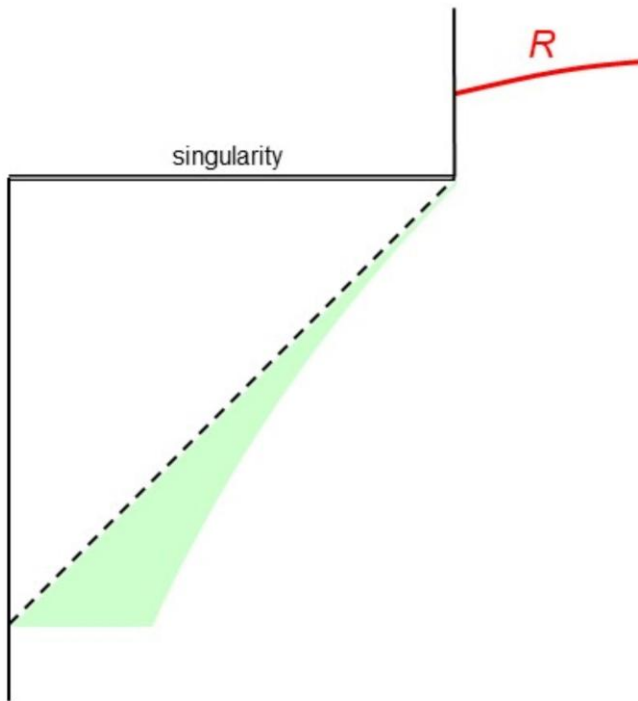
Unitarity of Hawking evaporation



$$S_R = -\text{Tr}[\rho_R \ln \rho_R] \quad (\rho_R = \text{Tr}_{\bar{R}}|\Psi\rangle\langle\Psi|)$$

~ the # of EPR particles in R whose partners are in \bar{R}

Unitarity of Hawking evaporation



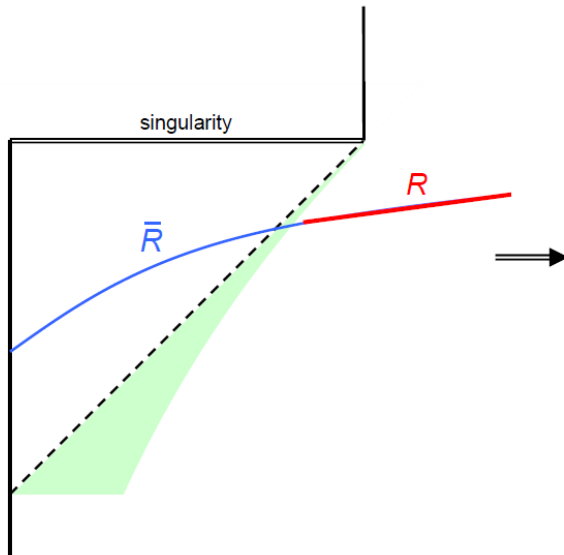
$$S_R = -\text{Tr}[\rho_R \ln \rho_R] \quad (\rho_R = \text{Tr}_{\bar{R}}|\Psi\rangle\langle\Psi|)$$

~ the # of EPR particles in R whose partners are in \bar{R}

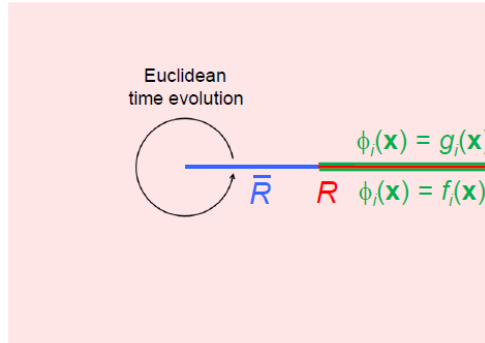
→ How to get this curve?

Page curve from replica wormholes

Penington ('19); Almheiri, Engelhardt, Marolf, Maxfield ('19); ...
 Penington, Shenker, Stanford, Yang ('19);
 Almheiri, Hartman, Maldacena, Shaghoulian, Tajdini ('19)



path integral

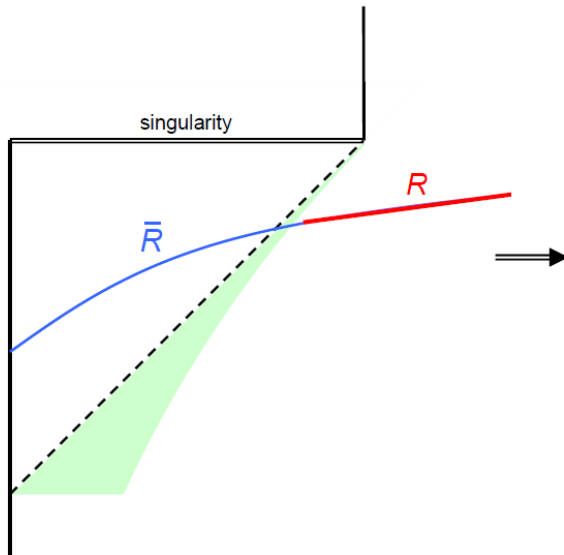


$\rightarrow \rho_R = \rho_R[f_i(\mathbf{x}), g_i(\mathbf{x})]$ (\sim coefficient of $|g_i(\mathbf{x})\rangle\langle f_i(\mathbf{x})|$)

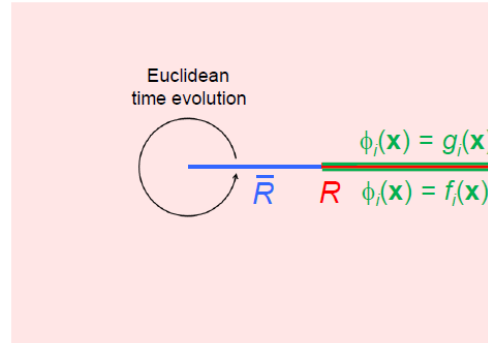
$$S_R \equiv -\text{Tr}[\rho_R \ln \rho_R] = \lim_{n \rightarrow 1} \frac{1}{1-n} \ln \text{Tr}[\rho_R^n]$$

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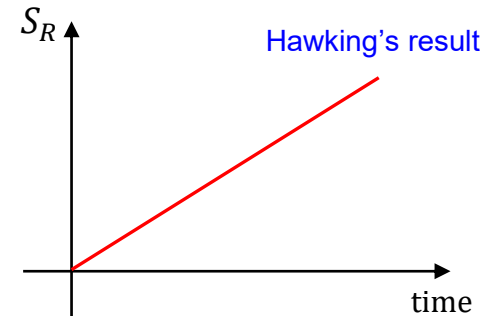
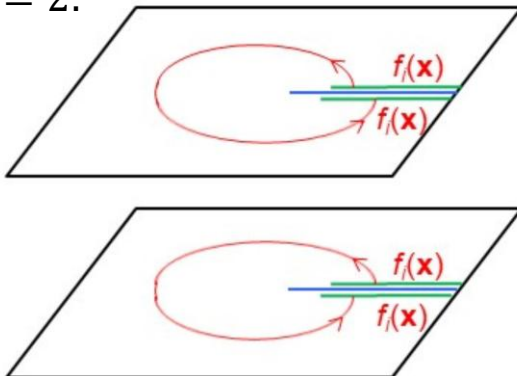
path integral



$$\rightarrow \rho_R = \rho_R[f_i(\mathbf{x}), g_i(\mathbf{x})] \quad (\sim \text{coefficient of } |g_i(\mathbf{x})\rangle\langle f_i(\mathbf{x})|)$$

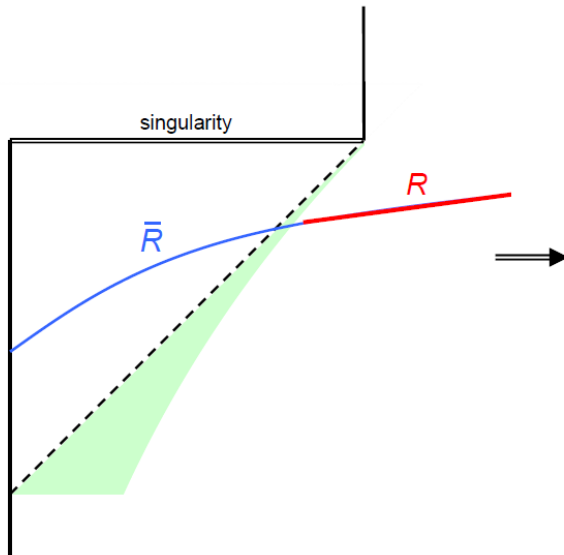
$$S_R \equiv -\text{Tr}[\rho_R \ln \rho_R] = \lim_{n \rightarrow 1} \frac{1}{1-n} \ln \text{Tr}[\rho_R^n]$$

$n = 2$:

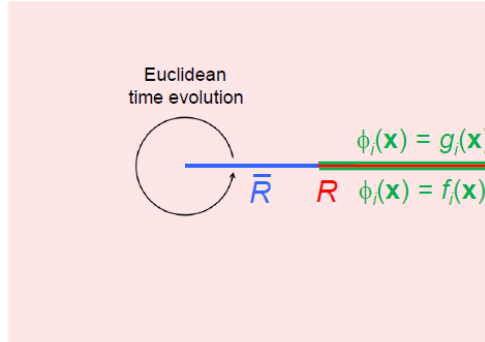


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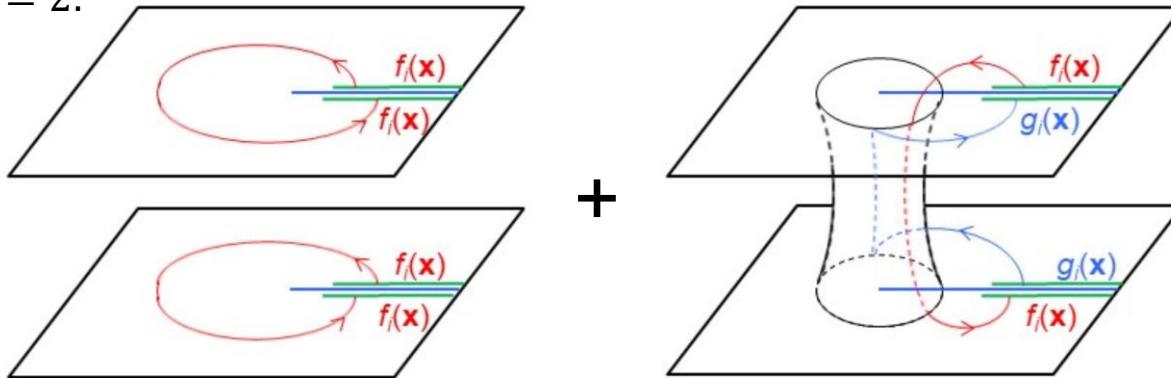
path integral



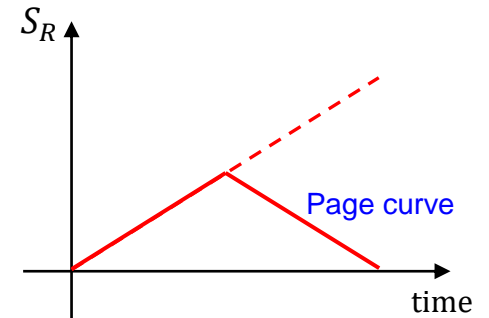
$$\rightarrow \rho_R = \rho_R[f_i(\mathbf{x}), g_f(\mathbf{x})] \quad (\sim \text{coefficient of } |g_f(\mathbf{x})\rangle\langle f_i(\mathbf{x})|)$$

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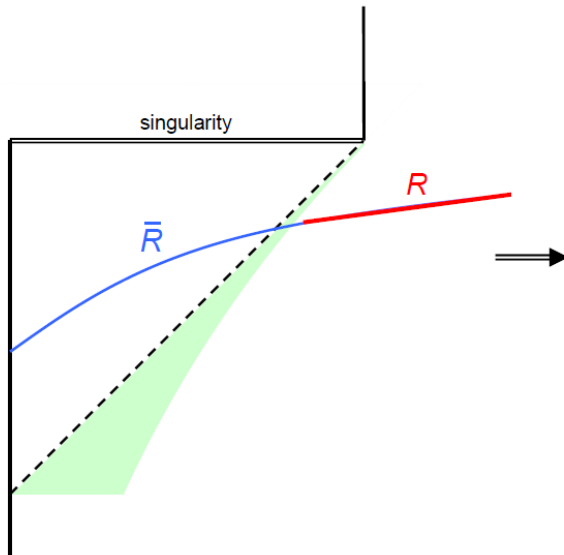


replica wormhole (nonperturbative effect)

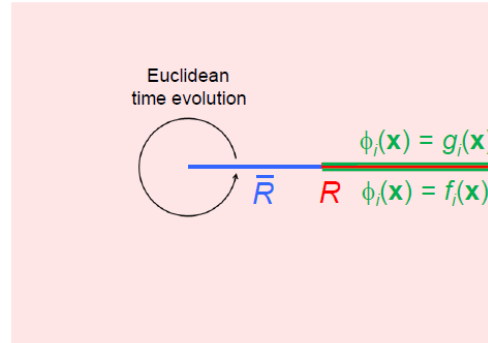


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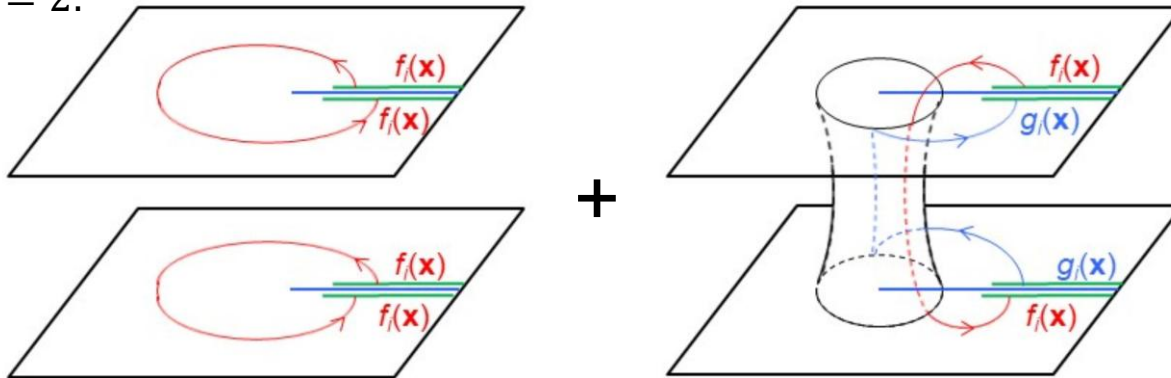
path integral



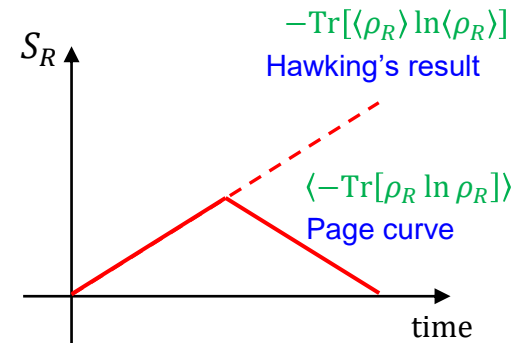
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$$S_R \equiv -\text{Tr}[\rho_R \ln \rho_R] = \lim_{n \rightarrow 1} \frac{1}{1-n} \ln \text{Tr}[\rho_R^n]$$

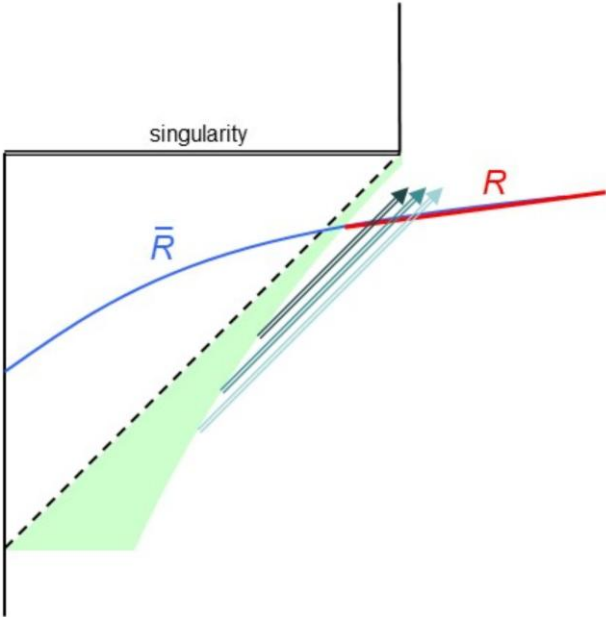
$n = 2$:



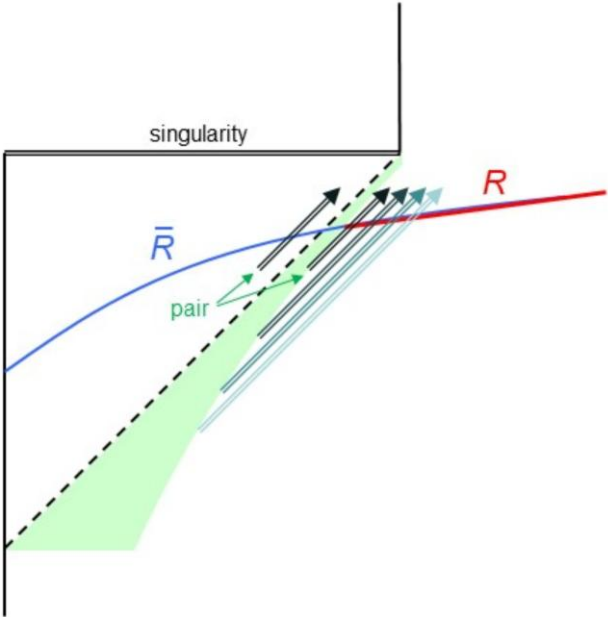
replica wormhole (nonperturbative effect)



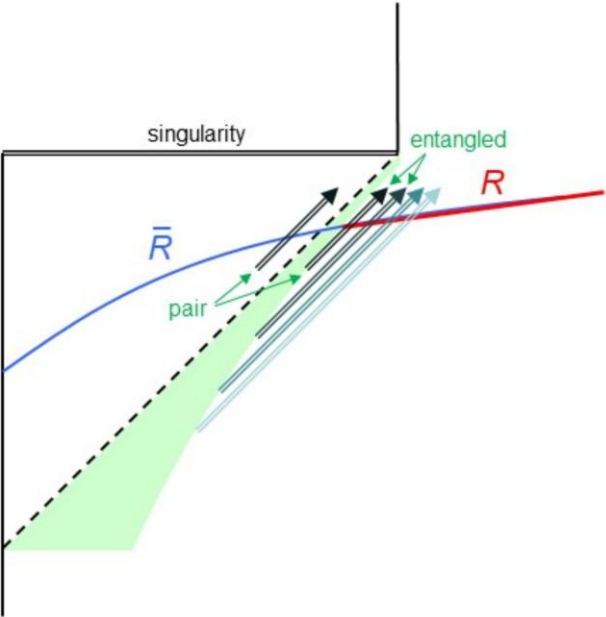
Redundancy in the Hawking process



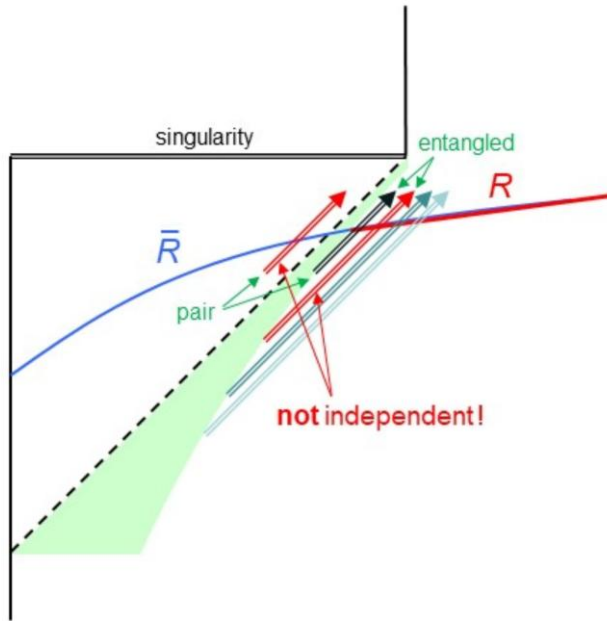
Redundancy in the Hawking process



Redundancy in the Hawking process



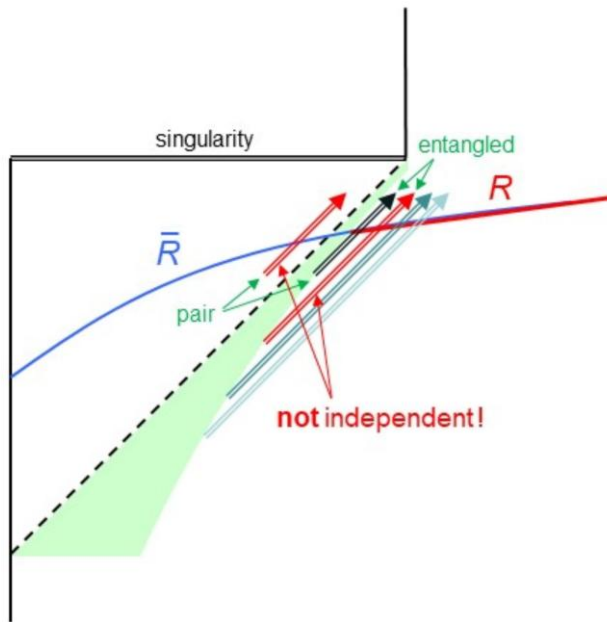
Redundancy in the Hawking process



→ Hawking radiation emitted earlier is
not independent of the interior d.o.f.s!

...; Maldacena, Susskind ('13); ...

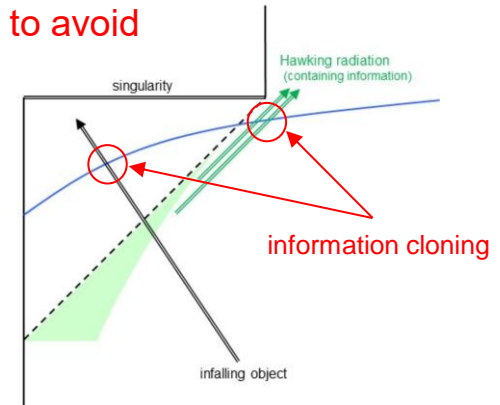
Redundancy in the Hawking process



→ Hawking radiation emitted earlier is **not** independent of the interior d.o.f.s!

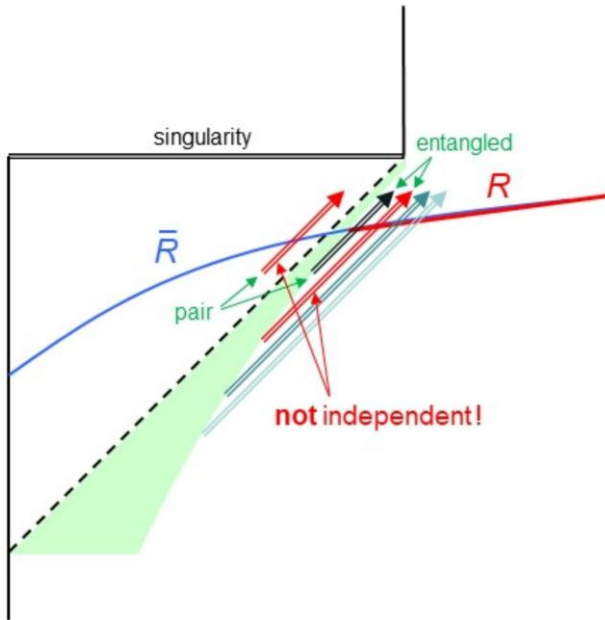
...; Maldacena, Susskind ('13); ...

- needed to avoid

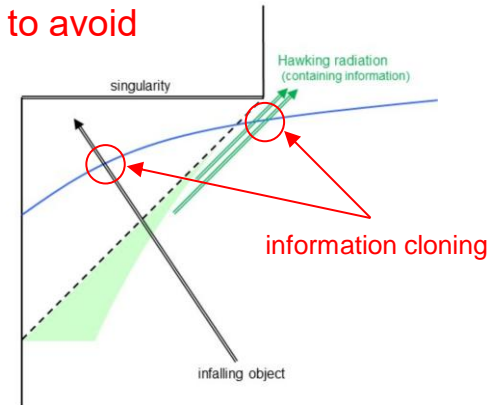


- consistent because of causality

Redundancy in the Hawking process



- needed to avoid



- consistent because of causality

→ Hawking radiation emitted earlier is **not** independent of the interior d.o.f.s!

...; Maldacena, Susskind ('13); ...

Global spacetime
(embracing the **interior**)

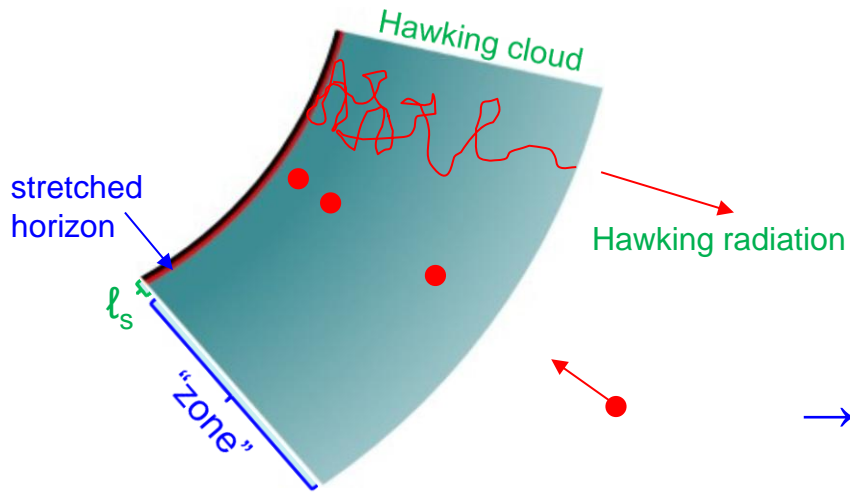
→
Replica wormholes
(nonperturbative effects of gravity)

Page curve
(signifying **unitarity**)

Recent Progress II

— unitary gauge construction —

Start with a “distant” (holographic) description



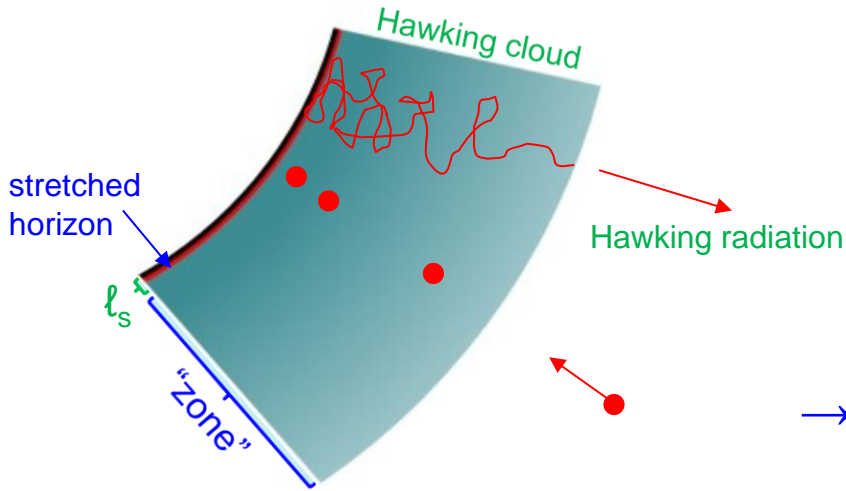
The d.o.f.s outside the horizon
comprise the **entire** system.

→ The evolution is unitary.

→ How does the “interior” emerge?

Papadodimas, Raju ('12-'15); Verlinde, Verlinde ('12-'13);
Y.N., Sanches, Varela, Weinberg ('12-'15); ...
Y.N. ('19, 20)

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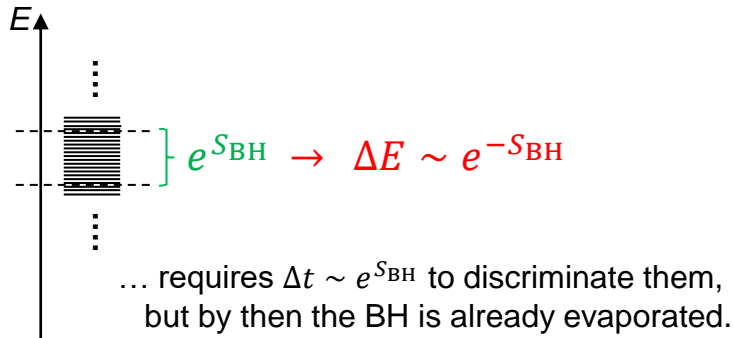
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Papadodimas, Raju ('12-'15); Verlinde, Verlinde ('12-'13);
Y.N., Sanches, Varela, Weinberg ('12-'15); ...
Y.N. ('19, 20)

Key features Y.N. ('19, 20)

— defining characteristics of BHs

(I) Exponentially dense spectrum



(II) Dynamics at the stretched horizon

$$T_{\text{local}} \sim M_{\text{string}}$$

... string dynamics $\left\{ \begin{array}{l} \bullet \text{ quantum chaos} \\ \bullet \text{ fast scrambling} \\ \bullet \text{ universal} \end{array} \right.$

Maldacena, Shenker, Stanford ('15)

Hayden, Preskill ('07); Sekino, Susskind ('08)

Banks, Seiberg ('10); ... Harlow, Ooguri ('18)

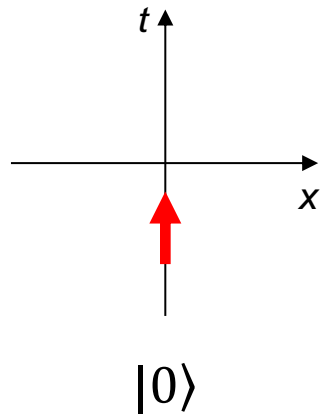
(e.g. no global symmetry)

→ “ultimate” thermalization

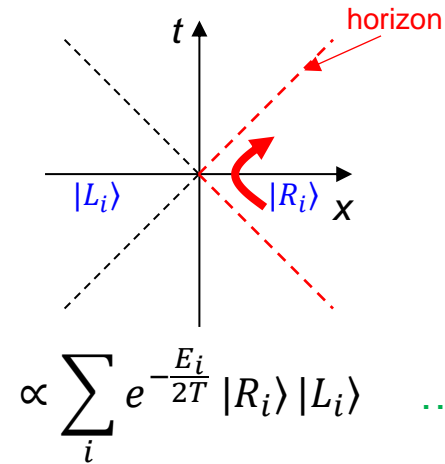
... **universal** across all low energy species

Emergence of the interior

Minkowski



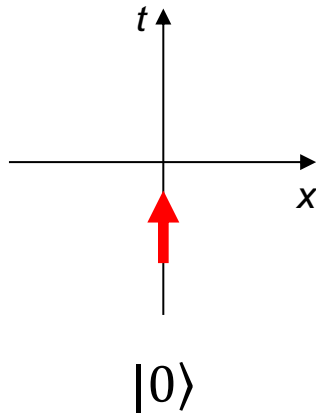
Rindler



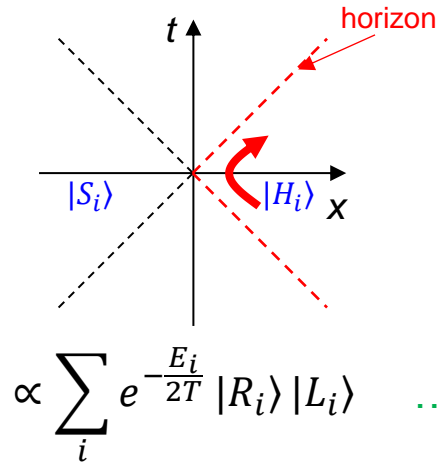
Fulling ('73); Davies ('75); Unruh ('76)

Emergence of the interior

Minkowski



Rindler



Near empty
Interior spacetime

(An object thrown "sees" interior spacetime)

frame change

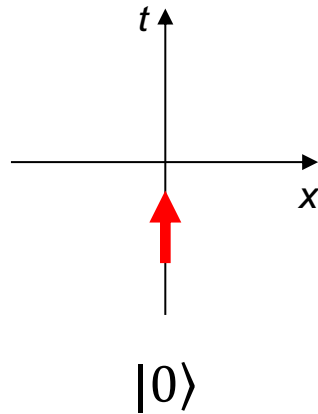
$$|\Psi_{\text{BH}}\rangle \propto \sum_i e^{-\frac{E_i}{2T_{\text{H}}}} |H_i\rangle |S_i\rangle$$

Hard mode states (pointing to $|H_i\rangle$)
Soft mode states (pointing to $|S_i\rangle$)
... play the role of the mirror partners
(representing their collective excitations)

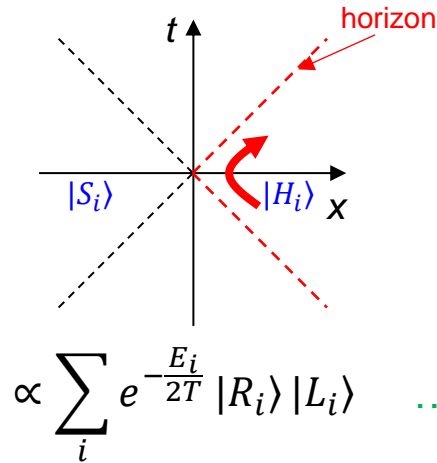
... universally thermal

Emergence of the interior

Minkowski



Rindler



Near empty
Interior spacetime

(An object thrown "sees" interior spacetime)

frame change

$$|\Psi_{\text{BH}}\rangle \propto \sum_i e^{-\frac{E_i}{2T_{\text{H}}}} |H_i\rangle |S_i\rangle$$

Hard mode states \downarrow $|H_i\rangle$ \uparrow Soft mode states $|S_i\rangle$

... play the role of the mirror partners

(representing their collective excitations)

... universally thermal

string dynamics

(not the case for the surface of regular material)

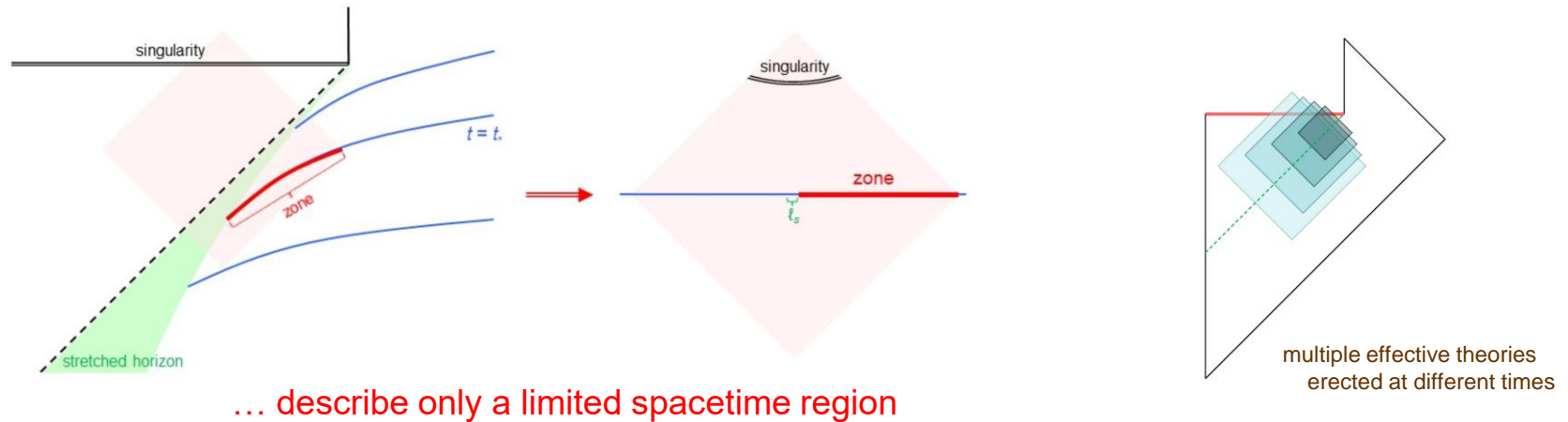
At late times, the BH is entangled with radiation

$$|\Psi_{\text{BH}}\rangle \propto \sum_i e^{-\frac{E_i}{2T_H}} |H_i\rangle \overbrace{|(S + R)_i\rangle}^{\text{Hard mode states}} \dots \text{play the role of the mirror partners}$$

Soft and far (radiation) mode states
(representing collective excitations of these modes)

... Interior d.o.f.s involve early Hawking radiation!

Effective theory of the interior



Distant description
(manifestly **unitary**)

Collective phenomena

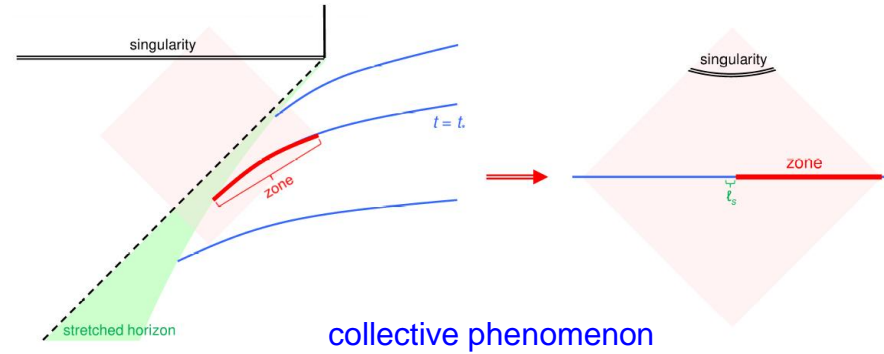
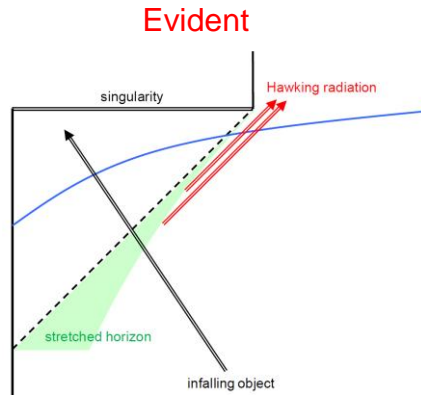
Interior spacetime
(effective emergence)

Structure of Quantum Gravity

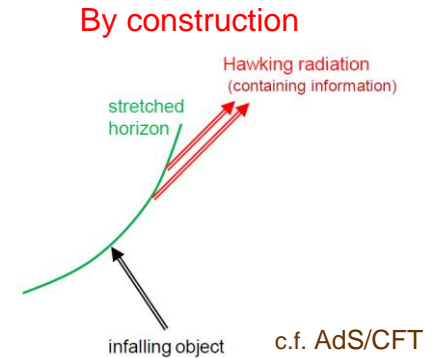
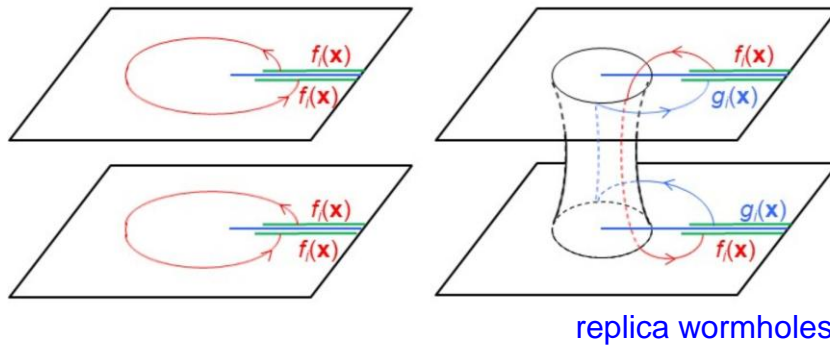
Global spacetime
— General relativity —

Unitary / holographic
— Quantum mechanics —

• Interior



• Unitarity



• Apparent violation of BH entropy

huge interior spatial volume at late times

Effective theory of the interior has a finite maximal volume.

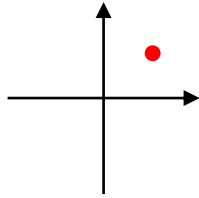
• Ensemble nature

semiclassically orthogonal states
in fact have $\langle \Psi_1 | \Psi_2 \rangle \sim e^{-S_{\text{BH}}/2}$
→ $e^{S_{\text{BH}}}$ states (+ null states)

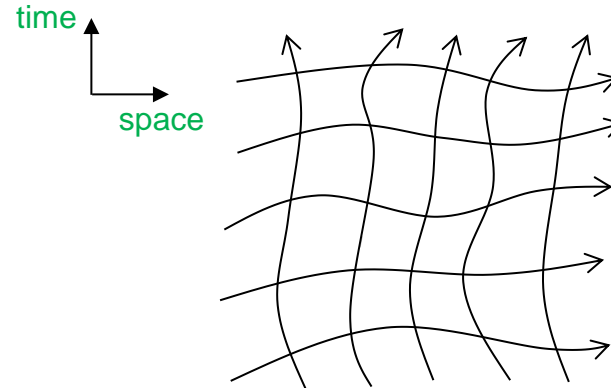
Hilbert space of dimension $e^{S_{\text{BH}}}$ can host $e^{S_{\text{BH}}}$ approximately orthogonal states.

Redundancies of the description

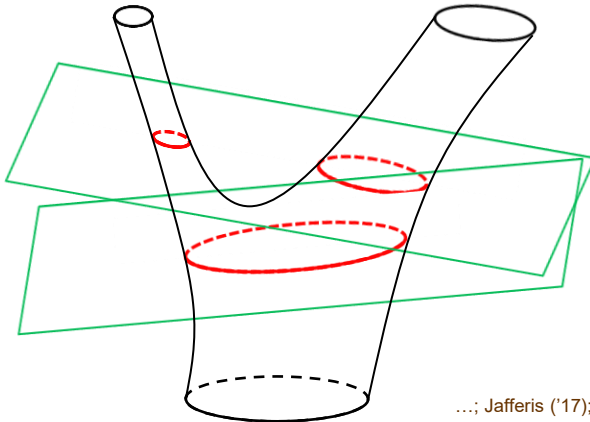
- General covariance (perturbative)



$$\begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} 1 \\ 1 \end{pmatrix} \leftrightarrow \begin{pmatrix} r \\ \theta \end{pmatrix} = \begin{pmatrix} \sqrt{2} \\ \pi/4 \end{pmatrix} \leftrightarrow \dots$$



- Nonperturbative redundancies



...; Jafferis ('17); Marolf, Maxfield ('20); ...

... allows for making (only) one of the two pillars manifest,
but the theory still accommodates both of them (QM + GR).

Summary

Black hole conundrum

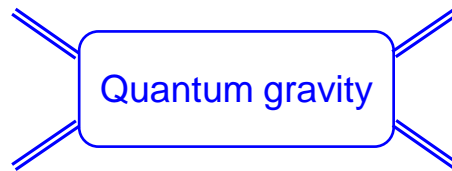


Structure of quantum gravity

⇒ Quantum mechanics & General relativity, but in a subtle manner!

High energy / Astro physics
(Black holes, gravity, ...)

Quantum information science
(Holography, ...)



Many-body physics (cond. matter & AMO)
(Chaos, fast scrambling, ...)

...

“From the Black Hole Conundrum to the Structure of Quantum Gravity”
Y.N., *Mod. Phys. Lett.* **A36** (2021) 2130007 [arXiv:2011.08707 [hep-th]]

“Complementarity for a Dynamical Black Hole”
B. Concepcion, Y.N., K. Ritchie, S. Weiss, arXiv:2405.15849 [hep-th]

Real Summary

Happy Birthday, Hitoshi!

