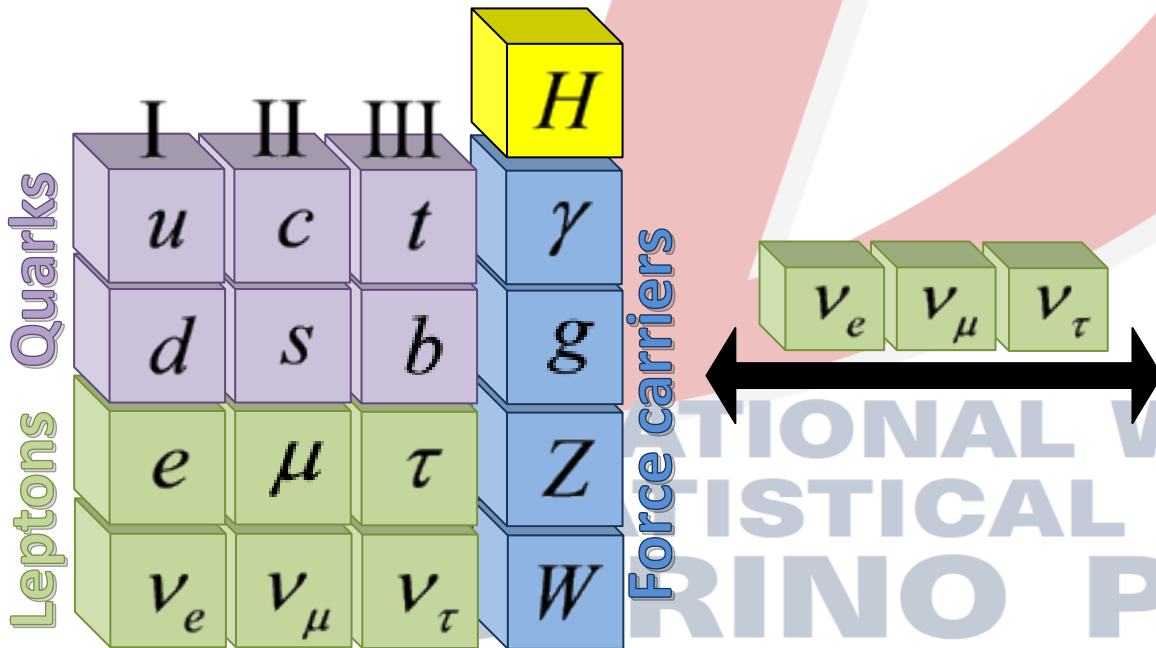


Cosmological Constraints on Neutrino Properties

Francisco Villaescusa-Navarro
OATS/INAF/INFN, Trieste, Italy



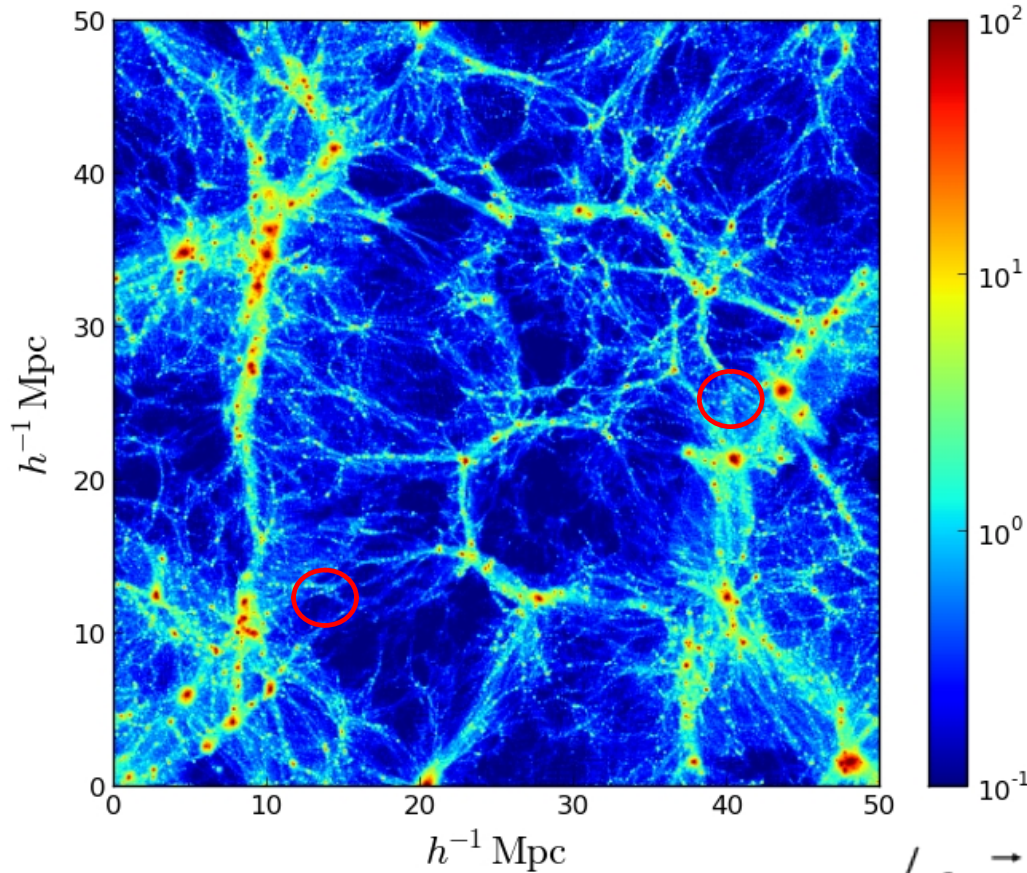
THE KAVLI INSTITUTE FOR THE PHYSICS AND
MATHEMATICS OF THE UNIVERSE, KASHIWA, JAPAN

30 MAY-1 JUNE 2016

Outline

- Cosmic neutrino background
- Effects on cosmological observables
- Statistical methods used to
constraint neutrino masses: challenges

Statistical estimators in cosmology



$$\delta(\vec{x}) = \frac{\rho(\vec{x}) - \langle \rho \rangle}{\langle \rho \rangle}$$

$$\xi(r) = \left\langle \delta(\vec{x}) \delta(\vec{x} + \vec{r}) \right\rangle$$

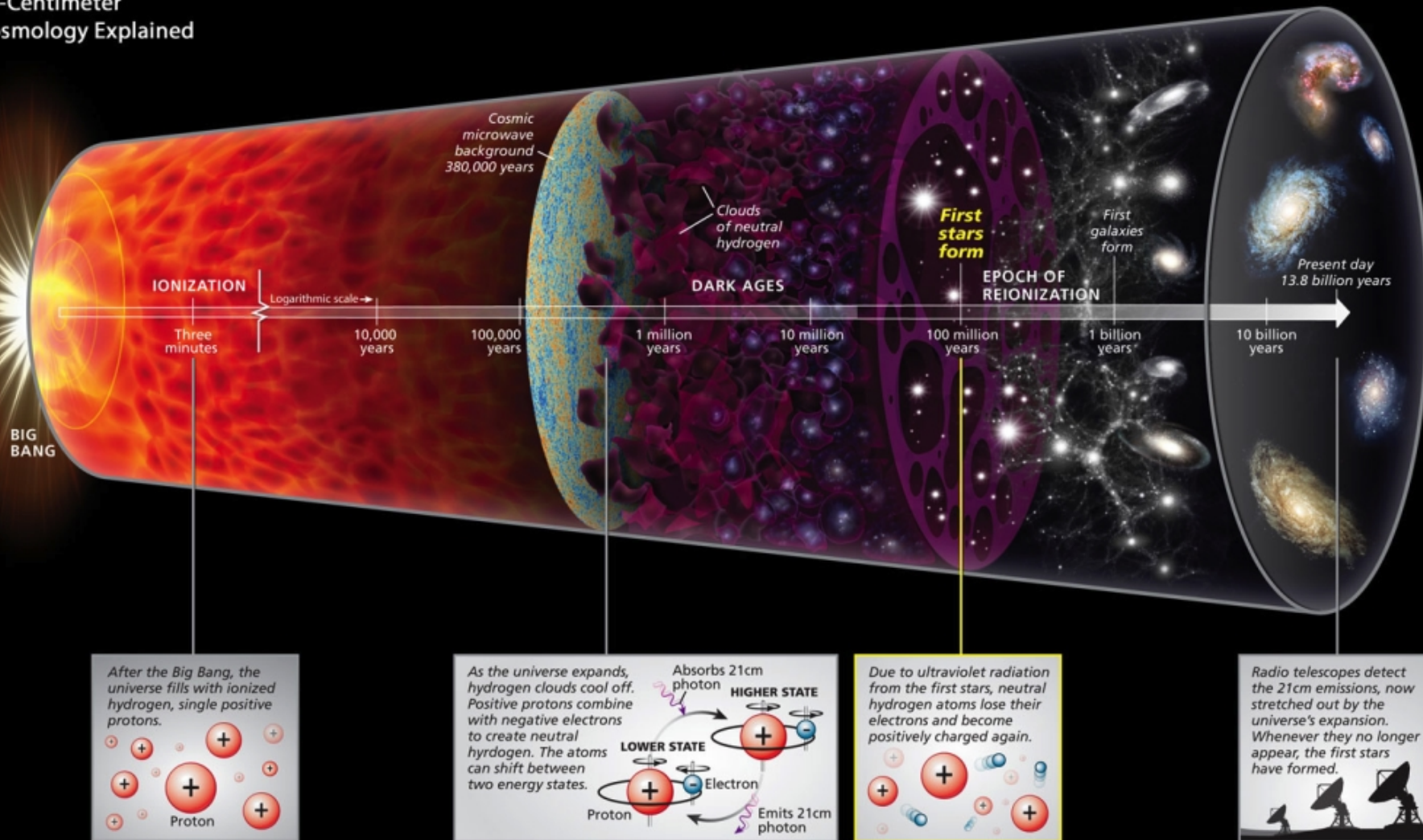
$$P(k) = \frac{1}{(2\pi)^3} \int d^3\vec{r} e^{-i\vec{k}\vec{r}} \xi(r)$$

$$\left\langle \delta(\vec{k}) \delta(\vec{k}') \right\rangle = (2\pi)^3 \delta_D(\vec{k} - \vec{k}') P(k)$$

Gaussian density field completely specified by its power spectrum/correlation function!!

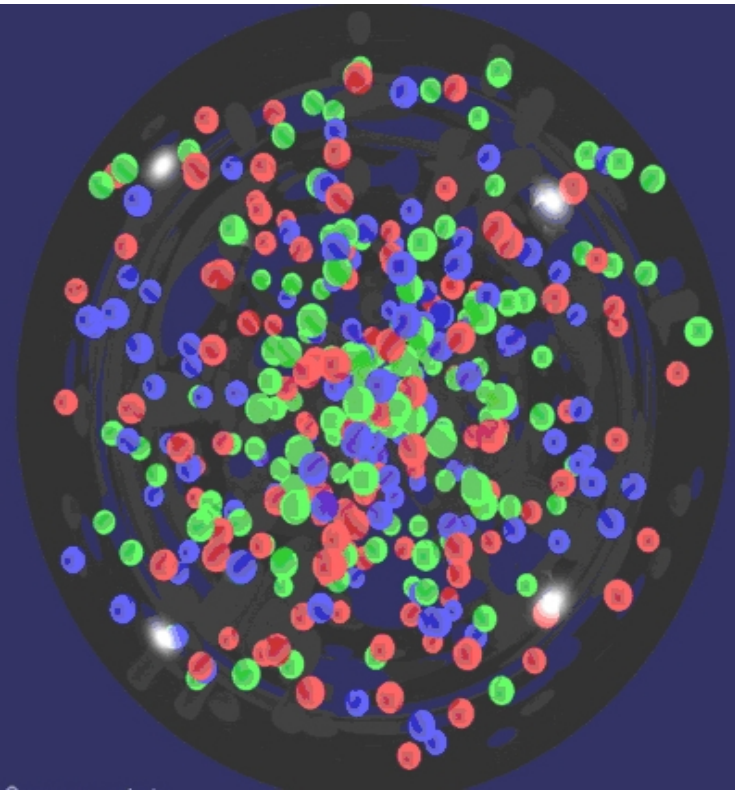
Universe history

21-Centimeter
Cosmology Explained



Cosmic neutrino background

$$T \approx 10^{10} \text{ K} \approx 1 \text{ MeV}$$



- Photons
- Electrons & Positrons
- Neutrinos & Antineutrinos

$$\nu + e^- \rightleftharpoons \nu + e^- \quad e^- + e^+ \rightleftharpoons \nu + \bar{\nu}$$

$$n_\nu(p, z) dp \cong \frac{4\pi g_\nu}{(2\pi\hbar c)^3} \left(\frac{p^2 dp}{e^{(p/k_B T_\nu(z))} + 1} \right)$$

$$T_{\nu,0} = \left(\frac{4}{11} \right)^{1/3} T_{\gamma,0} \approx 1.95 \text{ K}$$

$$T_\nu(z) = T_{\nu,0}(1+z)$$

Cosmic neutrino background

Properties

Redshift; measurement of time in cosmology

$z=0$ (today)

$z=\infty$ (Big Bang)

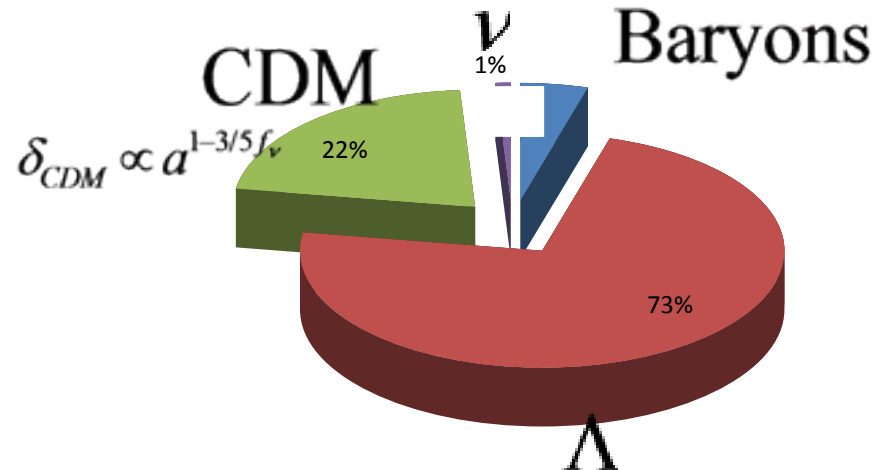
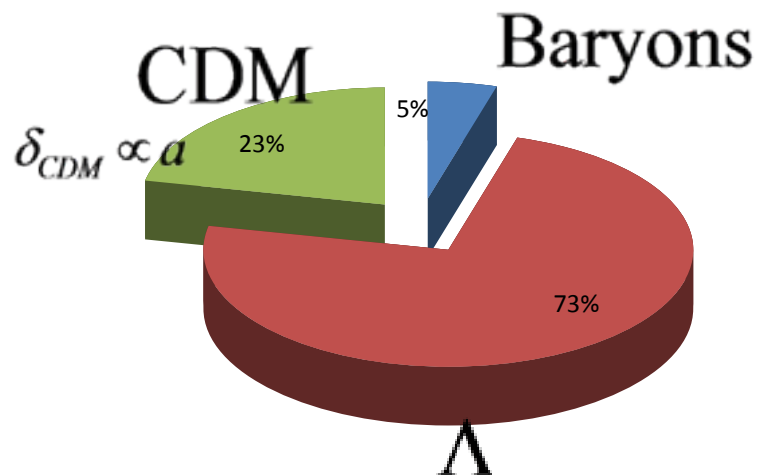
$$\bar{n}_\nu(z) \cong 113 (1+z)^3 \frac{(\nu + \bar{\nu})}{\text{cm}^3}$$

$$\bar{V}_\nu(z) \cong 160(1+z) \left(\frac{\text{eV}}{m_\nu} \right) \text{km/s}$$

- Relativistic at high-redshift (early Universe)
- Non-relativistic at low-redshift (late Universe)

Effects at linear order

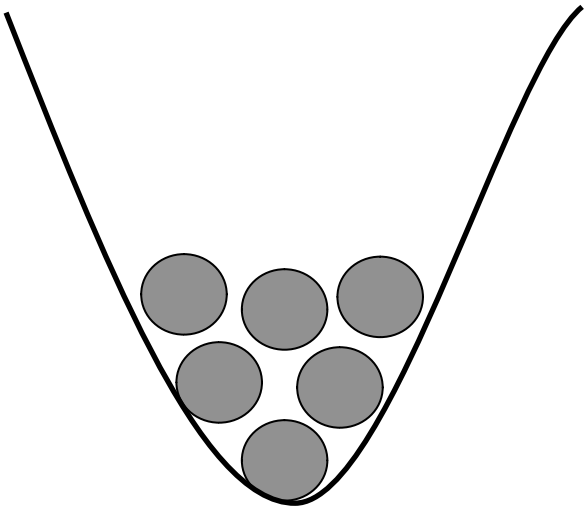
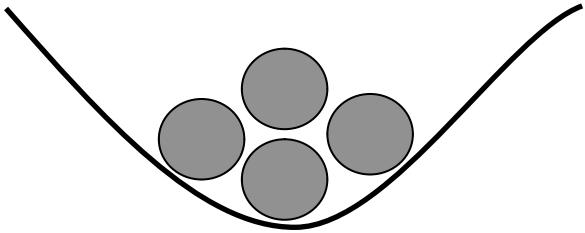
1. Modification of the Matter-Radiation equality time
1. Slow down the growth of matter perturbations



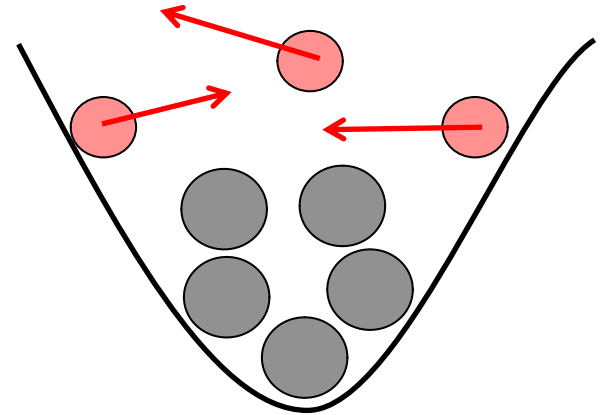
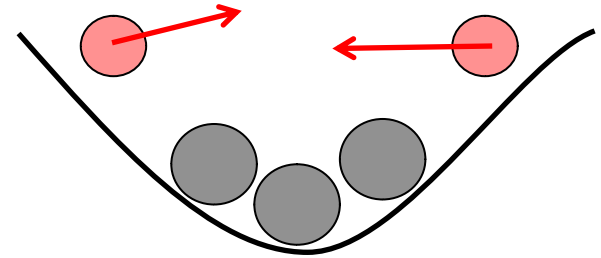
$$\Omega_{DM} = \Omega_{CDM} + \Omega_\nu = \text{fixed}$$

Effects at linear order

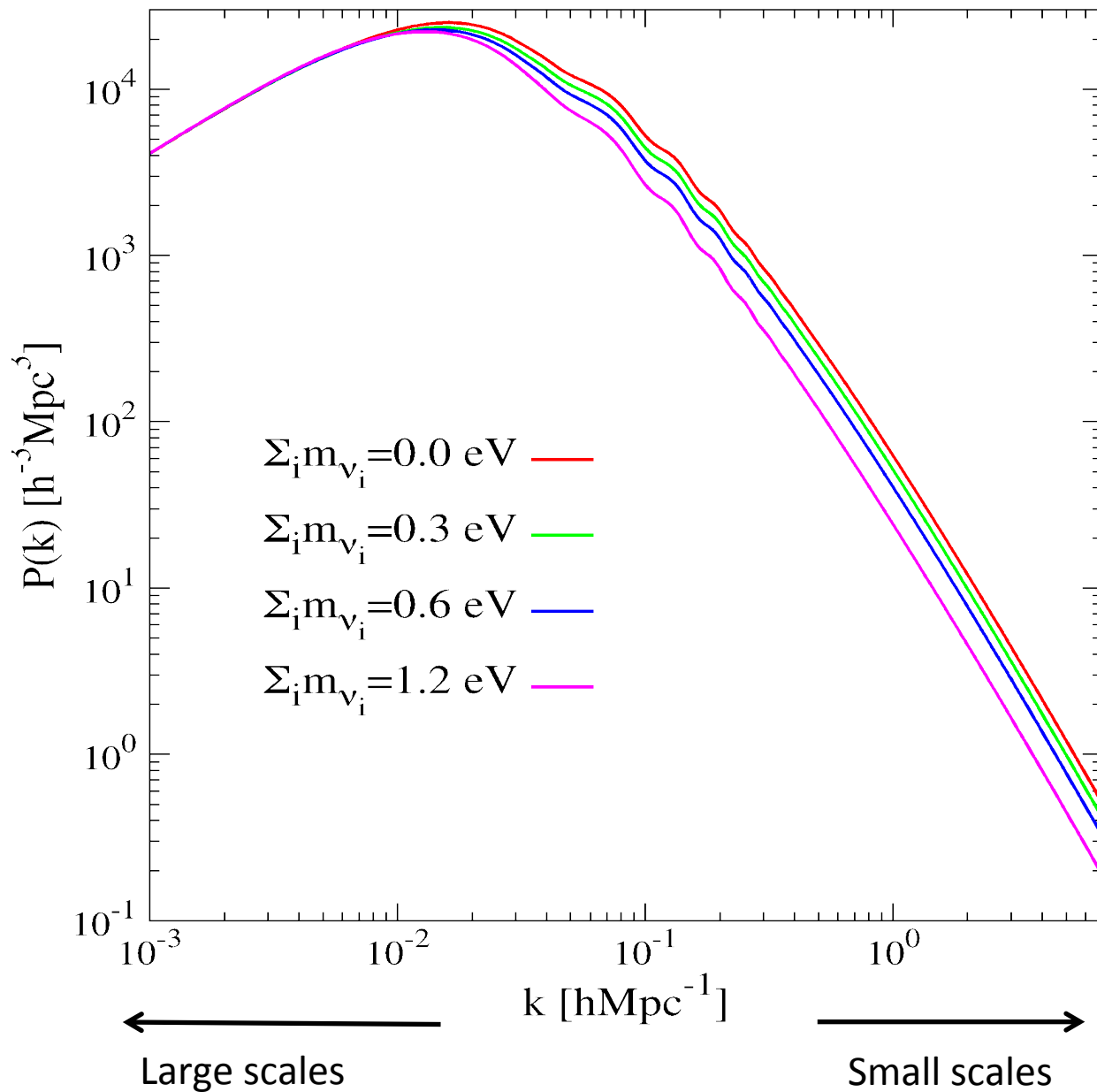
Massless neutrinos



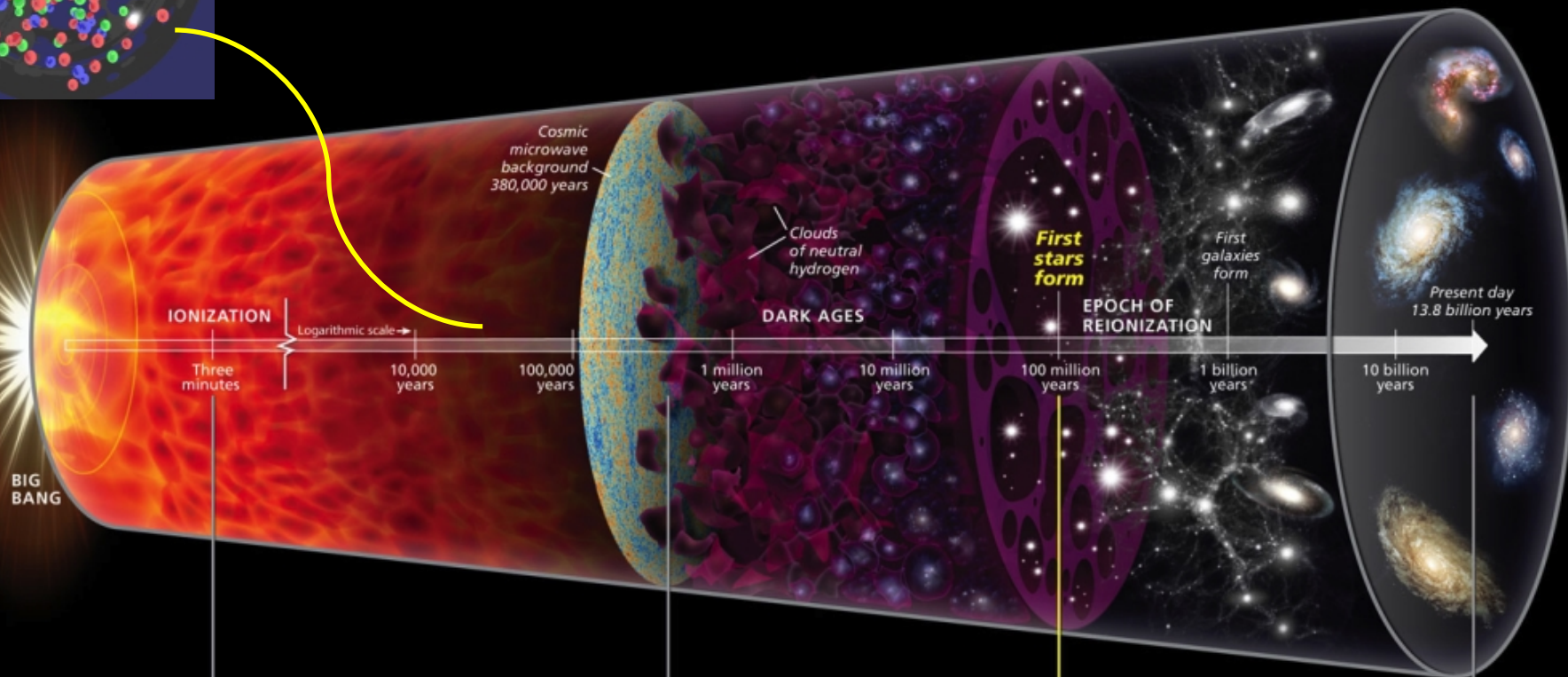
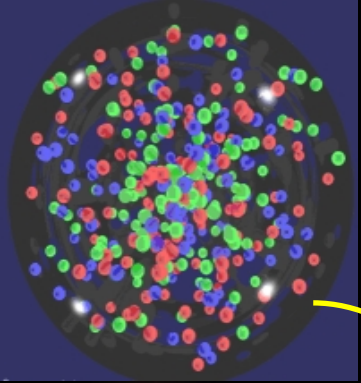
Massive neutrinos



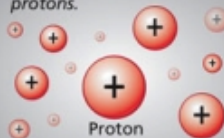
Effects at linear order



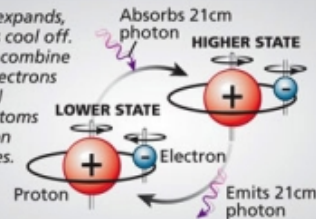
Cosmological observables



After the Big Bang, the universe fills with ionized hydrogen, single positive protons.



As the universe expands, hydrogen clouds cool off. Positive protons combine with negative electrons to create neutral hydrogen. The atoms can shift between two energy states.



Due to ultraviolet radiation from the first stars, neutral hydrogen atoms lose their electrons and become positively charged again.

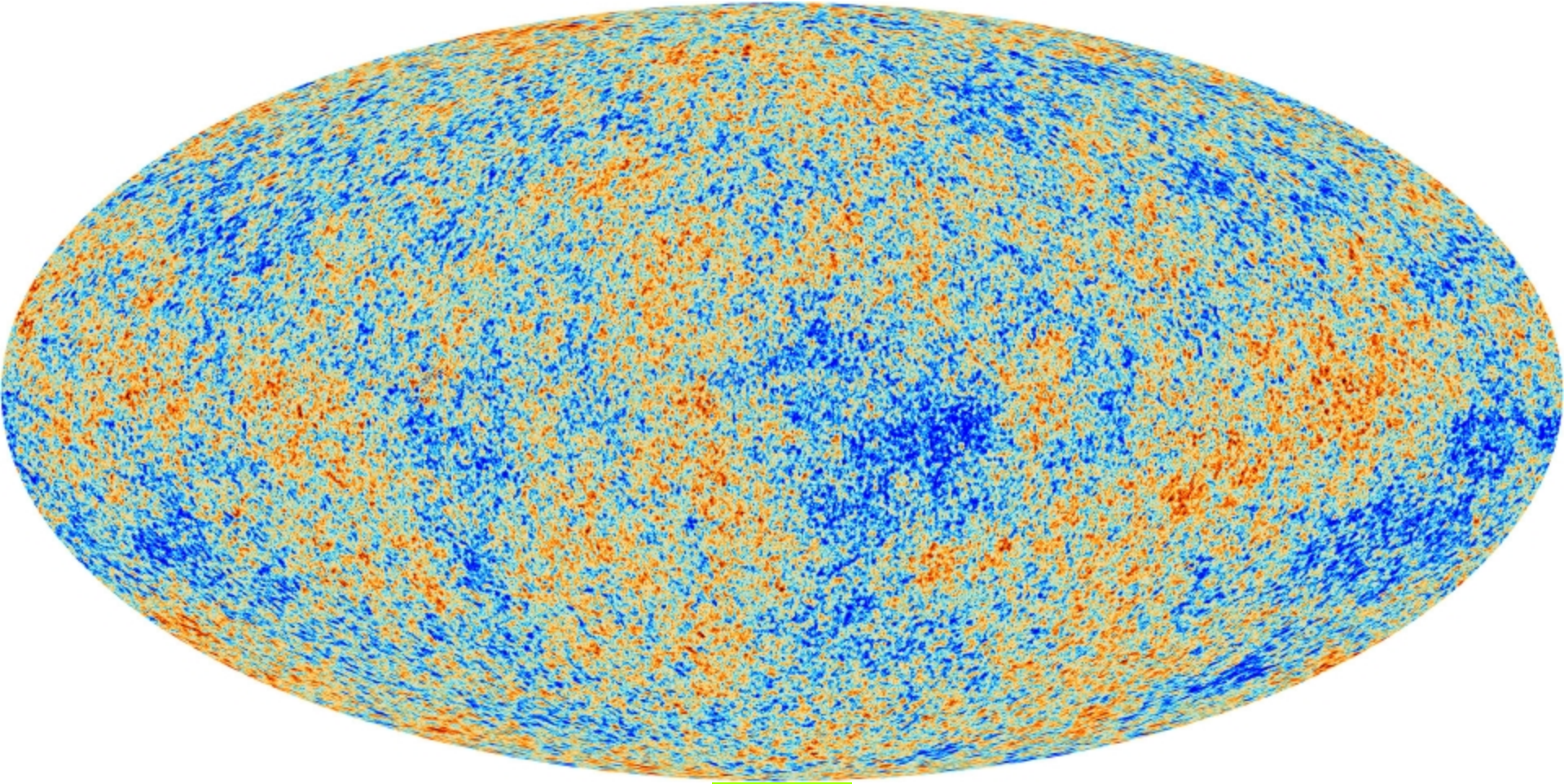


Radio telescopes detect the 21cm emissions, now stretched out by the universe's expansion. Whenever they no longer appear, the first stars have formed.



CMB anisotropies

Cosmological constraints are intrinsically statistical!!!

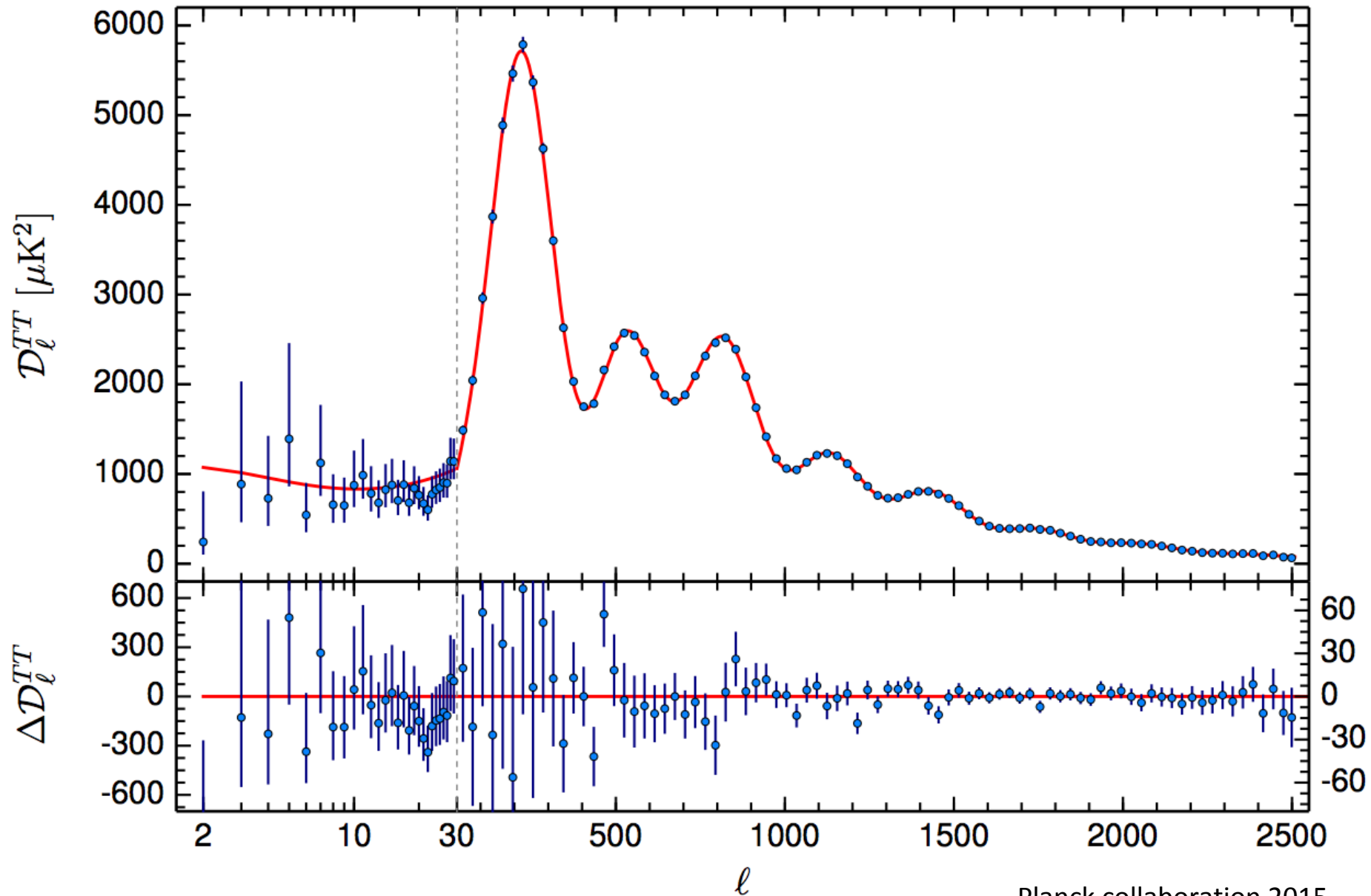


$$T_\gamma = 2.7255 \text{ K}$$

$$\frac{\Delta T_\gamma}{T_\gamma} \sim 10^{-5}$$

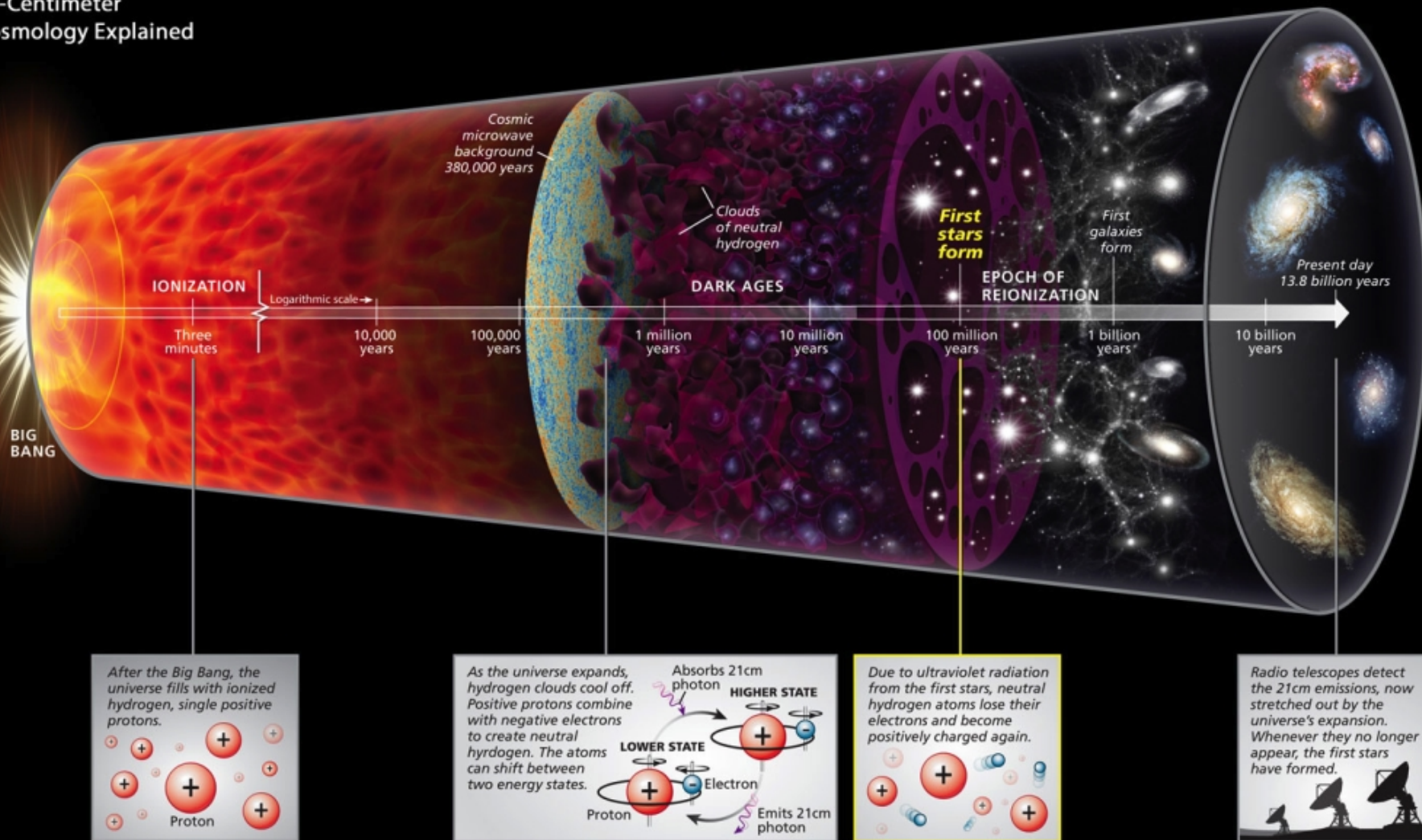
Planck collaboration 2015

CMB anisotropies power spectrum

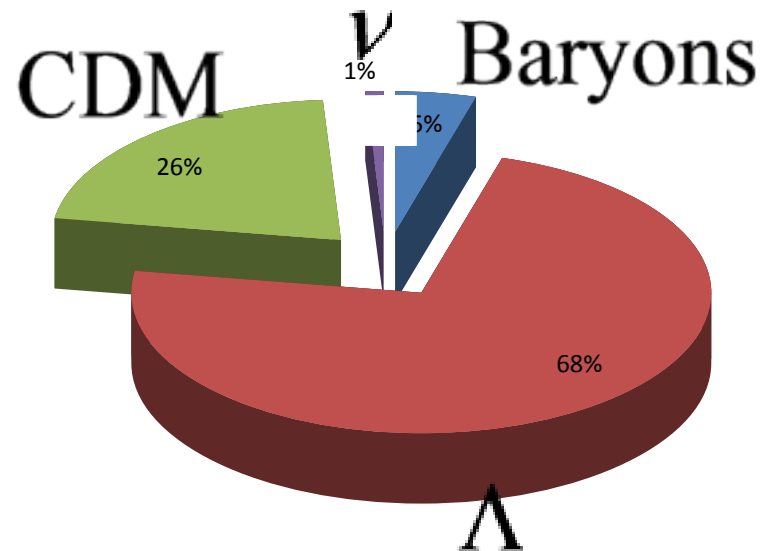
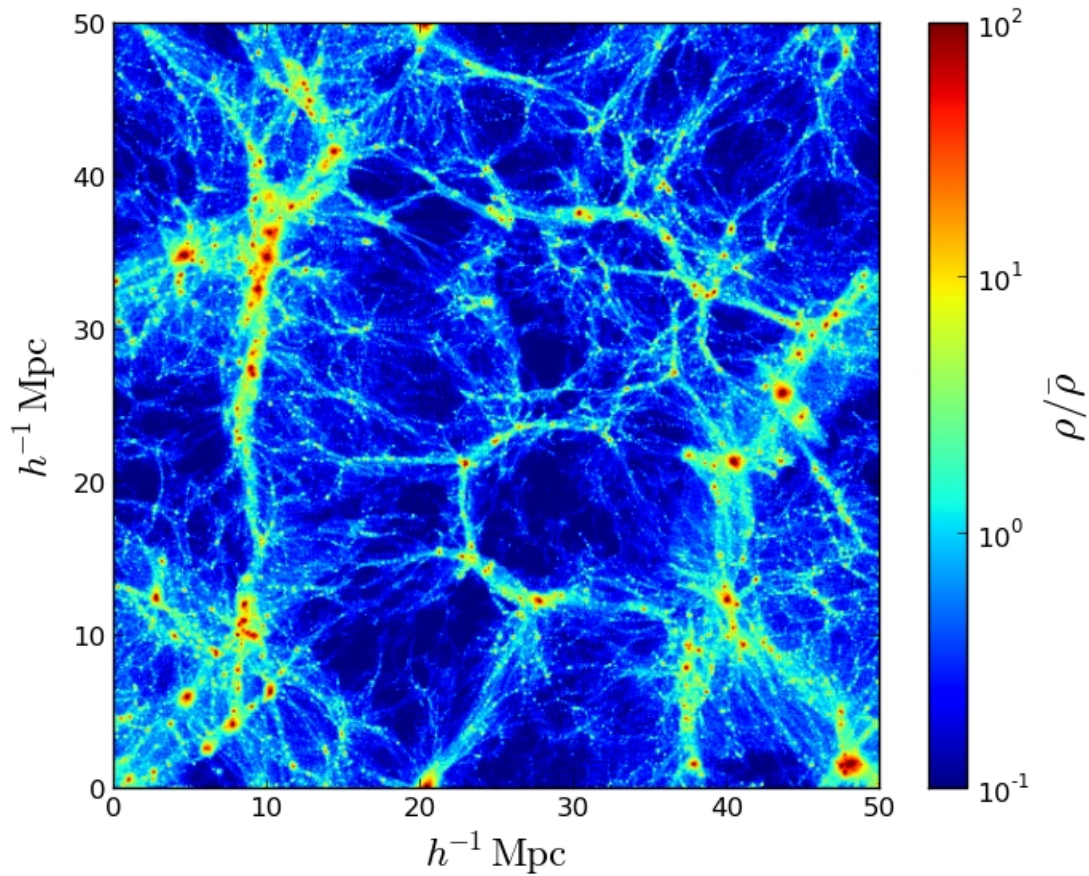


Cosmological observables

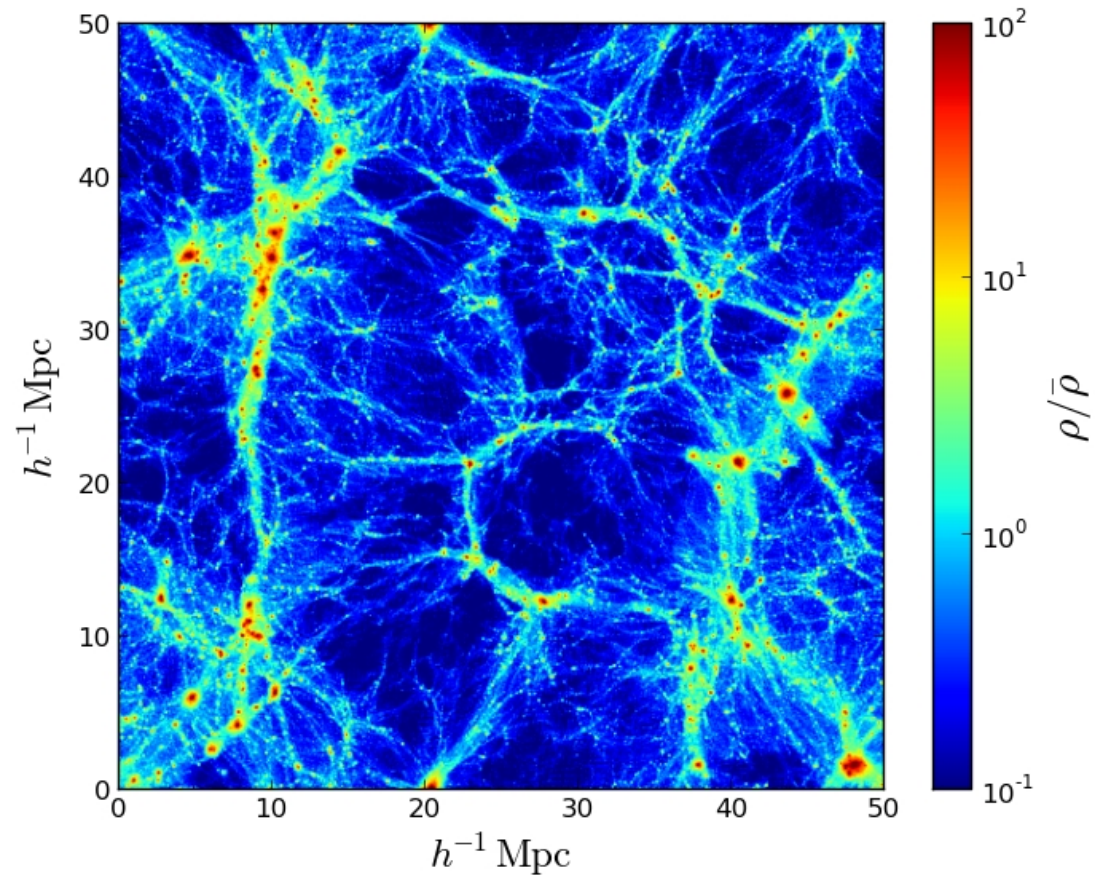
21-Centimeter
Cosmology Explained



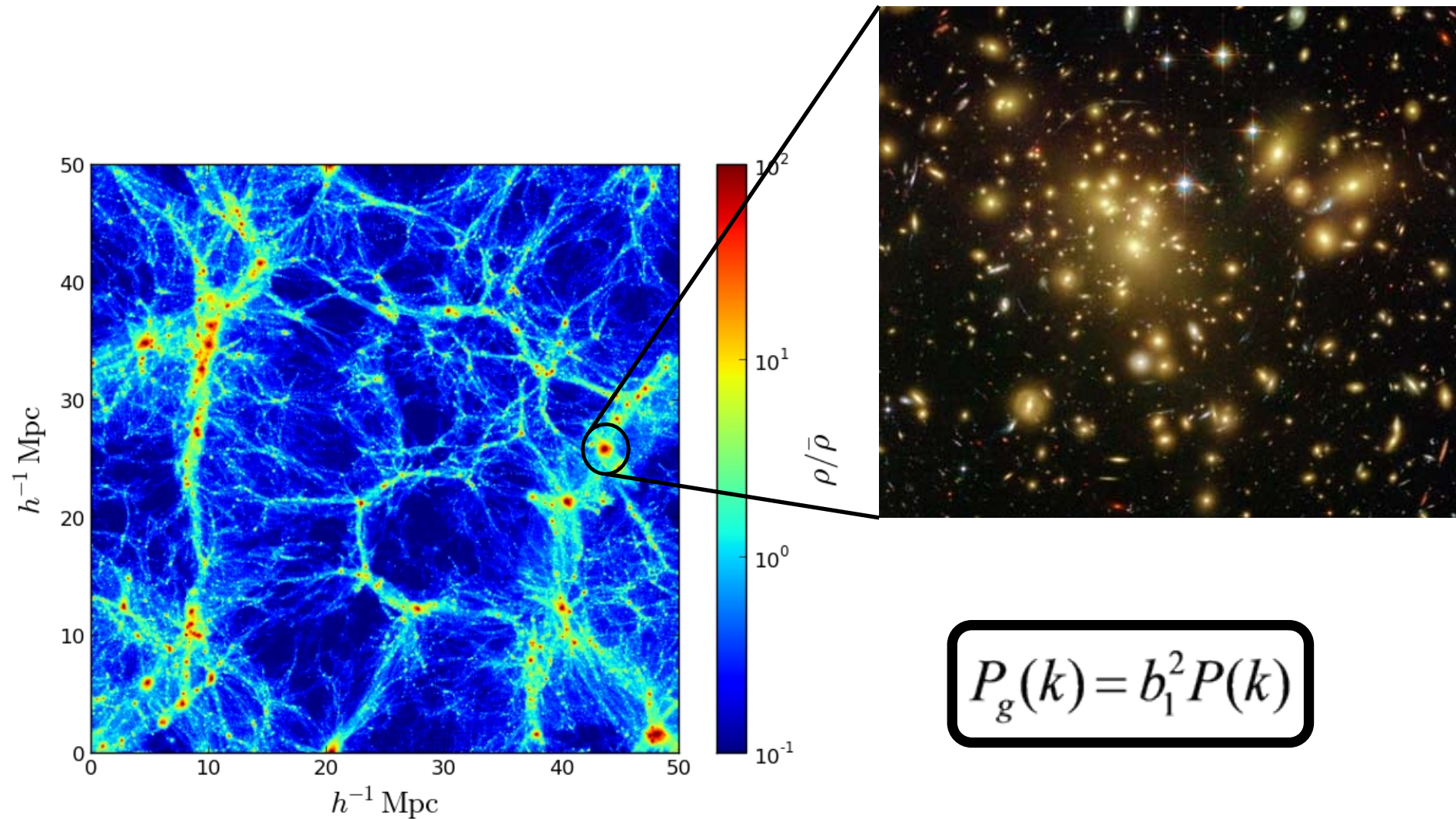
Density field tracers



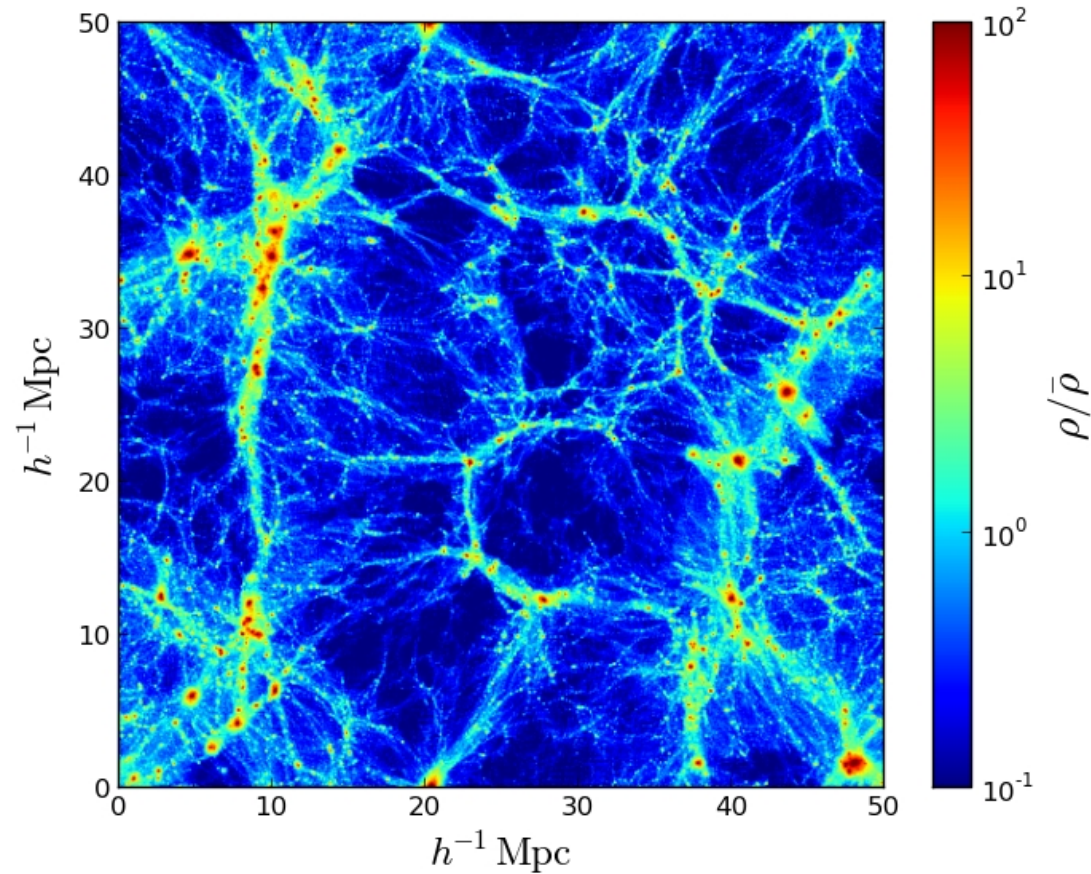
Density field tracers



Density field tracers



Density field tracers



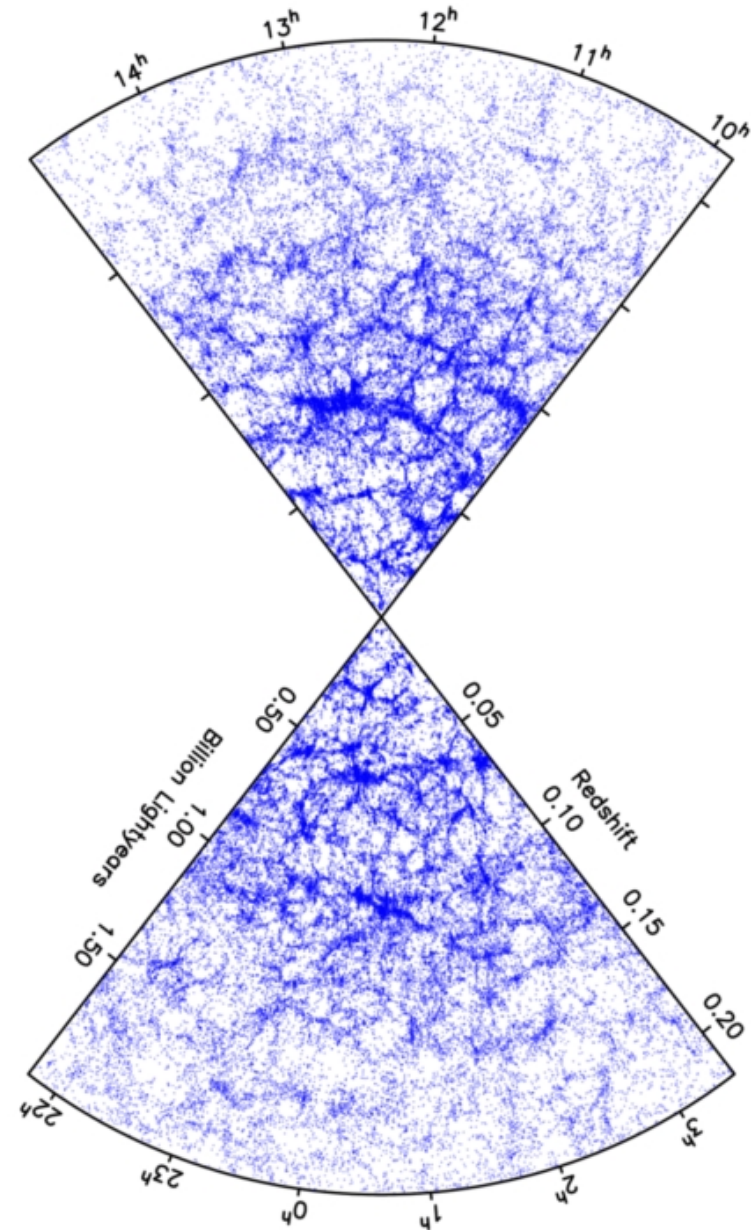
2dF survey

15000 degrees²

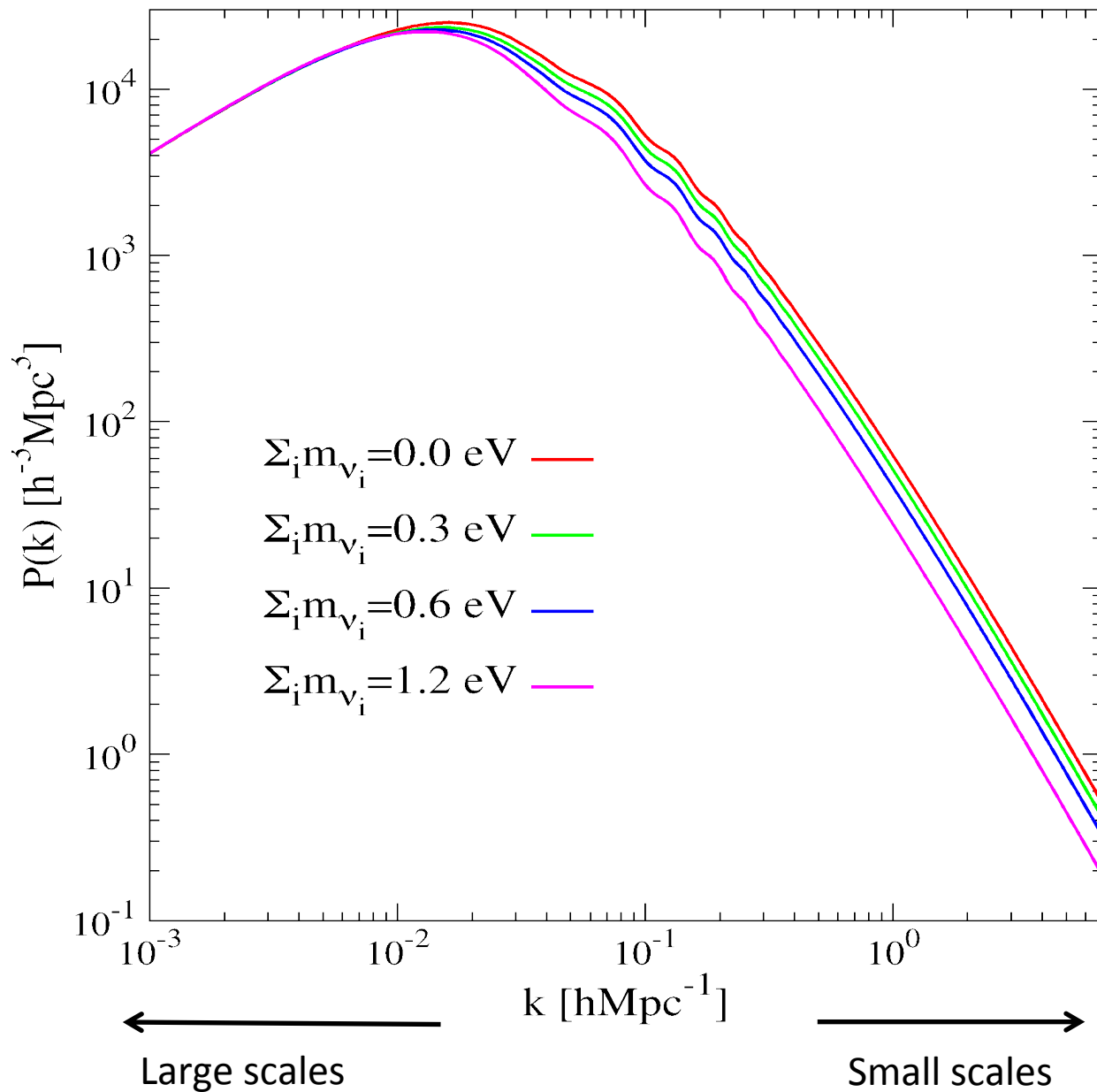
$V \sim 0.1 \text{ (Gpc/h)}^3$

~ 250000 spectroscopic redshifts

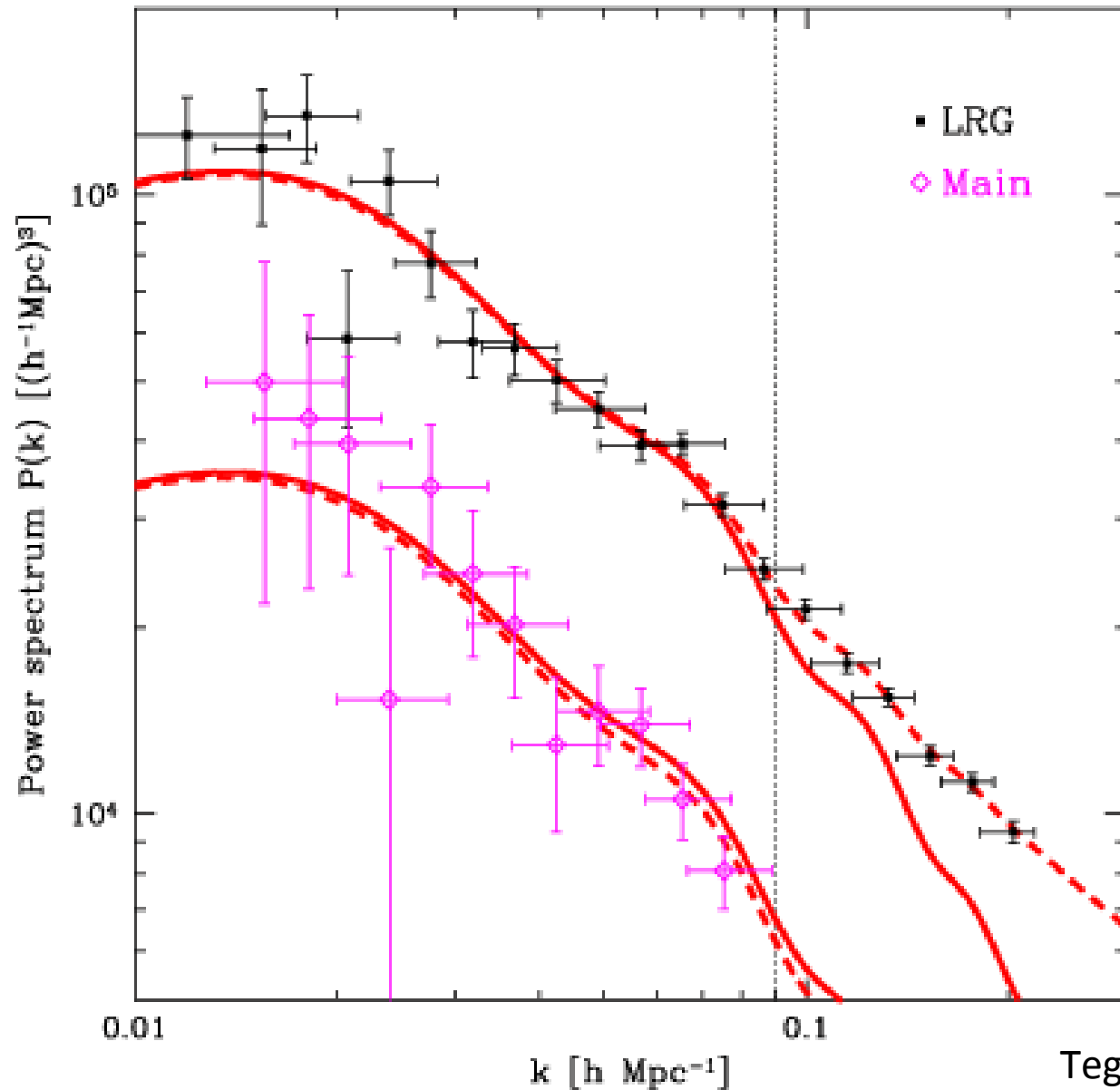
$$P_g(k) = b_1^2 P(k)$$



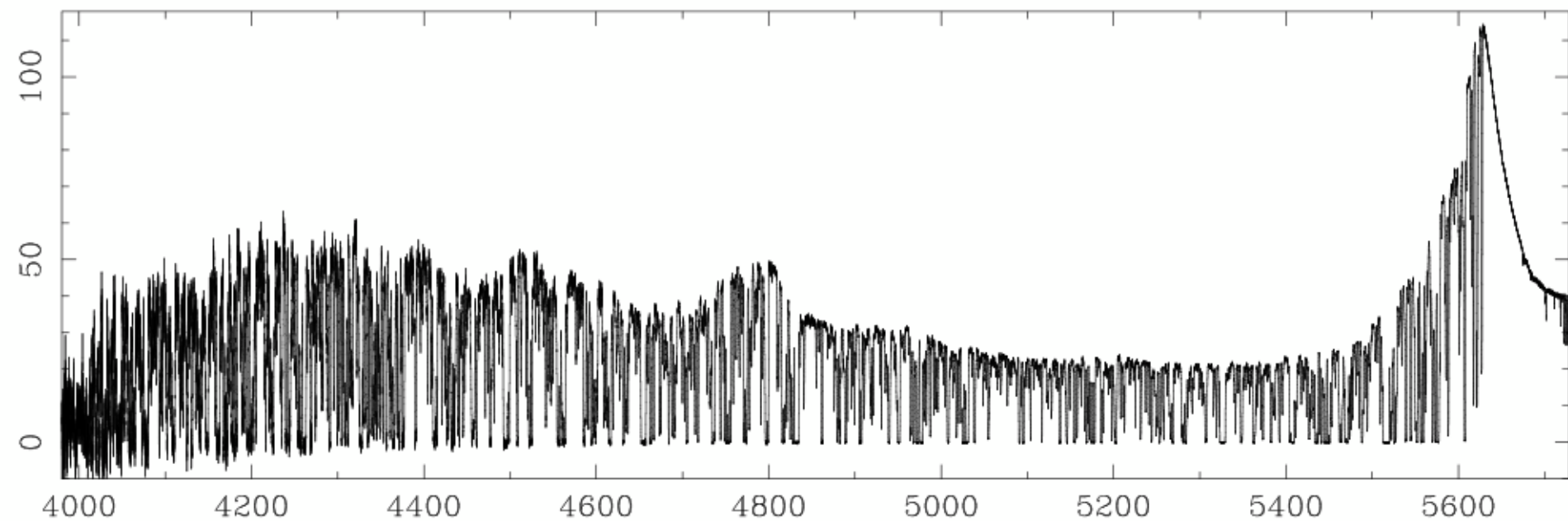
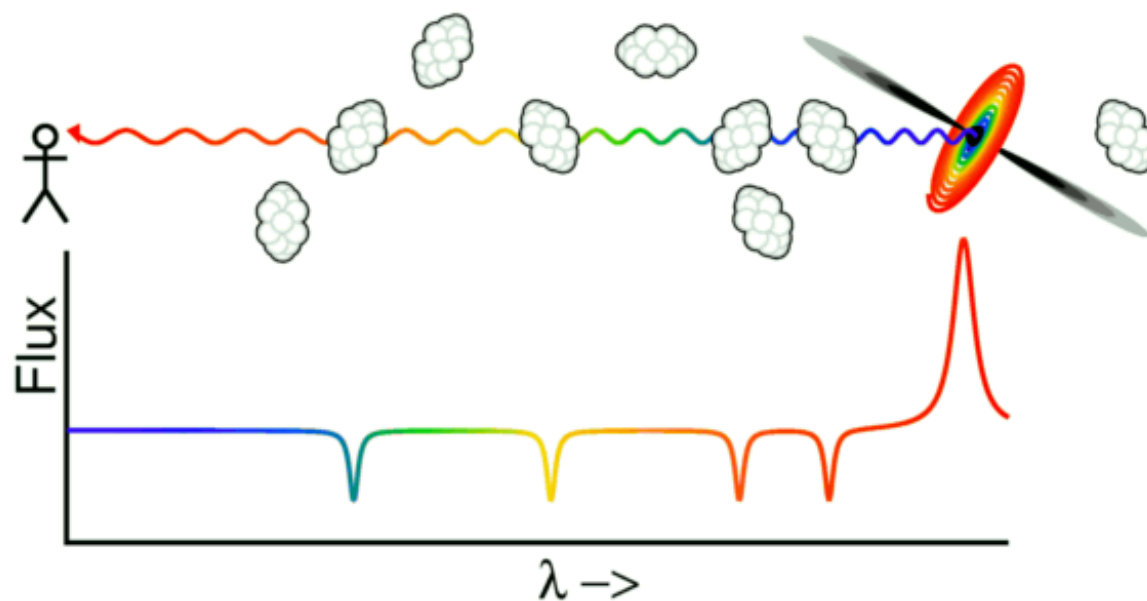
Density field tracers



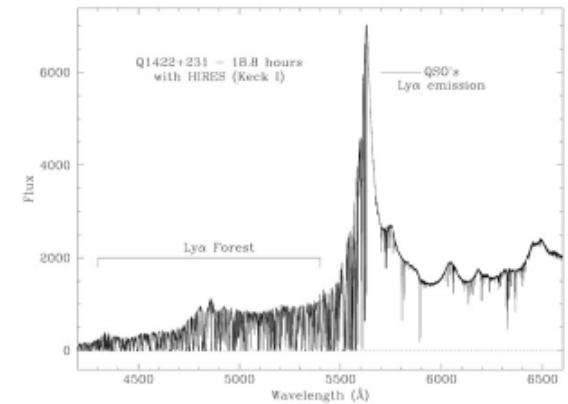
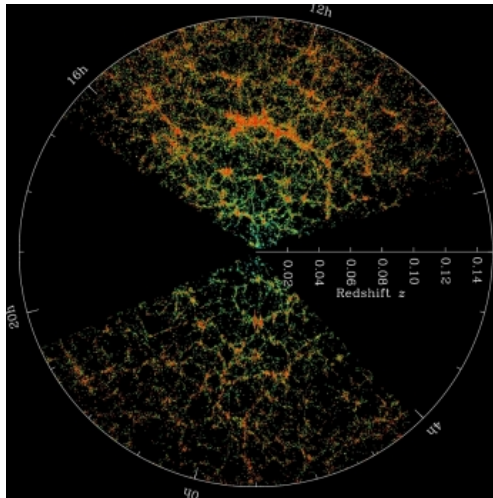
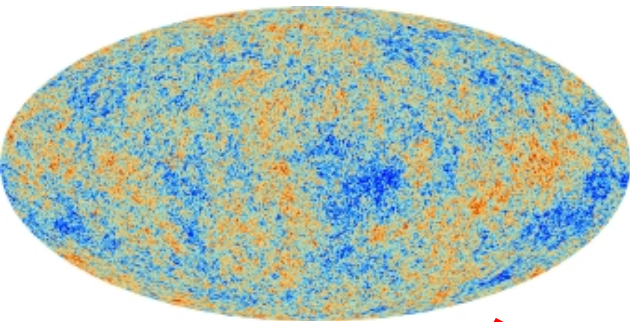
Density field tracers



Density field tracers: Ly α -forest



Cosmological constraints



$$\sum m_\nu < 0.72 \text{ eV (95\% C.L.)}$$

Planck TT + lowP
Planck collaboration 2015

$$\sum m_\nu \leq 0.21 \text{ eV (95\% C.L.)}$$

Planck TT + lowP + BAO
Planck collaboration 2015

$$\sum m_\nu \leq 0.13 \text{ eV (95\% C.L.)}$$

Planck TT + lowP + BAO + Ly α
Palanque-Delabrouille et al. 2015

Remember: model dependent!!!

Cosmological constraints

$$0.01 \text{ eV} < \sigma(\Sigma m_\nu) < 0.1 \text{ eV}$$

DESI – LSST – Euclid – Planck,...etc
Font-Ribera et al. 2014

$$\sigma(\Sigma m_\nu) \approx 0.01 \text{ eV}$$

Euclid clusters + Planck
Sartoris et al. 2015

$$\sigma(\Sigma m_\nu) \approx 0.03 \text{ eV}$$

SKA + Planck
Villaescusa-Navarro et al. 2015

Parameter estimation

$$\text{Data} \left\{ \begin{array}{l} \vec{x} = \{P(k_1, z_i), P(k_2, z_i), P(k_3, z_i), \dots, P(k_1, z_j), P(k_2, z_j), P(k_3, z_j), \dots\} \\ \vec{\theta} = \{\Omega_m, \Omega_b, \Omega_\Lambda, h, \sigma_8, n_s, \sum m_\nu, \dots\} \end{array} \right. \quad \text{Model parameters}$$

$$\mathcal{P}(\vec{\theta} | \vec{x}) = \frac{\mathcal{P}(\vec{\theta}) \mathcal{P}(\vec{x} | \vec{\theta})}{\mathcal{P}(\vec{x})}$$

Assumption!!!

Data are Gaussian distributed
Gaussian likelihood function

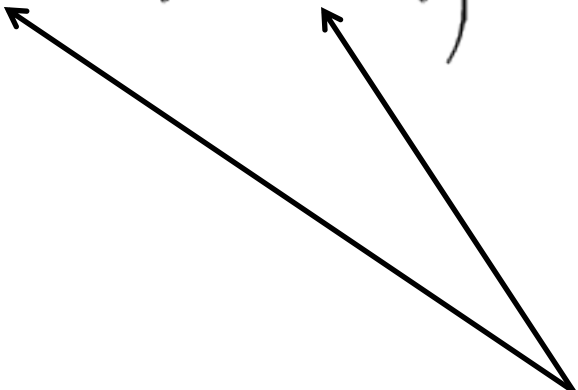
$$\mathcal{L} = \frac{1}{(2\pi)^{n/2} |C|^{1/2}} \exp \left(-\frac{1}{2} \sum_{ij} (\vec{x} - \vec{x}^{th})_i C_{ij}^{-1} (\vec{x} - \vec{x}^{th})_j \right)$$

Best fit and confident intervals typically found through MCMC

Parameter estimation

Assumption!!!

Gaussian likelihood function

$$\mathcal{L} = \frac{1}{(2\pi)^{n/2} |C|^{1/2}} \exp\left(-\frac{1}{2} \sum_{ij} (\vec{x} - \vec{x}^{th})_i C_{ij}^{-1} (\vec{x} - \vec{x}^{th})_j\right)$$


Most of the work in cosmology has been to accurately predict

Importance of the covariance?

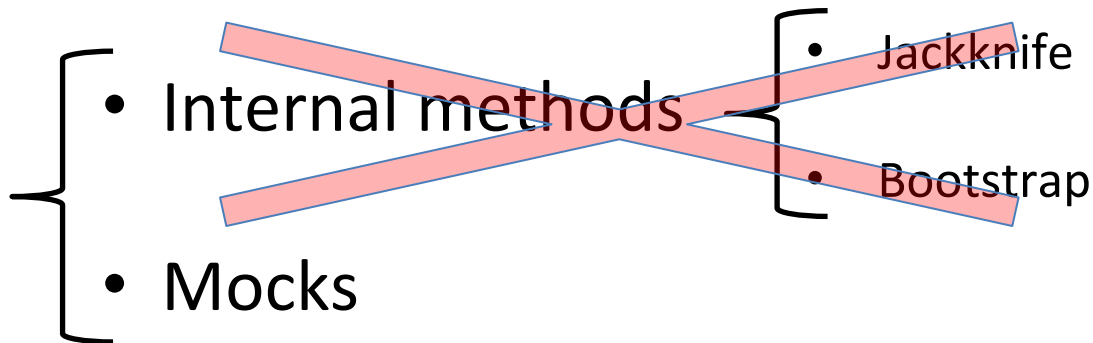
- Not known theoretically
- Need to be estimated somehow

Parameter estimation

Assumption!!!

Gaussian likelihood function

$$\mathcal{L} = \frac{1}{(2\pi)^{n/2} |C|^{1/2}} \exp\left(-\frac{1}{2} \sum_{ij} (\vec{x} - \vec{x}^{th})_i C_{ij}^{-1} (\vec{x} - \vec{x}^{th})_j\right)$$



Norberg et al. 2009

Parameter estimation

Assumption!!!

Gaussian likelihood function

$$\mathcal{L} = \frac{1}{(2\pi)^{n/2} |C|^{1/2}} \exp\left(-\frac{1}{2} \sum_{ij} (\vec{x} - \vec{x}^{th})_i C_{ij}^{-1} (\vec{x} - \vec{x}^{th})_j\right)$$

$$\hat{C} = \frac{1}{N-1} \sum_{i=1}^N \left(P(k_i) - \langle P(k_i) \rangle \right) \left(P(k_i) - \langle P(k_i) \rangle \right)^T$$

$$\hat{C}^{-1} =$$

$$\hat{C}^{-1}$$

Hartlap et al. 2007

Underestimation of the size
of the confidence regions

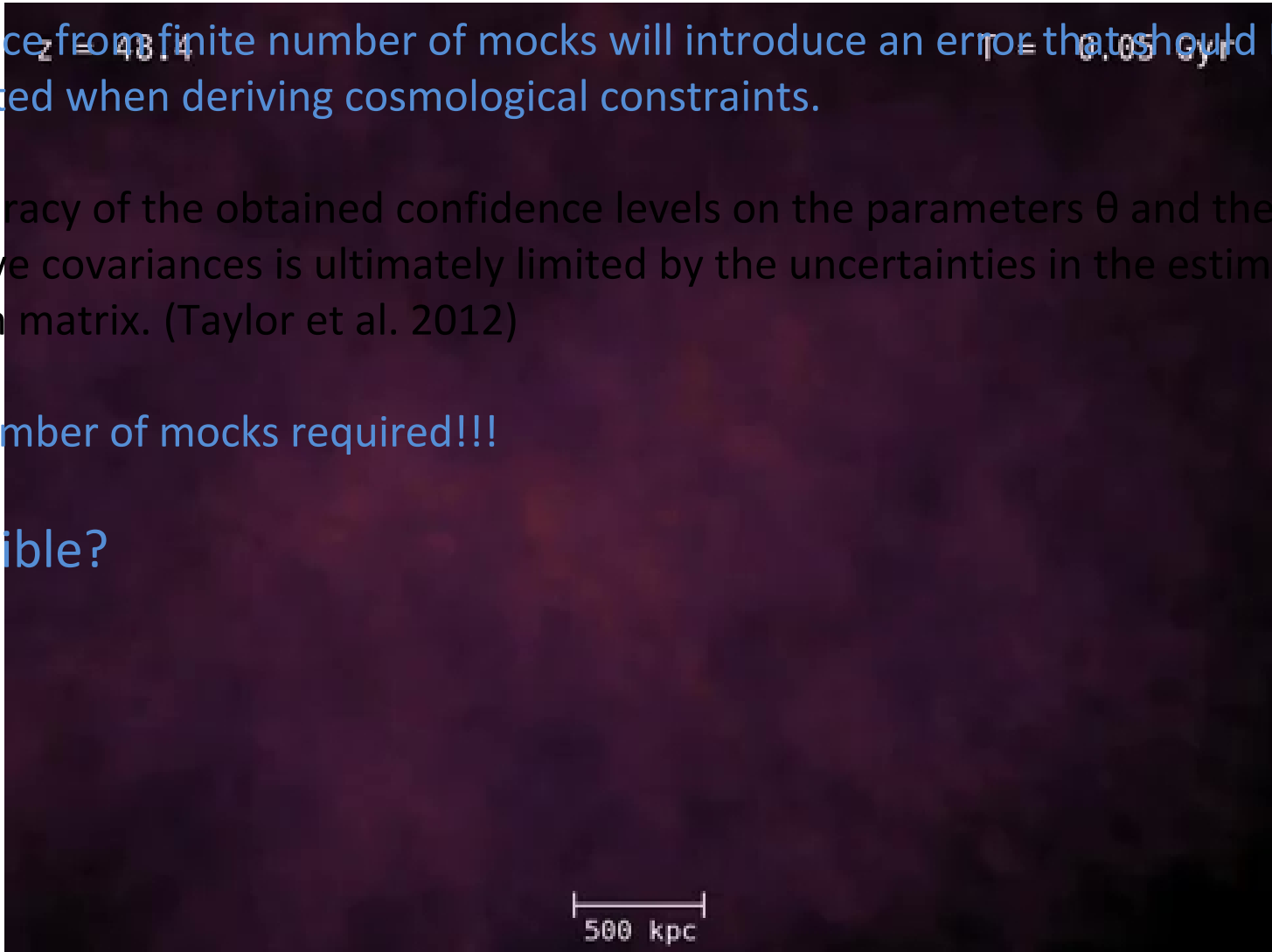
Parameter estimation

Covariance from finite number of mocks will introduce an error that should be propagated when deriving cosmological constraints.

The accuracy of the obtained confidence levels on the parameters θ and their respective covariances is ultimately limited by the uncertainties in the estimated precision matrix. (Taylor et al. 2012)

Large number of mocks required!!!

- Feasible?



Parameter estimation

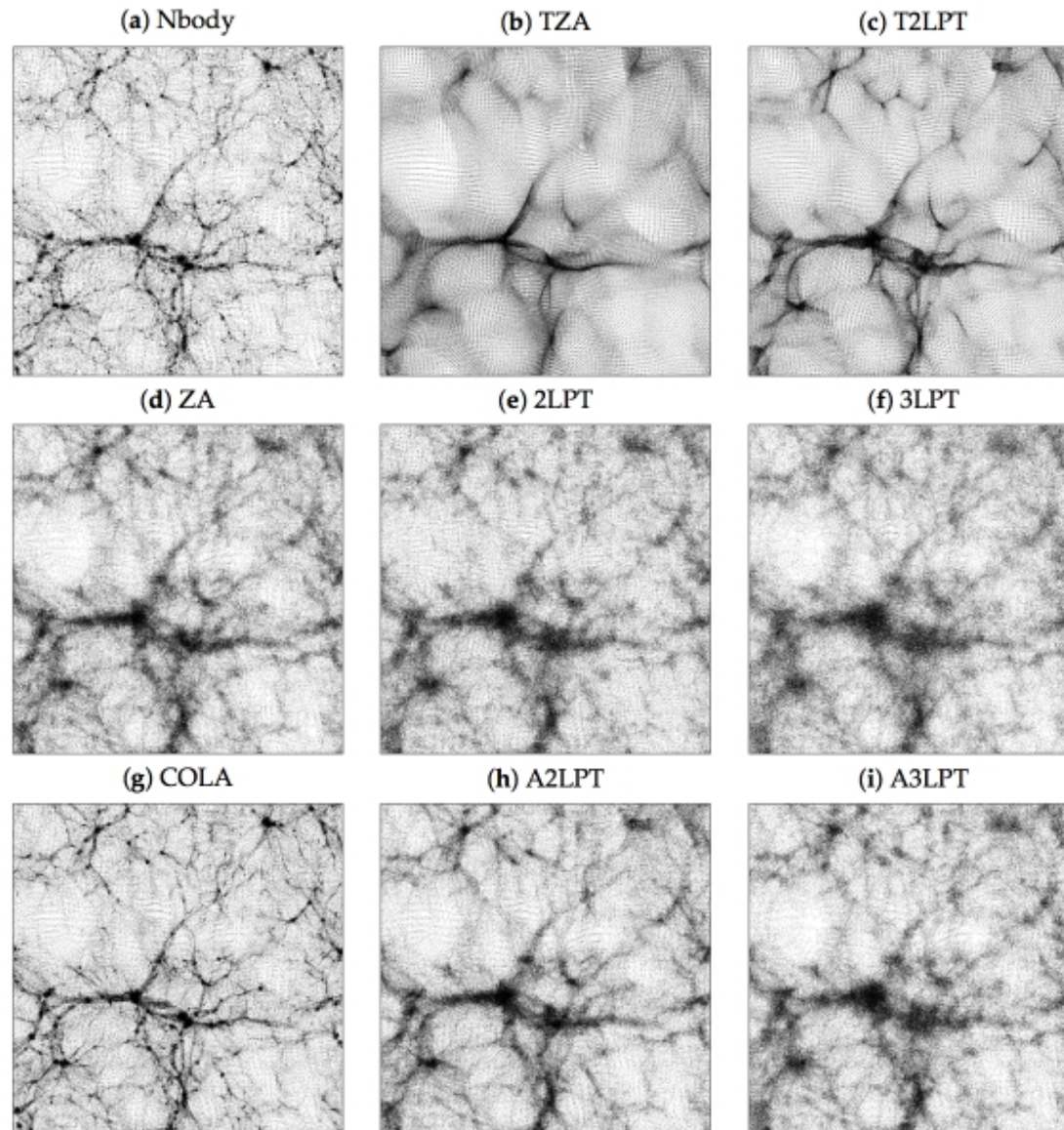


Figure 2. Large-scale structure in a $20 h^{-1}$ Mpc thick slice of a square of $128 h^{-1}$ Mpc, as recovered by different methods: (a) N-body simulation; (b) truncated ZA; (c) truncated 2LPT; (d) ZA; (e) 2LPT; (f) 3LPT; (g) COLA; (h) augmented 2LPT; (i) augmented 3LPT.

Parameter estimation

Assumption!!!

Gaussian likelihood function

$$\mathcal{L} = \frac{1}{(2\pi)^{n/2} |C|^{1/2}} \exp\left(-\frac{1}{2} \sum_{ij} (\vec{x} - \vec{x}^{th})_i C_{ij}^{-1} (\vec{x} - \vec{x}^{th})_j\right)$$

$$\hat{C} = \frac{1}{N-1} \sum_{i=1}^N \left(P(k_i) - \langle P(k_i) \rangle \right) \left(P(k_i) - \langle P(k_i) \rangle \right)^T$$

$$\hat{C}^{-1} = \frac{N-p-2}{N-1} \hat{C}^{-1}$$

Hartlap et al. 2007

Underestimation of the size
of the confidence regions

$$\mathcal{L} = \frac{\bar{c}_p |C|^{-1/2}}{\left[1 + \frac{\sum_{ij} (\vec{x} - \vec{x}^{th})_i C_{ij}^{-1} (\vec{x} - \vec{x}^{th})_j}{N-1} \right]^{N/2}}$$

Estimated covariance matrix follows a Wishart distribution. Compute the likelihood by marginalising over the inverse-Wishart distribution of the true covariance matrix Σ , conditioned on C. This gives a modified multivariate t-distribution

Sellentin & Heavens 2016

Parameter estimation

$$\mathcal{L} = \frac{1}{(2\pi)^{n/2} |C|^{1/2}} \exp\left(-\frac{1}{2} \sum_{ij} (\vec{x} - \vec{x}^{th})_i C_{ij}^{-1} (\vec{x} - \vec{x}^{th})_j\right)$$

- We expect that the covariance matrix depends on the parameters!!!
- That dependence is neglected in current analysis
- It is feasible to estimate it on the whole parameter space via MCMC?

Parameter estimation

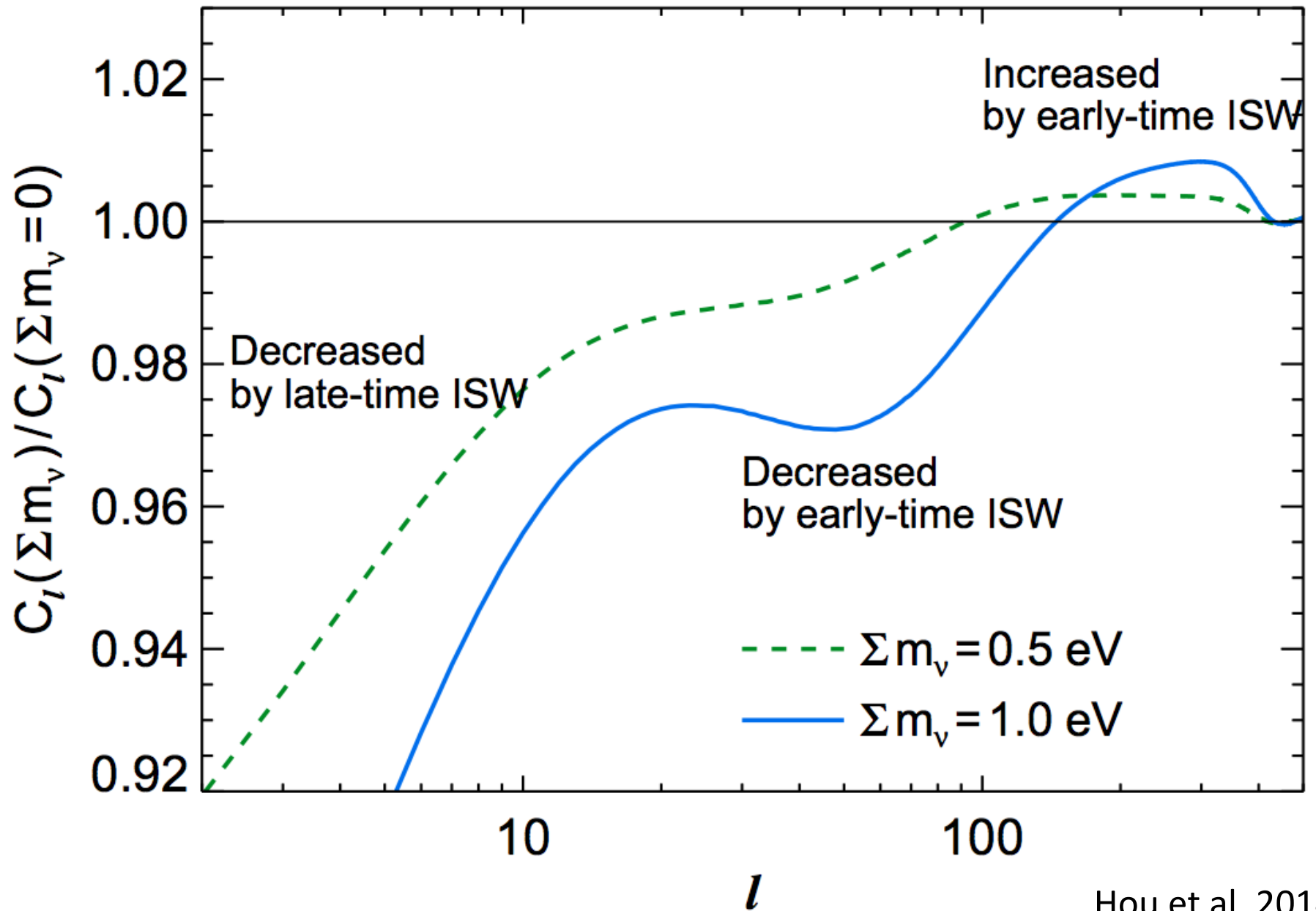
$$\mathcal{L} = \frac{1}{(2\pi)^{n/2} |C|^{1/2}} \exp\left(-\frac{1}{2} \sum_{ij} (\vec{x} - \vec{x}^{th})_i C_{ij}^{-1} (\vec{x} - \vec{x}^{th})_j\right)$$

1. Can we keep describing the likelihood as a multivariate Gaussian?
2. Can we compute covariance matrices to percent accuracy on non-diagonal elements?
3. Can we neglect the dependence of covariance on parameters?

Conclusions

- Big Bang theory predicts existence of a CνB
- Massive neutrino effects well understood at linear order
- Use cosmological observables to weigh neutrinos
- Challenges for subpercent cosmology:
covariance matrices

Neutrino effects on CMB



$$\mathcal{G}(\vec{x}|\vec{\mu}(\vec{p}), C) = \frac{1}{\sqrt{|2\pi C|}} \exp \left(-\frac{1}{2}(\vec{x} - \vec{\mu})^T C^{-1}(\vec{x} - \vec{\mu}) \right)$$

$$S = \frac{1}{N-1} \sum_{i=1}^N (\vec{x}_i - \bar{\vec{x}})(\vec{x}_i - \bar{\vec{x}})^T$$

→ C is **not** a random object – the estimator S **is**

→ $C^{-1} = \alpha \langle S^{-1} \rangle$ where $\alpha = \frac{N-p-2}{N-1}$

→ Hartlap et al. (2007): Gaussian with $C^{-1} \rightarrow \alpha S^{-1}$

But: Removing $\langle \rangle$ means ignoring the randomness of S!

Hartlap et al.'s method from (2007):

$$C^{-1} \rightarrow \alpha S^{-1}$$

$$L(\vec{x}|\vec{\mu}, S) = \int \mathcal{G}(\vec{x}|\vec{\mu}, C) \underline{\mathcal{P}(C|S, N)} dC \leftarrow \begin{array}{l} \text{matrix integration} \\ \text{matrix-valued prob. dens. func.} \end{array}$$

- $\mathcal{P}(S|C) = \text{Wishart distribution}$
- $\pi(C) \propto |C|^{-(p+1)/2}$

$$\rightarrow \mathcal{T}(\vec{x}|\vec{\mu}, S) \propto \frac{|S|^{-1/2}}{\left[1 + \frac{(\vec{x}-\vec{\mu})^T S^{-1} (\vec{x}-\vec{\mu})}{N-1}\right]^{\frac{N}{2}}}$$

Eq. (12) in arXiv 1511:05969

(changes 3 lines of codes)

