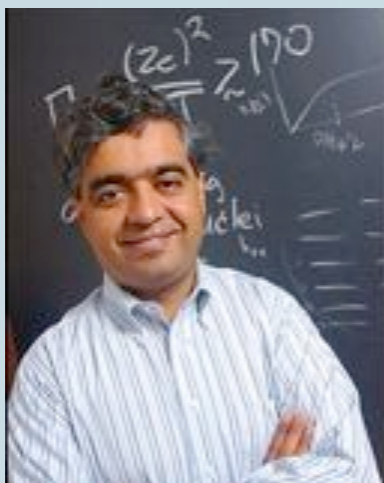


Dynamical & entanglement properties of CFTs [2+1D]

William Witczak-Krempa
Perimeter Institute, Canada

AdS/CFT & CM workshop, IPMU, 29/05



S. Sachdev
@Harvard / PI



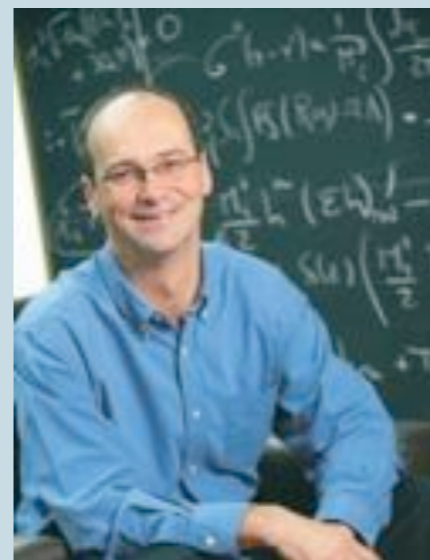
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Punch lines

focus:
CFTs

- ❖ Super-universality of **entanglement entropy** from corners
- ❖ **Non-perturbative** results for dynamics at $T > 0$
- ❖ **Asymptotics** and **sum rules** of observables (conductivity)
- ❖ Concrete statements for the **superfluid insulator** quantum critical point in 2+1D
[incl. large-scale Monte Carlo]

Plan

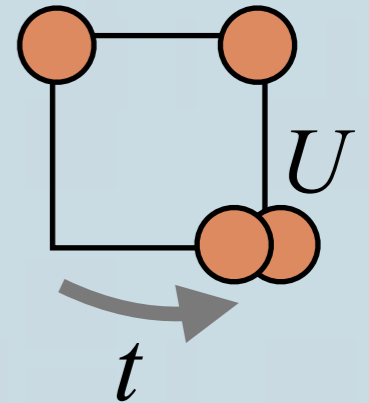
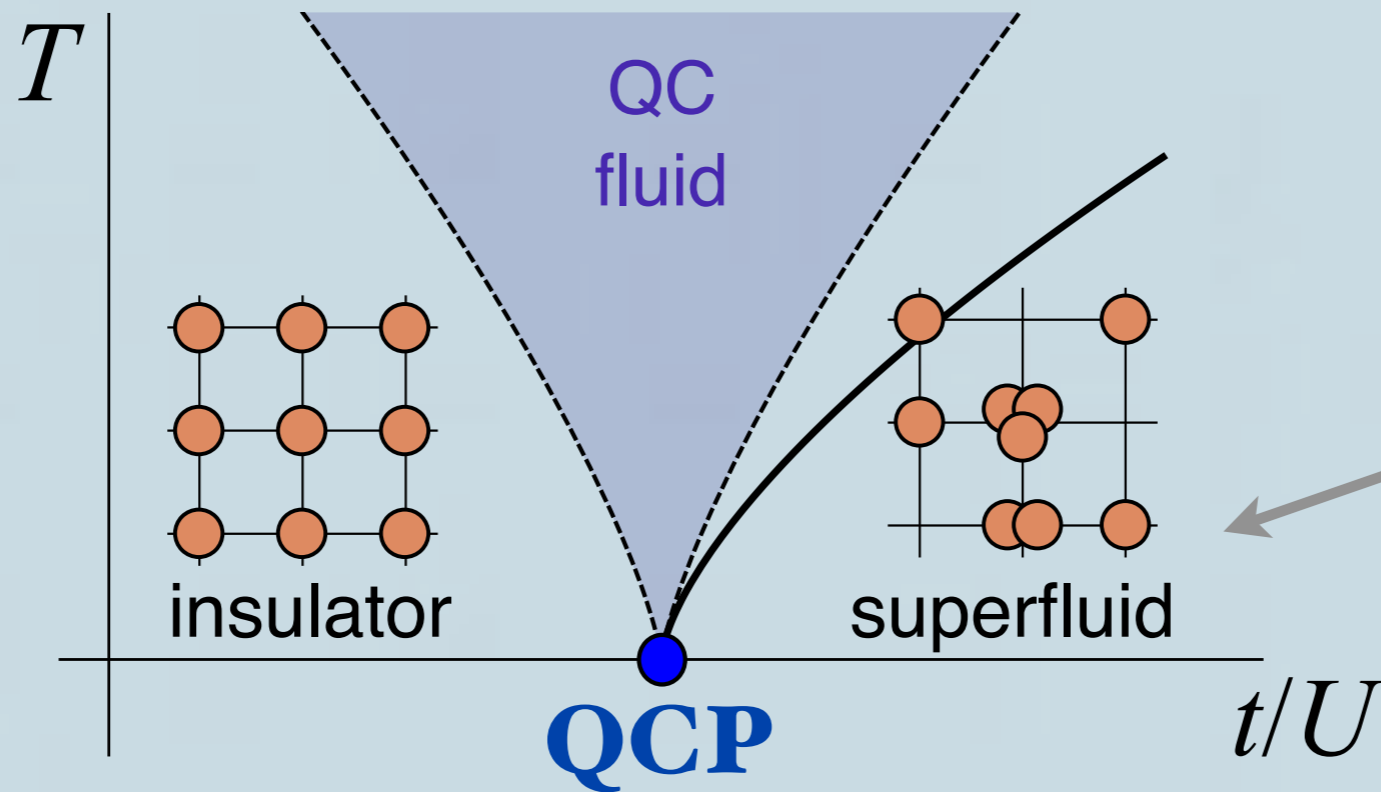
- ❖ Brief intro to quantum criticality
- ❖ Entanglement entropy (corner contribution)
- ❖ Quantum dynamics at $T > 0$

Quantum critical fluids

- (quasi)particles: ∞ lifetime (poles)
- Systems *without* quasiparticles:
 - ▶ Quantum critical phase transition
 - ▶ Strongly coupled gauge theories (e.g. gapless spin liquids), etc
- Focus on **Conformal Field Theories** ($z = 1$)

Superfluid-insulator transition

$$H = -t \sum_{\langle i,j \rangle} b_i^\dagger b_j + U \sum_i n_i (n_i - 1)$$



order param.
 ρ_s

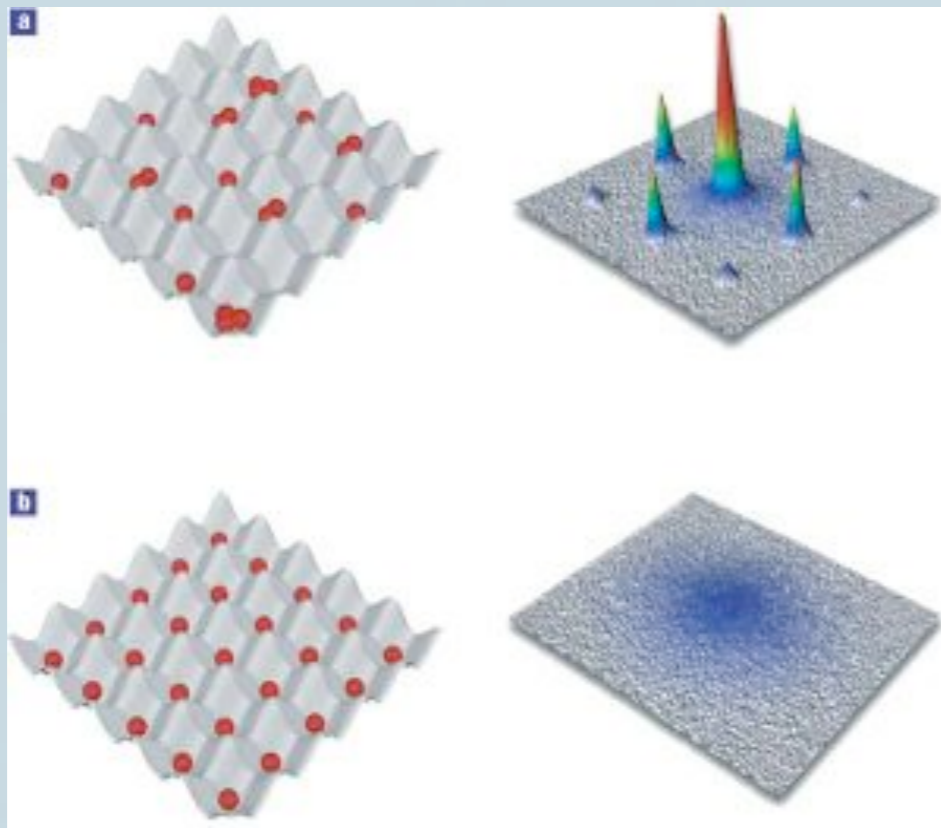
- ❖ **O(2) universality class:** XY spins, arrays of Josephson junctions, cold atoms, etc

❖ SSB of O(2) order parameter

$$\mathcal{L} = (\partial_t \vec{\phi})^2 + (\nabla \vec{\phi})^2 + m^2 \phi^2 + u \phi^4$$

strongly coupled fixed point
in 2+1D (Wilson-Fisher)

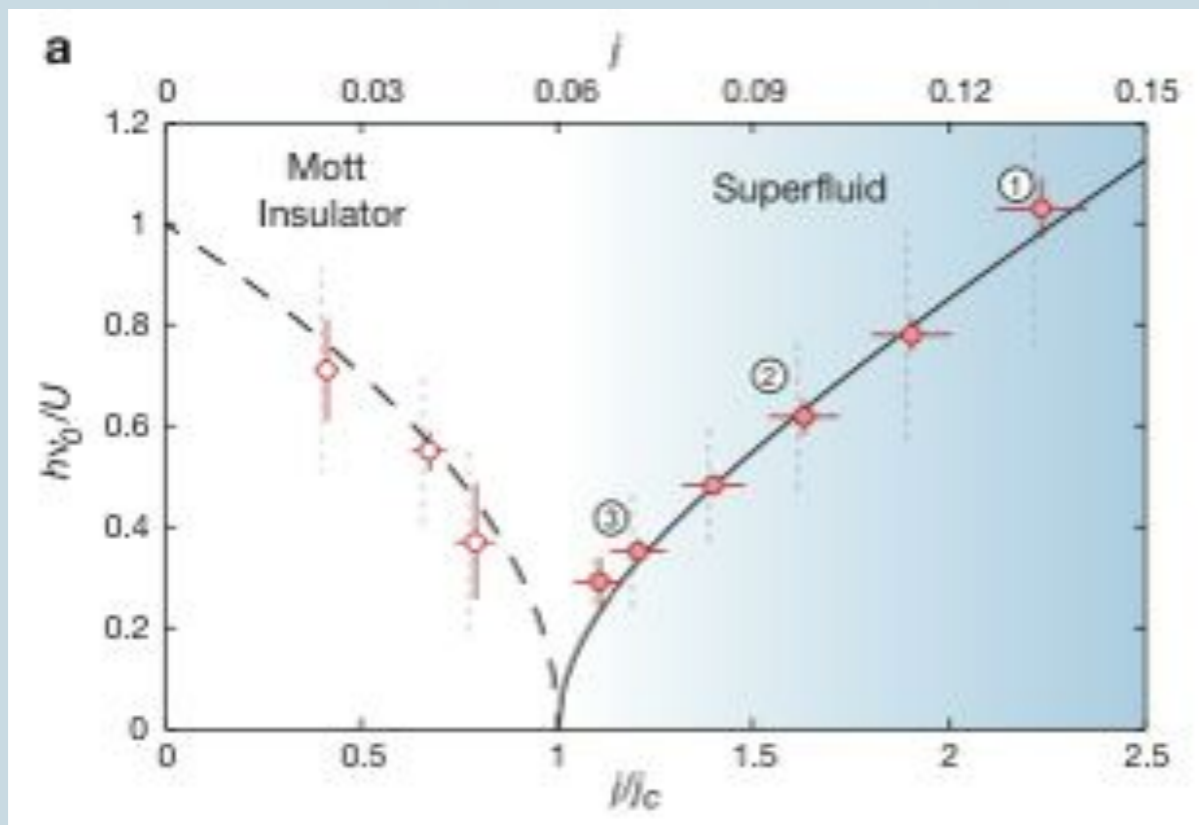




Superfluid

[Bloch]

Insulator



^{87}Rb

[Endres *et al*]

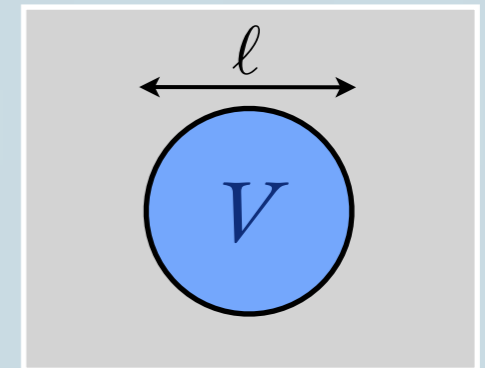
Entanglement

$$T = 0$$

P. Bueno, R. Myers, WK, 1505.04804

Counting non-quasiparticles

$$S = -\text{Tr}(\rho_V \ln \rho_V)$$



❖ Entanglement entropy

❖ RG monotone for CFTs (analog of c of 2D CFTs):
[Myers *et al* ; Casini, Huerta]

$$S = B \frac{\ell}{\delta} - F + O(\delta/\ell)$$

❖ Hard to compute in many-body simulations,
e.g. $O(2)$ QCP

CORNER entanglement

❖ V has corner w/ angle θ

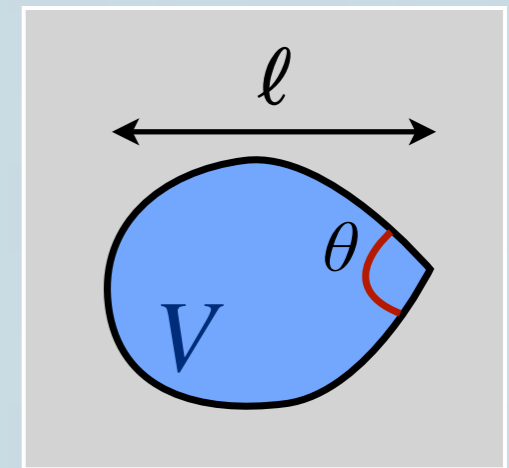
❖ EE has **universal log**:

$$S = B \frac{\ell}{\delta} - a(\theta) \ln(\ell/\delta) + \text{const}$$

❖ $a(\theta)$ good measure of dof

- free CFTs: $a = N a_{\text{scalar}}(\theta)$ [Casini, Huerta]
- large- N super-conformal gauge th.: $a(\theta) \propto N^{3/2}$
[Hirata, Takayanagi]

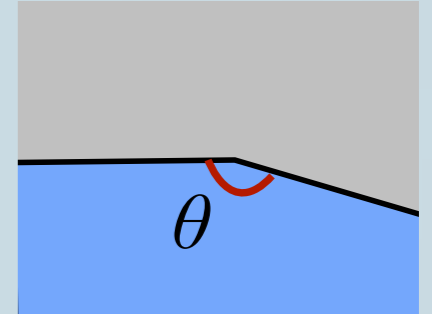
❖ Numerics for $O(N)$ Wilson-Fisher QCPs
[Melko *et al*, Helmes *et al*; Devakul *et al*; etc]



Corner super-universality

- ❖ **Smooth** limit, $\theta \simeq \pi$

$$a(\theta \rightarrow \pi) = \sigma (\theta - \pi)^2$$



- ❖ Super-universality

$$\sigma = \frac{\pi^2}{24} C_T$$

$$\langle T_{\mu\nu}(x) T_{\lambda\rho}(0) \rangle = \frac{C_T}{x^6} \mathcal{I}_{\mu\nu\lambda\rho}(x)$$

- ❖ smooth limit determined by **local** correlator

Evidence for conjecture

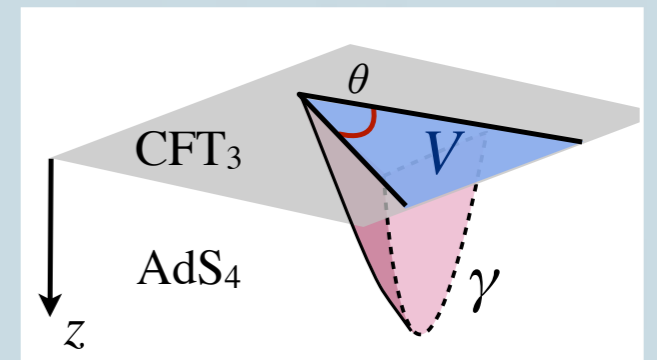
$$\sigma = \frac{\pi^2}{24} C_T$$

- ❖ Holds exactly for free theories
→ predict new result (numerically checked)

$$\sigma_{\text{scalar}} = \frac{1}{256} = \sigma_{\text{Dirac}}/2$$

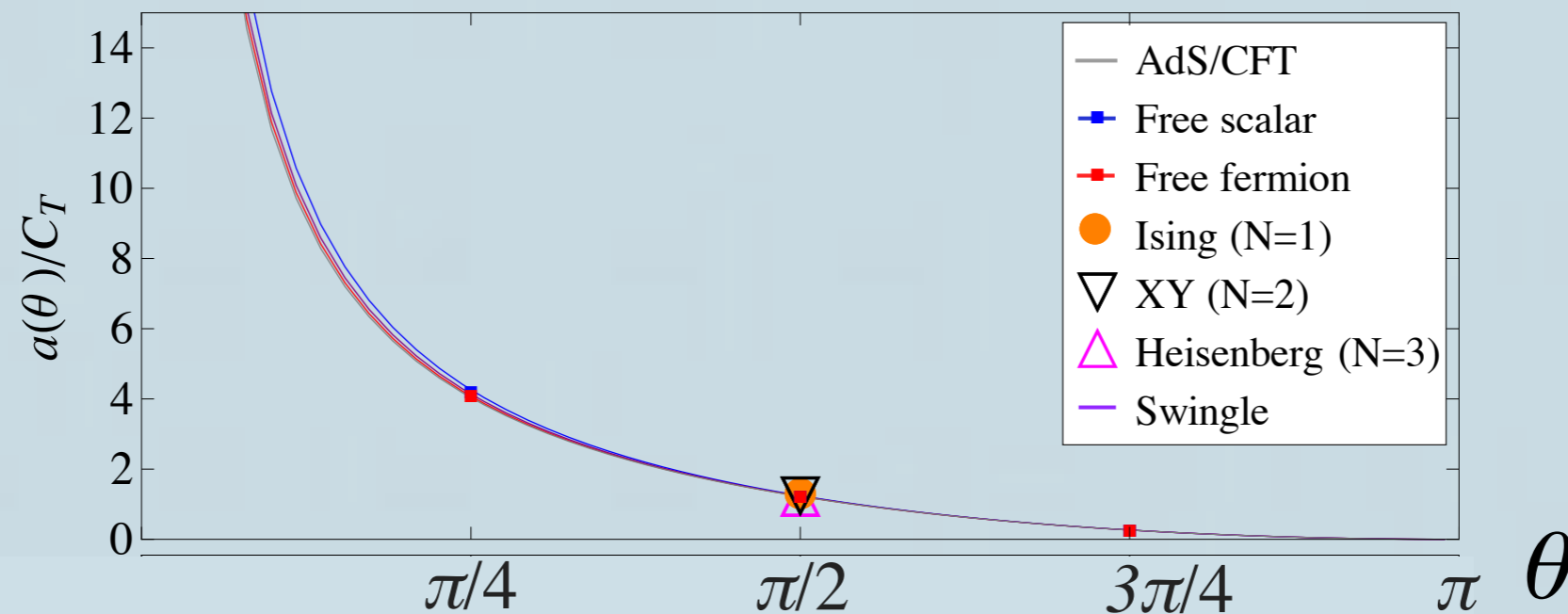
- ❖ Holds in holographic models with higher curvature terms [*à la* Ryu-Takayanagi]

$$S_{\text{bulk}} = \int_x \left[\frac{6}{L^2} + R + \lambda_1 R \mathcal{X}_4 + \lambda_2 \mathcal{X}_4^2 \right]$$



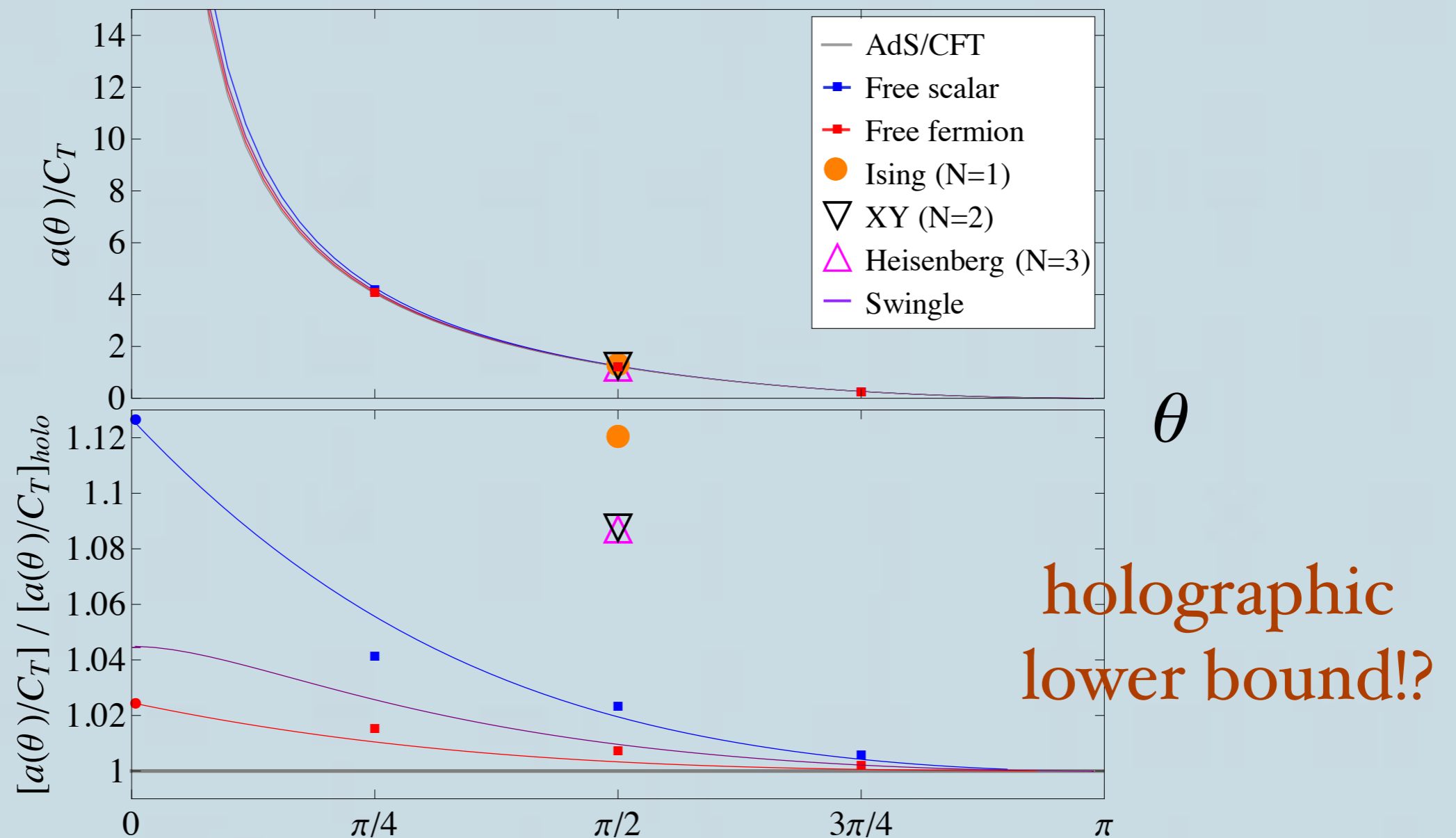
Beyond smooth limit

- ❖ Near super-universal behavior for many different CFTs



Beyond smooth limit

- ❖ Near super-universal behavior for many different CFTs



Quantum dynamics at $T > 0$

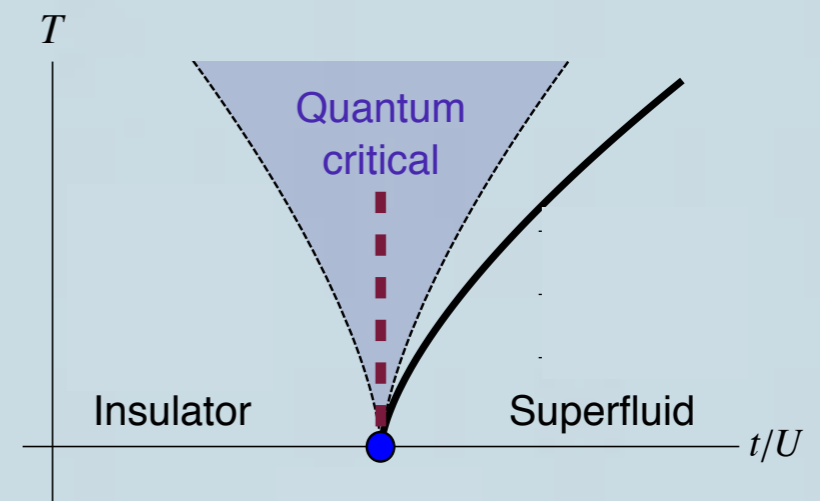
Universal conductivity

$$\partial_\mu J^\mu = 0$$

$$\sigma(\omega) = \frac{1}{i\omega} \langle J_x(\omega) J_x(-\omega) \rangle$$

Universal scaling function:

$$\sigma(\omega) = \Phi\left(\frac{\omega}{T}\right)$$



[early work: Damle, Sachdev]

Short & long times

- ❖ Characteristic time scale: $1 / T$
- ❖ **Short** times $\omega \gg T$: probing near vacuum
- ❖ **Long** times $\omega \ll T$: excitations interact strongly with thermal background

Short times via Operator Product Expansion

- ❖ A scalar primary op $O(x)$, w/ scaling dim Δ :

$$\langle O(x) O(0) \rangle = 1 / x^{2\Delta}$$

- ❖ OPE:


The diagram illustrates the Operator Product Expansion (OPE) for two scalar primary operators. On the left, two operators are shown: $\mathcal{O}(x)$ with a black dot and $\mathcal{O}(0)$ with a blue dot. A red arrow points to the right, where the expansion is given as a sum over n of $\frac{\mathcal{O}_n(0) + \text{desc.}}{x^{2\Delta - \Delta_n}}$. A blue dot is placed below the denominator to indicate the expansion point.


$$\mathcal{O}(x) \mathcal{O}(0) \rightarrow \sum_n \frac{\mathcal{O}_n(0) + \text{desc.}}{x^{2\Delta - \Delta_n}}$$

[Wilson ; Polyakov; Ferrara *et al* ; etc]

Current OPE

$$J_\mu(x) J_\nu(0) = \frac{I_{\mu\nu} 1}{x^{2\cdot 2}} + \mathcal{C}_{JJ\mathcal{O}} \frac{\tilde{I}_{\mu\nu} \mathcal{O}(0)}{x^{4-\Delta}} + \mathcal{C}_{JJT} \frac{T_{\mu\nu}(0)}{|x|} + \dots$$


O(N) CFT:
relevant
scalar $\mathcal{O} \sim \varphi^2$


Stress tensor

❖ Get OPE coefficients from $\langle JJ\mathcal{O} \rangle_{T=0}$

Asymptotic conductivity

[WK, Sorensen, Sachdev, Katz]

❖ Fourier & take expectation value of OPE:

thermal average $\langle \mathcal{O} \rangle_T = BT^\Delta$

$$\sigma(i\omega_n) \stackrel{\omega_n \gg T}{=} \sigma_\infty + b_O \left(\frac{T}{\omega_n} \right)^\Delta + b_T \left(\frac{T}{\omega_n} \right)^3 + \dots$$

↑
relevant
scalar \mathcal{O}

↑
stress tensor

Dominant operator?

- ❖ $O(N)$ Wilson-Fisher, it's the **scalar** O

$$O(x) = \varphi^2(x) \quad \text{since} \quad \Delta = 3 - 1/\nu$$

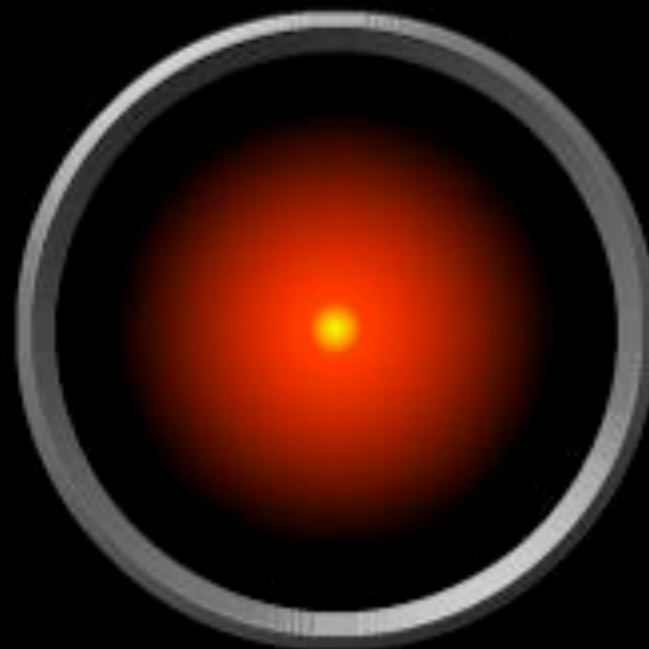
- ❖ $N = \infty$: $\Delta = 2$

- $N = 2$: $\Delta = 1.51$

- ❖ All CFTs have $1/\omega^3$ term because *there's always a stress tensor*
 - sometimes dominates (Dirac CFT, QED3, etc)

Ask the computer

I'm sorry Dave,
I'm afraid I can't do that.



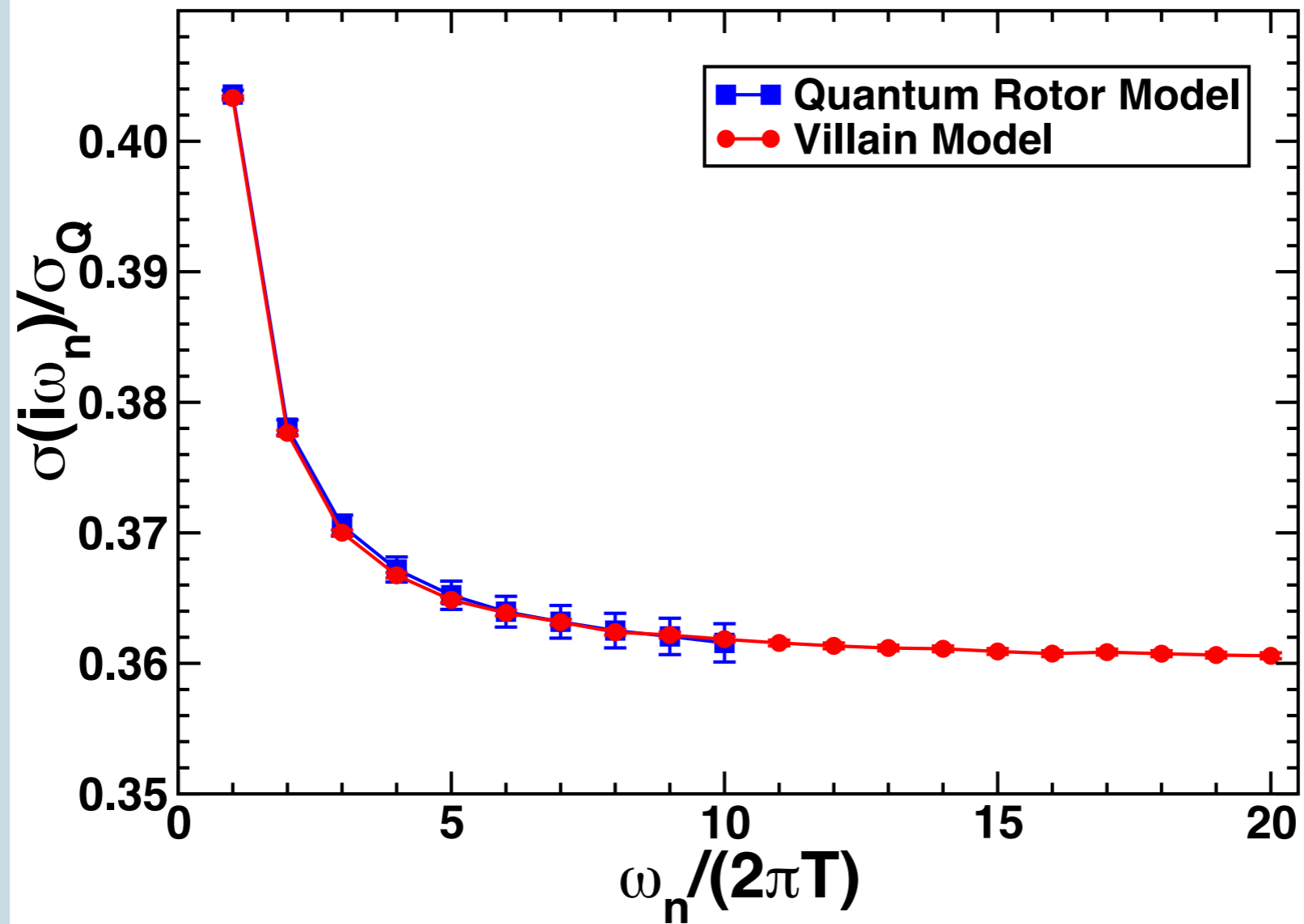
Quantum Monte Carlo

[WK, **Sorensen**, Sachdev, Nat. Phys. 2014]

- ❖ Large-scale simulations of $O(2)$ QCP
 - ❖ loop-current model (Villain)
 - ❖ quantum rotors
- ❖ Finite- T but **imaginary time...**
- ❖ Analytic continuation difficult!

1st case of universal quantum critical dynamics!

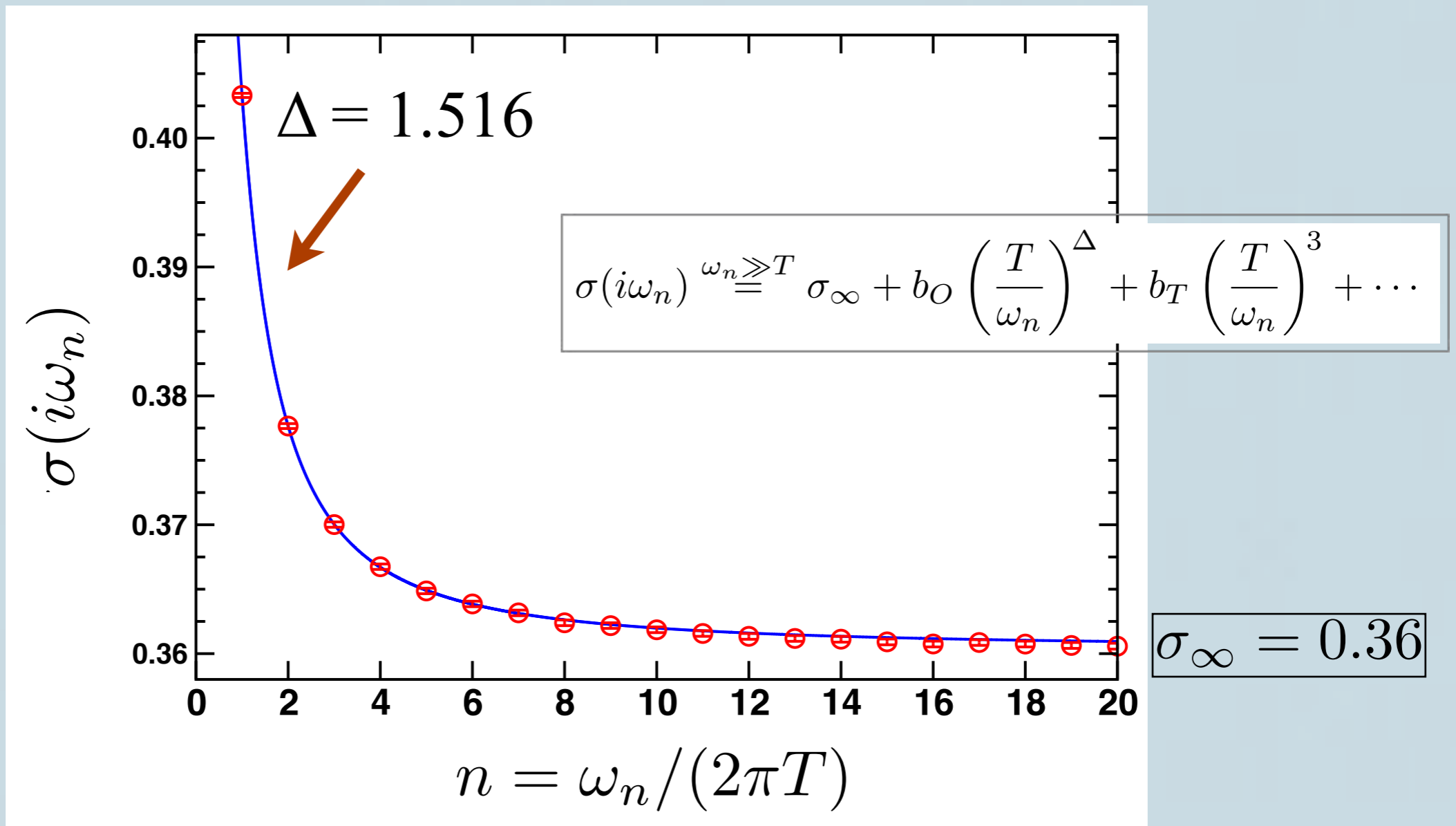
[WK, **Sorensen**, Sachdev, Nat. Phys. 2014]



[also Chen *et al*]

Quantum Monte Carlo II

[WK, **Sorensen**, Sachdev, Katz]



Presence of Quantum Diffusion in Two Dimensions: Universal Resistance at the Superconductor-Insulator Transition

Matthew P. A. Fisher and G. Grinstein

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S. M. Girvin

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$$\sigma_{\infty} = 0.36$$

- ❖ Determined after 25 years!
- ❖ Conformal bootstrap : $\sigma_{\infty} = 0.3554(8)$
[Kos *et al* 2015]

[also Chen *et al*;
Gazit *et al*]

Sum rules

Sum rules

$$\int_0^{\infty} d\omega \operatorname{Re} \sigma(\omega) = \frac{\pi n e^2}{2m}$$

independent of T ,
even of superconductivity!
[Kramers; Ferrell-Glover-Thinkham]

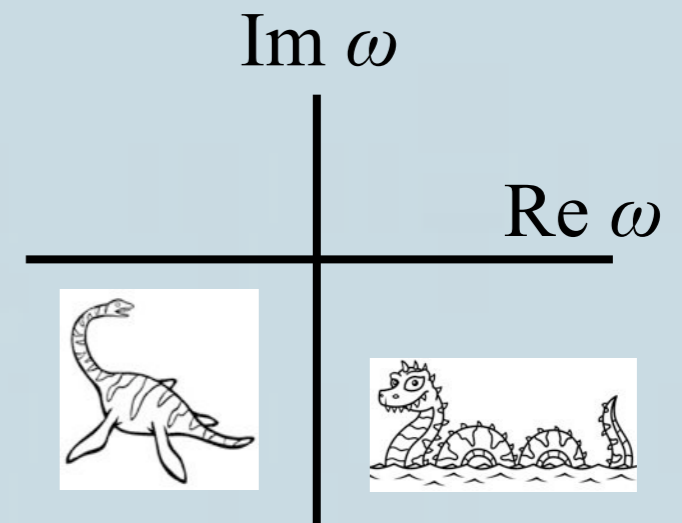
Robust, super-universal statements useful to:

- ❖ Assess approximations
- ❖ Constrain numerical results
- ❖ Help with the analysis of experimental data

Building sum rules

❖ 2 ingredients:

1. Kramers-Kronig (retarded)



2. **Asymptotics** via OPE

Sum rule w/out quasiparticles

[WK, Sachdev ; Gulotta *et al* ;
Katz, Sachdev, Sorensen, WK]

$$\int_0^{\infty} d\omega [\text{Re } \sigma(\omega) - \sigma_{\infty}] = 0$$

❖ General proof via OPE

$$\sigma(i\omega_n) \stackrel{\omega_n \gg T}{\cong} \sigma_{\infty} + b_O \left(\frac{T}{\omega_n}\right)^{\Delta} + b_T \left(\frac{T}{\omega_n}\right)^3 + \dots$$

❖ Check:

- ✓ Maximally SUSY in large- N
- ✓ $O(N)$ CFT, $N=\infty$
- ✓ Dirac CFT



*Long times &
analytic continuation*

★ Use **AdS/CFT** to generate a family of physically motivated (constrained) variational functions $\sigma(\omega/T)$

Why dabble w/ black holes?

- ❖ **Physical** properties:

- ◆ ***Tailored*** holographic description to match asymptotics

- ◆ **Sum rules**

$$\int_0^\infty d\omega \operatorname{Re}[\sigma(\omega/T) - \sigma(\infty)] = 0$$

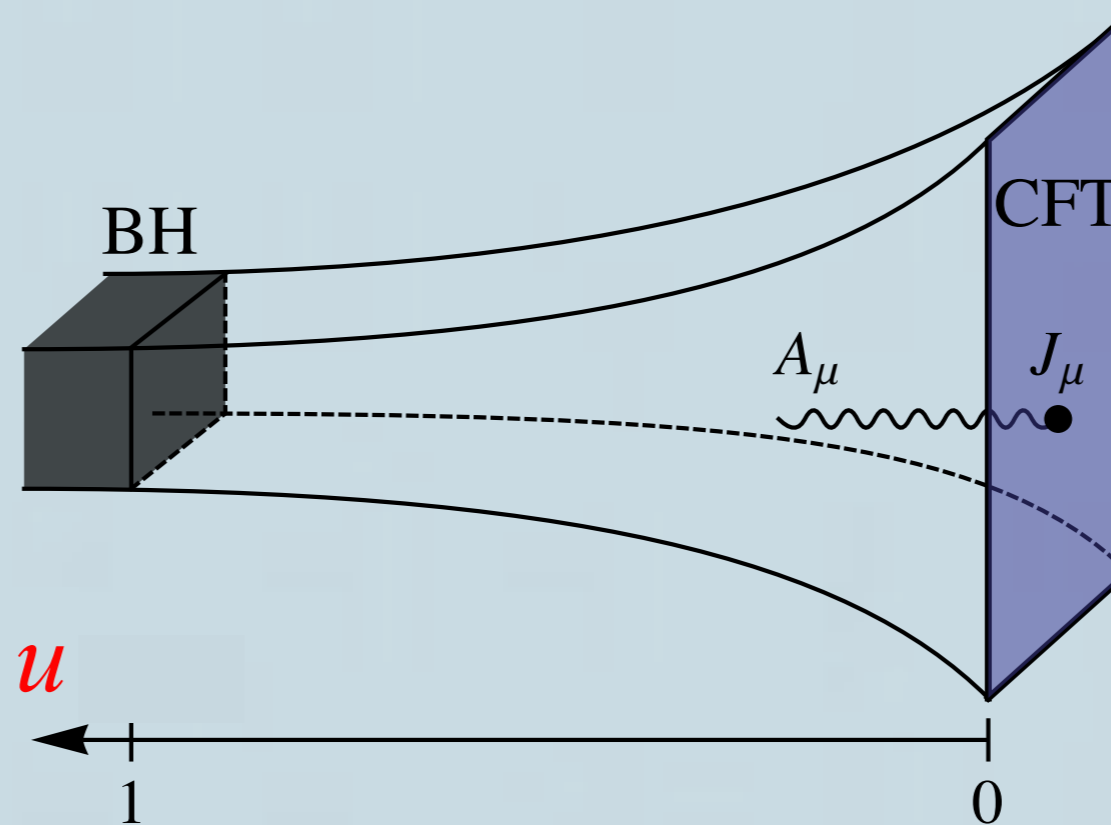
- ◆ **Analytic properties** in common with field theory results

- ◆ **Practical:** real-time, small number of params, fast numerics

σ via AdS/CFT

[Maldacena *etc*]

spacetime:
**B-Hole in
AdS₄**



$$A_\mu(t, x, y; u) \leftrightarrow J_\mu^{\text{CFT}}(t, x, y)$$

❖ Solve **classical** EoM for $A_\mu \rightarrow$ get J -correlator in **CFT**

$$\sigma(\omega/T) = \frac{T}{\omega} \left. \frac{\partial_u A_y(\omega, \vec{0}; u)}{A_y(\omega, \vec{0}; u)} \right|_{u=0}$$

Bulk action

$$S_{\text{bulk}}[A_\mu] = \int_x \frac{1}{g_4^2} [1 + \alpha \varphi(x)] F_{ab} F^{ab}$$

- ❖ Add **scalar field** (dilaton)

CFT	AdS
\mathcal{O}	φ

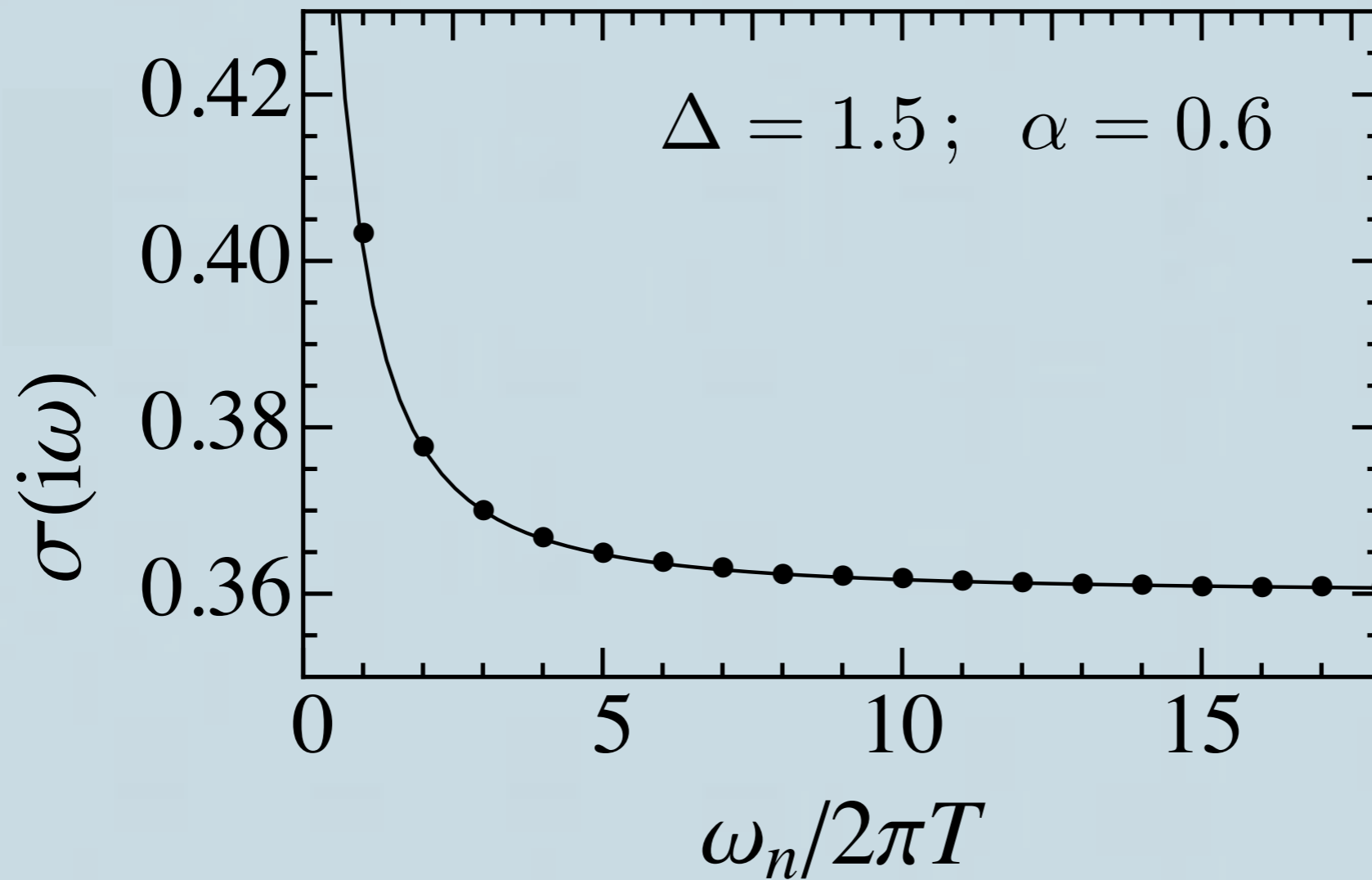
- ❖ Simplest **Ansatz**:
fix profile using OPE of $O(2)$ Wilson-Fisher CFT

$$\varphi(u) = u^\Delta$$

[Katz, Sachdev, Sorensen, WK]

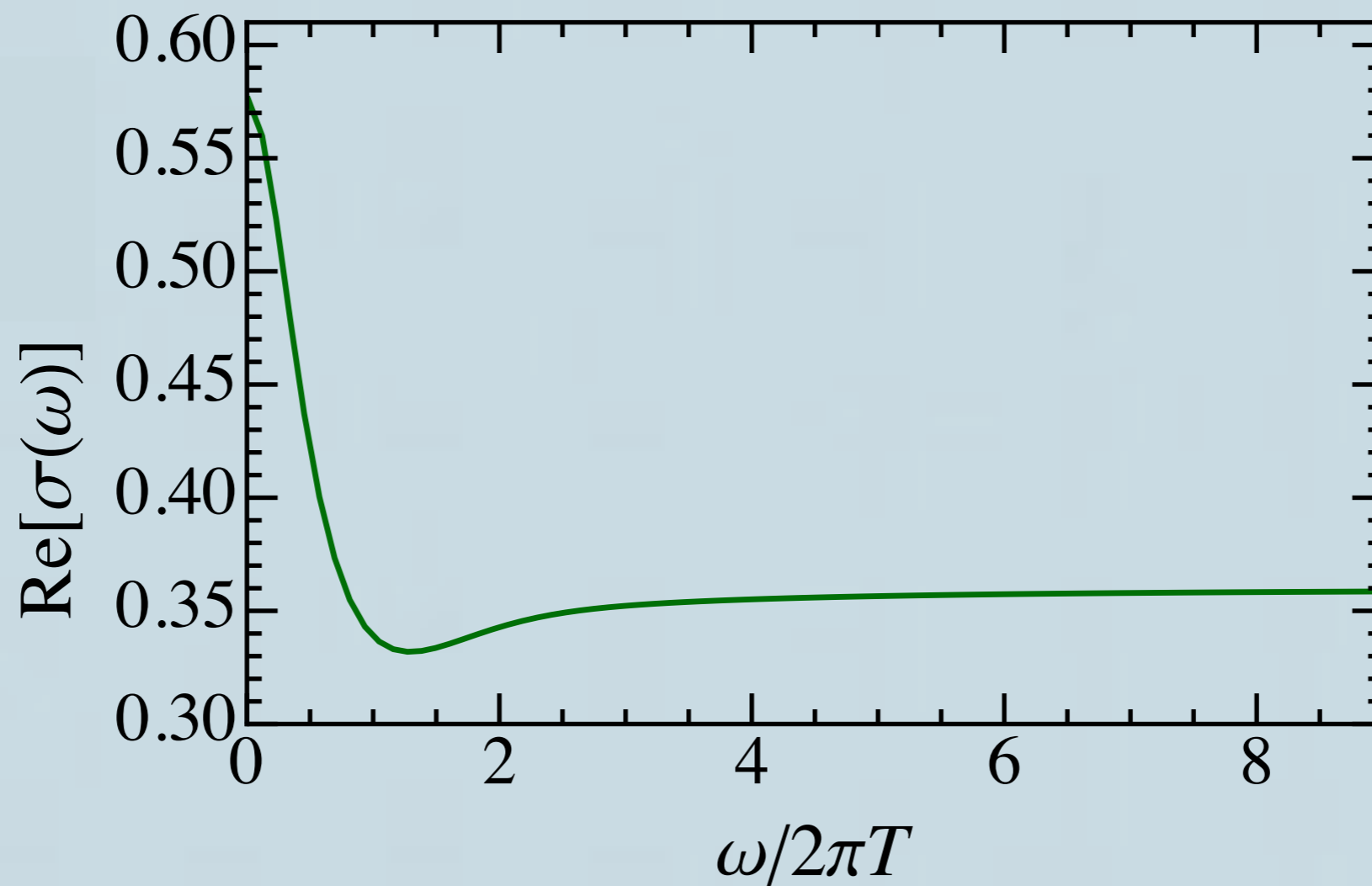
Holographic fit

[Katz, Sachdev, Sorensen, WK]



Real frequencies

[Katz, Sachdev, Sorensen, WK]



Conclusions

- ❖ **Corner** entanglement is a useful measure of dof
- ❖ In **smooth limit**, $a(\theta)$ fixed by $T_{\mu\nu}$ central charge

$$a(\theta \rightarrow \pi) = \sigma (\theta - \pi)^2$$

$$\sigma = \frac{\pi^2}{24} C_T$$

- ❖ Does pure Einstein holography give a **lower bound** for $a(\theta)/C_T$?

Conclusions

- ❖ Quantum critical dynamics at finite T , (CFTs in 2+1D)
- ❖ OPE to constrain *short time* physics

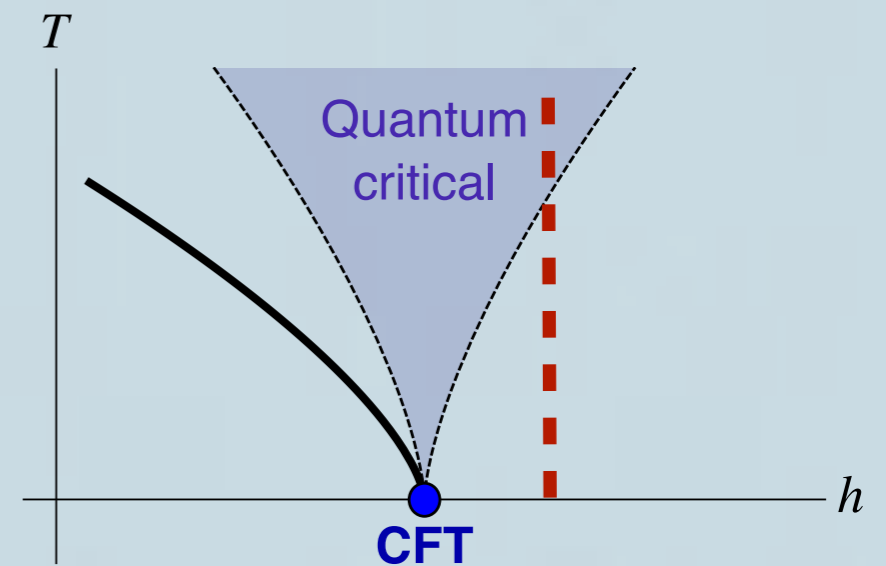
- ❖ large ω conductivity

$$\sigma(i\omega_n) \stackrel{\omega_n \gg T}{\approx} \sigma_\infty + b_1 \left(\frac{T}{\omega_n}\right)^\Delta + b_2 \left(\frac{T}{\omega_n}\right)^3 + \dots$$

- ❖ Input OPE data of CFT into simple holographic ansatz
- ❖ Use holography to analytically continue Monte Carlo data for conductivity at superfluid/insulator QCP

Outlook

- ❖ Apply this program (OPE, sum rules) to other correlators & CFTs [WK, Phys. Rev. Lett. 2015]
- ❖ Go beyond simplest holographic Ansatz [in prog. w/ Sierens & Myers]
- ❖ Finite detuning from QCP [in prog. w/ Gazit & Podolsky]



ありがとう

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