

Anomaly polynomial of general 6d SCFTs

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Based on arXiv:1408.5572 with Ohmori, Tachikawa, Yonekura

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Review: What are 't Hooft anomalies?

- 't Hooft anomaly: obstruction to gauging global symmetries.

Couple the theory to background gauge fields $A_\mu, g_{\mu\nu}$

→ effective action fails to be gauge invariant:

$$\delta W_d[g_{\mu\nu}, A_\mu] = \int I_d^{(1)}[g_{\mu\nu}, A_\mu].$$

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- Descent equation and anomaly polynomial:

$$I_{d+2} = dI_{d+1}^{(0)}, \quad \delta I_{d+1}^{(0)} = dI_d^{(1)}.$$

Anomaly polynomial I_{d+2} : polynomial of characteristic classes

$$p_1(T) = -\frac{1}{8\pi^2} \text{tr } R^2, \quad p_2(T) = \frac{1}{128\pi^2} ((\text{tr } R^2)^2 - 2 \text{tr } R^4) \text{ etc...}$$

Anomalies and 6d SCFTs

- 6d SCFT: self-dual 2-form gauge fields and tensionless strings.

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- Cancellation of gauge anomalies: strong constraint for 6d SCFTs.

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“Atomic classification” [Bhardwaj '15] [Heckman, Morrison, Rudelius, Vafa '15]
 - 't Hooft anomalies for global symmetries:
 - N^3 scaling law of d.o.fs
 - Central charges of compactified theory [Benini, Tachikawa, Wecht '09]
 - RG flow between 6d SCFTs [Heckman Morrison Rudelius Vafa '15]
- etc...

How to calculate 't Hooft anomalies of 6d SCFTs?

- Gravitational calculation:

embed 6d SCFT into M-theory and use anomaly inflow.

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- Field theory calculation:

deform 6d SCFT to free field theory and use anomaly matching.

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In this short talk, I will explain the field theoretical calculation of anomaly polynomial of 6d $\mathcal{N} = (2, 0)$ theory.

Anomaly polynomial of 6d $\mathcal{N}=(2,0)$ theory

- G -type $\mathcal{N}=(2,0)$ theory: IIB on ADE orbifold \mathbb{C}^2/Γ_G .

Anomaly polynomial conjecture: [Intriligator '00]

$$I_G^{\mathcal{N}=(2,0)} = \frac{h_G^\vee d_G}{24} p_2(N) + r_G I^{\mathcal{N}=(2,0) \text{ tensor}},$$

$$I^{\mathcal{N}=(2,0) \text{ tensor}} = \frac{1}{48} \left(p_2(N) - p_2(T) + \frac{1}{4} (p_1(T) - p_1(N))^2 \right).$$

$$h_{\text{SU}(k)}^\vee d_{\text{SU}(k)} = k^3 - k \quad N: \text{SO}(5) \text{ R-symmetry}$$

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- When $G = A_{k-1}, D_k$, gravitational calculation is available:

$G = A_{n-1} \rightarrow$ coincident k M5-branes, [(Freed,) Harvey, Minasian, Moore '98]

$G = D_n \rightarrow$ coincident k M5-branes + orientifold. [Yi '01]

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 $G = D_n \rightarrow$ coincident k M5-branes + orientifold. [Yi '01]
- Field theory calculation is valid for any G . [Ohmori, HS, Tachikawa, Yonekura]

Tensor branch RG flow

- Tensor branch RG flow: giving tension to self-dual strings.

G -type $\mathcal{N}=(2,0)$ theory $\rightarrow r_G$ free $\mathcal{N}=(2,0)$ tensor multiplets.

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IR

$(B_2, \phi^{i=1\cdots 5}, \text{fermions})$

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- Anomaly matching on tensor branch:

$$I_G^{\mathcal{N}=(2,0)} = (\text{anomaly matching term}) + r_G I^{\mathcal{N}=(2,0)} \text{ tensor}.$$

Origin of anomaly matching term?

Green-Schwarz mechanism for anomaly matching

- Integrate out massive strings: induce electric/magnetic coupling for B_2

$$\Delta L = \int_{X_6} \Omega^{ij} B_i I_j \quad \text{and} \quad \underbrace{dH_i = I_i}_{\text{Bianchi identity}} \quad i = 1 \cdots r_G.$$

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- Additional contribution from **Green-Schwarz mechanism**:

$$I^{\text{GS}} = \frac{1}{2} \Omega^{ij} I_i I_j$$

$$I_G^{\mathcal{N}=(2,0)} = I^{\text{GS}} + r_G I^{\mathcal{N}=(2,0)} \text{ tensor}.$$

How to determine I_i ?

S^1 compactification of 6d $\mathcal{N}=(2,0)$ theory

- G -type 6d $\mathcal{N}=(2,0)$ theory on $S^1 \rightarrow$ 5d G $\mathcal{N}=2$ SYM.

6d tensor branch \rightarrow 5d Coulomb branch: $G \rightarrow U(1)^{r_G}$.

$\langle \phi_{i=1\dots 4}^a \rangle = 0$ and $\langle \phi_5^a \rangle = v^a$, then $SO(5)_R \rightarrow SO(4)_R$.

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- Massive strings \rightarrow massive $\mathcal{N}=2$ vectors Φ_α (α : roots of G) w/ mass $v \cdot \alpha$.

Integrating out Φ_α : **induced Chern-Simons terms**

$$S^{CS} = \Omega^{ij} A_i I_j, \quad A_i : \mathrm{U}(1)_i \text{ gauge field.}$$

Reduce to ordinary 1-loop calculation!

Calculation

- Induced Chern-Simons term for A_i $i = 1 \cdots r_G$:

$$\frac{1}{2} \sum_{\alpha > 0} (\alpha \cdot A) \left[\underbrace{\left(c_2(L) + \frac{2}{24} p_1(T) \right)}_{\text{pos real mass fermions}} - \underbrace{\left(c_2(R) + \frac{2}{24} p_1(T) \right)}_{\text{neg real mass fermions}} \right]$$
$$= \rho \cdot A (c_2(L) - c_2(R)),$$

$$\rho = \frac{1}{2} \sum_{\alpha > 0} \alpha: \text{Weyl vector}, \quad \text{SO}(4)_R \sim \text{SU}(2)_L \times \text{SU}(2)_R.$$

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- **F-theory geometry**: determine both massless spectrum and Green-Schwarz coupling → **anomaly polynomial of general 6d SCFTs.**

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- We established the field theoretical way to calculate them.
- Anomaly matching on tensor branch:
Need massless spectrum/Green-Schwarz coupling on tensor branch.
- **F-theory geometry**: determine both massless spectrum and Green-Schwarz coupling \rightarrow **anomaly polynomial of general 6d SCFTs**.
- These polynomials may be used to prove a-theorem for 6d SCFTs/
investigate compactification etc.....