

Radiative seesaw in minimal 3-3-1 model

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The Standard Model is the best theory of describing the nature of particle physics, which is in excellent agreement with almost of all current experiments.

Problems

neutrino mass hierarchy

3 generations

Dark matter

...

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Neutrino mass hierarchy

→ radiative seesaw mechanism

3 generations

→ 3-3-1 model

Neutrino mass hierarchy

→ radiative seesaw mechanism

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→ 3-3-1 model

→ Radiative seesaw in 3-3-1 model

Neutrino hierarchy

Quark sector

$$m_u = 2.3 \times 10^{-3} \text{ GeV}, \quad m_c = 1.3 \text{ GeV}, \quad m_t = 173 \text{ GeV}$$
$$m_d = 4.8 \times 10^{-3} \text{ GeV}, \quad m_s = 9.5 \times 10^{-2} \text{ GeV}, \quad m_b = 4.2 \text{ GeV}$$

Lepton sector

$$m_e = 5.1 \times 10^{-4} \text{ GeV}, \quad m_\mu = 0.11 \text{ GeV}, \quad m_\tau = 1.8 \text{ GeV}$$

$$m_\nu \sim 10^{-11} \text{ GeV}$$

Neutrino masses are very light

Seesaw mechanism

Dimension 5 Operator

$$\frac{f}{\Lambda} h h \nu_L \nu_L$$

$$m_\nu \sim \mathcal{O}(0.1) eV \xrightarrow{\text{blue arrow}} \frac{f}{\Lambda} \sim \mathcal{O}(10^{-14}) GeV^{-1}$$

$$f \sim \mathcal{O}(1) \qquad \Lambda \sim \mathcal{O}(10^{14}) \text{GeV}$$

Radiative seesaw mechanism

$$m_\nu = \frac{1}{(16\pi^2)^n} \frac{fv^2}{\sqrt{2}\Lambda}$$

Neutrino masses are suppressed by loop factor

$$\begin{aligned} \frac{f}{\Lambda} &\sim (16\pi^2)^n \times \mathcal{O}(10^{-14}) GeV^{-1} \\ &\sim \mathcal{O}(10^{2n-14}) GeV^{-1} \end{aligned}$$

Particle contents

	Lepton Fields	Scalar Fields			
	$L_L = (\nu_L, e_L, e_R^c)$	S	η	ρ	χ
$SU(3)_L$	3	6	3	3	3
$U(1)_X$	0	0	0	1	-1

$$\mathcal{L}_Y = y_{\ell_1} \sum_{i,j,k=1-3} \bar{L}_{Li} (L_L)_j^c \eta_k^* \epsilon^{ijk} + y_{\ell_2} \text{Tr}[\bar{L}_L S (L_L)^c] + \text{h.c.}$$

Anomaly

$$L_L^a = \begin{bmatrix} \nu_L^a \\ e_L^a \\ e_R^{c,a} \end{bmatrix} \sim (1, 3, 0)$$

$$Q = \text{diag}(0, -1, 1) + X$$

$$Q_L^i = \begin{bmatrix} u_L^i \\ d_L^i \\ J_L^i \end{bmatrix} \sim (3, 3, 2/3) \quad u_R^i \sim (3, 1, 2/3) \quad J_R^i \sim (3, 1, 5/3)$$

$$d_R^i \sim (3, 1, -1/3)$$

$$Q_L^j = \begin{bmatrix} d_L^j \\ u_L^j \\ J_L^j \end{bmatrix} \sim (3, 3^*, -1/3) \quad u_R^j \sim (3, 1, 2/3) \quad J_R^j \sim (3, 1, -4/3)$$

$$d_R^j \sim (3, 1, -1/3)$$

$$a = 1, 2, 3$$

$$i = 1, j = 2, 3$$



All gauge anomalies cancel

Why sextet?

$$\sum_{i,j,k=1-3} y_{\ell_1}^{ij} \bar{L}_{Li} (L_L)_j^c \eta_k^* \epsilon^{ijk}$$

Mass eigenvalues are 0, -M, M

Why sextet?

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$$\bar{L}_{Li}(L_L)_j^c \sim (1, 3 + 6^*, 0)$$

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$$y_{\ell_2} \text{Tr}[\bar{L}_L S(L_L)^c]$$

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Vacuum Expectation Values

$$\langle S \rangle = \begin{bmatrix} 0 & 0 & 0 \\ 0 & 0 & \frac{v_S}{\sqrt{2}} \\ 0 & \frac{v_S}{\sqrt{2}} & 0 \end{bmatrix}, \quad \langle \eta \rangle = \begin{bmatrix} \frac{v_\eta}{\sqrt{2}} \\ 0 \\ 0 \end{bmatrix}, \quad \langle \rho \rangle = \begin{bmatrix} 0 \\ \frac{v_\rho}{\sqrt{2}} \\ 0 \end{bmatrix}, \quad \langle \chi \rangle = \begin{bmatrix} 0 \\ 0 \\ \frac{v_\chi}{\sqrt{2}} \end{bmatrix},$$

$$v_\chi \gg v = \sqrt{v_S^2 + v_\eta^2 + v_\rho^2}$$

Symmetry breakings

$$SU(3)_c \times SU(3)_L \times U(1)_X$$

 $\langle \chi \rangle$

$$SU(3)_c \times SU(2)_L \times U(1)_Y + W'^{\pm}, Z', U^{\pm\pm}$$

 $\langle \rho \rangle, \langle \eta \rangle, \langle S \rangle$

$$SU(3)_c \times U(1)_{em} + W'^{\pm}, Z', U^{\pm\pm}$$

 $+ W^{\pm}, Z$

Bounds

Q parameter Pisano, Pleitez 1992

$$\frac{M_Z^2}{M_W^2} = \frac{1 + 4t^2}{1 + 3t^2} \quad t = \frac{g_X}{g_{SU(3)_L}}$$

$$t^2 = \frac{s_W^2}{1 - 4s_W^2} \quad \longrightarrow \quad \rho = 1$$

Bounds

$M_{W'} > 230 GeV$ $\mu \rightarrow e\nu_e\bar{\nu}_\mu$ Ng 1994

$M_{Z'} > 2.2 TeV$

LHC

$v_\chi > 3.6 TeV$

Coutinho,
Guimares,
Nepomuceno
2013

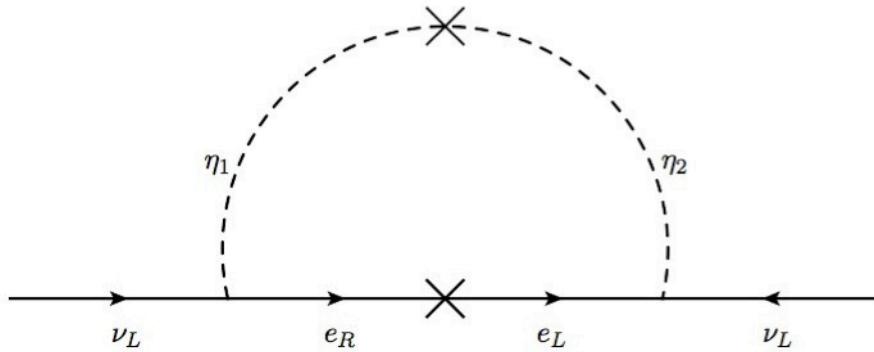
Neutrino oscillation

LFV($\mu \rightarrow e\gamma, \dots$)

Muon g-2



Neutrino oscillation

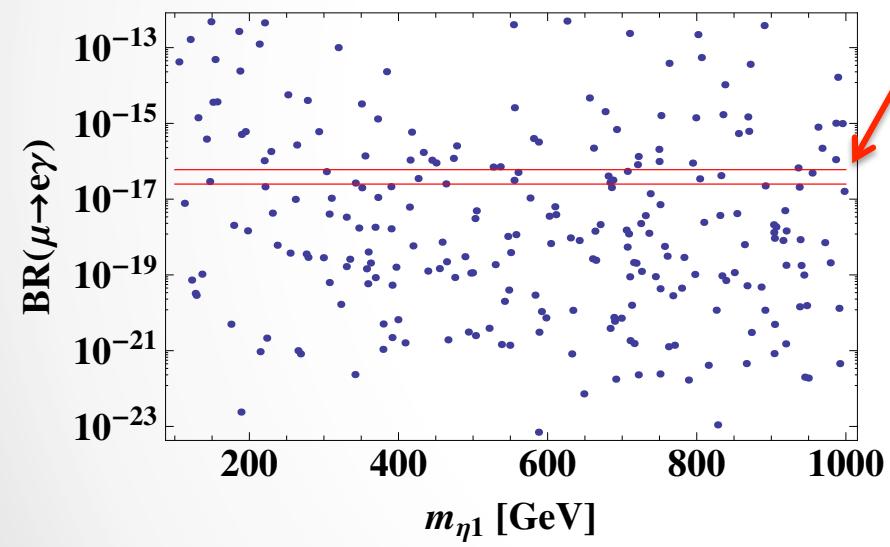


$$-(\mathcal{M}_\nu^{\text{th}})_{ab} \approx \frac{\left[y_{\ell_1} \left(m_\ell^\dagger + m_\ell^* \right) y_{\ell_1}^T \right]_{ab}}{(4\pi)^2} \left[\frac{\delta m_{\eta_1 \eta_2}^{+2}}{m_{\eta_1^+}^2 - m_{\eta_2^+}^2} \right] \ln \left[\frac{m_{\eta_1^+}^2}{m_{\eta_2^+}^2} \right] - \frac{\left[y_{\ell_2} \left(m_\ell^\dagger + m_\ell^* \right) y_{\ell_2}^T \right]_{ab}}{(4\pi)^2} \left[\frac{\delta m_{h_1 \eta_2}^{+2}}{m_{h_1^+}^2 - m_{h_2^+}^2} \right] \ln \left[\frac{m_{h_1^+}^2}{m_{h_2^+}^2} \right]$$

$$- \frac{\left(y_{\ell_1} m_\ell^* y_{\ell_2}^T + y_{\ell_2} m_\ell^\dagger y_{\ell_1}^T \right)_{ab}}{2(4\pi)^2} \left[\frac{\delta m_{\eta_2 h_1}^{+2}}{m_{\eta_2^+}^2 - m_{h_1^+}^2} \right] \ln \left[\frac{m_{\eta_2^+}^2}{m_{h_1^+}^2} \right] - \frac{\left(y_{\ell_2} m_\ell^* y_{\ell_1}^T + y_{\ell_1} m_\ell^\dagger y_{\ell_2}^T \right)_{ab}}{2(4\pi)^2} \left[\frac{\delta m_{\eta_1 h_2}^{+2}}{m_{\eta_1^+}^2 - m_{h_2^+}^2} \right] \ln \left[\frac{m_{\eta_1^+}^2}{m_{h_2^+}^2} \right].$$

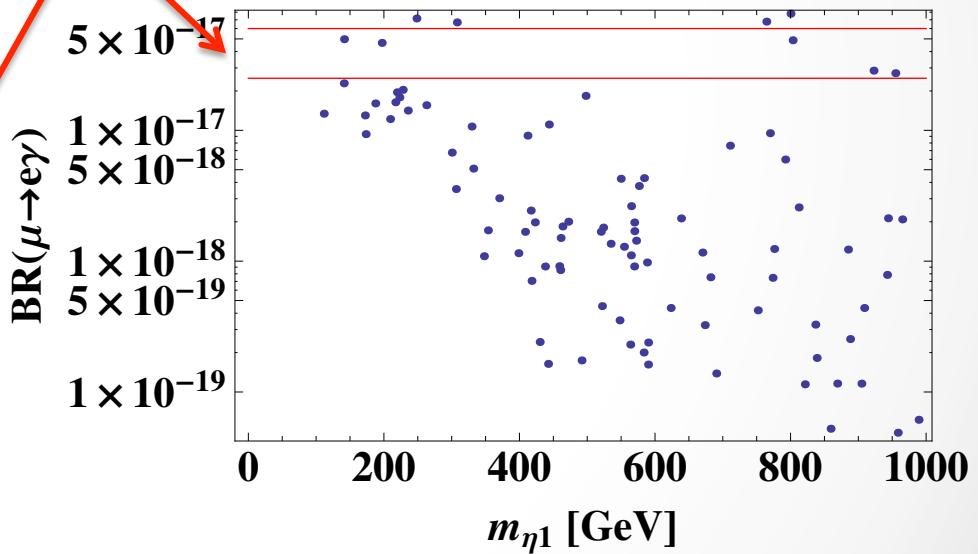
$\mu \rightarrow e\gamma$

NH



IH

Mu2e



Radiative seesaw in 3-3-1 model

neutrino mass hierarchy
3 generations

Buckup

• • •

3-3-1 model

gauge group

$$SU(3)_c \times SU(3)_L \times U(1)_X$$

$$\xrightarrow{\hspace{1cm}} SU(3)_c \times SU(2)_L \times U(1)_Y$$

$$\xrightarrow{\hspace{1cm}} SU(3)_c \times U(1)_{em}$$

$$Q = T_3 - \sqrt{3}T_8 + X$$

$$T_3 = \frac{1}{2} diag(1, -1, 0) \quad \quad T_8 = \frac{1}{2\sqrt{3}} diag(1, 1, -2)$$

Parameter regions

NH

$$1 \text{ GeV} \leq (v_\eta, v_\sigma) \leq 100 \text{ GeV}, \quad 0.1 \text{ GeV}^2 \leq (\delta m_{\eta_1 h_1}^{+2}, \delta m_{\eta_2 h_1}^{+2}) \leq 10 \text{ GeV}^2,$$
$$100 \text{ GeV} \leq m_{\eta_{1,2}^+} \leq 1000 \text{ GeV}, \quad -1 \leq (s_{eL(R)ij}, c_{eL(R)ij}) \leq 1,$$

IH

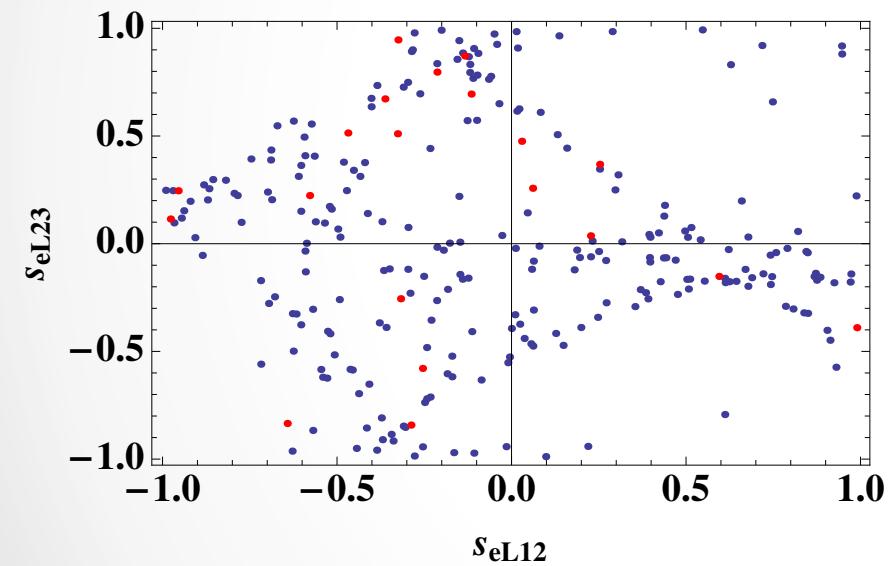
$$33 \text{ GeV} \leq v_\eta \leq 37 \text{ GeV}, \quad 44 \text{ GeV} \leq v_\sigma \leq 48 \text{ GeV}, \quad 0.1 \text{ GeV}^2 \leq (\delta m_{\eta_1 h_1}^{+2}, \delta m_{\eta_2 h_1}^{+2}) \leq 1 \text{ GeV}^2,$$
$$100 \text{ GeV} \leq m_{\eta_{1,2}^+} \leq 1000 \text{ GeV}, \quad -1 \leq s_{eL12} \leq 0, \quad 0 \leq s_{eL23} \leq 0.5, \quad -1 \leq s_{eL13} \leq -0.7,$$
$$-0.4 \leq s_{eR12} \leq -0.1, \quad 0 \leq s_{eR23} \leq 0.2, \quad 0.4 \leq s_{eR13} \leq 0.7,$$

Our parameter

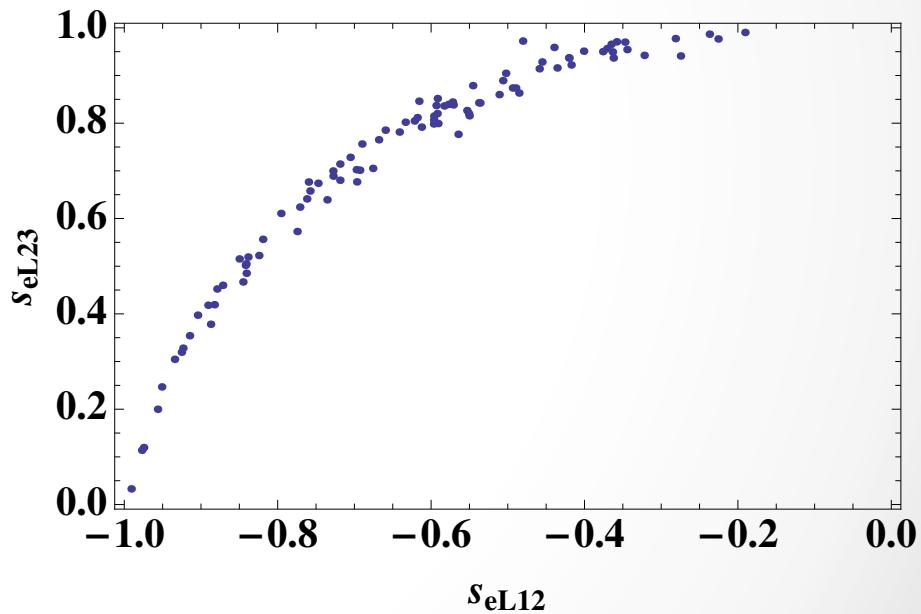
$$(\bar{e}_L)_a (m_\ell)_{ab} (e_R)_b \equiv (\bar{e}_L)_a \left[\frac{(y_{\ell_1})v_\eta}{\sqrt{2}} + \frac{(y_{\ell_2})v_\sigma}{\sqrt{2}} \right]_{ab} (e_R)_b = (\bar{e}_L)_a (V_{eL}^\dagger)_{ai} (m_\ell^{diag})_i (V_{eR})_{ib} (e_R)_b$$

$$V_{eL(R)} \equiv \begin{bmatrix} 1 & 0 & 0 \\ 0 & c_{eL(R)23} & s_{eL(R)23} \\ 0 & -s_{eL(R)23} & c_{eL(R)23} \end{bmatrix} \begin{bmatrix} c_{eL(R)13} & 0 & s_{eL(R)13} \\ 0 & 1 & 0 \\ -s_{eL(R)13} & 0 & c_{13}^{eL(R)} \end{bmatrix} \begin{bmatrix} c_{eL(R)12} & s_{eL(R)12} & 0 \\ -s_{eL(R)12} & c_{eL(R)12} & 0 \\ 0 & 0 & 1 \end{bmatrix}$$

NH



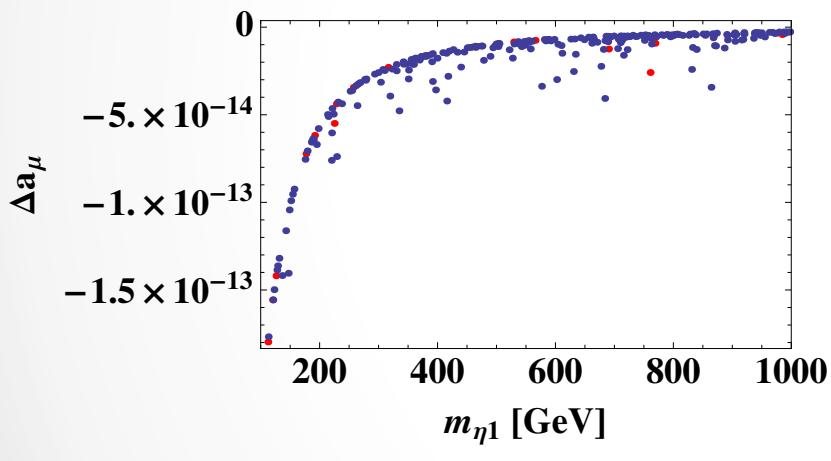
IH



The red points do not satisfy the LFV constraints

Muon g-2

NH



IH

