

# **Discussions on Shallow and deep inelastic scattering**

**Shunzo Kumano**

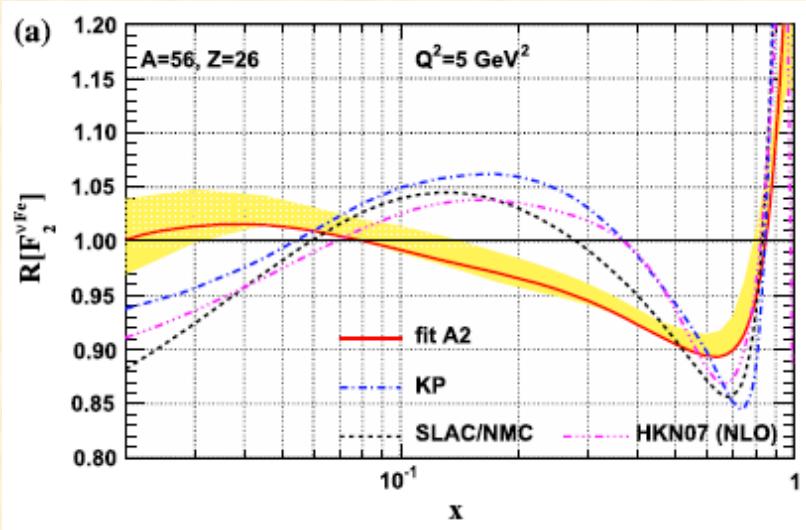
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J-PARC Center (J-PARC)**  
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**<http://research.kek.jp/people/kumanos/>**

**10th International workshop on Neutrino-nucleus interactions in the few-GeV region,  
16-21 November 2015, Suita Campus of Osaka University, Japan**  
**<http://indico.ipmu.jp/indico/conferenceDisplay.py?ovw=True&confId=46>**

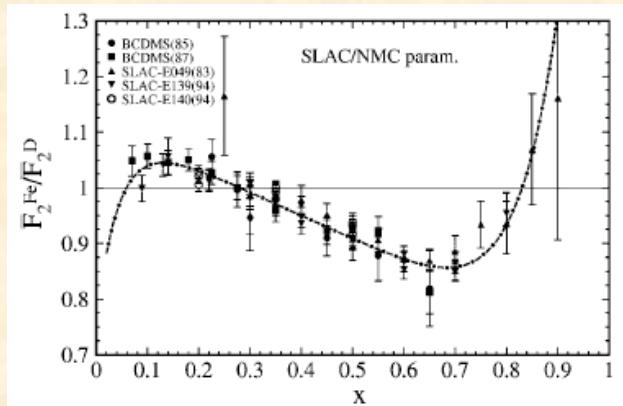
**November 19, 2015**

1. Nuclear modification difference between  $\ell^\pm$  and  $\nu$
2. “Isoscalar ion” data
3. Nuclear antiquark distributions at  $x \sim 0.1$  and pion contributions

# Nuclear modification differences between $\ell^\pm$ and $\nu$

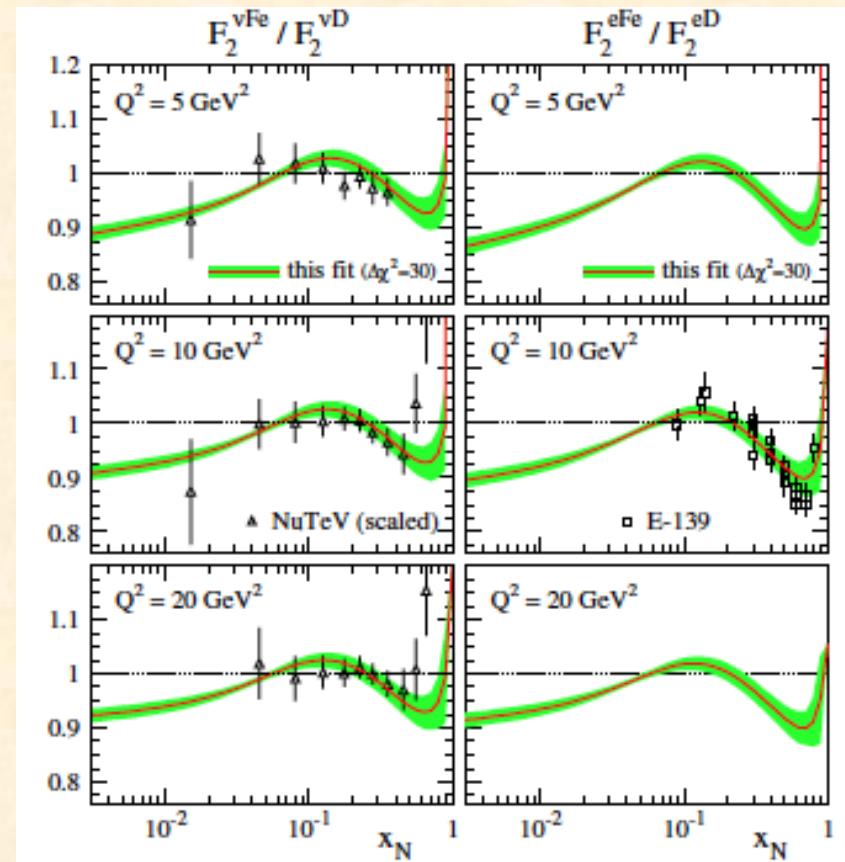


I. Schienbein *et al.*,  
 PRD 77 (2008) 054013

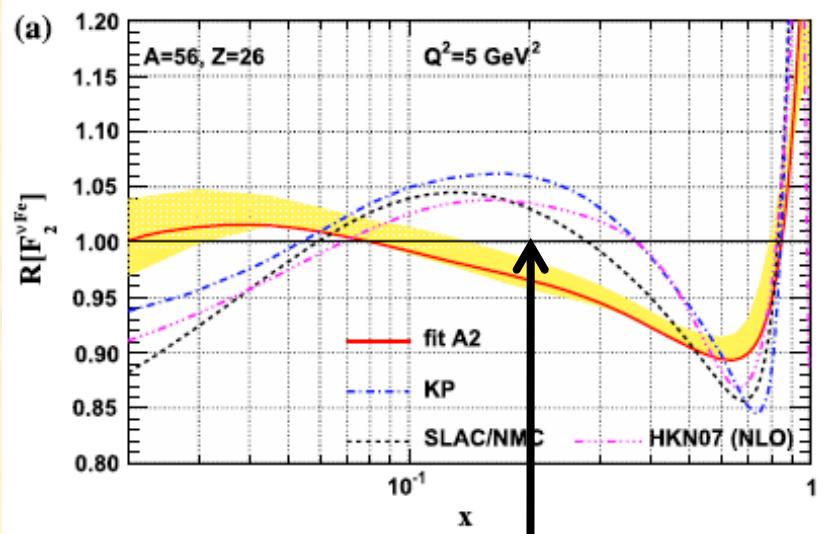


Charged-lepton scattering

D. de Florian, R. Sassot,  
 P. Zurita, M. Stratmann,  
 Phys. Rev. D 85 (2012) 074028.



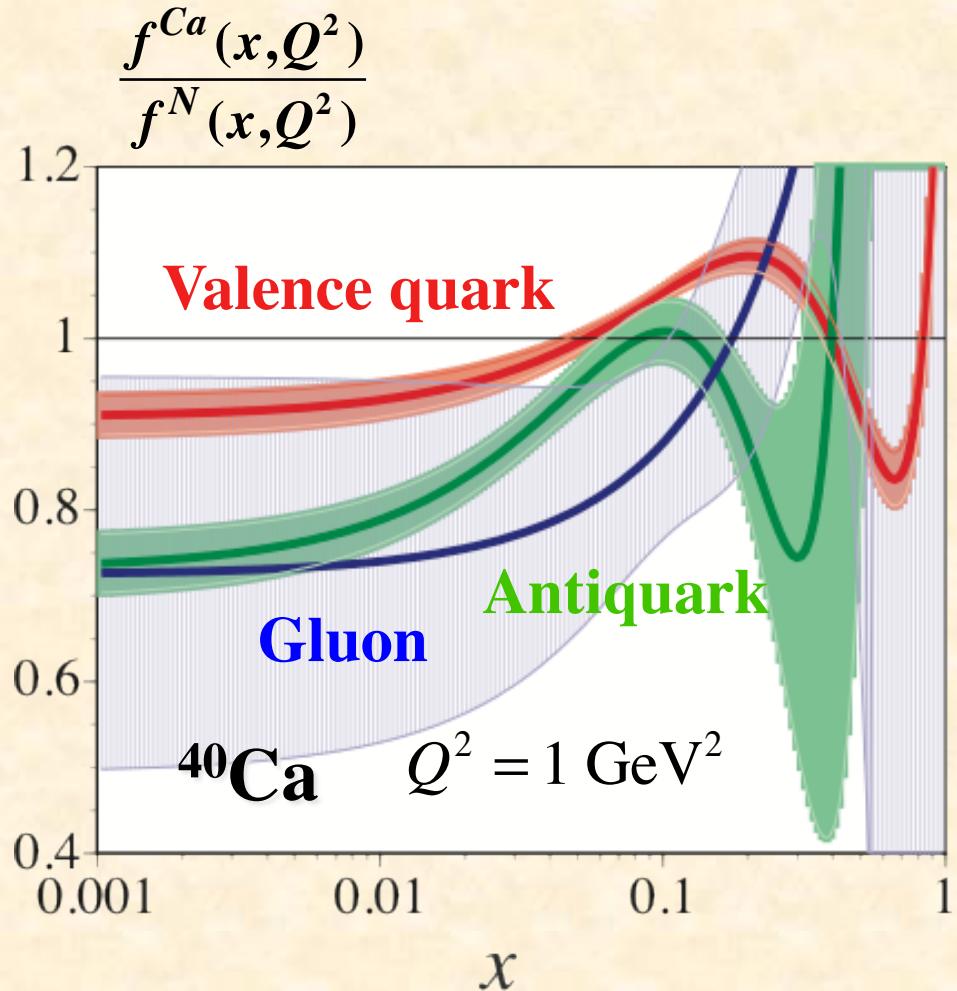
It is (almost) impossible to obtain  
 nuclear corrections for PDFs only from neutrino DIS.



The issue could be solved experimentally by

- (1) Fermilab-E906:  $\frac{\sigma_{pA}}{\sigma_{pD}} \approx \frac{\bar{q}_A}{\bar{q}_D},$
- (2) Minerva:  $\frac{F_3^A}{F_3^{A'}} \approx \frac{q_v^A}{q_v^{A'}}$

in the near future if the data are precise.



HKN2007, M. Hirai *et al.*,  
PRC 76 (2007) 065207.

# “Isoscalar iron” data

Experiment	Target	$\nu$ energy (GeV)
CCFR	Fe	30-360
CDHSW	Fe	20-212
CHORUS	Pb	10-200
NuTeV	Fe	30-500

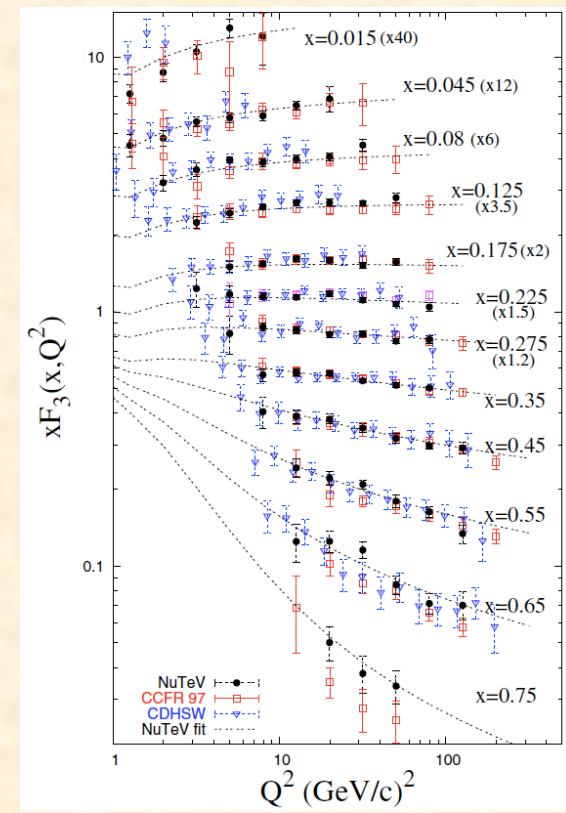
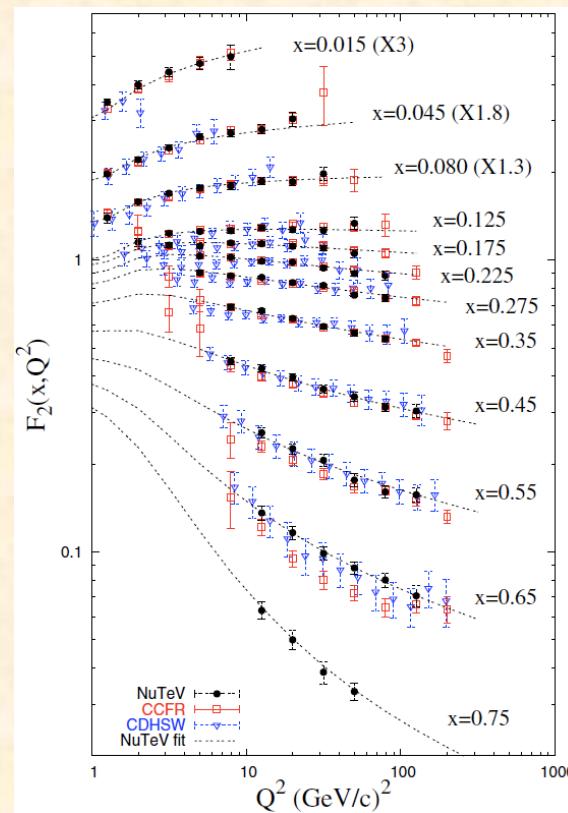
MINERvA (He, C, Fe, Pb), ...

Structure-function data  
for “isoscalar iron”



Often quoted as  
neutrino-nucleon  
scattering data!?

M. Tzanov *et al.* (NuTeV),  
PRD74 (2006) 012008.



# Isoscalar correction

Communications with U.-K. Yang in 2014

**Issue: Experimental data are shown for "isoscalar iron".**

$$\text{CC: } \frac{d\sigma_{v,\bar{v}}^A}{dx dy} (\text{isoscalar}, v_f = 0) \equiv \frac{\frac{d\sigma_{v,\bar{v}}^A}{dx dy} (\text{Simulation with } v_f = 0)}{\frac{d\sigma_{v,\bar{v}}^A}{dx dy} (\text{Simulation with } v_f = (N - Z)/A)} \frac{d\sigma_{v,\bar{v}}^A}{dx dy} (\text{real data})$$

$$\frac{d\sigma_{v,\bar{v}}^A}{dx dy} = \frac{G_F^2 ME}{\pi} \left[ \left\{ y^2 + \left( 1 - y - \frac{M x y}{2 E} \right) \frac{1+R}{1+(2Mx)^2/Q^2} 2 \right\} 2x F_1^A \pm 2x y \left( 1 - \frac{y}{2} \right) F_3^A \right]$$

$$2x F_1^{vA} = [(1 - v_f) x d_v + (1 + v_f) x u_v + x S]_A$$

$$2x F_1^{\bar{v}A} = [(1 - v_f) x u_v + (1 + v_f) x d_v + x S]_A$$

$$x F_3^{vA} = [(1 - v_f) x d_v + (1 + v_f) x u_v + v_f x (\bar{u} - \bar{d}) + x (s - \bar{s})]_A$$

$$x F_3^{\bar{v}A} = [(1 - v_f) x u_v + (1 + v_f) x d_v - v_f x (\bar{u} - \bar{d}) + x (c - \bar{s})]_A$$

For example, for  $vA$  reaction

$$\begin{aligned} \frac{d\sigma_v^A}{dx dy} &= \frac{G_F^2 ME}{2\pi} \left[ \left\{ y^2 + \left( 1 - y - \frac{M x y}{2 E} \right) \frac{1+R}{1+(2Mx)^2/Q^2} 2 \right\} [(1 - v_f) x d_v + (1 + v_f) x u_v + x S]_A \right. \\ &\quad \left. + 2 y \left( 1 - \frac{y}{2} \right) [(1 - v_f) x u_v + (1 + v_f) x d_v - v_f x (\bar{u} - \bar{d}) + x (c - \bar{s})]_A \right] \\ &\quad \left\{ y^2 + \left( 1 - y - \frac{M x y}{2 E} \right) \frac{1+R}{1+(2Mx)^2/Q^2} 2 \right\} [x d_v + x u_v + x S]_{A'} \end{aligned}$$

$$\begin{aligned} \frac{d\sigma_v^{A'}}{dx dy} (v_f = 0) &= +2 y \left( 1 - \frac{y}{2} \right) [x u_v + x d_v + x (c - \bar{s})]_{A'} \\ \frac{d\sigma_v^A}{dx dy} (v_f \neq 0) &= \frac{\left\{ y^2 + \left( 1 - y - \frac{M x y}{2 E} \right) \frac{1+R}{1+(2Mx)^2/Q^2} 2 \right\} [(1 - v_f) x d_v + (1 + v_f) x u_v + x S]_A}{+2 y \left( 1 - \frac{y}{2} \right) [(1 - v_f) x u_v + (1 + v_f) x d_v - v_f x (\bar{u} - \bar{d}) + x (c - \bar{s})]_A} \end{aligned}$$

# Request to experimentalists

$$\frac{d\sigma_{v,\bar{v}}^A}{dx dy} \text{(isoscalar, } v_f = 0) \equiv \frac{\frac{d\sigma_{v,\bar{v}}^A}{dx dy} \text{(Simulation with } v_f = 0)}{\frac{d\sigma_{v,\bar{v}}^A}{dx dy} \text{(Simulation with } v_f = (N - Z) / A)} \frac{d\sigma_{v,\bar{v}}^A}{dx dy} \text{(real data)}$$

## Request to experimentalists:

- It is fine to publish  $\frac{d\sigma_{v,\bar{v}}^A}{dx dy}$ (isoscalar,  $v_f = 0$ ).
- However, please provide also  $\frac{d\sigma_{v,\bar{v}}^A}{dx dy}$ (real data)

because we theorists cannot handle  $\frac{d\sigma_{v,\bar{v}}^A}{dx dy}$ (isoscalar,  $v_f = 0$ ) to be precise.

If one is not worried about details of nuclear corrections, one can use the isoscalar data.

# Nuclear antiquark distributions and Drell-Yan

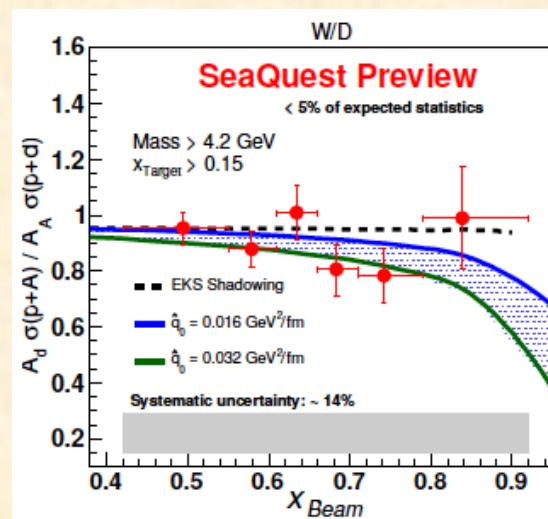
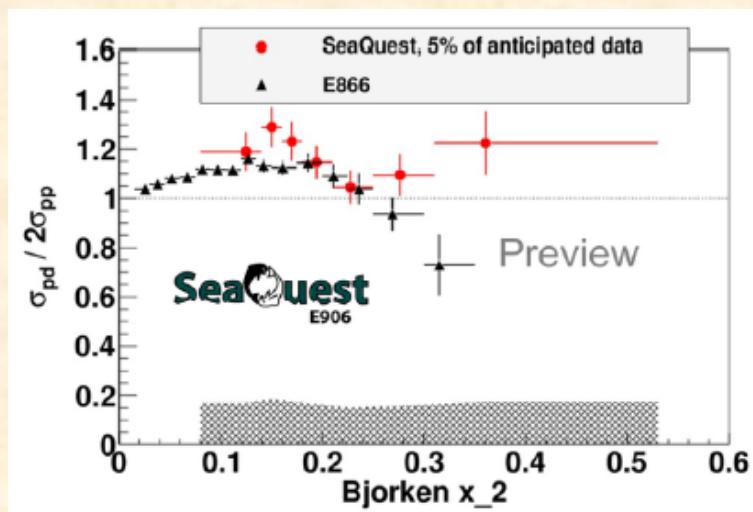
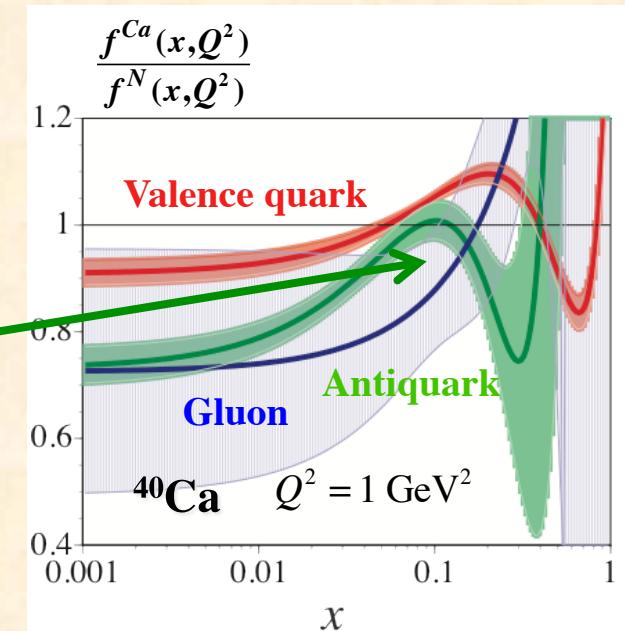
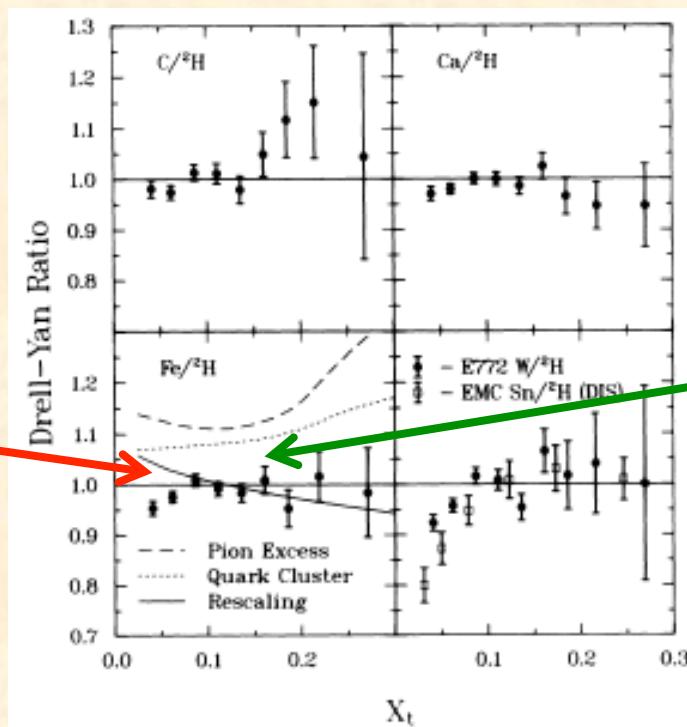
D. M. Alde *et al.*,  
PRL. 64, 2479 (1990).

$$\frac{\sigma_{pA}}{\sigma_{pD}} \approx \frac{\bar{q}_A}{\bar{q}_D}$$

No nuclear effects  
from pion contributions

E. L. Berger, F. Coester, R. B. Wiringa,  
PRD 29, 398 (1984)

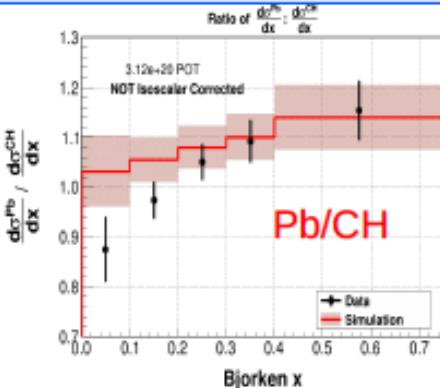
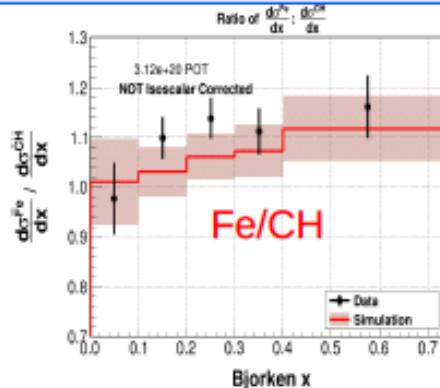
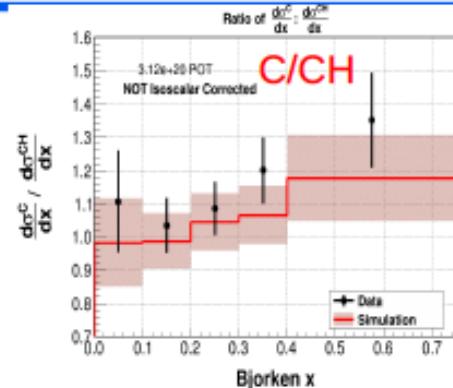
Fermilab-E906 in progress!



K. Nakano  
@Pacific-Spin2015

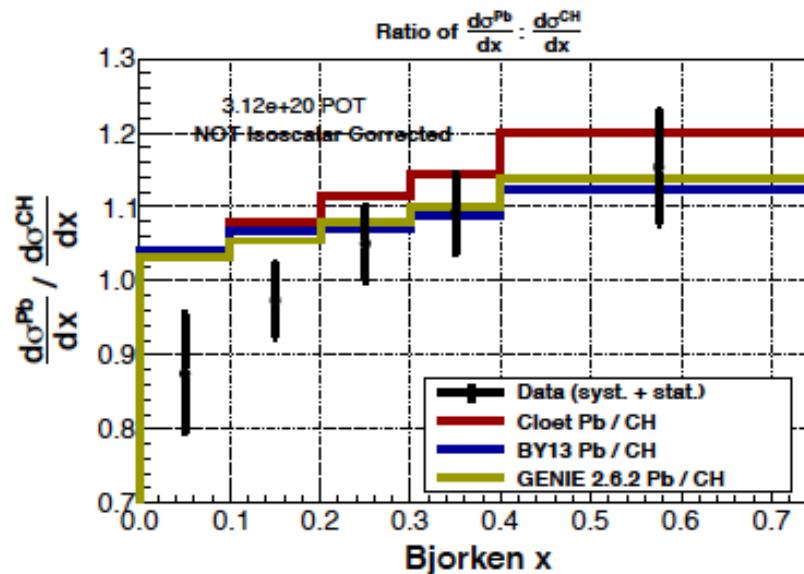
# We Now Have A New DIS Player - What does MINERvA see?

## DIS Cross Section Ratios – $d\sigma/dx$



J. Mousseau

- The shape of the data at low  $x$ , especially with lead is consistent with additional nuclear shadowing. **at an  $\langle x \rangle$  (0.07) &  $\langle Q^2 \rangle$  (2 GeV $^2$ ) - where negligible shadowing is expected with  $l^\pm$ .**



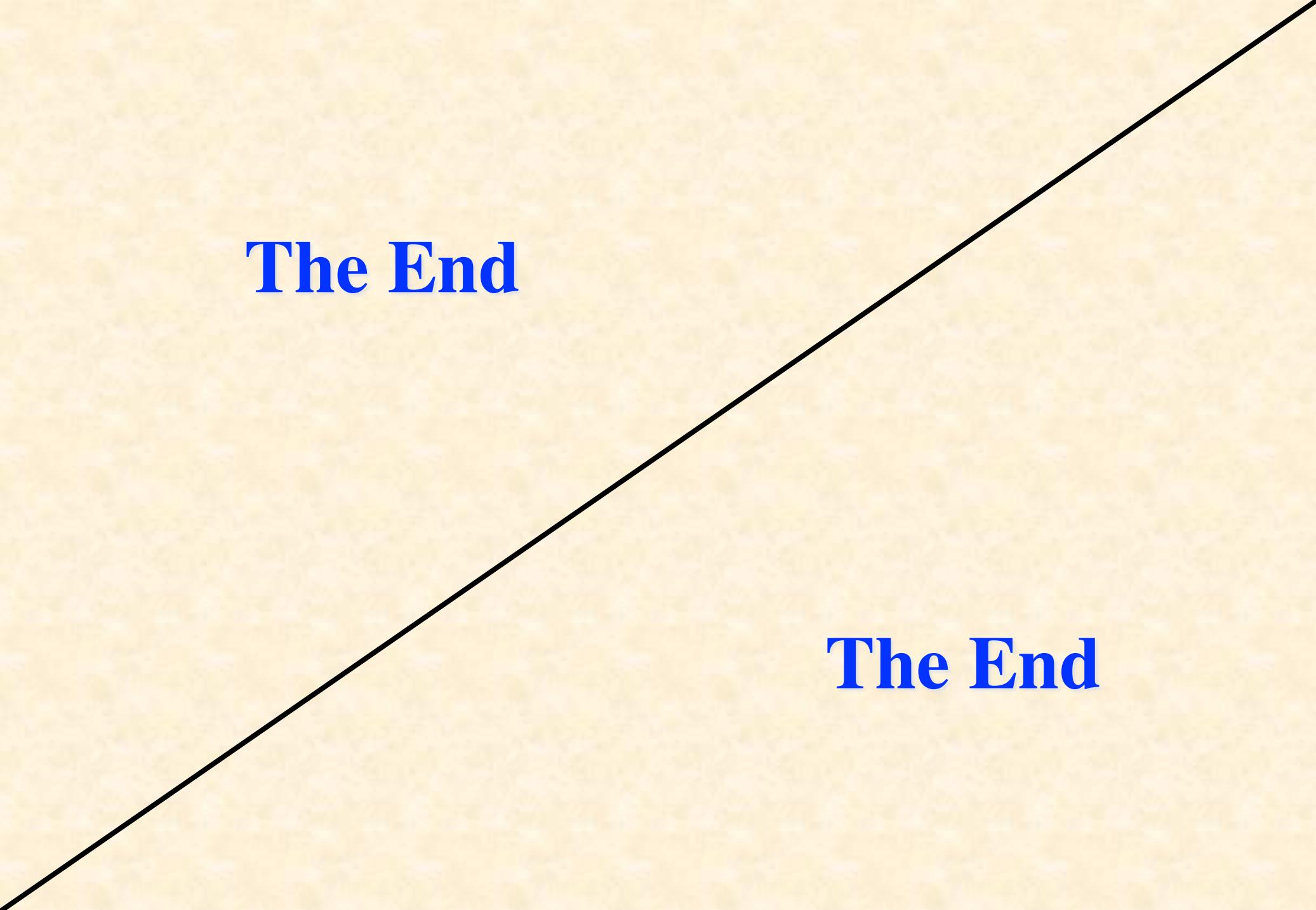
## Shadowing - continued

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- ◆ Why low  $x$ ?
- ◆ The lifetime of the hadronic fluctuation has to be sufficient to allow for these multiple diffractive scatters:

$$t_c = 2E_{\text{had}} / (Q^2 + m^2)$$

- ◆ For a given  $Q^2$  need large  $E_{\text{had}}$  to yield sufficient  $t_c$  which implies small  $x$ .
- ◆  $m$  is larger for the vector current than the axial vector current → for a given  $Q^2$  you need more  $E_{\text{had}}$  for the vector current than the axial vector current to have sufficient  $t_c$ .
- ◆ This implies you can have shadowing at higher  $x$  with neutrinos than with charged leptons



**The End**

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