

# Why B-mode?

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**Yes!**

Let's take a look...



$$H \equiv \frac{\dot{a}}{a} \quad [a \text{ is the scale factor}]$$

# Inflation, defined

- **Accelerated expansion during the early universe**

$$\frac{\ddot{a}}{a} = \dot{H} + H^2 > 0 \quad \longrightarrow \quad \epsilon \equiv -\frac{\dot{H}}{H^2} < 1$$

- For inflation to explain flatness of our observable universe, a **sustained** period of acceleration is required
- This implies  $\epsilon = O(N^{-1})$  [or smaller], where  $N$  is the number of e-fold of expansion counted from the end of inflation:

$$N \equiv \ln \frac{a_{\text{end}}}{a} = \int_t^{t_{\text{end}}} dt' \, H(t') \approx 50$$

# Have we found inflation?

- I.e., *have we found  $\varepsilon \ll 1$ ?*
- To show  $\varepsilon \ll 1$ , **we need to map  $H(t)$** 
  - In other words, we need to draw the “Hubble diagram” during inflation

$$\epsilon \equiv -\frac{\dot{H}}{H^2}$$

# We measure distortions in space

- A distance between two points in space

$$d\ell^2 = a^2(t)[1 + 2\zeta(\mathbf{x}, t)][\delta_{ij} + h_{ij}(\mathbf{x}, t)]dx^i dx^j$$

- $\zeta$ : “curvature perturbation” (scalar mode)
  - Perturbation to the determinant of the spatial metric
- $h_{ij}$ : “gravitational waves” (tensor mode)
  - Perturbation that does not change the determinant (area)

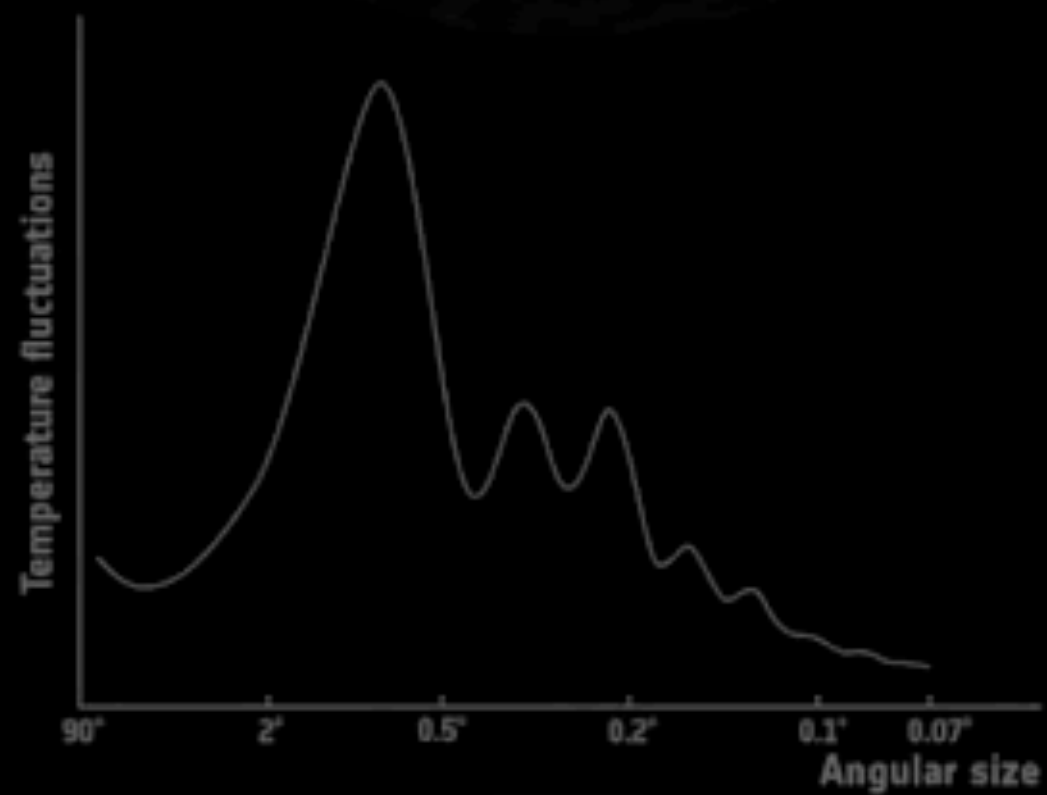


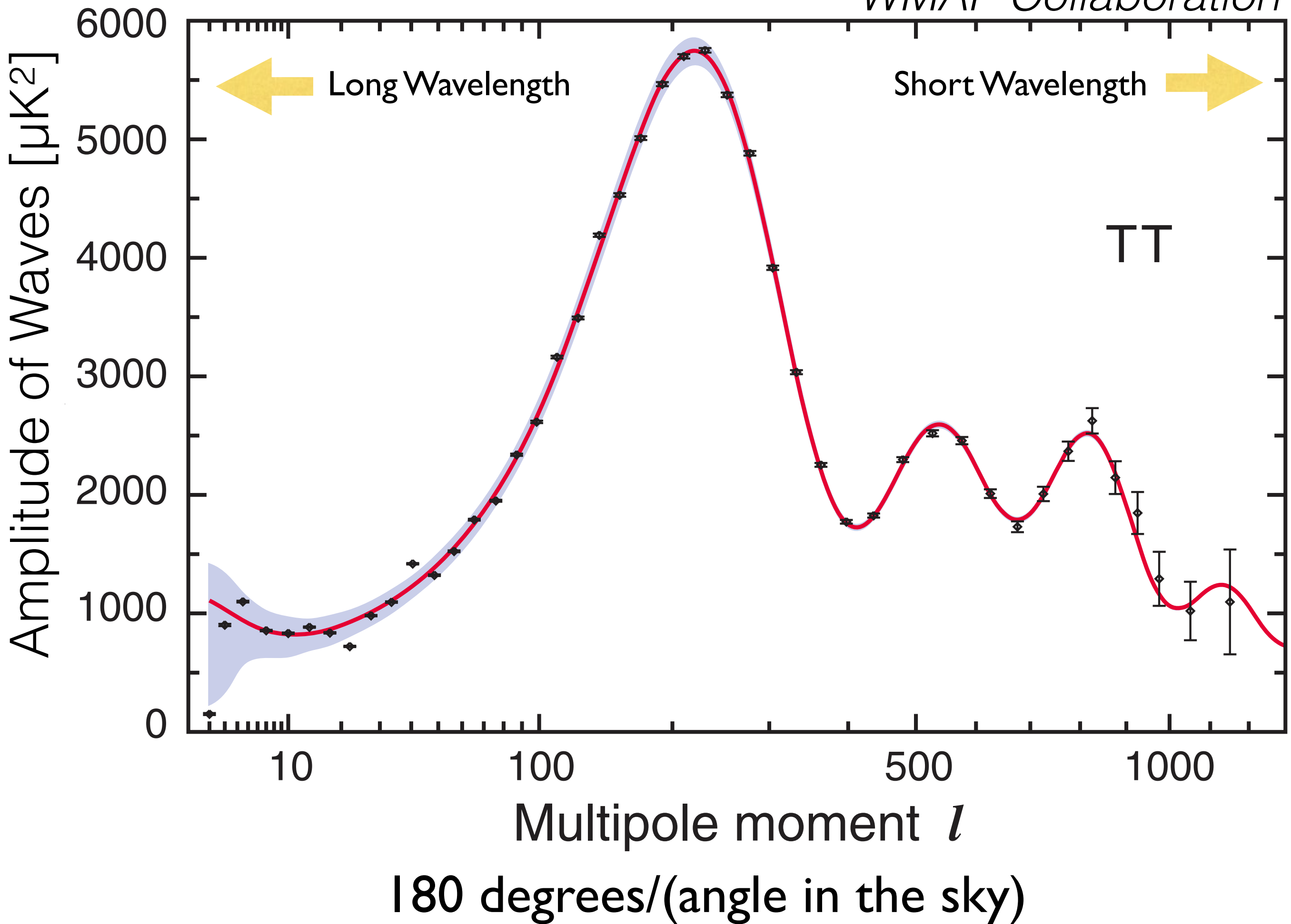
$$\sum_i h_{ii} = 0$$

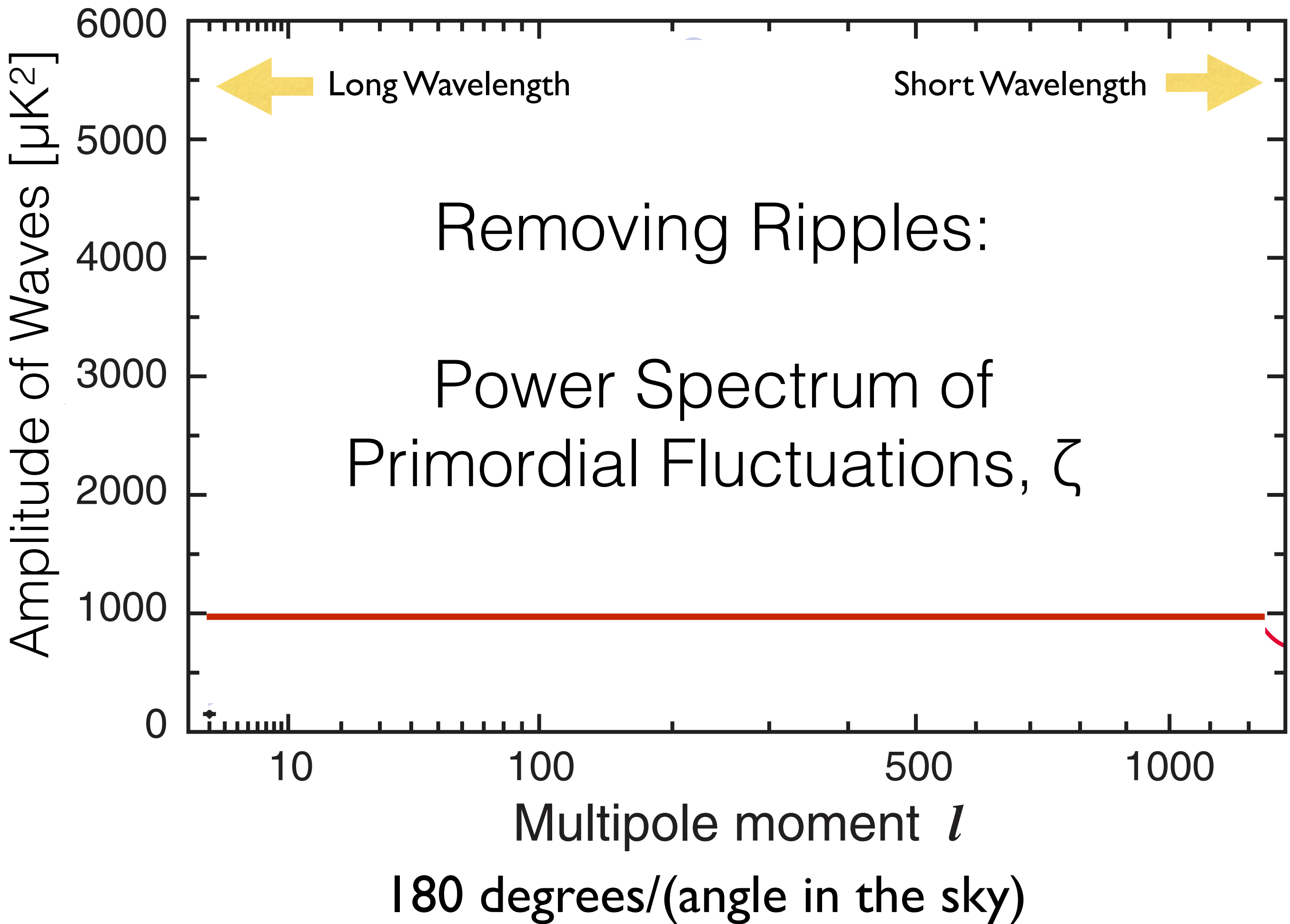


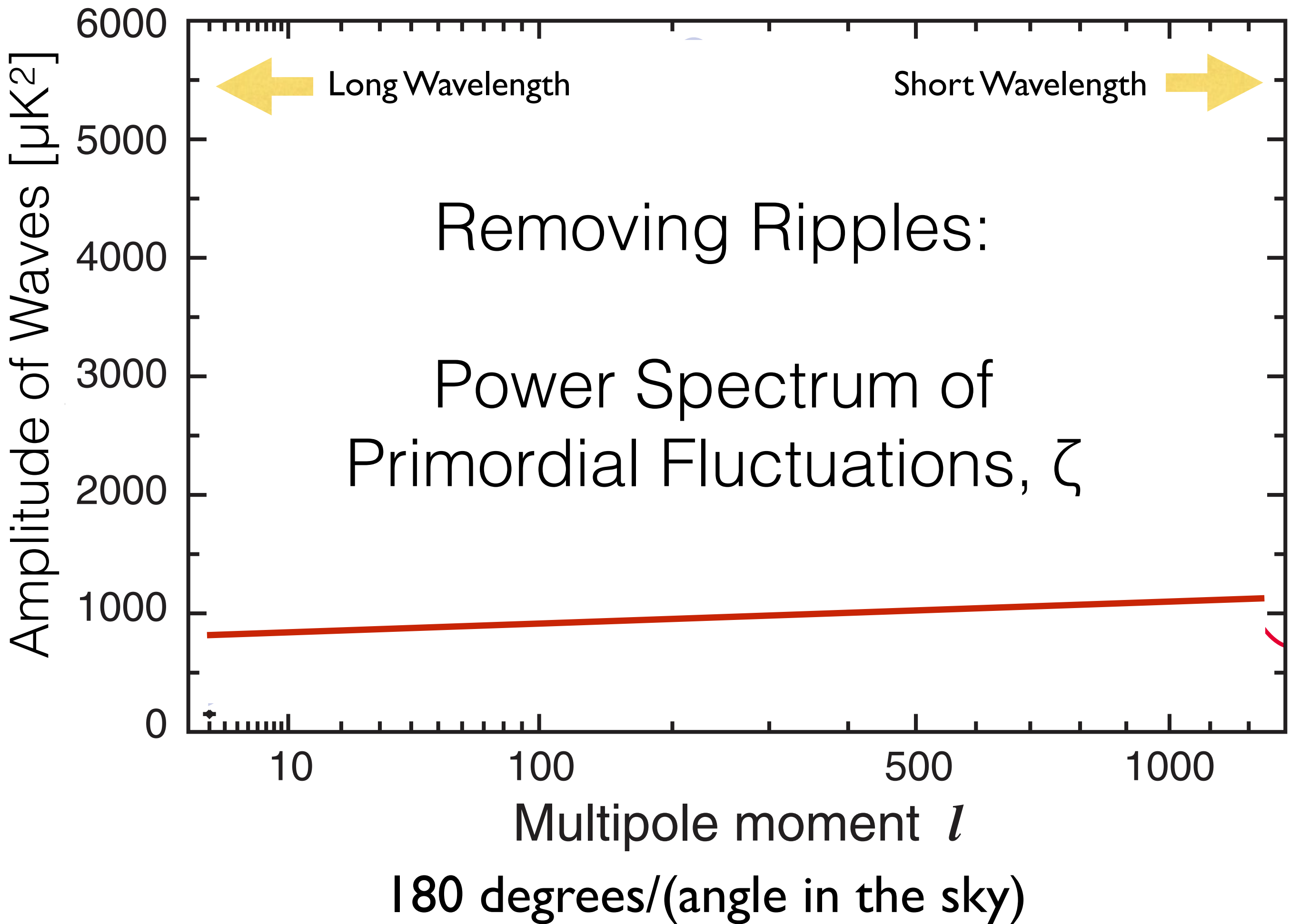
# Fluctuations are proportional to $H$

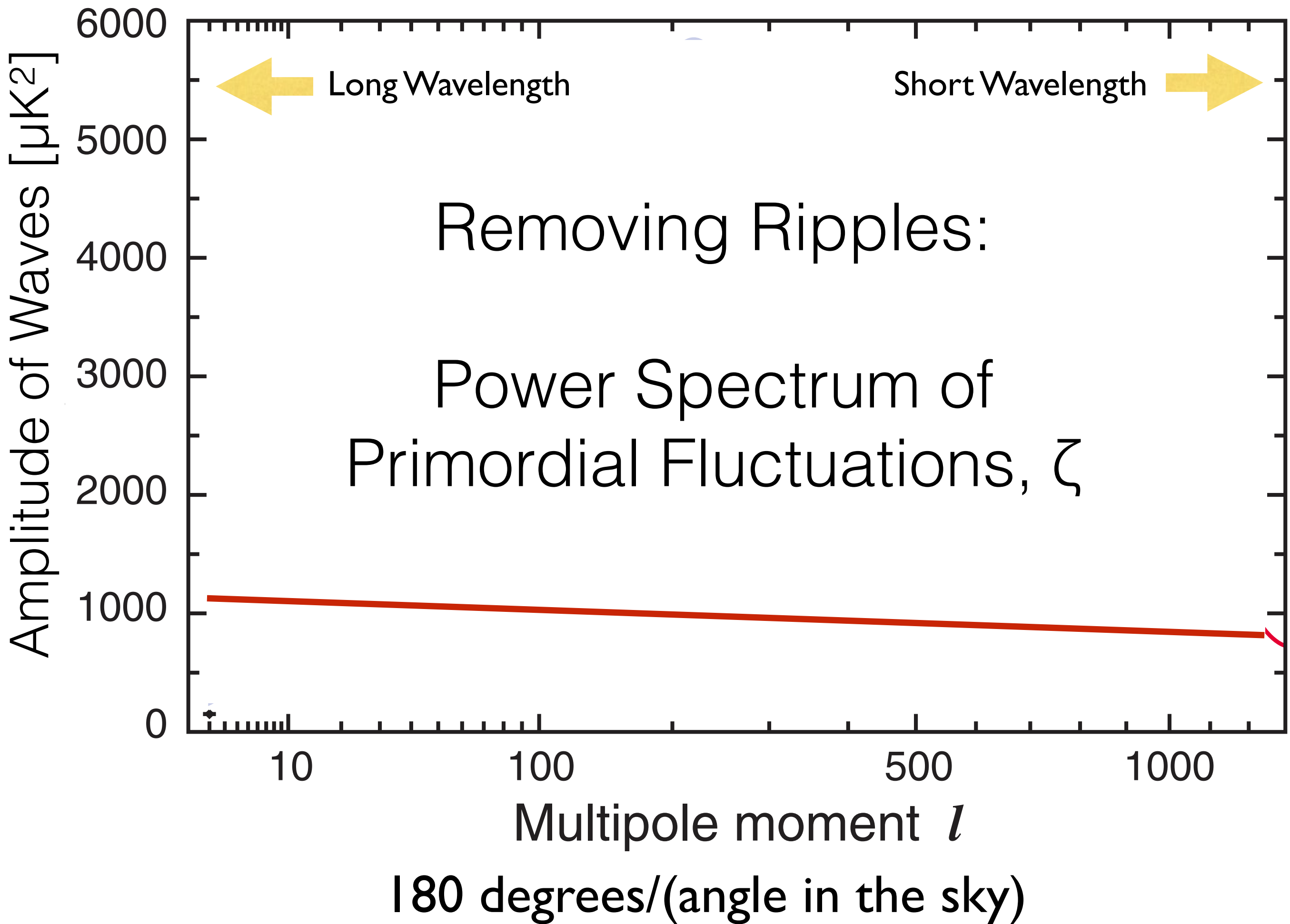
- Uncertainty Principle:
  - *[Energy you can borrow] x [Time you borrow] = constant*
- $H \equiv \frac{\dot{a}}{a}$  [This has units of 1/time]
- Then, **both  $\zeta$  and  $h_{ij}$  are proportional to  $H$**
- Earlier the fluctuations are generated, the bigger the angles they subtend in the sky. **We can map  $H(t)$  by measuring fluctuations over a wide range of angles**

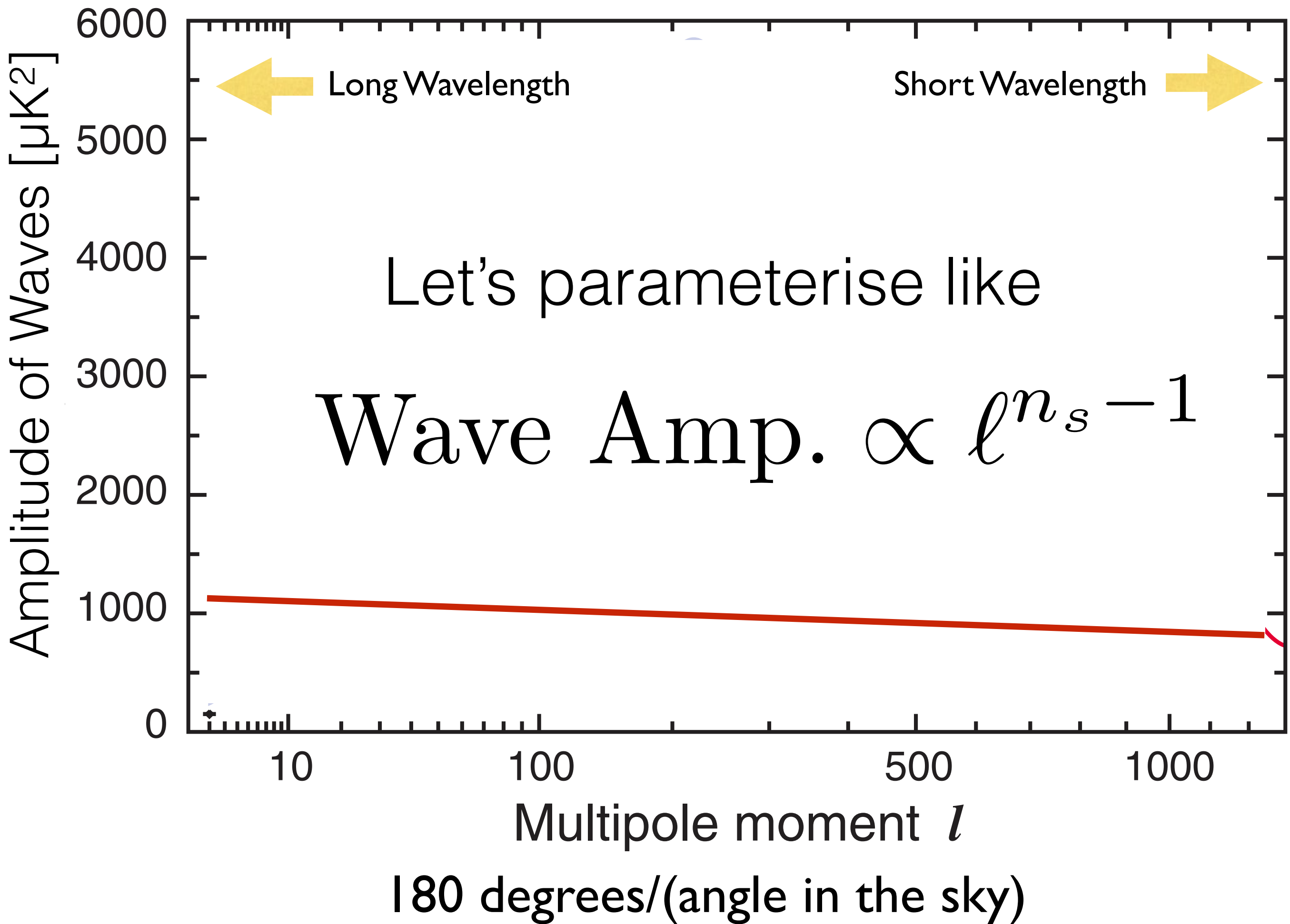


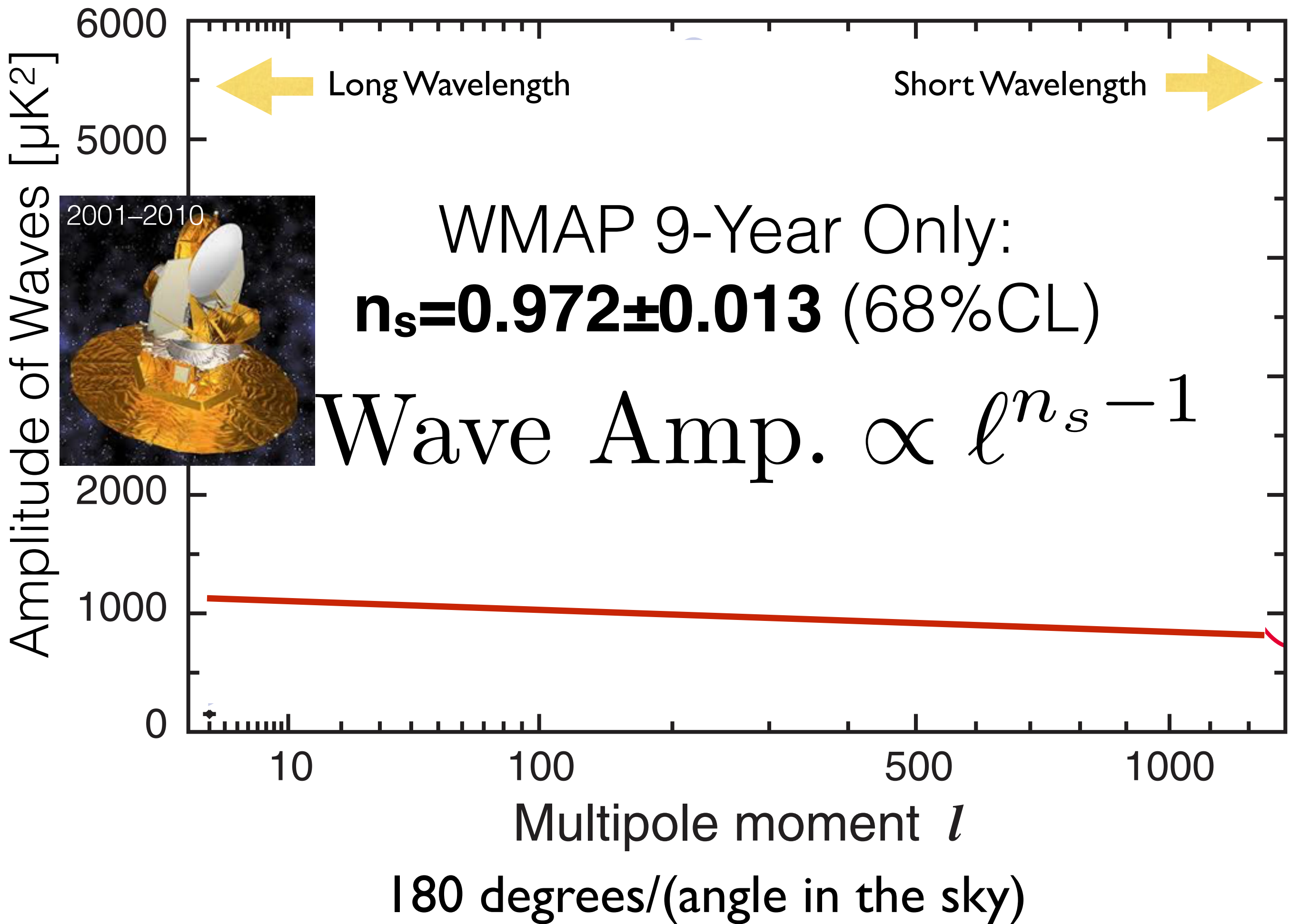




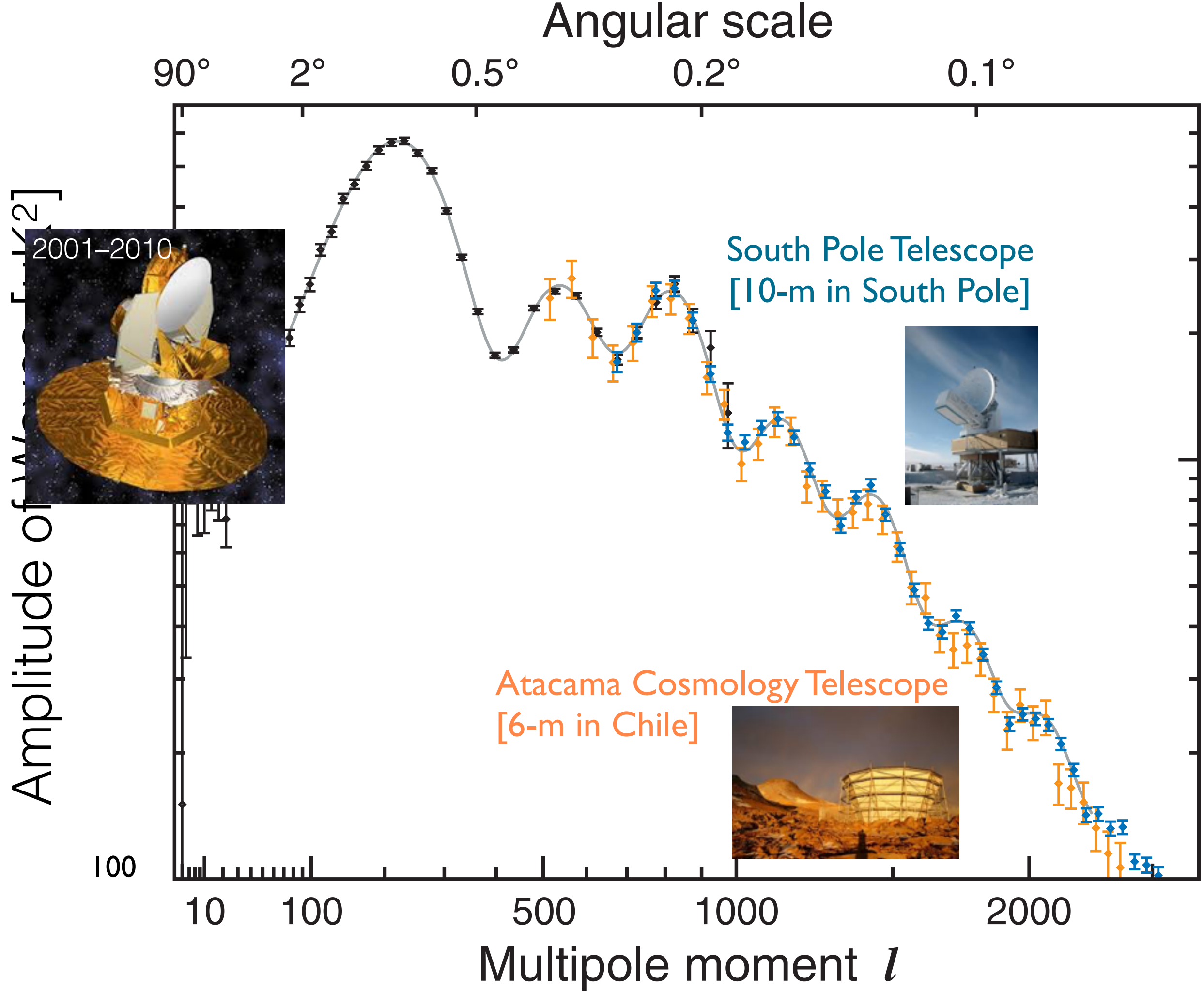


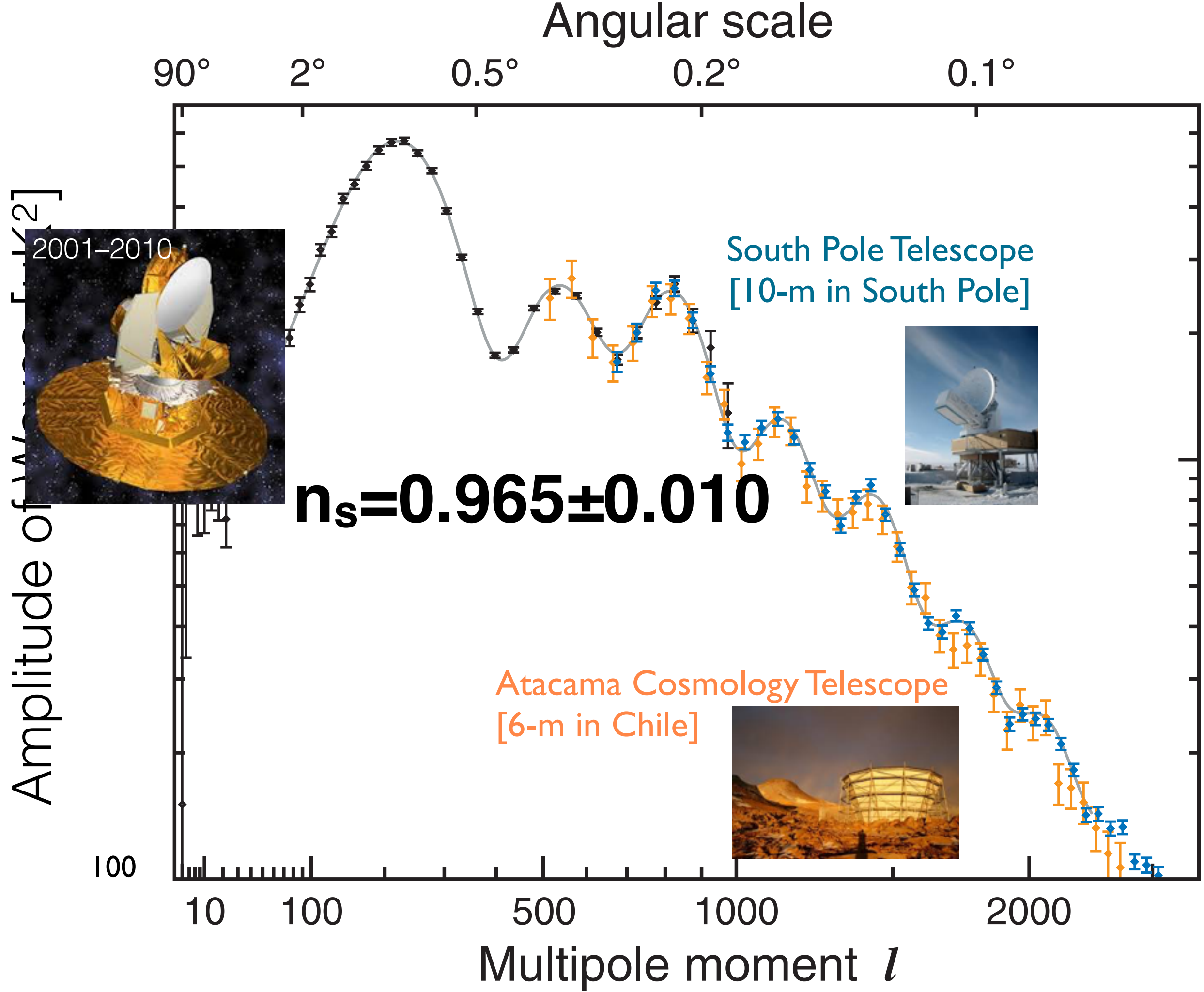


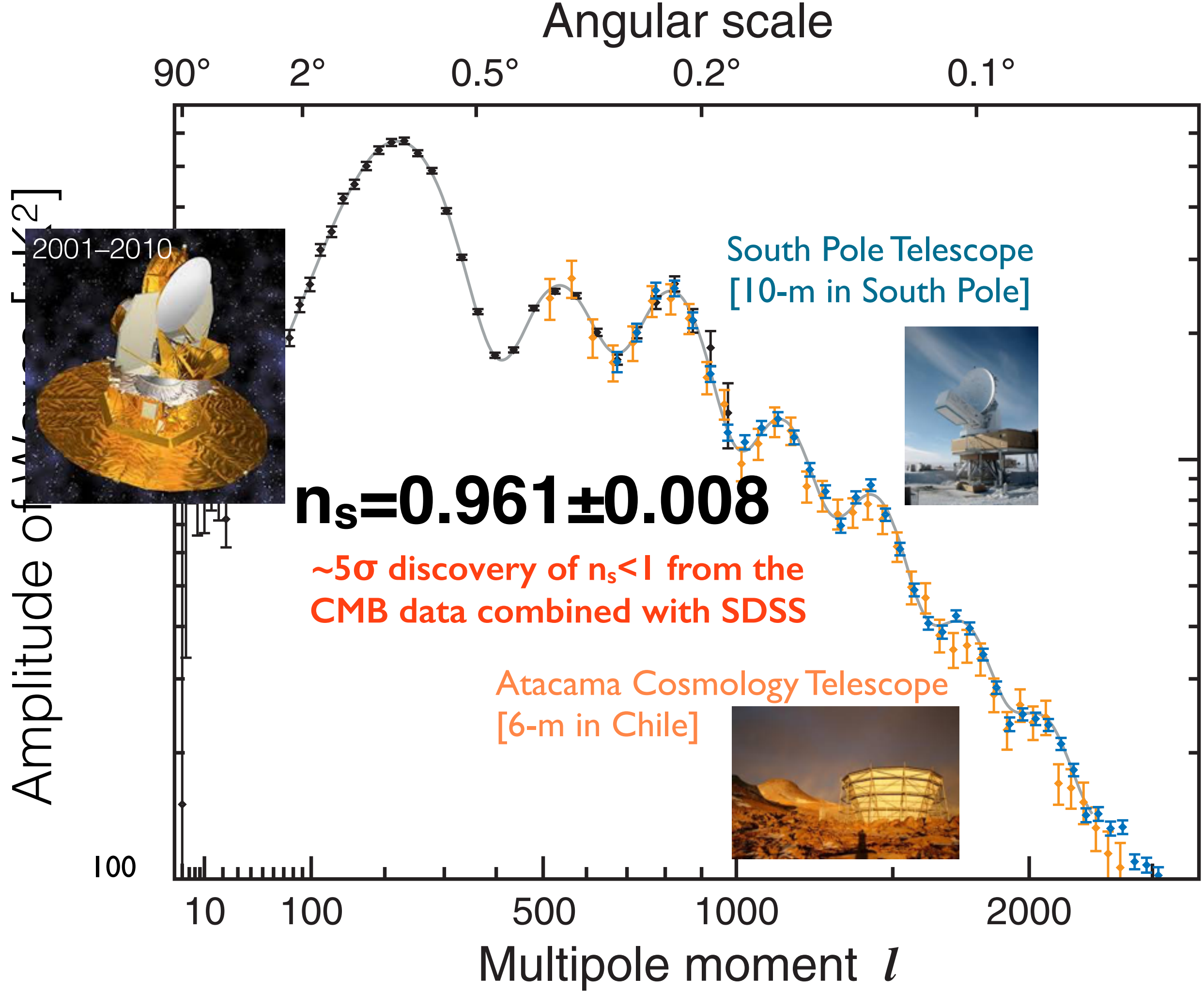








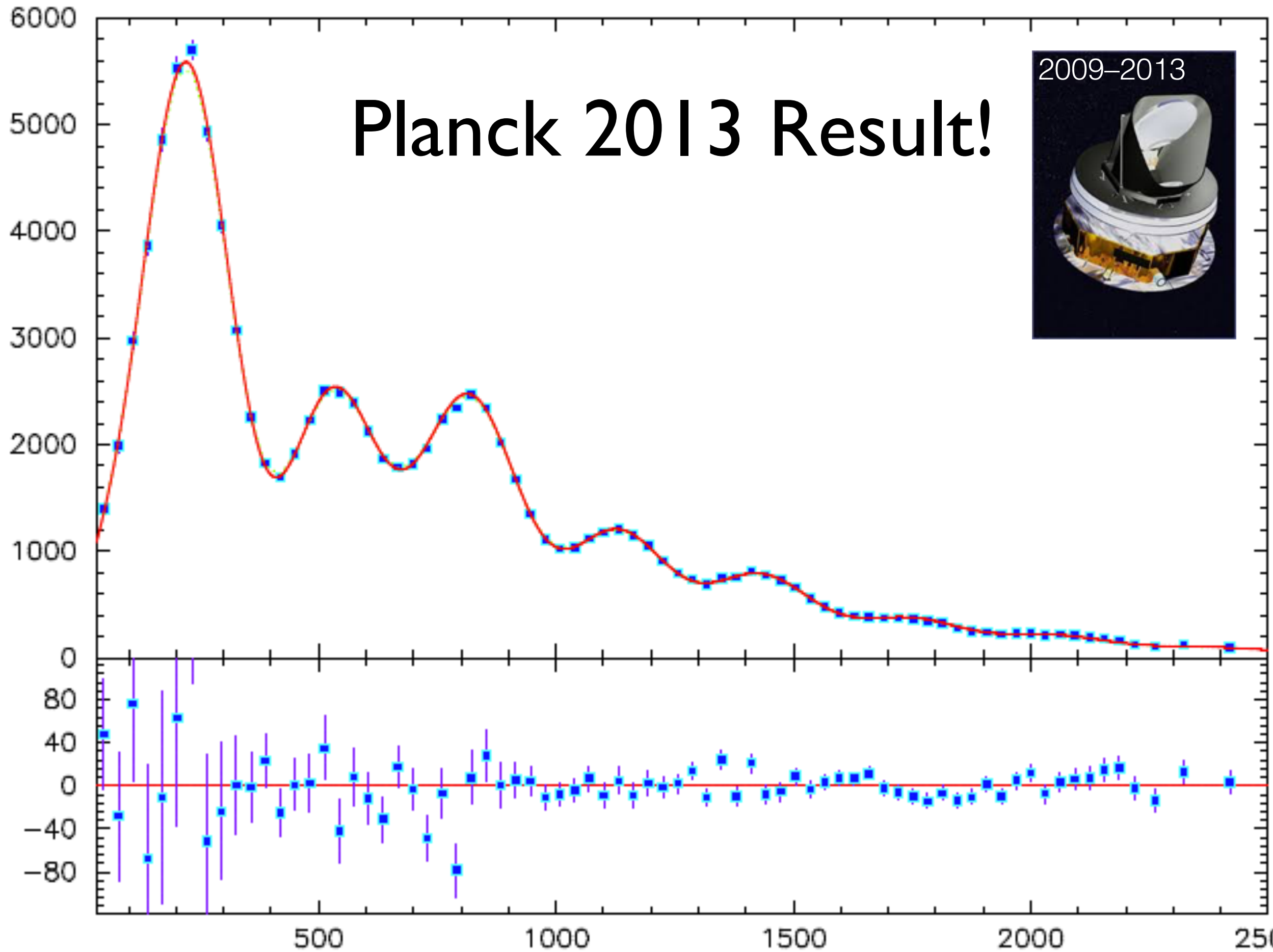
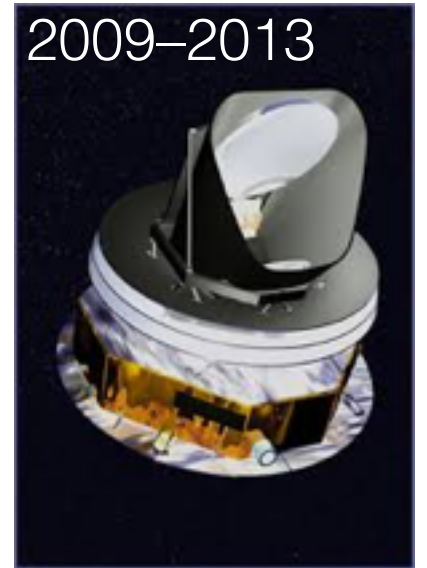




Residual Amplitude of Waves [ $\mu\text{K}^2$ ]

# Planck 2013 Result!

2009–2013



$l$  80 degrees/(angle in the sky)



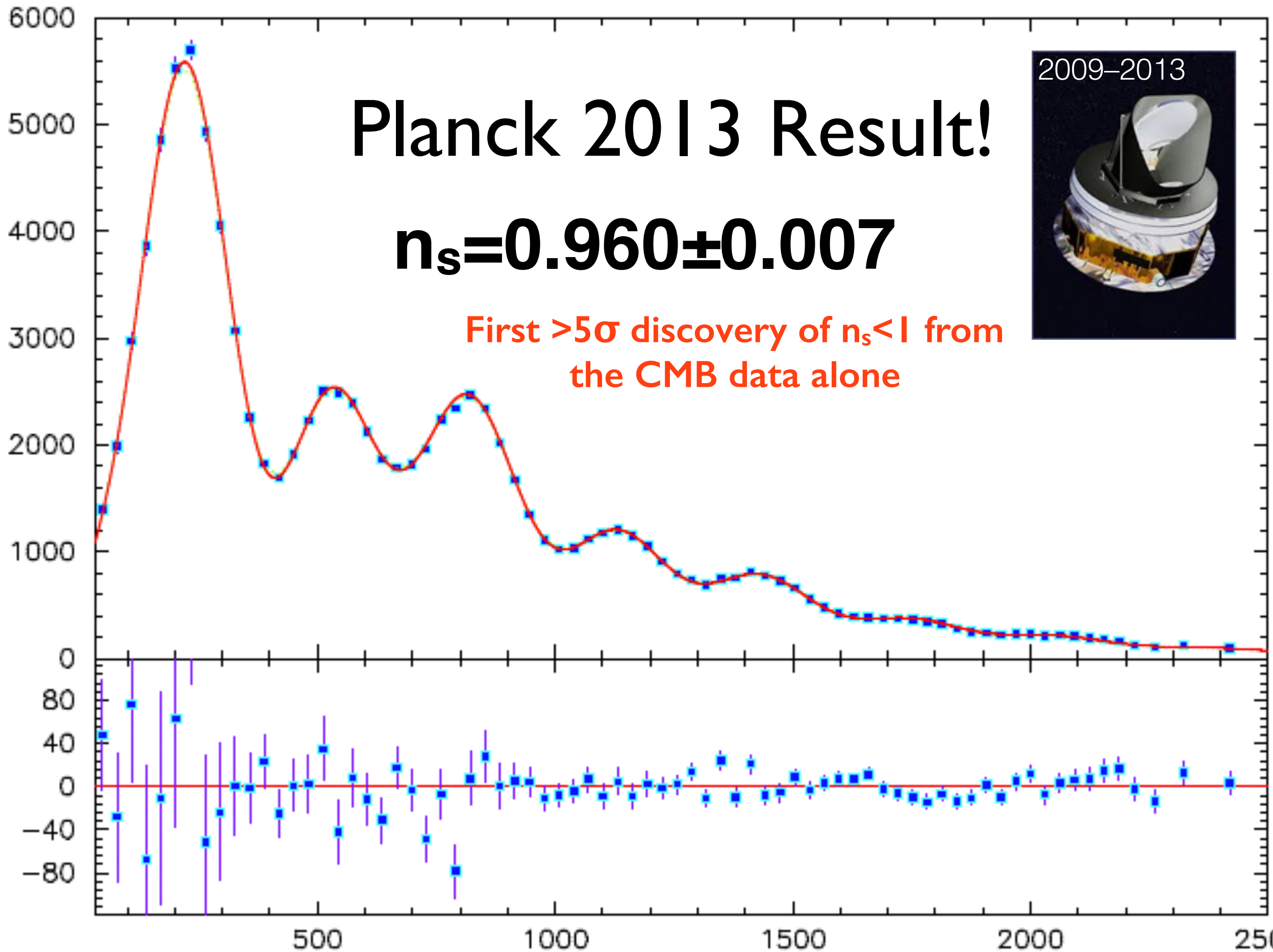
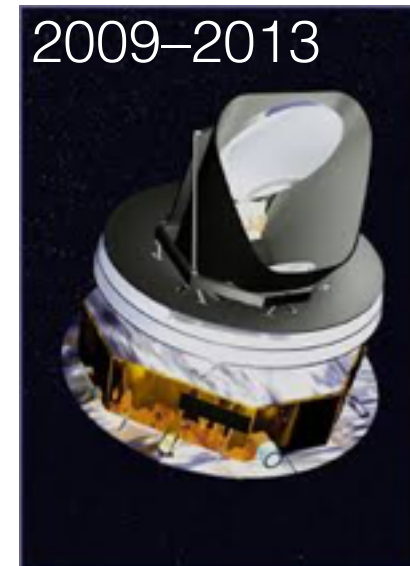
Residual Amplitude of Waves [ $\mu\text{K}^2$ ]

# Planck 2013 Result!

$$n_s = 0.960 \pm 0.007$$

First  $>5\sigma$  discovery of  $n_s < 1$  from  
the CMB data alone

2009–2013



$l$  80 degrees/(angle in the sky)

$$1 - n_s \ll 1:$$

Have we seen  $\varepsilon \ll 1$ ?

- **Not quite.**  $\zeta$  is basically proportional to  $H$ , but the pre-factor can depend on time. If there was only one dominant field in the early universe:

$$\zeta = (2\varepsilon c_s)^{-1/2} \times H$$

*propagation speed  
of a scalar field*

- Even if  $H(t)$  varies rapidly,  $n_s \sim 1$  can still be achieved if  $\varepsilon$  or  $c_s$  or both also vary, canceling the variation in  $H(t)$ .  **$\zeta$  does not give  $H(t)$  directly**

# In general:

- $\zeta$  does not probe the expansion history during inflation directly because **its behaviour depends very much on properties of matter** present in the universe
- The connection to  $H(t)$  can be more complicated for multi-field (multi-matter) models
- We need a probe which maps  $H(t)$  **more directly**

# Here comes gravitational waves

- Gravitational waves are not coupled to (scalar) matter. Thus, it directly probes  $H(t)$  via

$$h_{ij}^{\text{prim}} = \frac{\sqrt{2}e_{ij}}{M_{\text{Pl}}} \times H$$

*independent of time!*



# Has Inflation Occurred?

- We must see [near] scale invariance of the gravitational wave power spectrum:

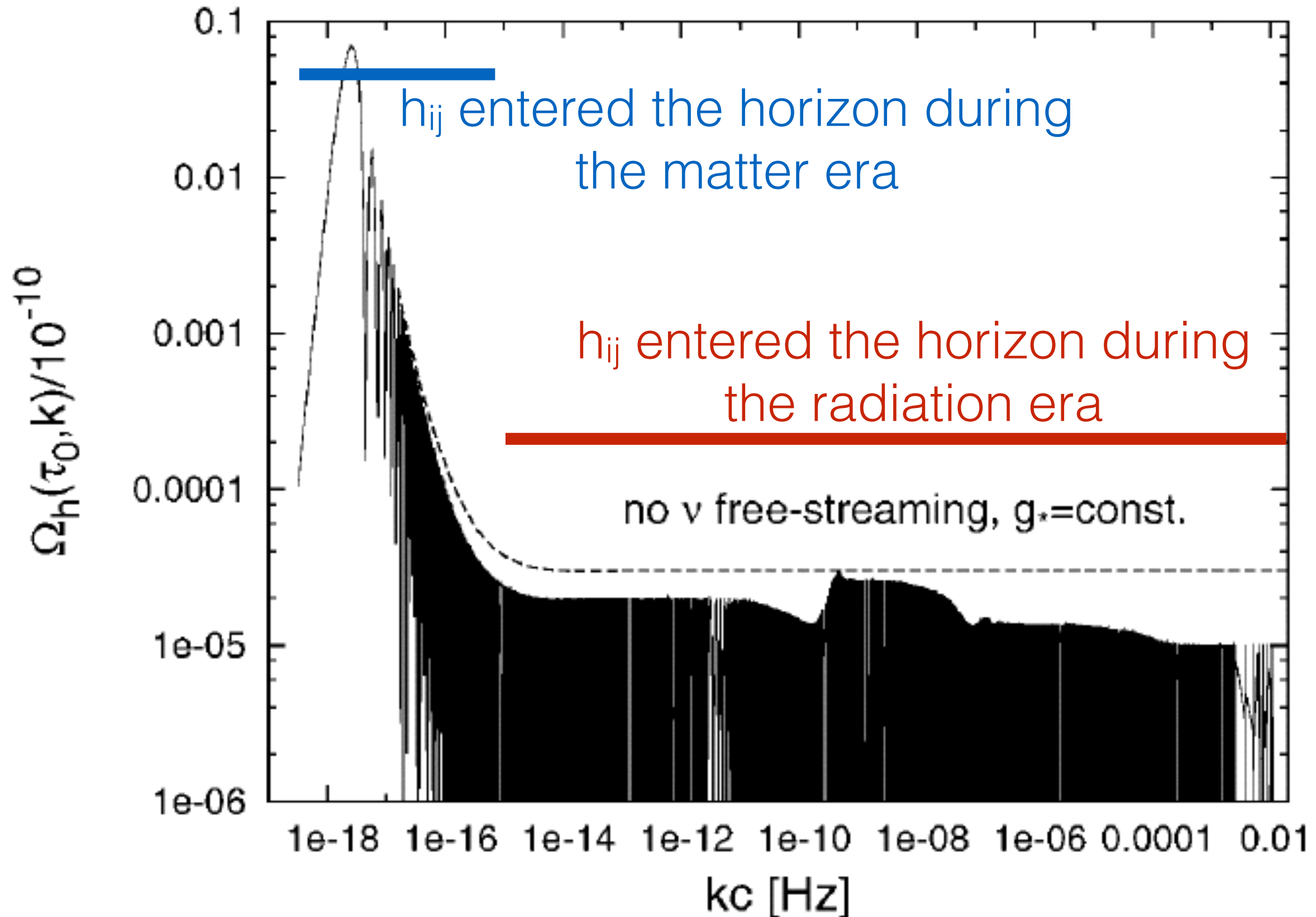
$$\langle h_{ij}^{\text{prim}}(\mathbf{k}) h_{\text{prim}}^{ij,*}(\mathbf{k}) \rangle \propto k^{n_t}$$

with

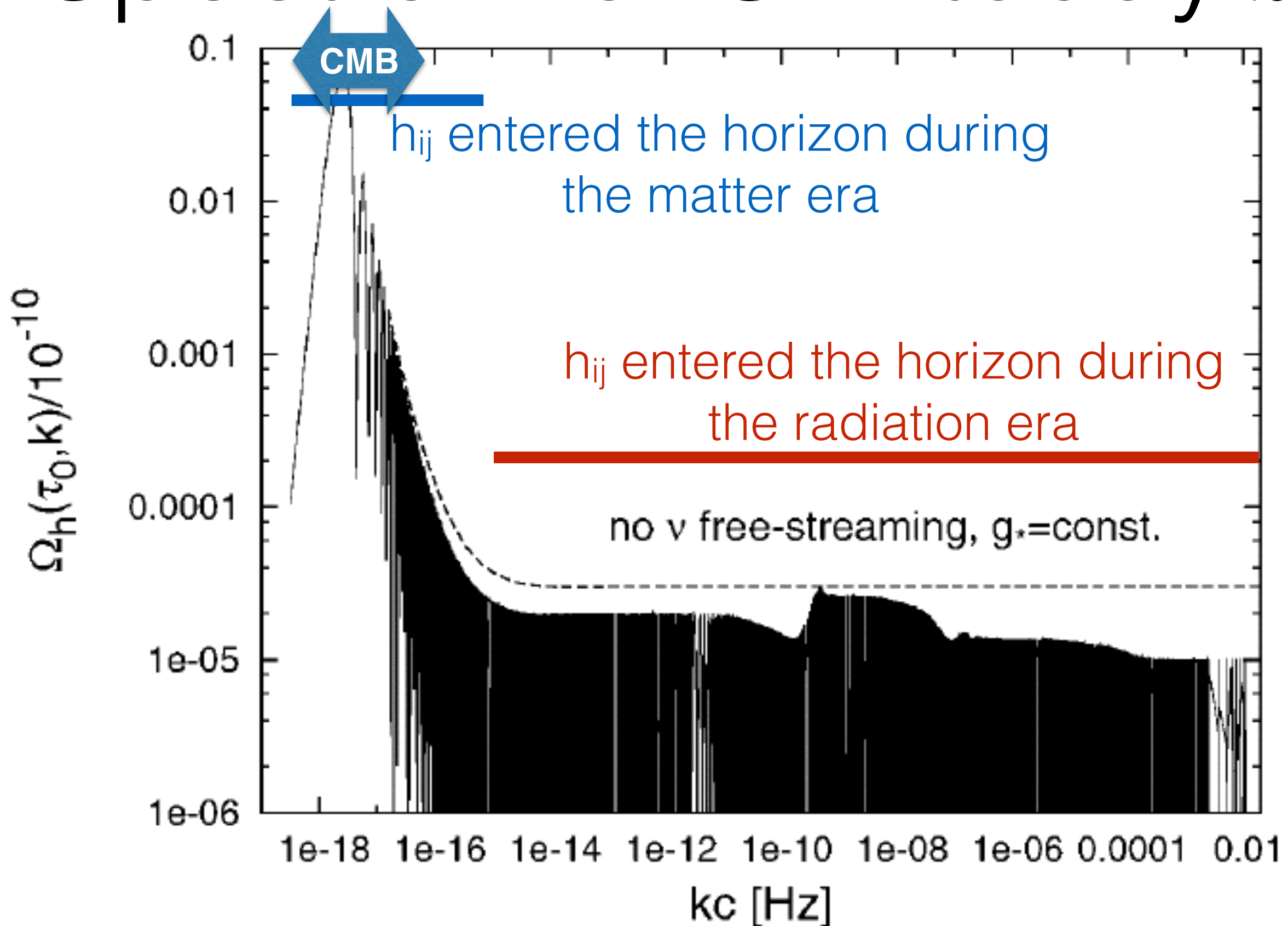
$$|n_t| \ll 1$$

In most models,  
 $n_t = -2\epsilon < 0$

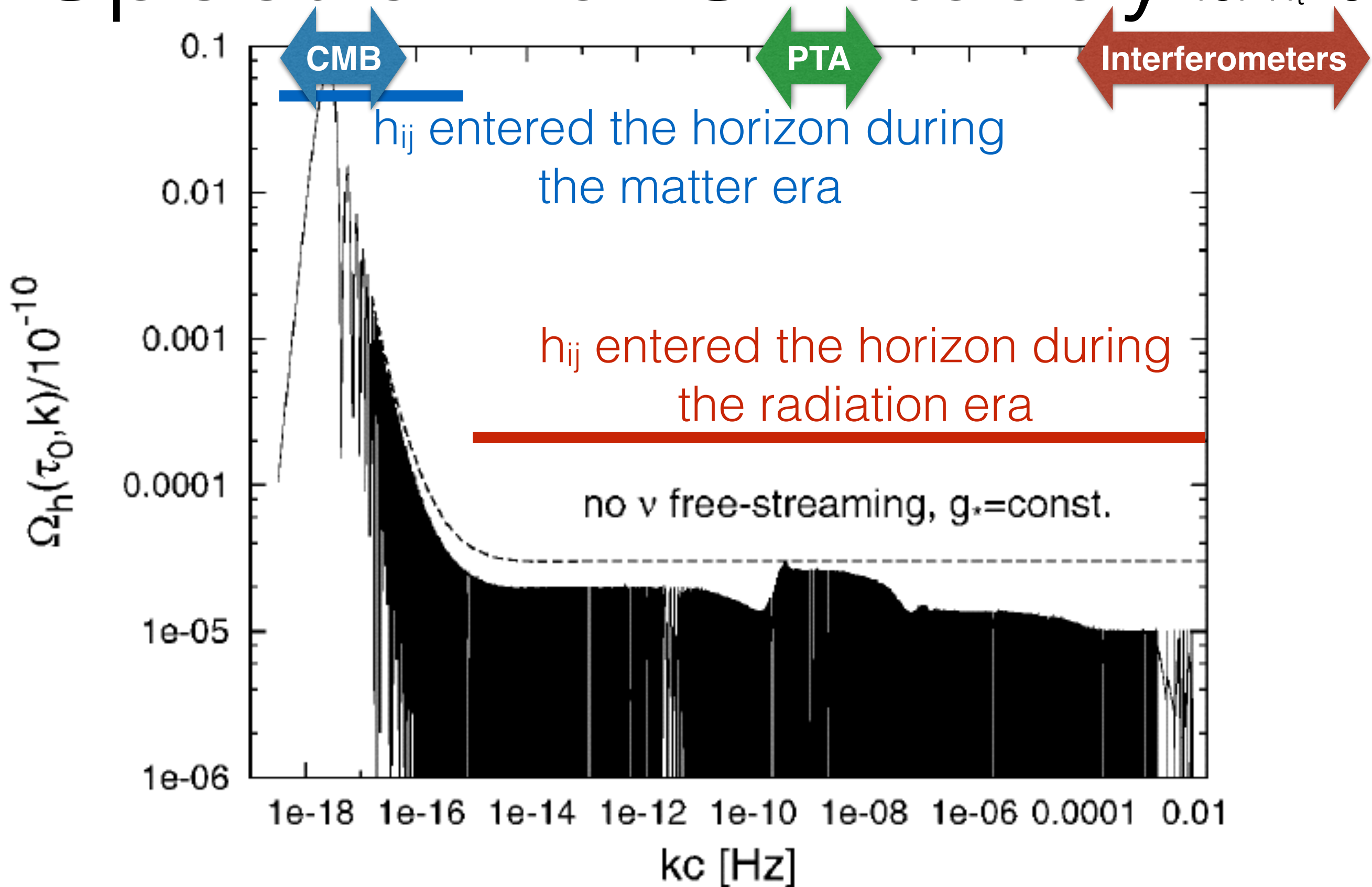
# Spectrum of GW today for $n_t=0$



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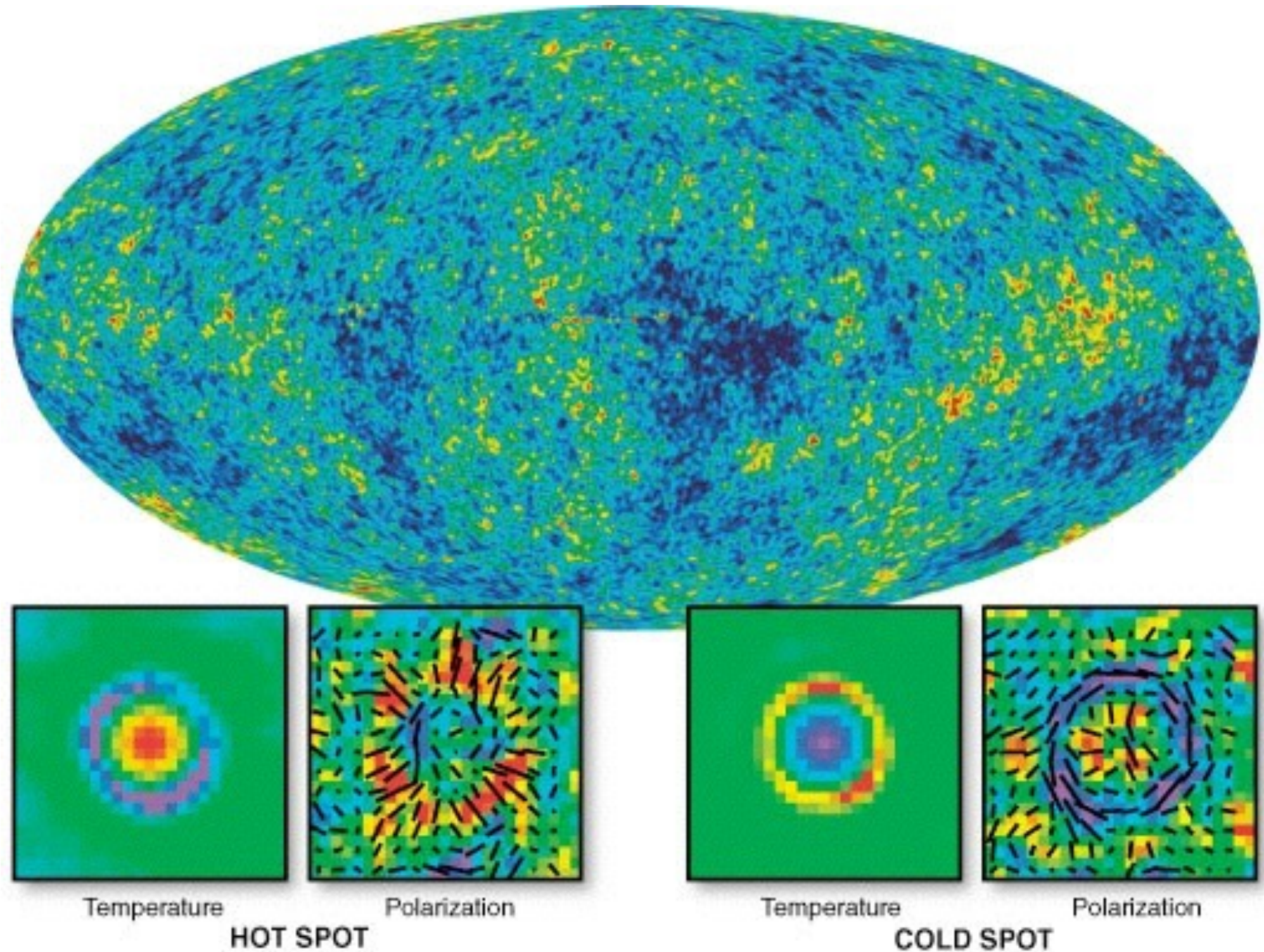


# Spectrum of GW today for $n_t=0$



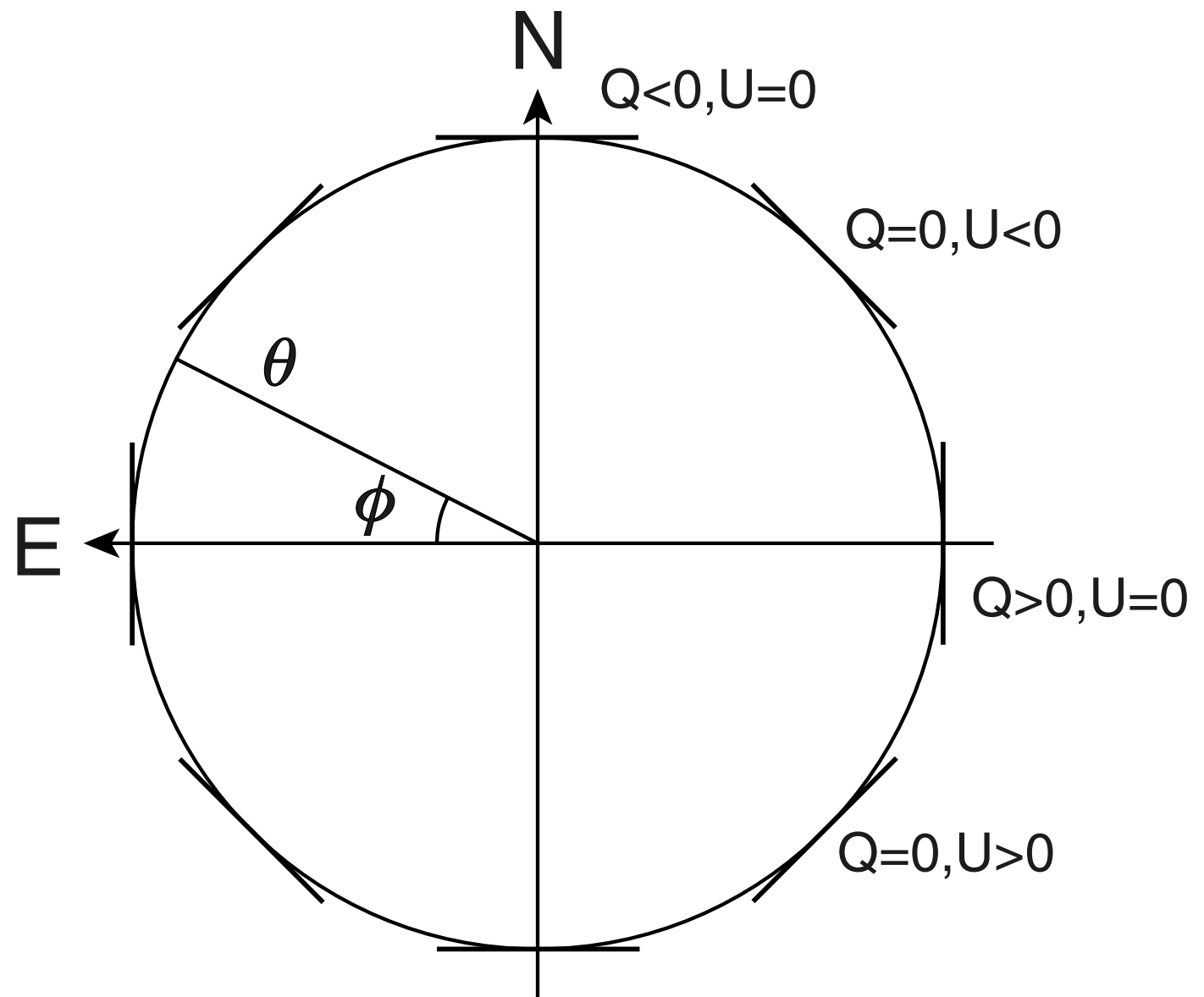
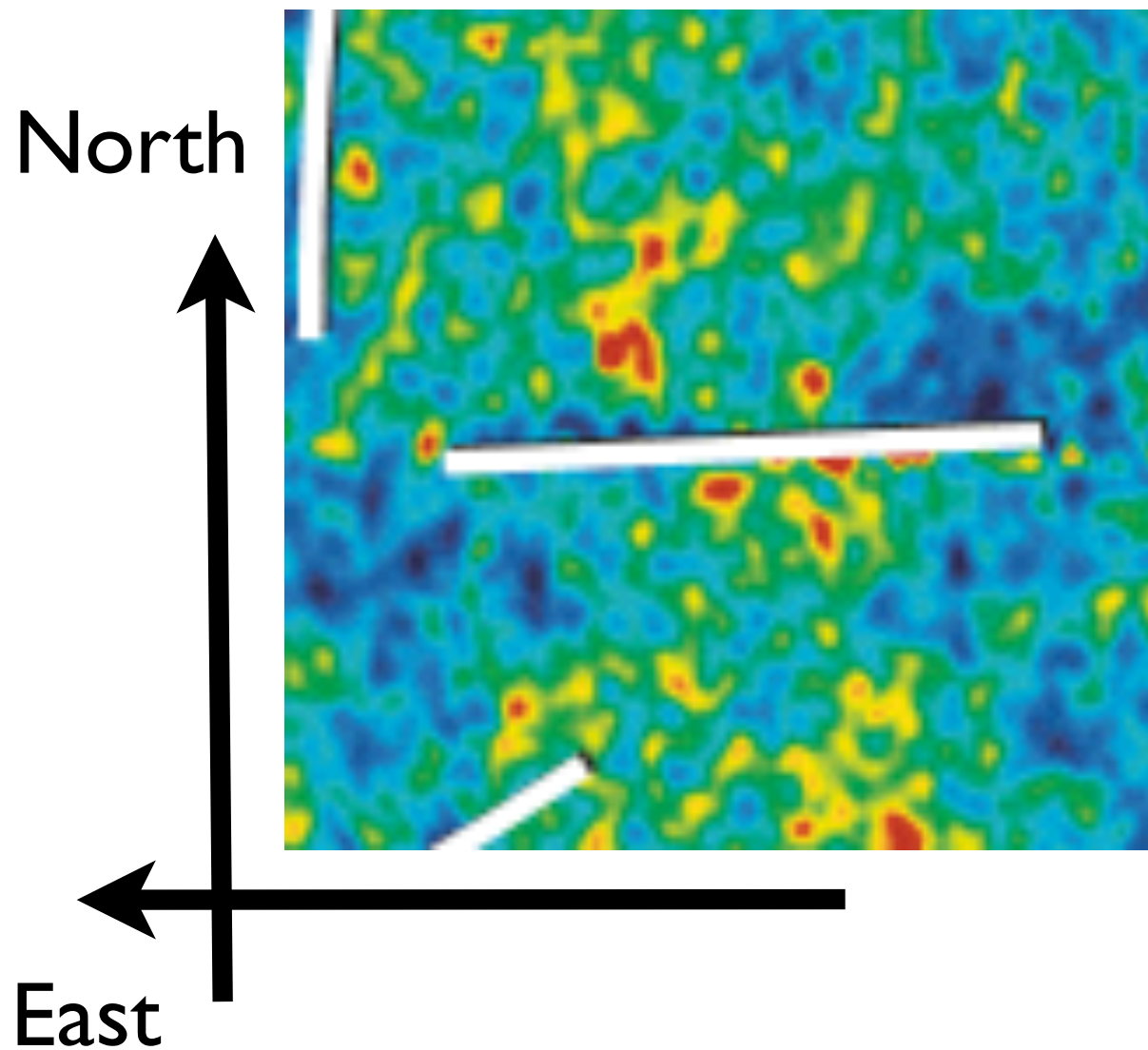


# CMB Polarisation



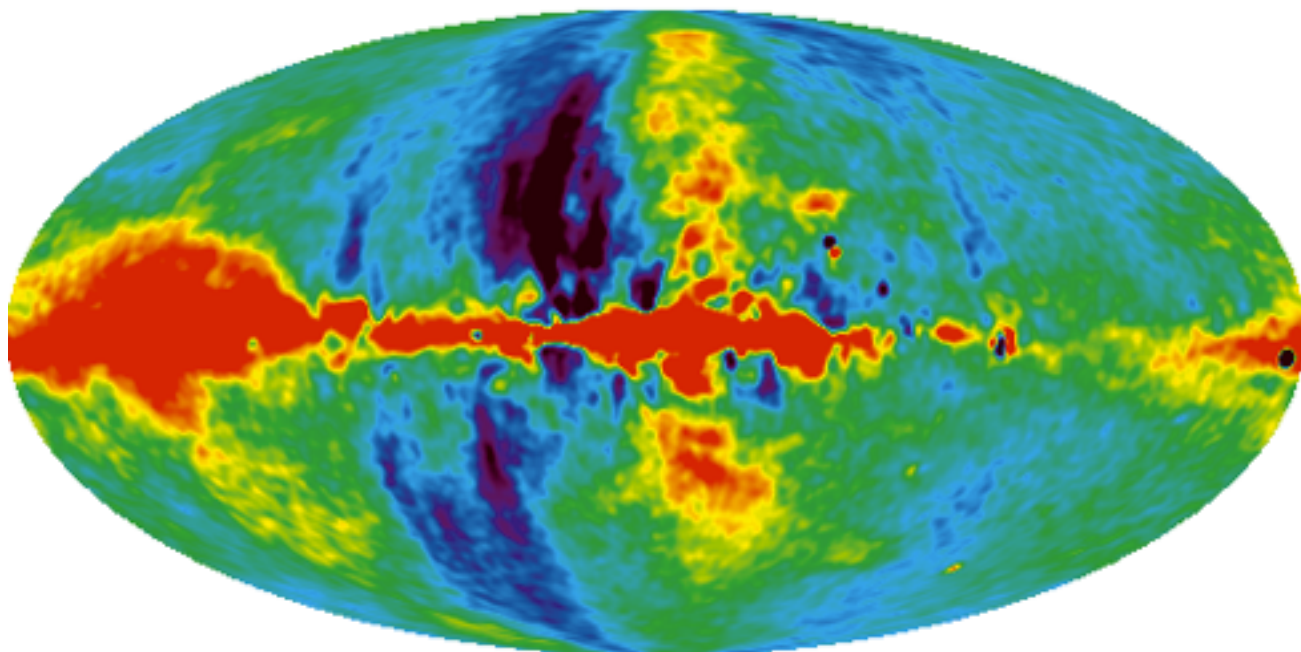
- CMB is [weakly] polarised!

# Stokes Parameters

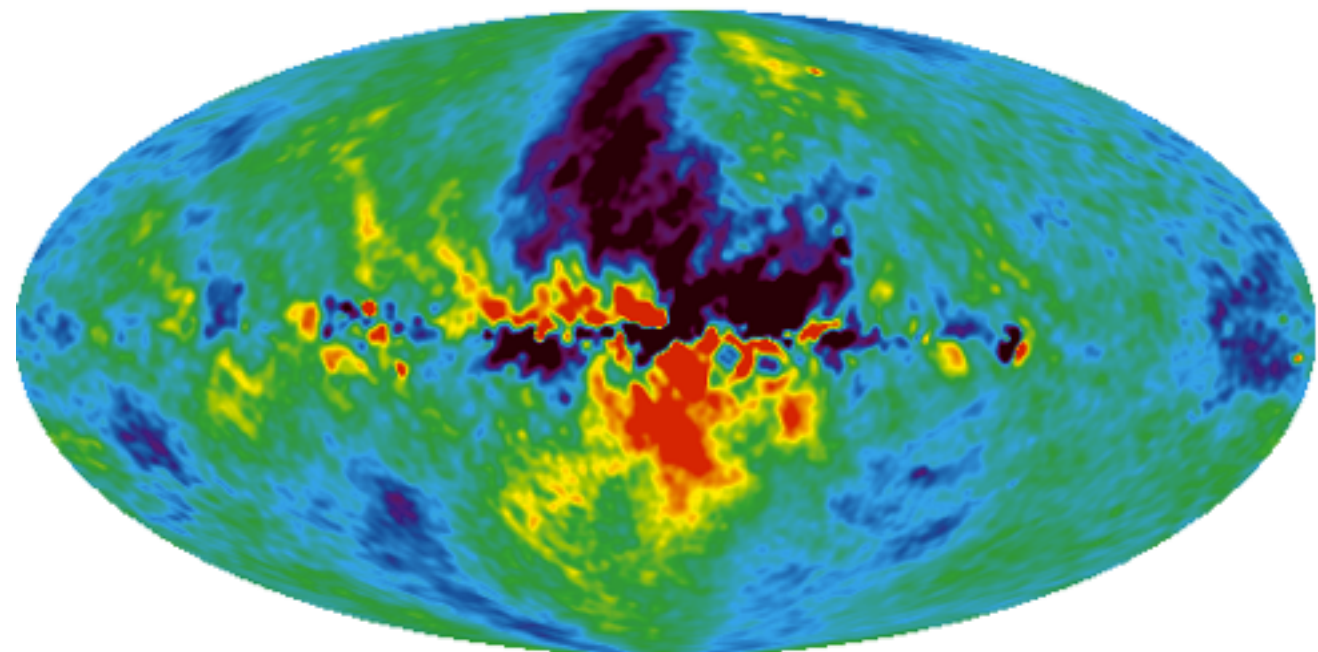




23 GHz

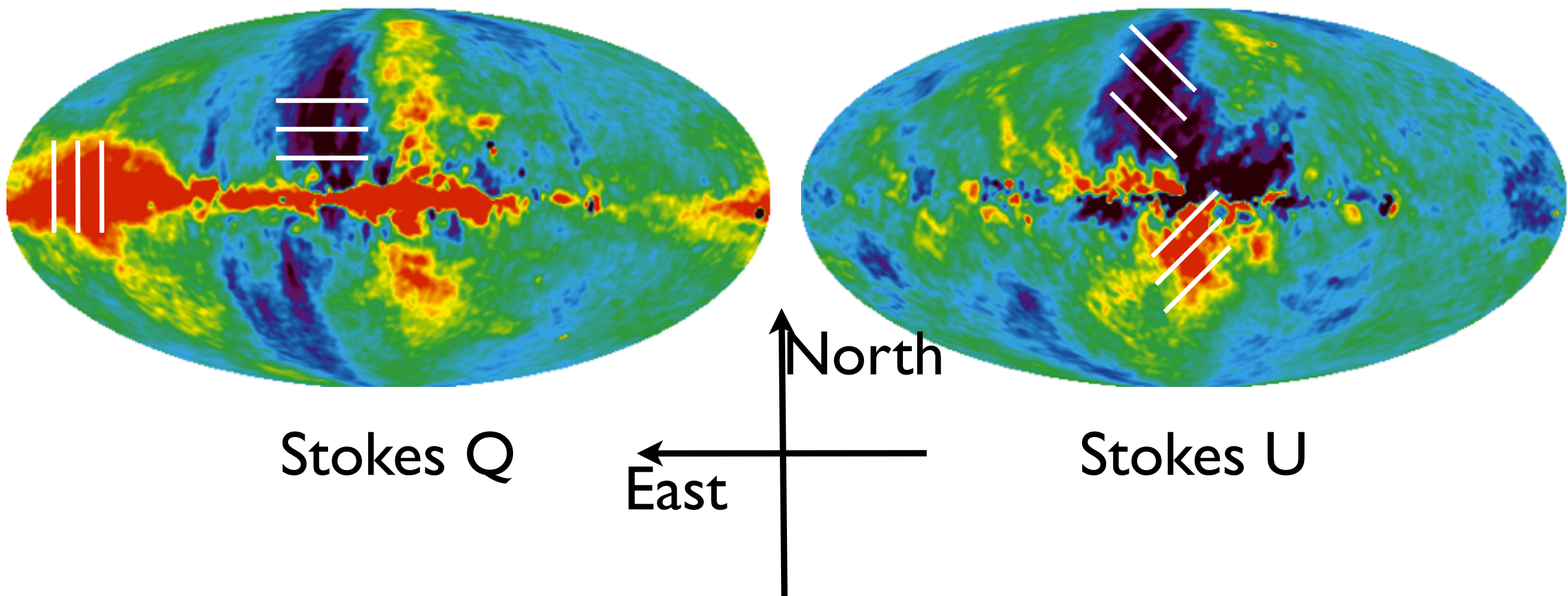


Stokes Q



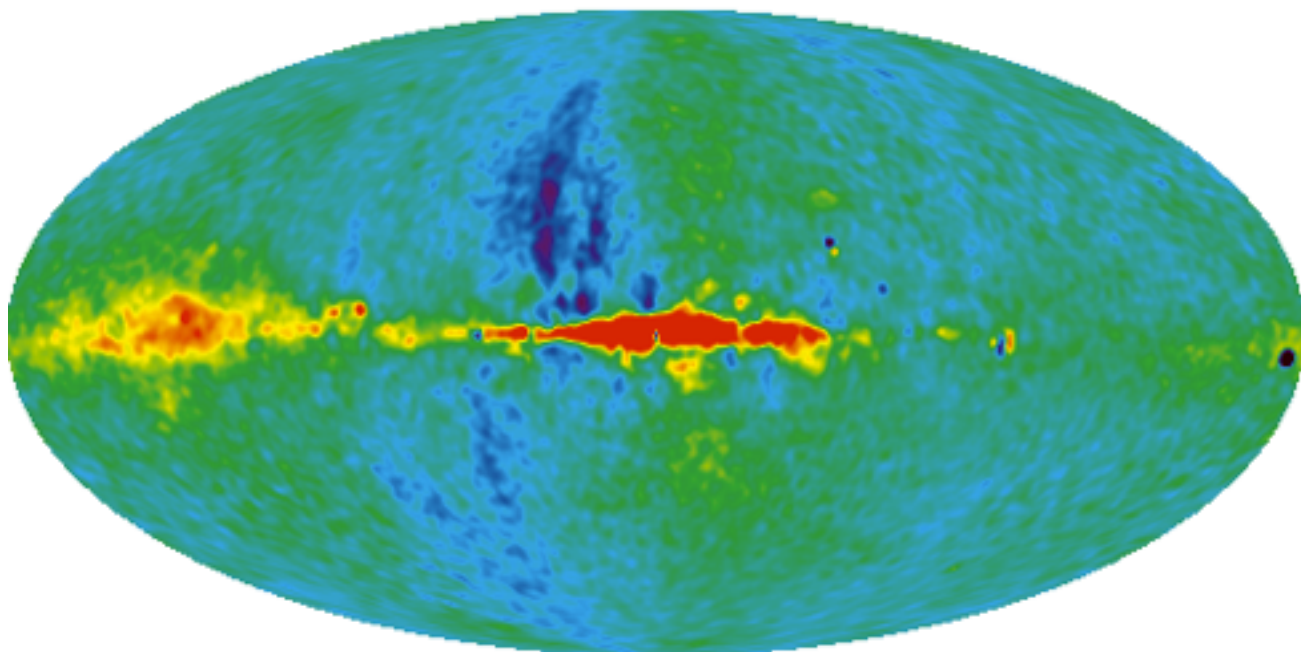
Stokes U

23 GHz

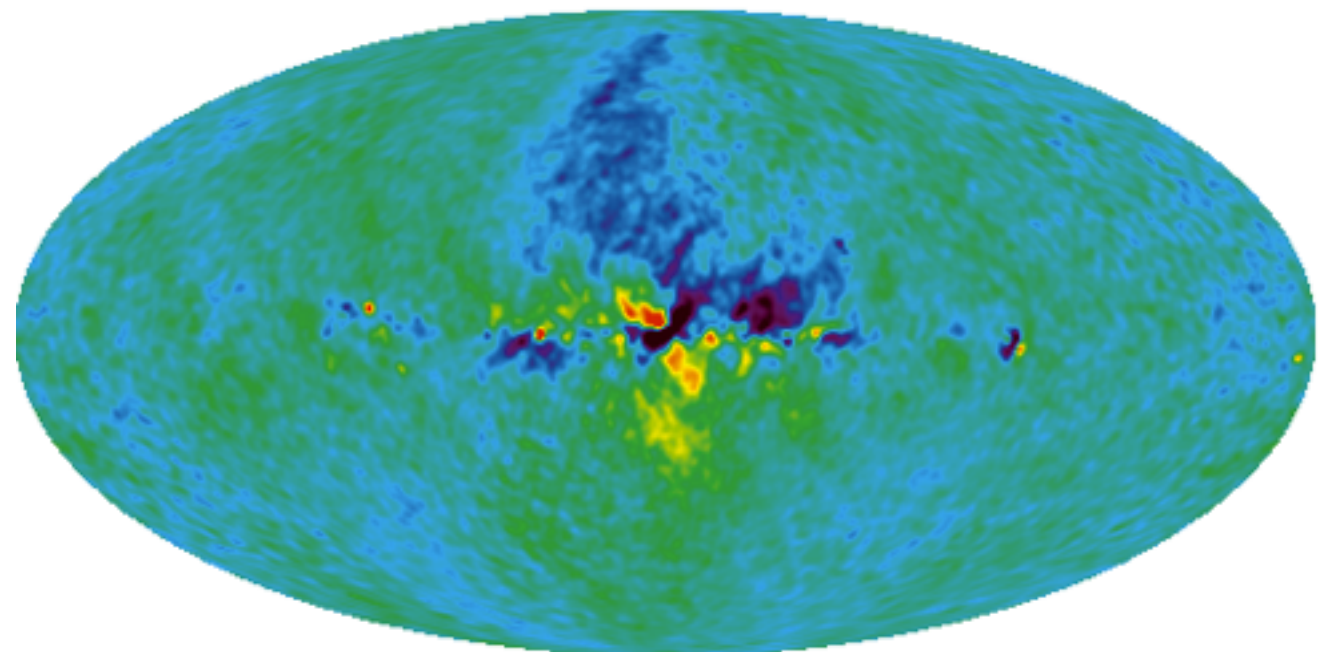




33 GHz

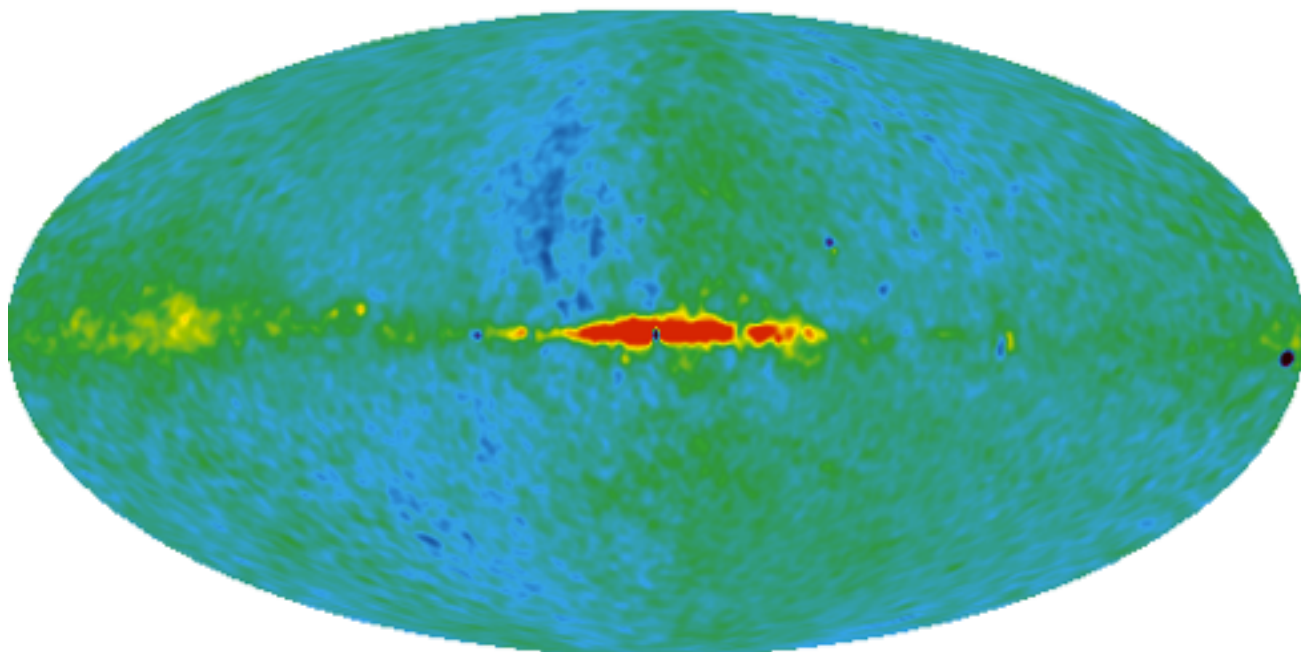


Stokes Q

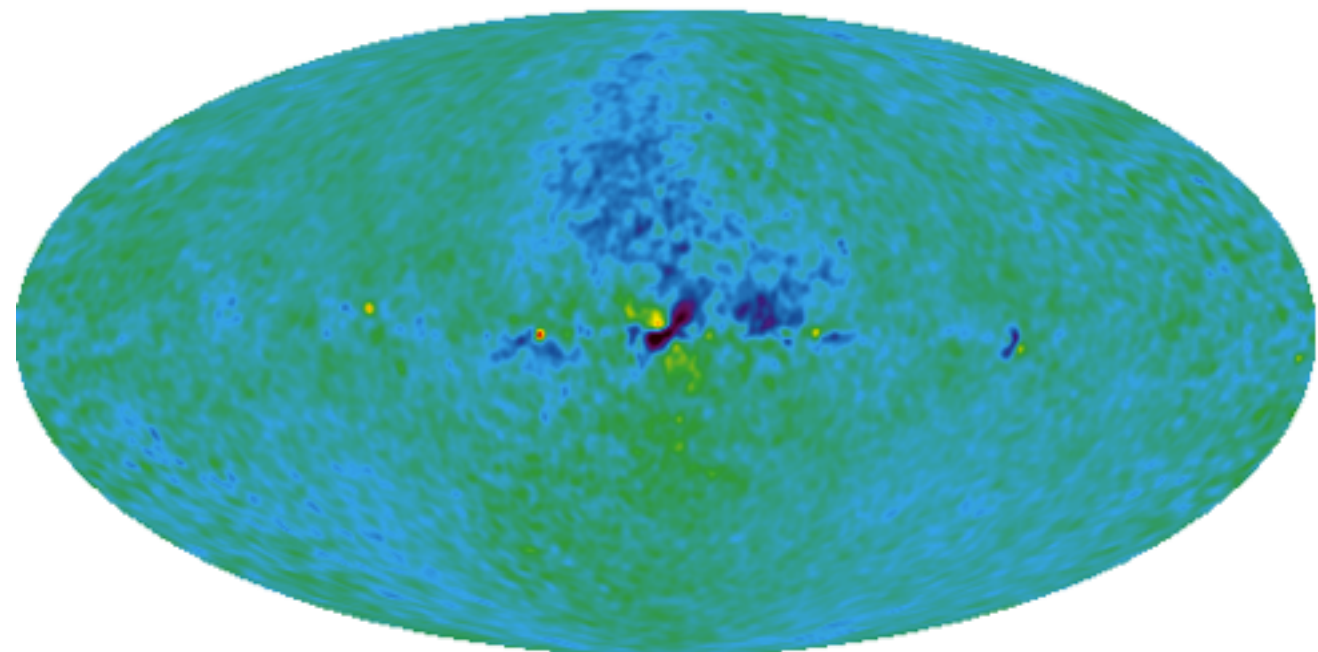


Stokes U

41 GHz

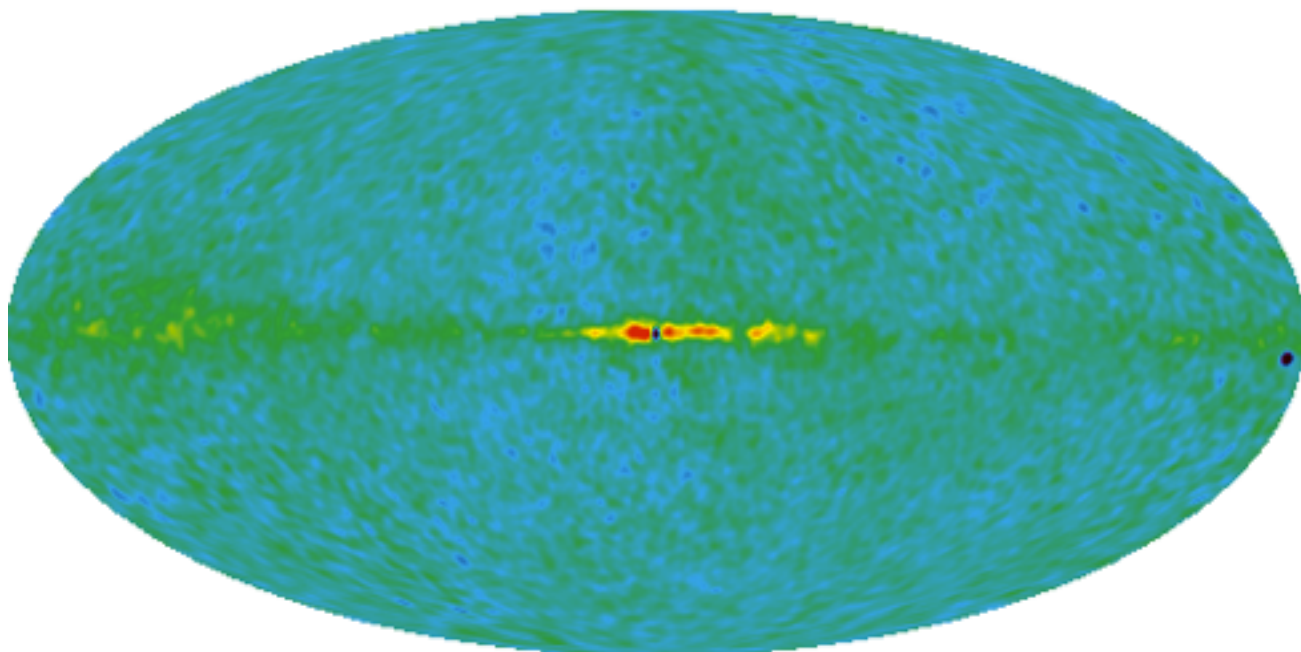


Stokes Q

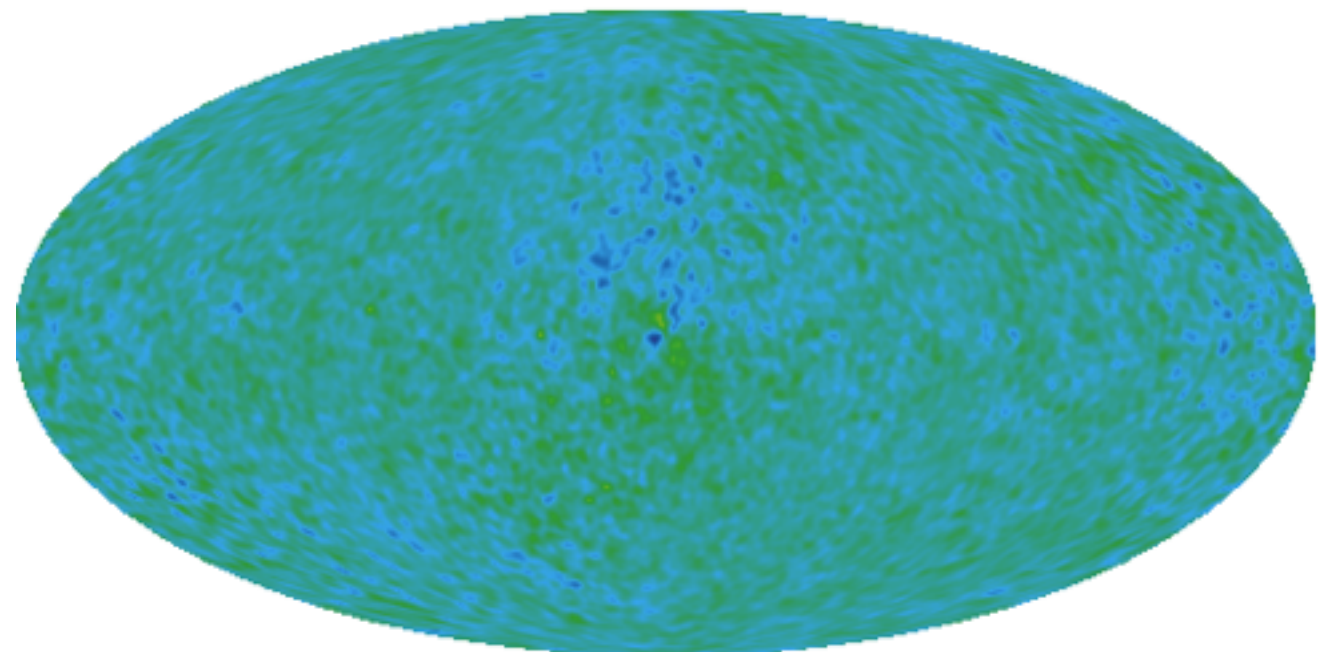


Stokes U

61 GHz



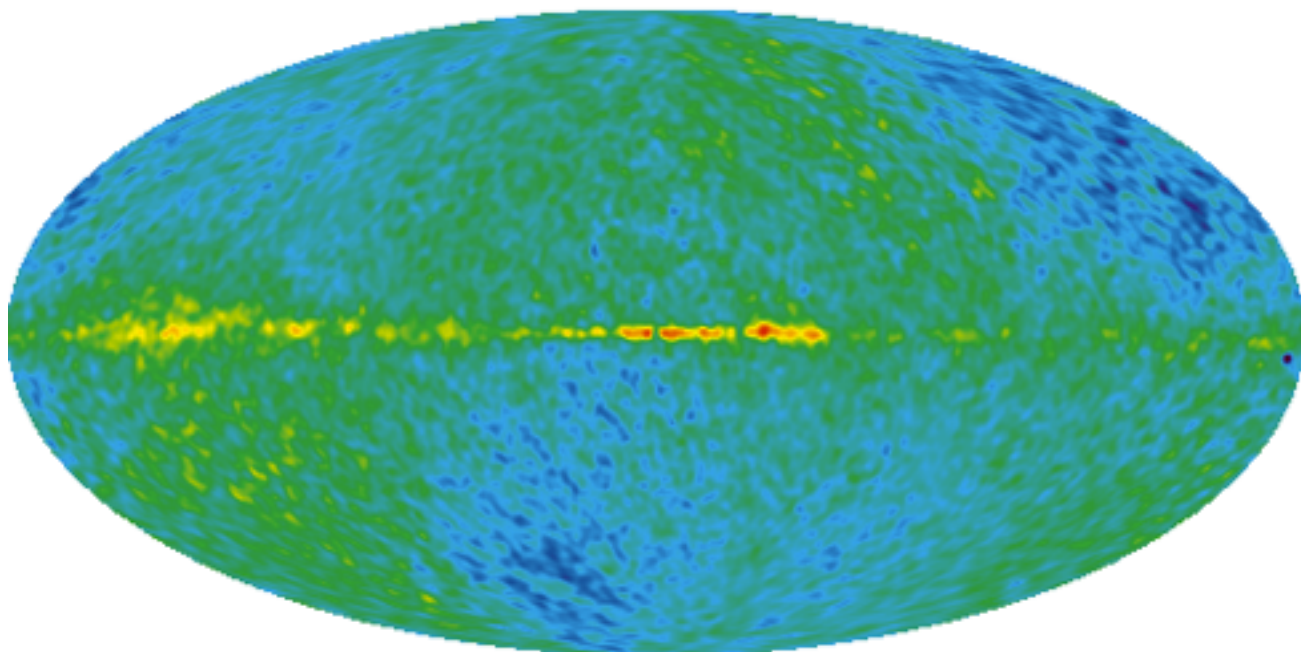
Stokes Q



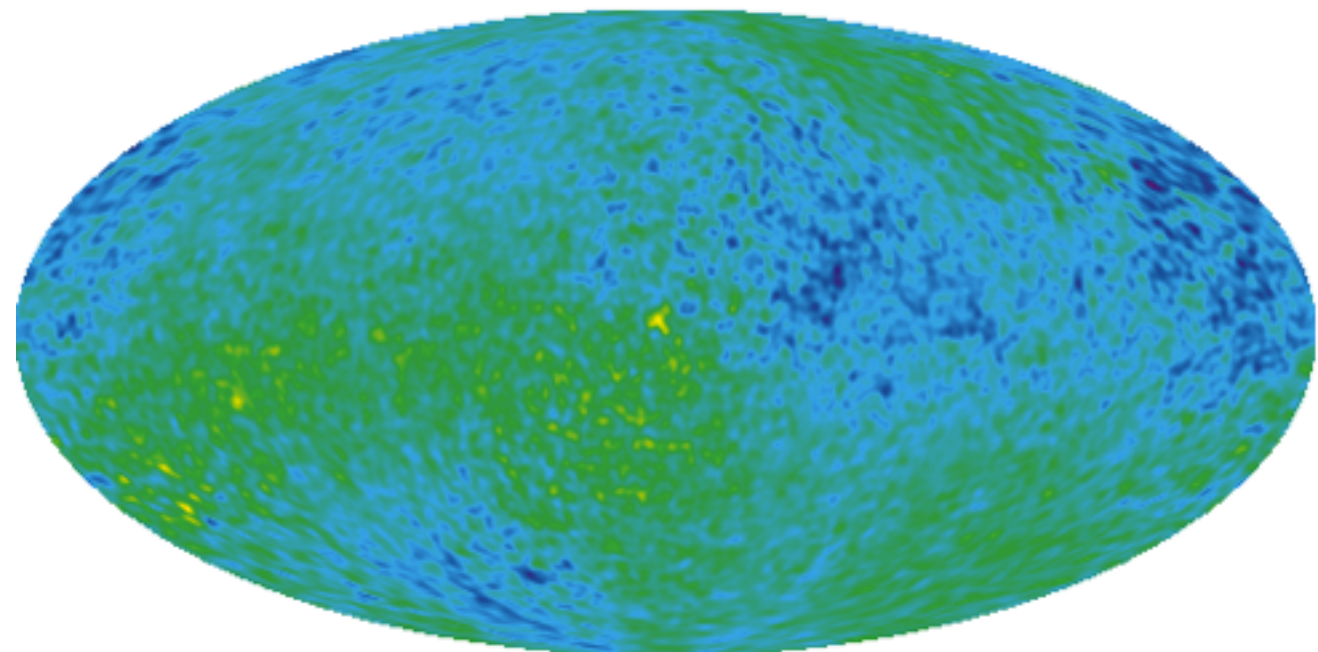
Stokes U



94 GHz



Stokes Q



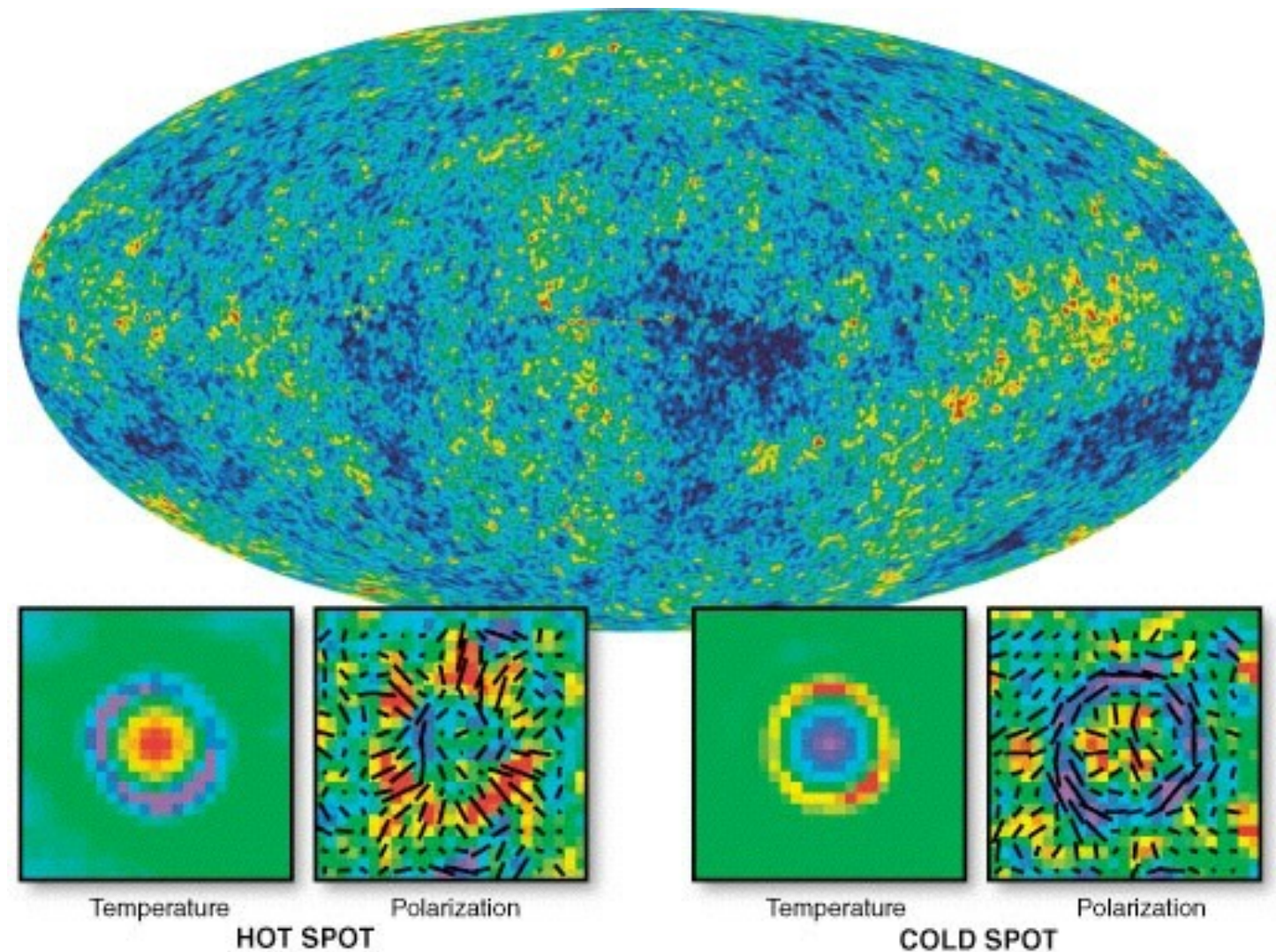
Stokes U

# How many components?

- CMB:  $T_\nu \sim \nu^0$
- Synchrotron:  $T_\nu \sim \nu^{-3}$
- Dust:  $T_\nu \sim \nu^2$
- Therefore, we need **at least** 3 frequencies to separate them

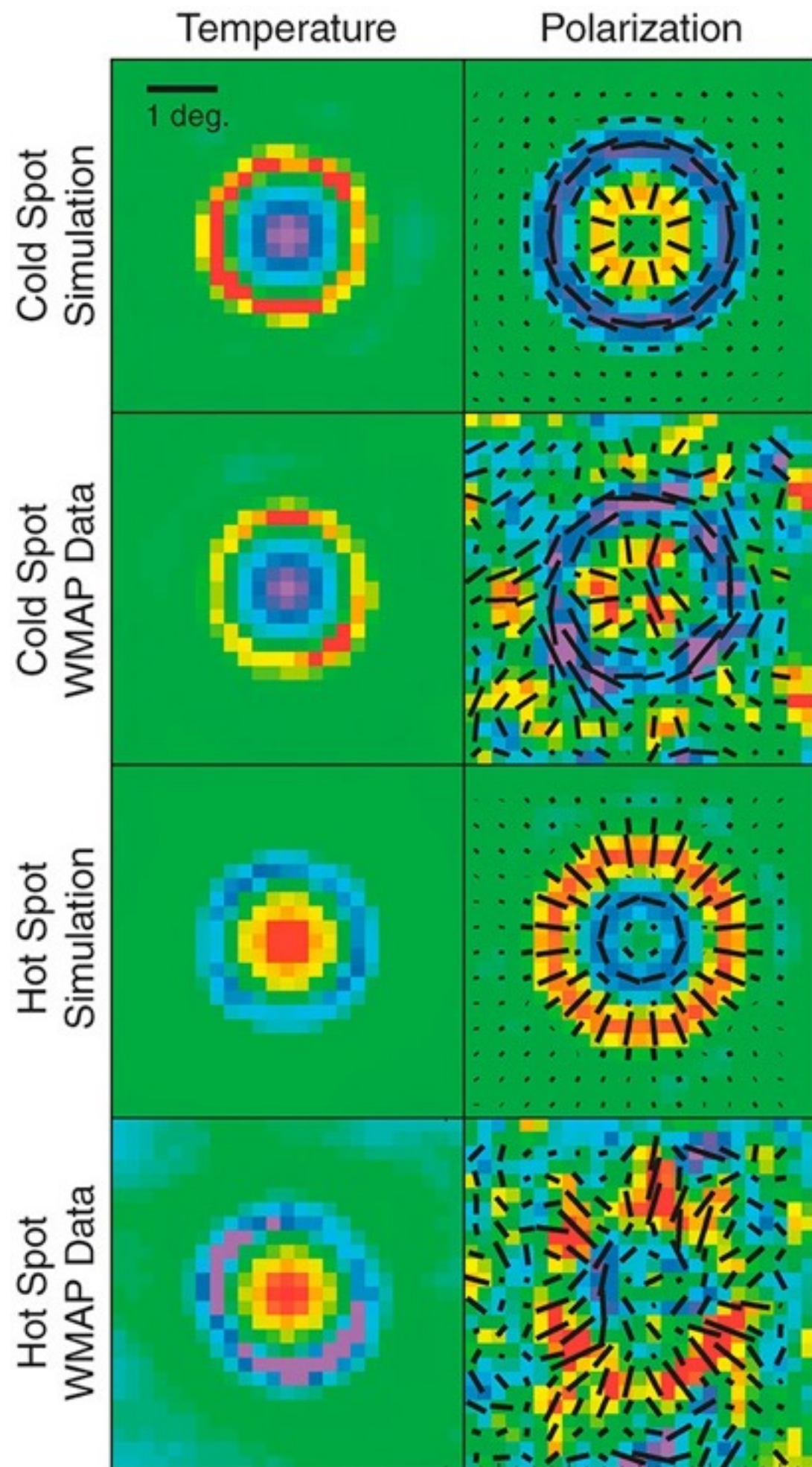
# Seeing polarisation in the WMAP data

- Average polarisation data around cold and hot temperature spots
- Outside of the Galaxy mask [not shown], there are 11536 hot spots and 11752 cold spots
- Averaging them beats the noise down





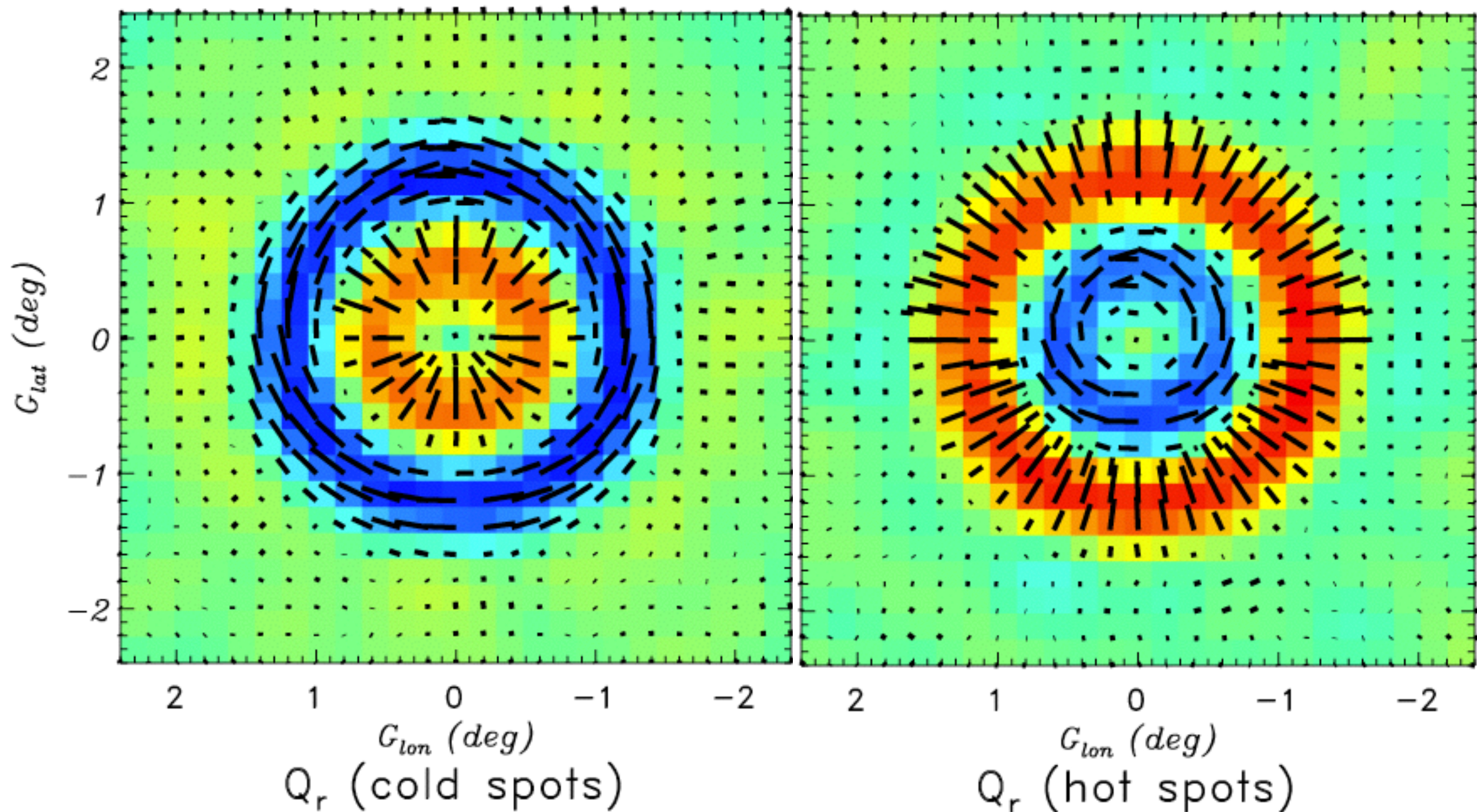
# Radial and tangential polarisation around temperature spots



- This shows polarisation generated by the plasma flowing into gravitational potentials
- Signatures of the “scalar mode” fluctuations in polarisation
- These patterns are called “**E modes**”

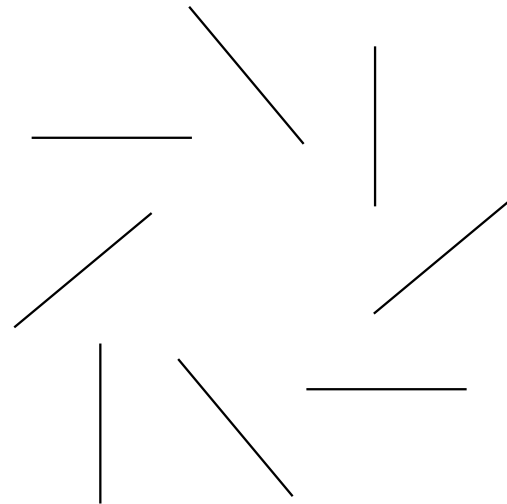
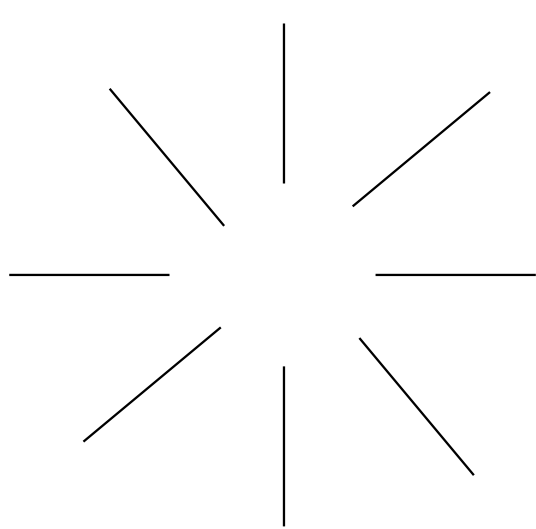


# Planck Data!

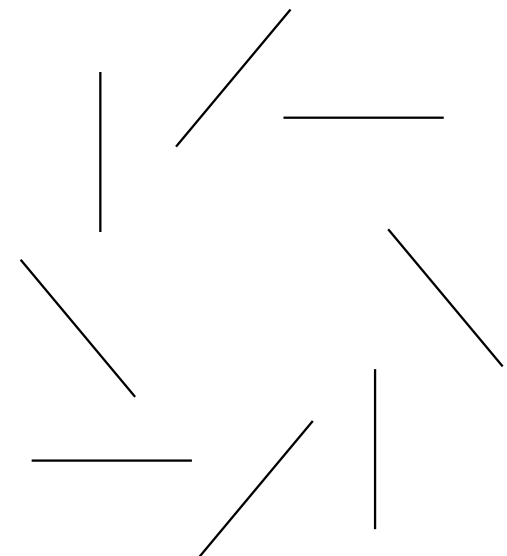
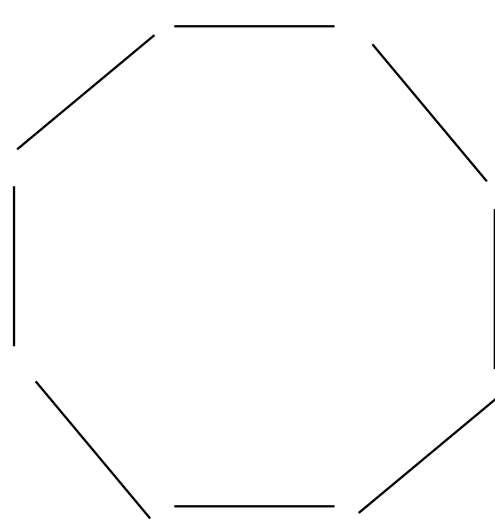




# E and B modes



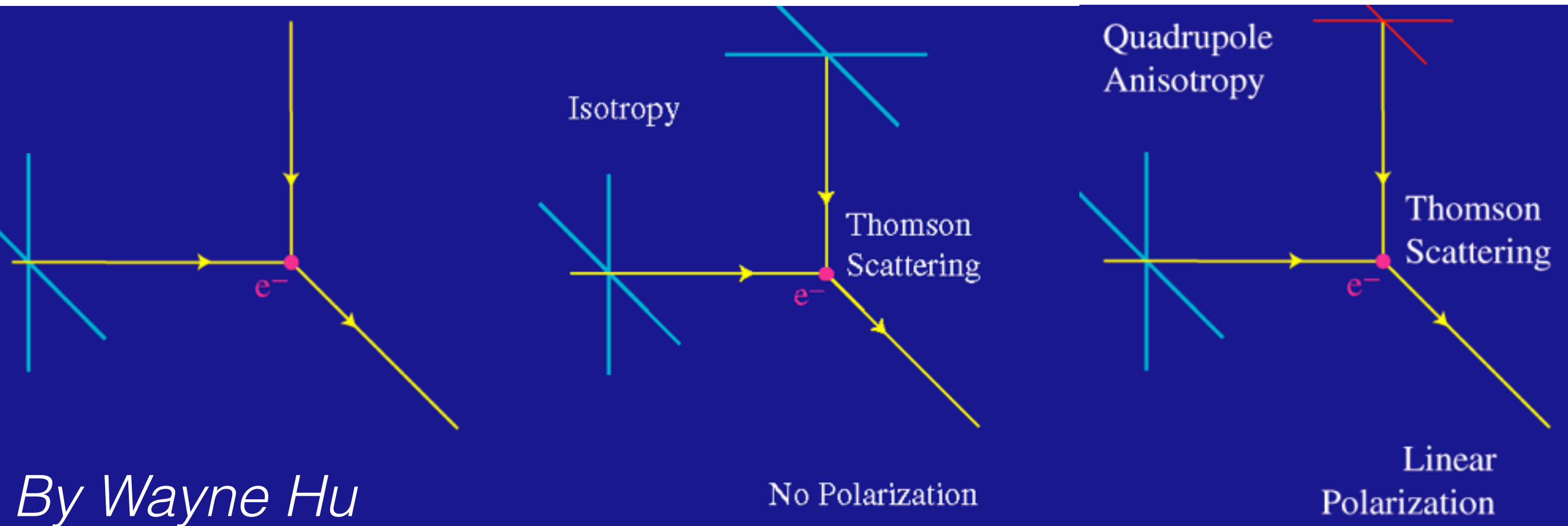
- Density fluctuations [scalar modes] can only generate E modes
- Gravitational waves can generate both E and B modes



E mode

B mode

# Physics of CMB Polarisation

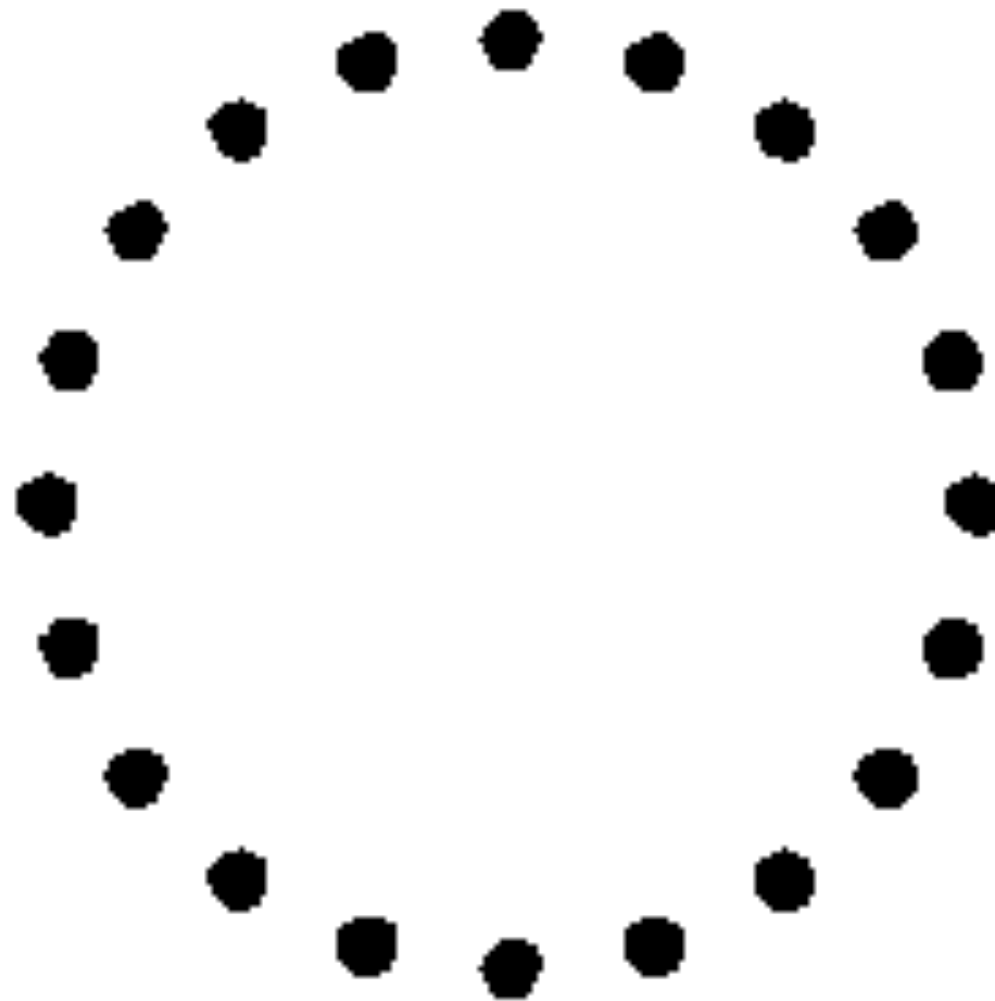


- Necessary and sufficient conditions for generating polarisation in CMB:
  - Thomson scattering
  - Quadrupolar temperature anisotropy around an electron

# Origin of Quadrupole

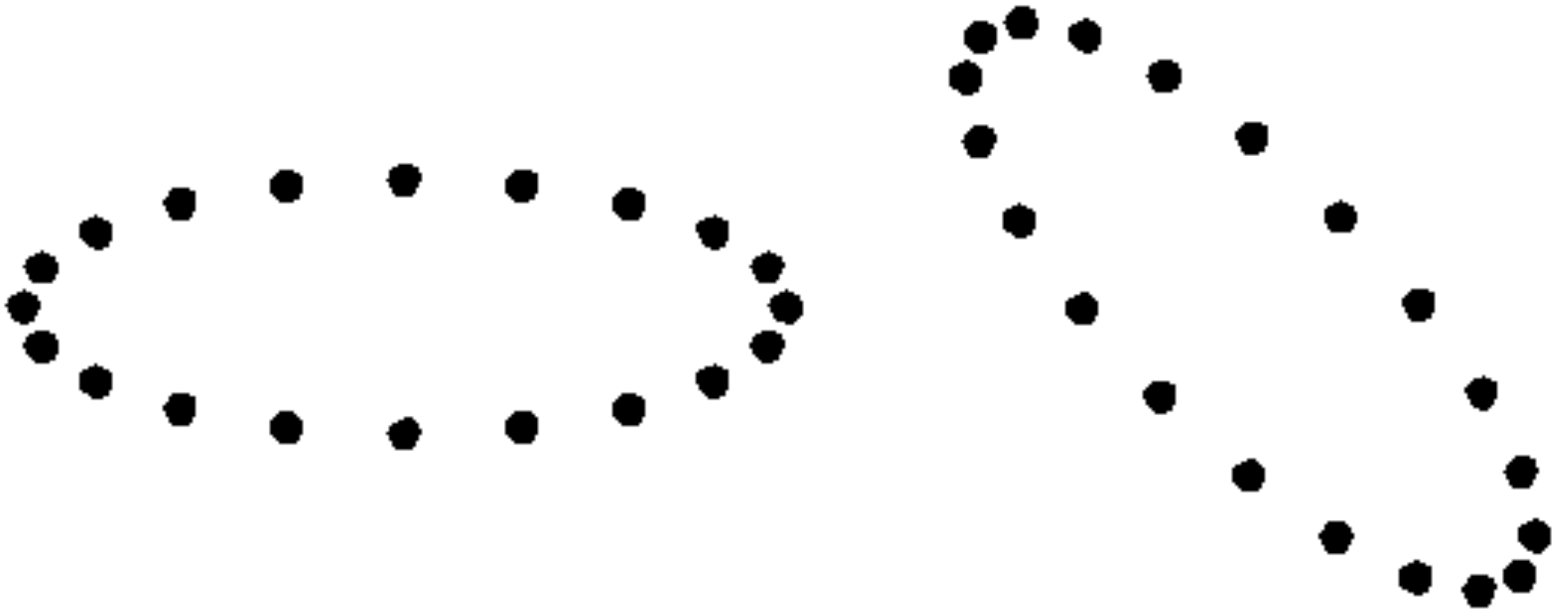
- **Scalar perturbations:** motion of electrons with respect to photons
- **Tensor perturbations:** gravitational waves

# Gravitational waves are coming toward you!



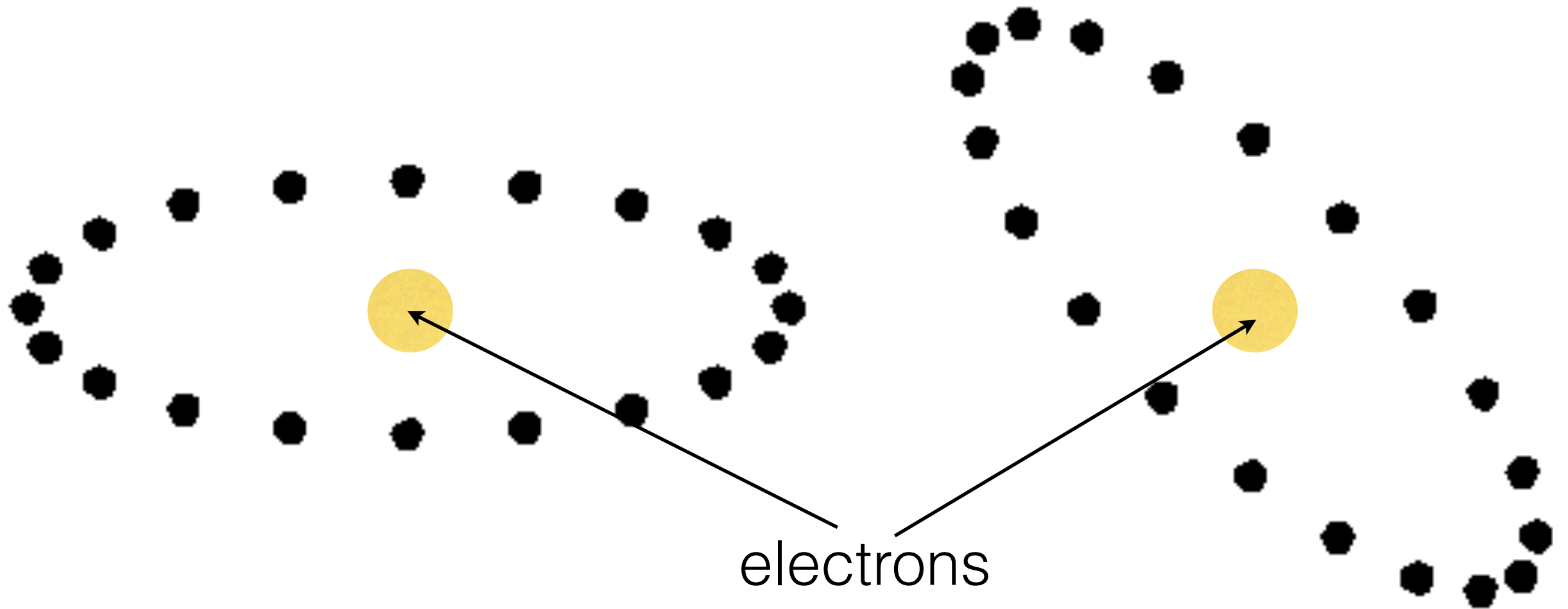
- What do they do to the distance between particles?

# Two GW modes

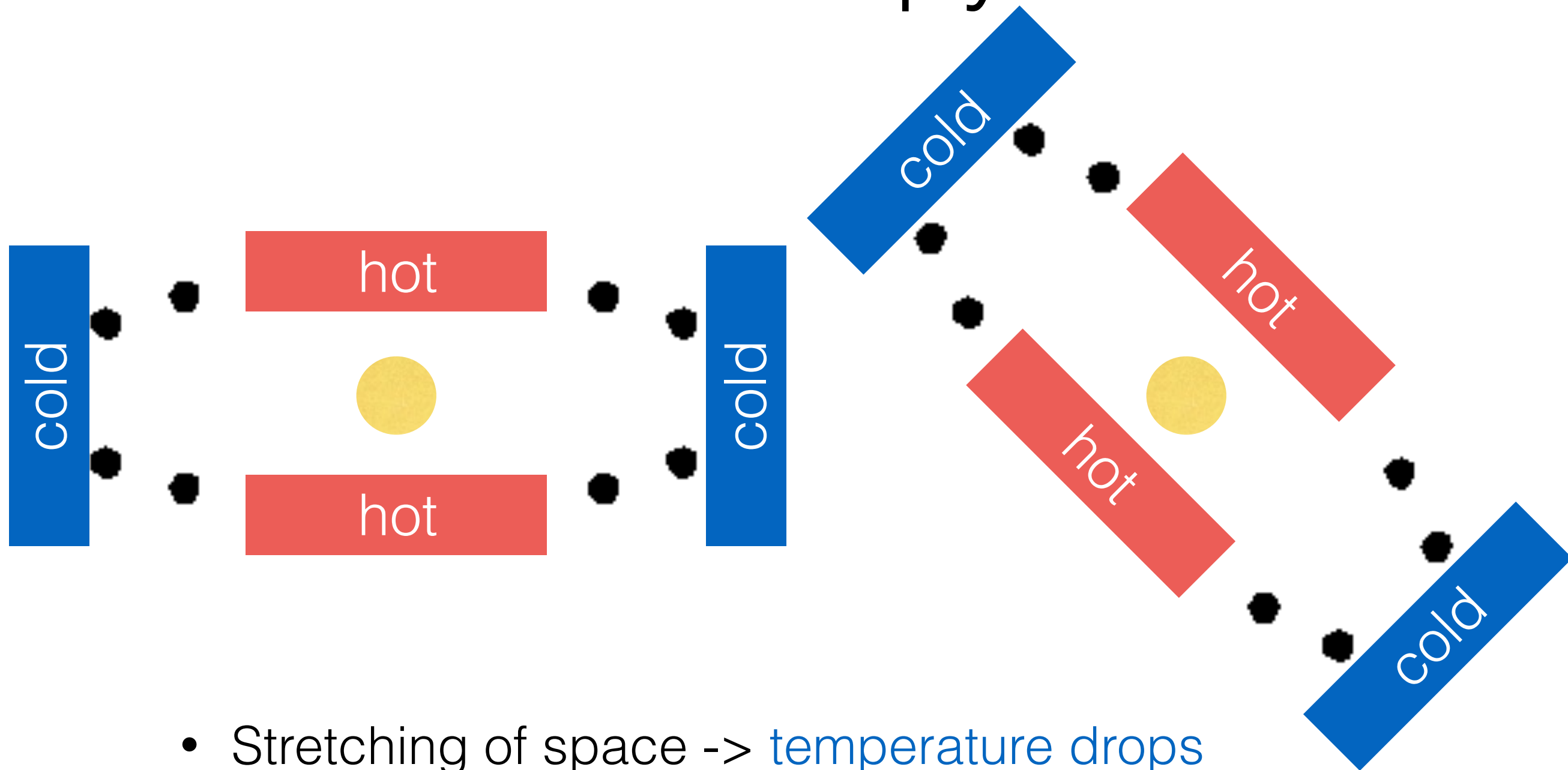


- Anisotropic stretching of space generates quadrupole temperature anisotropy. How?

# GW to temperature anisotropy

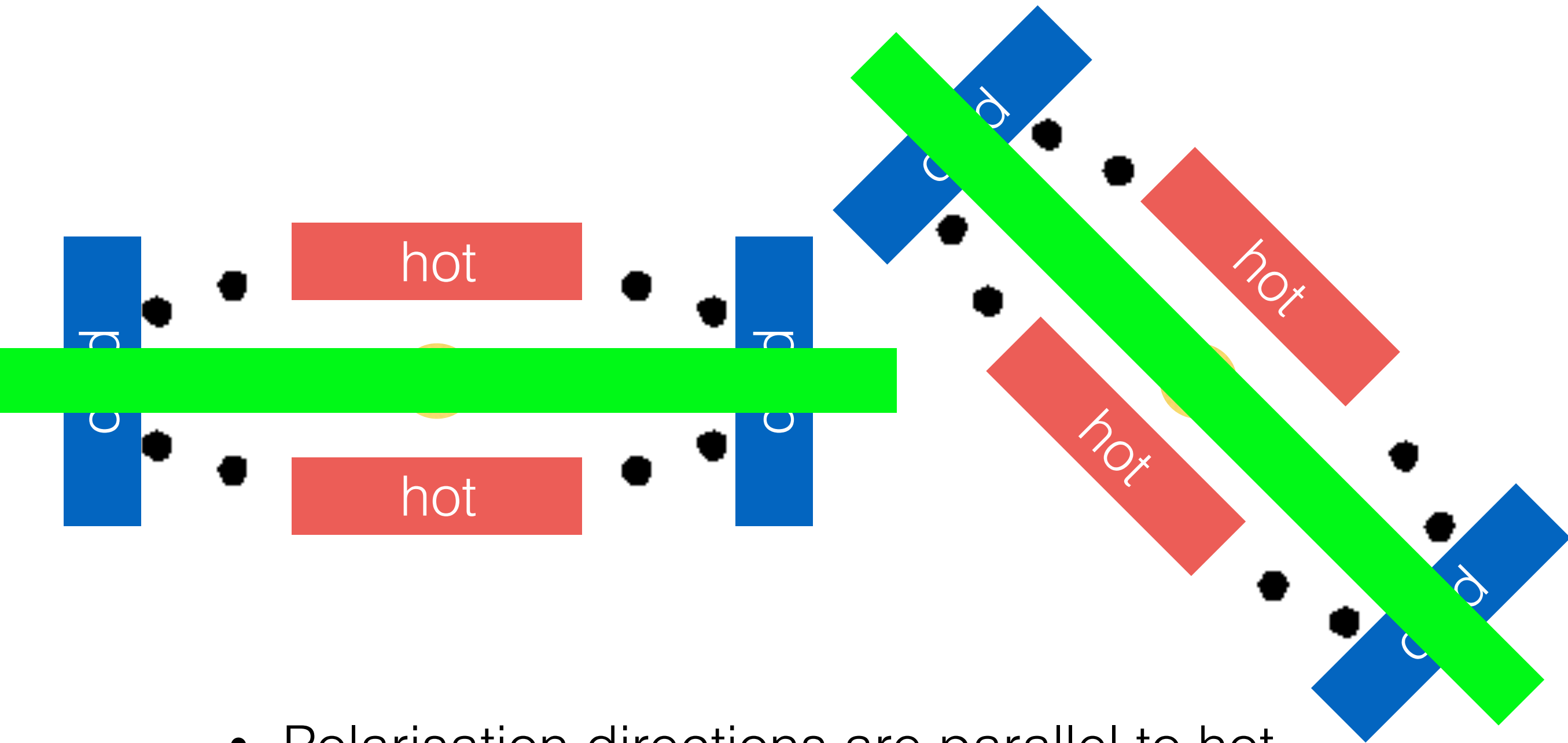


# GW to temperature anisotropy



- Stretching of space -> temperature drops
- Contraction of space -> temperature rises

# Then to polarisation!



- Polarisation directions are parallel to hot regions



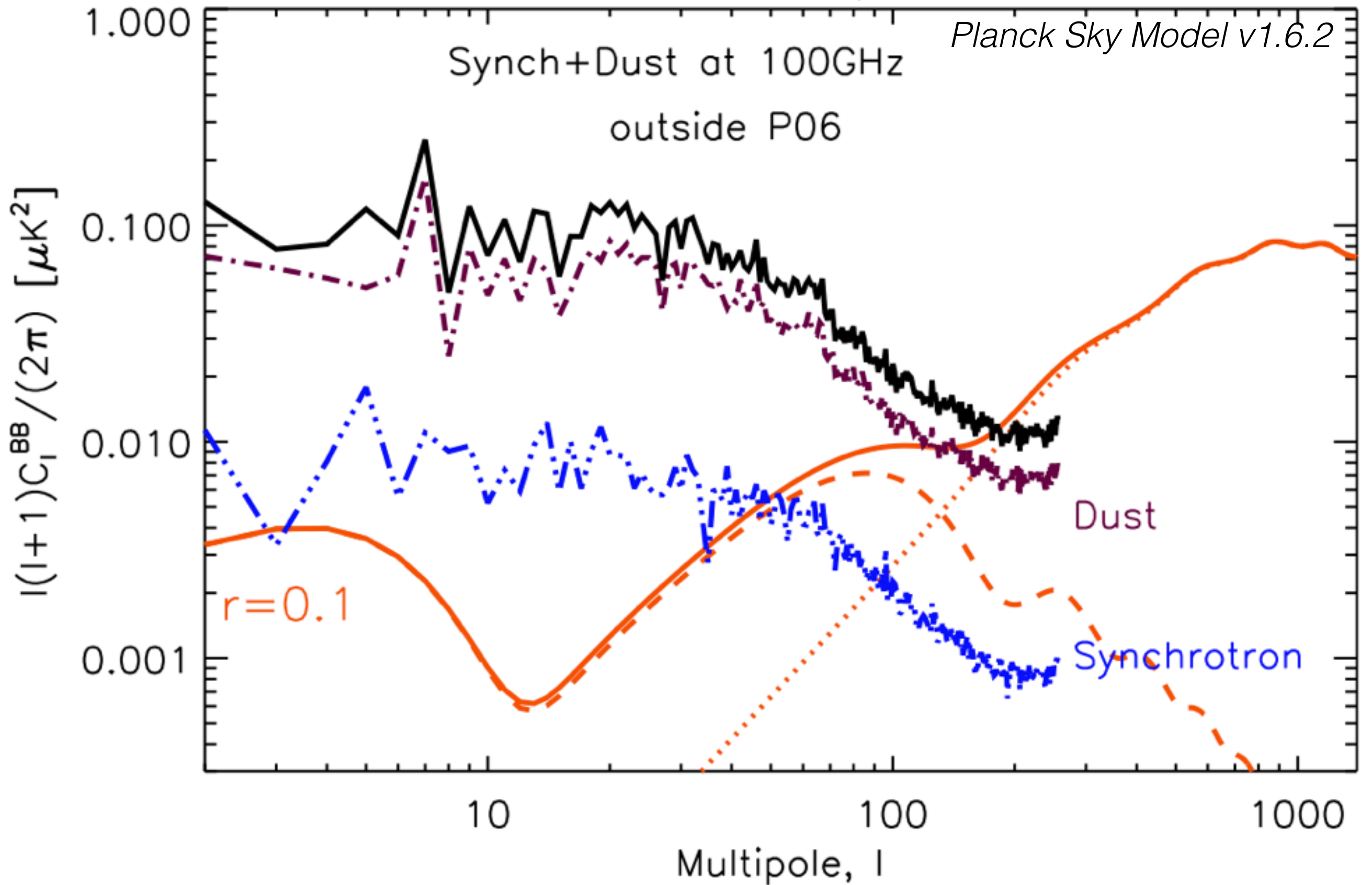
# Important note:

- Definition of  $h_+$  and  $h_x$  depends on coordinates, but definition of E- and B-mode polarisation does not depend on coordinates
- Therefore,  $h_+$  does not always give E;  $h_x$  does not always give B
- The important point is that  **$h_+$  and  $h_x$  always coexist**. When a linear combination of  $h_+$  and  $h_x$  produces E, another combination produces B

# Tensor-to-scalar Ratio

$$r \equiv \frac{\langle h_{ij} h^{ij} \rangle}{\langle \zeta^2 \rangle}$$

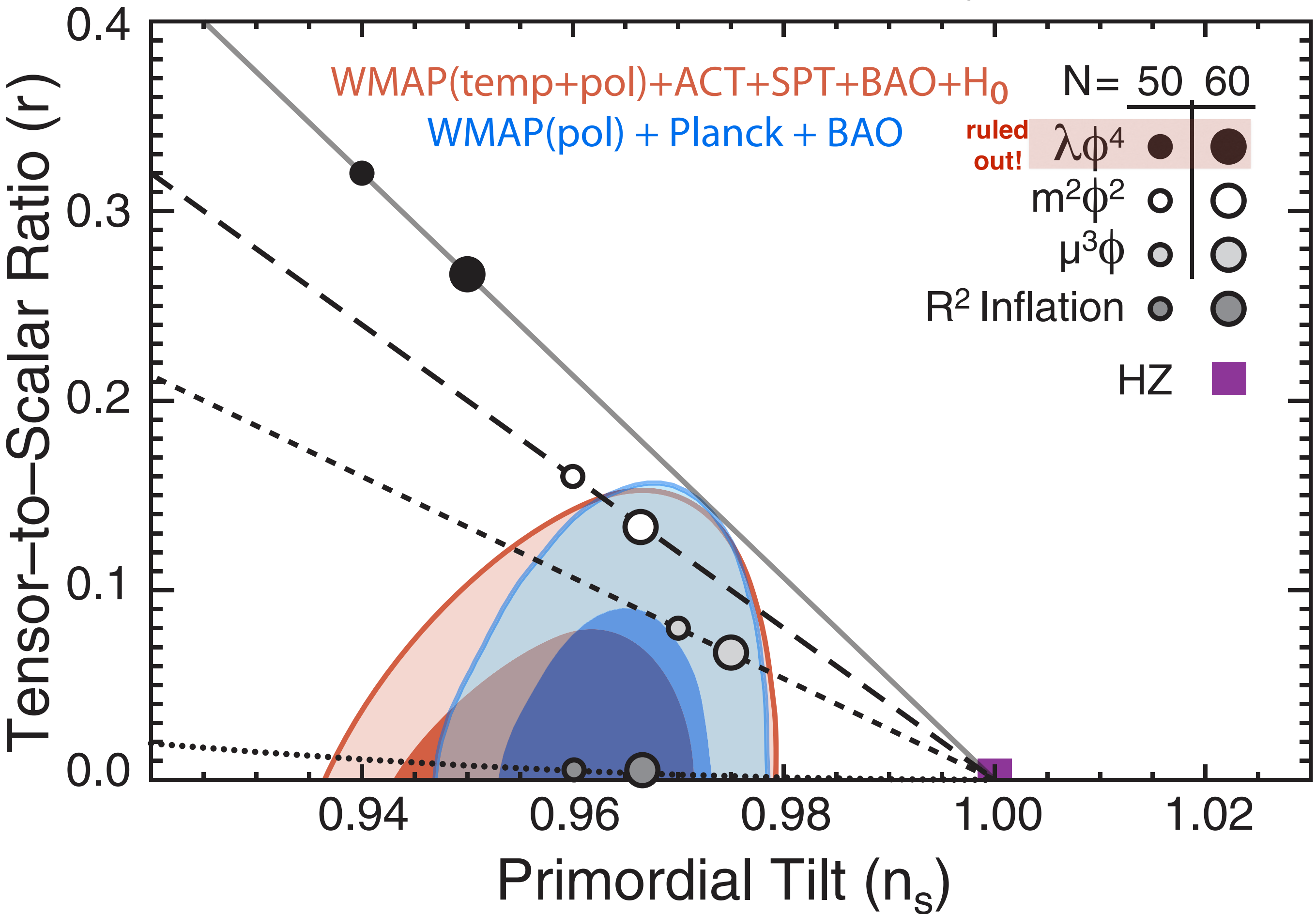
- We really want to find this quantity!
- **The upper bound from the temperature anisotropy data:  $r < 0.1$  [WMAP & Planck]**



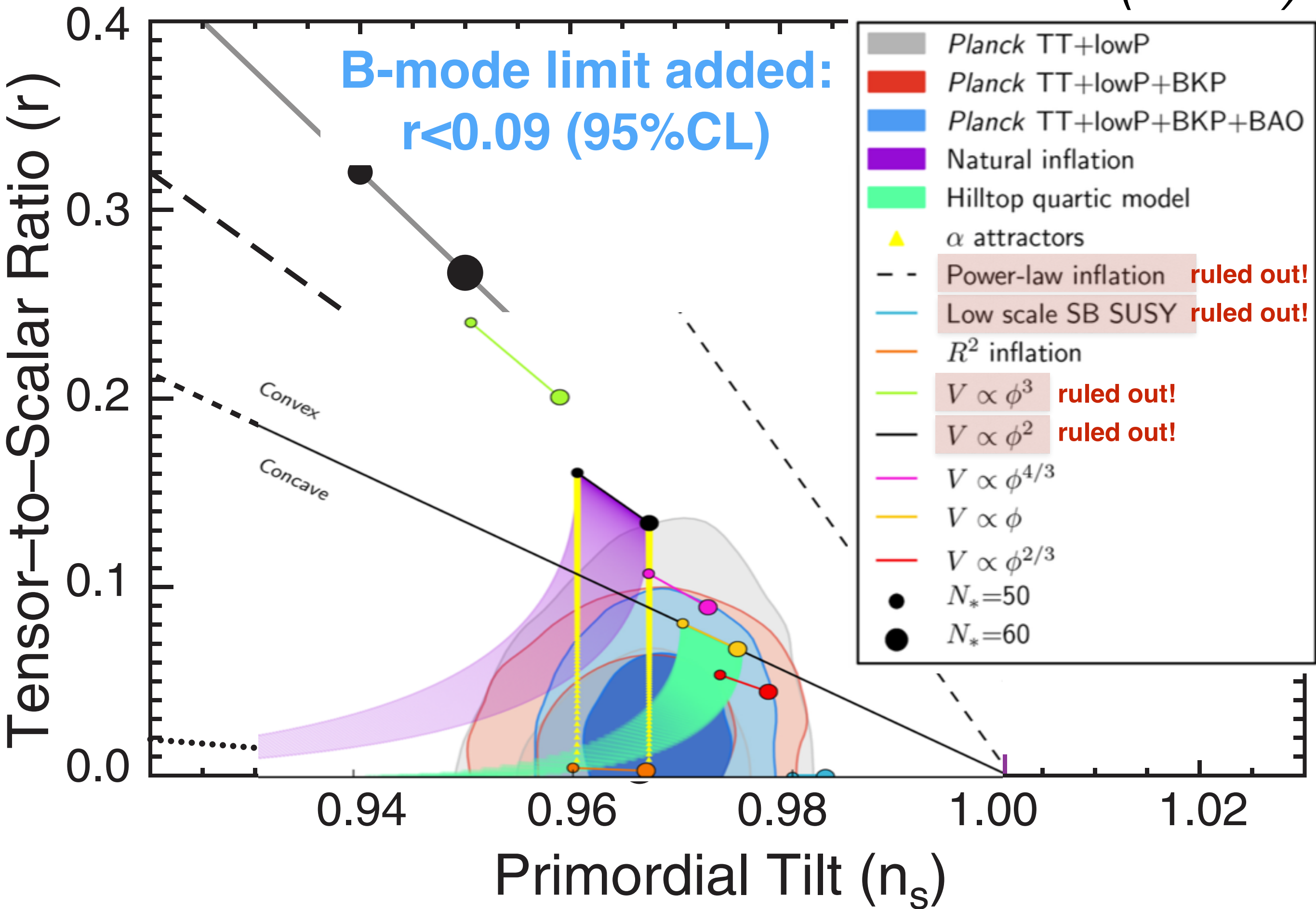
- At 100 GHz, the total foreground emission is a couple of orders of magnitude bigger in power at  $l < 10$

# If B-mode is found...

- Then what?
- The next step is to nail the specific model of inflation



# Planck Collaboration (2015)



# Summary

- *Why B-mode?*
  - Definitive evidence for inflation by showing  $dH/dt$  is small, hence **the accelerating universe!**
- We must show that the primordial gravitational waves have a **near scale-invariant power spectrum**
- Of course we need to find it first. **Challenges: foreground and lensing.** Both topics will be discussed extensively throughout this workshop!