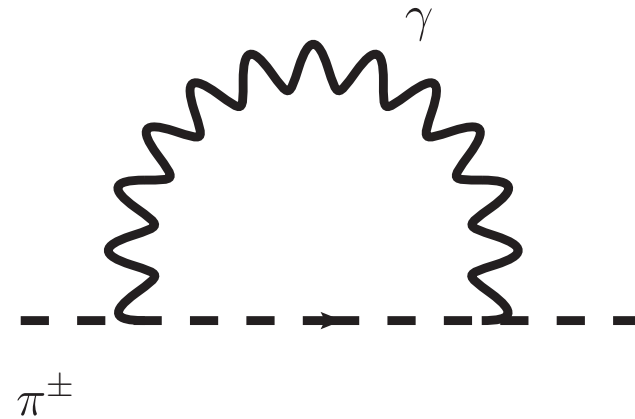
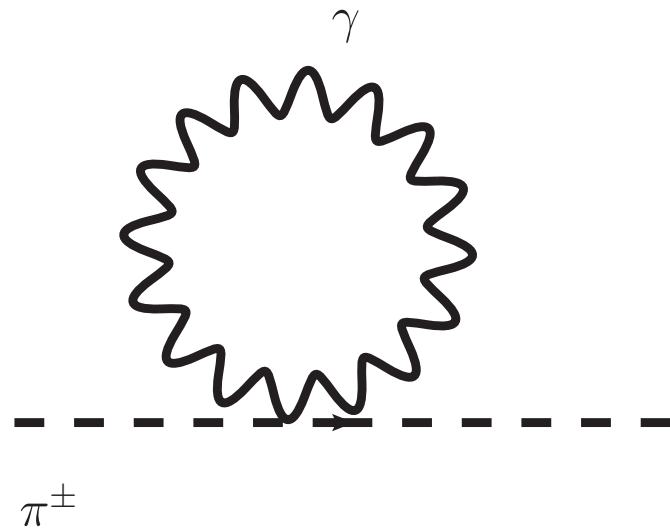


Composite Spin-1 resonance at the LHC (and beyond)

LianTao Wang
University of Chicago

IPMU Sept. 11, 2015

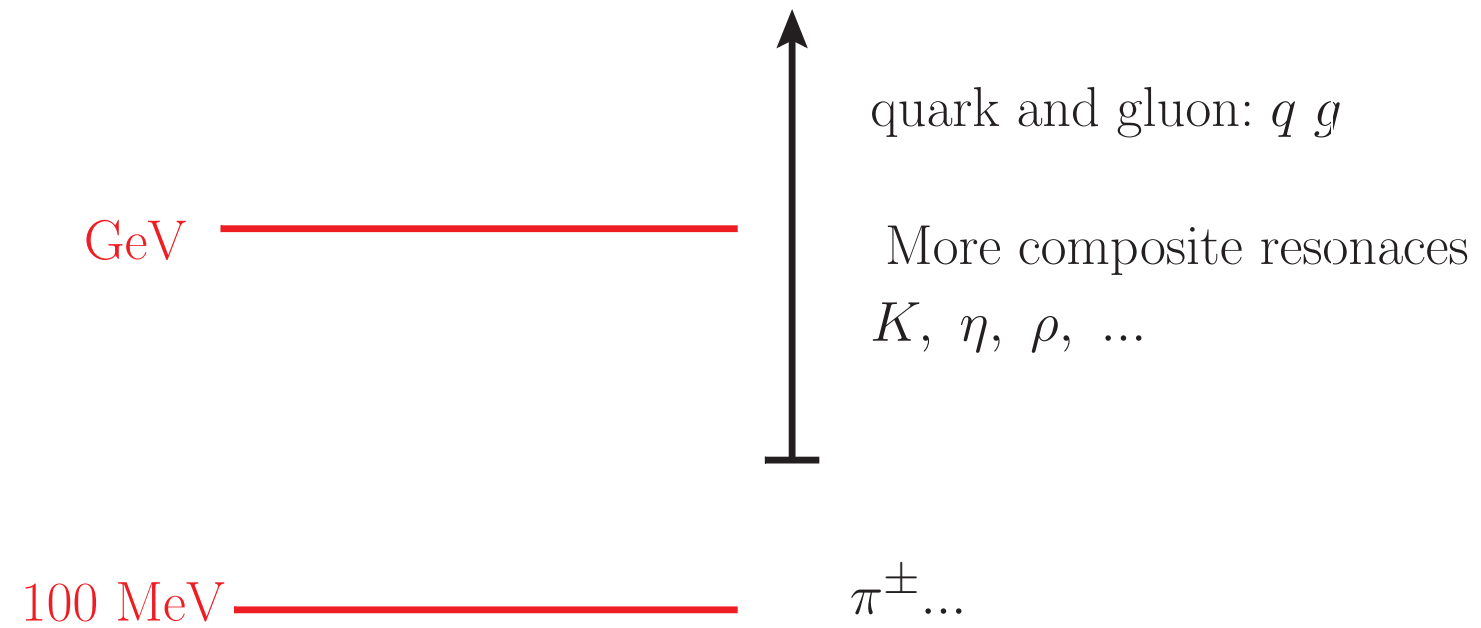
Naturalness in nature?



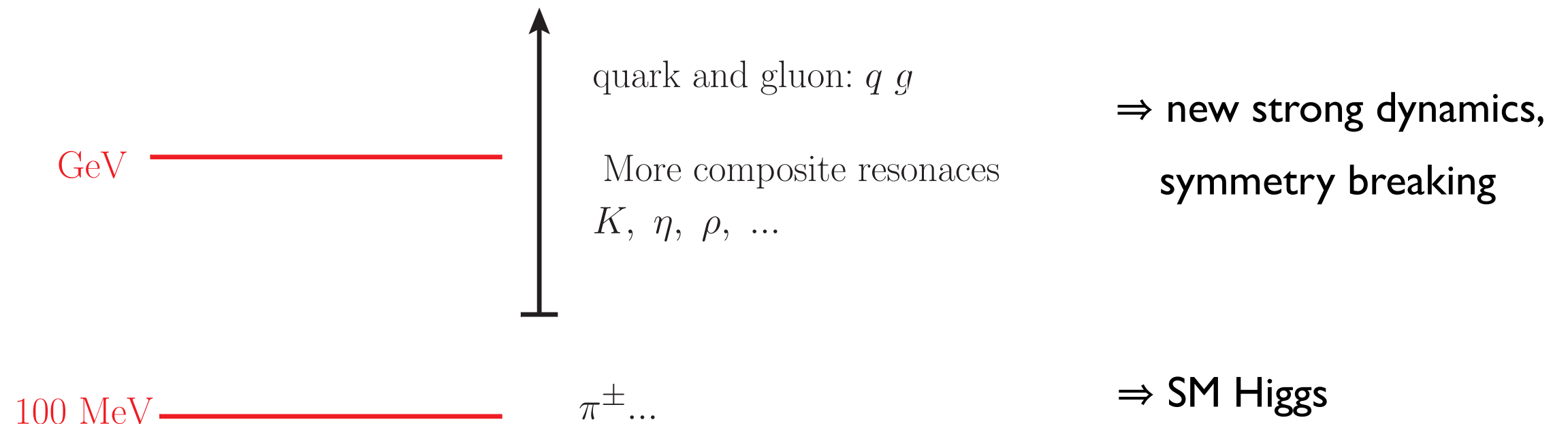
$$\delta m_{\pi^\pm}^2 \simeq \frac{e^2}{16\pi^2} \Lambda^2$$

- Example: low energy QCD resonances: pion
- $m_\pi \sim 100$ MeV.
- Naturalness requires $\Lambda \approx \text{GeV}$.
 - Indeed, at GeV, QCD $\Rightarrow$ theory of quark and gluon
 - Pion is not elementary.

“Learning” from QCD

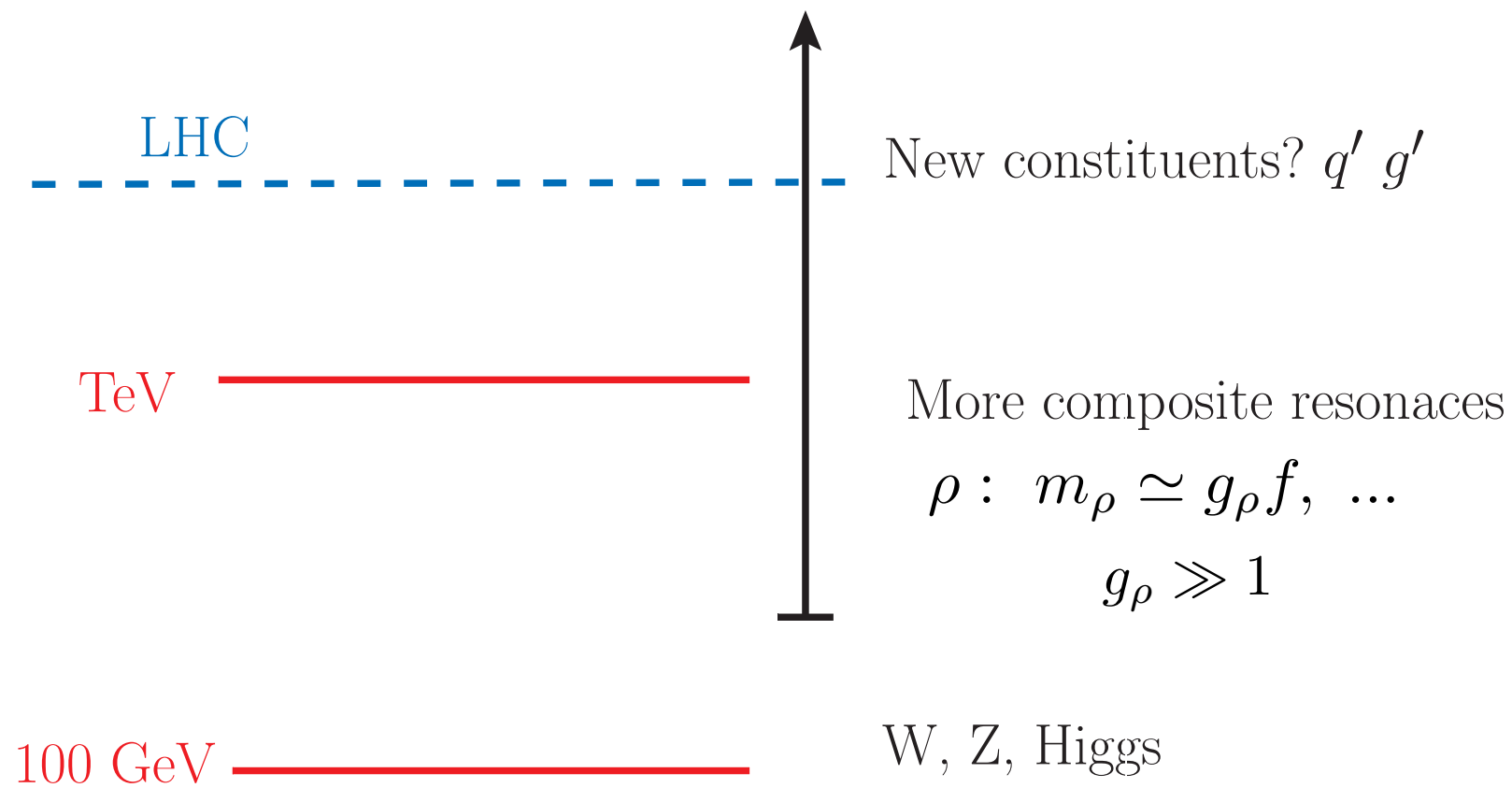


“Learning” from QCD



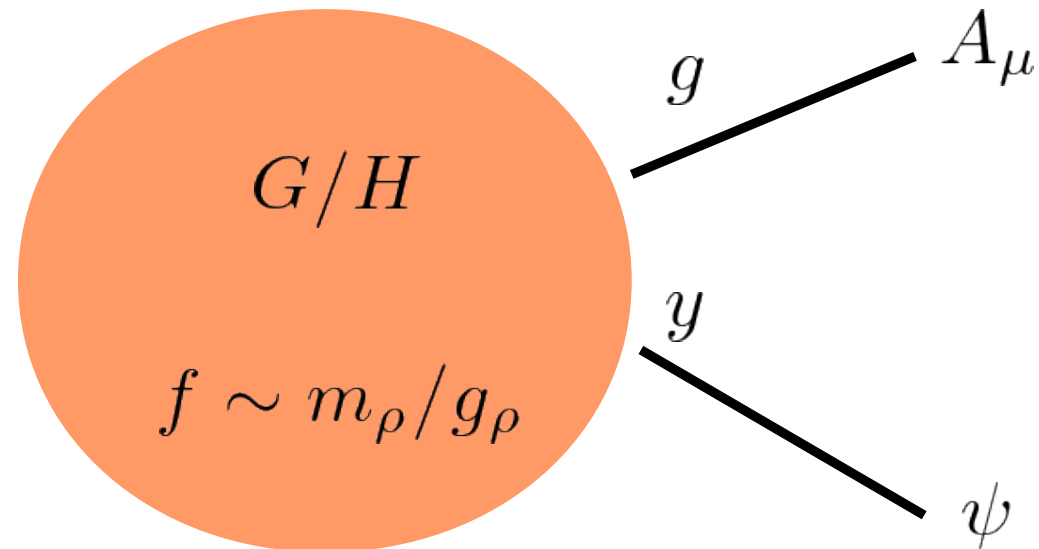
- Construct a new strong dynamics in which the low lying states will be the SM Higgs.
- Composite Higgs models. Still a natural theory.
- Nature may be more interesting, but it could also just repeat itself.

Composite Higgs



- Many models in this class.
- Similar scenarios: Randall–Sundrum...

Composite Higgs



Higgs (and W/Z goldstones) are part of the strong sector

The external fields are the SM quarks and (transverse) gauge bosons

- Higgs boson (and W_L Z_L) NGB of symmetry breaking G/H .
- Small explicit symmetry breaking (involving external fields) generates Higgs potential. (NGB $\rightarrow$ pNGB).

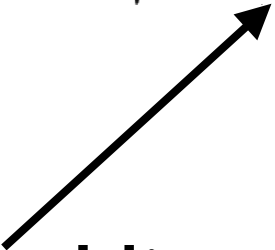
Minimal composite

Agashe, Contino, Pomarol

$$SO(5)/SO(4) \rightarrow 4 \text{ GBs}$$
$$U = \exp \left(i\sqrt{2} \frac{\pi^a}{f} T^a \right) \quad T^a = \begin{pmatrix} 0 & X_a \\ -X_a & 0 \end{pmatrix} \quad \Sigma_i = U_{i5}$$

$$\frac{1}{2}(D_\mu \Sigma)^2 \supset \frac{1}{2}(\partial h)^2 + \frac{1}{2}m_V^2 V_\mu^2 \left(1 + 2\sqrt{1 - \frac{v^2}{f^2} \frac{h}{v}} + \dots \right)$$

First prediction: deviation in Higgs coupling



Spin-1 resonances

$$G/H = SO(5)/SO(4)$$

Spin-1 resonance in 6 of $SO(4)$

$$m_\rho \sim g_\rho f$$

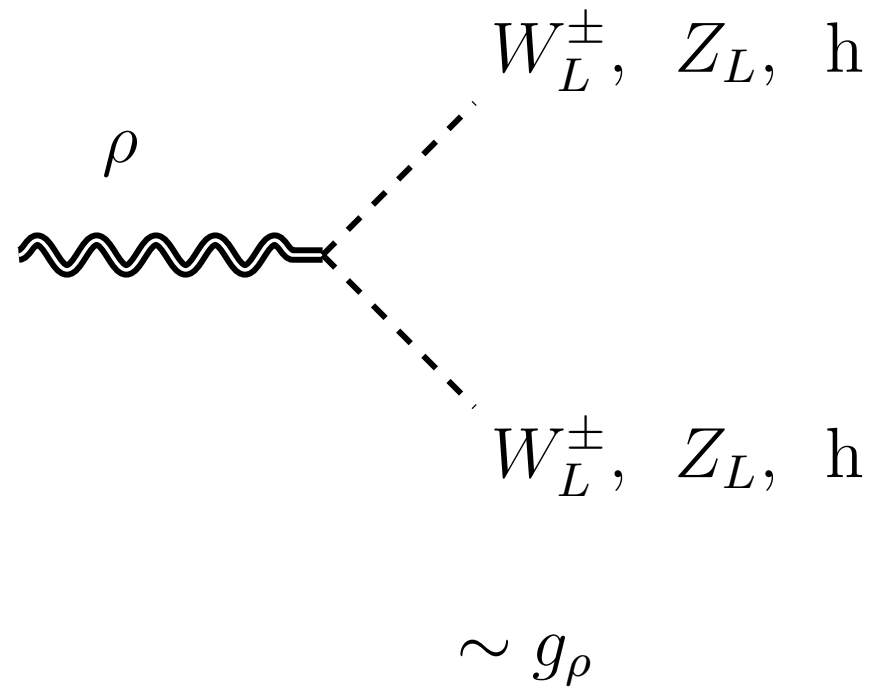
$$SO(4) \times U(1)_X = SU(2)_L \times SU(2)_R \times U(1)_X \quad Y = T_R^3 + X$$

Under this, spin-1 res. decompose as

$$\begin{array}{c} \mathbf{6} \rightarrow \mathbf{3}_0 + \mathbf{1}_0 + \mathbf{1}_\pm \\ \nearrow \quad \uparrow \quad \nwarrow \\ \rho_L : \rho_0, \rho^\pm \quad \rho_B \quad \rho_C^\pm \end{array}$$

ρ coupling to h, W_L, Z_L

$$\mathcal{L} \supset ig_\rho c_H \rho_\mu^a (H^\dagger \tau^a D^\mu H - (D^\mu H)^\dagger \tau^a H)$$



- h, W_L, Z_L composite, large coupling to ρ .

Partial compositeness

$$\mathcal{L}_{\text{pc}} = -m_\Psi \bar{\Psi} \Psi - y_L f(\bar{q}_L \Psi + h.c.) - y_R f(\bar{u}_R \Psi + h.c.)$$

A Feynman diagram representing a fermion self-energy correction. It consists of a horizontal fermion line. On the left, a double line with an arrow pointing right is labeled Ψ . This line enters a loop structure. The loop is formed by two lines: a solid line and a dashed line. The solid line continues to the right as a single line with an arrow, labeled q_L, u_R, d_R . The dashed line forms the upper part of the loop. The fermion line is labeled $\sim y$ below it.

$$y_f = \frac{m_\Psi}{f} \sin \phi_L^f \sin \phi_R^f$$

$$\sin \phi_{L,R} \equiv \frac{y_{L,R}}{\sqrt{(m_\Psi/f)^2 + y_{L,R}^2}}.$$

- Mixing angles not completely fixed.
- For example, for top quark, the mixing should be large, $O(1)$.
- ▶ Top quark heavy because it is composite.

Partial compositeness

composite fermion with the same gauge quantum numbers as SM fermion.

$$\mathcal{L}_{\text{pc}} = -m_\Psi \bar{\Psi} \Psi + y_L f (\bar{q}_L \Psi + h.c.) - y_R f (\bar{u}_R \Psi + h.c.)$$

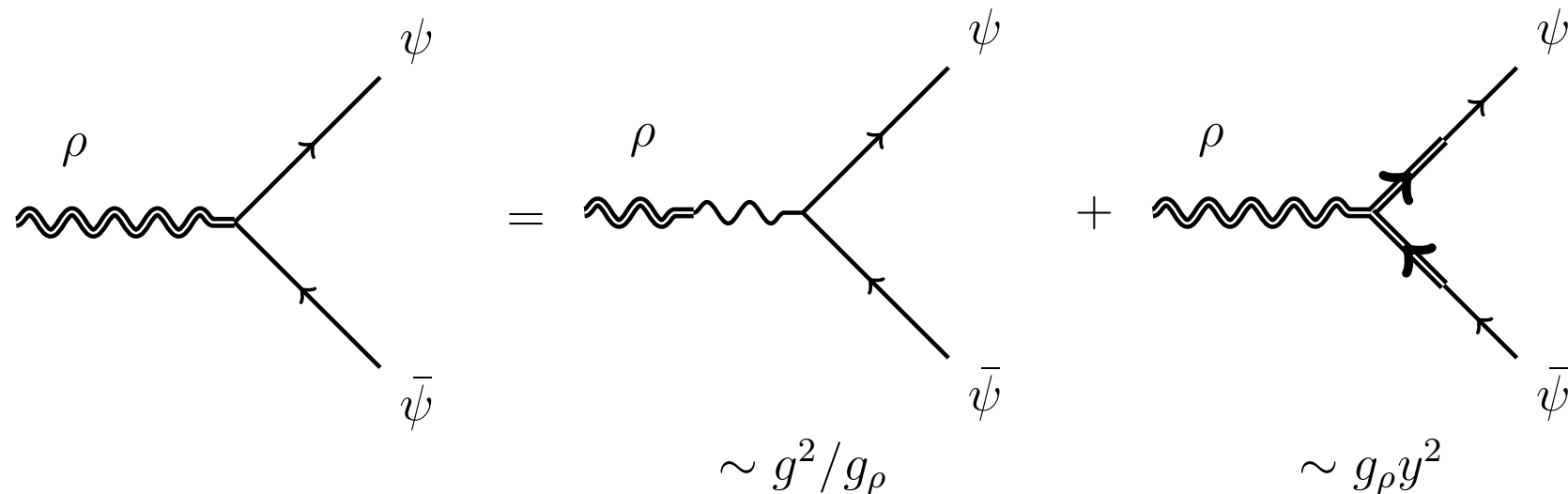
$$\begin{array}{c} \Psi \\ \hline \longrightarrow \quad q_L, u_R, d_R \\ \sim y \end{array}$$

$$y_f = \frac{m_\Psi}{f} \sin \phi_L^f \sin \phi_R^f$$

$$\sin \phi_{L,R} \equiv \frac{y_{L,R}}{\sqrt{(m_\Psi/f)^2 + y_{L,R}^2}}.$$

- Mixing angles not completely fixed.
- For example, for top quark, the mixing should be large, $O(1)$.
 - Top quark heavy because it is composite.

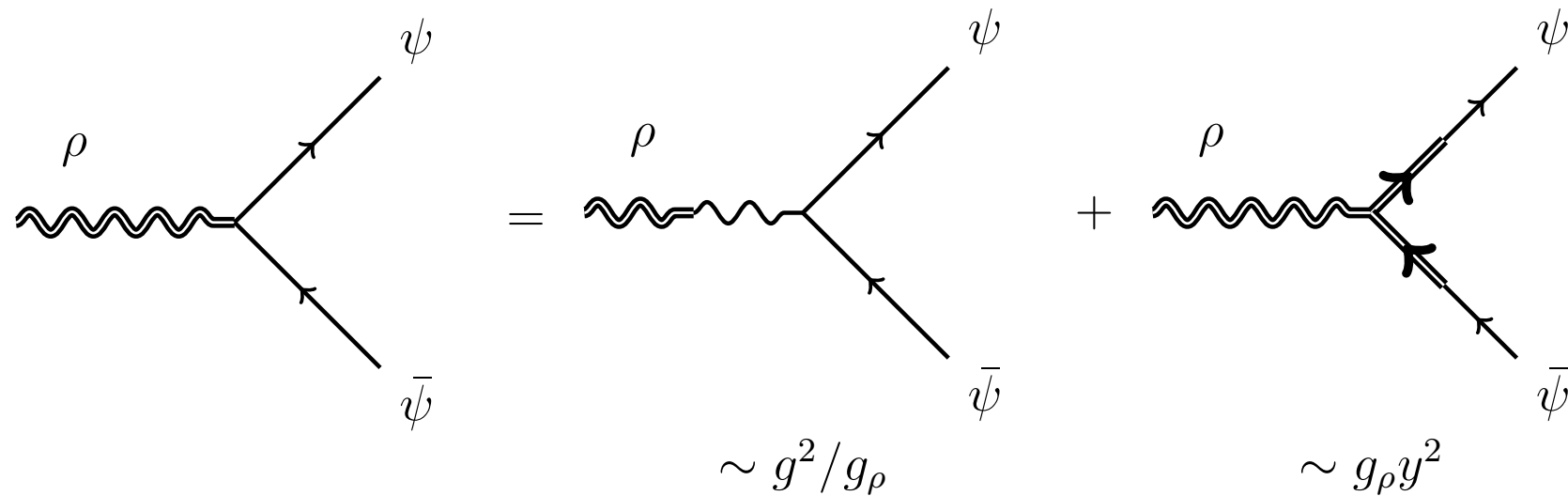
ρ coupling to SM fermion



ρ mixes with W/Z
mixing angle: g/g_ρ

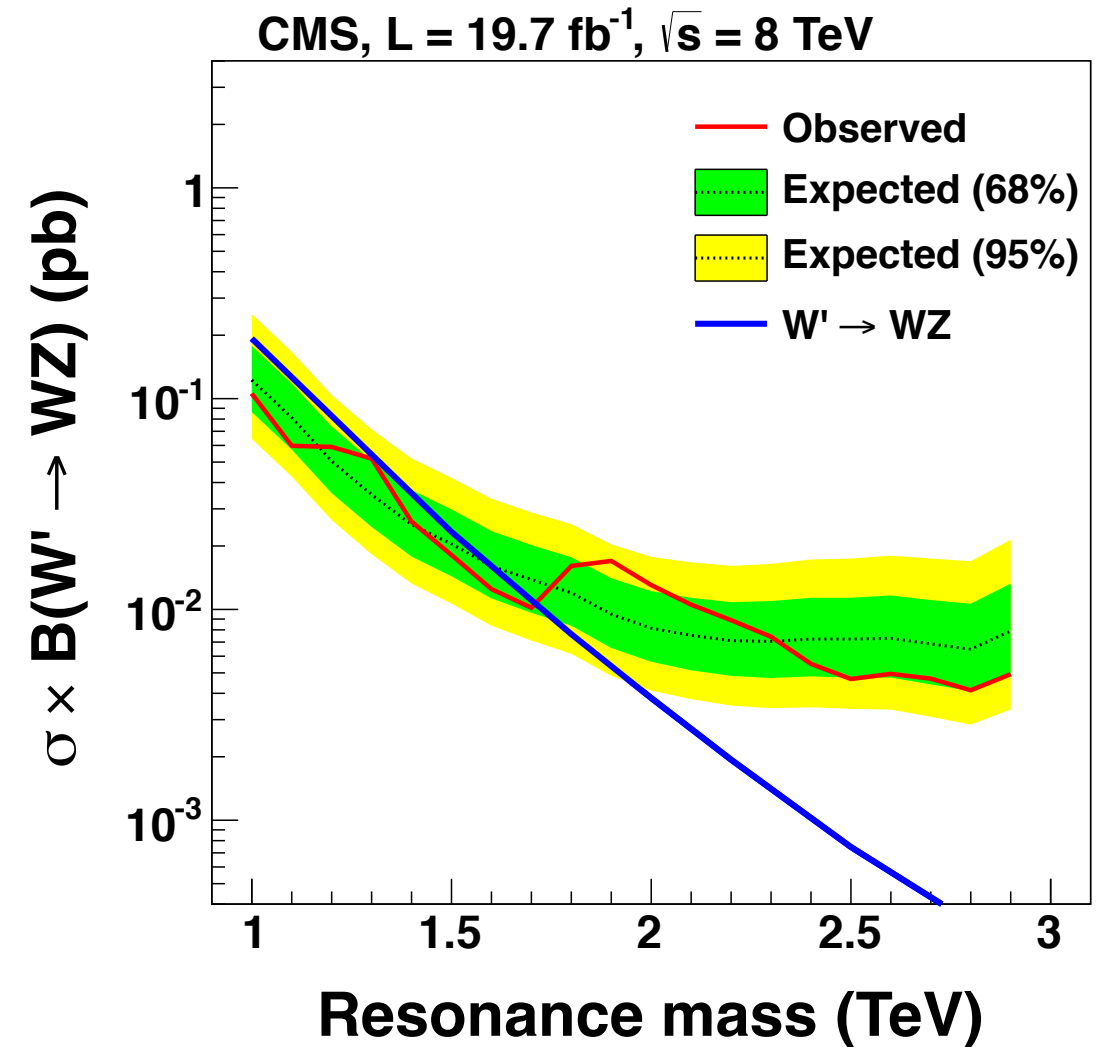
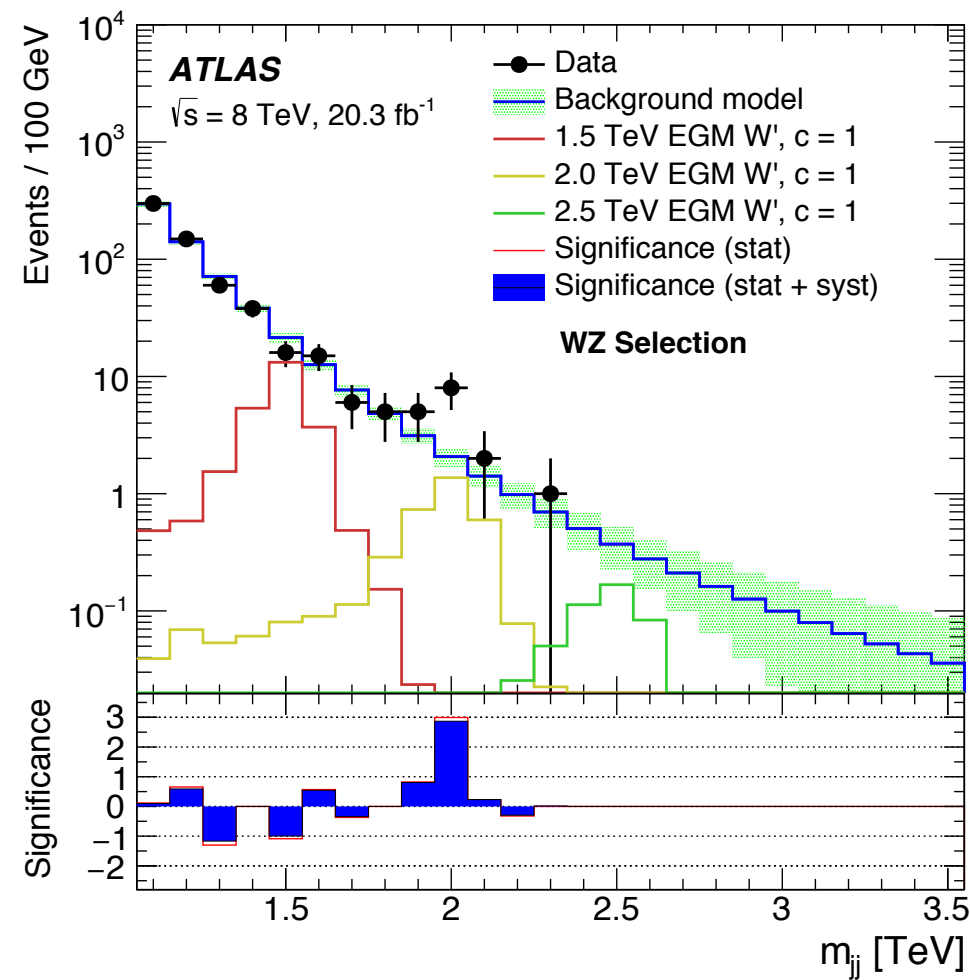
ρ couples to composite
fermion first, which mixes
with SM fermion.

ρ coupling to SM fermion



	VV, Vh	$\bar{q}_L \gamma^\mu q_L$	$\bar{u}_R \gamma^\mu u_R$	$\bar{d}_R \gamma^\mu d_R$	$\bar{\ell}_L \gamma^\mu \ell_L$	$\bar{e}_R \gamma^\mu e_R$
$\rho^{0,\pm}$	g_ρ	$-\frac{g^2}{g_\rho} (1 - a_L \frac{g_\rho^2}{g^2} s_{L,q}^2) \tau^a$	—	—	$-\frac{g^2}{g_\rho} \tau^a$	—
ρ_B^0	g_ρ	$-\frac{1}{6} \frac{g'^2}{g_\rho} (1 + 3a_L \frac{g_\rho^2}{g'^2} s_{L,q}^2)$	$-\frac{2}{3} \frac{g'^2}{g_\rho}$	$\frac{1}{3} \frac{g'^2}{g_\rho}$	$\frac{1}{2} \frac{g'^2}{g_\rho}$	$\frac{g'^2}{g_\rho}$
$\rho_C^\pm$	g_ρ	—	—	—	—	—

Excess around 2 TeV?



Composite spin-1 vector?

Thamm, Torre, Wulzer
Bian, Liu, Shu
Low, Tesi, LTW

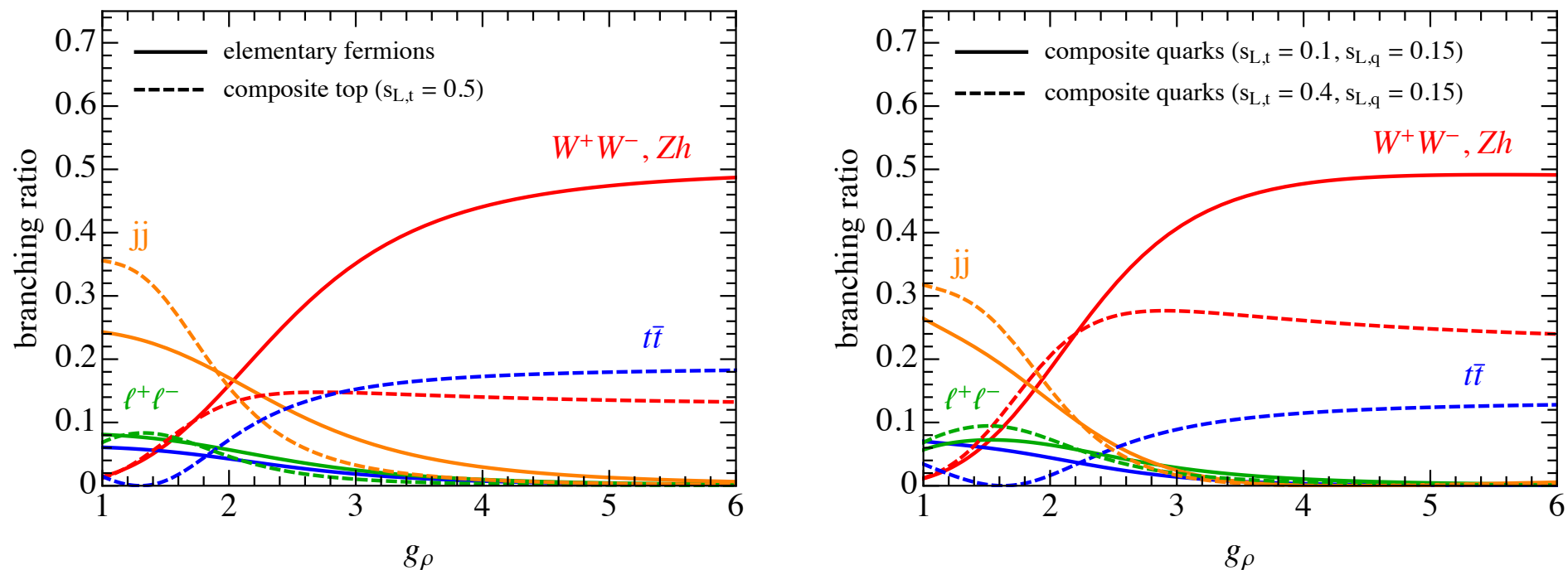
...

Some features of the “excess”

- Data is confusing, no clear picture.
- Large-ish rate. 5–10 fb.
- Not seen else where. Diboson is the leading channel.
- “Usual” gauge boson.
 - ▶ Di-lepton and single lepton limits < 1 fb.
 - ▶ L-R models.
- Scalar?
 - ▶ Rate small. BR to WZ small.

Decay of composite spin-1 res.

m_ρ fixes to be 2 TeV



- BR to diboson is large. Suppressed fermion coupling. Could have large rate.
- c.f., usual gauge boson, small BR to diboson.
- Suppressed fermion coupling $\rightarrow$ suppress di-lepton mode.

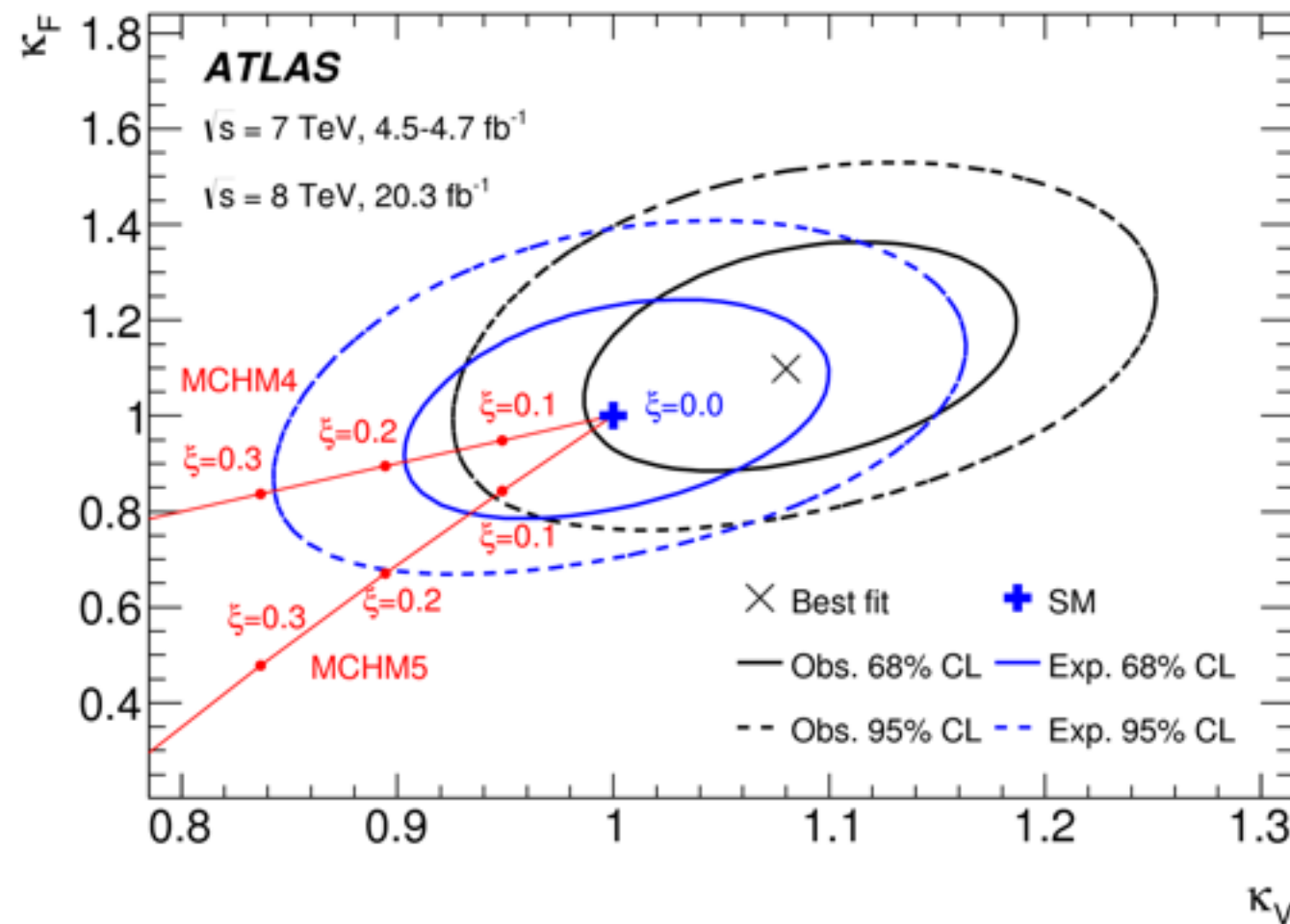
$$\frac{\text{BR}(\rho^0 \rightarrow \ell^+\ell^-)}{\text{BR}(\rho^0 \rightarrow W^+W^-)} = \frac{8}{c_H^2} \frac{g^4}{g_\rho^4}$$

Constraints other than di-boson

final state		ATLAS		CMS	
Can still be relevant	$\ell^+ \ell^-$	0.2 fb	[3]	0.25 fb	[4]
	$\ell^\pm \cancel{E}_T$	0.9 fb	[66]	0.4 fb	[67]
Not as strong	$t\bar{b}$	120 fb	[68]	100 fb	[69]
	$t\bar{t}$	50 fb	[70]	20 fb	[71]
	$j\bar{j}$	130 fb	[72]	100 fb	[73]

Indirect constraints

Higgs coupling.



$$c_{hVV} = \sqrt{1 - \frac{v^2}{f^2}} c_{hVV}^{\text{SM}}$$

— Higgs couplings.

- $f > 500\text{--}600 \text{ GeV} > v$. Some fine tuning seems unavoidable.

EW precision.

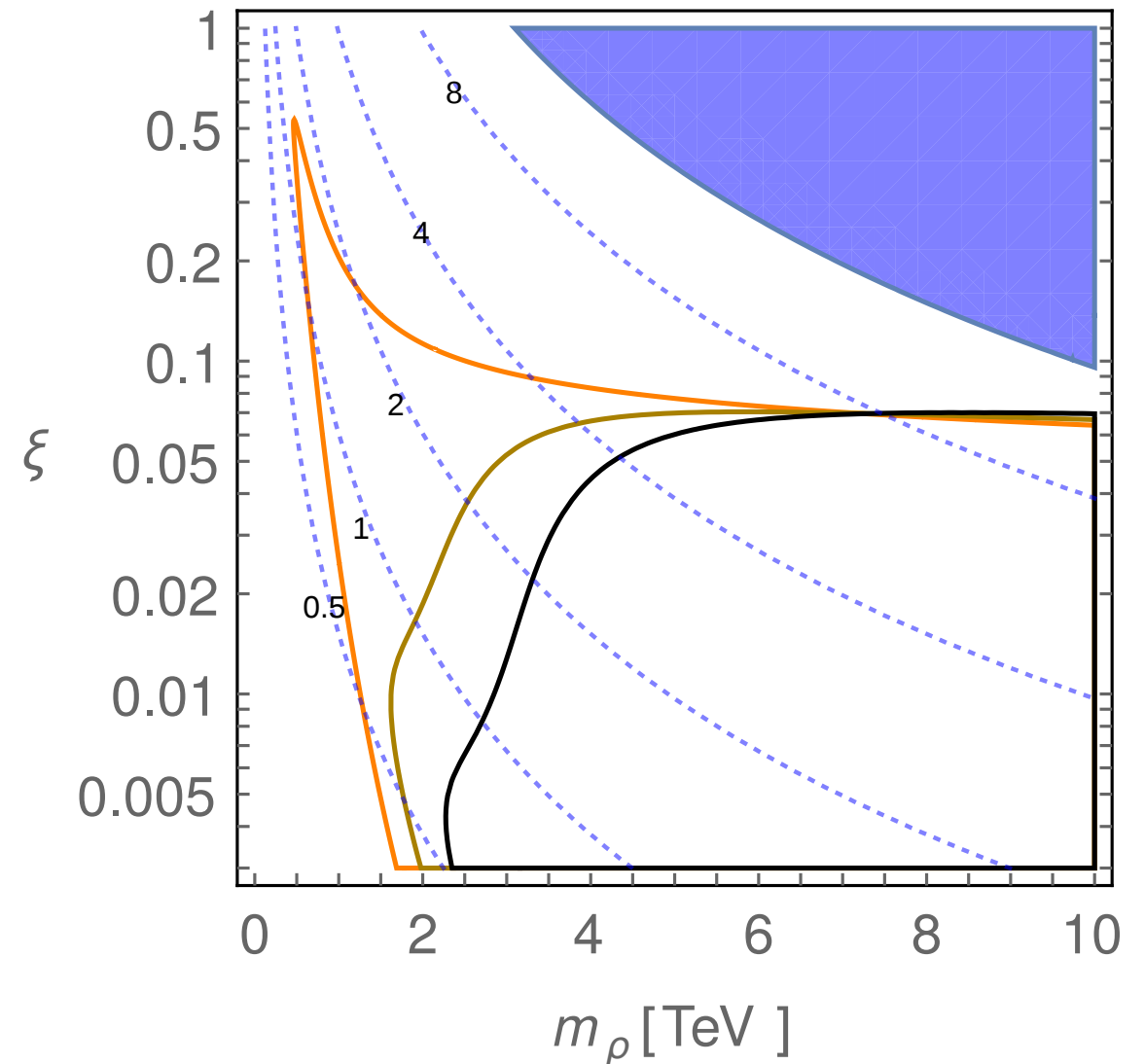
Tree level S-parameter

$$S \sim \frac{4\pi v^2}{m_\rho^2}$$

One loop to S and T

$$\frac{1}{16\pi^2} \frac{v^2}{f^2} \log(m_\rho/m_h)$$

Additional UV contribution
parameterized by general CCWZ



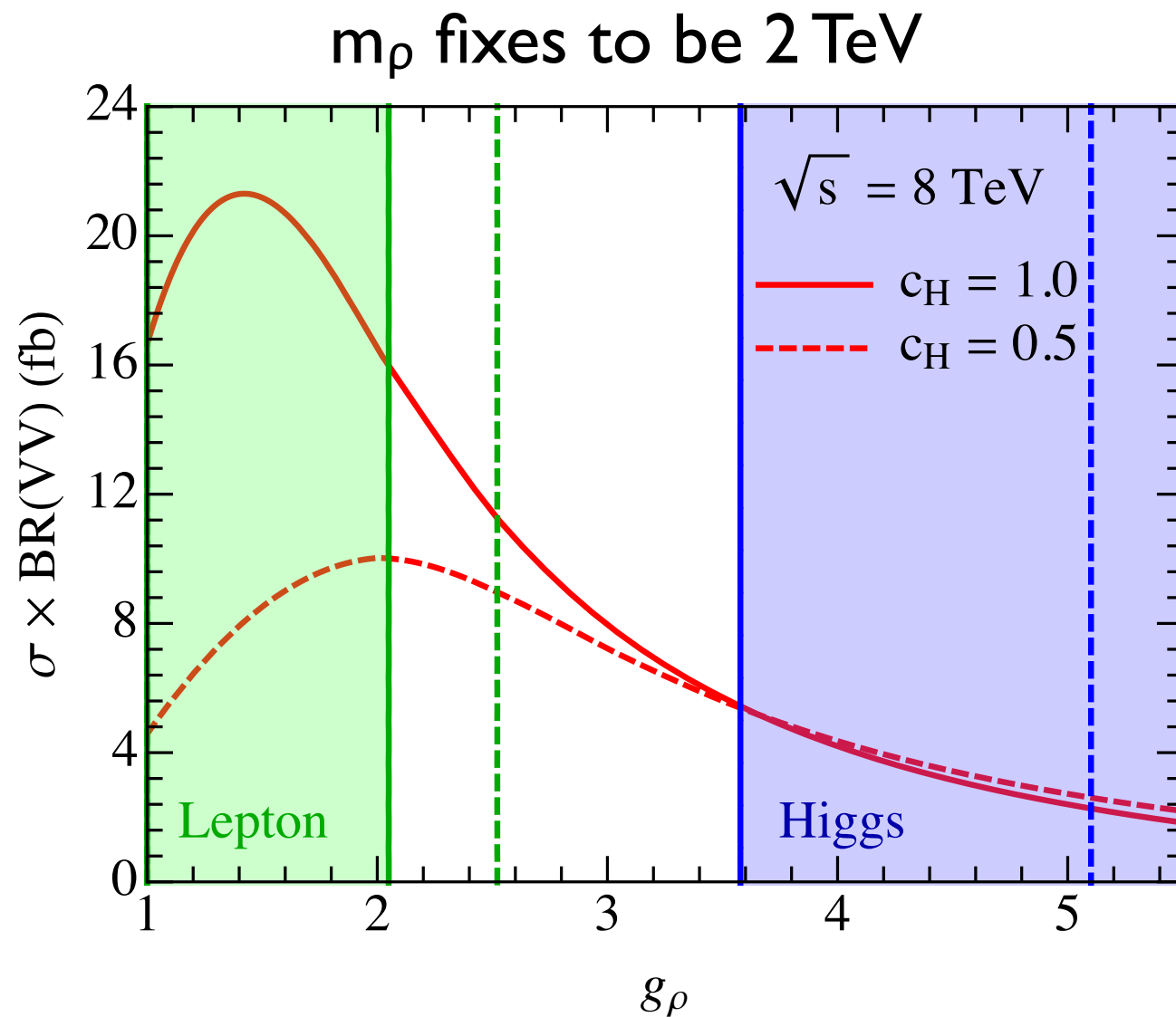
Contino, Salvarezza, 2015

In addition, constraints from Z decay. Relevant for partially composite u_L

$$R_h : \quad \delta g_{u_L} = \frac{1}{4} \frac{v^2}{f^2} s_{L,u}^2 < 0.5 \times 10^{-3}$$

Parameter space for 2 TeV
resonance

(over)simplified model.



final state	ATLAS		CMS	
$\ell^+ \ell^-$	0.2 fb	[3]	0.25 fb	[4]
$\ell^\pm \cancel{E}_T$	0.9 fb	[66]	0.4 fb	[67]
$t\bar{b}$	120 fb	[68]	100 fb	[69]
$t\bar{t}$	50 fb	[70]	20 fb	[71]
jj	130 fb	[72]	100 fb	[73]

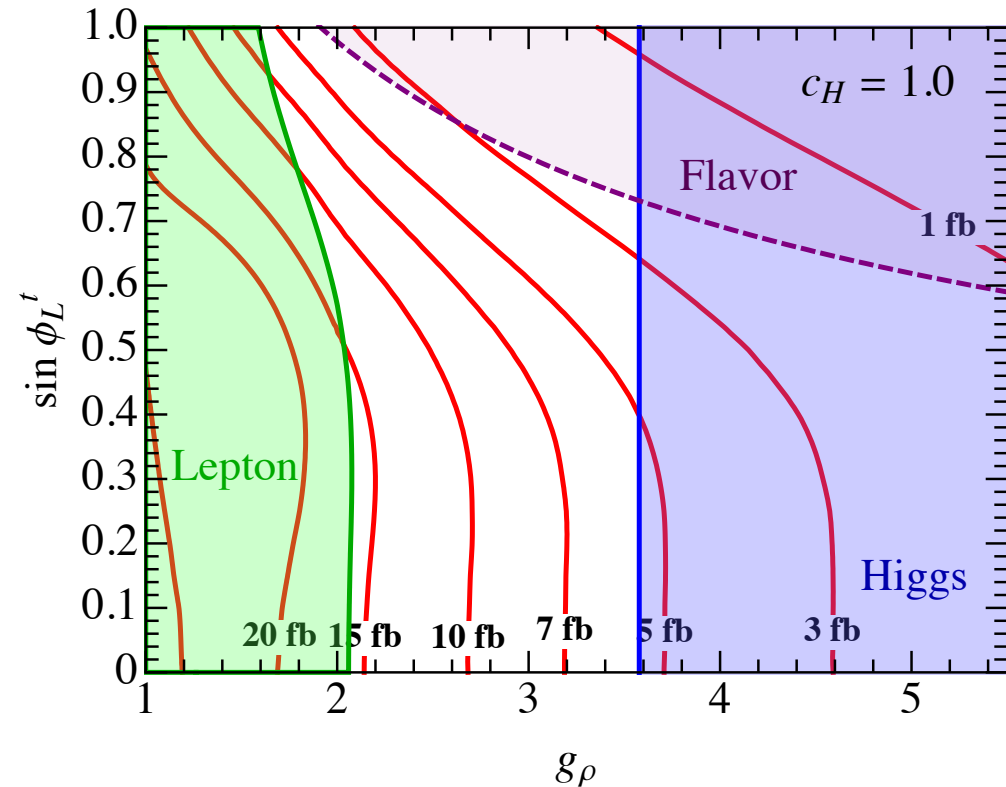
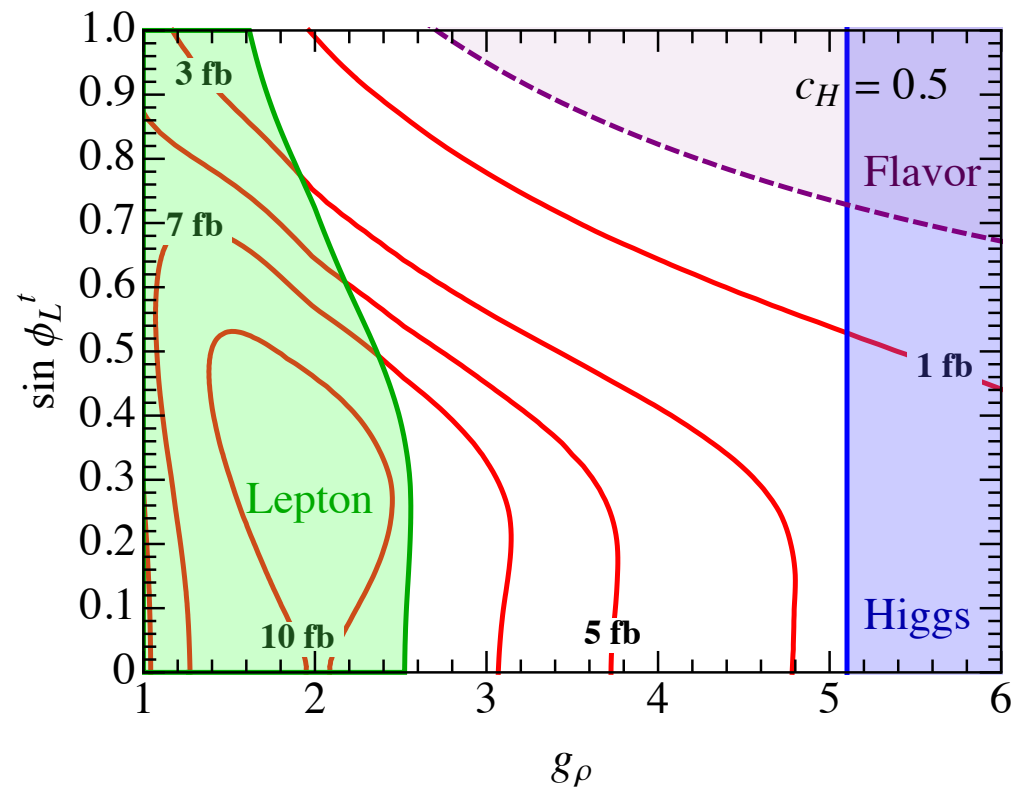
$$\mathcal{L} \supset \underline{i g_\rho c_H \rho_\mu^a} (H^\dagger \tau^a D^\mu H - (D^\mu H)^\dagger \tau^a H)$$

$$m_\rho = \sqrt{\underline{c_H}} g_\rho f$$

— Elementary fermion. Toy model.

Composite top.

m_ρ fixes to be 2 TeV

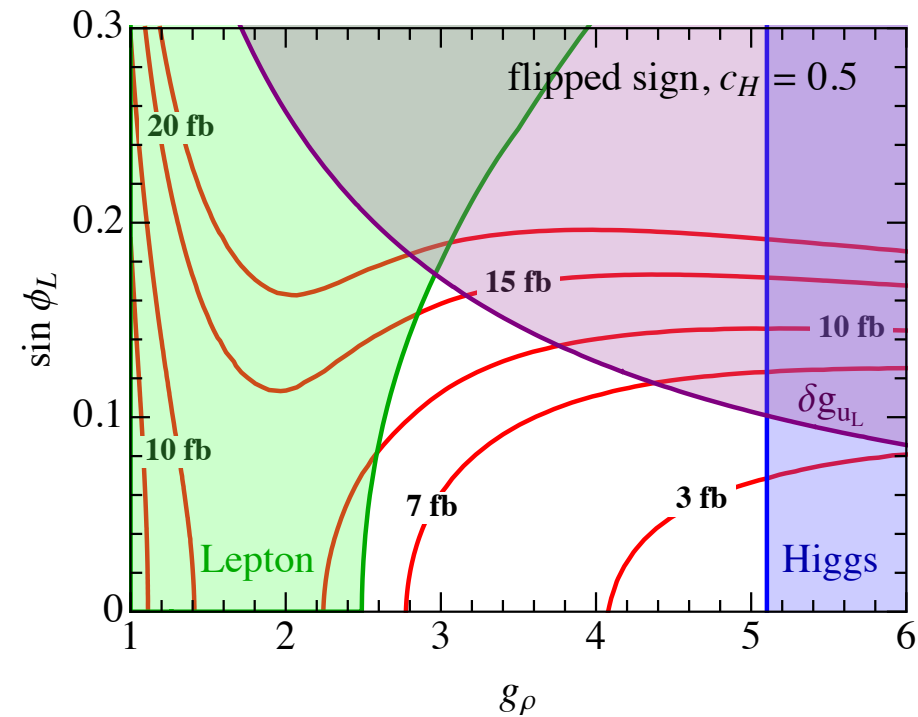
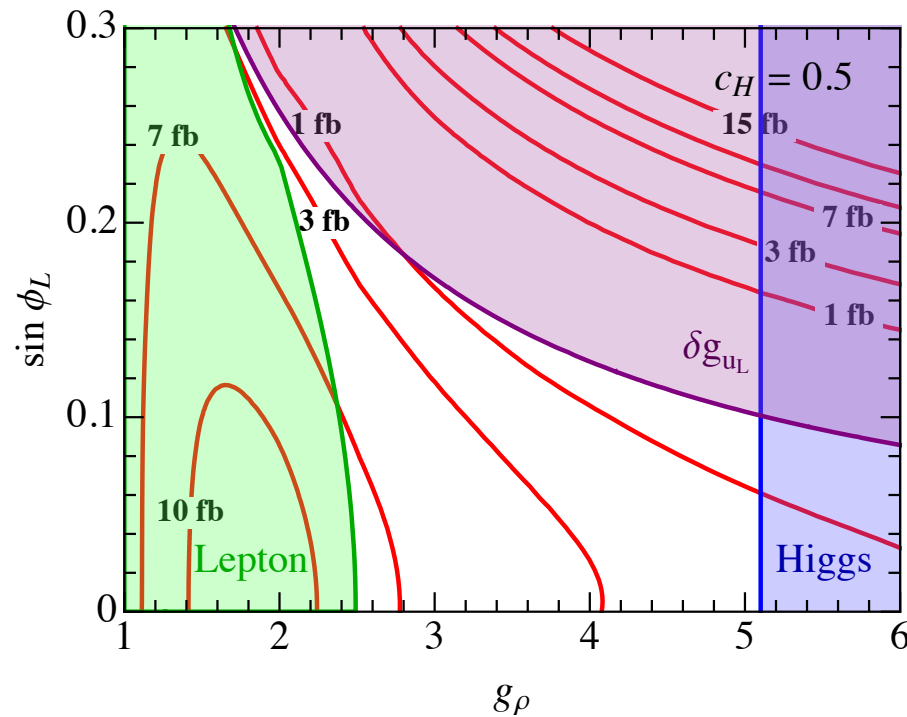


- Vary top compositeness.
- 1st and 2nd generations elementary.

Composite quarks.

m_ρ fixes to be 2 TeV

$\sin \phi_L^t = 0.4$



	VV, Vh	$\bar{q}_L \gamma^\mu q_L$
$\rho^{0,\pm}$	g_ρ	$-\frac{g^2}{g_\rho} (1 - a_L \frac{g_\rho^2}{g^2} s_{L,q}^2) \tau^a$

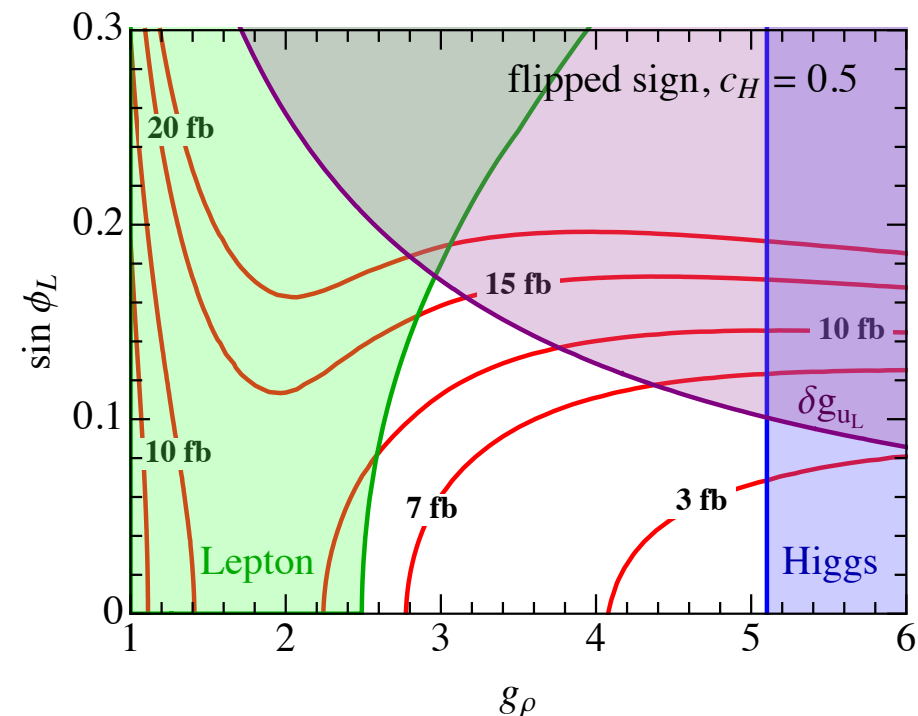
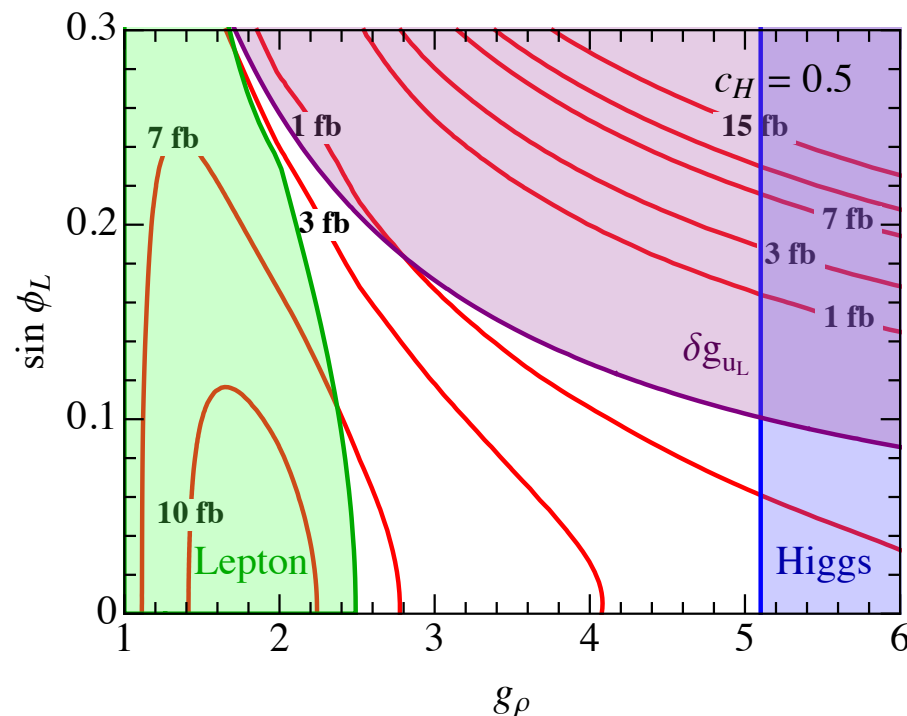
flipped sign: $a_L = -1$

- Vary compositeness of light quarks.

Composite quarks.

m_ρ fixes to be 2 TeV

$\sin \phi_L^t = 0.4$



	VV, Vh	$\bar{q}_L \gamma^\mu q_L$
$\rho^{0,\pm}$	g_ρ	$-\frac{g^2}{g_\rho} \left(1 - a_L \frac{g_\rho^2}{g^2} s_{L,q}^2\right) \tau^a$

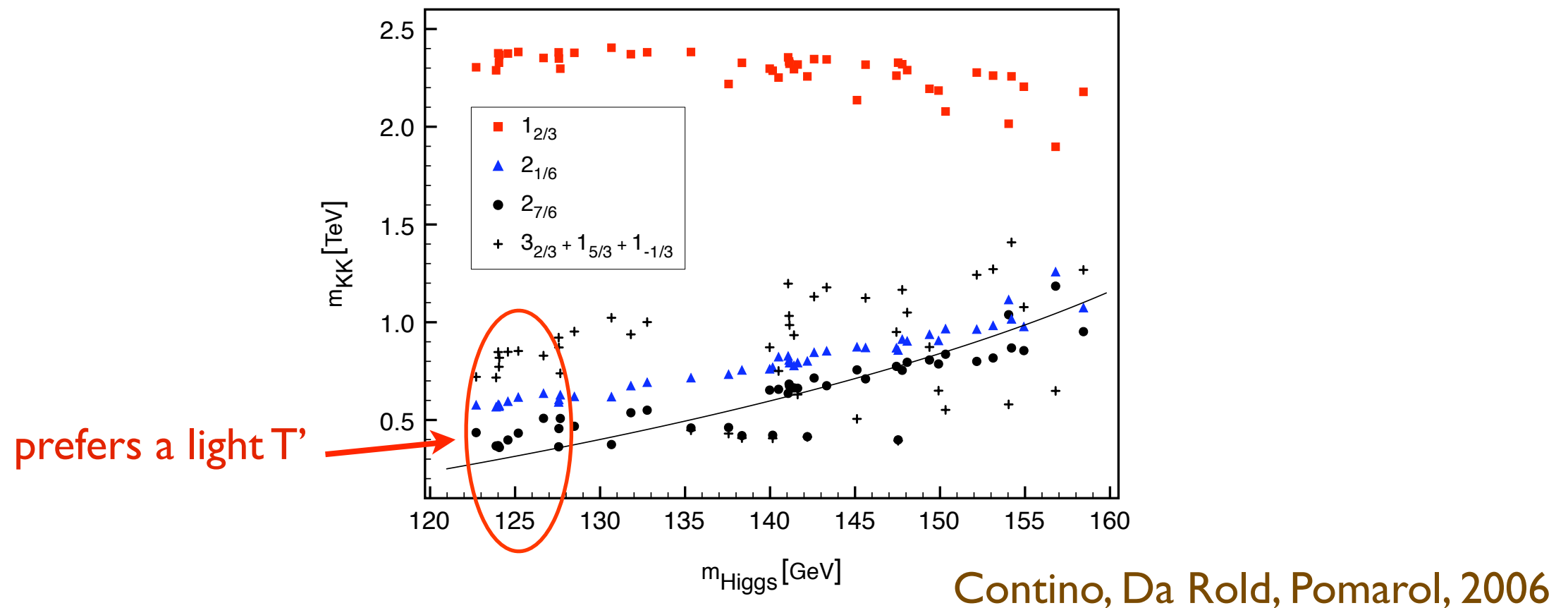
flipped sign: $a_L = -1$

- Vary compositeness of light quarks.

Bottom line

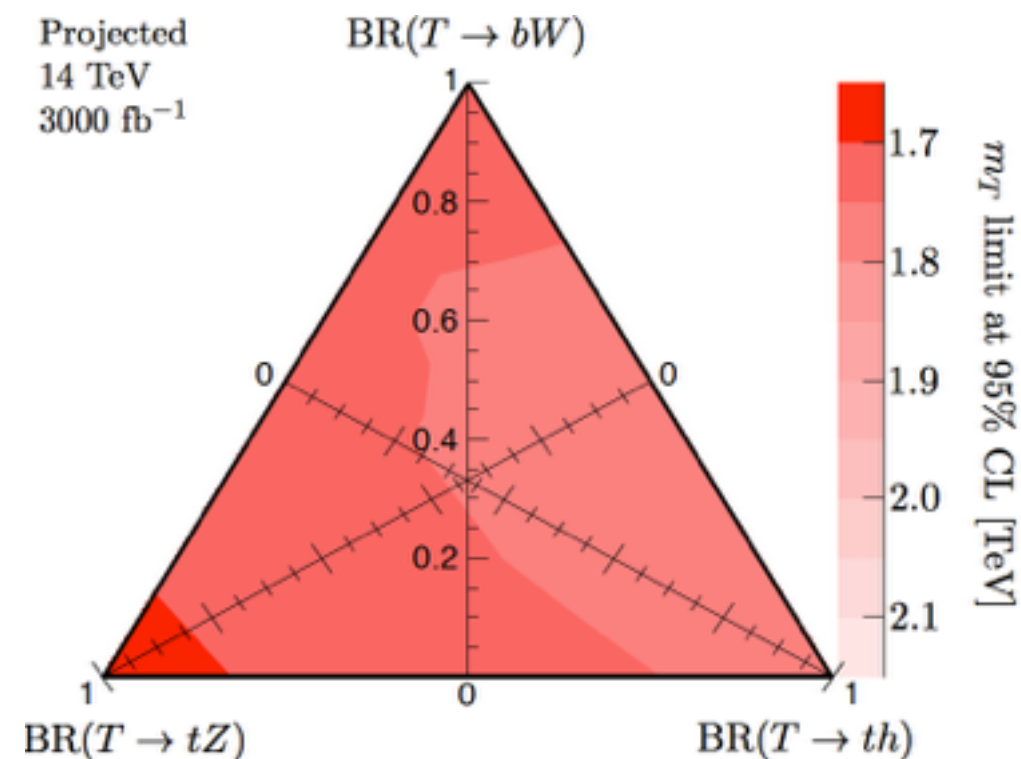
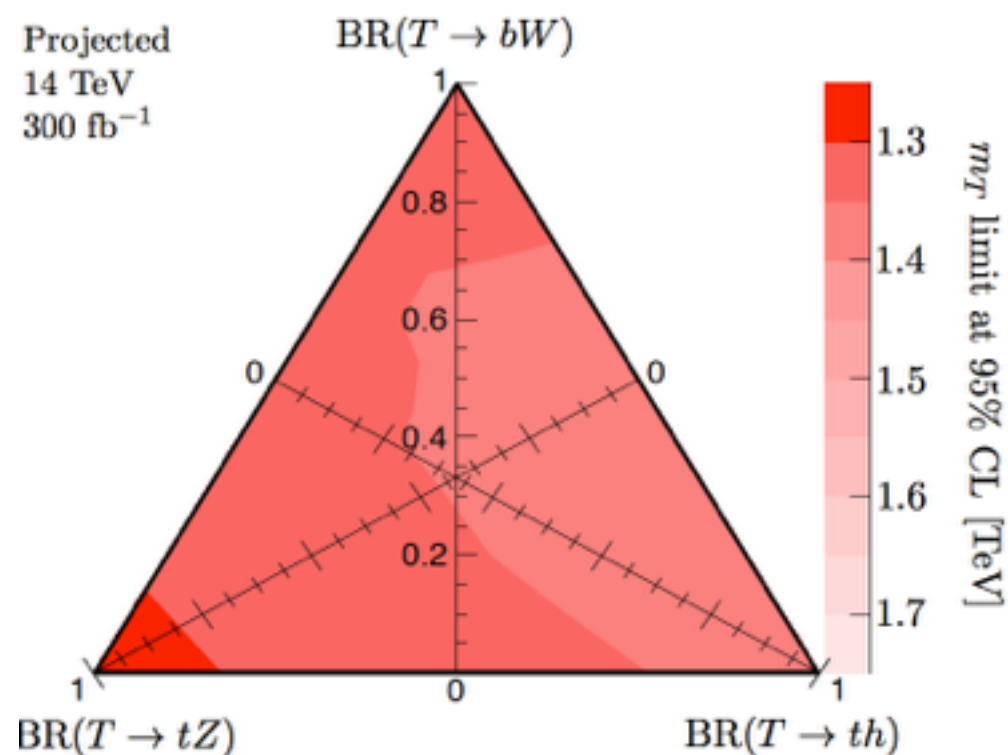
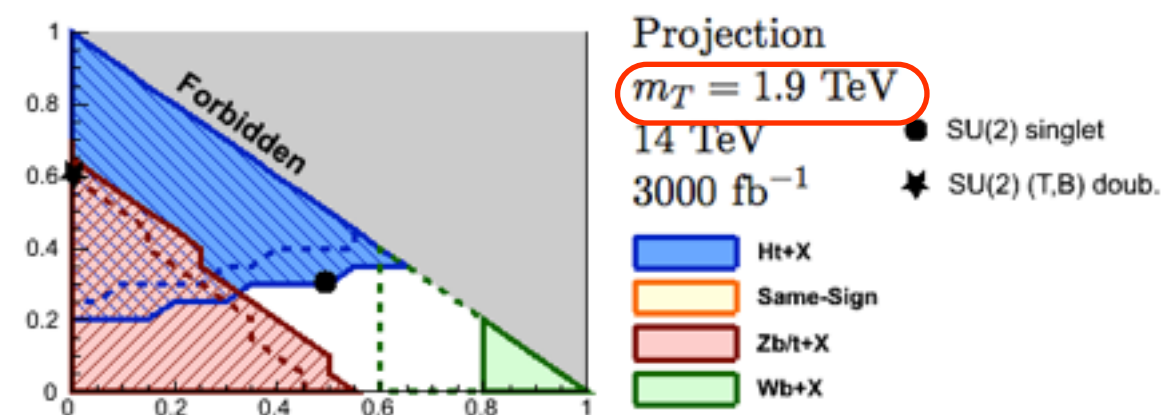
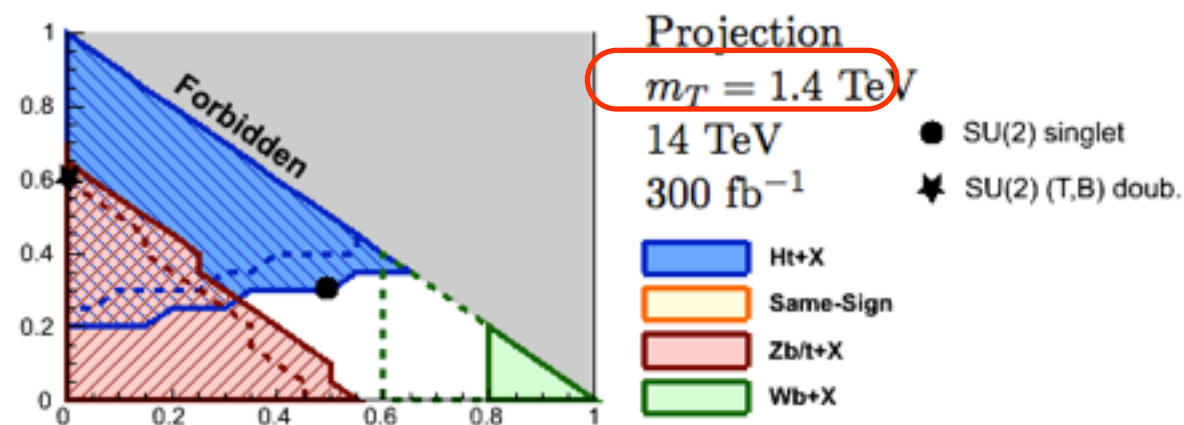
Composite spin 1 resonance can fit the excess and satisfy all constraints without too much effort.

Compositeness and top partner

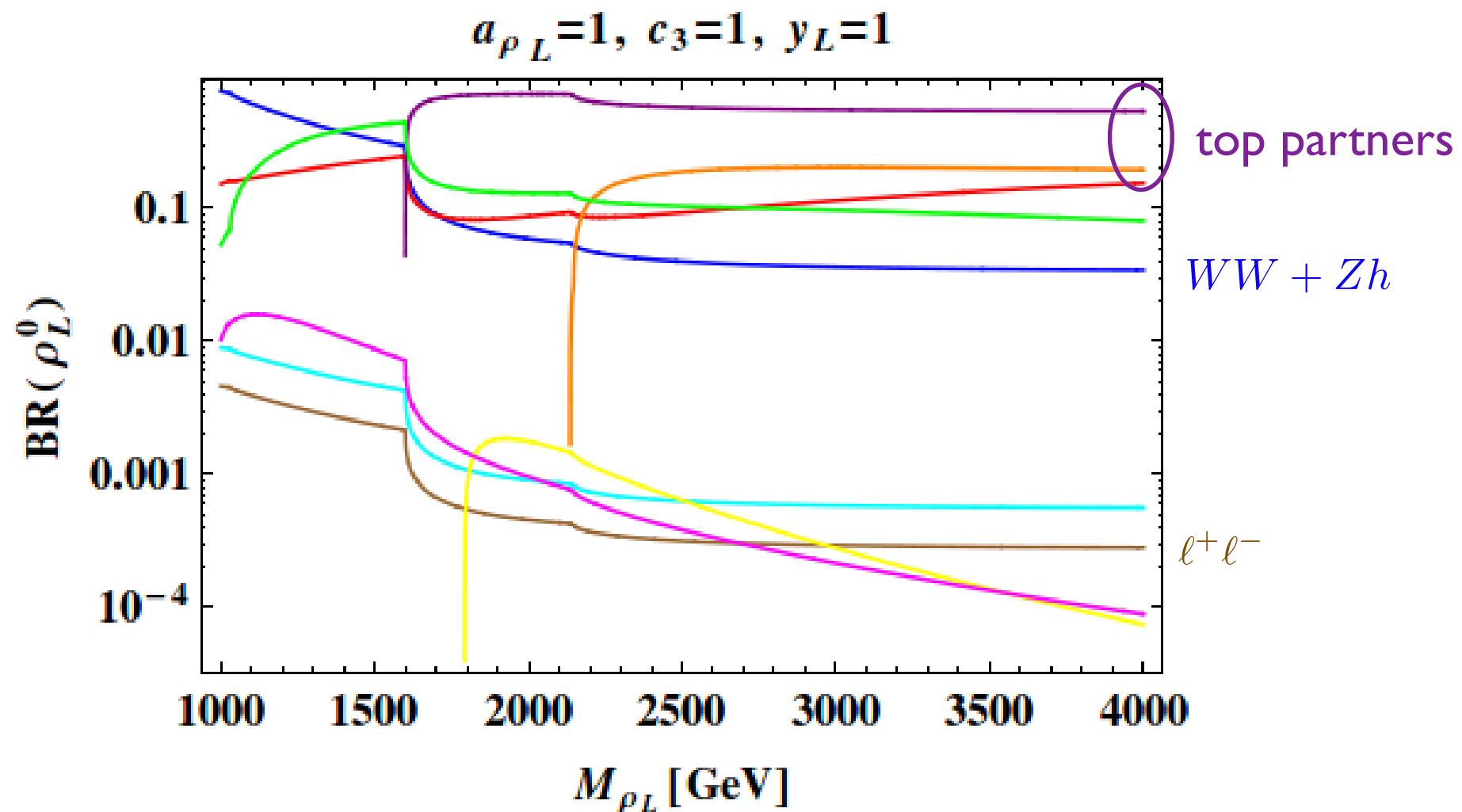


- Light top partner (ψ which mixes with top) could be less than TeV.
- Plays a crucial role in EWSB.

LHC 14 should cover (most of) it.



ρ decaying into top partners



- $\text{BR}(\rho \rightarrow \psi\psi) \sim \mathcal{O}(1)$. Would dominate if allowed.
- Diboson would not be the leading channel.
- Can assume ψ heavy, more fine-tuning.

Top partner

Integrating out top partner $\rightarrow$ Higgs potential

$$V \simeq \frac{N_c}{16\pi^2} m_\Psi^4 \left[a \frac{y_t f}{m_\Psi} F_a + b \left(\frac{y_t f}{m_\Psi} \right)^2 F_b \right]$$

$F_{a,b}$: function of $\frac{h}{f}$
 m_Ψ : mass of top partner

$$m_h^2 \simeq b \frac{N_c y_t^2 v^2}{2\pi^2} \frac{m_\Psi^2}{f^2}$$

Light Higgs $\rightarrow$ light top partner

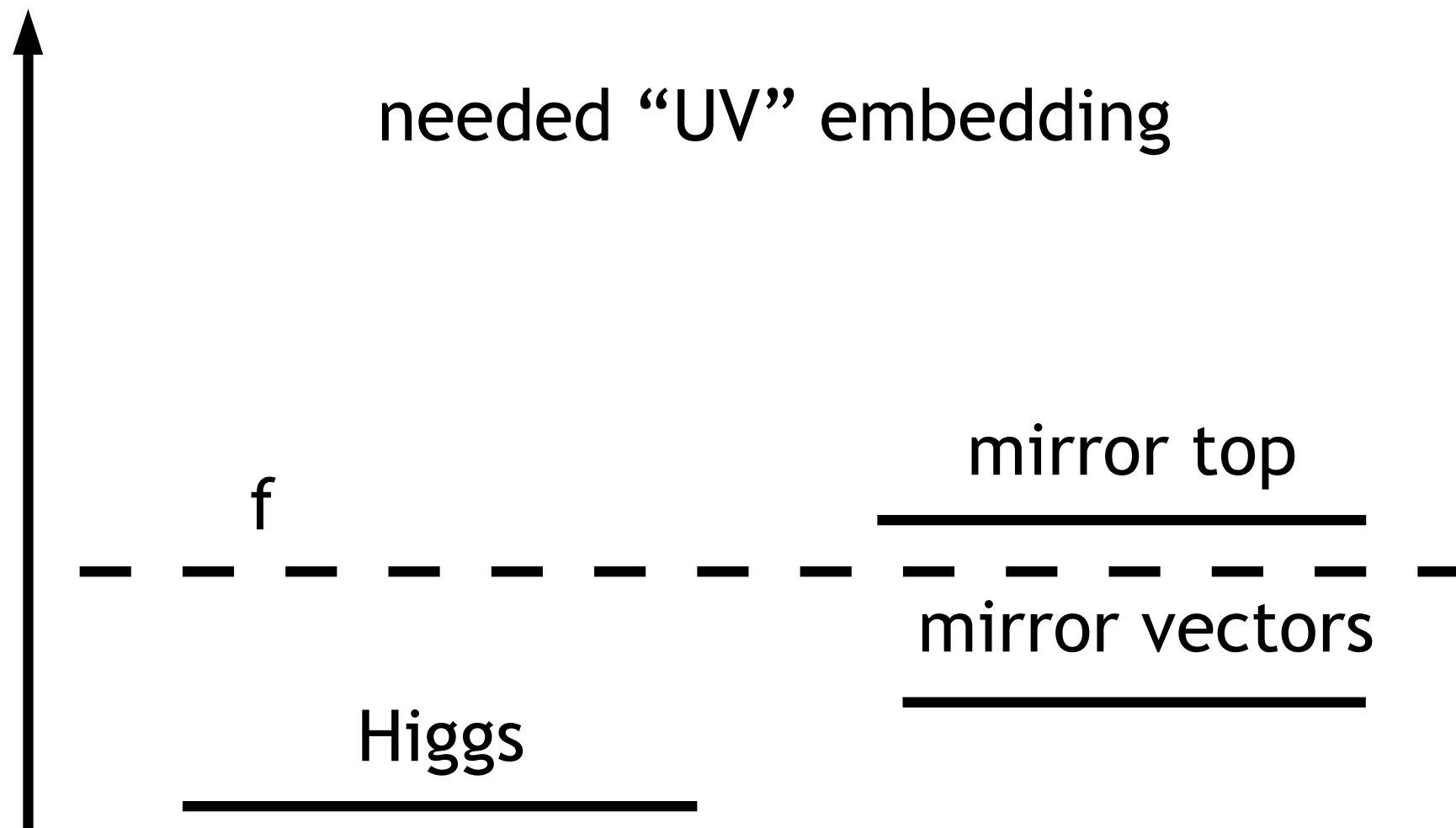
Top partner colored.

LHC searches already constrains top partner mass.

Run 2 will completely cover the simplest cases.

There is an exception: twin Higgs. Singlet top partner.

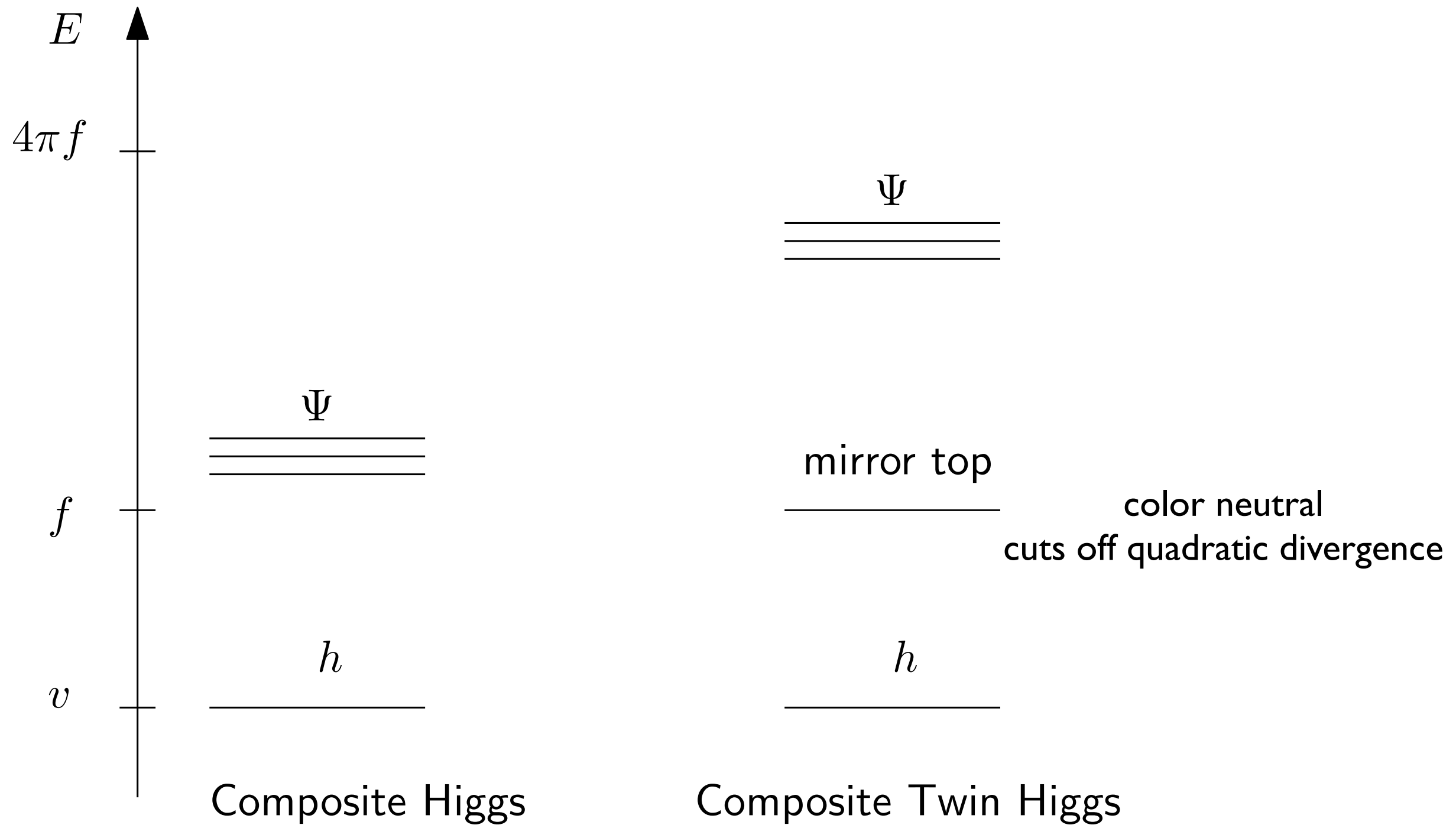
Twin Higgs.



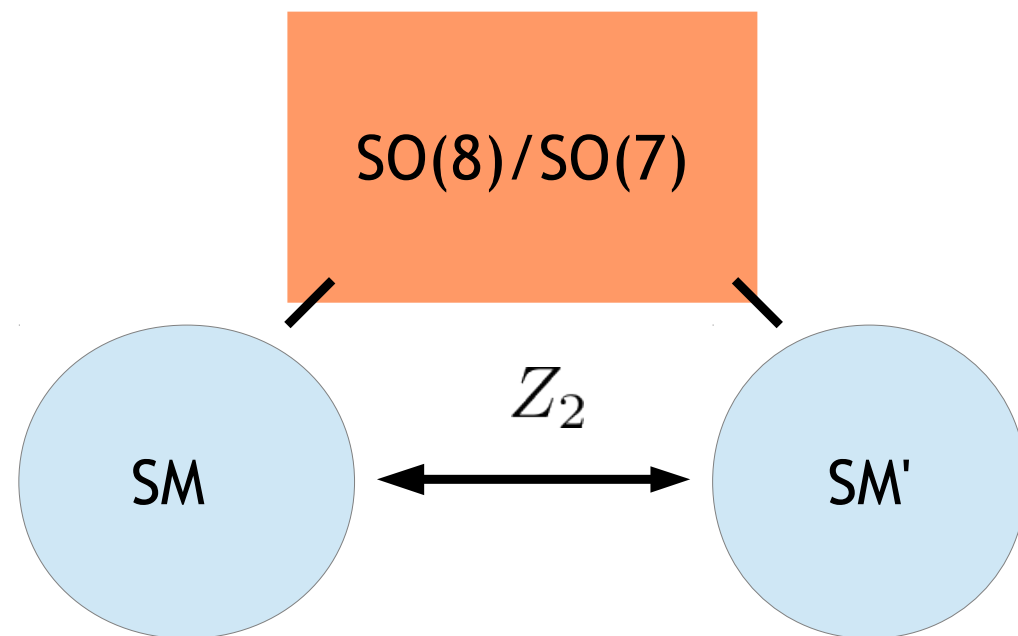
Twin composite Higgs.

Low, Tesi, LTW

Barbieri, Greco, Rattazzi, Wulzer



Twin composite Higgs



The gauging of the EW part of SM and SM' is given by

$$\left(\begin{array}{c|c} g \cdot SO(4) & 0 \\ \hline 0 & g' \cdot SO(4)' \end{array} \right)$$

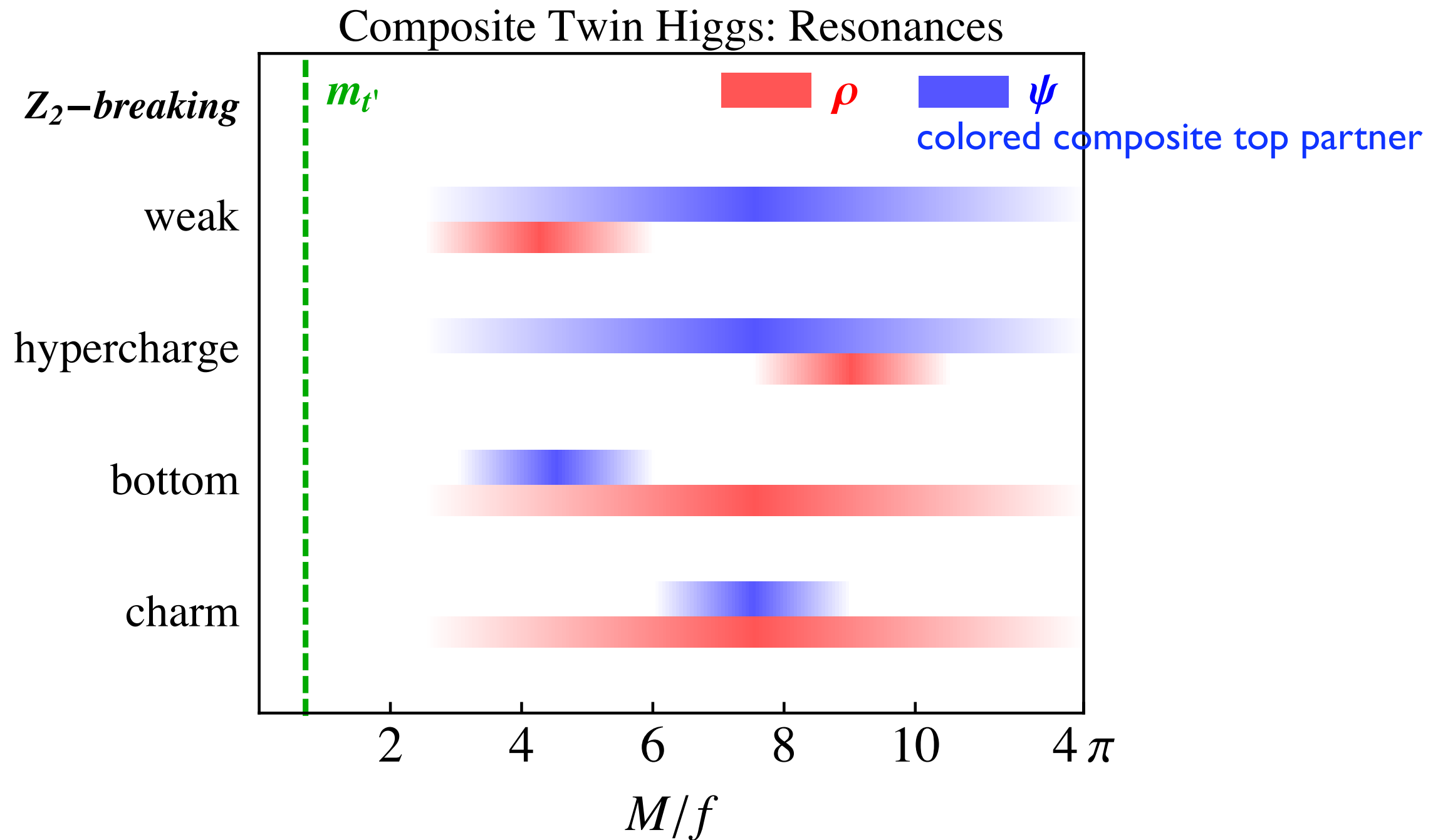
If Z_2 is exact, then SM and SM' has the same scale, $v \simeq f$.

Must introduce Z_2 breaking.

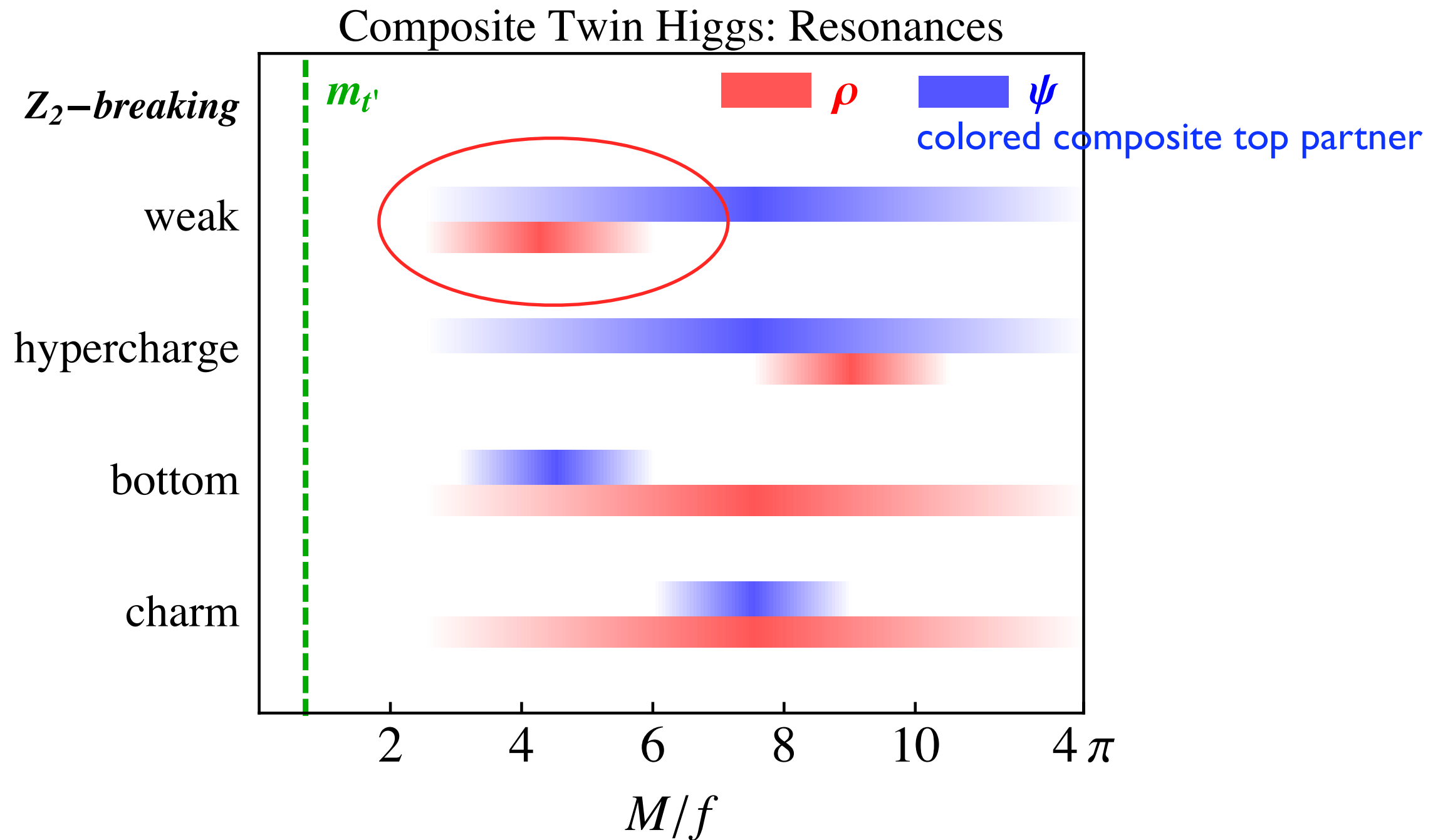
Freedom in choosing how to do it:

Difference in gauge or Yuk interaction between SM and SM'

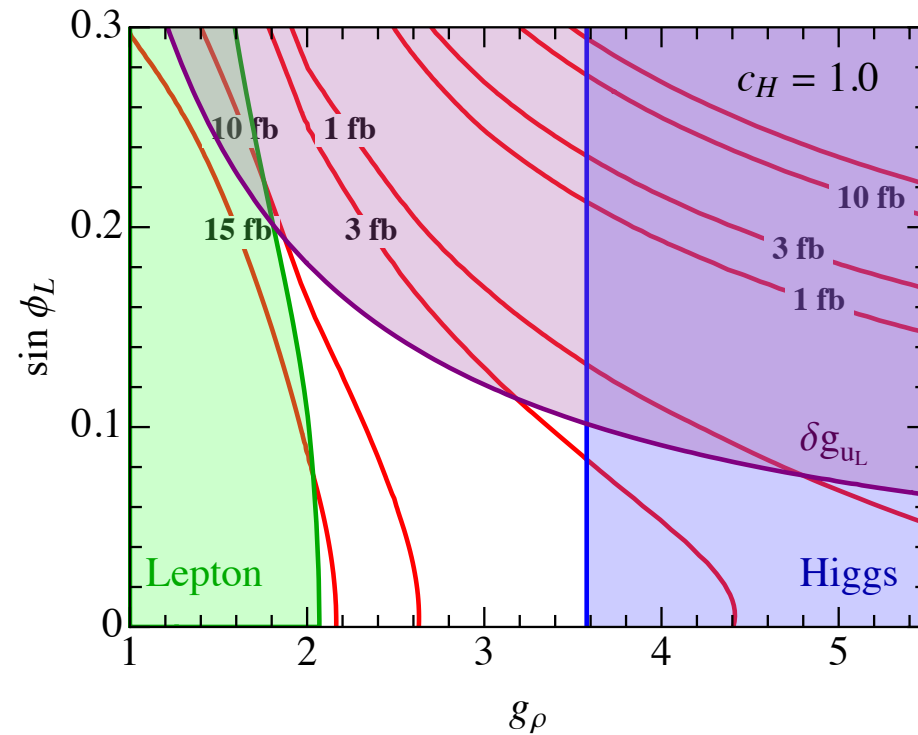
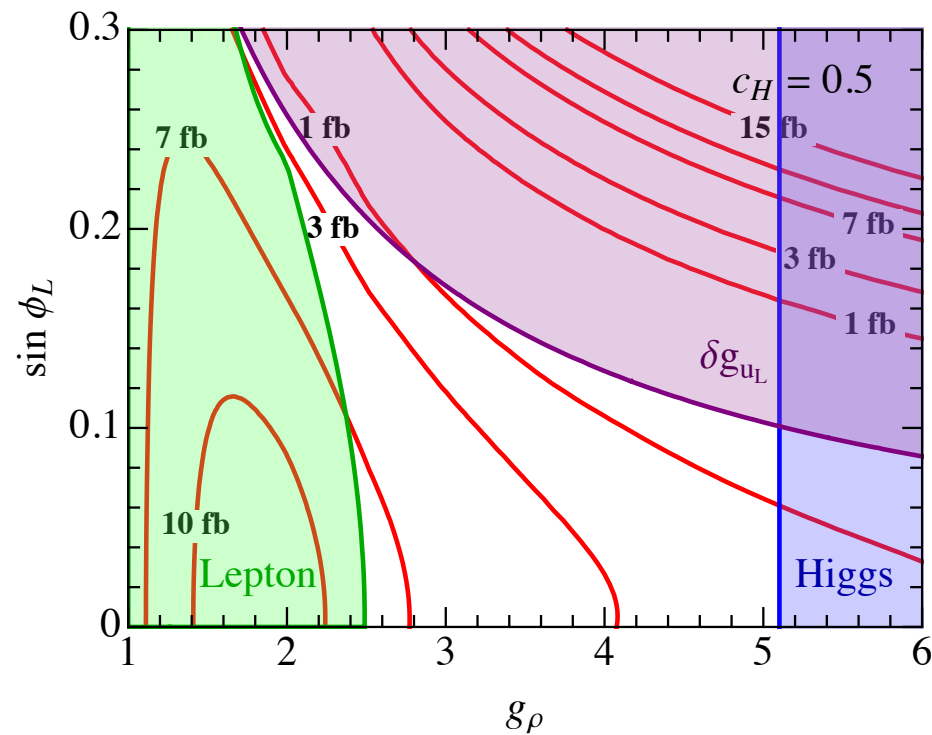
Z_2 breaking and spectrum



Z_2 breaking and spectrum



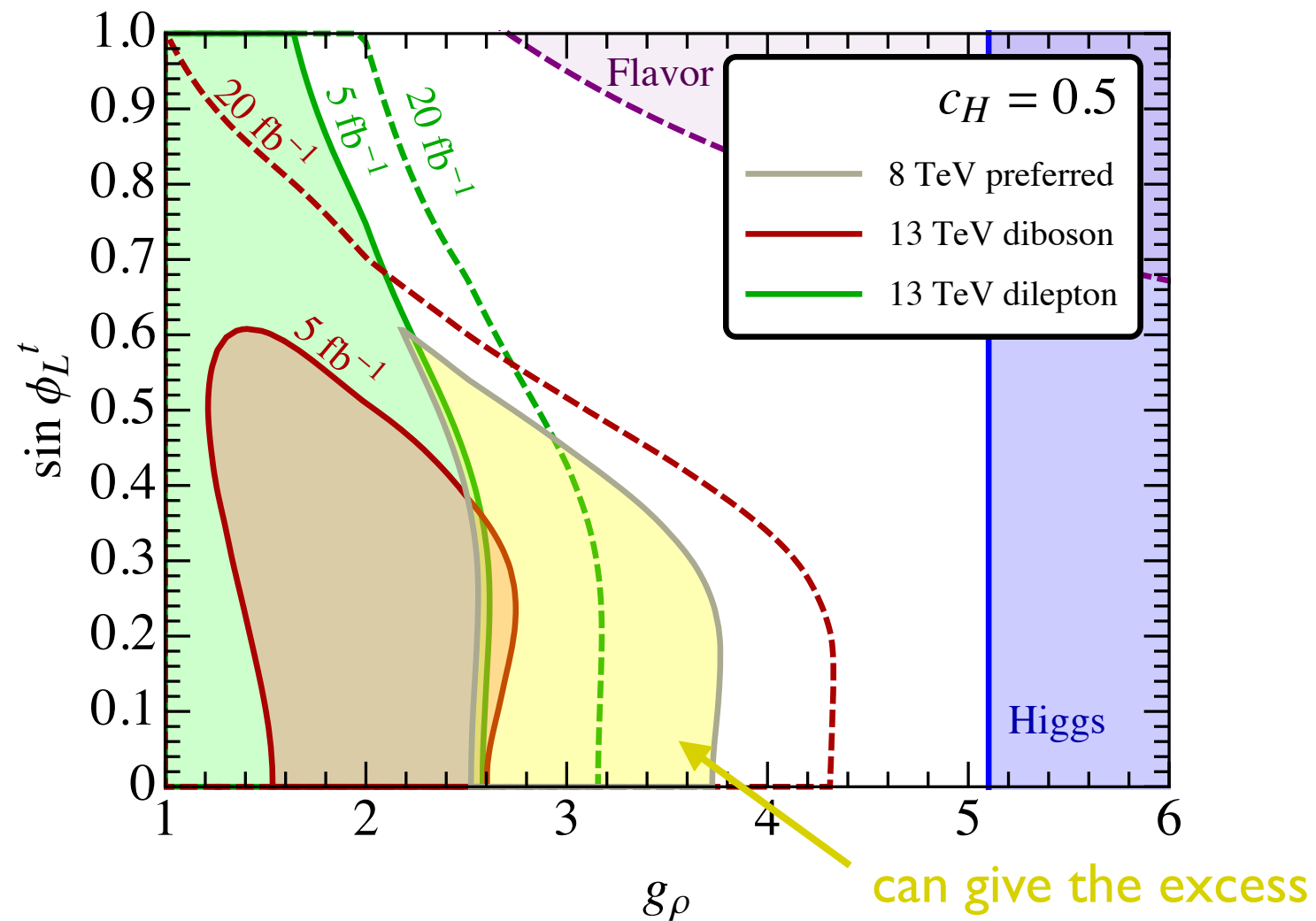
$$\sin \phi_L^t = 0.1$$



$$y_f = \frac{m_\Psi}{f} \sin \phi_L^f \sin \phi_R^f$$

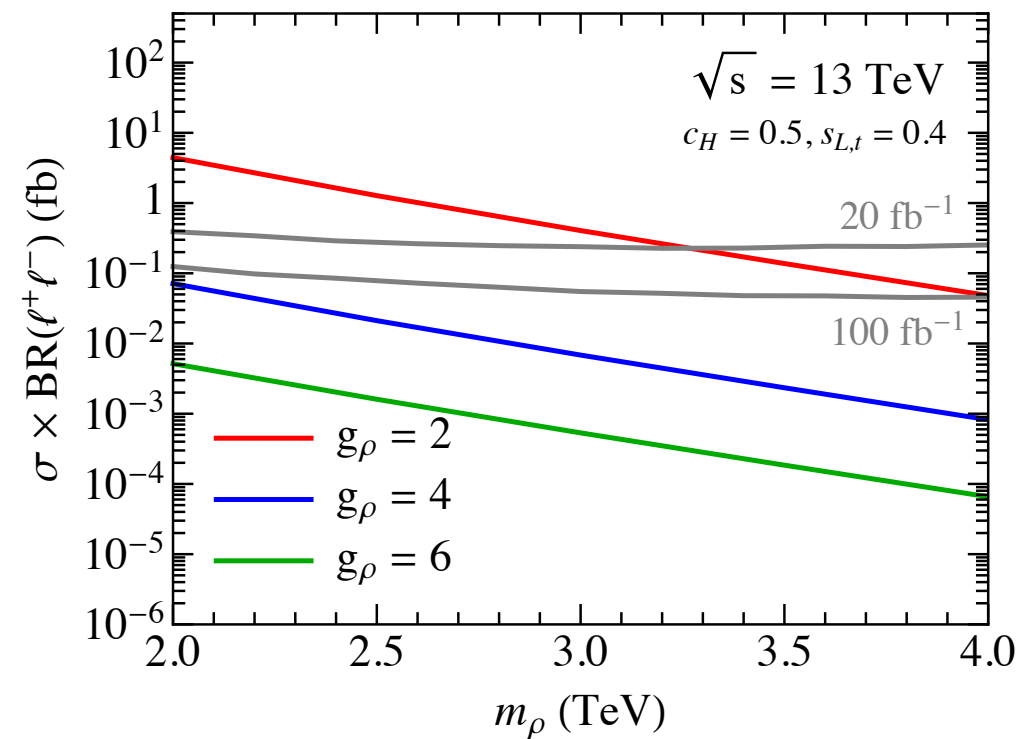
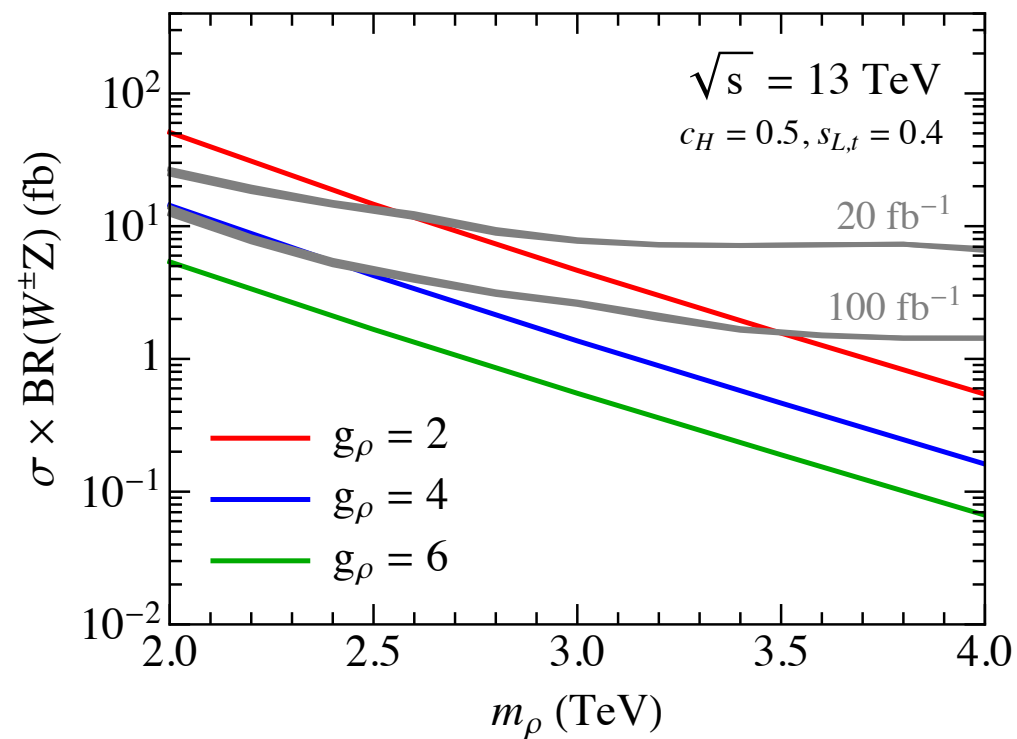
- For a twin composite scenario, composite top can be very heavy.
- Mixing can be smaller.

Run 2 projection



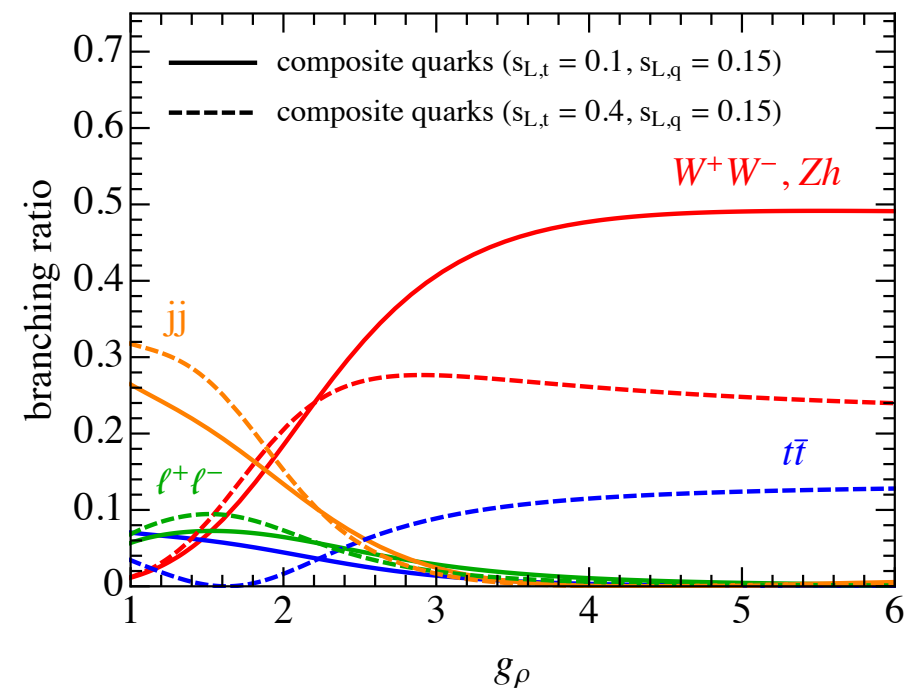
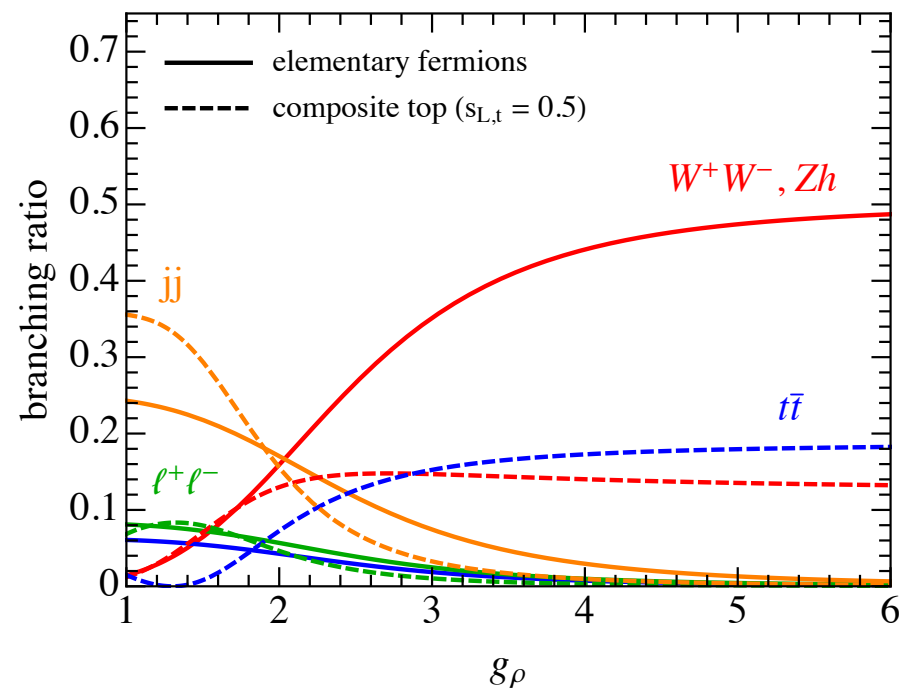
- Can confirm or rule out with modest luminosity.

Reach of Run 2, di-boson and di-lepton



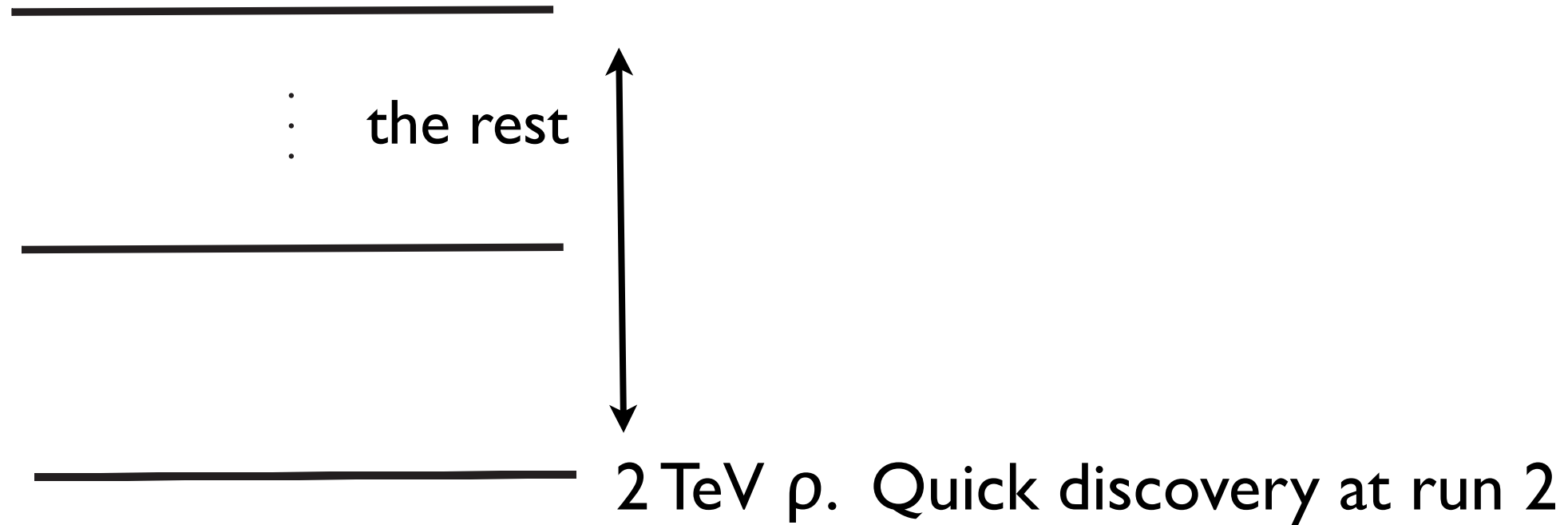
— At most 4 TeV.

Additional channels

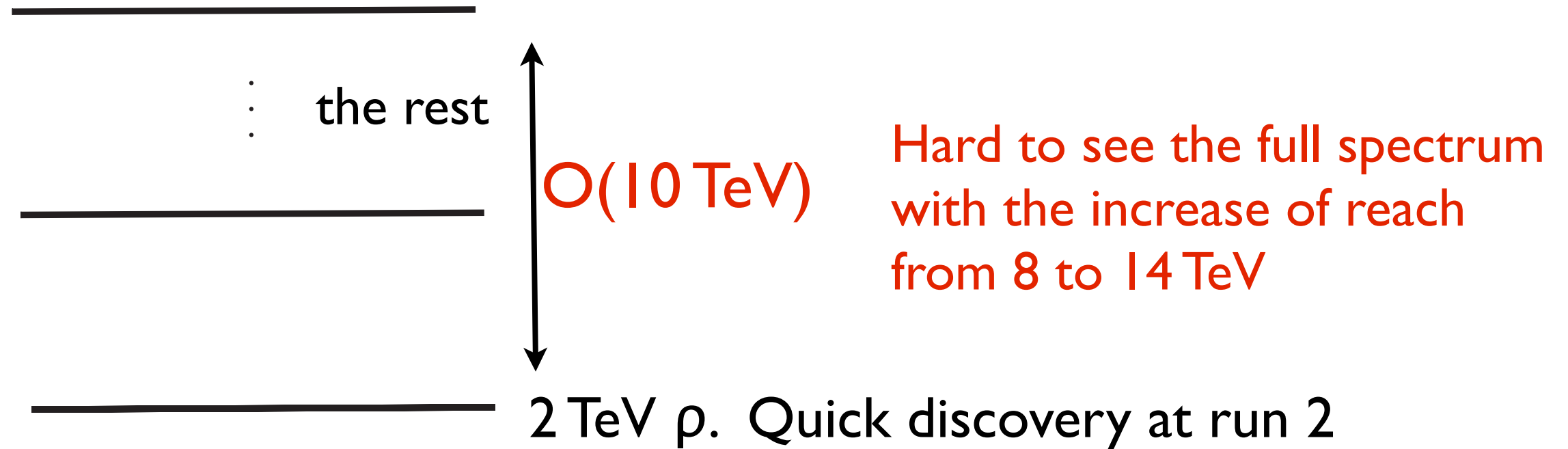


- If it is there, should be able to see it in
 - Wh, Zh
 - $t\bar{t}, t\bar{b}$

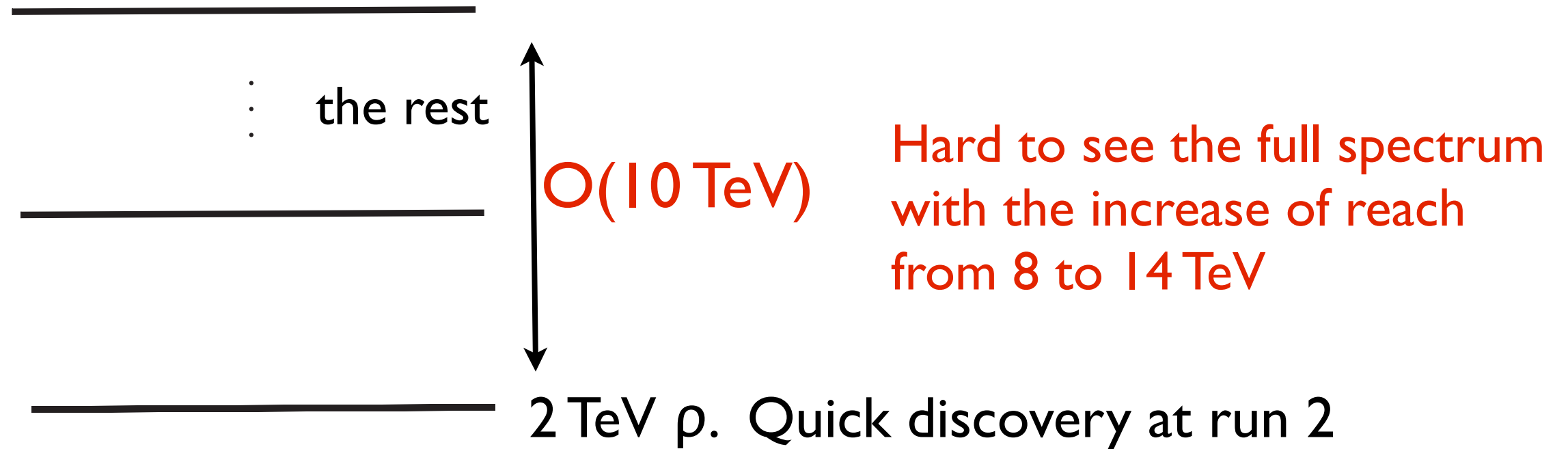
Composite Higgs



Composite Higgs



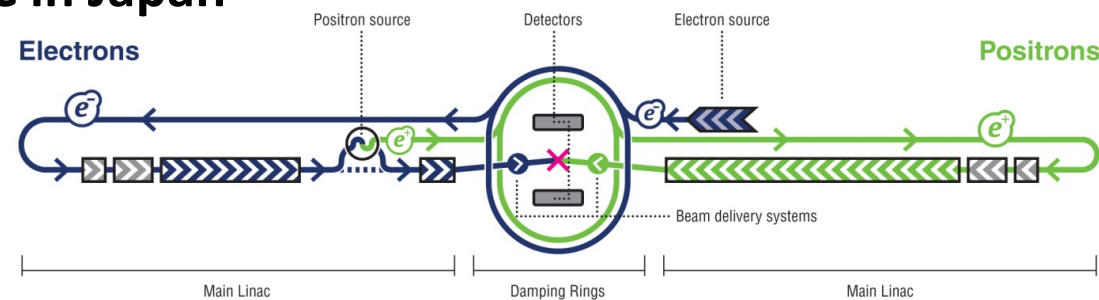
Composite Higgs



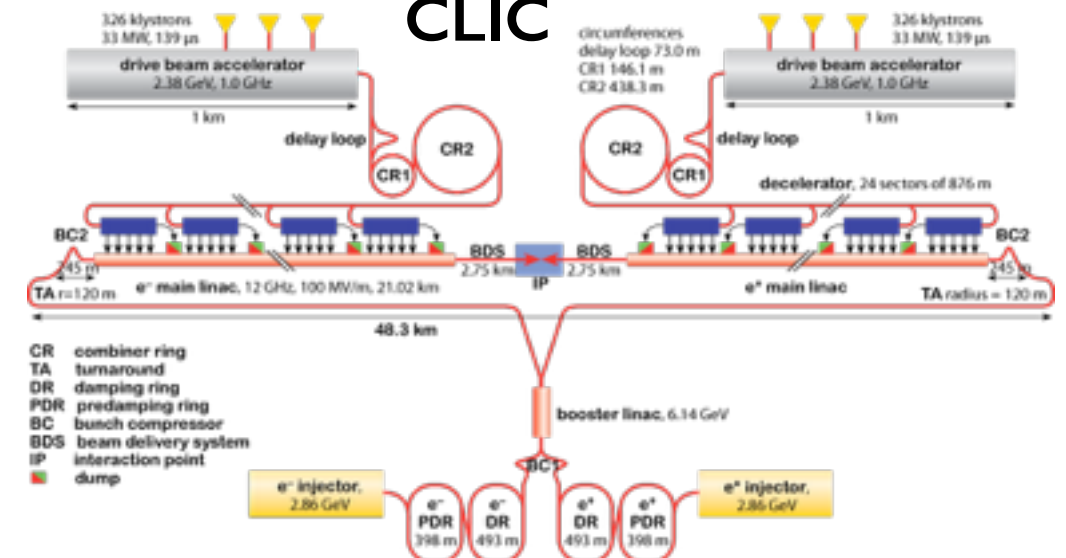
Need to go beyond!

Beyond the LHC, future facilities

ILC in Japan



CLIC



Circular. “Scale up” LEP+LHC

~100 TeV

pp collider

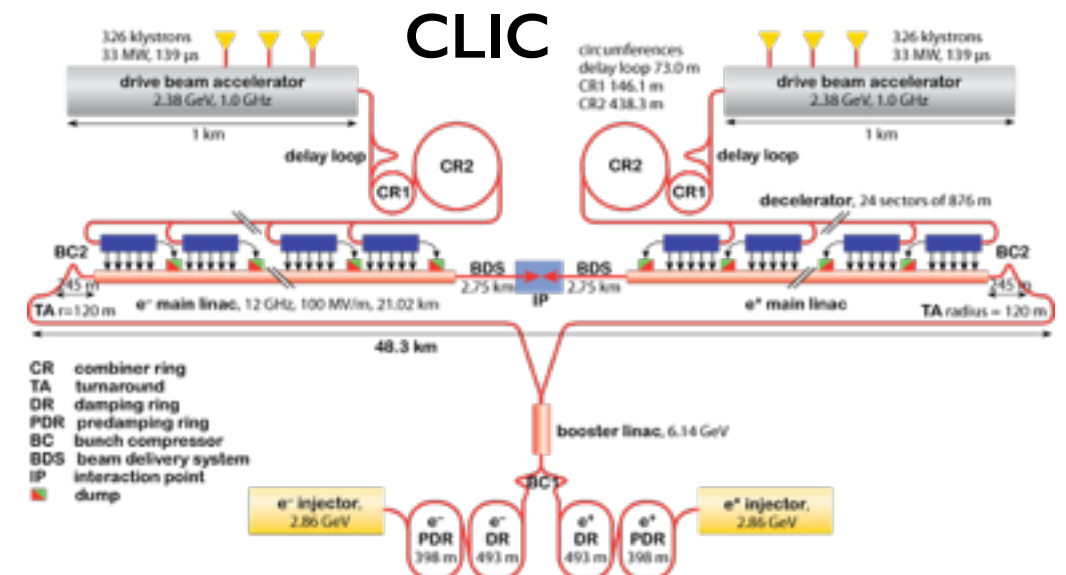
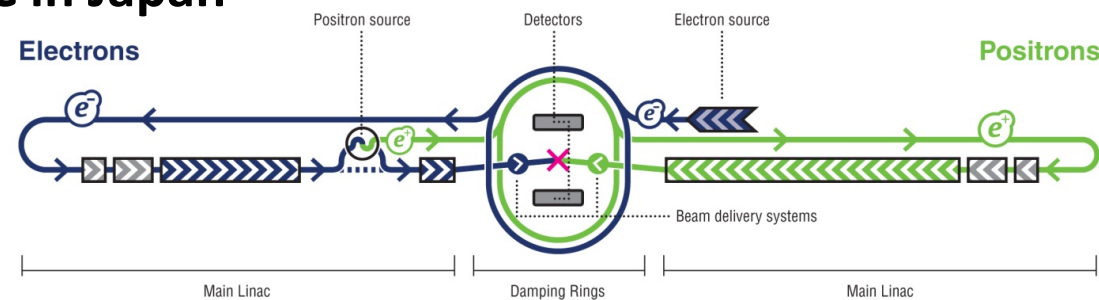
FCC-hh (CERN), SppC(China)

250 GeV **e^-e^+ Higgs Factory**

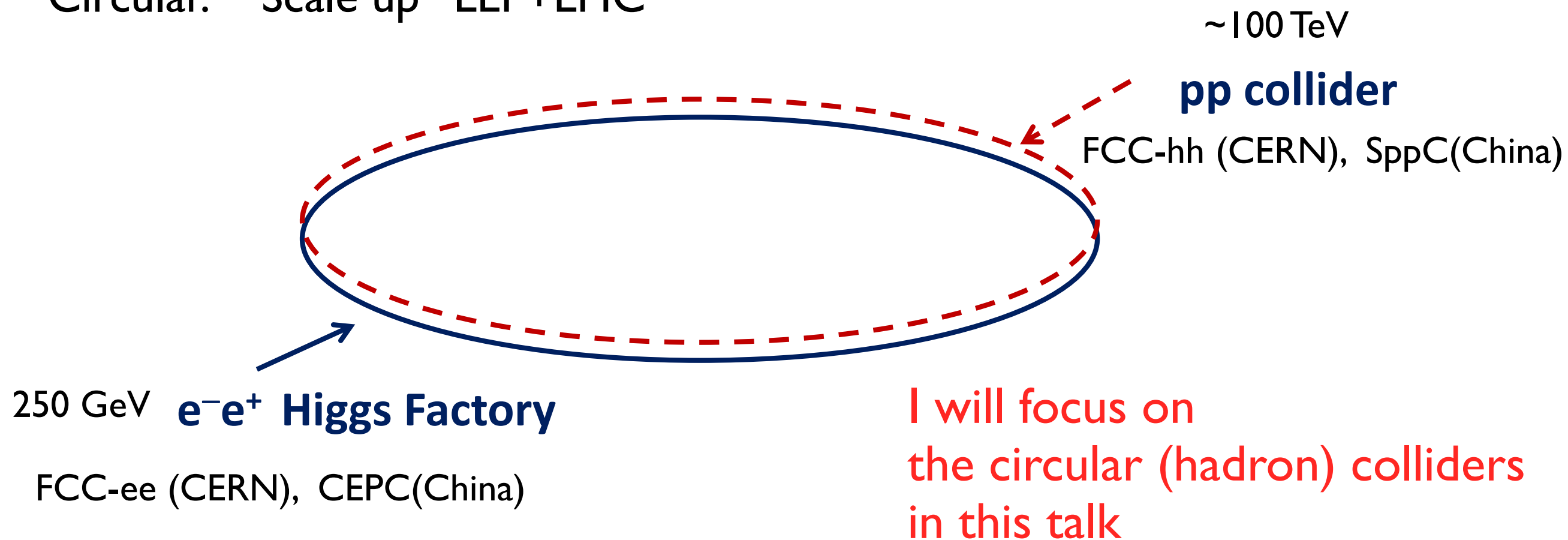
FCC-ee (CERN), CEPC(China)

Beyond the LHC, future facilities

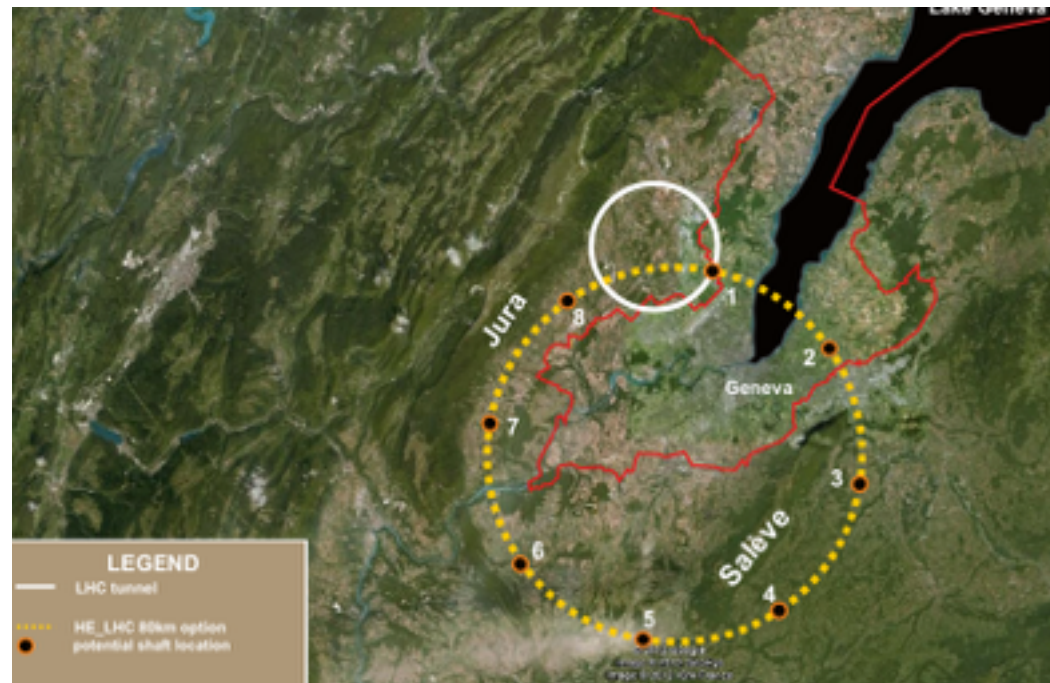
ILC in Japan



Circular. “Scale up” LEP+LHC



Future circular colliders



CERN

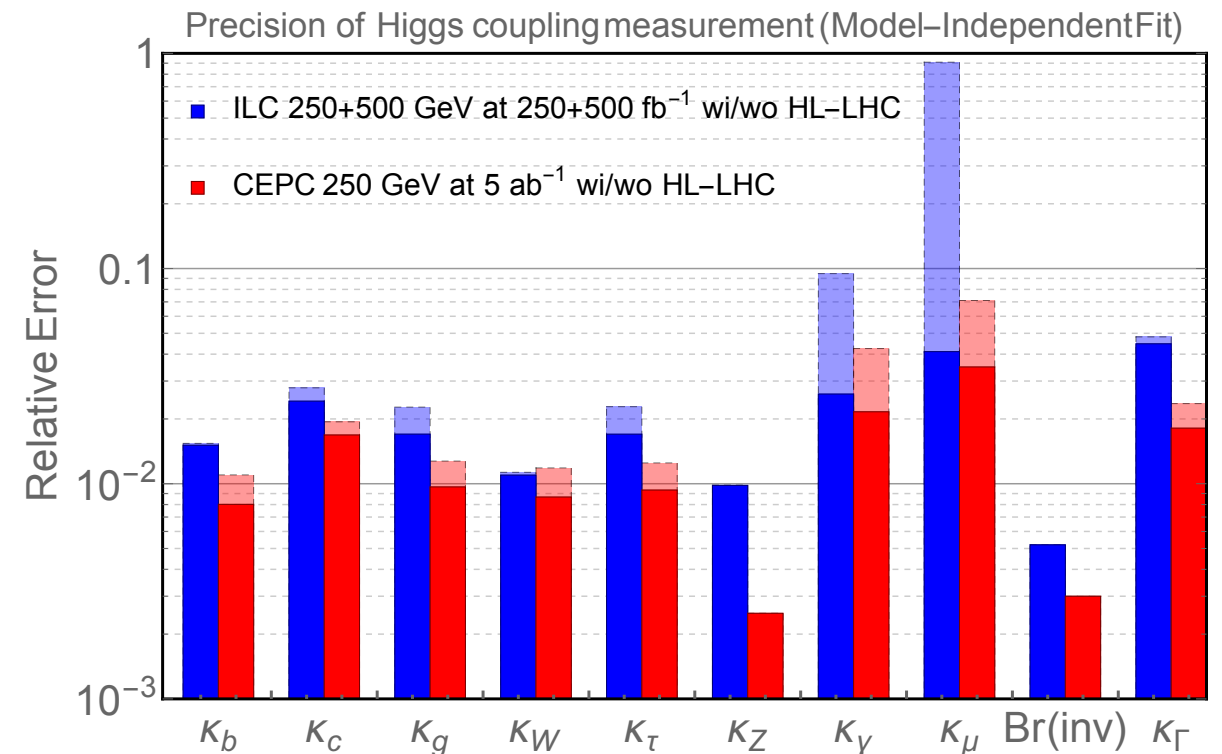
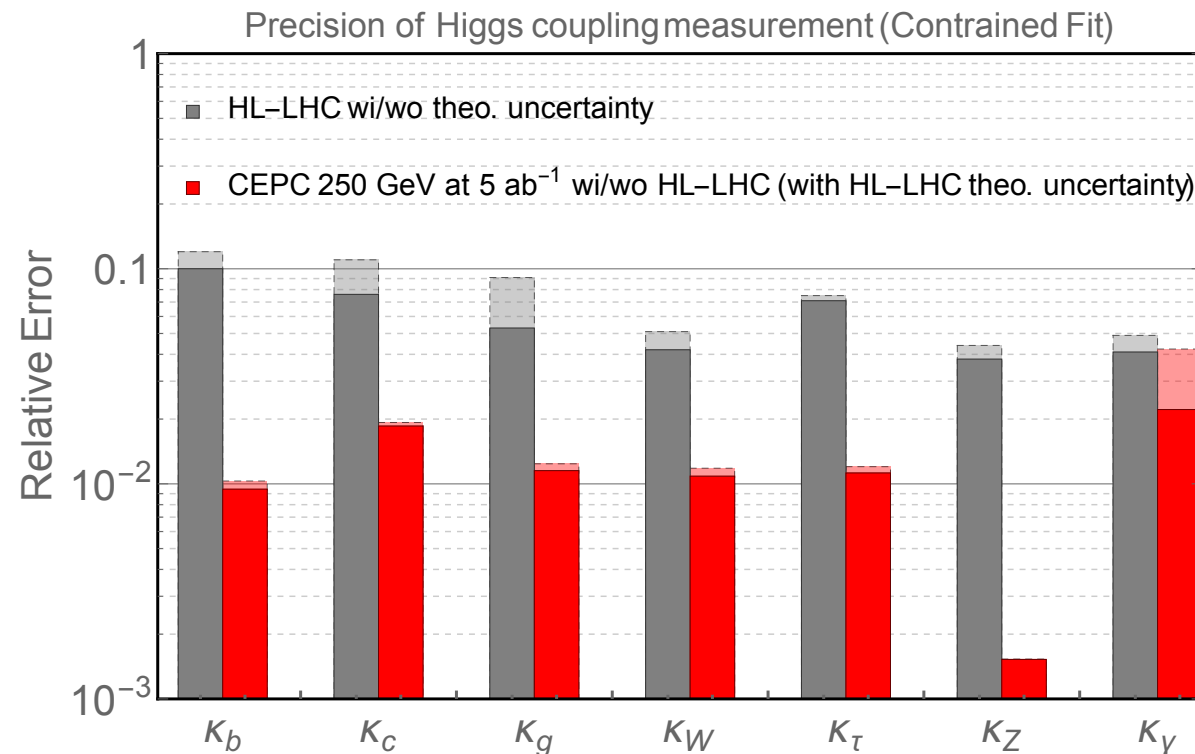
Higgs factory: FCC-ee
pp Collider: FCC-hh



China.

Higgs factory: CEPC
pp Collider: SppC

Higgs coupling at lepton colliders



$$\kappa_X = \frac{\text{Measured Higgs-X coupling}}{\text{Standard Model Higgs-X coupling}}$$

Highlights:

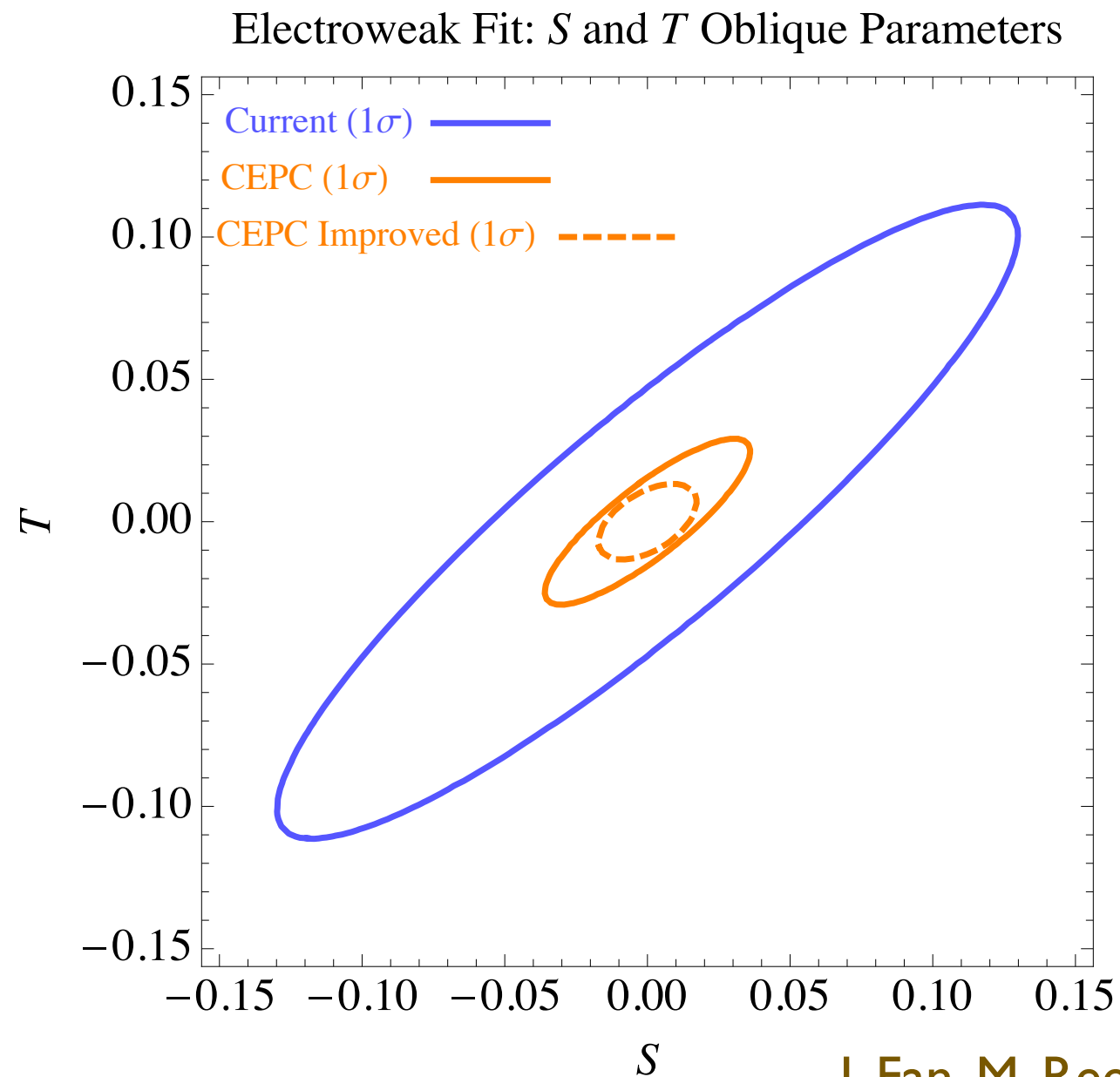
HZ coupling to sub-percent level.

Many couplings to percent level.

Model independent measurement of total width.

Sensitive to the triple Higgs coupling: 20-30%

Electroweak precision at CEPC



J. Fan, M. Reece, LT Wang, I411.1054

- A big step beyond the current precision.

Composite Higgs at lepton collider

Higgs is not (quite) elementary, will have deviations in Higgs couplings.

$$\delta W_h \sim \delta Z_h \sim \frac{v^2}{f^2}$$

Composite resonances couples to W and Z. Will give rise to deviation in EW precision observables.

$$S \simeq \frac{N}{4\pi} \frac{v^2}{f^2}$$

Experiment	S (68%)	f (GeV)
ILC	0.012	1.1 TeV
CEPC (opt.)	0.02	880 GeV
CEPC (imp.)	0.014	1.0 TeV
TLEP-Z	0.013	1.1 TeV
TLEP- t	0.009	1.3 TeV

A clear big step above the LHC.

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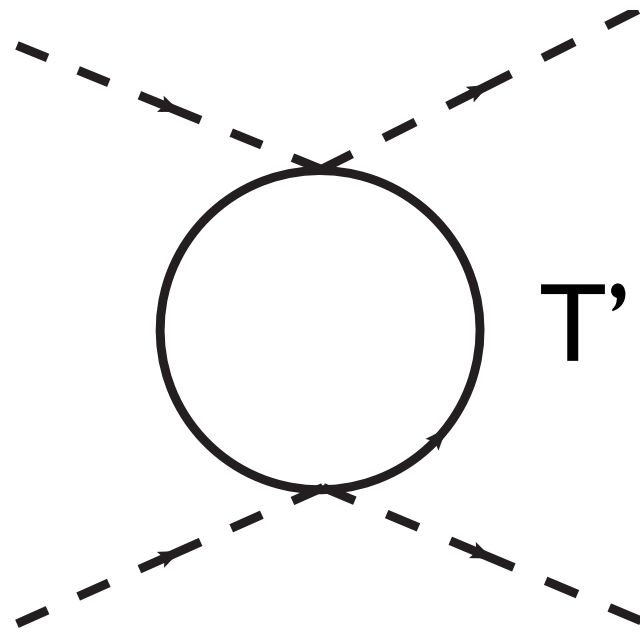
Experiment	κ_Z (68%)	f (GeV)
HL-LHC	3%	1.0 TeV
ILC500	0.3%	3.1 TeV
ILC500-up	0.2%	3.9 TeV
CEPC	0.2%	3.9 TeV
TLEP	0.1%	5.5 TeV

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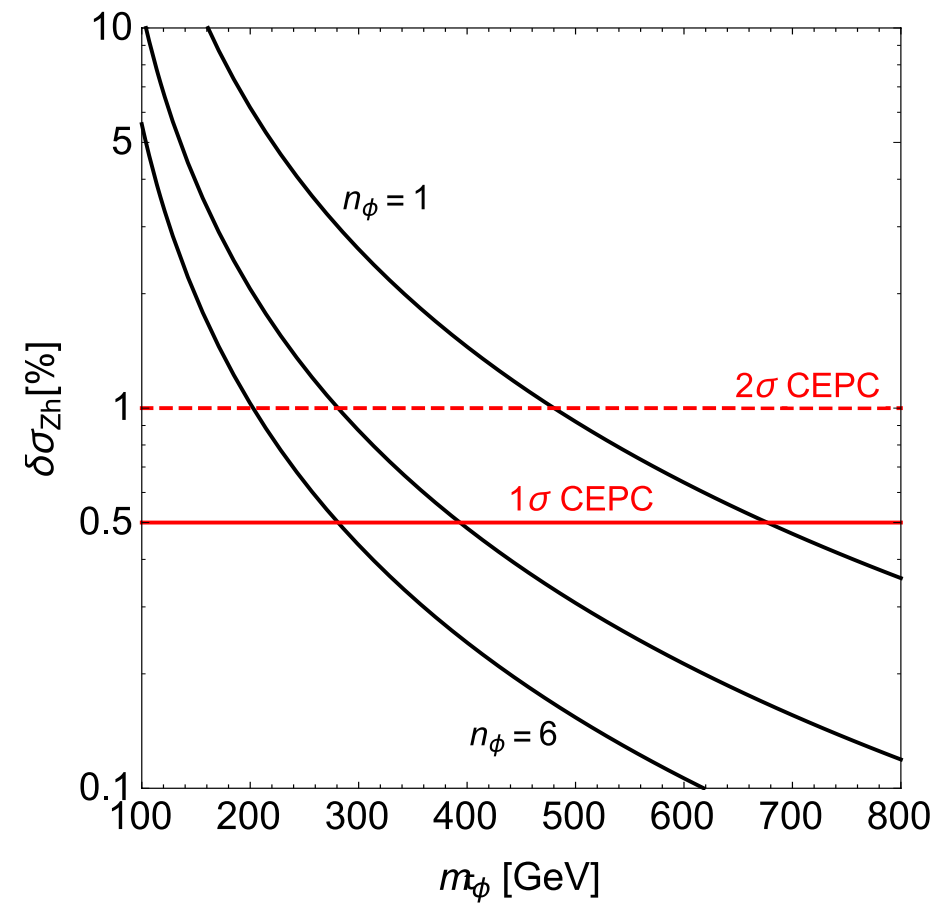
A clear big step above the LHC.

Neutral naturalness

Twin Higgs. Chacko et al. Talk by Craig



Top partner only couple to Higgs.
Wavefunction renormalization
Induce shift in Higgs coupling.

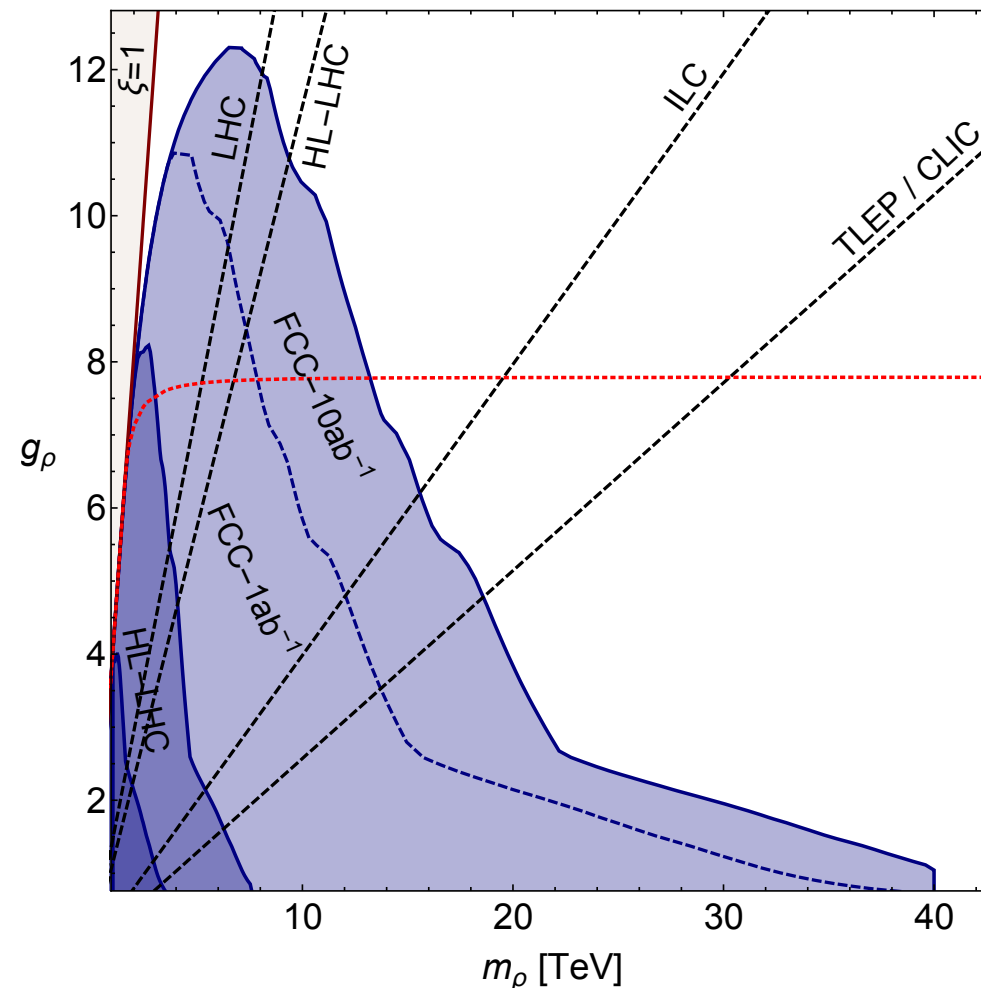


Craig, Englert, McCullough, 2013

- LHC reach poor. Theory can be completely natural.
- Higgs factory can test this.

At 100 TeV collider

Pappadopulo, Thamm, Torre, Wulzer, 2014



- tune proportional to $(m_{\text{NP}})^2$.
 - Much better test than LHC, by orders of magnitude!
 - Potential for discovery (would be a victory for naturalness).

Conclusions.

- Composite Higgs is a plausible solution to the naturalness problem.
- Could give rise to the excess near 2 TeV.
 - ▶ Will confirm or rule out soon at Run 2.
- Twin version can have large mass of colored top partner.
 - ▶ Di-boson will be the leading discovery channel.
- Whether the excess is there or not, spin-1 resonance should be a main target for LHC searches.

Conclusions.

- LHC would not be able to cover full composite Higgs spectrum, since we have not seen anything yet.

At most a couple of lower lying states.

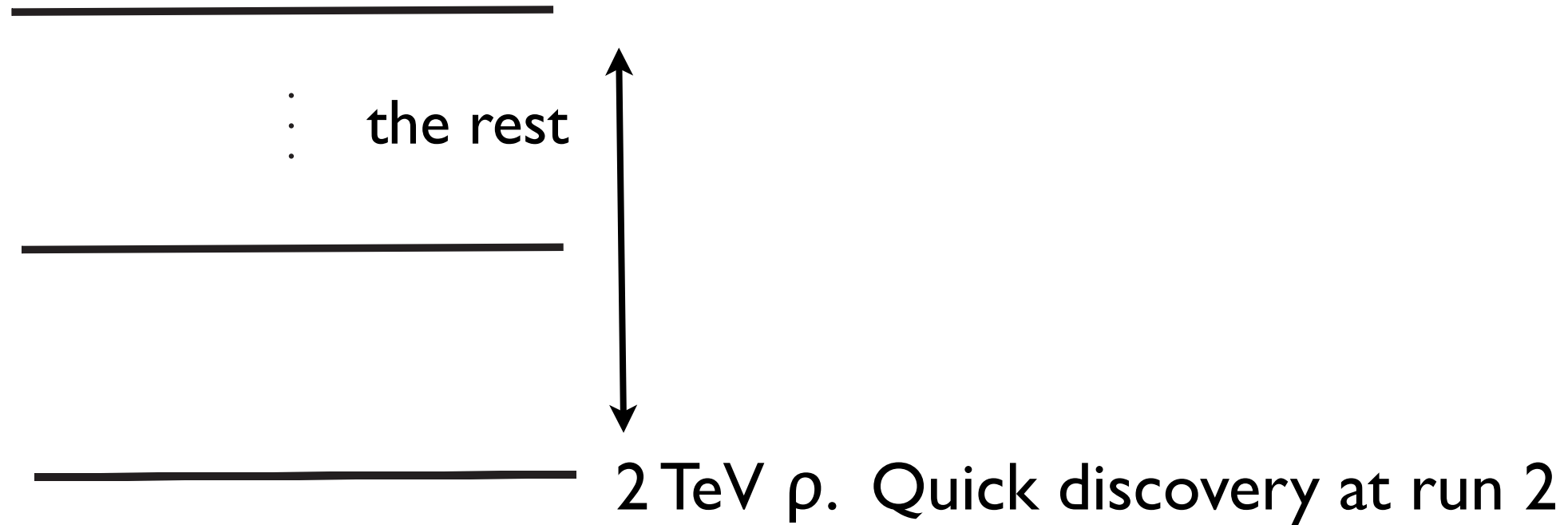
- Need to go beyond.
 - ▶ Higgs factory + 100 TeV pp collider can do a good job.



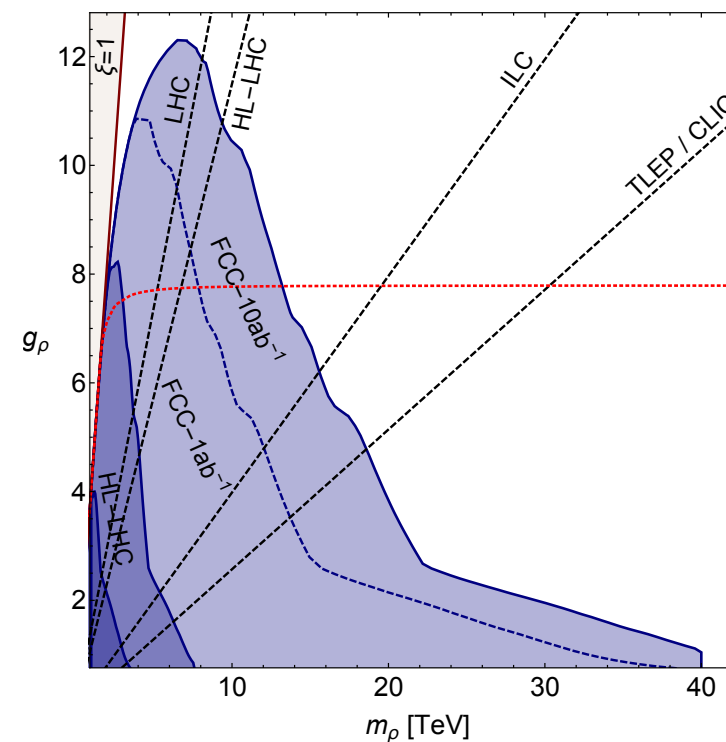
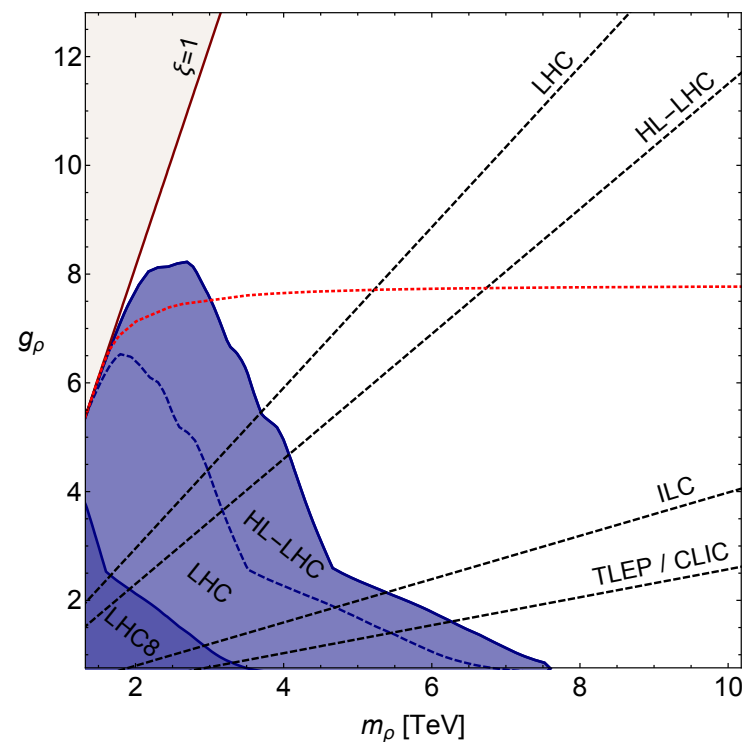
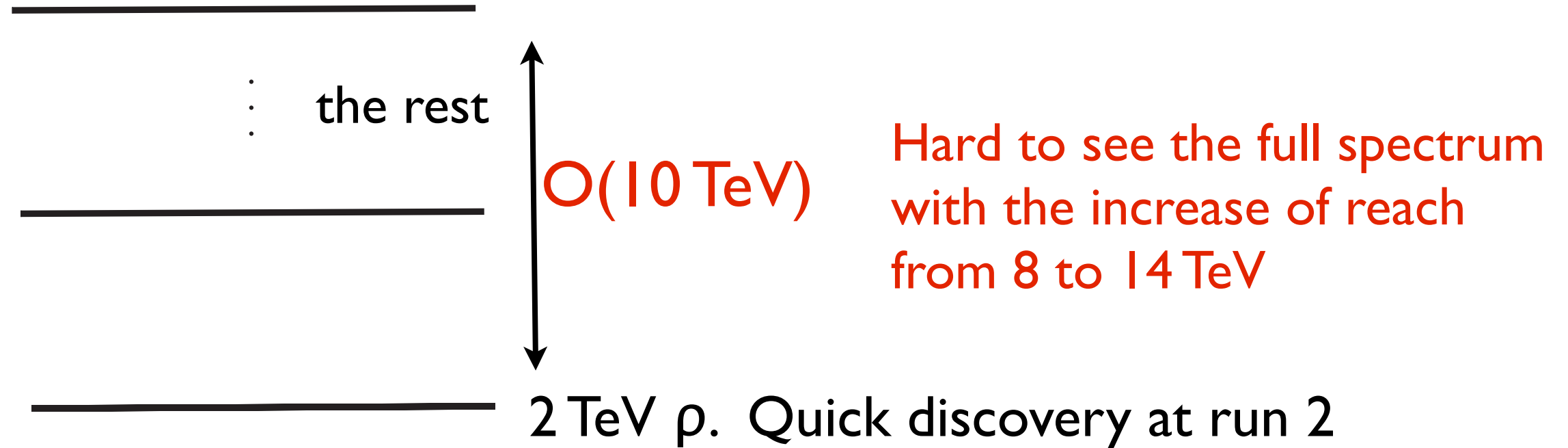
A lot to look forward to!

extras

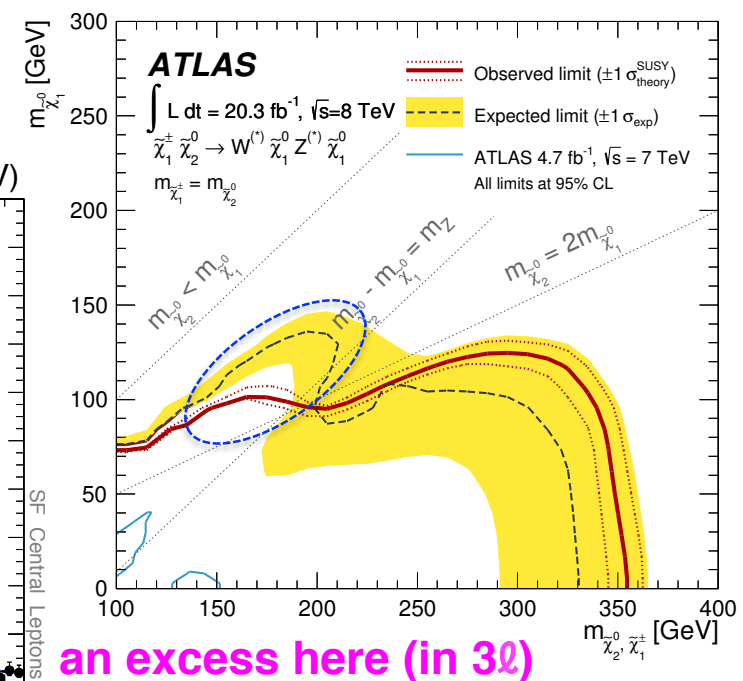
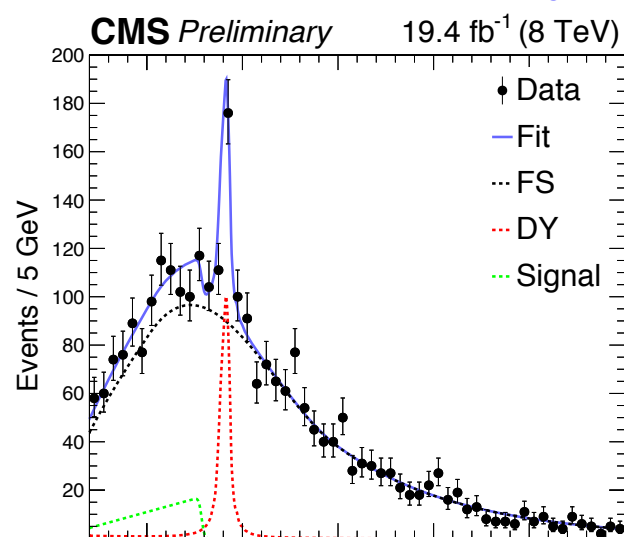
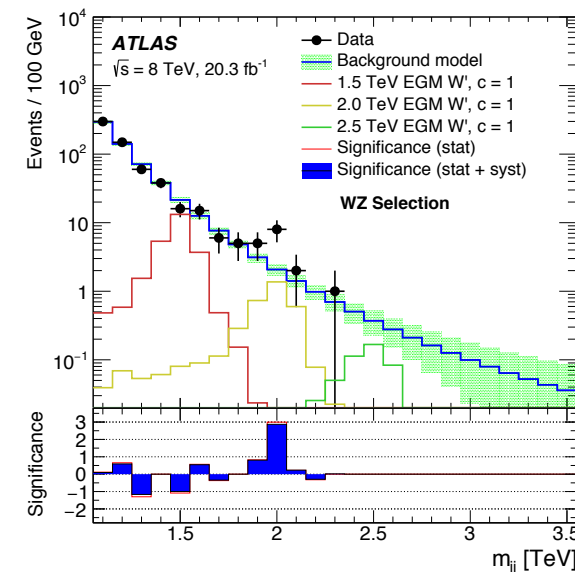
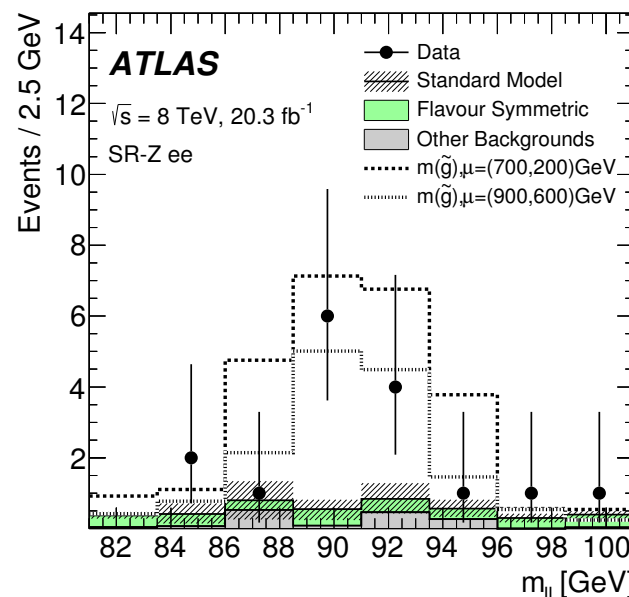
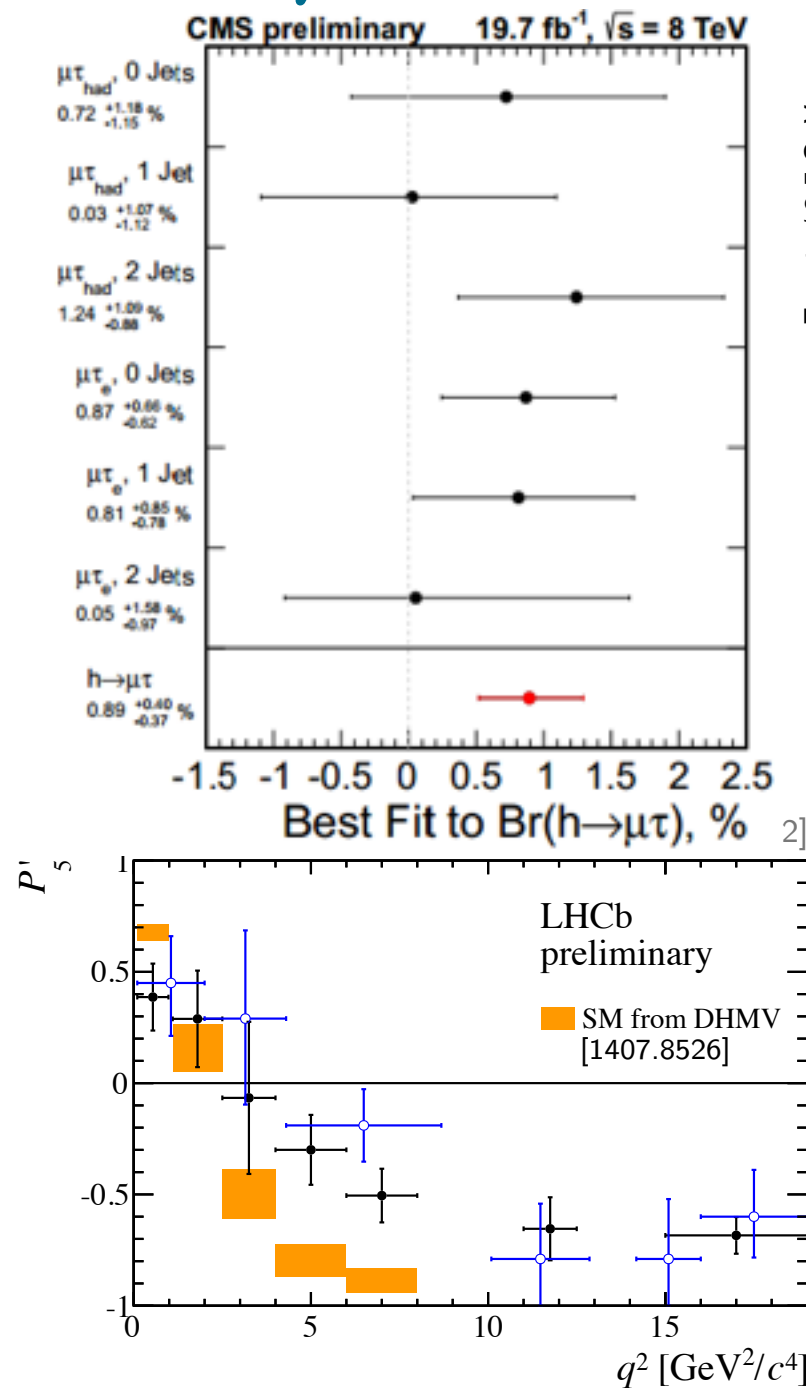
Composite Higgs



Composite Higgs



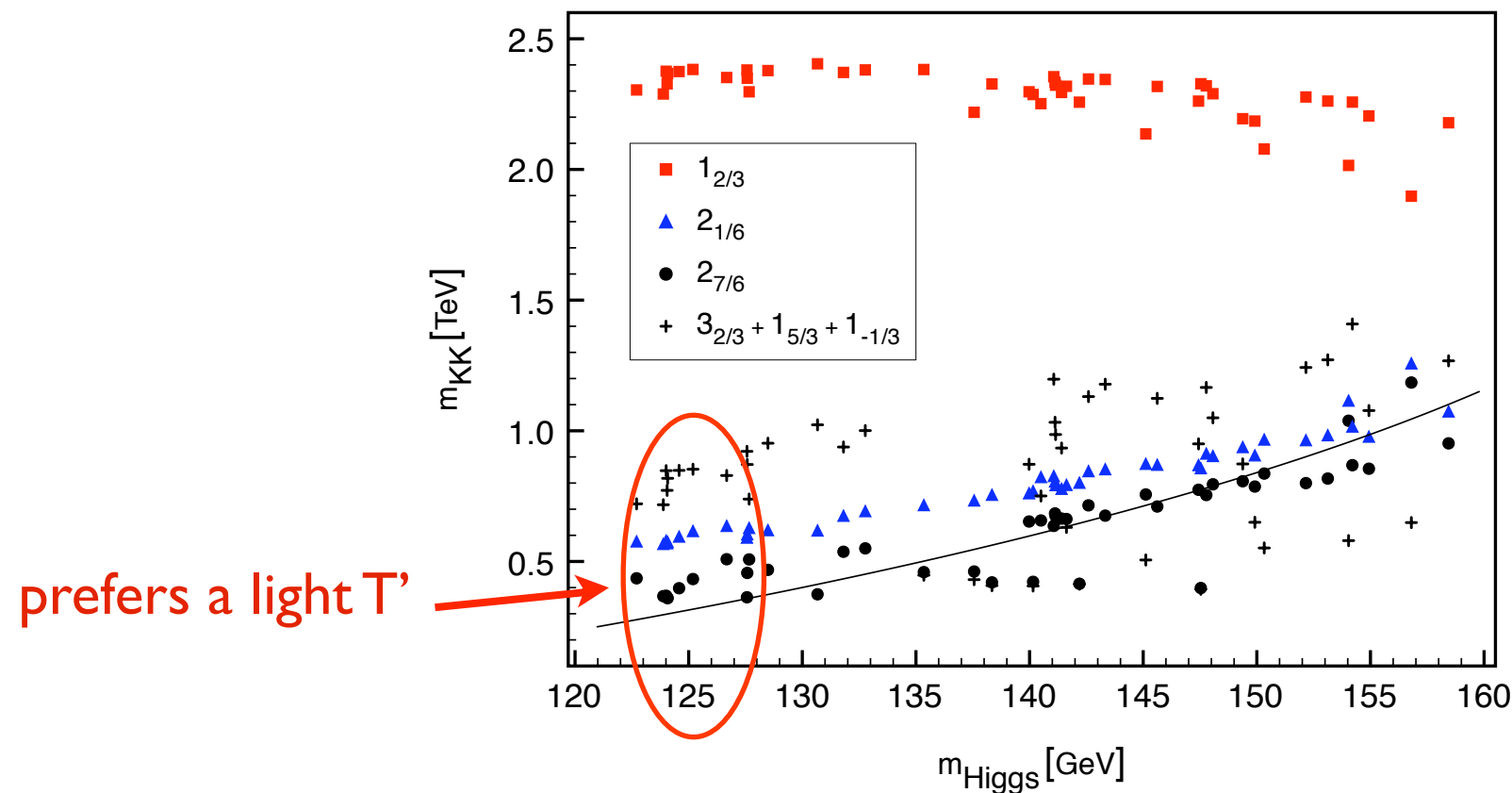
Maybe one of these?



Or something else just around the corner?
Spectacular early discovery!

Compositeness and top partner

Wulzer's talk

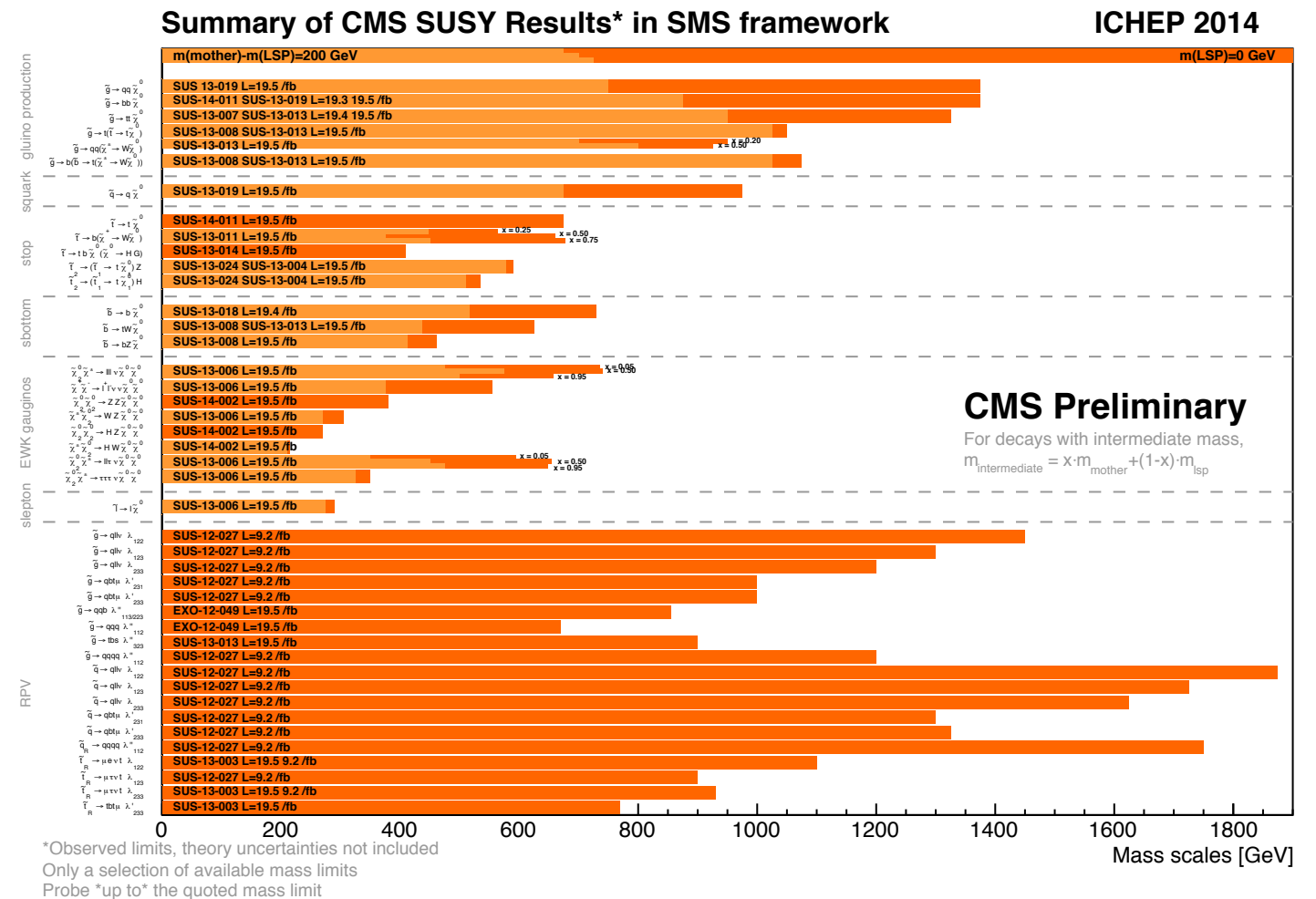
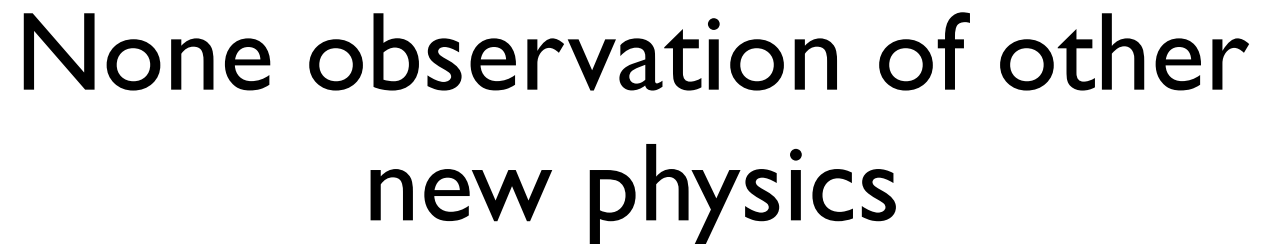


Contino, Da Rold, Pomarol, 2006

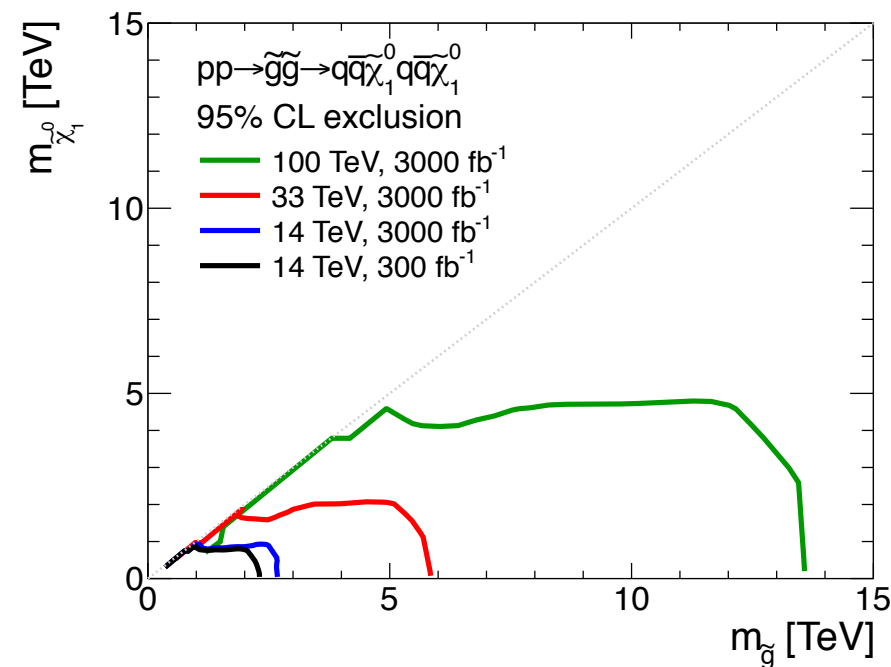
- Plays a crucial role in EWSB.

For a comprehensive discussion, see
De Simone, Matsedonskyi, Rattazzi, Wulzer, 1211.5663

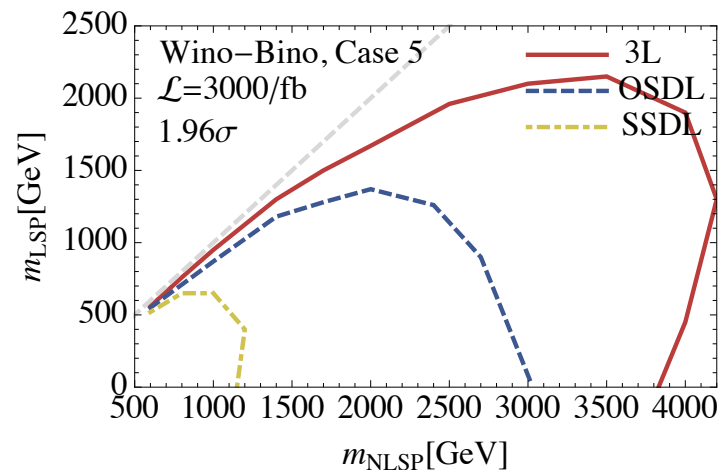
Found Higgs



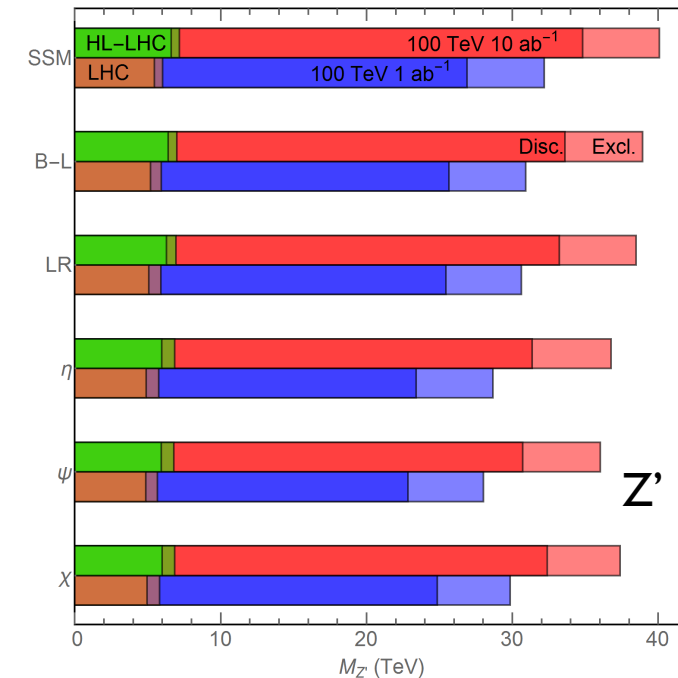
A big step forward in the energy frontier



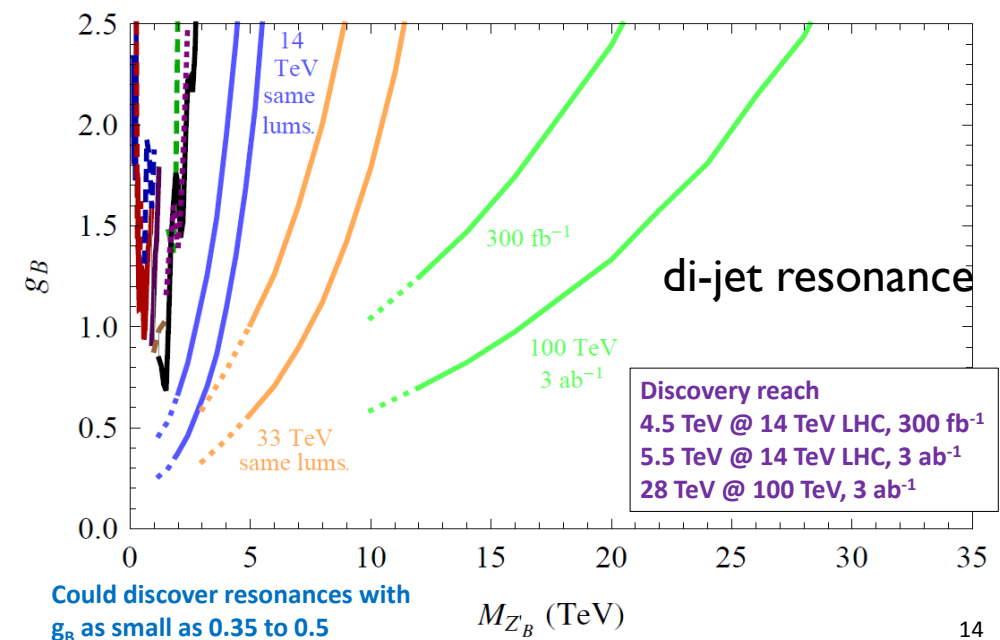
Cohen et al, 2013



Gori, Jung, LTW, Wells, 2014



Han, Langacker, Liu, LTW, to appear



Felix Yu, 2013

cross the board: x 5(more) improvement, into (10)TeV regime