#### Exotic LFV events at the LHC PRD91(2015)015001 – 1508.04623 – 1508.05074

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## **Motivations**



- Different mixing pattern from CKM,  $\nu$  lightness  $\stackrel{?}{\leftarrow}$  Majorana  $\nu$
- SM: no 
  *ν* mass term, lepton flavour is conserved
   ⇒ need new Physics
  - Radiative models
  - Extra dimensions
  - R-parity violation in supersymmetry
  - Seesaw mechanism  $\rightarrow \nu$  mass at tree-level
    - + BAU through leptogenesis
- Low-scale seesaw mechanisms
  - $\rightarrow$  no naturalness problem from heavy neutrinos



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### The inverse seesaw mechanism

• Inverse seesaw  $\Rightarrow$  Consider fermionic gauge singlets  $\nu_{Ri}$ (L = +1) and  $X_i$  (L = -1) [Mohapatra and Valle, 1986]

$$\mathcal{L}_{inverse} = -Y_{\nu}^{ij}\overline{L_{i}}\tilde{H}\nu_{Rj} - M_{R}^{ij}\overline{\nu_{Ri}^{C}}X_{j} - \frac{1}{2}\mu_{X}^{ij}\overline{X_{i}^{C}}X_{j} + \text{h.c.}$$

with 
$$m_D = Y_{\nu} v, M^{\nu} = \begin{pmatrix} 0 & m_D & 0 \\ m_D^T & 0 & M_R \\ 0 & M_R^T & \mu_X \end{pmatrix}$$



$$m_{\nu} \approx \frac{1}{m_D^2 + M_R^2}$$
  
 $m_{N_1,N_2} \approx \mp \sqrt{m_D^2 + M_R^2} + \frac{M_R^2 \mu_X}{2(m_D^2 + M_R^2)}$ 

 $m_D^2 \mu_X$ 

2 scales:  $\mu_X$  and  $M_R$ 

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# A rich phenomenology

• Inverse seesaw:  $Y_{\nu} \sim \mathcal{O}(1)$  and  $M_R \sim 1 \text{ TeV}$  $\Rightarrow$  testable at the LHC and low energy experiments

#### LHC/LC signatures

- single lepton + dijet + missing energy [Das and Okada, 2013]
- di-lepton + missing p<sub>T</sub> [Bhupal Dev et al., 2012, Bandyopadhyay et al., 2013]
- LFV di-lepton + dijet [1508.05074]
- tri-lepton + missing E<sub>T</sub> [Mondal et al., 2012, Das et al., 2014]...
- invisible Higgs decays [Banerjee et al., 2013]
- Low energy:
  - deviations from lepton universality [Abada, Teixeira, Vicente and CW, 2014]
  - charged lepton flavour violation (LFV) [Bernabéu et al., 1987]...
  - neutrinoless double beta decay [Awasthi et al., 2013]...
  - charged lepton anomalous magnetic moment [Abada et al., 2014]



## An aside on LFV Higgs decays

•  $h \rightarrow \tau \mu$ : 2.4 $\sigma$  excess in CMS, Br = 0.84<sup>+0.39</sup><sub>-0.37</sub>% [Khachatryan et al., 2015] 1.3 $\sigma$  excess in ATLAS, Br = 0.77 ± 0.62% [1508.03372]

• Very stringent constraints from other LFV processes, e.g.

$$BR(\mu \to e\gamma) \leq 5.7 \times 10^{-13} \text{ [MEG, 2013]}$$
$$BR(\tau \to e\gamma) \leq 3.3 \times 10^{-8} \text{ [BaBar, 2010]}$$
$$BR(\tau \to \mu\gamma) \leq 4.4 \times 10^{-8} \text{ [BaBar, 2010]}$$

• Approximate formulas for the ISS with large  $Y_{\nu}$ 

[Arganda, Herrero, Marcano and CW, PRD91(2015)015001]

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$$\begin{aligned} &\mathsf{BR}_{l_m \to l_k \gamma} \simeq 8 \times 10^{-17} \frac{m_{l_m}^5}{\Gamma_{l_m}} \left| \frac{v^2}{2M_R^2} (Y_\nu Y_\nu^\dagger)_{km} \right|^2 \\ &\mathsf{BR}_{h \to \mu \bar{\tau}} \simeq 10^{-7} \frac{v^4}{M_R^4} \left| (Y_\nu Y_\nu^\dagger)_{23} - 5.7 (Y_\nu Y_\nu^\dagger Y_\nu Y_\nu^\dagger)_{23} \right|^2 \end{aligned}$$



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## Producing large $h \rightarrow \tau \mu$ rates

• Textures with 
$$(Y_{\nu}Y_{\nu}^{\dagger})_{12}=0$$
 and  $\frac{|Y_{\nu}^{ij}|^2}{4\pi}<1.5$ 

$$Y_{\tau\mu}^{(1)} = f \begin{pmatrix} 0 & 1 & -1 \\ 0.9 & 1 & 1 \\ 1 & 1 & 1 \end{pmatrix}, \ Y_{\tau\mu}^{(2)} = f \begin{pmatrix} 0 & 1 & 1 \\ 1 & 1 & -1 \\ -1 & 1 & -1 \end{pmatrix}, \ Y_{\tau\mu}^{(3)} = f \begin{pmatrix} 0 & -1 & 1 \\ -1 & 1 & 1 \\ 0.8 & 0.5 & 0.5 \end{pmatrix}$$

#### Flavour composition of the heavy neutrinos:



- 3 very different flavour patterns
- Heavy neutrino mixing of  $\tau \mu$  type is always present



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# $h \rightarrow \tau \mu$ in the (SUSY) ISS

•  $f = \sqrt{6\pi}$ ,  $M_R$  real, degenerate,  $\nu$  oscillations reproduced by  $\mu_X$ 



[PRD91(2015)015001]

- Dotted: excluded by  $\tau \rightarrow \mu \gamma$ Solid: allowed by LFV, LUV...
- Br<sup>max</sup>  $(h \rightarrow \mu \bar{\tau}) \sim 10^{-5}$
- Same with hierarchical heavy neutrinos



[arXiv:1508.04623]

- $\times$ : excluded by  $\tau \rightarrow \mu \gamma$ 
  - ▲: allowed
- $\operatorname{Br}^{\max}(h \to \mu \overline{\tau}) = \mathcal{O}(1\%)$
- SUSY loops could explain the TE CMS and ATLAS excess



# Heavy neutrinos production and decays at the LHC

[1508.05074]



- Main production channel: Drell-Yan
- $\tau \mu$  mixing in N  $\Rightarrow \mu^{\pm} \tau^{\mp} j j$  signal with no  $\not\!\!\!E_T$
- Wγ fusion relevant at large M<sub>R</sub> [Dev et al., 2014, Alva et al., 2015]
- Contribute to μ<sup>±</sup>τ<sup>∓</sup>jj signal if extra-jets are soft or collinear → p<sub>T</sub> < p<sub>T</sub><sup>max</sup>
- Numerics done with MadGraph5 and NNPDFQED, M<sub>R</sub> real and degenerate
- Similar results with other PDF sets

### Production and decays at LHC14



# $pp \rightarrow \mu \tau j j$ events at LHC14 for $Y_{\tau \mu}^{(3)}$



- Lower line: production only from Drell-Yan Shaded regions:  $W\gamma$  fusion added with  $p_T^{\text{max}} = 10, 20, 40 \text{ GeV}$ (darker to lighter)
- Up to  $\mathcal{O}(100)$  events, naively background free



# $pp \rightarrow \mu \tau jj$ events at LHC14 for $Y_{\tau \mu}^{(1)}$





- Lower line: production only from Drell-Yan Shaded regions:  $W\gamma$  fusion added with  $p_T^{\text{max}} = 10, 20, 40 \text{ GeV}$  (darker to lighter)
- Up to  $\mathcal{O}(10)$  events, naively background free



# $pp \rightarrow \mu \tau j j$ events at LHC14 for $Y_{\tau \mu}^{(2)}$





- Lower line: production only from Drell-Yan Shaded regions:  $W\gamma$  fusion added with  $p_T^{\text{max}} = 10, 20, 40 \text{ GeV}$  (darker to lighter)
- Up to  $\mathcal{O}(200)$  events, naively background free



#### Conclusions

- (SUSY) Inverse seesaw: specific examples of low-scale seesaw mechanisms
- New physics at the TeV scale with large couplings: rich phenomenology
- SUSY loops could explain the CMS and ATLAS excess in  $h \rightarrow \tau \mu$
- 10-200 events would be expected at LHC14
- Next step: Consider  $e\tau jj$  and  $e\mu jj$  final states
  - Estimate the signal/background ratio via detector level studies



# **Backup slides**



# Diagrams for the ISS

#### (PRD91(2015)015001)

• In the Feynman-'t Hooft gauge, same as [Arganda et al., 2005]:









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• Formulas adapted from [Arganda et al., 2005]

- Diagrams 1, 8, 10  $\rightarrow$  dominate at large  $M_R$
- Enhancement from: - $\mathcal{O}(1) Y_{\nu}$  couplings -TeV scale  $n_i$

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## Most relevant constraints

• Neutrino data  $\rightarrow$  Use specific parametrizations (modified Casas-Ibarra [Casas and Ibarra, 2001] or  $\mu_X$  parametrization)

$$vY_{\nu}^{T} = V^{\dagger} \operatorname{diag}(\sqrt{M_{1}}, \sqrt{M_{2}}, \sqrt{M_{3}}) R \operatorname{diag}(\sqrt{m_{1}}, \sqrt{m_{2}}, \sqrt{m_{3}}) U_{PMNS}^{\dagger}$$
$$M = M_{R} \mu_{X}^{-1} M_{R}^{T}$$
$$OR$$

$$\mu_X = M_R^T m_D^{-1} U_{\text{PMNS}}^* m_\nu U_{\text{PMNS}}^\dagger m_D^{T-1} M_R$$

- Charged lepton flavour violation  $\rightarrow$  For example: Br( $\mu \rightarrow e\gamma$ ) < 5.7 × 10<sup>-13</sup> [MEG, 2013]
- Lepton universality violation: less contraining than  $\mu \rightarrow e\gamma$
- Electric dipole moment: 0 with real PMNS and mass matrices
- Invisible Higgs decays:  $M_R > m_H$ , does not apply



# cLFV Higgs decays from SUSY loops

(arXiv:1508.04623)

In the Feynman-'t Hooft gauge, same as [Arganda et al., 2005]:



- Formulas adapted from [Arganda et al., 2005]
- Enhancement from: - $\mathcal{O}(1) Y_{\nu}$  couplings -TeV scale  $\tilde{\nu}$



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