



A02

# Dark Matter and Large-Scale Structure

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@Kavli IPMU

Fuminobu Takahashi  
(Tohoku U)

# Members

Masahiro Kawasaki (ICRR)

CDM/baryon isocurv. 21cm, CMB

Kazunori Kohri (KEK)

Decaying DM, CHAMP, MaVaN

Atsushi Taruya (Kyoto)

Structure formation history

Non-linear gravitational evolution

Fuminobu Takahashi

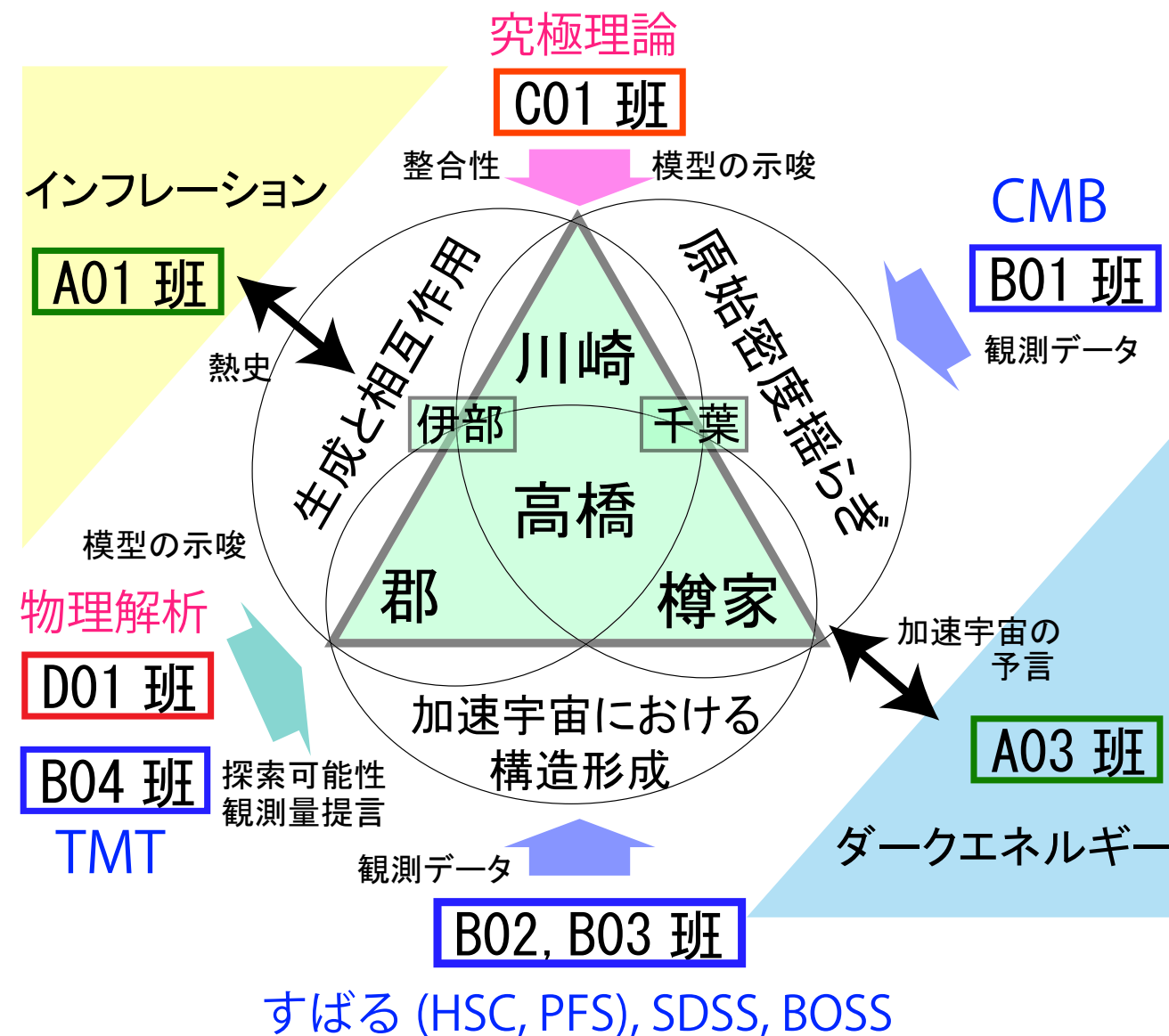
(Tohoku) Axion, sterile nu, dark photons

連携研究者：

Masashi Chiba (Tohoku)

DM distribution

Masahiro Ibe (ICRR) DM model building

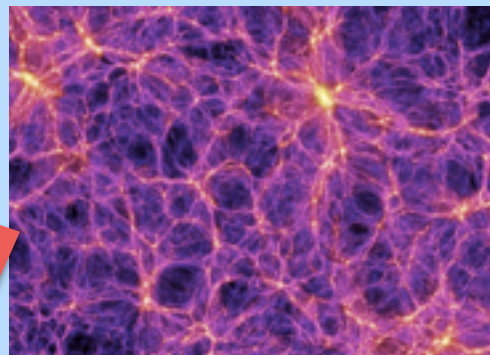
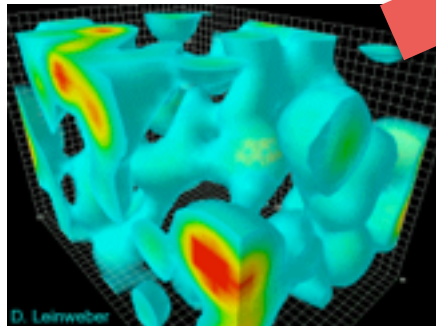


Scale factor

$$\ddot{a} > 0$$

$$\ddot{a} < 0$$

$$\ddot{a} > 0$$



Time

Accelerated  
expansion

Decelerated  
expansion

Accelerated  
expansion

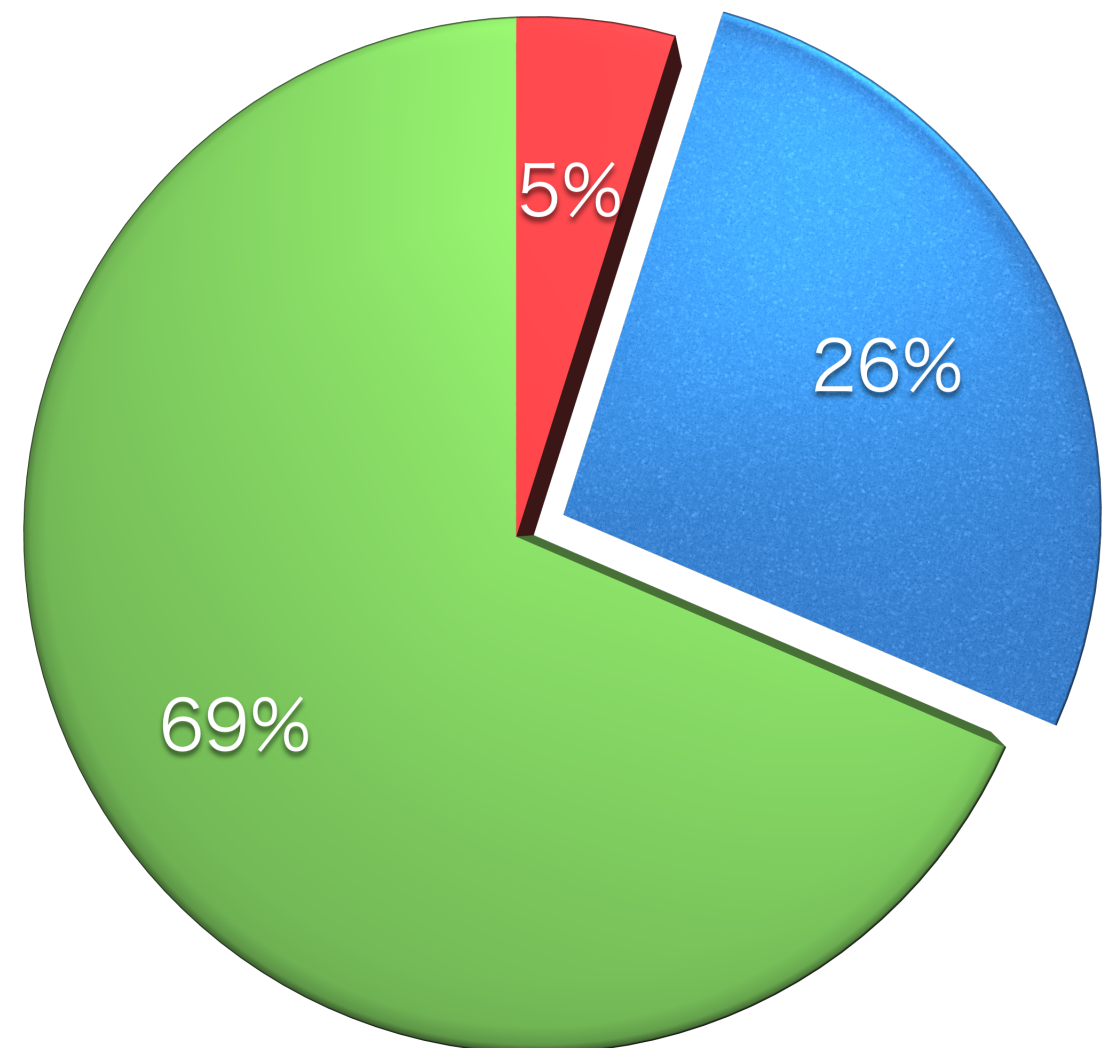
# Dark Matter

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The presence of DM has been firmly established.

$$\Omega_{\text{DM}} \simeq 0.26$$

DM may be made of an as-yet-undiscovered particle.



- Baryon
- Dark matter
- Dark energy



# What we know

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- ✓ DM gravitates.
- ✓ DM is electrically neutral.
- ✓ DM has only (very) weak interactions with the SM particles.
- ✓ DM is non-relativistic.
- ✓ DM is long-lived.
- ✓ DM played a crucial role for LSS formation.

# What we want to know

## ✓ What is DM?

- Why long-lived?
- How produced? Abundance?
- Cold or warm?
- Non-gravitational couplings?

## ✓ DM density perturbations

- Isocurvature mode/non-Gaussianity?

## ✓ Non-linear gravitational evolution

## ✓ DM distribution in dSphs, galaxies, and clusters. See talk by Taruya

- Any tension with LCDM and N-body simulations?

See talk by Chiba

# Research Themes

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## (1) Dark matter density perturbations

Kawasaki

Isocurvature perturbations, non-Gaussianity

Takahashi

## (2) Production and properties of dark matter

Ibe

Non-thermal production, self-interacting/decaying DM

## (3) Structure formation and DM distribution

Taruya

Chiba

Non-linear gravitational evolution, halo bias

Matter power spectrum modified by WDM or CHAMP

Kohri

## (4) Unification of dark matter and dark energy

Mass-Varying Neutrinos, Ultra-light axions

Kohri

## (5) Anything about DM!

# Why long-lived?

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## 1) Symmetry

- e.g. R-parity, KK parity: LSP, LKP.

## 2) Light mass

- e.g. modulus lifetime:  $\tau \simeq \frac{M_p^2}{m_{\text{DM}}^3}$

## 3) Very weak interactions

- e.g. Hidden sector, gravity sector

The light mass and/or weak int. may be due to symmetry (e.g. chiral, shift sym, SUSY), extra dim or compositeness.

These are not exclusive: e.g. axion [2&3], gravitino [1,2,3]

# How produced?

1) From the standard model thermal plasma

- DM was in thermal equilibrium **Neutralino LSP**

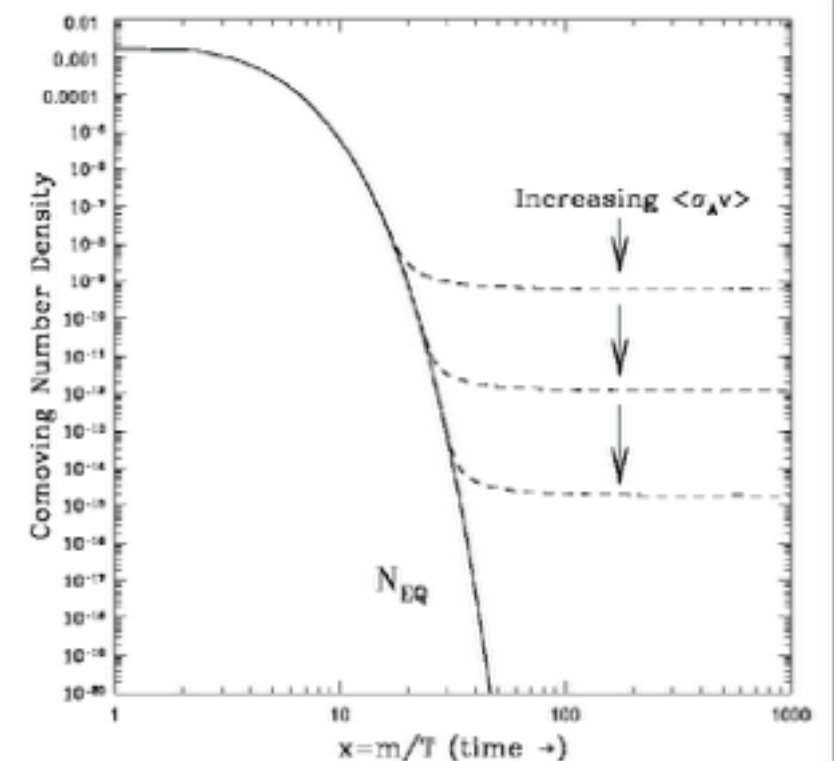
## “WIMP miracle”

- Thermal relic abundance of WIMPs of mass  $O(100)\text{GeV}$  is close to the observed DM density.

$$\Omega_{\text{WIMP}} \simeq \frac{0.3}{\langle \sigma v \rangle / \text{pb}}$$

$$\langle \sigma v \rangle \simeq 3 \times 10^{-26} \text{cm}^3/\text{sec}$$

Sounds reasonable if there is new physics around TeV scale.





# How produced?

## 1) From the standard model thermal plasma

- DM was in thermal equilibrium      **Neutralino**
- Non-thermal process; decay, oscillations, scattering.  
   **SuperWIMP, sterile nu, gravitino, etc.**

## 2) From other sources

- decay of inflaton or some other heavy particle
- scalar condensate      **QCD axion, moduli**      **Gravitino**
- hidden plasma (thermal or non-thermal)
- gravitational production
- etc.

# Cold or Warm

Free streaming length (sterile neutrino)

$$\lambda_{\text{FS}} \simeq 1 \text{ Mpc} \left( \frac{m_{\text{DM}}}{\text{keV}} \right)^{-1} \left( \frac{\langle p_{\text{DM}}/T \rangle}{3.15} \right)$$

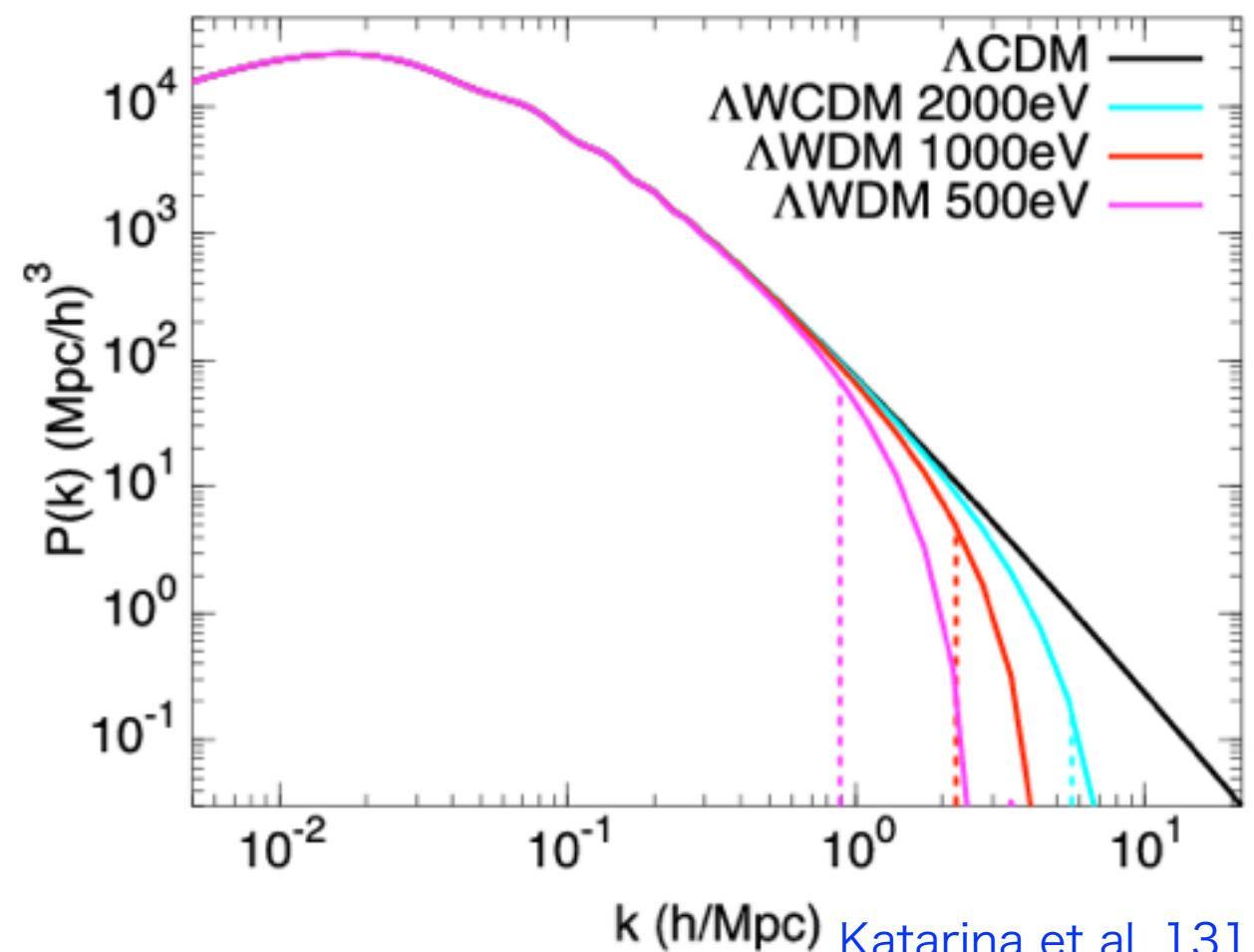
$$M_{\text{FS}} \simeq 3 \times 10^7 M_{\odot} \left( \frac{m_{\text{DM}}}{\text{keV}} \right)^{-3} \left( \frac{\langle p_{\text{DM}}/T \rangle}{3.15} \right)^3$$

DM mass

Production process

WDM suppresses power on small scales.

Similar suppression is possible with long-lived CHAMP.



cf. Kohri & T.Takahashi, 0909.4610

Katarina et al, 1311.5223

# Sterile neutrino DM

Let us consider SM + 3 RH neutrinos.

$$-\mathcal{L} = \frac{1}{2} \kappa_I M \bar{N}_I^c N_I + \lambda_{I\alpha} \bar{N}_I L_\alpha H + \text{h.c.},$$

$\alpha = e, \mu, \tau$

$$\kappa_3 > \kappa_2 > \kappa_1 \quad M_I = \kappa_I M \quad M : \text{B-L breaking scale}$$

The seesaw formula for the light neutrino mass:

Minkowski '77, Yanagida '79, Ramond '79, Glashow '80

$$(m_\nu)_{\alpha\beta} = \lambda_{\alpha I} \lambda_{I\beta} \frac{v^2}{M_I}, \quad v = \langle H^0 \rangle \simeq 174 \text{ GeV}$$

$$\Delta m_{\text{atm}}^2 \simeq 2.3 \times 10^{-3} \text{ eV}^2 \quad \Delta m_{\odot}^2 \simeq 8 \times 10^{-5} \text{ eV}^2$$

The B-L breaking is  $M \sim 10^{15} \text{ GeV}$  for  $\lambda_{I\alpha} \sim \kappa_I \sim 1$

# Sterile neutrino DM

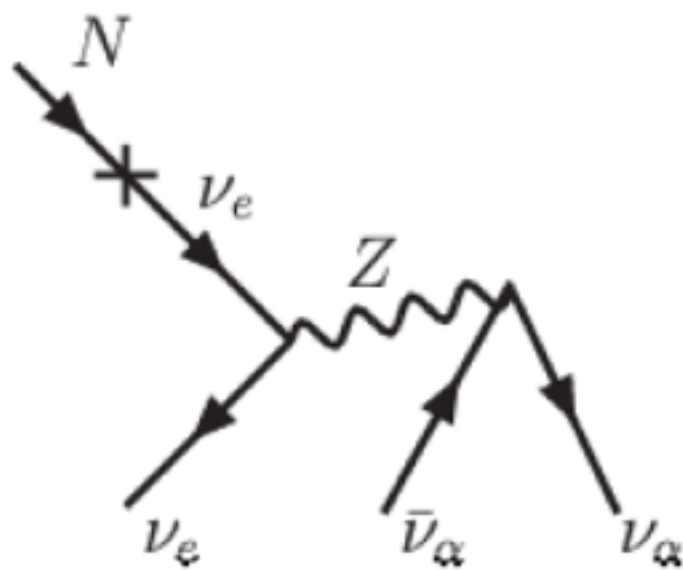
Suppose that the lightest one is very light.

$$-\mathcal{L} = \frac{1}{2} \kappa_I M \bar{N}_I^c N_I + \lambda_{I\alpha} \bar{N}_I L_\alpha H + \text{h.c.},$$

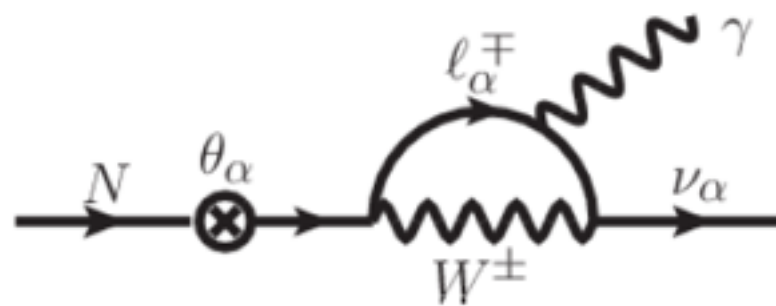
$\alpha = e, \mu, \tau$

$$\kappa_3 > \kappa_2 \gg \kappa_1 \quad M_I = \kappa_I M \quad M : \text{B-L breaking scale}$$

Sterile neutrino decays thru mixings with active nu.



$$N \rightarrow \nu \nu \bar{\nu}$$

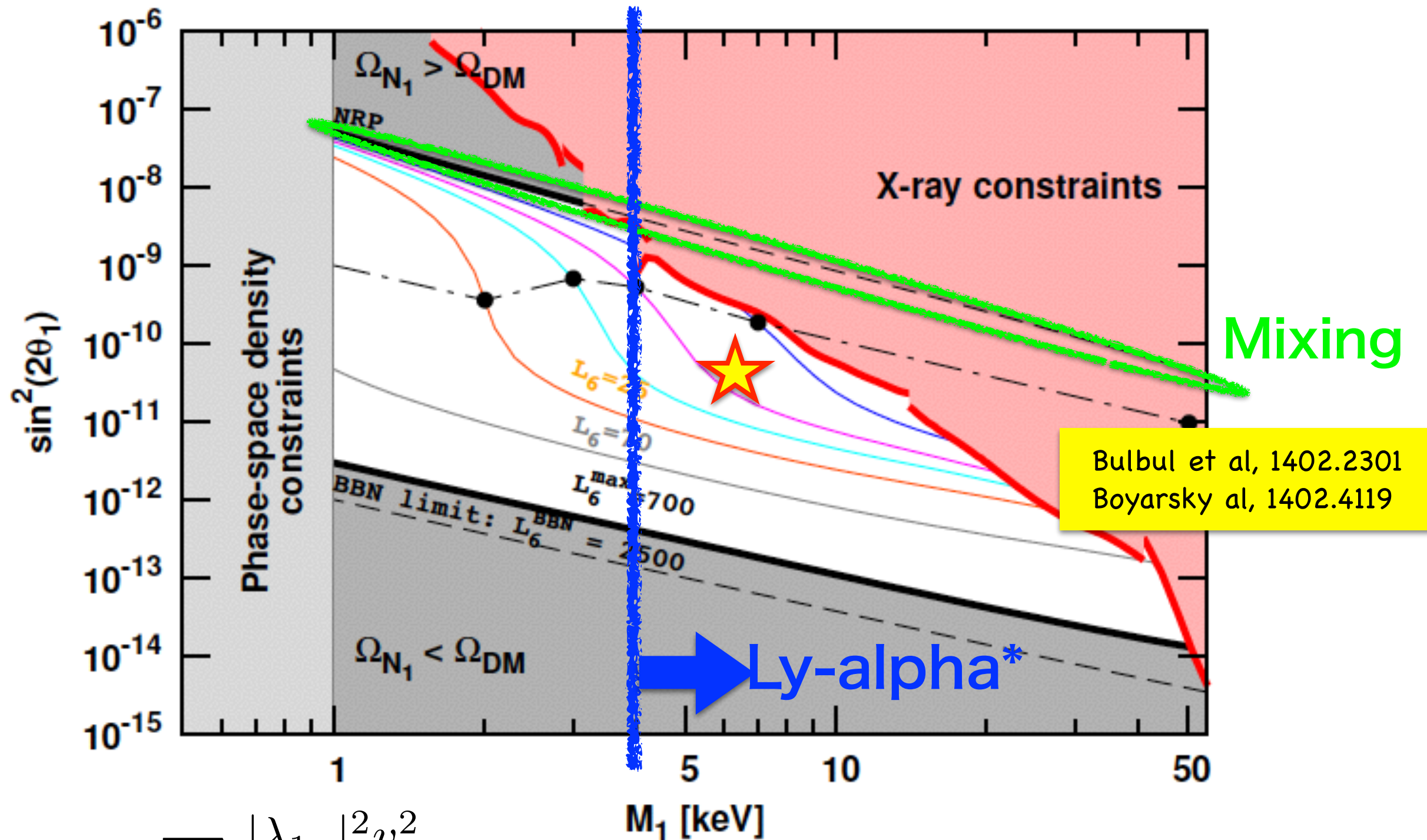


$$N \rightarrow \nu \gamma$$

$$\theta^2 \equiv \sum_{\alpha} \frac{|\lambda_{1\alpha}|^2 v^2}{M_1^2}$$

# Cosmological bounds

$$\Gamma^{-1} > 200 \text{ Gyr}$$



$$\theta^2 \equiv \sum_{\alpha} \frac{|\lambda_{1\alpha}|^2 v^2}{M_1^2}$$

From "Light Sterile Neutrinos: A White Paper" 1204.5379

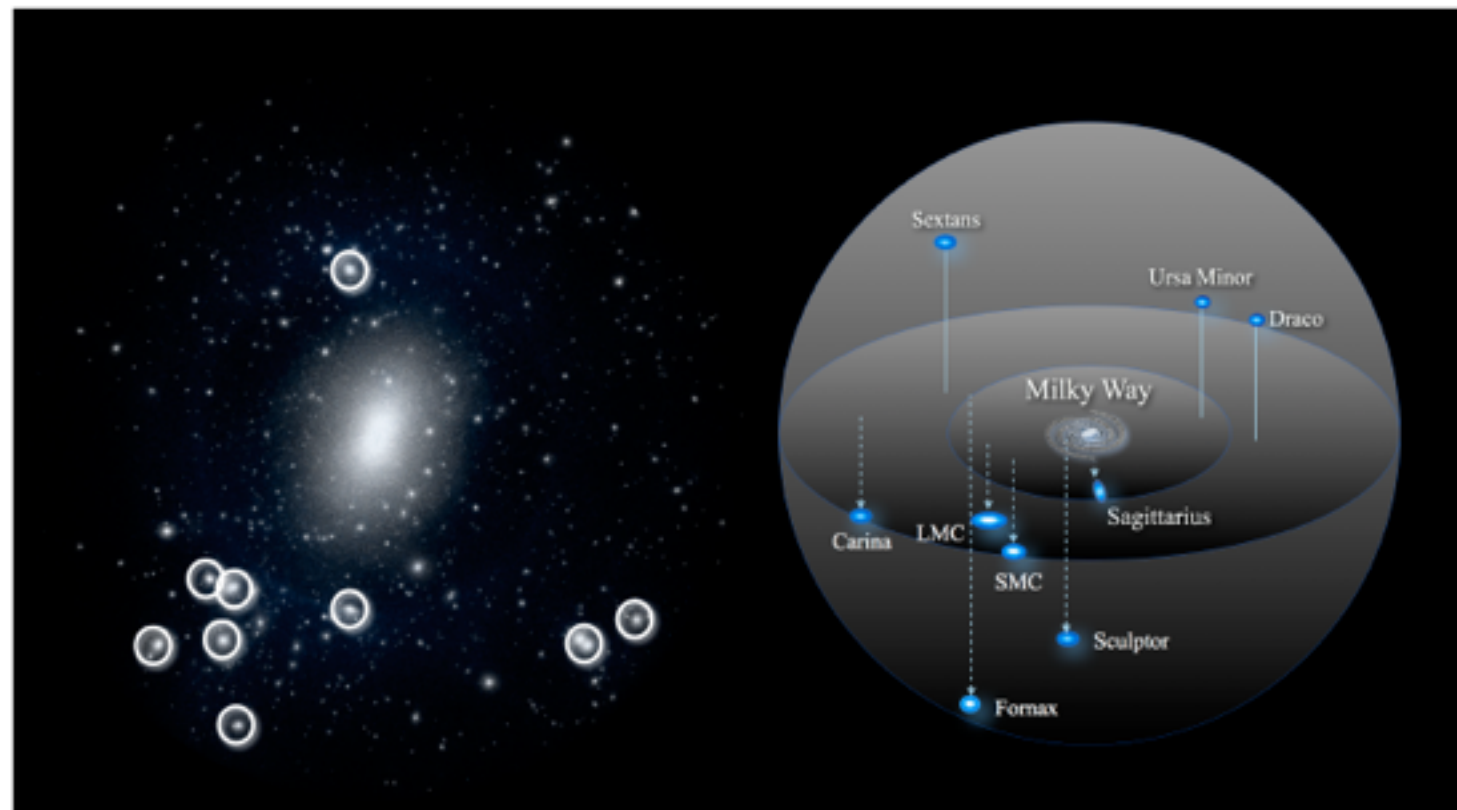
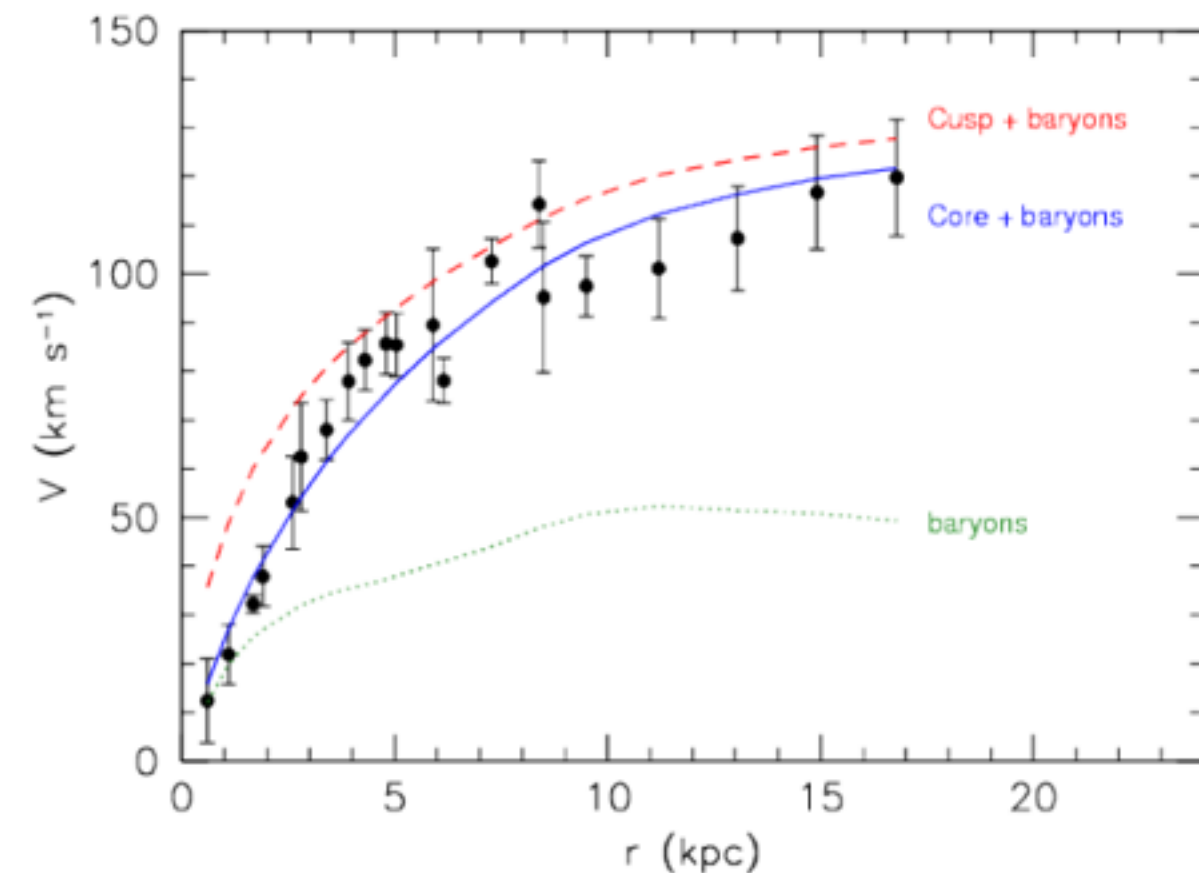
\*Viel et al, 0709.0131



# Crisis of LCDM?

See talk by Chiba

- Core-cusp problem
- Too-big-to-fail problem
- Missing satellite problem



Taken from H. Wlenberg et al, 1309.0913

N-body simulation with collisionless CDM predicts more DM in the central regions of typical galaxies.

Complex baryon physics or DM properties (WDM, self-coupling?)

# Abundance

The baryon-DM coincidence problem:

**Why  $\Omega_{\text{DM}} \sim 5\Omega_B$  ?**

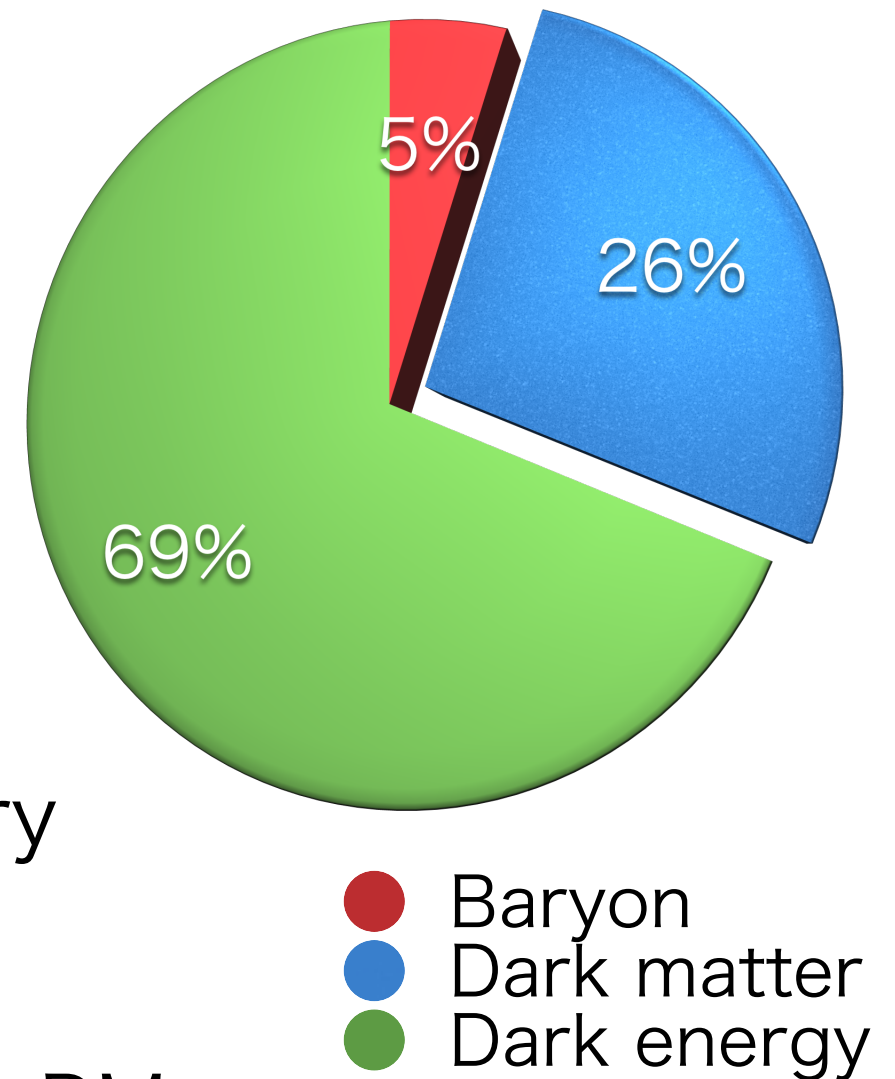
$$\Omega_{\text{DM}} \simeq 0.26 \quad \Omega_B \simeq 0.05$$

It is a puzzle if DM and baryon asymmetry have a different origin.

Many works on this puzzle. e.g. asymmetric DM

(Kaplan, Luty, Zurek 0901.4117 and many others. )

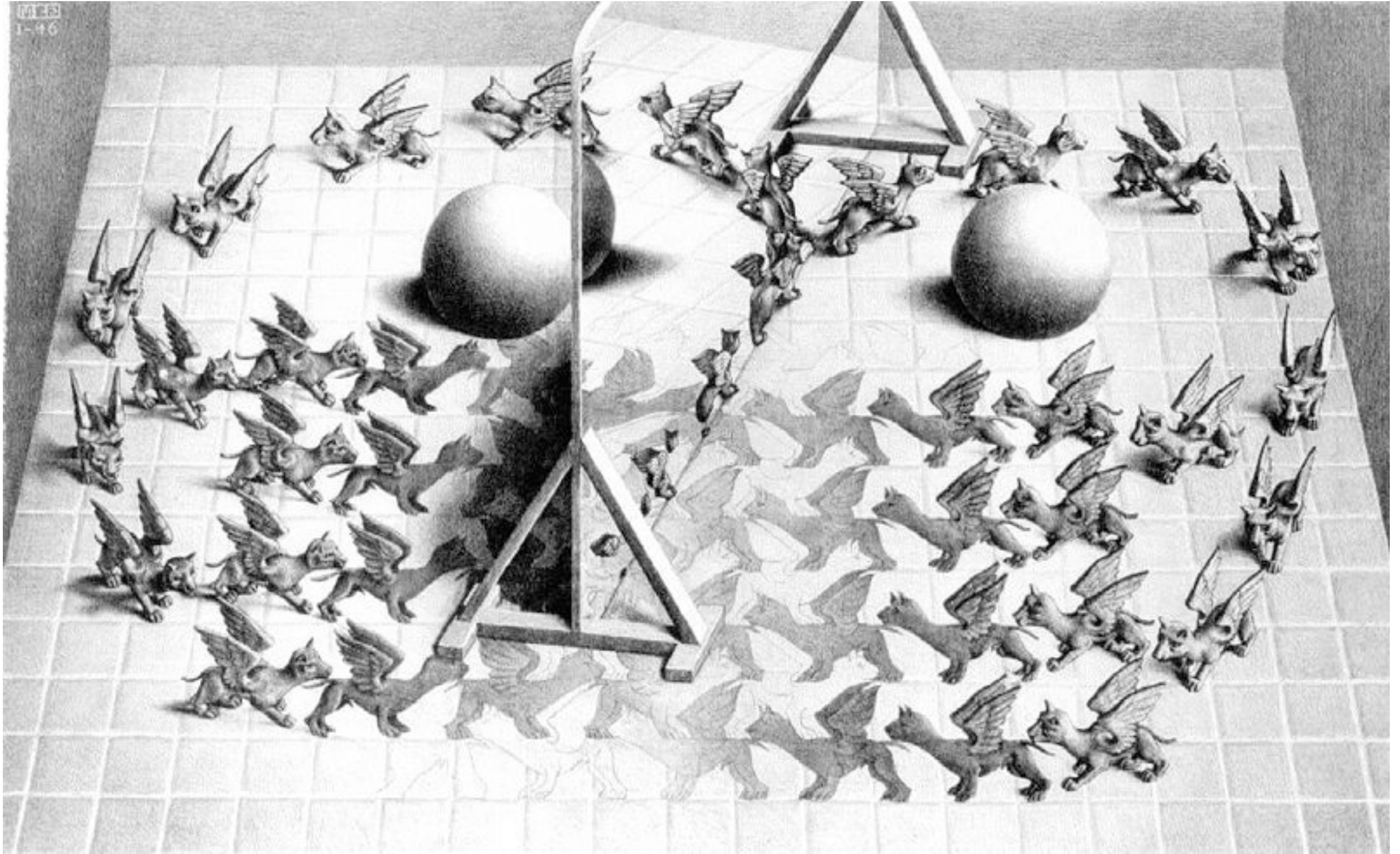
In order to solve the coincidence problem, however, **the DM must know somehow the proton mass.**





# Parallel World

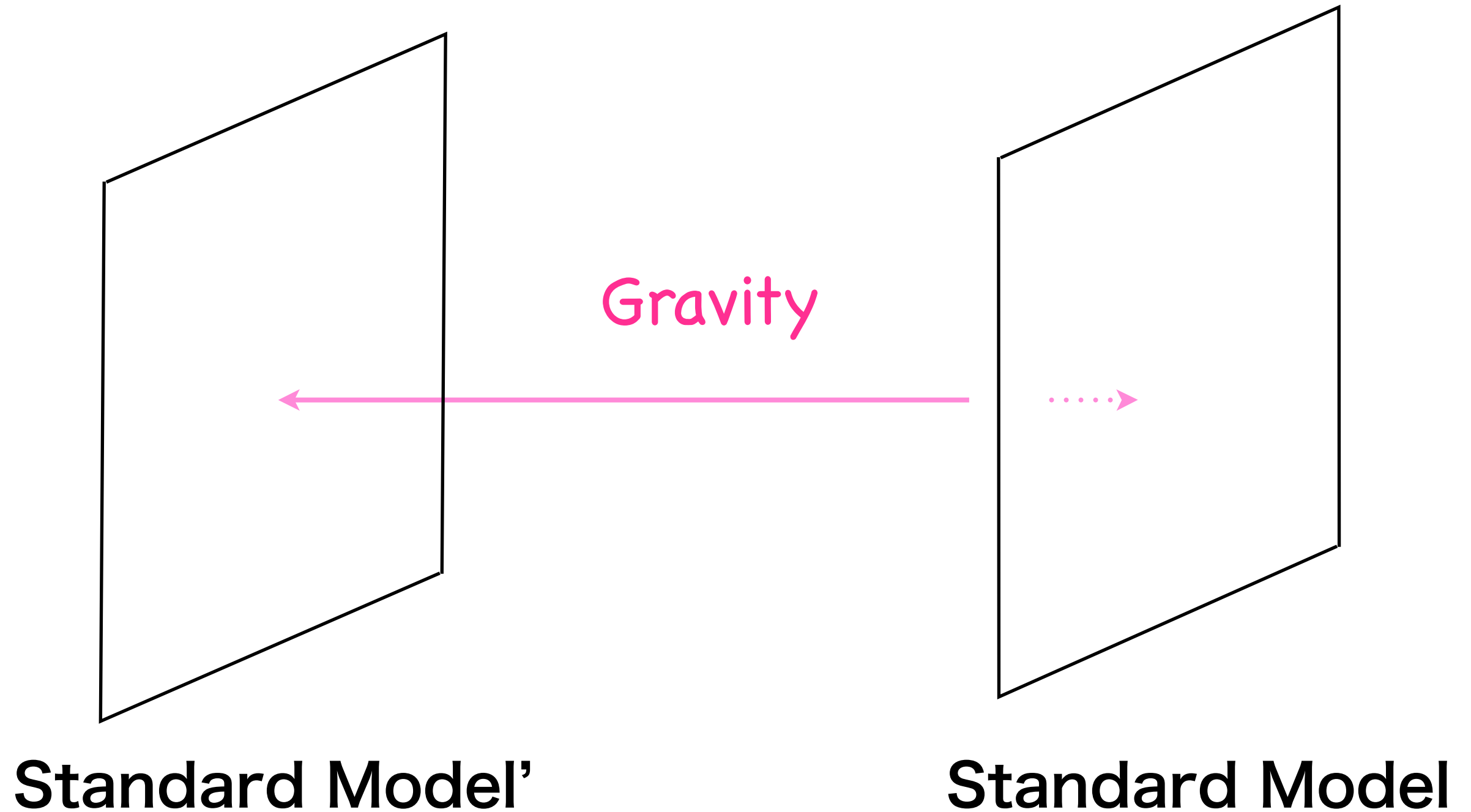
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"Magic mirror" by Escher

# Parallel World

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# Parallel World

- Has the same gauge structure  $SU(3) \times SU(2) \times U(1)$  and the particle content of the SM.

e.g. the counterpart of photon =  $\gamma'$ , para-photon, or simply p-photon.

- The parameters in the Lagrangian can be different by a  $O(1)$  factor.

cf. The so-called “mirror world” assumes an exact  $Z_2$  parity.

(Lee & Yang '56, Blinnikov & Khlopov '82 and many others)

- Interacts with the SM (almost) thru gravity.

\*There are many variations: Higgs portal,  $U(1)$  kinetic mixing, neutrino portal, etc.



# Parallel World

- Dark matter = p-baryon (p-neutron)

Self-interacting DM

- Dark radiation = p-photon and p-neutrino(s)

Dark radiation

- Sterile neutrino = massive p-neutrino(s)

HDM

The amount of DR and sterile neutrinos satisfy

a certain relation:  $N_s^{\frac{4}{3}} = \kappa \Delta N_{\text{eff}}$  [Higaki, Jeong and FT, 1302.2516](#)

$\kappa$  depends on the neutrino mass hierarchy

$N_s$  : effective # of sterile neutrinos

# Self-Interacting Dark Matter

Spiegel and Steinhardt '00.

- DM self-coupling is constrained by the halo shape of galaxy clusters:

$$\frac{\sigma_{\text{DM}}}{m_{\text{DM}}} \lesssim \mathcal{O}(0.1) \text{ cm}^2/\text{g}$$

Rocha et al, 1208.3025  
Peter et al, 1208.3026

*f.*  $\sigma/m = 0.6 \text{ cm}^2/\text{g}$  can solve the central density problem of dwarf spheroidals.

Zavala et al, 1211.6426

For comparison, the neutron-neutron cross section is

$$\frac{\sigma_{nn}}{m_n} \sim 30 \text{ cm}^2/\text{g}$$

The above bound is satisfied if the QCD' scale is several times higher than the QCD scale.

# Dark matter density perturbations

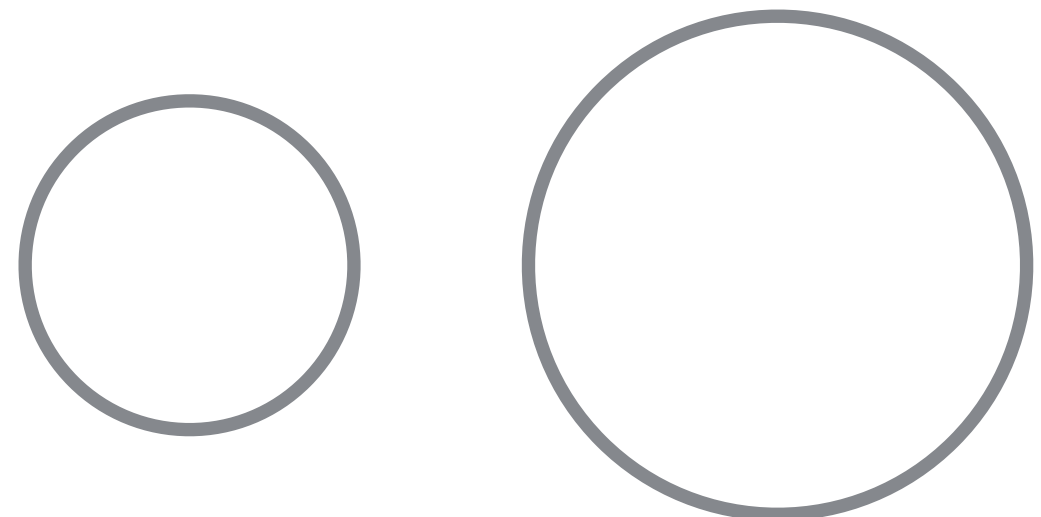
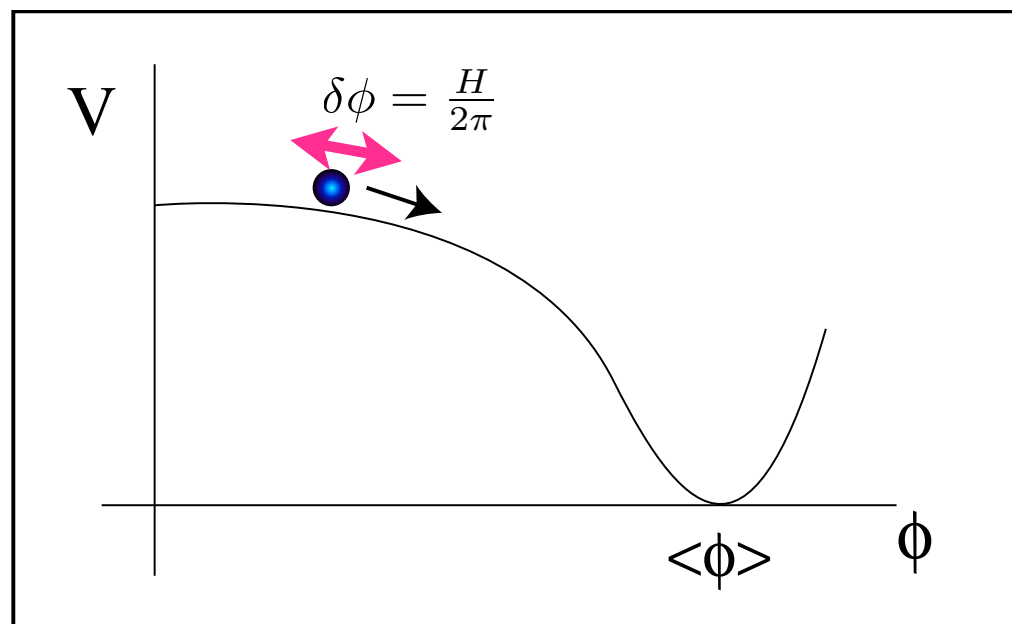
## Adiabatic perturbations :

(usually) produced by quantum fluctuations of the inflaton. Fluctuation of time and volume.

Gravitational potential

Curvature perturbations

$$ds^2 = -(1 + 2\Phi)dt^2 + a(t)^2(1 + 2\Psi)d\mathbf{x}^2$$

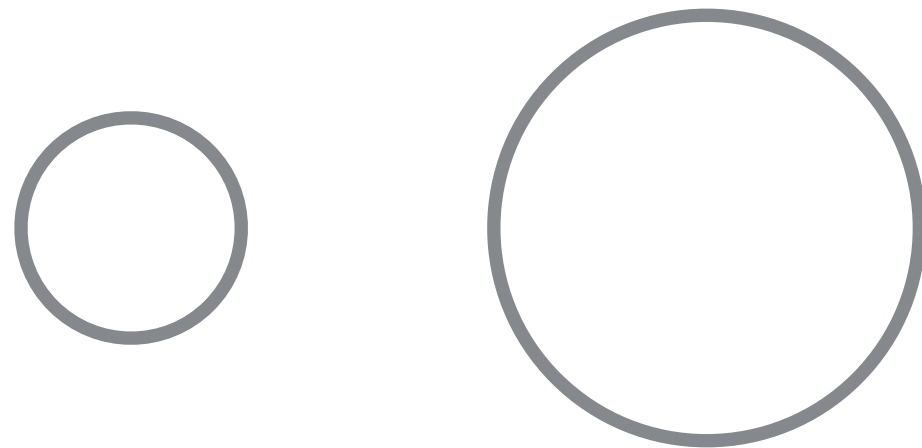


# Dark matter density perturbations

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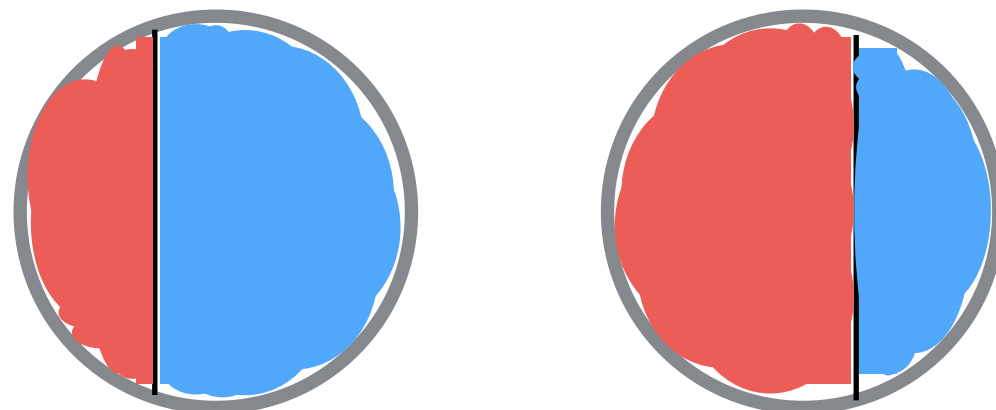
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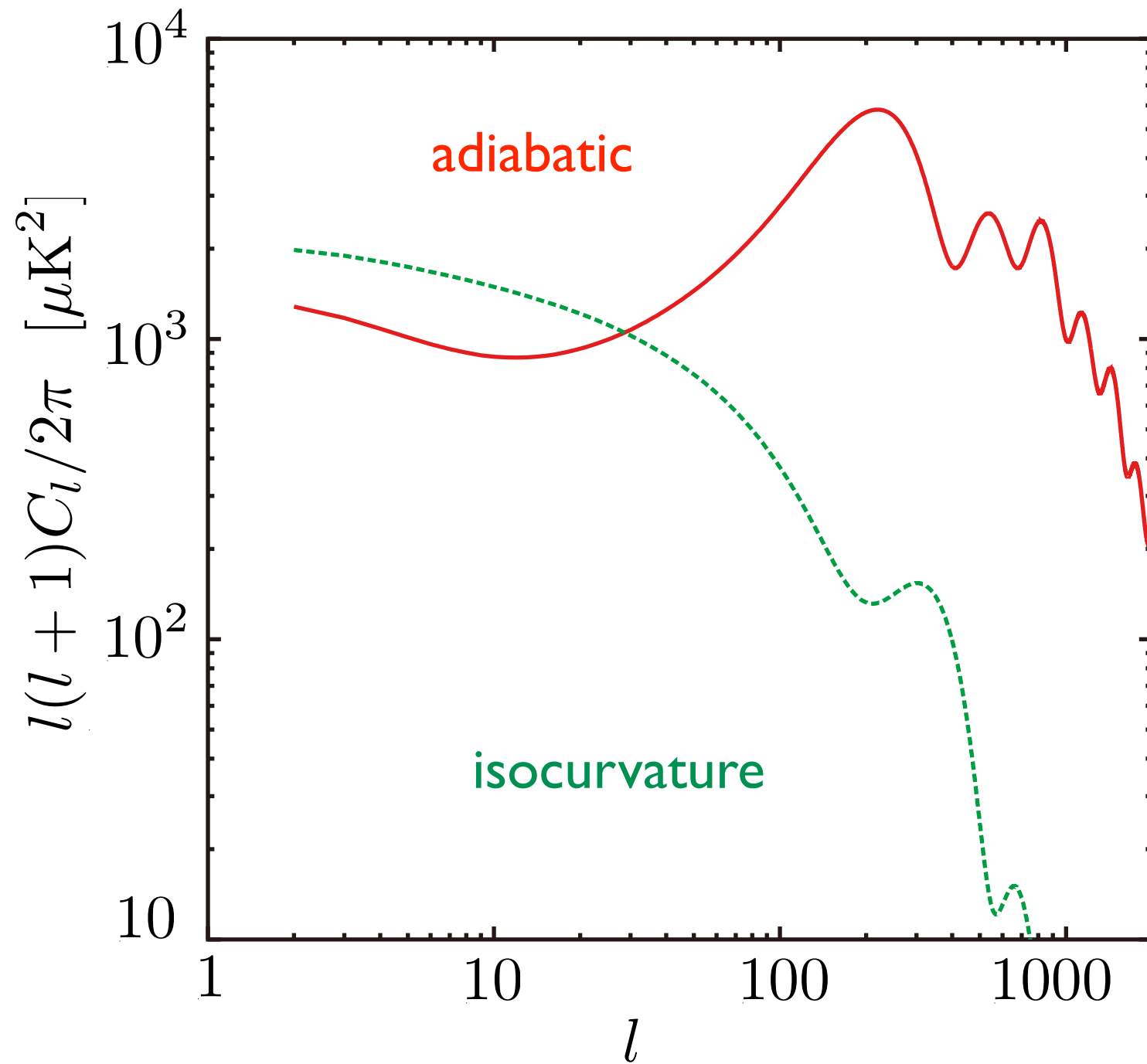


## Isocurvature perturbations:

Fluctuations between components



# Adiabatic/isocurvature perturbations



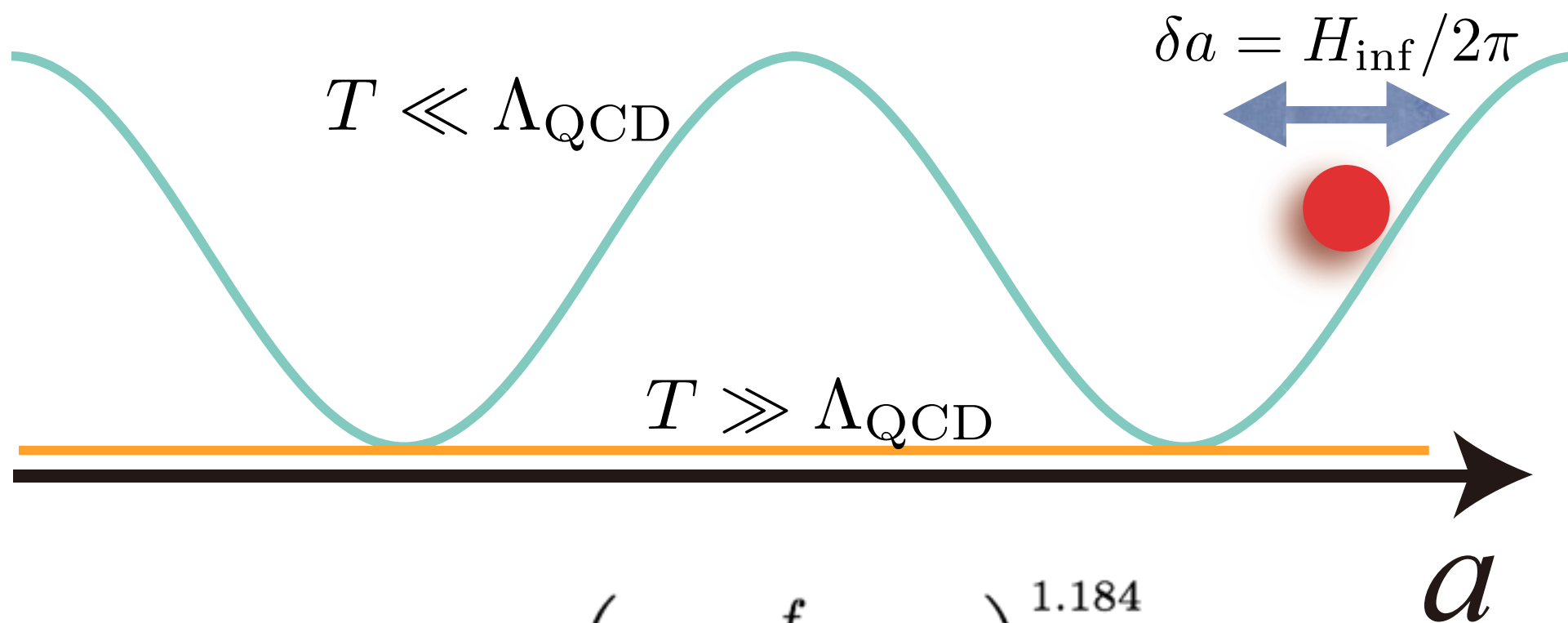
$$\beta_{\text{iso}}(k) = \frac{\mathcal{P}_{II}(k)}{\mathcal{P}_{\mathcal{R}\mathcal{R}}(k) + \mathcal{P}_{II}(k)} < 0.038 \quad (95\% \text{ CL}) \quad (\text{uncorrelated})$$

@  $k = 0.05 \text{ Mpc}^{-1}$



# QCD Axion

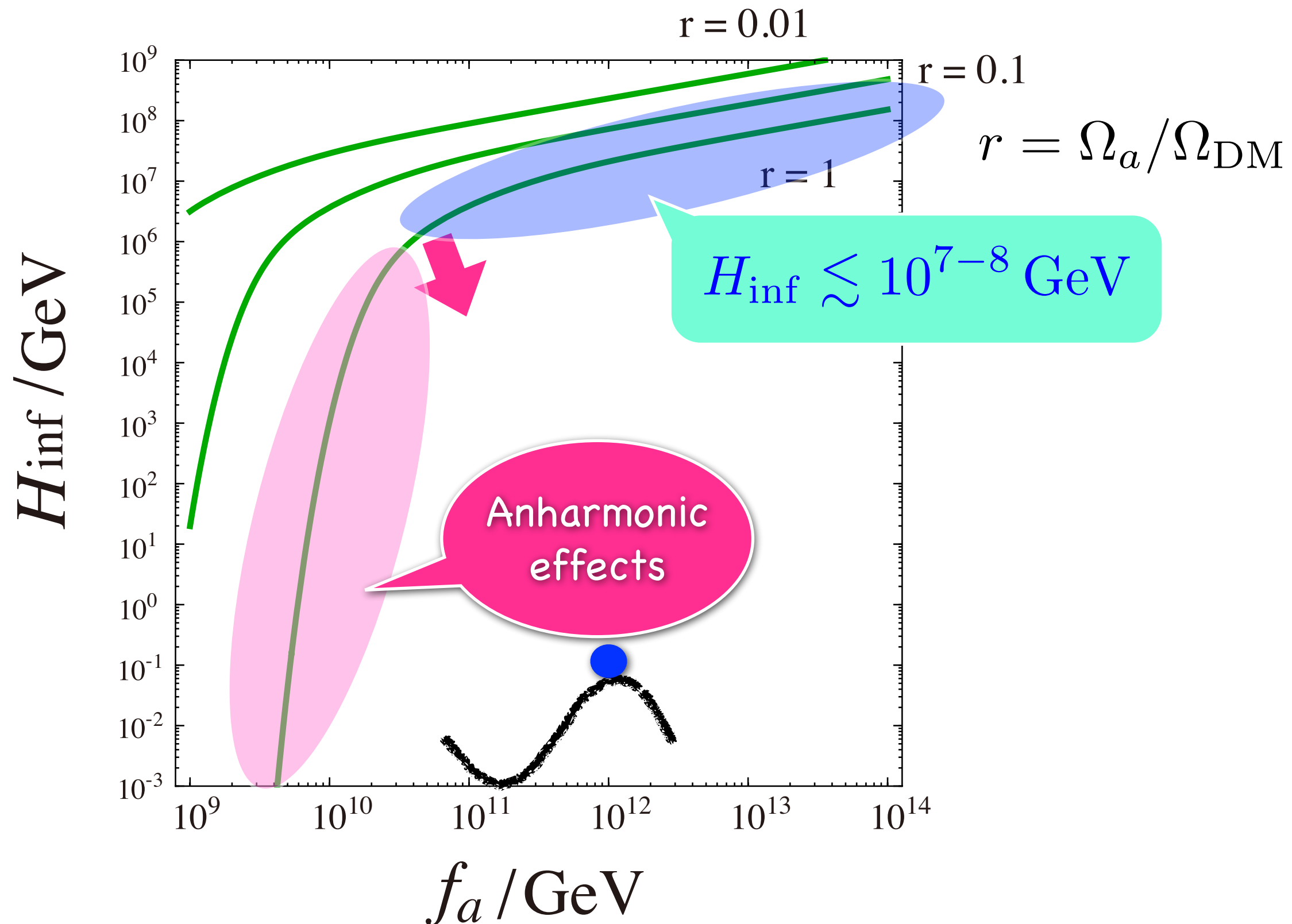
$$\mathcal{L} = \left( \theta + \frac{a}{f_a} \right) \frac{g_s^2}{32\pi^2} G^{\mu\nu} \tilde{G}_{\mu\nu}$$



$$\Omega_a h^2 \simeq 0.11 \left( \frac{f_a}{5 \times 10^{11} \text{ GeV}} \right)^{1.184} F(\theta_*) \theta_*^2$$

$$S \simeq \left( \frac{\Omega_a}{\Omega_{\text{DM}}} \right) \left[ 2 \frac{\delta a}{a_*} + \left( \frac{\delta a}{a_*} \right)^2 \right]$$

# Isocurvature constraint on $H_{\text{inf}}$



# Isocurvature constraint on $H_{\text{inf}}$

The tension with high-scale inflation can be avoided if

1. PQ symmetry restored during/after inflation

Linde and Lyth '90

2. Larger effective decay constant during inflation.

Linde '91

3. Temporarily enhanced PQ symmetry breaking.

Dine, Anisimov '05,  
Jeong, FT '13, FT, Yamada '15.

# Summary

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The decelerated expansion itself may be reasonable as the gravitational force is attractive, but there are many puzzles about DM:

- ✓ **What is DM?**

- Why long-lived? How produced? Abundance? Cold or warm?  
Non-gravitational couplings?

- ✓ **DM density perturbations**

- Isocurvature mode/non-Gaussianity?

- ✓ **Non-linear gravitational evolution** See talk by Taruya

- ✓ **DM distribution** in dSphs, galaxies, and clusters.

- Any tension with LCDM and N-body simulations? See talk by Chiba

- ✓ **Unified picture of DM and DE?**

- Ultralight axion DE and saxion DM, Mass-varying neutrinos, etc.