

Workshop on the Physics and the Mathematics of the Universe @ IPMU

Can we prove inflation?

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CSIC
CONSEJO SUPERIOR DE INVESTIGACIONES CIENTÍFICAS

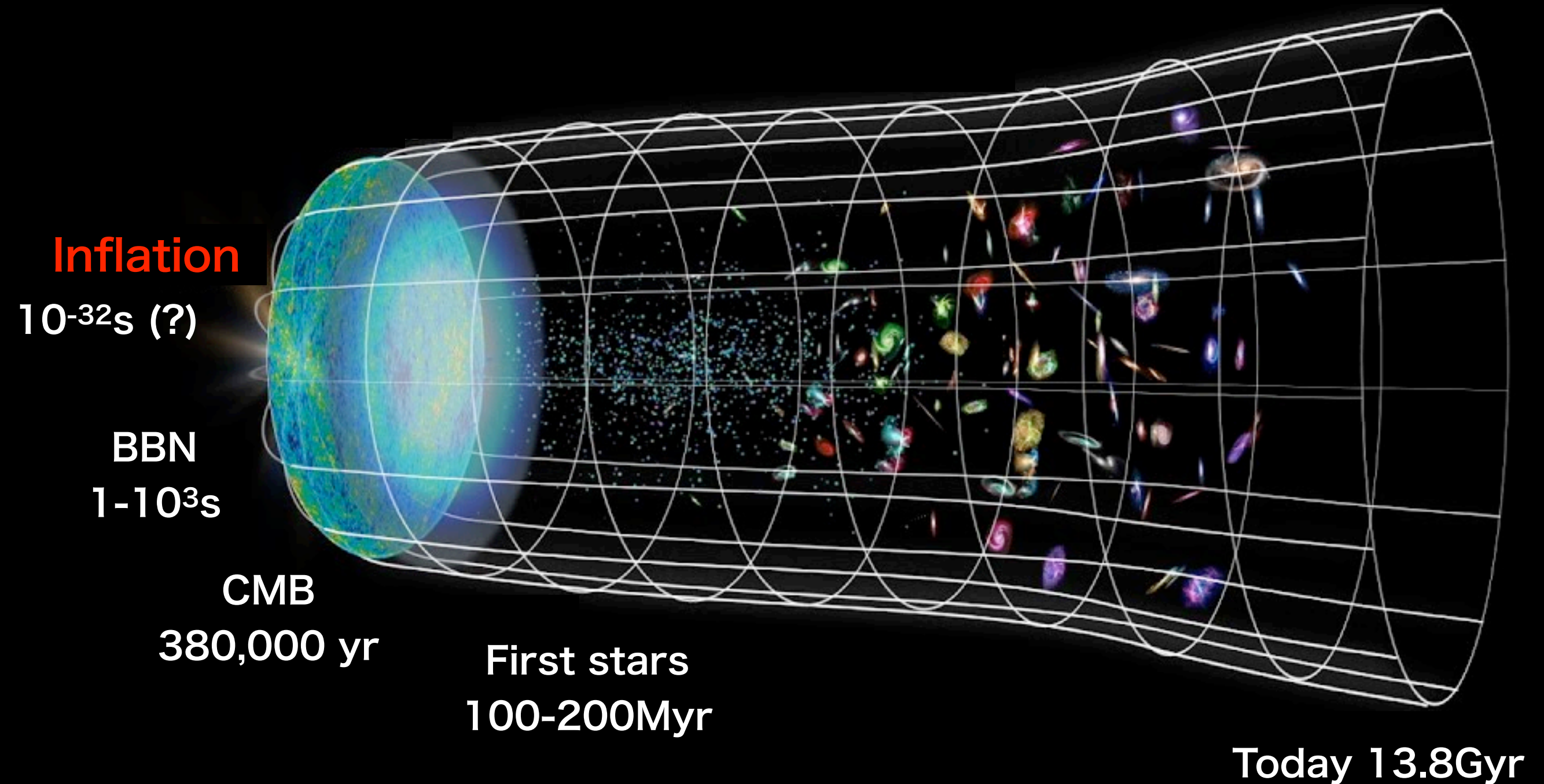


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Inflation

Early accelerated expansion of the Universe



No direct evidence, but supported by many observational facts.

How can we prove inflation?

prove



verb (used with object)

proved, proved, proven, proving.

- 1 to establish the truth or genuineness of, as by evidence or argument.
| *to prove one's claim.*
Synonyms: [verify](#), [substantiate](#), [confirm](#), [demonstrate](#)
Antonyms: [disprove](#)
- 2 *Law.* to establish the authenticity or validity of (a will); probate.
- 3 to give demonstration of by action.
- 4 to subject to a test, experiment, comparison, analysis, or the like, to determine quality, amount, acceptability, characteristics, etc..
| *to prove ore.*
- 5 to show (oneself) to have the character or ability expected of one, especially through one's actions.
- 6 *Mathematics.* to verify the correctness or validity of by mathematical demonstration or arithmetical proof.

How can we ~~prove~~ inflation?

Predictions of Inflation Theory

confirm/test

1. Accelerated expansion of the Universe

$$\ddot{a} > 0$$

$a(t)$: scale factor of the Universe
 t : cosmic time

2. Nearly scale-invariant SCALAR perturbations

$$\mathcal{P}_{\text{S,prim}}(k) \sim \text{const.} \quad k : \text{wavenumber}$$

3. Nearly scale-invariant TENSOR perturbations

$$\mathcal{P}_{\text{T,prim}}(k) \sim \text{const.}$$

How can we ~~prove~~ inflation?

Predictions of Inflation Theory

confirm/test

Special Thanks to
Yashar Akrami (IFT)

1. Accelerated expansion of the Universe

Only a direct measurement of the accelerated expansion could serve as proof of cosmic inflation!



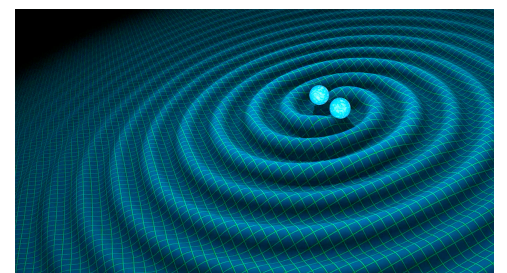
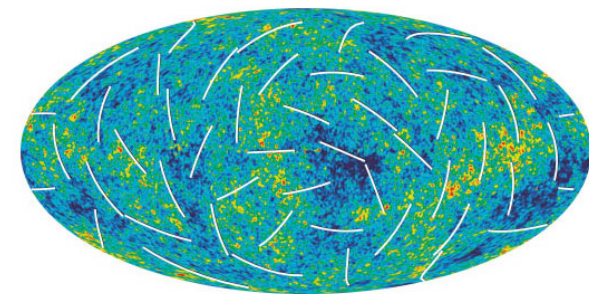
2. Nearly scale-invariant SCALAR perturbations ✓

- Cosmic microwave background (CMB)
- Galaxy survey



3. Nearly scale-invariant TENSOR perturbations

- CMB B-mode polarization
- Direct detection of gravitational waves (GWs)



→ **This talk**

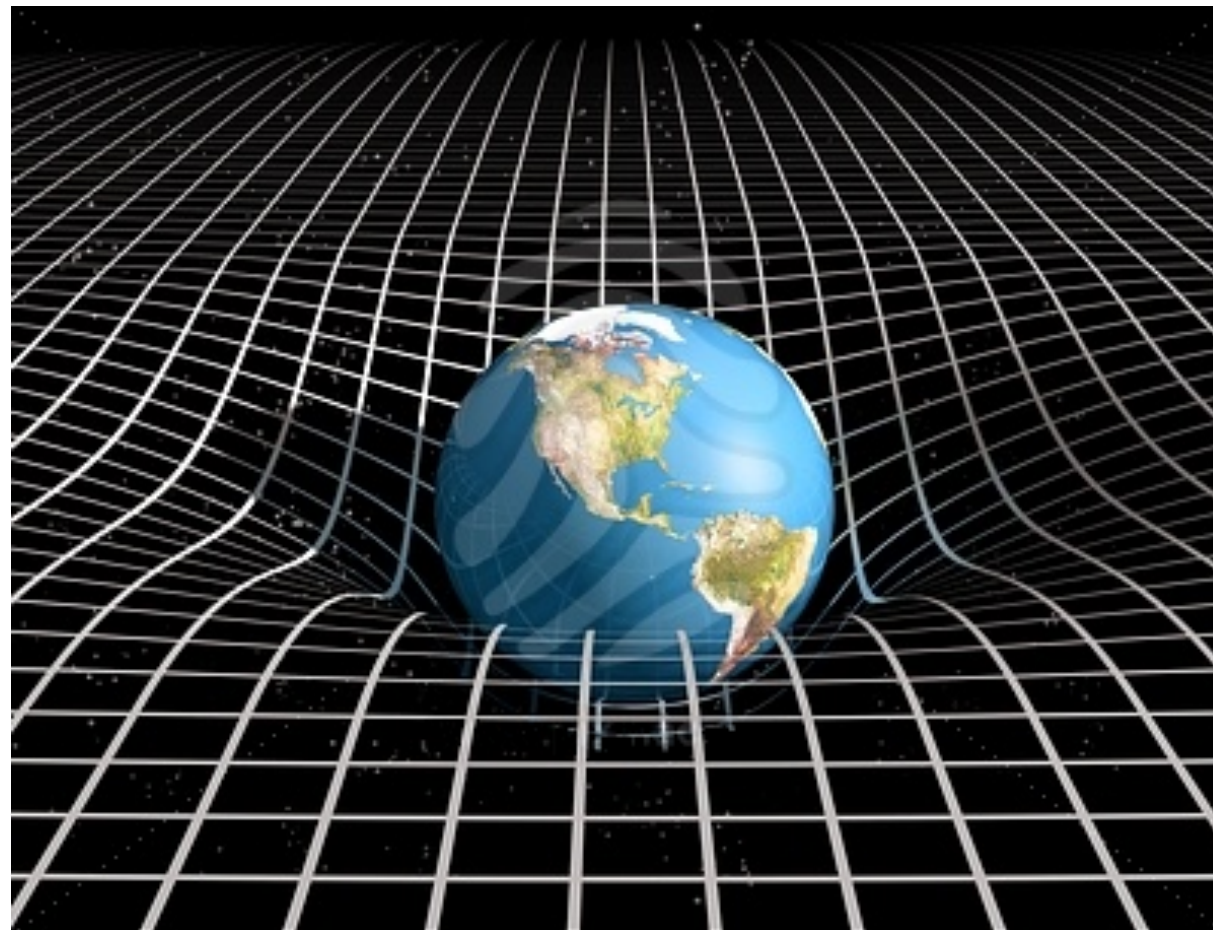
Basics of cosmology

Einstein equation

$$G_{\mu\nu} = 8\pi G T_{\mu\nu}$$

Geometry
= Gravity

Matter



Equation for an expanding Universe

Einstein equation

$$G_{\mu\nu} = 8\pi G T_{\mu\nu}$$

$$T_{\mu\nu} = \begin{pmatrix} -\rho & 0 & 0 & 0 \\ 0 & p & 0 & 0 \\ 0 & 0 & p & 0 \\ 0 & 0 & 0 & p \end{pmatrix}$$

ρ : density p : pressure

Friedmann-Lemaître-Robertson-Walker (FLRW) metric)

$$ds^2 = -dt^2 + a^2(t) \left[\frac{dr^2}{1 - Kr^2} + r^2(d\theta^2 + \sin^2 \theta d\phi^2) \right]$$

$a(t)$: scale factor

K : curvature $< 10^{-2}$ from observation

Friedmann equation:

$$\left(\frac{\dot{a}}{a} \right)^2 = \frac{8\pi G}{3} \rho - \frac{K}{a^2}$$

Hubble expansion rate

$$\text{Friedmann equation: } H^2 = \frac{8\pi G}{3} \rho - \frac{K}{a^2}$$

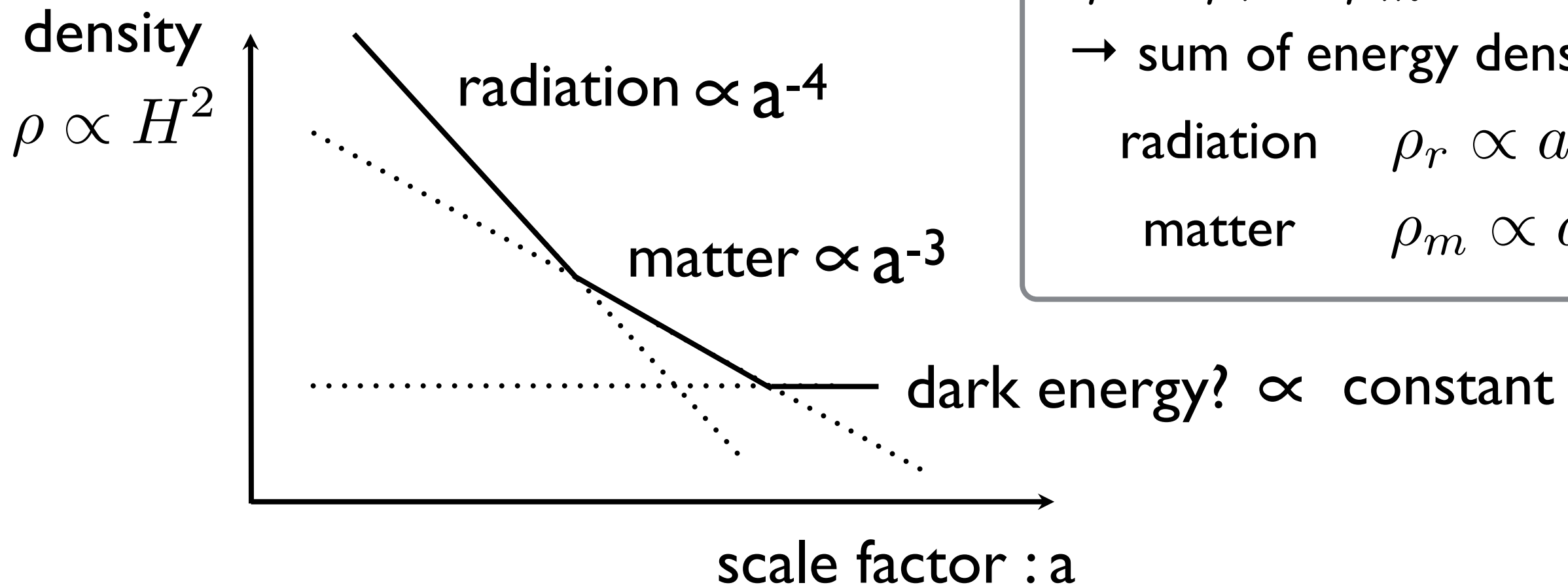
$$H \equiv \frac{\dot{a}}{a} : \text{Hubble parameter}$$

$$\rho = \rho_r + \rho_m + \dots$$

→ sum of energy densities

radiation $\rho_r \propto a^{-4}$

matter $\rho_m \propto a^{-3}$



Motivation to consider inflation

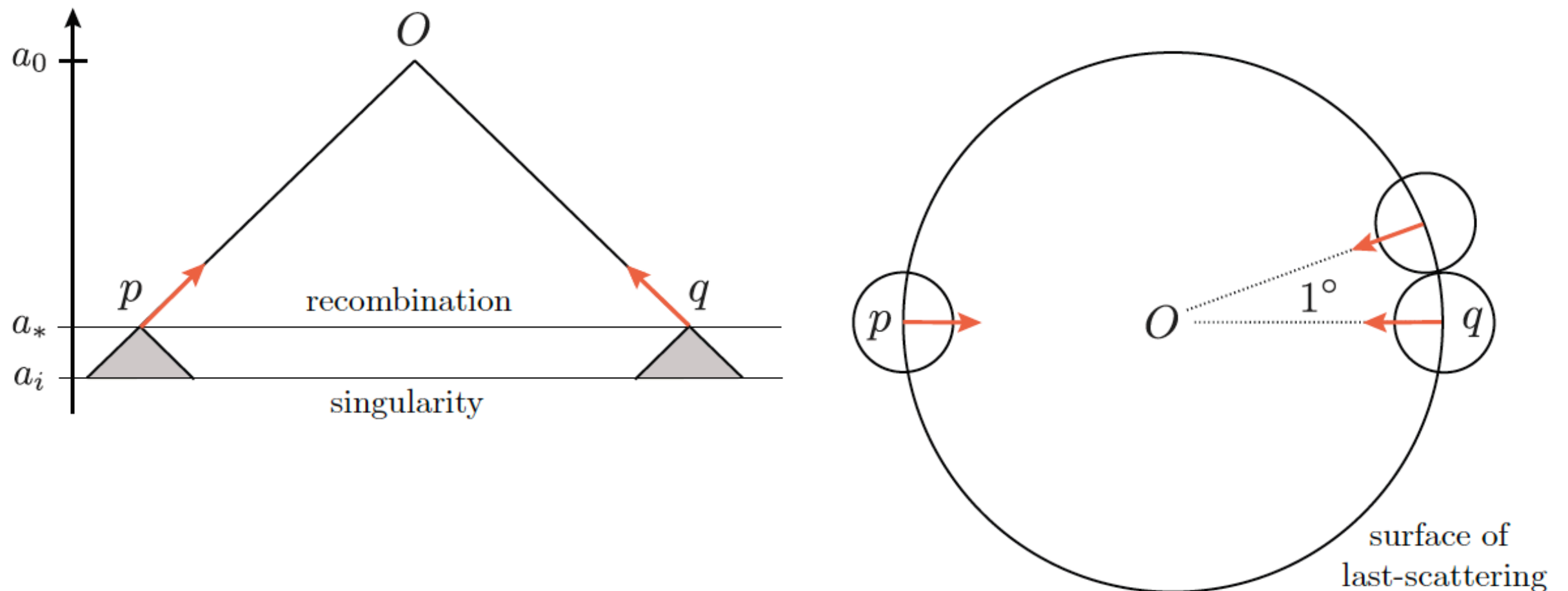
Horizon / Flatness / Monopole problems

Inflation explains why...

the Universe is homogeneous

the Universe is flat

we do not observe magnetic monopole



Horizon problem

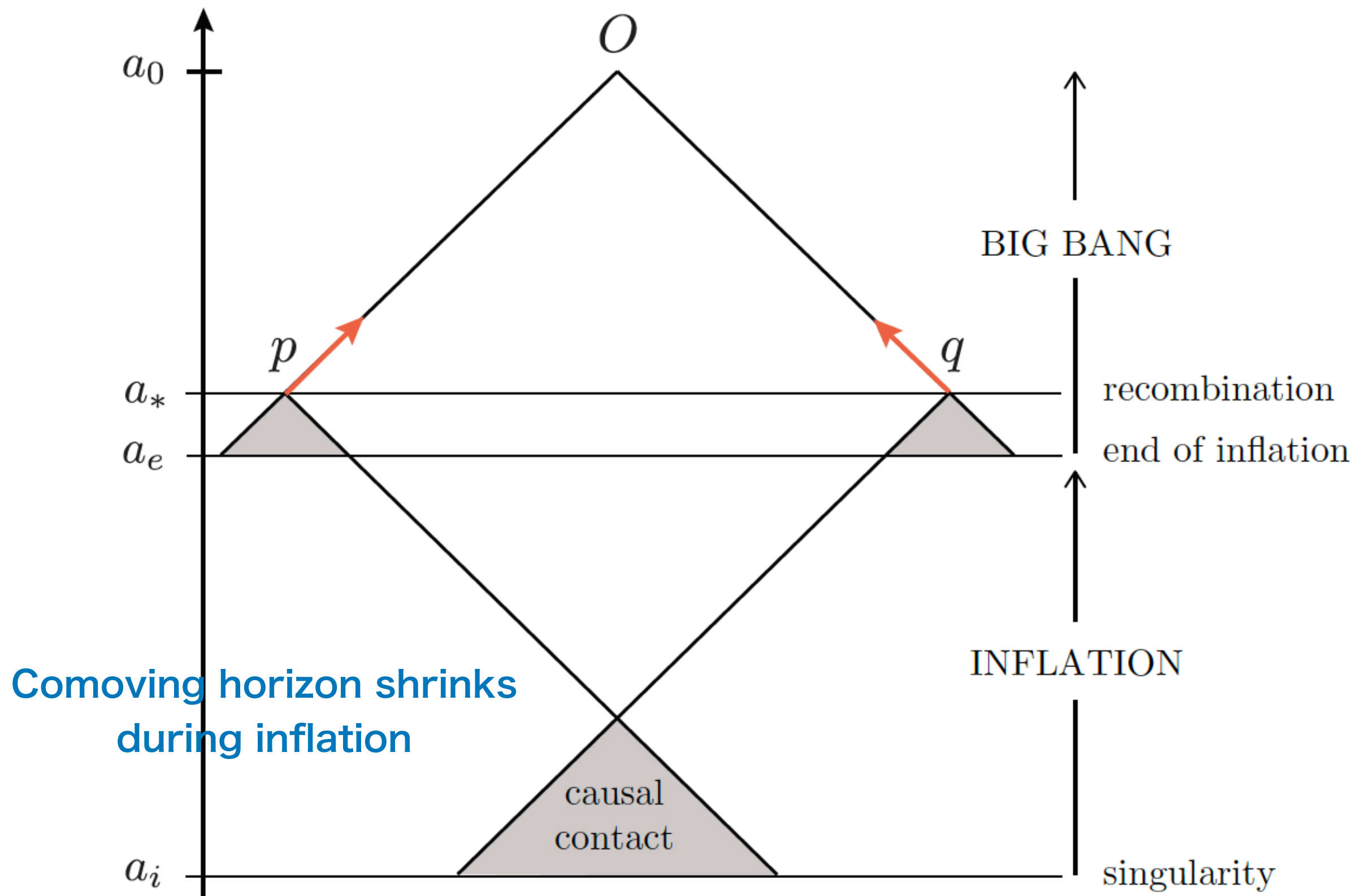


Figure from Baumman, "TASI Lectures on Primordial Cosmology", arXiv:1807.03098

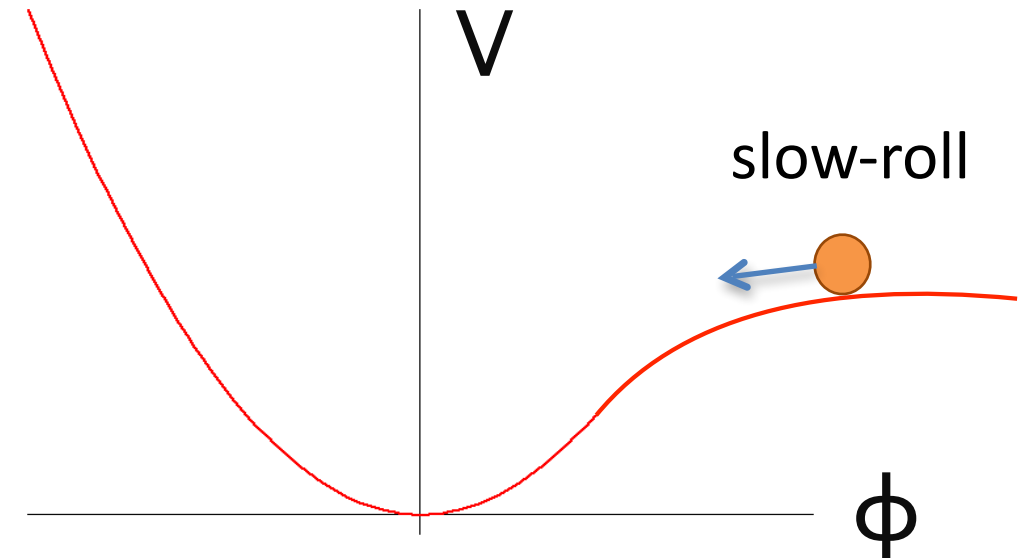
What drives inflation?

Most popular scenario

Inflation is driven by a scalar field slowly rolling down in its potential

Why use a scalar field?

→ Because it easily produces an isotropic universe.



Energy density of a scalar field

$$\rho_\phi = \dot{\phi}^2/2 + V$$

Equation of Motion

$$\ddot{\phi} + 3H\dot{\phi} + V' = 0$$

Friedmann equation

$$H^2 = \frac{8\pi}{3m_{\text{Pl}}^2} \left(\frac{1}{2} \dot{\phi}^2 + V(\phi) \right)$$

slow-roll approximation

$$\dot{\phi}^2/2 \ll V$$

$$\longrightarrow H = \text{const.} \propto V$$

$$\longrightarrow a \propto \exp(Ht)$$

Exponential expansion

Why do we want to confirm inflation?

because it's about the origin of the universe
and could be linked to physics beyond standard models

What causes the accelerated expansion?

Einstein equation

$$G_{\mu\nu} = 8\pi G T_{\mu\nu}$$

Geometry
= Gravity

modification of gravity

Matter

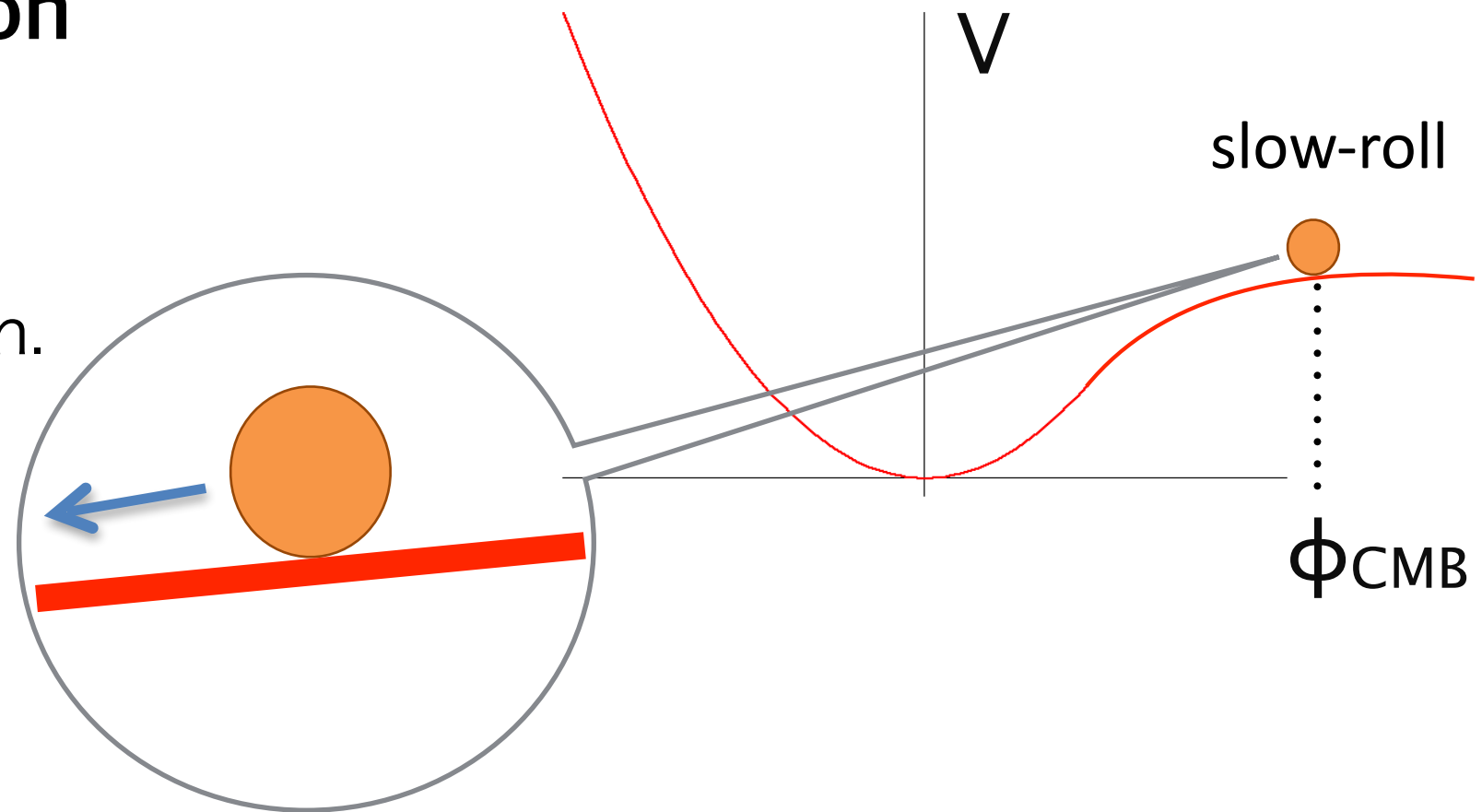
new (scalar) particle

→ new hint for the fundamental law of the universe

Predictions of inflation theory

Slow-roll parametrization

Since the scalar field rolls slowly, its position does not change much during inflation.



slow-roll parameters $\ll 1$

$$\epsilon \equiv \frac{M_{\text{Pl}}^2}{2} \left(\frac{V'}{V} \right)^2 \quad \eta \equiv M_{\text{Pl}}^2 \frac{V''}{V} \quad \xi_V^2 \equiv M_{\text{Pl}}^4 \frac{V' V'''}{V^2}$$

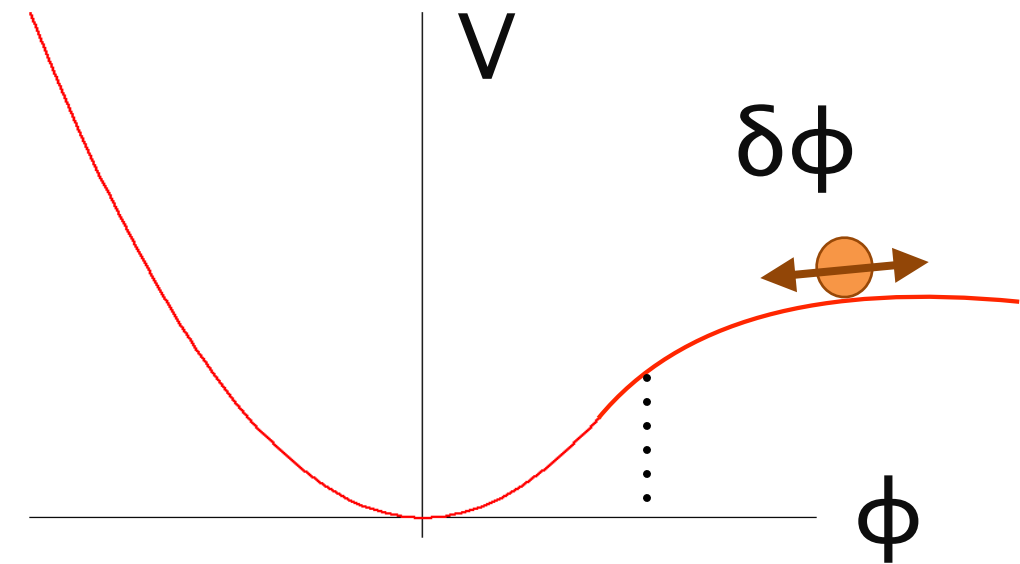
→ A quantity that characterizes the potential V around $\phi = \phi_{\text{CMB}}$ (' denotes a derivative with respect to ϕ)

Creation of scalar perturbations

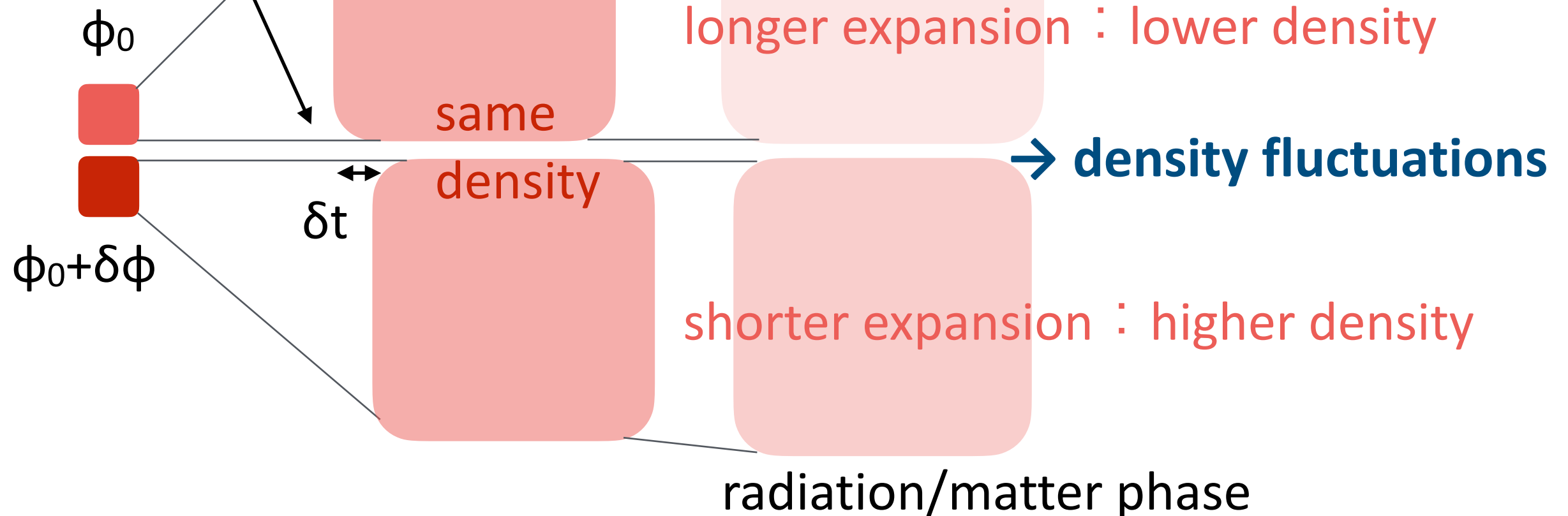
Inflaton scalar field has quantum fluctuations

→ Average ϕ_0 , Fluctuation $\delta\phi$

End of inflation is deviated by $\delta t = \frac{\delta\phi}{\dot{\phi}}$



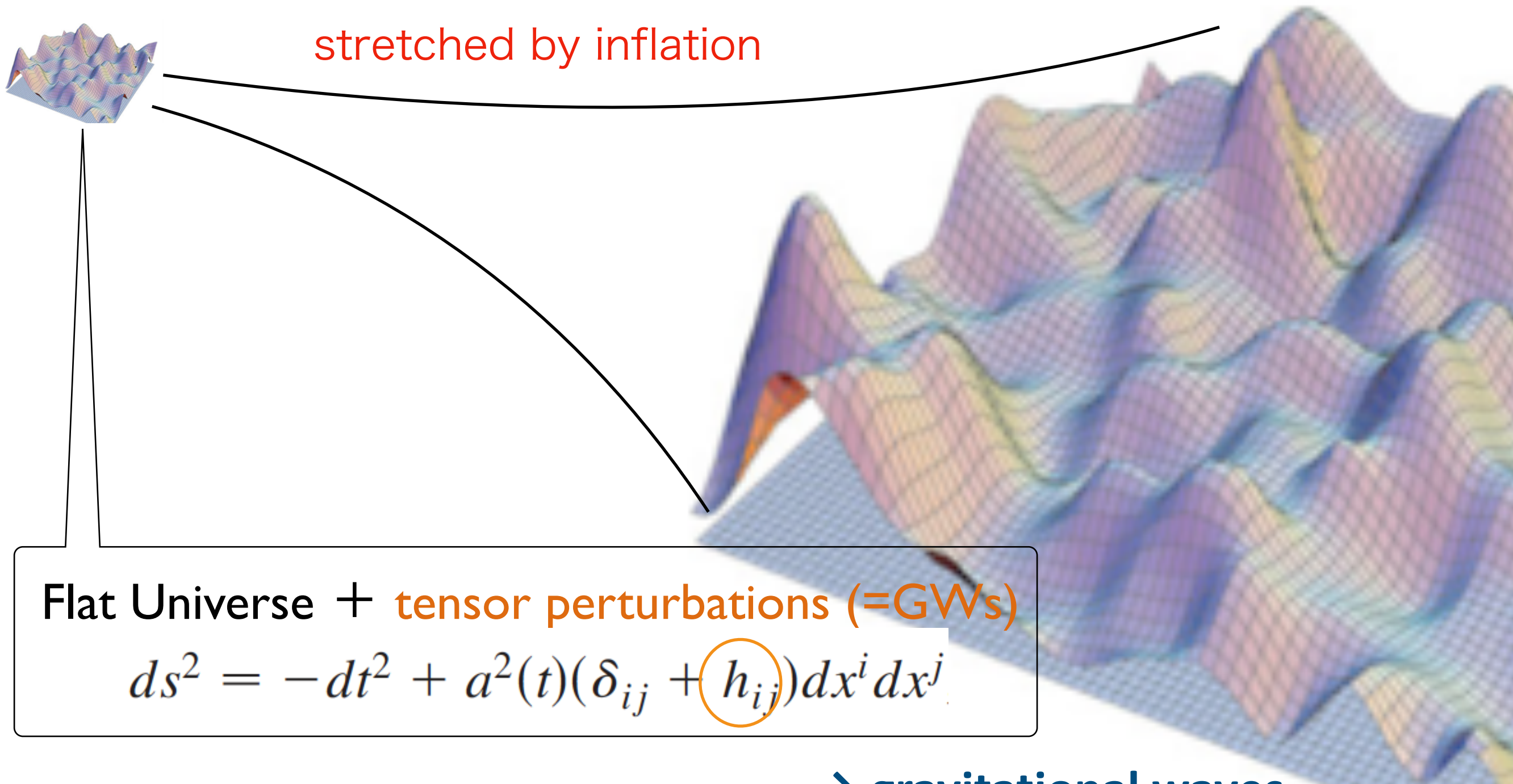
inflation ends when $\epsilon = 1$



Creation of tensor perturbations

Space-time itself exhibits
quantum fluctuations

stretched by inflation



Flat Universe + tensor perturbations (=GWs)

$$ds^2 = -dt^2 + a^2(t)(\delta_{ij} + h_{ij})dx^i dx^j$$

→ gravitational waves

Predictions of inflation theory

Scalar power spectrum

$$\mathcal{P}_{S,\text{prim}}(k) = \frac{1}{\pi\epsilon} \left(\frac{H}{m_{\text{Pl}}} \right)^2 \Big|_{k=aH}$$

→ observed as density perturbations

Spectral tilt

$$n_S(k) - 1 \equiv \frac{d \ln \mathcal{P}_{S,\text{prim}}(k)}{d \ln k} \simeq -6\epsilon + 2\eta$$

Running index

$$\alpha_S(k) \equiv \frac{dn_S(k)}{d \ln k} \simeq -16\epsilon\eta + 24\epsilon^2 + 2\xi^2$$

Tensor power spectrum

$$\mathcal{P}_{T,\text{prim}}(k) = \frac{16}{\pi} \left(\frac{H}{m_{\text{Pl}}} \right)^2 \Big|_{k=aH}$$

→ observed as gravitational waves

Tensor-to-scalar ratio

$$r \equiv \frac{\mathcal{P}_{T,\text{prim}}(k)}{\mathcal{P}_{S,\text{prim}}(k)} \simeq 16\epsilon$$

Spectral tilt

$$n_T(k) \equiv \frac{d \ln \mathcal{P}_{T,\text{prim}}(k)}{d \ln k} \simeq -2\epsilon$$

Predictions of inflation theory

Scalar power spectrum

$$\mathcal{P}_{S,\text{prim}}(k) = \frac{1}{\pi\epsilon} \left(\frac{H}{m_{\text{Pl}}} \right)^2 \Big|_{k=aH}$$

→ observed as density perturbations

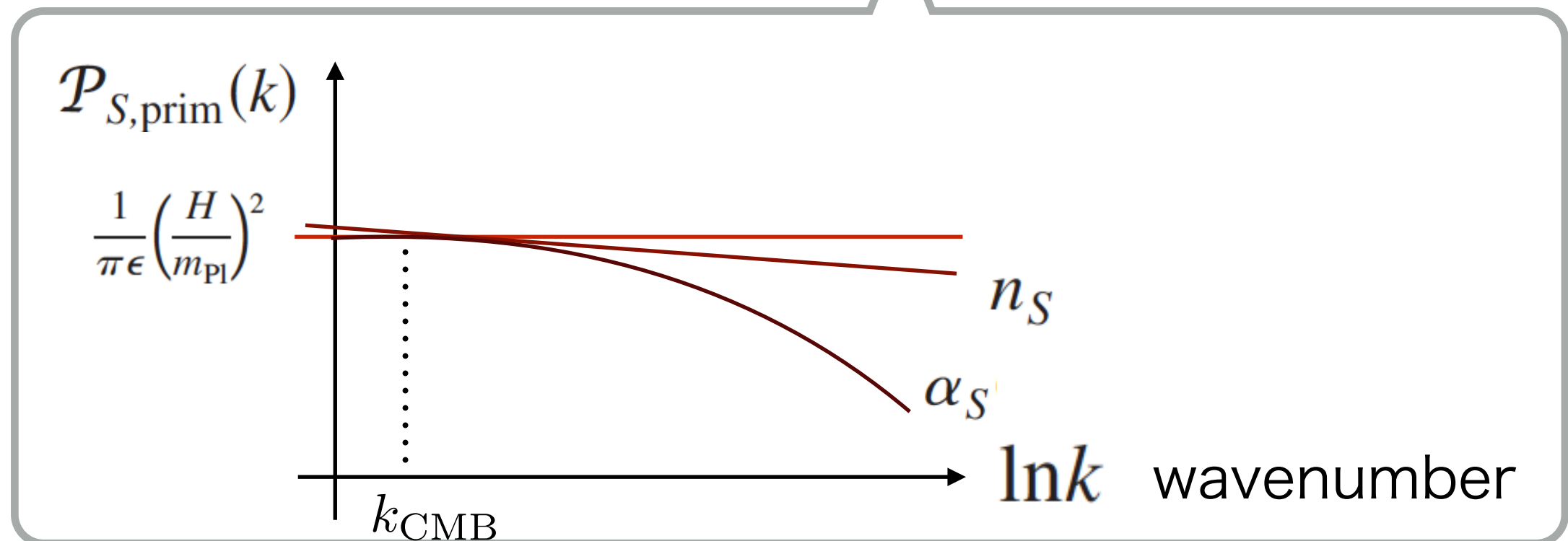
Spectral tilt

$$n_S(k) - 1 \equiv \frac{d \ln \mathcal{P}_{S,\text{prim}}(k)}{d \ln k} \simeq -6\epsilon + 2\eta$$

Running index

$$\alpha_S(k) \equiv \frac{dn_S(k)}{d \ln k} \simeq -16\epsilon\eta + 24\epsilon^2 + 2\xi^2$$

suppressed

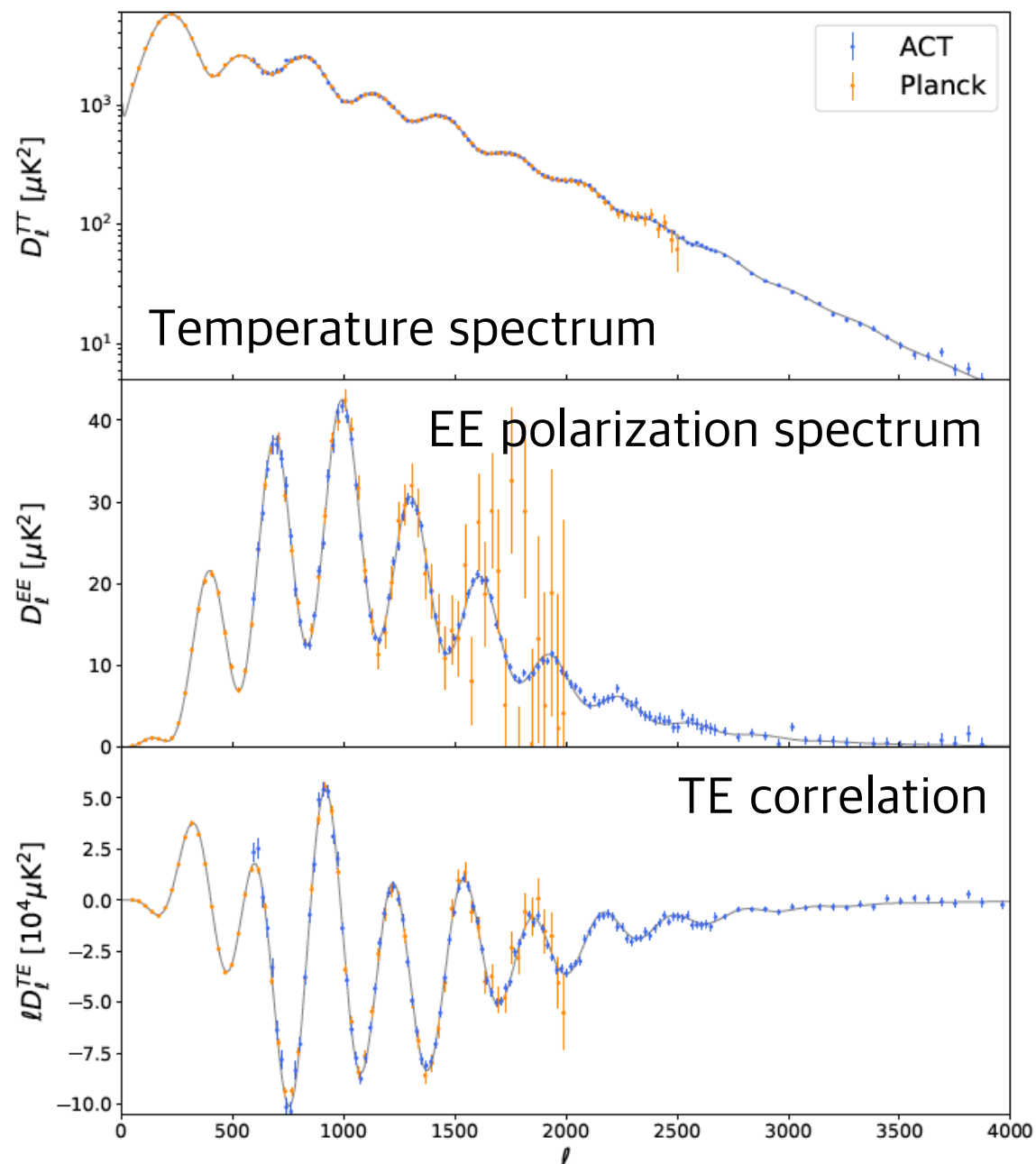
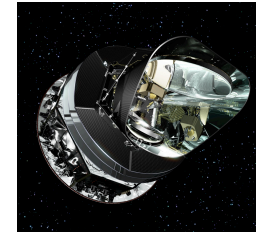


Observational support so far

2. Nearly scale-invariant SCALAR perturbations

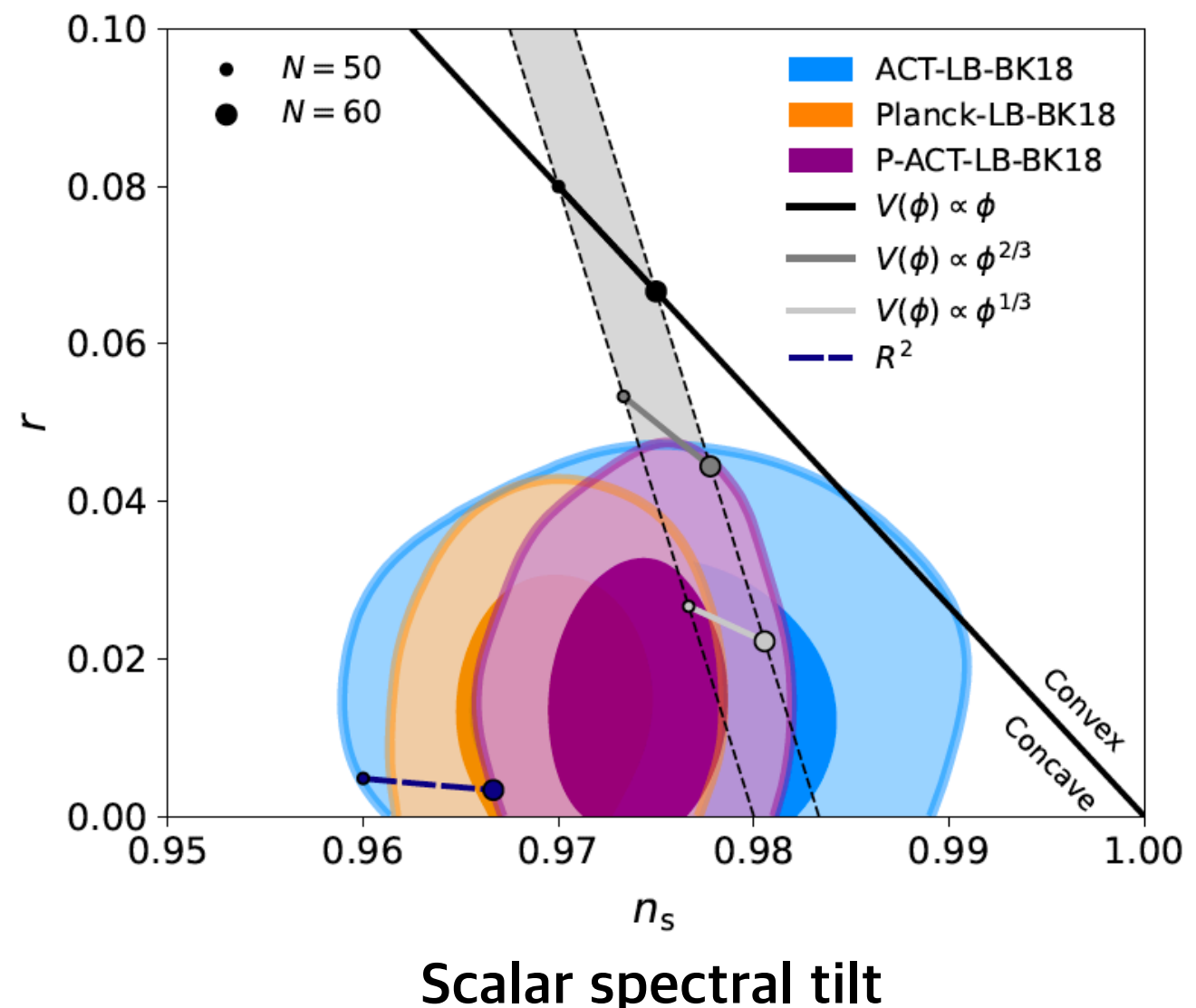
strongly supported by the the CMB observations

ACT Collaboration, arXiv:2503.14452, arXiv:2503.14454



Small deviation from
the scale invariant spectrum ($n_s=1$)

Tensor-to-scalar ratio

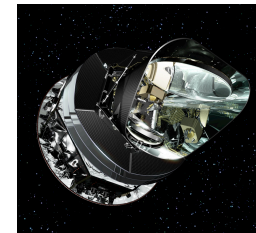


Observational support so far

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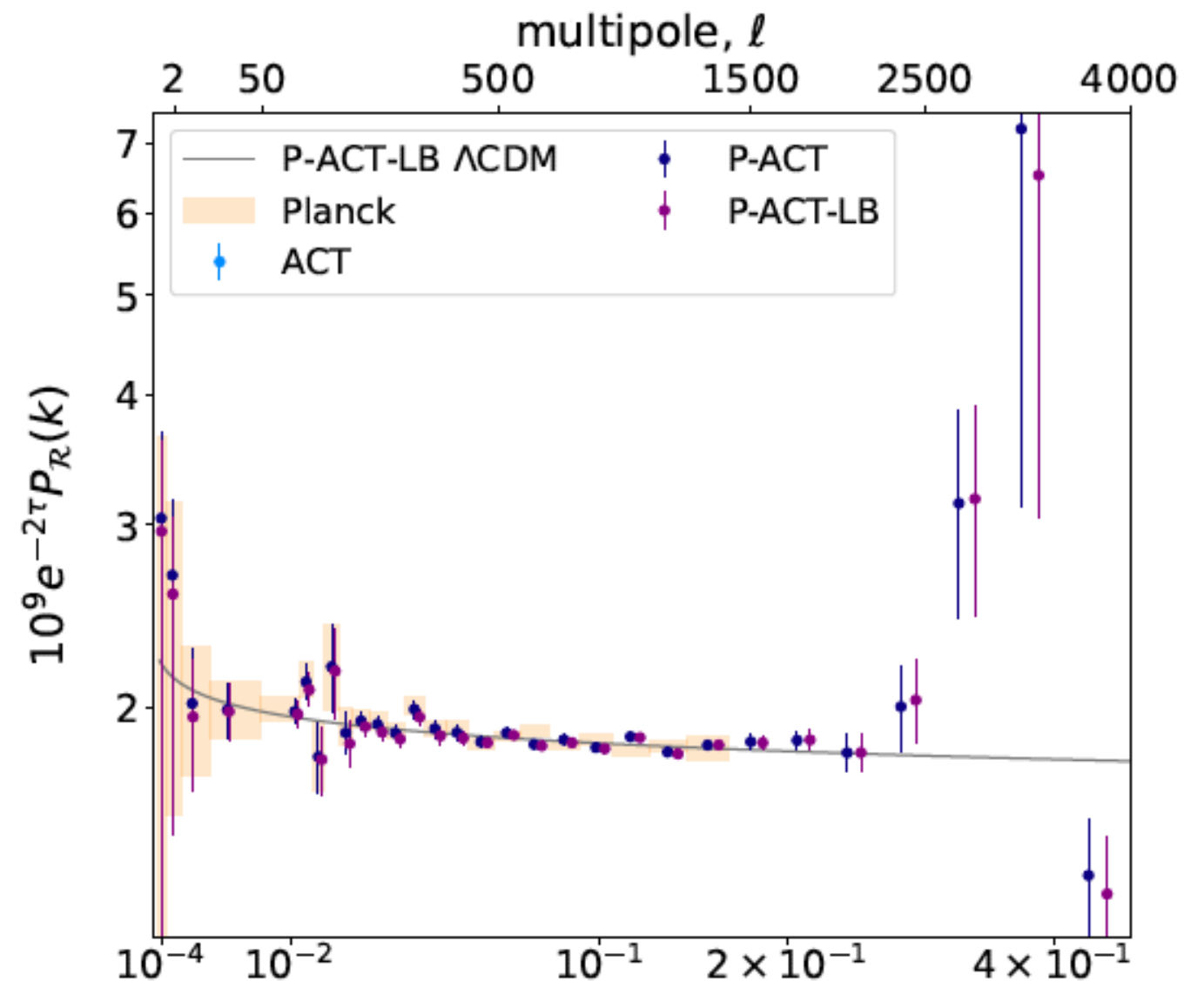
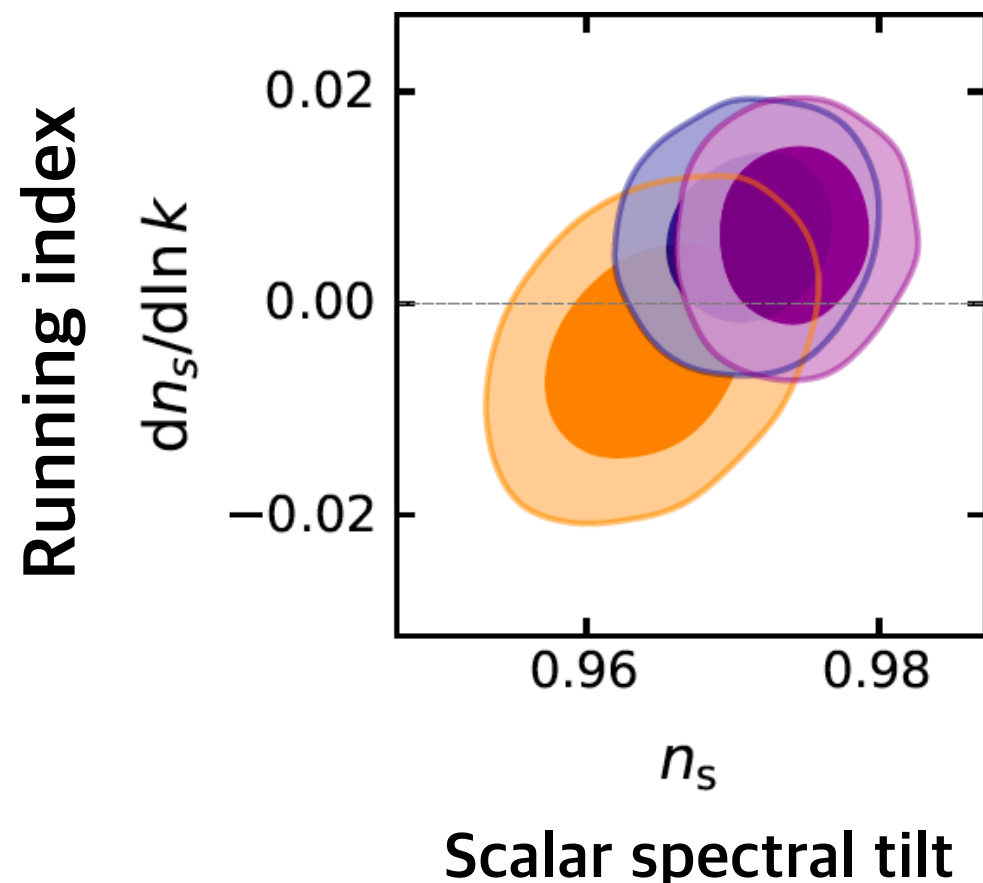
strongly supported by the the CMB observations

ACT Collaboration, arXiv:2503.14452, arXiv:2503.14454



Reconstruction of the power spectrum

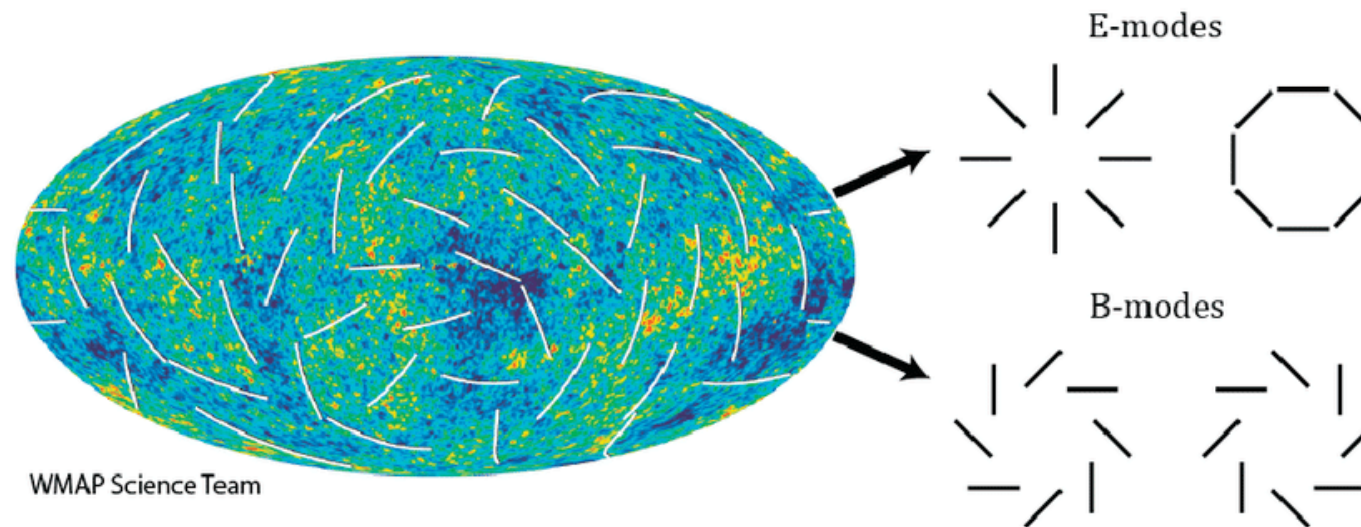
Scalar running is
consistent with zero



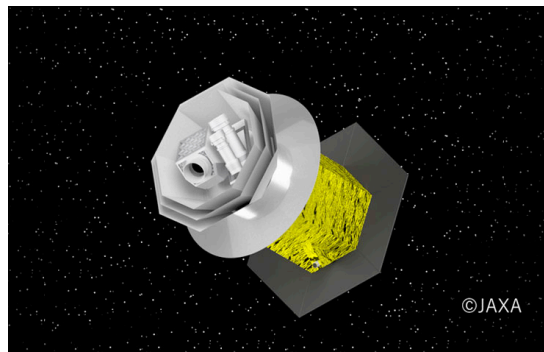
Future observations

3. Nearly scale-invariant TENSOR perturbations

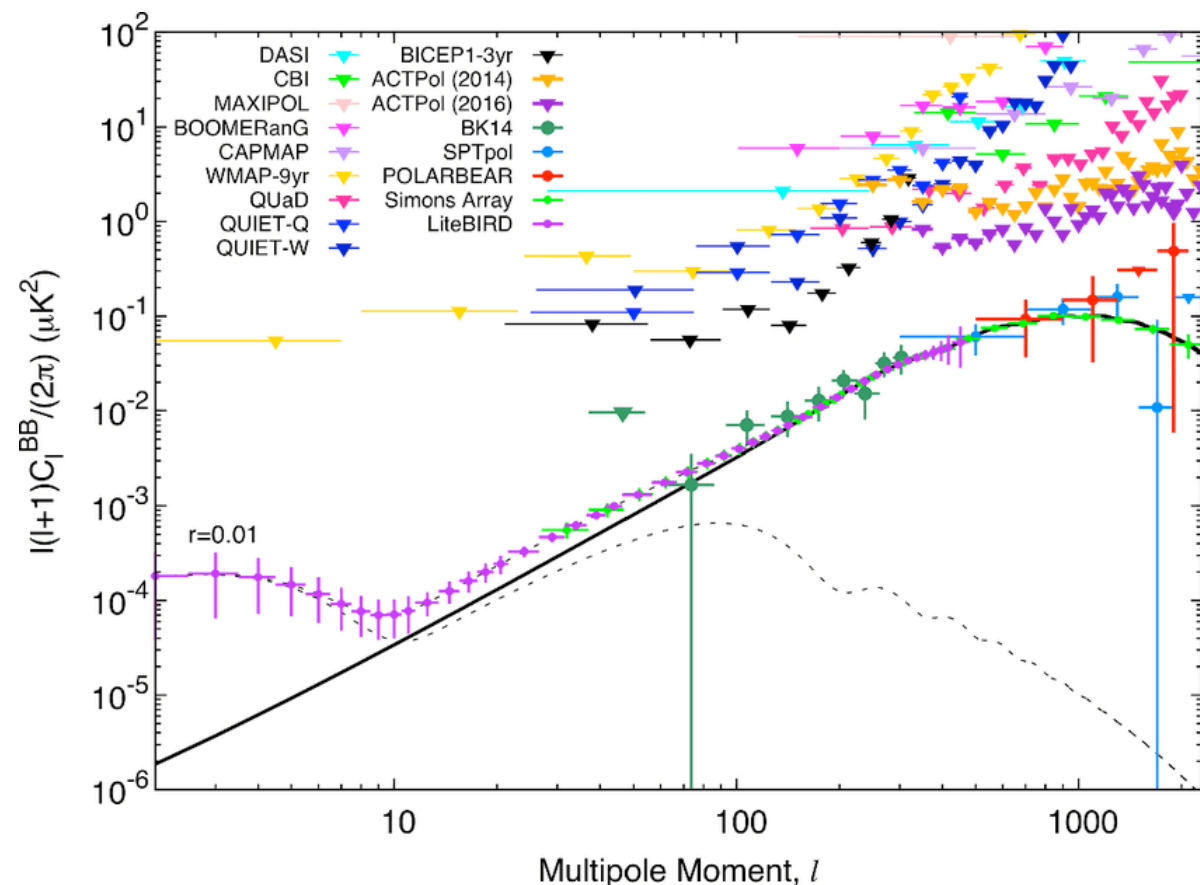
CMB B-mode polarization is a unique probe for inflation!



LiteBIRD



Simons
Observatory

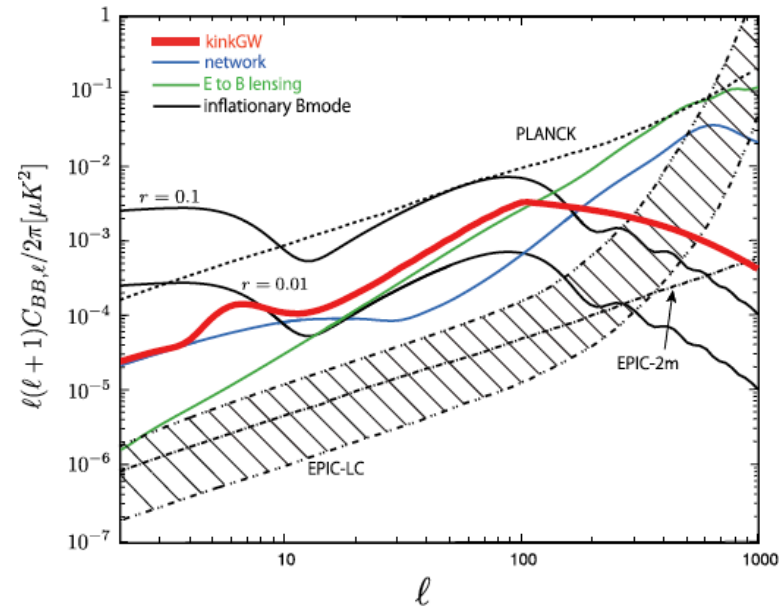


But... Is this enough?

No, because not only inflation theory predicts B-mode polarization!

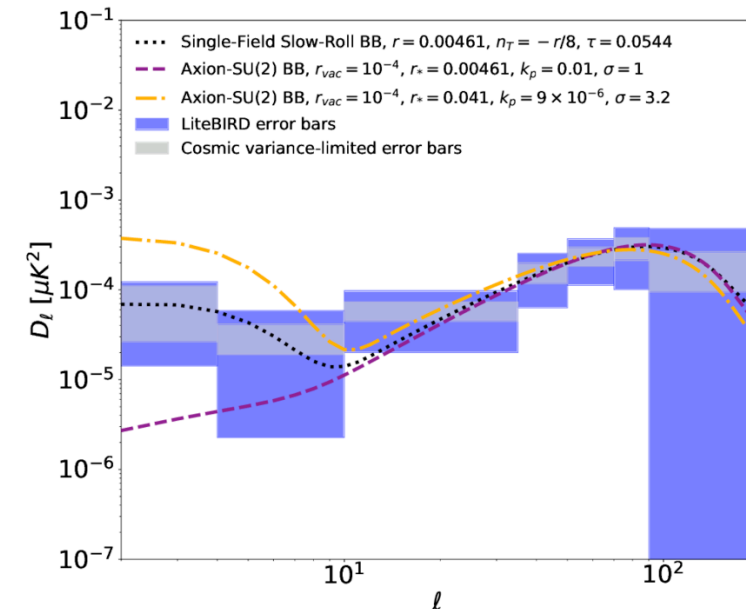
- **Cosmic strings**

Kawasaki et al., PRD 82, 103504 (2010)



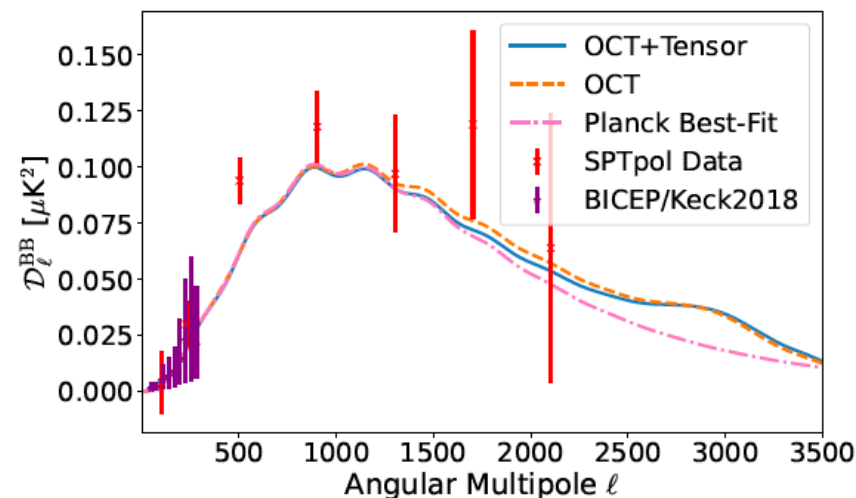
- **SU(2) gauge fields**

LiteBIRD Collaboration, PTEP 042F01 (2023)



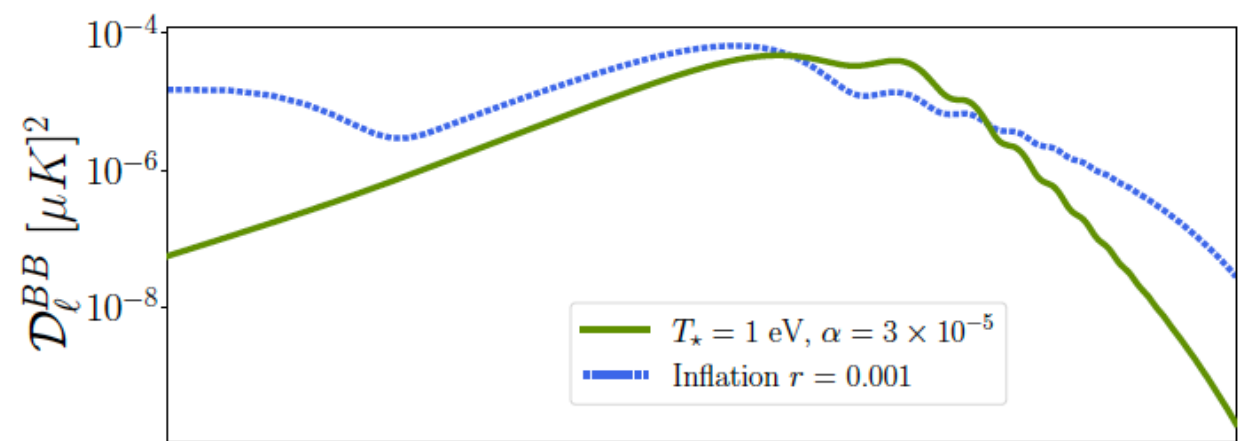
- **Magnetic fields**

Khalife & Pitrou, arXiv:2410.03612



- **Phase transitions**

Greene et al., arXiv:2410.23348



Any further tests?

→ **Consistency relation!**

$$r = -8n_T$$

Predictions of inflation theory

Scalar power spectrum

→ observed as density perturbations

$$\mathcal{P}_{S,\text{prim}}(k) = \frac{1}{\pi\epsilon} \left(\frac{H}{m_{\text{Pl}}} \right)^2 \Big|_{k=aH}$$

Spectral tilt

$$n_S(k) - 1 \equiv \frac{d \ln \mathcal{P}_{S,\text{prim}}(k)}{d \ln k} \simeq -6\epsilon + 2\eta$$

Running index

$$\alpha_S(k) \equiv \frac{dn_S(k)}{d \ln k} \simeq -16\epsilon\eta + 24\epsilon^2 + 2\xi^2$$

Tensor power spectrum

→ observed as gravitational waves

$$\mathcal{P}_{T,\text{prim}}(k) = \frac{16}{\pi} \left(\frac{H}{m_{\text{Pl}}} \right)^2 \Big|_{k=aH}$$

Tensor-to-scalar ratio

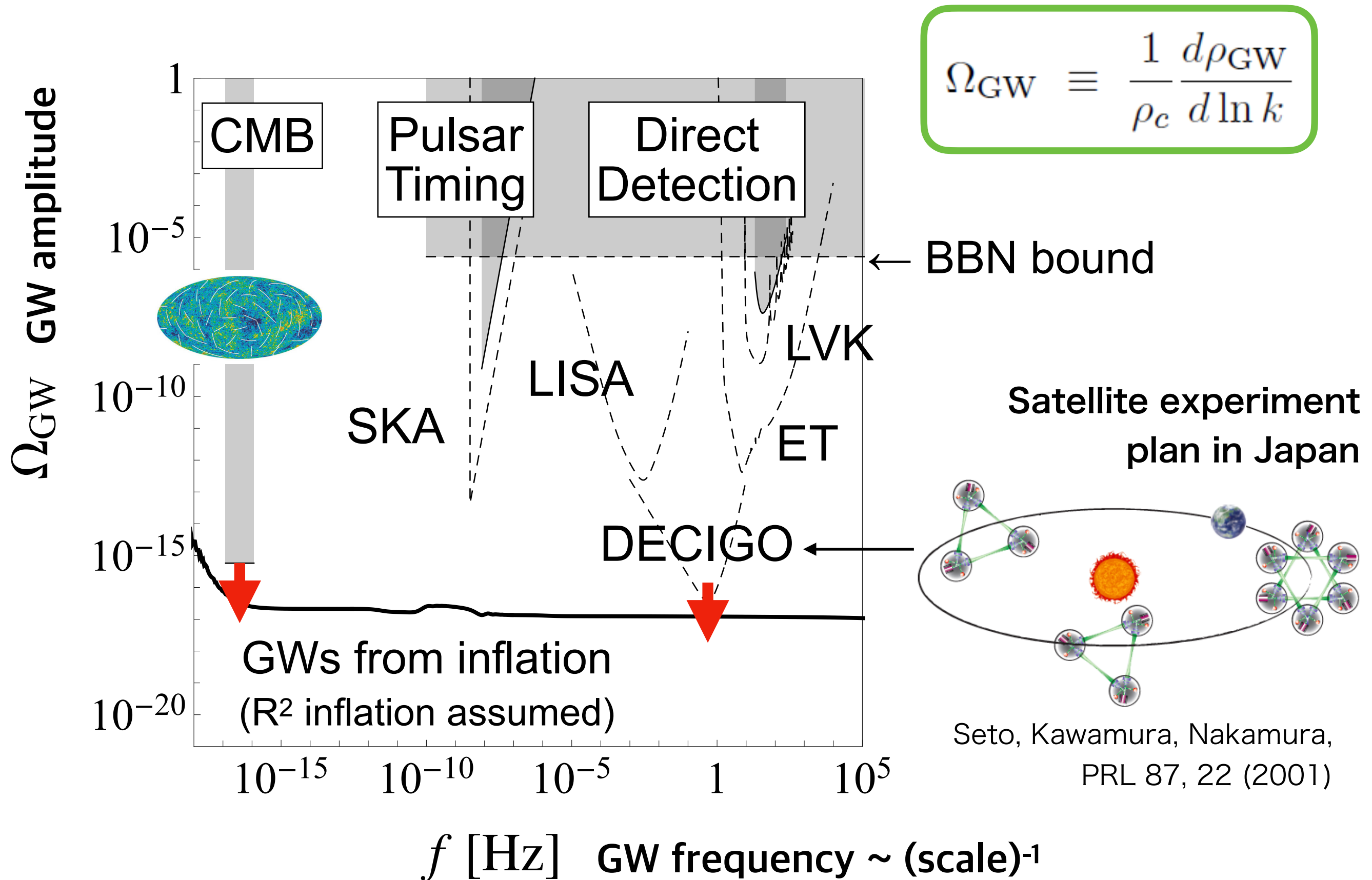
$$r \equiv \frac{\mathcal{P}_{T,\text{prim}}(k)}{\mathcal{P}_{S,\text{prim}}(k)} \simeq 16\epsilon$$

$r = -8n_T$

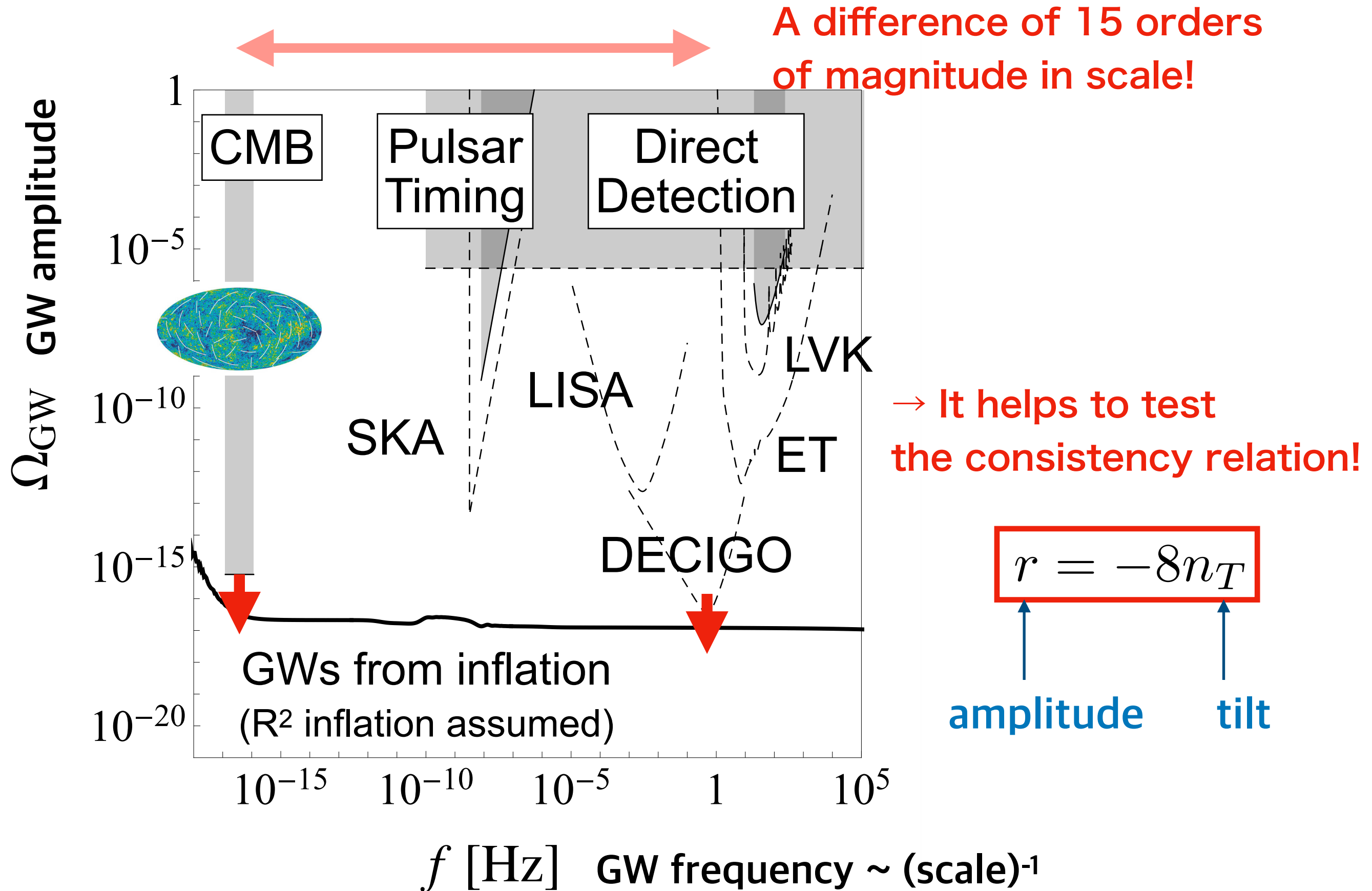
Spectral tilt

$$n_T(k) \equiv \frac{d \ln \mathcal{P}_{T,\text{prim}}(k)}{d \ln k} \simeq -2\epsilon$$

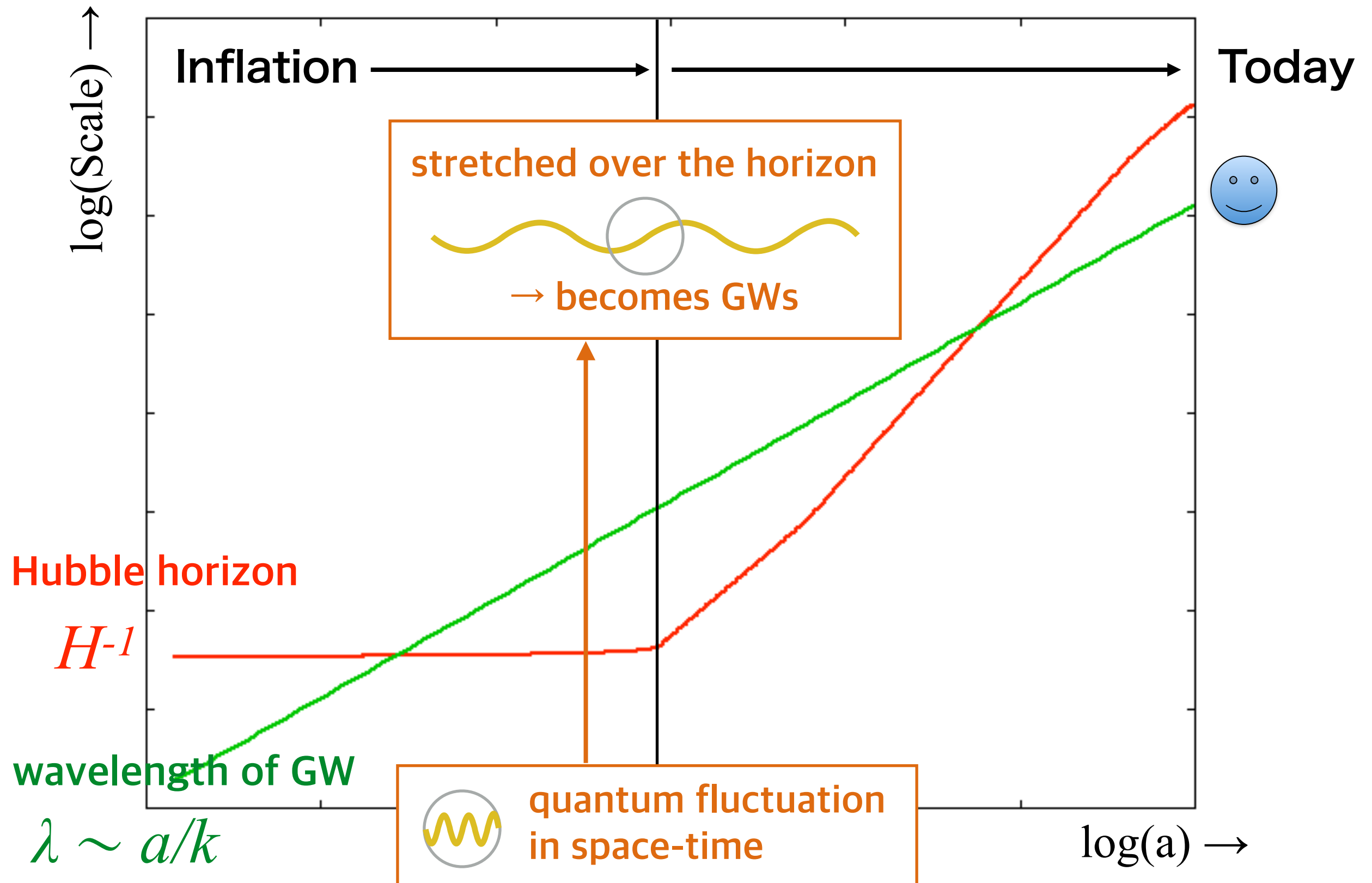
Direct detection of GWs



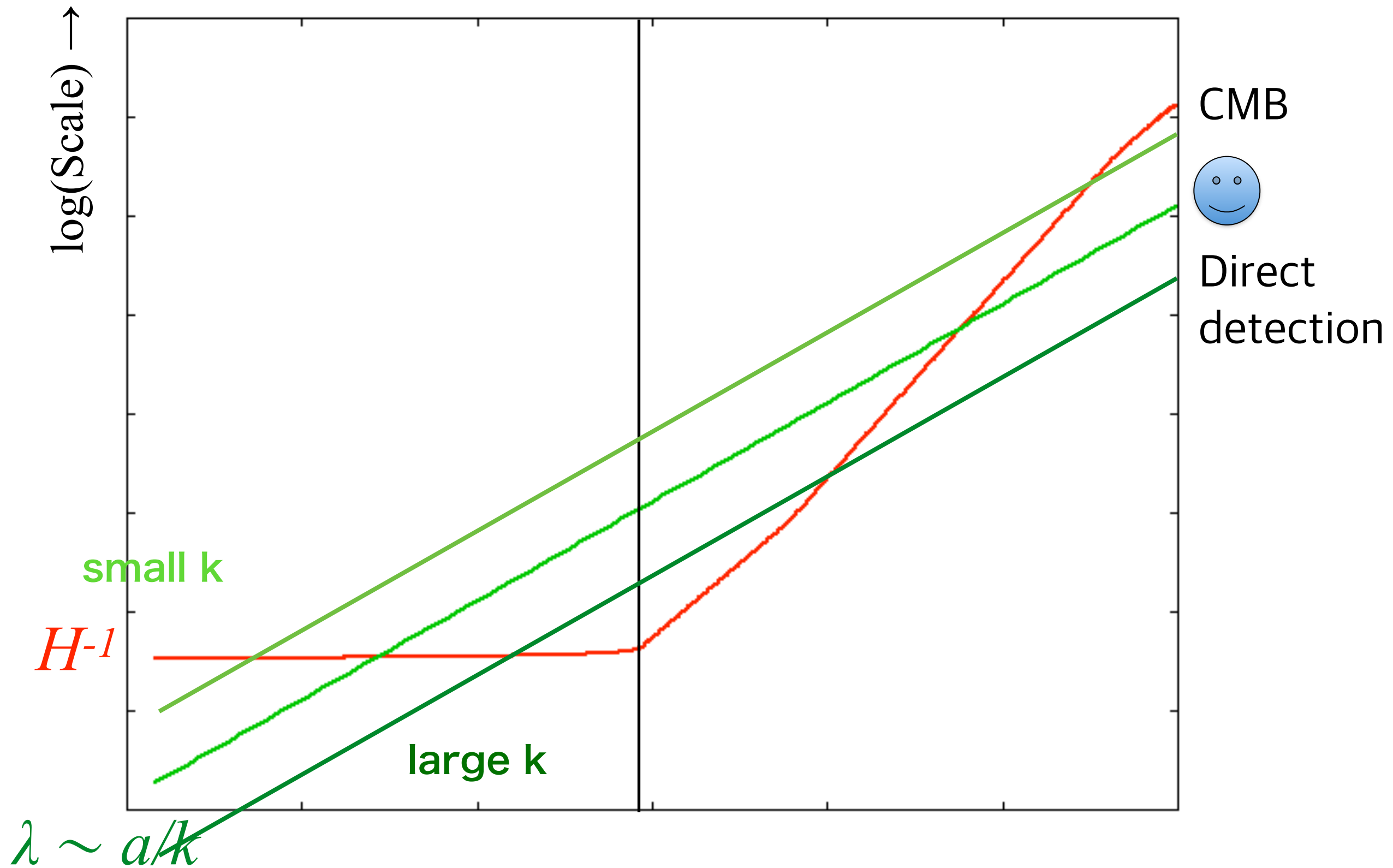
Measuring GWs at two different scales!



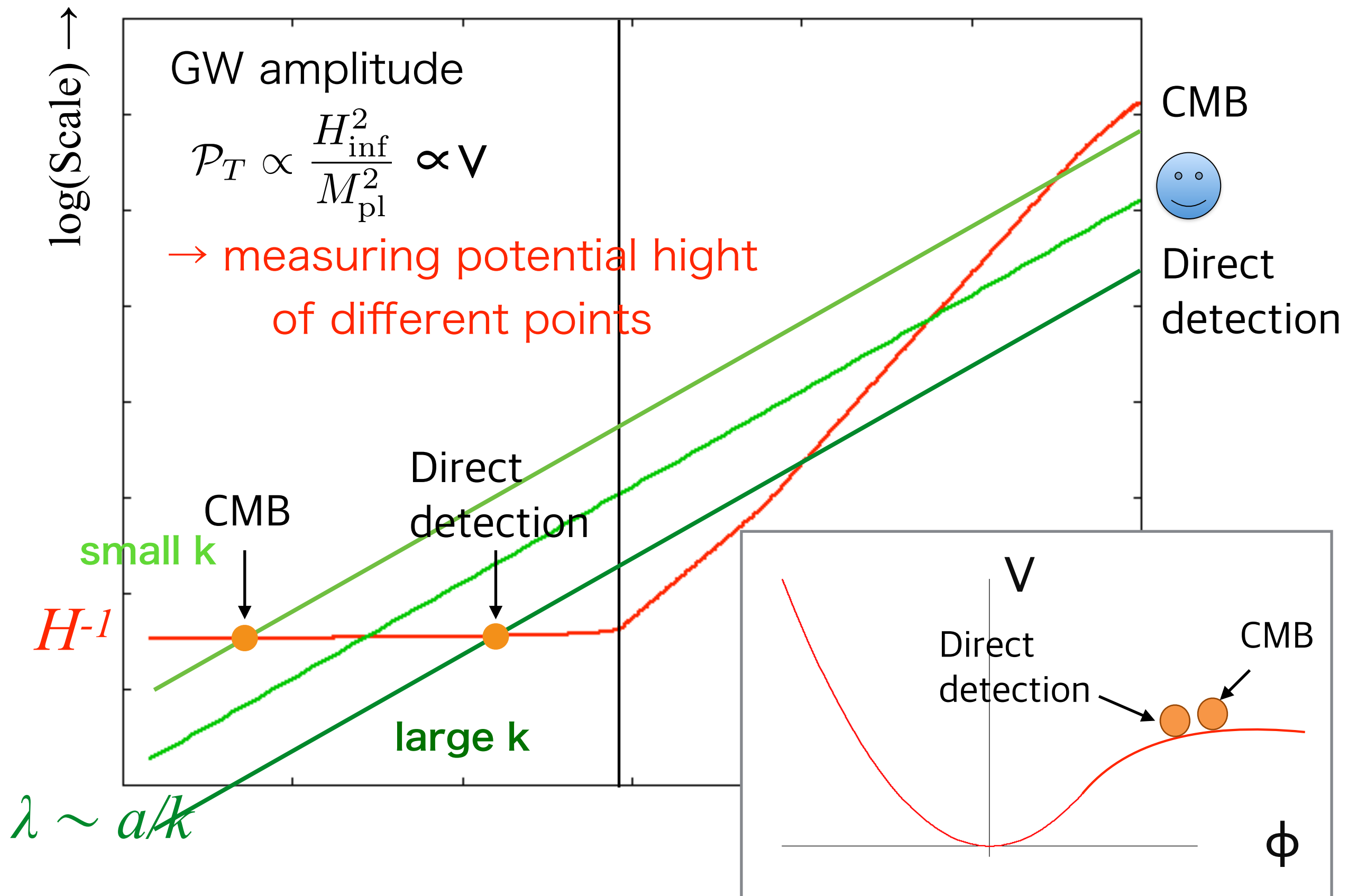
Primordial GWs



Importance of measuring at two scales

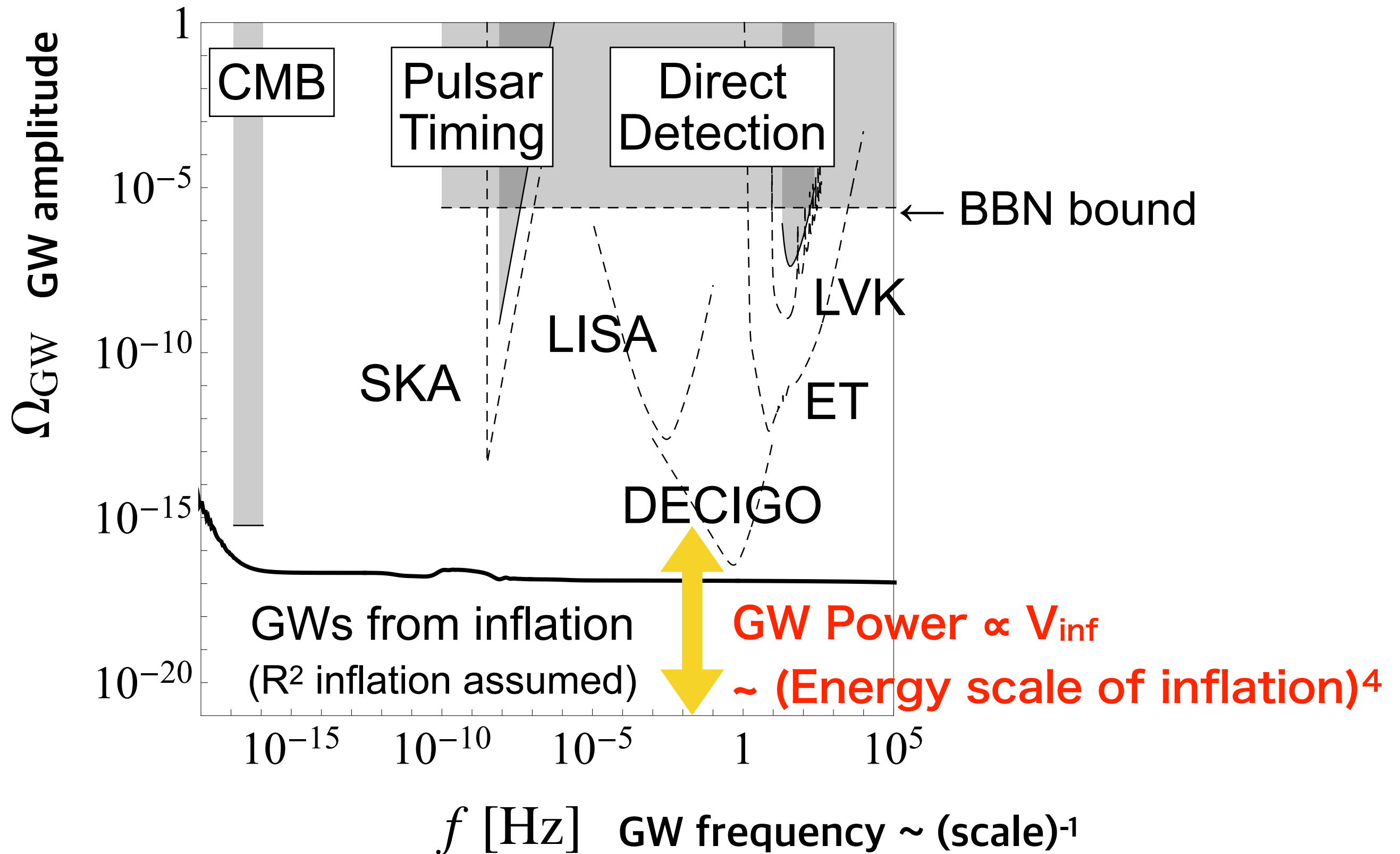


Importance of measuring at two scales



Difficulty I

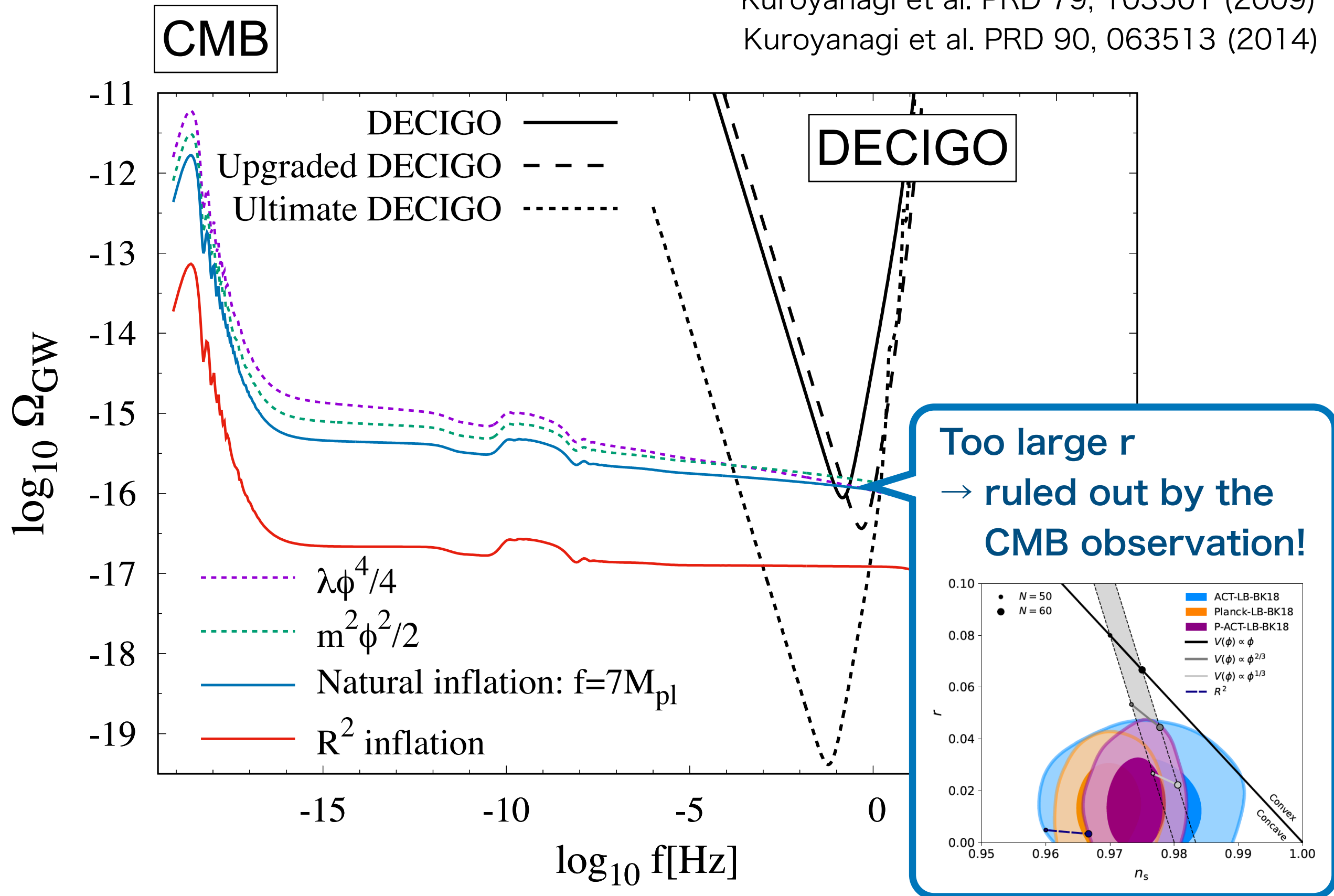
Amplitude is typically very small



Dependence on inflation models

Kuroyanagi et al. PRD 79, 103501 (2009)

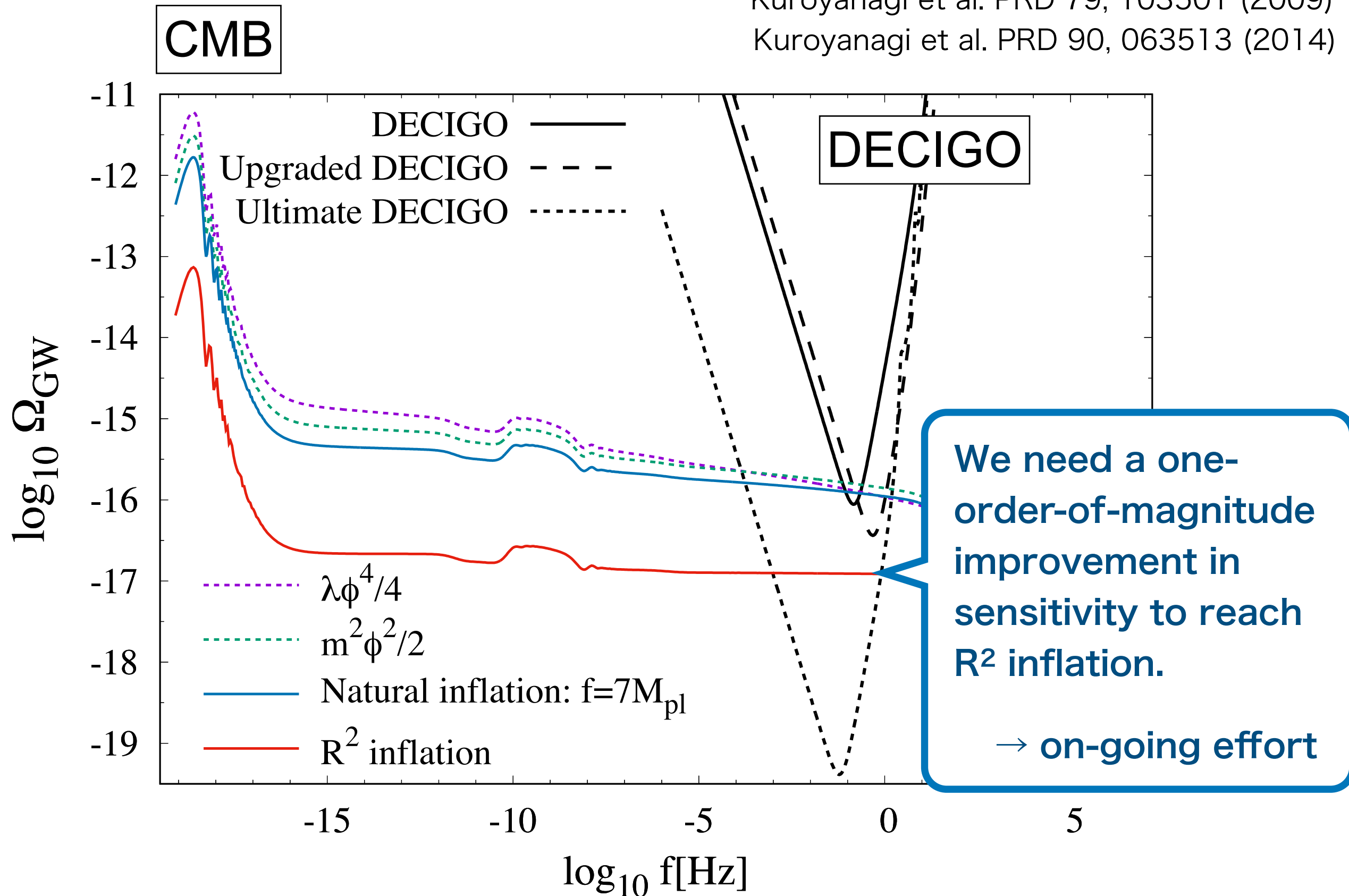
Kuroyanagi et al. PRD 90, 063513 (2014)



Dependence on inflation models

Kuroyanagi et al. PRD 79, 103501 (2009)

Kuroyanagi et al. PRD 90, 063513 (2014)



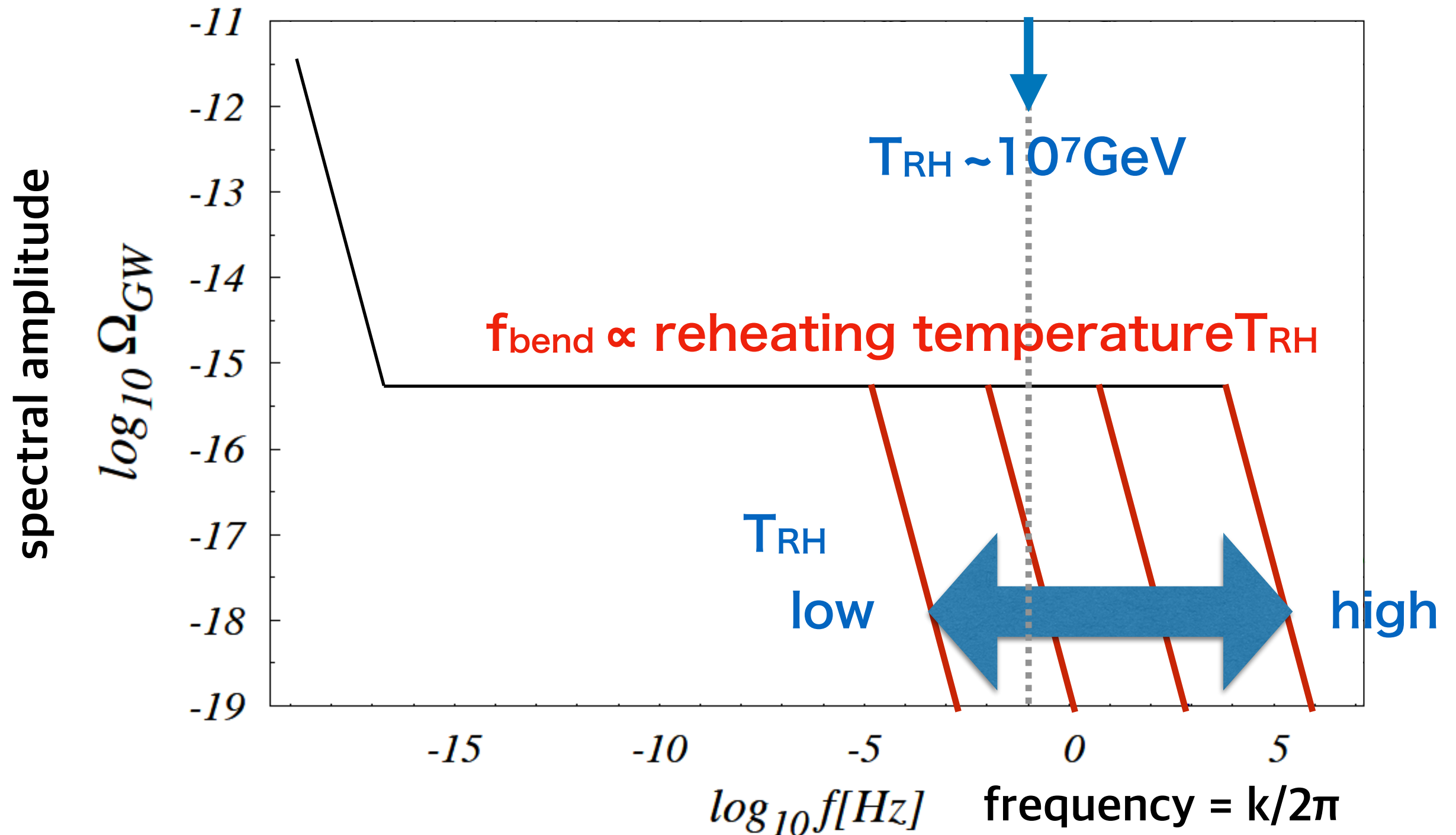
Difficulty 2

High frequency modes may be suppressed by reheating

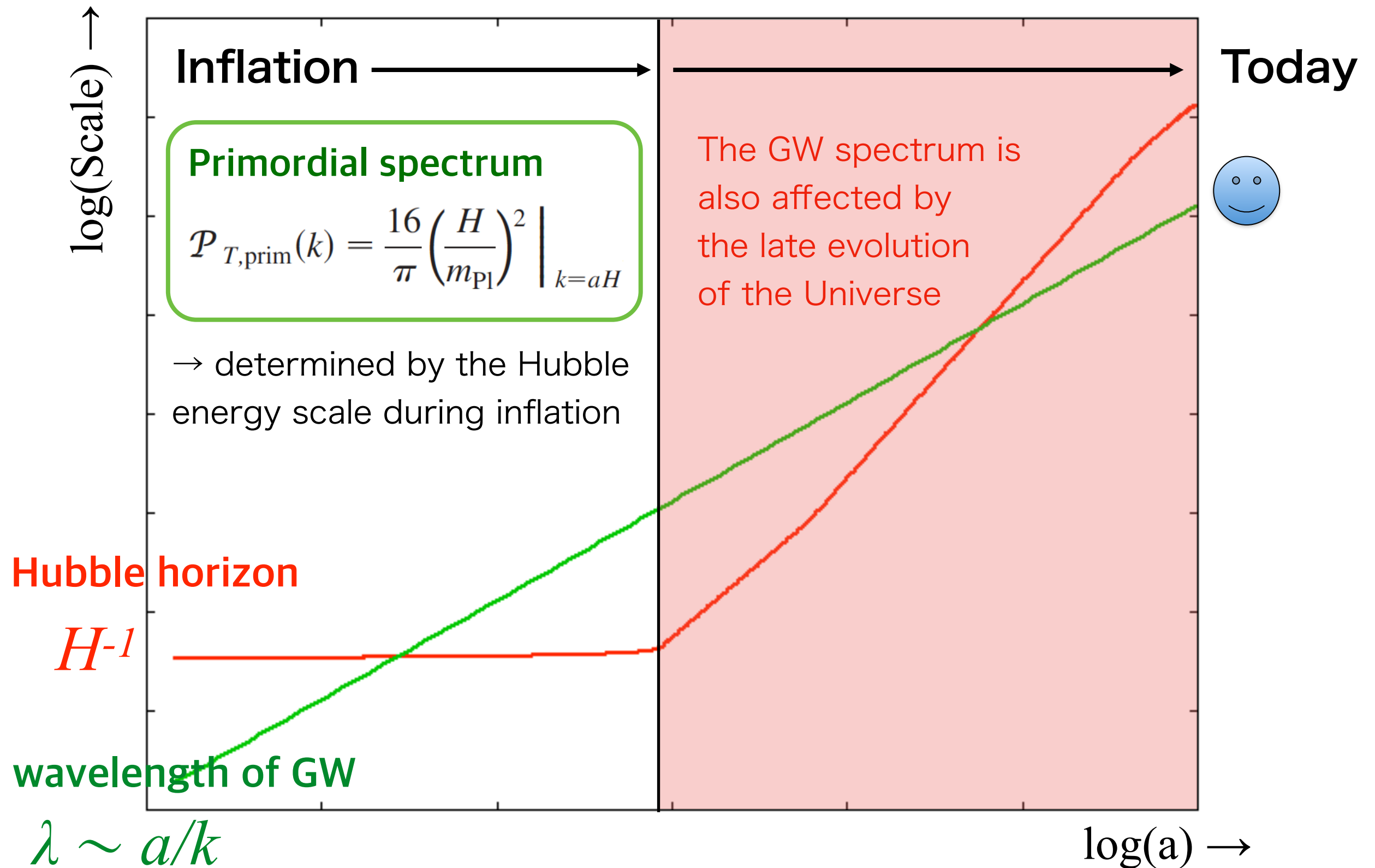
Seto & Yokoyama, JPSJ 72 (2003) 3082-3086

Nakayama et al., JCAP 06, 020 (2008), PRD 77, 124001 (2008)

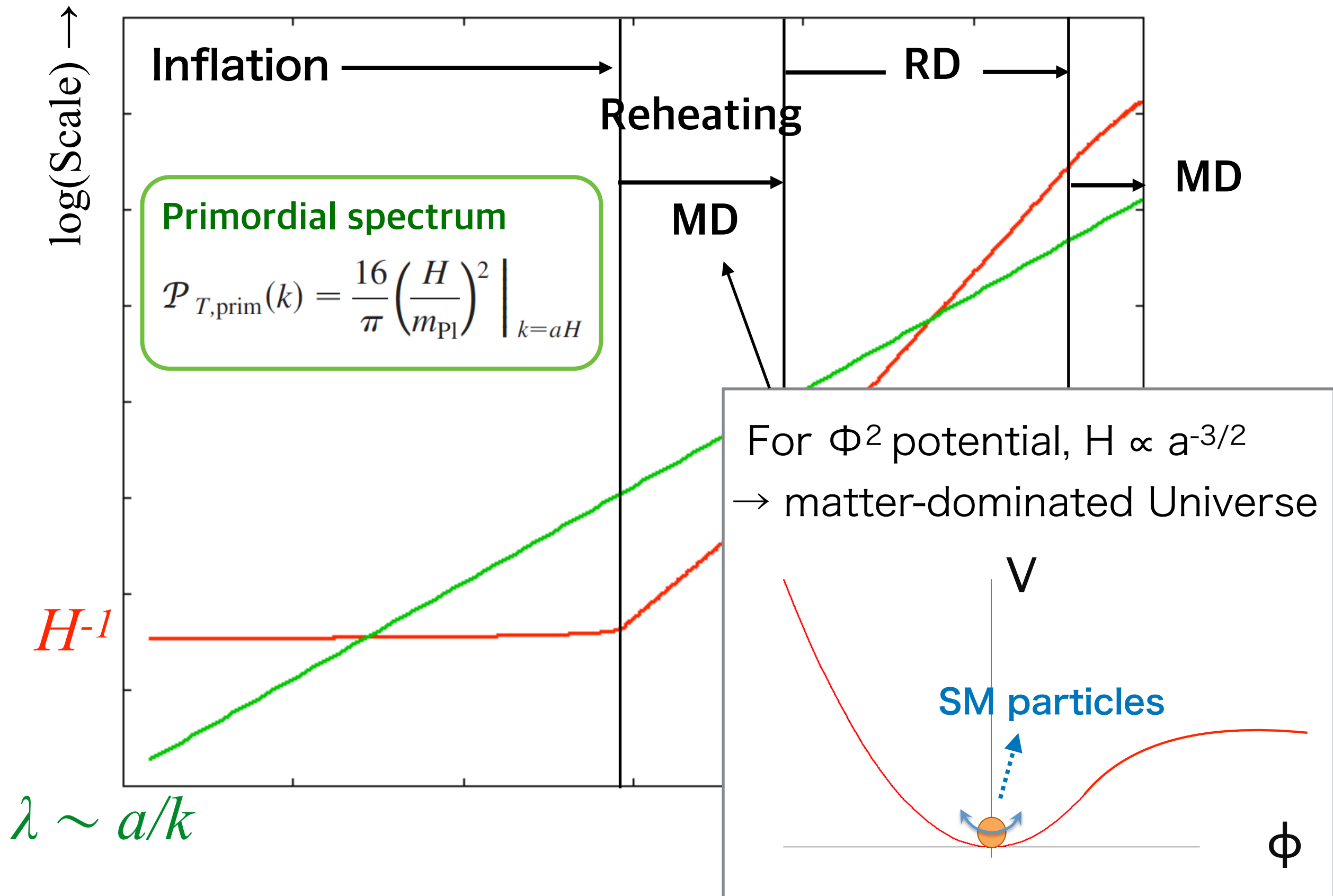
DECIGO's target frequency
0.1 Hz



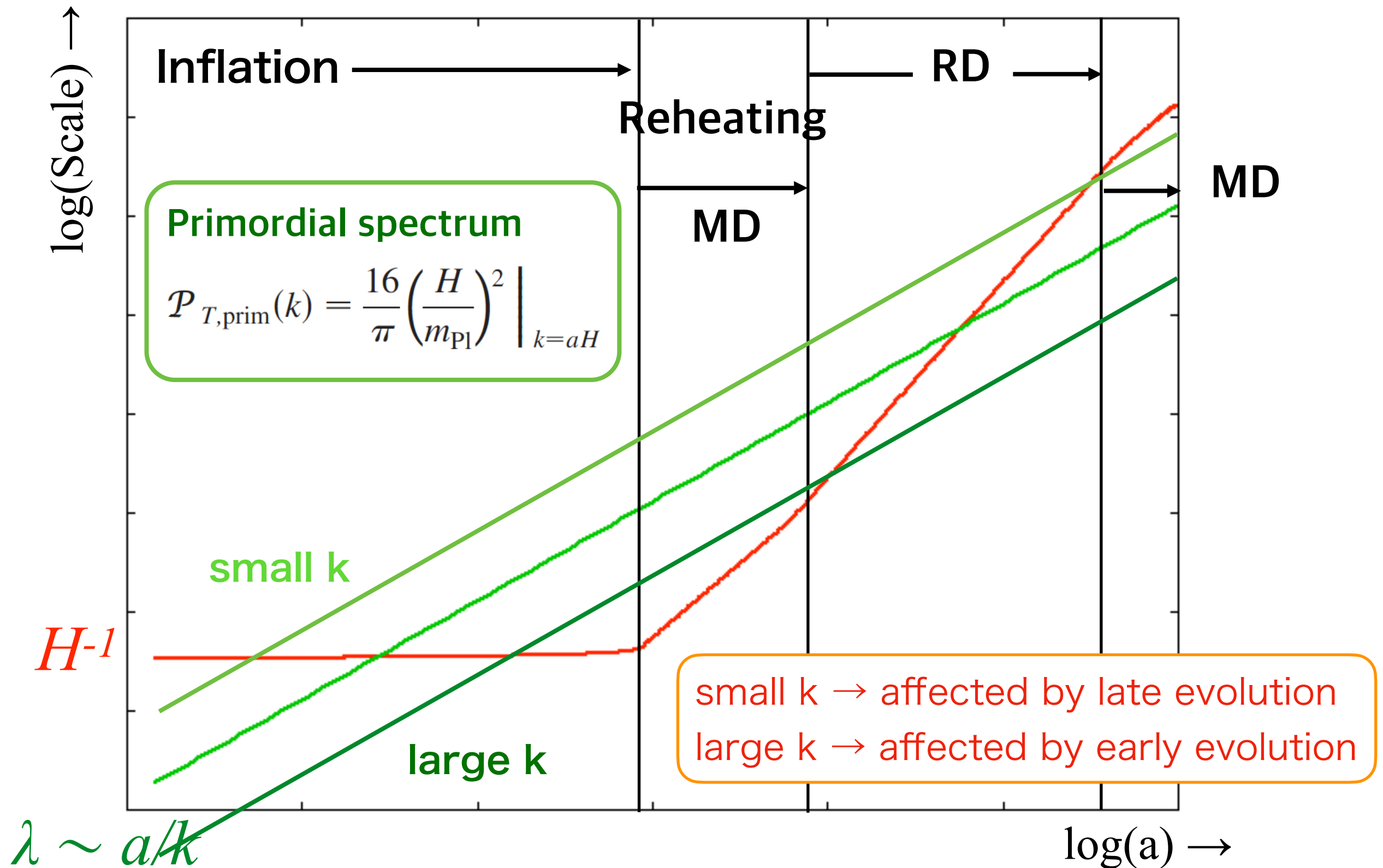
Early history of the Universe



Early history of the Universe

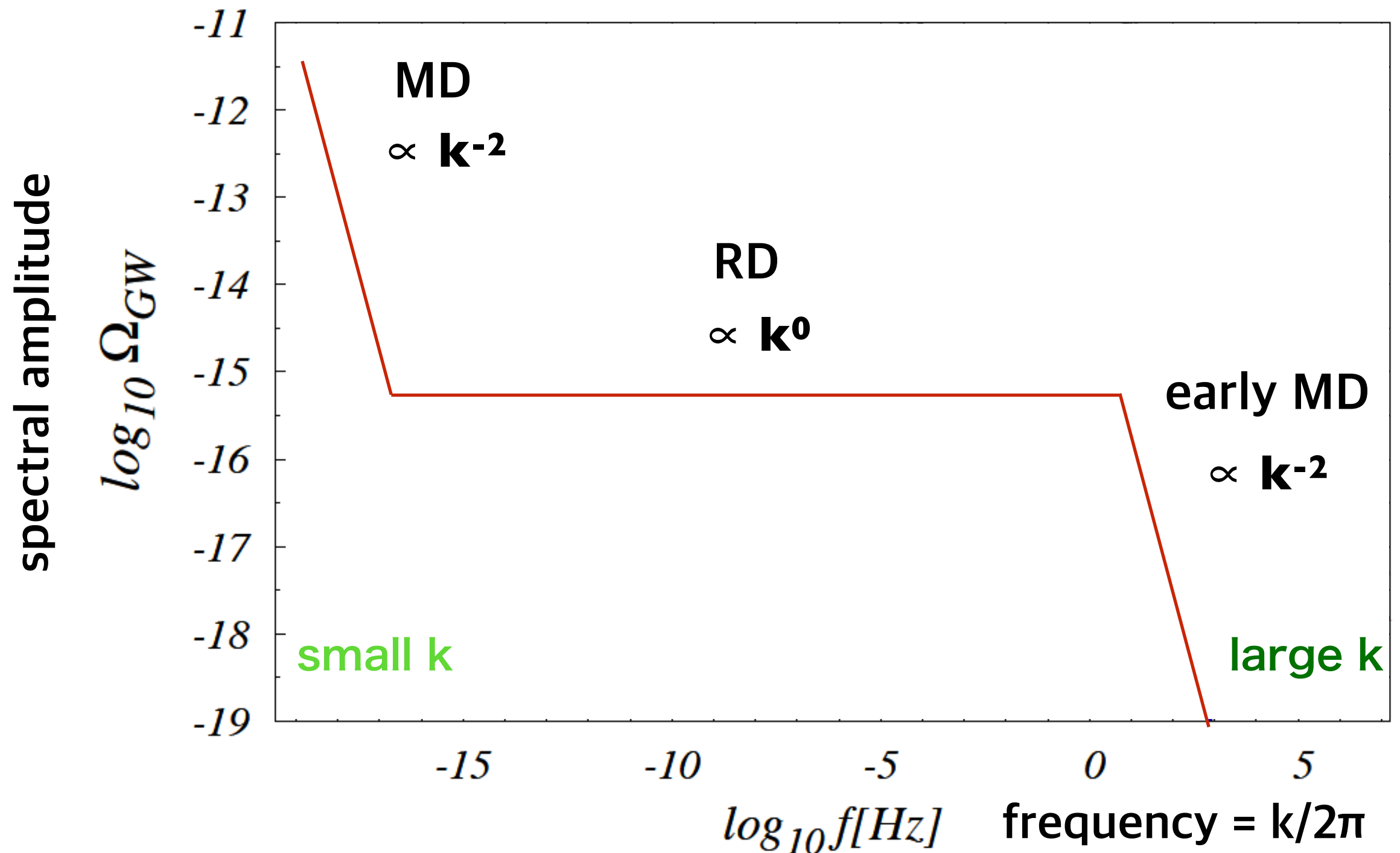


Early history of the Universe



Effect of the early matter phase

We need $T_{\text{RH}} > 10^7 \text{ GeV}$ to have the signal at the DECIGO frequency!

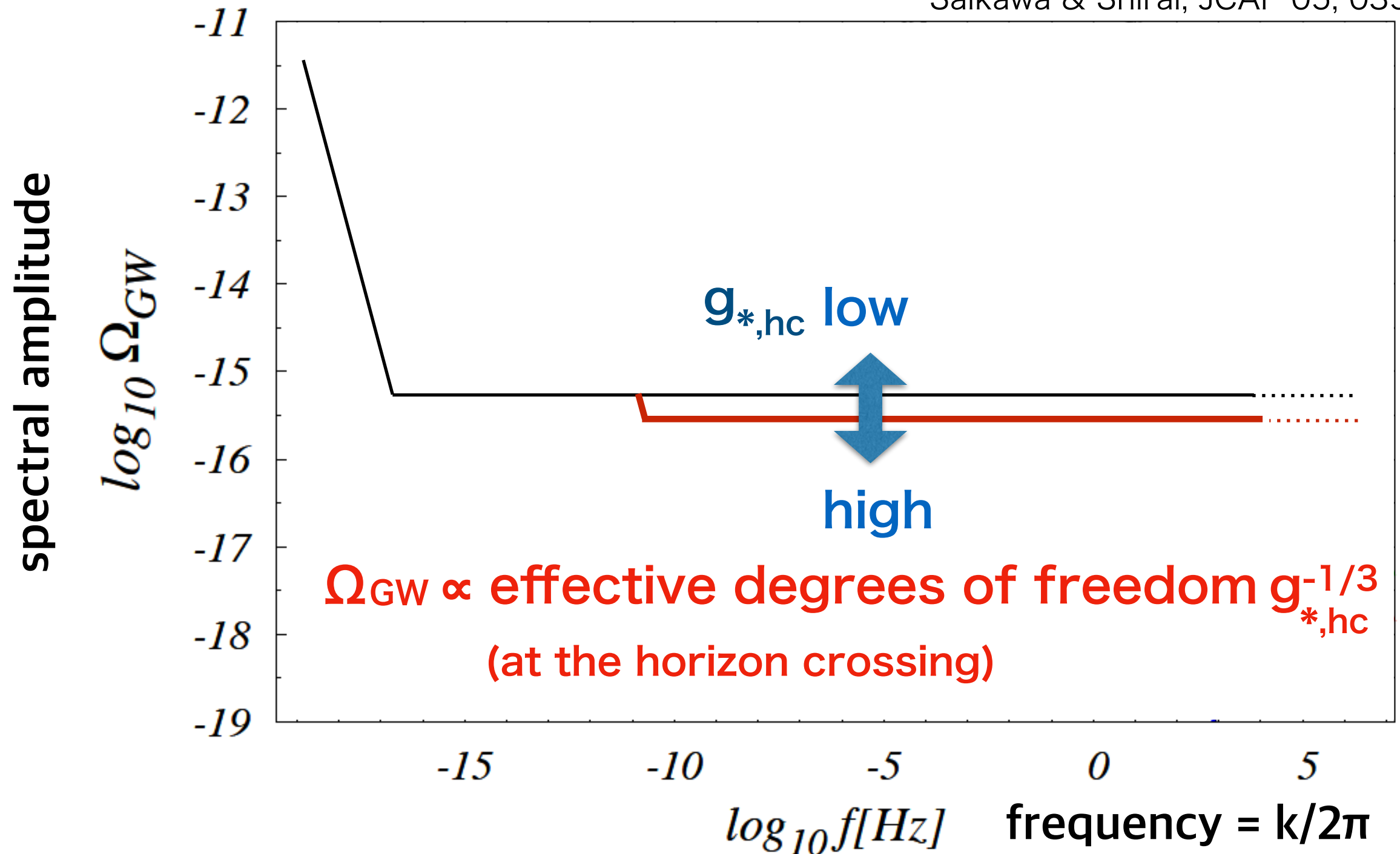


Difficulty 3

Ambiguity in predicting the amplitude
due to **the effective degrees of freedom**

Watanabe & Komatsu, PRD 73, 123515 (2006)

Saikawa & Shirai, JCAP 05, 035 (2018)



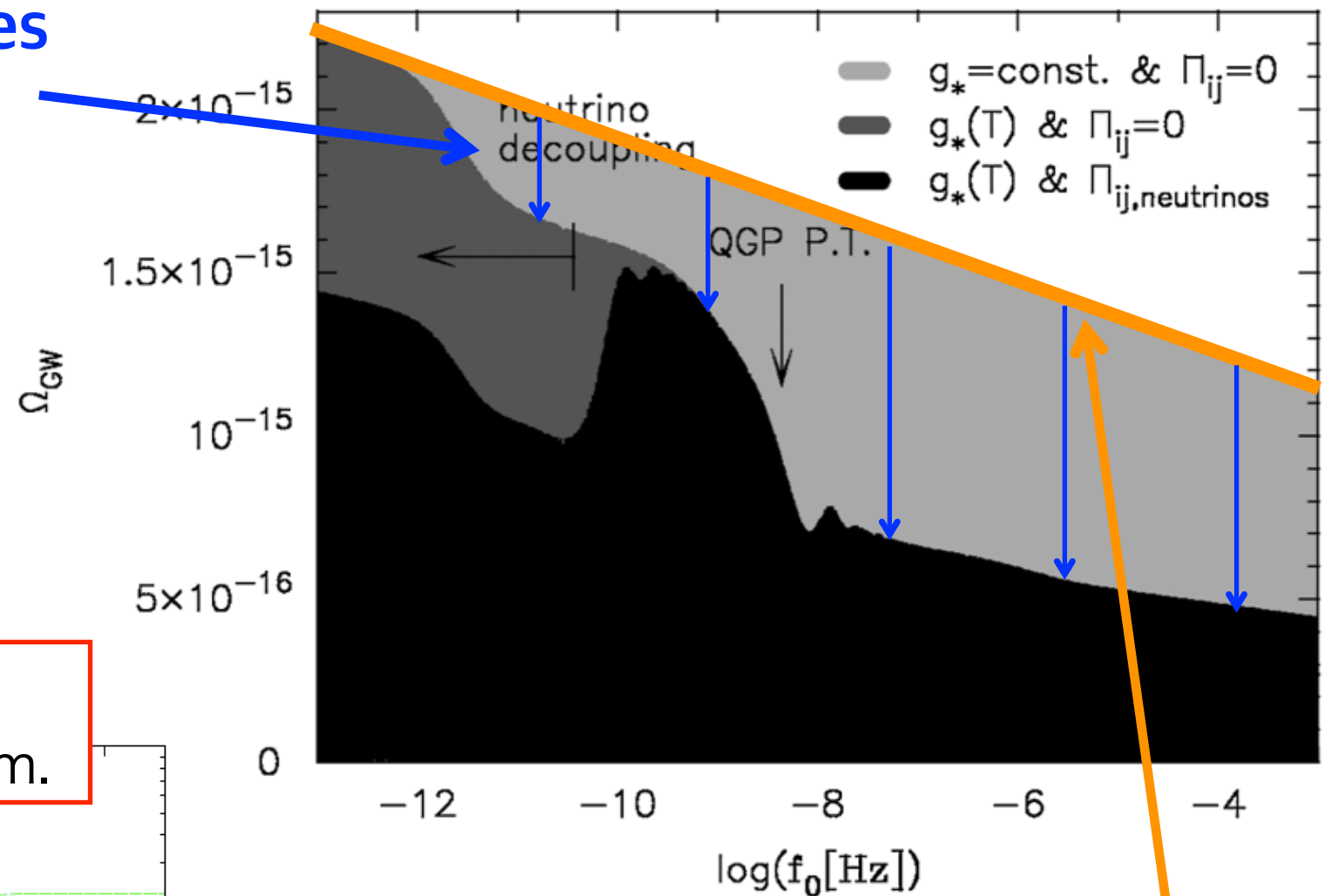
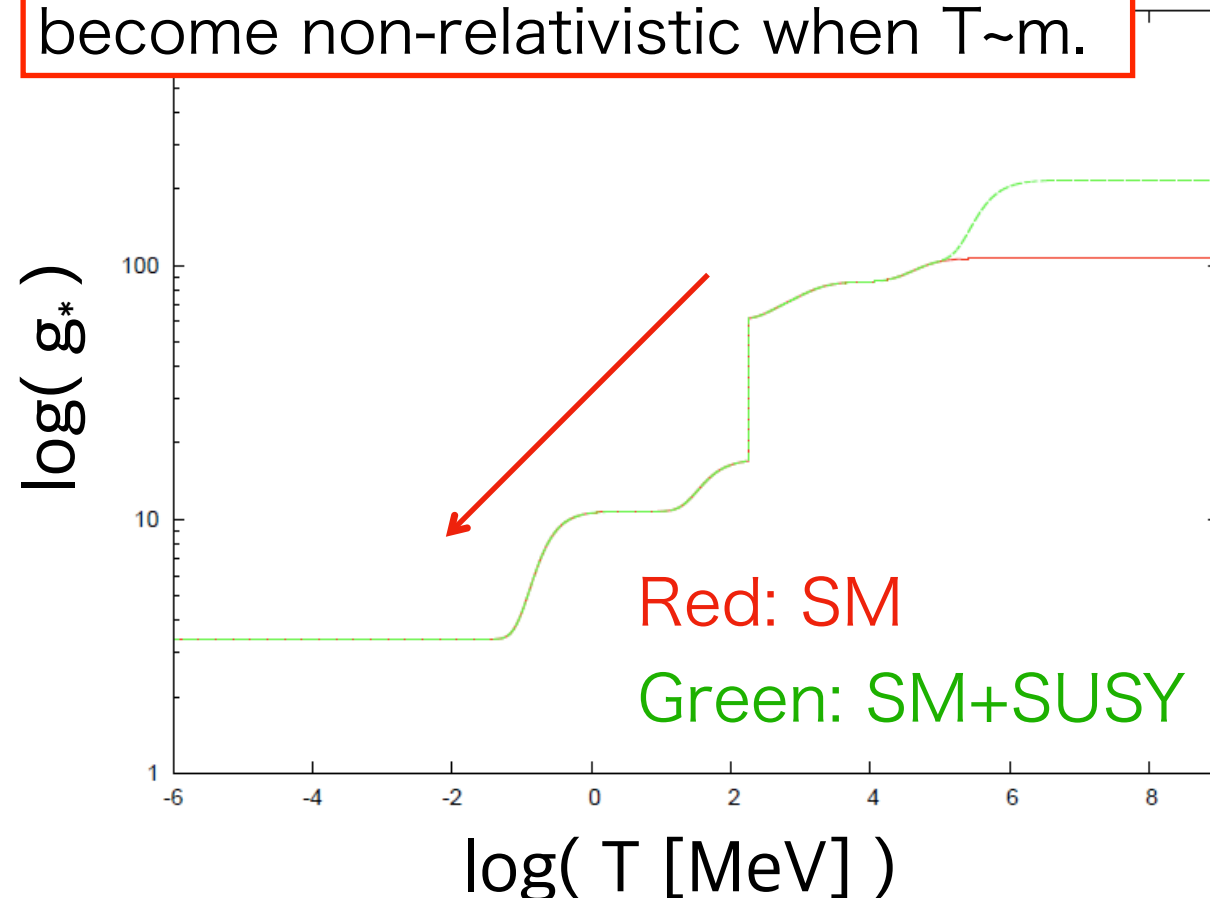
Effective number of degrees of freedom

Damping due to the changes in effective number of degrees of freedom g_*

$$\text{density } \rho(T) = \frac{\pi^2}{30} g_*(T) T^4,$$

$$\text{entropy } s(T) = \frac{2\pi^2}{45} g_{*s}(T) T^3$$

As T decreases, each particles become non-relativistic when $T \sim m$.



primordial spectrum with tilt

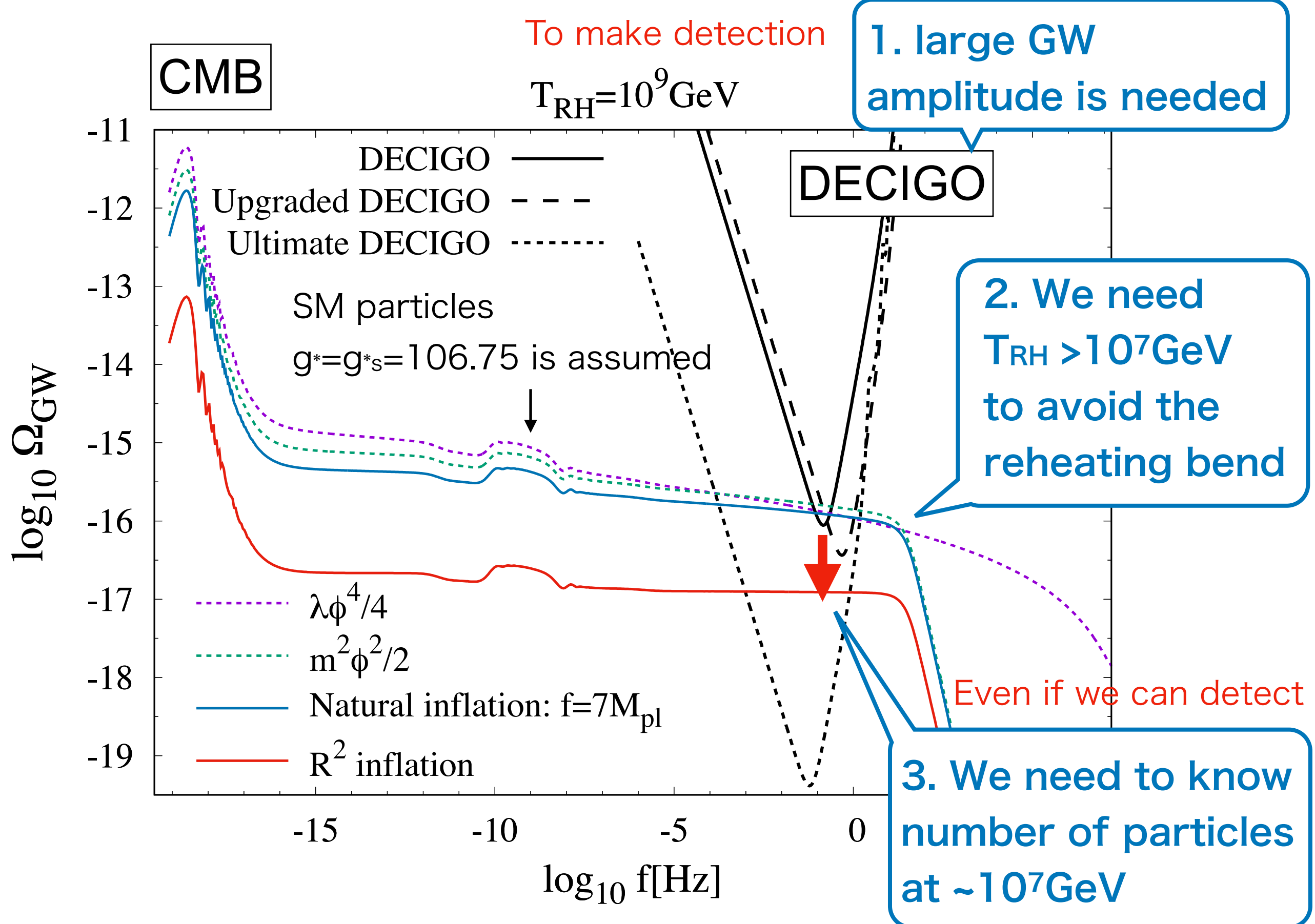
For SM particles, $g^* = g_{*s} = 106.75$

If we have non-SM particles,

$$g^*, g_{*s} > 106.75$$

→ More damping in the GW amplitude

Three challenges to measure n_T



Summary

Several observations can be used to test inflation

1. Accelerated expansion of the Universe

- Required to solve the horizon / flatness / monopole problems

2. Nearly scale-invariant SCALAR perturbations

- Almost confirmed by CMB & large scale observations!

3. Nearly scale-invariant TENSOR perturbations

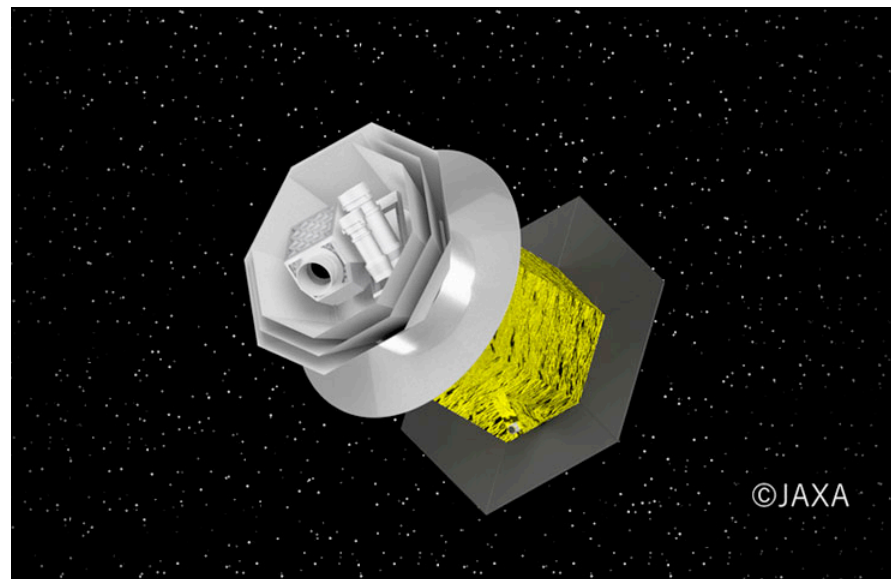
- Key observations to confirm inflation
- Testing **the consistency relation** through CMB B-modes and direct gravitational wave detection would offer further, stronger evidence

Summary

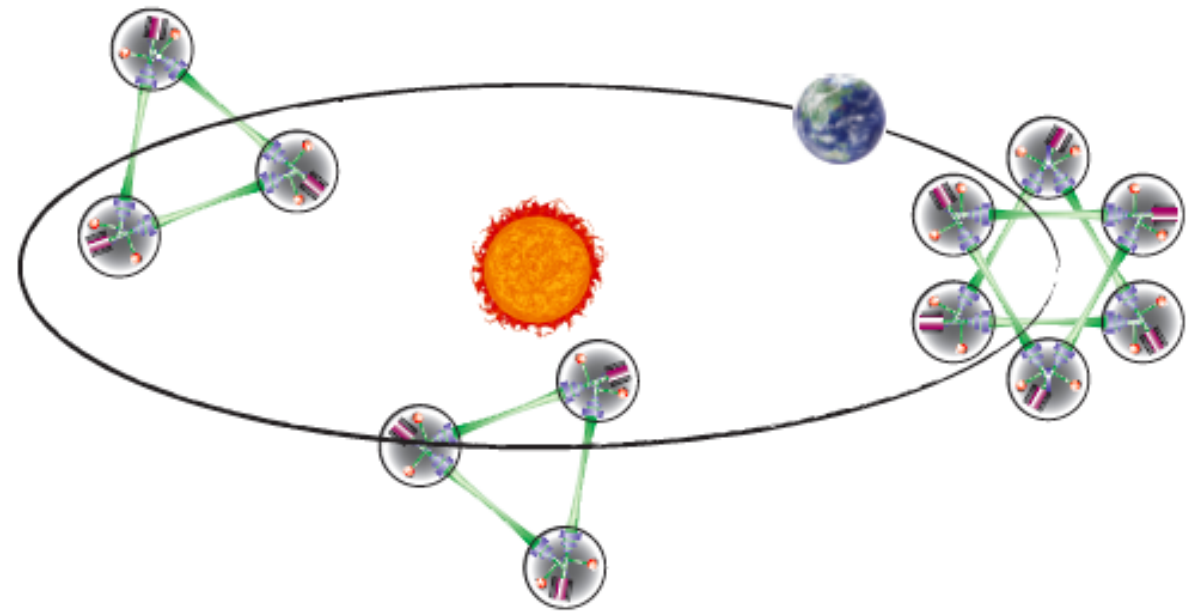
Two Japanese experiments are the key to confirm inflation

Let's push it forward!

LiteBIRD



DECIGO



3. Nearly scale-invariant TENSOR perturbations

- Key observations to confirm inflation
- Testing **the consistency relation** through CMB B-modes and direct gravitational wave detection would offer further, stronger evidence

Possible discussion

Further observational tests?

1. Accelerated expansion of the Universe

- Any way to measure the expansion directly???

2. Nearly scale-invariant SCALAR perturbations

- Gaussianity (Non-Gaussianity) of perturbations
- Quantum nature of perturbations

3. Nearly scale-invariant TENSOR perturbations

- Gaussianity (Non-Gaussianity) of perturbations
- Quantum nature of perturbations

Possible discussion

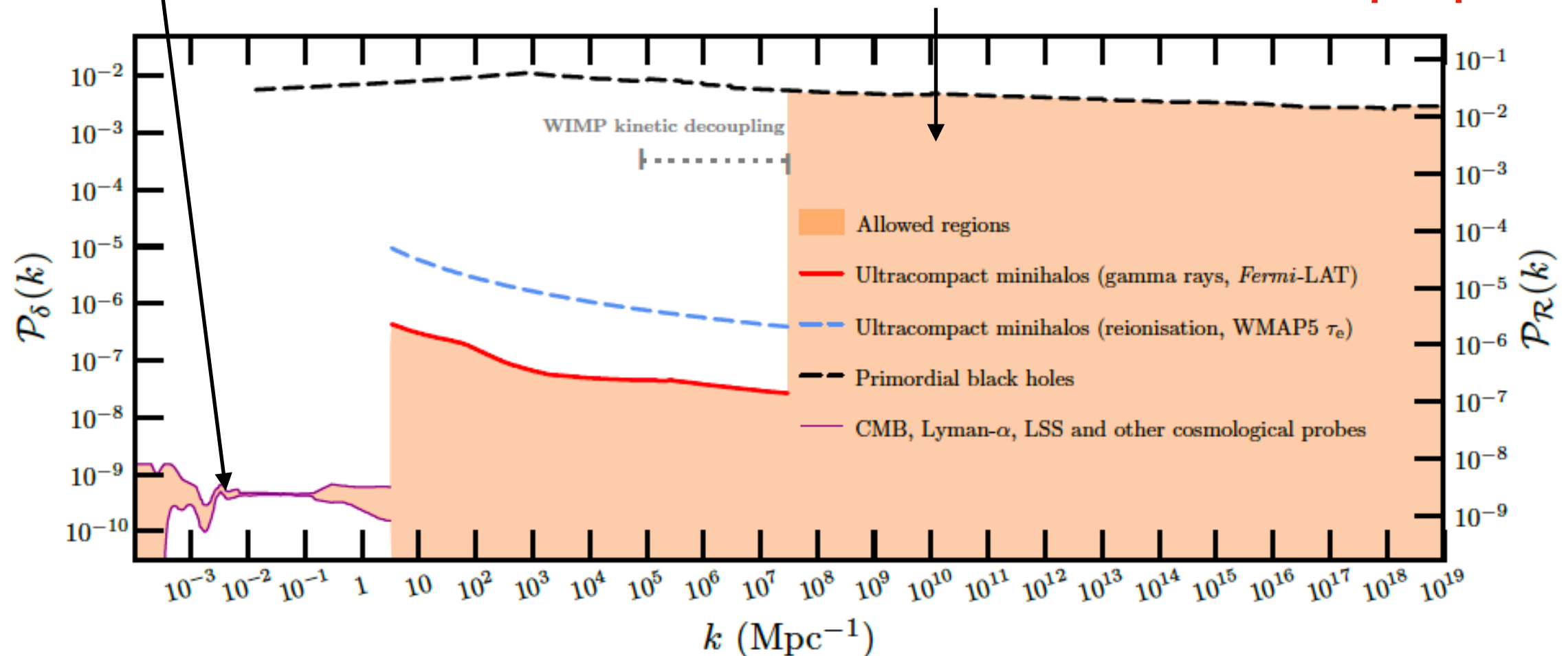
Non-linearity appears when we add extra ingredients

- Gauge field during inflation
- Formation of primordial black holes

$P_s \sim 2 \times 10^{-9}$ @ CMB scale

can be much larger at small scales

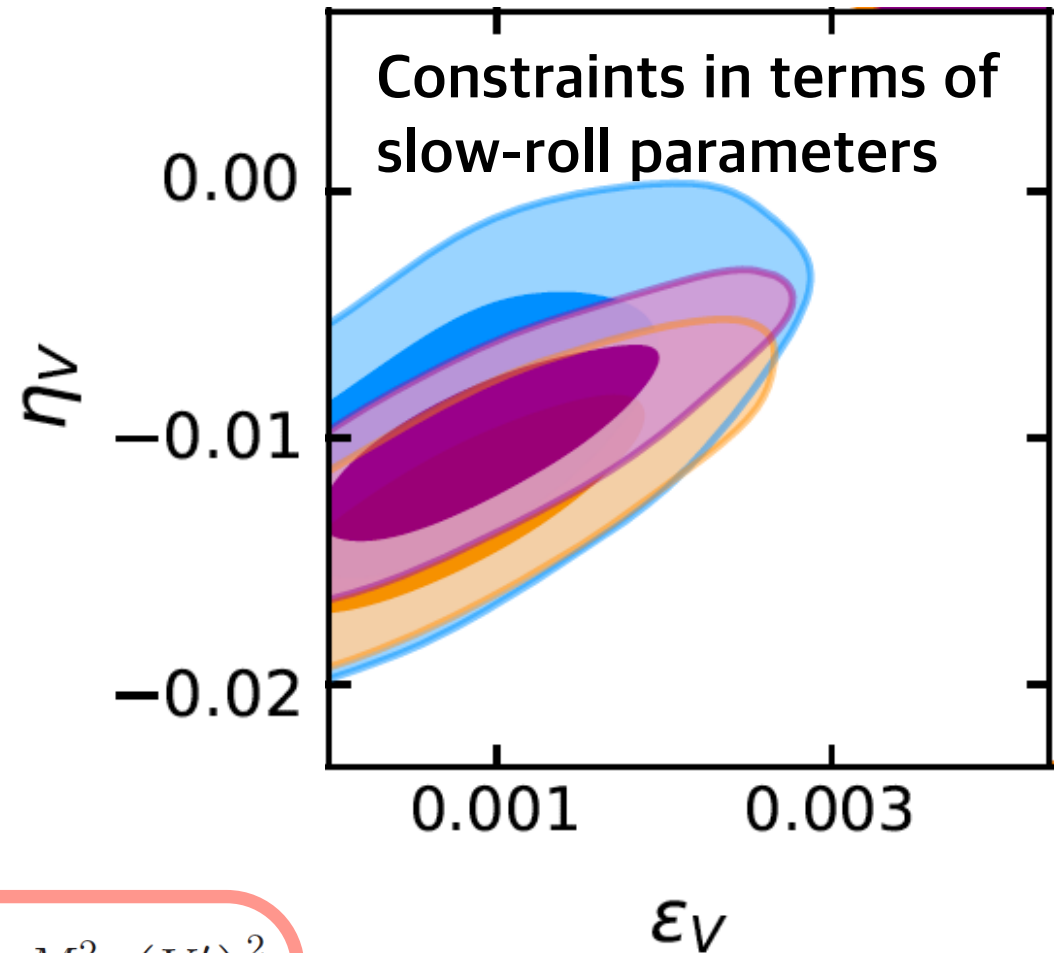
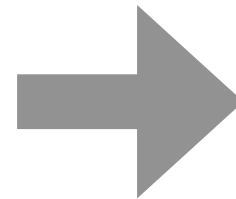
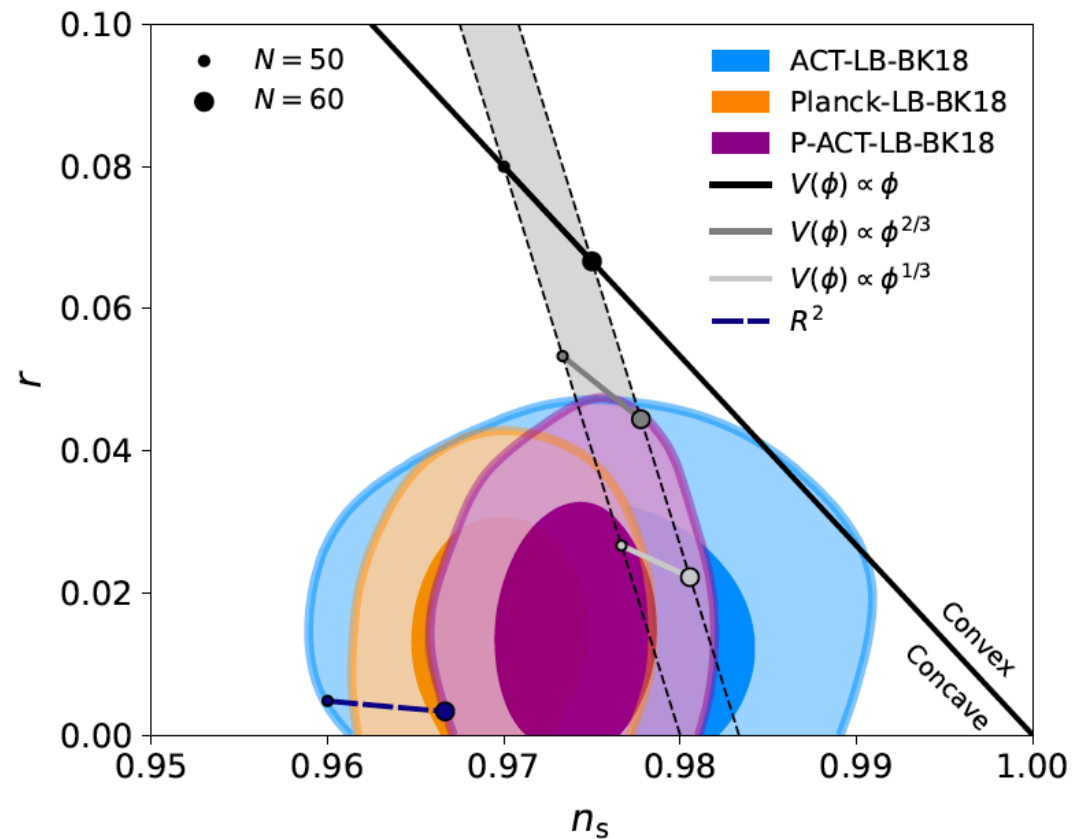
2nd order GWs is a unique probe



→ **developing a precise analytic formalism is still an active area of research**

How is it likely to have large r ?

(added after a comment from Yanagida-san)



Spectral tilt $n_s(k) - 1 \equiv \frac{d \ln \Delta_{\zeta, \text{prim}}^2(k)}{d \ln k} \simeq -6\epsilon + 2\eta$ $\epsilon_V \equiv \frac{M_{\text{Pl}}^2}{2} \left(\frac{V'}{V} \right)^2$

Tensor-to-scalar ratio $r \equiv \frac{\mathcal{P}_{T, \text{prim}}(k)}{\mathcal{P}_{S, \text{prim}}(k)} \simeq 16\epsilon$ $\eta_V \equiv M_{\text{Pl}}^2 \frac{V''}{V}$

A small ϵ (small r) indicates a large negative η , which leads $\epsilon \ll |\eta|$

My personal optimistic view

The idea of 'naturalness' can vary from person to person, but if we assume that it is natural to have $\epsilon \sim |\eta|$, then we would not expect ϵ to be extremely small.

→ **Hope for observable tensor mode!**