



Spatial Modulation and Topological Current in Holographic QCD Matter



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Outline



***p*-wave pion condensation in nuclear matter**

An old but vital idea of inhomogeneity in nuclear physics, but...

Landau-Migdal interaction in a Fermi liquid

Almost abandoned... and revived recently in quark matter

Interaction controlled by topological current in B – *Chiral Magnetic Effect* –

Effect of the topological current induced by the magnetic field

Analysis in the Sakai-Sugimoto model

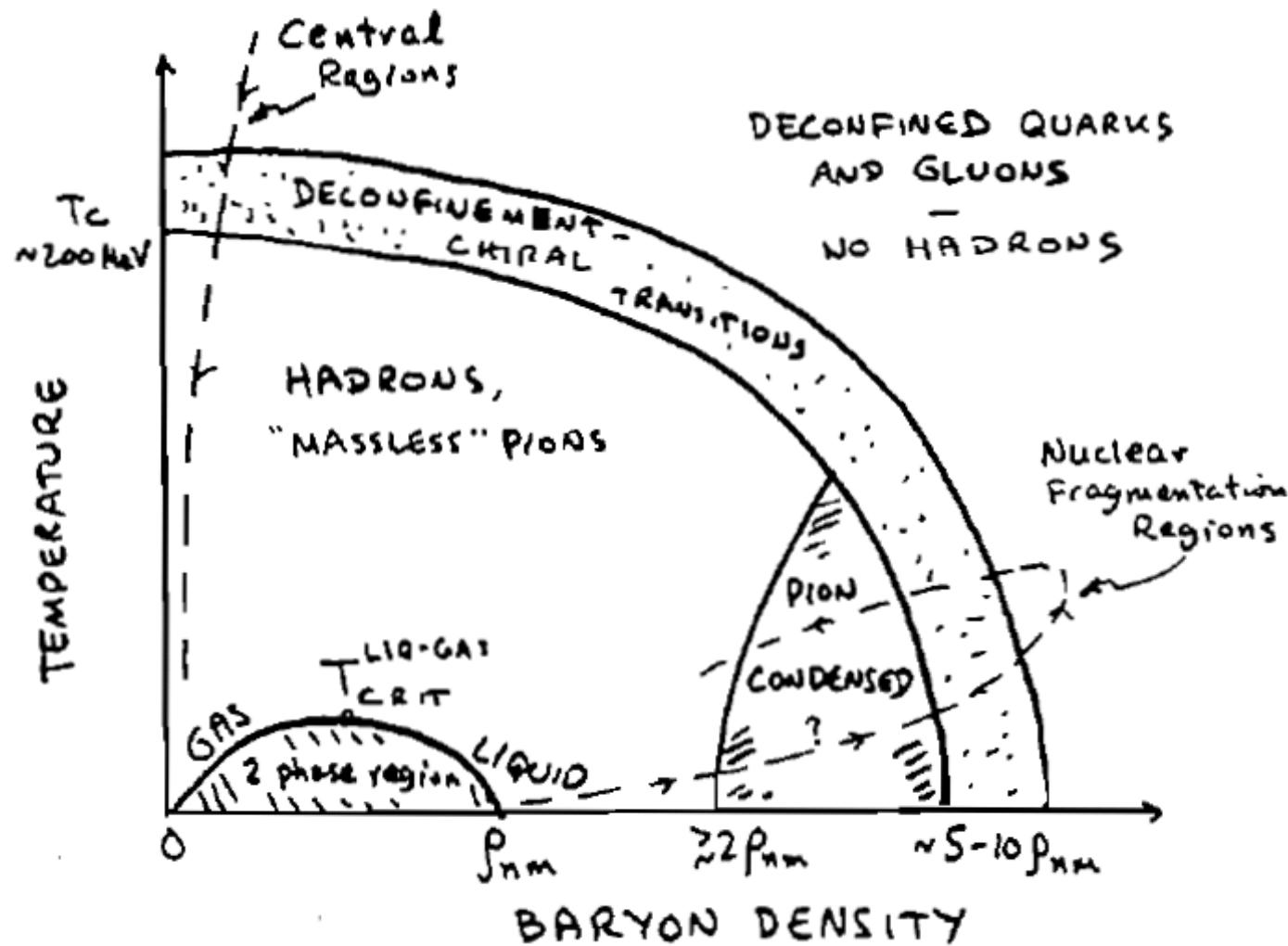
Inhomogeneous states favored or disfavored by B ???

Results consistent with what happened to the pion condensation

Historical Phase Diagram of QCD

Baym (1983)

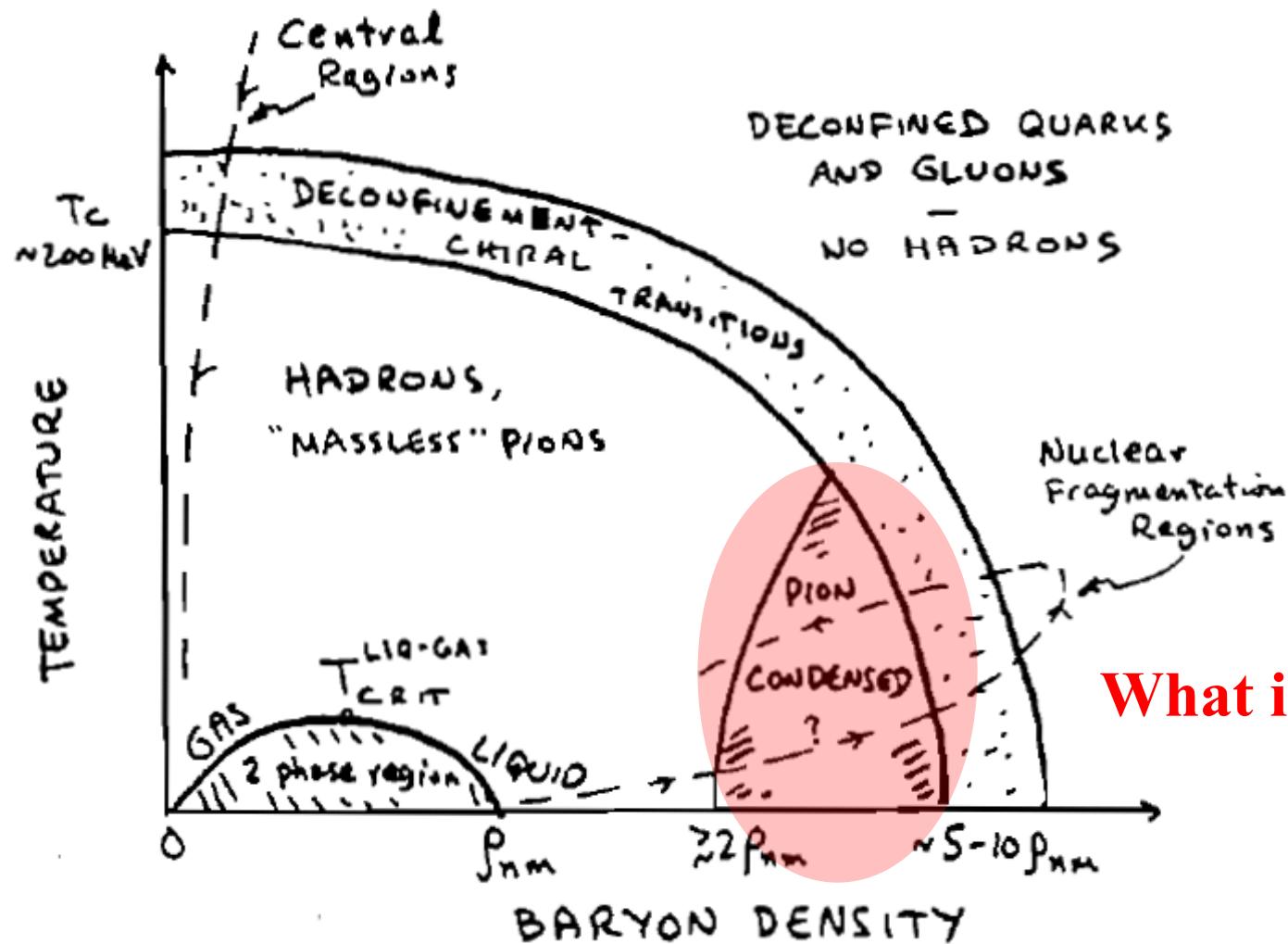
PHASE DIAGRAM OF NUCLEAR MATTER.



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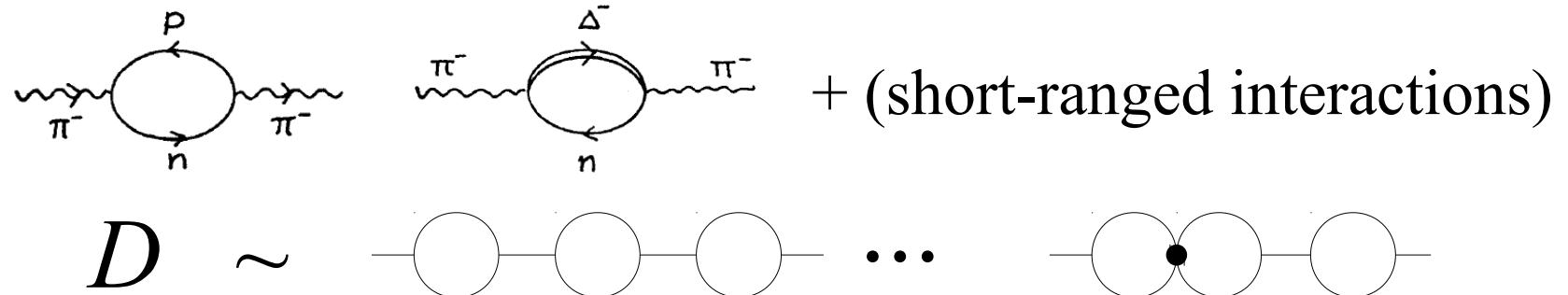


p-wave Pion Condensation



Sawyer-Scalapino, Migdal (1972)

$$\Pi(\omega, k) \rightarrow D^{-1}(\omega=0, k=k_c) = 0 \text{ at } \rho=\rho_c$$



One Pion Exchange Potential (OPEP)

$$V = \frac{m_\pi^2}{3} \frac{g^2}{4\pi} \boldsymbol{\tau}_1 \cdot \boldsymbol{\tau}_2 \left[\boldsymbol{\sigma}_1 \cdot \boldsymbol{\sigma}_2 \frac{e^{-m_\pi r}}{r} + S_{12} \left(1 + \frac{3}{m_\pi r} + \frac{3}{(m_\pi r)^2} \right) \frac{e^{-m_\pi r}}{r} \right] - \frac{g^2}{3} (\boldsymbol{\sigma}_1 \cdot \boldsymbol{\sigma}_2) (\boldsymbol{\tau}_1 \cdot \boldsymbol{\tau}_2) \delta(r)$$

Landau-Migdal (short-ranged) Interactions

$$f + g \boldsymbol{\sigma}_1 \cdot \boldsymbol{\sigma}_2 + f' \boldsymbol{\tau}_1 \cdot \boldsymbol{\tau}_2 + g' (\boldsymbol{\sigma}_1 \cdot \boldsymbol{\sigma}_2) (\boldsymbol{\tau}_1 \cdot \boldsymbol{\tau}_2)$$

p-wave Pion Condensation



Why *p*-wave?

$$\left| \begin{array}{l} \text{Relative angular-momentum must be 1 (p-wave)} \\ -\text{-- } \pi \text{ (negative parity)} \\ N \quad \langle \sigma \rangle \sim \chi \cos(2qz) \quad \langle \pi^0 \rangle \sim \chi \sin(2qz) \end{array} \right.$$

Landau-Migdal Parameter

g' is sensitive to the spin-isospin collective excitation
Gamow-Teller resonance \rightarrow large $g' \rightarrow$ **No pion condensation?**

Most of nuclear physicists consider no pion condensation up to a few times normal nuclear density

p-wave Pion Condensation



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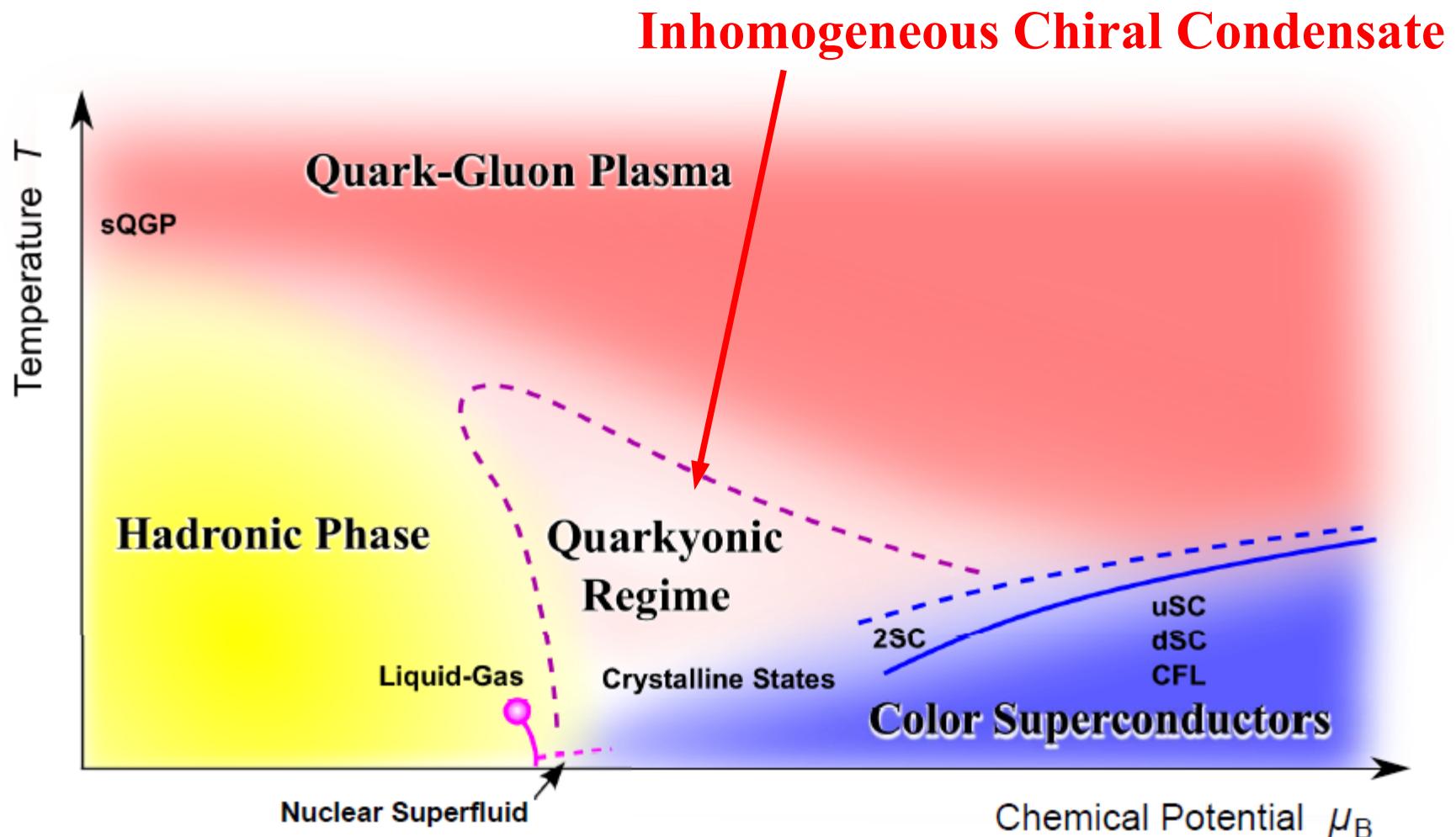
Most of nuclear physicists consider no pion condensation up to a few times normal nuclear density

History repeats itself (Nuclear \rightarrow Quark Matter)

One Possible QCD Phase Diagram



Fukushima-Sasaki (2013)



Energy Gain by Spiral



Chiral spiral (DGR Ansatz)

Deryagin-Grigoriev-Rubakov (1992)

$$\psi(x) = e^{i \gamma_5 \tau_3 q z} \psi'(x) \text{ with } \chi = \langle \bar{\psi}' \psi' \rangle$$

$$\langle \sigma \rangle \sim \langle \bar{\psi} \psi \rangle = \chi \cos(2qz)$$

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Quasi-particle dispersion relation

$$\omega = \sqrt{p_{\perp}^2 + (\sqrt{p_z^2 + M^2} \pm q)^2}$$

Vacuum

favors large M



High Density

favors small M

Vacuum + High Density
favors large $q \sim M$ most!

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Vacuum
favors large M

\longleftrightarrow
Competitive

High Density
favors small M

Counterpart of the Landau-Migdal int?
Can be the true ground state with strong B

Landau Quantization



Energy dispersion relation in B

$$\omega^2 = p_z^2 + \frac{2|gB|(n+1/2)}{m^2 - 2s g B}$$

Transverse motion = Harmonic Oscillator

- **Light fermions ($s=1/2$) have zero mode.**
Pseudo-(1+1) dimensional system of fermions
- **Light vector bosons have (Nielsen-Olesen) instability.**
Gluons in the chromo- B / ρ in a superstrong B (Chernodub)
- **Charged scalar bosons are all massive.**
 π^+ , π^- , ... Explicit breaking of isospin symmetry
Etc, etc...

Pseudo-(1+1) dimensional System



Dirac Lagrangian in (1+1) dimensions

$$L = \bar{\Psi} [(\partial_4 + \mu) \gamma^4 + \partial_3 \gamma^3] \Psi = \bar{\Psi}' [\partial^4 \gamma_4 + \partial_3 \gamma^3] \Psi'$$

$$\psi = e^{-\mu \gamma^3 \gamma^4 x_3} \psi' \quad \bar{\psi} = \bar{\psi}' e^{-\mu \gamma^3 \gamma^4 x_3}$$

Thermodynamic potential

$$\Omega/V = - \int_{-\Lambda+\mu}^{\Lambda-\mu} \frac{dp}{2\pi} \frac{|\varepsilon(p)|}{2} - \int_{-\Lambda-\mu}^{\Lambda+\mu} \frac{dp}{2\pi} \frac{|\varepsilon(p)|}{2} + \dots$$

$$= \Omega(\mu=0)/V - \frac{\mu^2}{2\pi}$$

Surface integral: Anomaly origin
c.f. CS term in Sakai-Sugimoto

No mass suppression as compared to the homogeneous case
Energy gain by spiral maximized in (1+1) dimensions!

Pseudo-(1+1) dimensional System



If the zero-density system has a condensate:

$$\langle \bar{\psi} \psi \rangle = (\text{homogeneous chiral condensate})$$

Rotated system has the same condensate: $\langle \bar{\psi}' \psi' \rangle$

(1+1)-dimensional system forms a “spiral”

$$\langle \bar{\psi} \psi \rangle = \langle \bar{\psi}' \psi' \rangle \cos(2\mu z)$$

$$\langle \bar{\psi} \gamma^3 \gamma^4 \psi \rangle = \langle \bar{\psi}' \psi' \rangle \sin(2\mu z)$$

Chiral Magnetic Spiral

Basar-Dunne-Kharzeev (2010)

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Chiral Magnetic Spiral

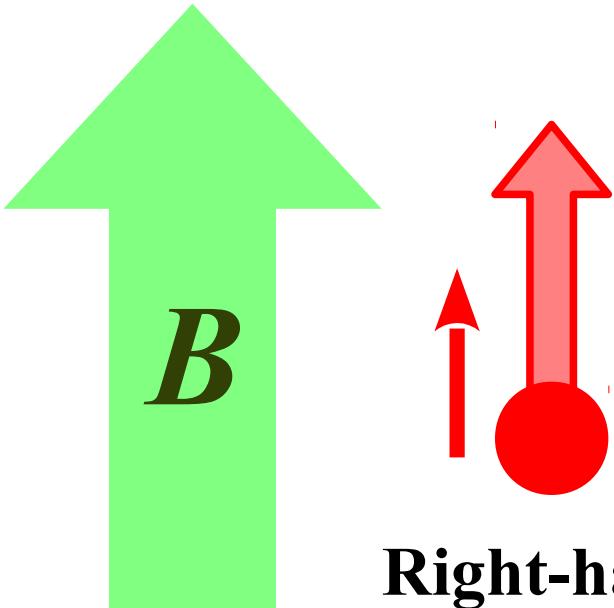
Basar-Dunne-Kharzeev (2010)

$B + \mu$ causes another interesting phenomenon!

Chiral Magnetic (Separation) Effect

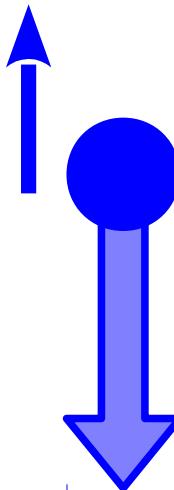


Classical Picture



Right-handed Quarks
= momentum parallel to spin

Left-handed Quarks
= momentum anti-parallel to spin



$$J_5 \neq 0 \text{ if } \mu \text{ or } N \neq 0$$

Kharzeev-McLerran-Warringa (2007)
Fukushima-Kharzeev-Warringa (2008)

Quantum Formula



Chiral Magnetic (Separation) Effect

$$j_5 = \frac{e^2 \mu}{2\pi^2} B \quad (j^\mu = \epsilon^{\mu\nu\rho\sigma} \partial_\nu \varphi F_{\rho\sigma})$$

Vector = Axial-Vector ($\gamma^5 = \gamma^0 \gamma^1$) in (1+1)

$$\gamma^\mu \gamma^5 = -\epsilon^{\mu\nu} \gamma_\nu \quad + \quad j_V^\mu = \bar{\Psi} \gamma^\mu \Psi, \quad j_A^\mu = \bar{\Psi} \gamma^\mu \gamma^5 \Psi$$

$$\longrightarrow \quad j_V^1 = j_A^0, \quad j_A^1 = j_V^0$$

Question



At high density

Chiral spiral is expected?

At strong B

Chiral spiral is favored by low dimensionality...



Competitive

Axial-current is strengthened leading to –

$$g'(\sigma_1 \cdot \sigma_2)(\tau_1 \cdot \tau_2)$$

that disfavors the chiral spiral...

What we need is...



Spatial Inhomogeneity

Ooguri-Park, Chuang-Dai-Kawamoto-Lin-Yeh (2010)

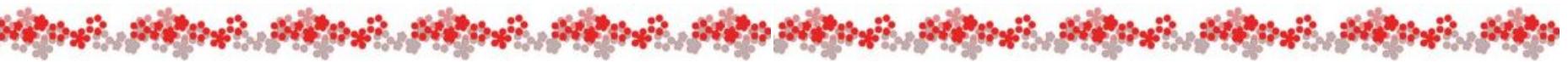
Topological Current

Yee, Rebhan-Schmitt-Stricker (2009)

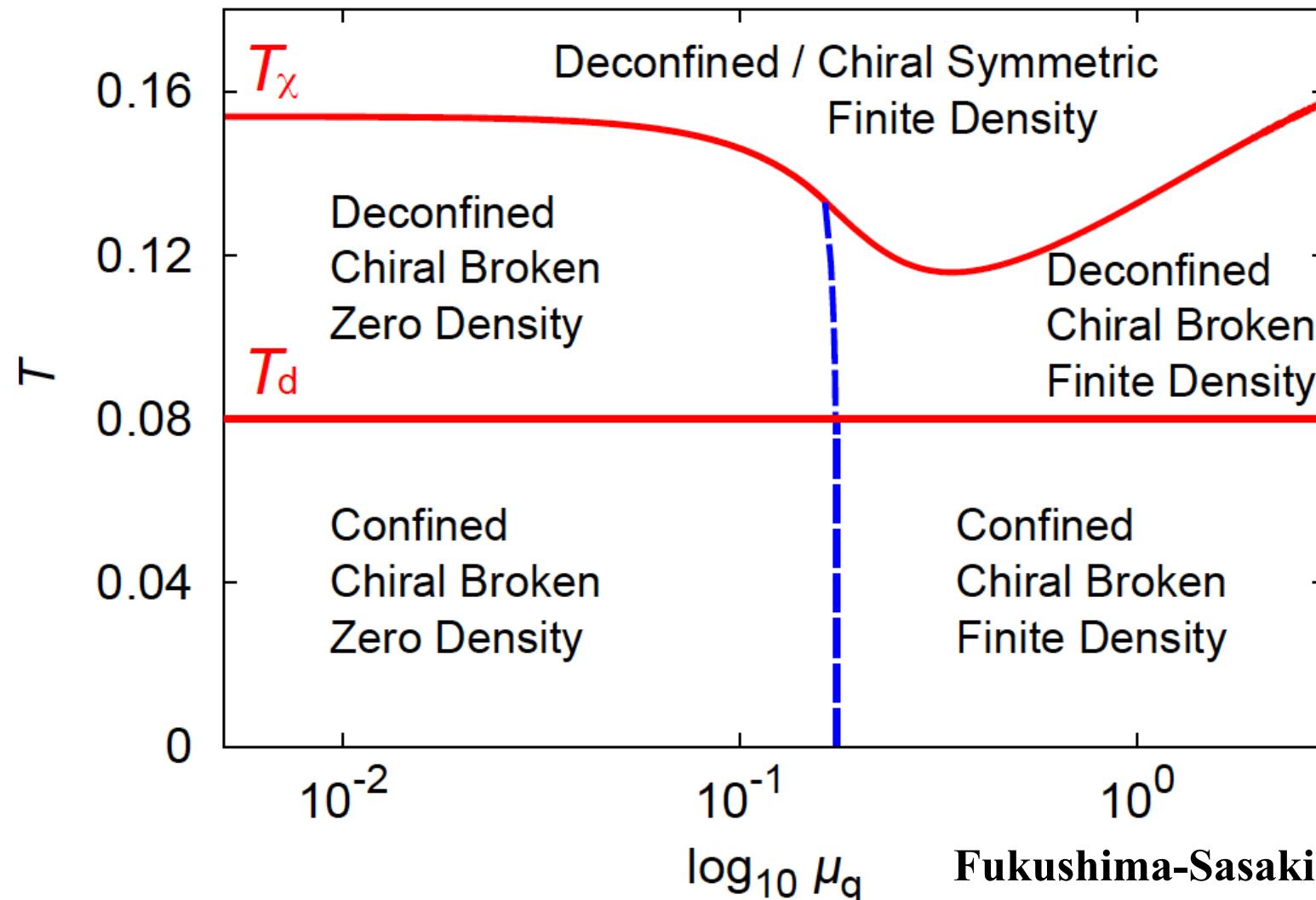
Non-perturbative Method → Sakai-Sugimoto Model

Holographic QCD Model with $D4 + D8 + \bar{D8}$ Sakai-Sugimoto (2004)

Phase Diagram in SSM



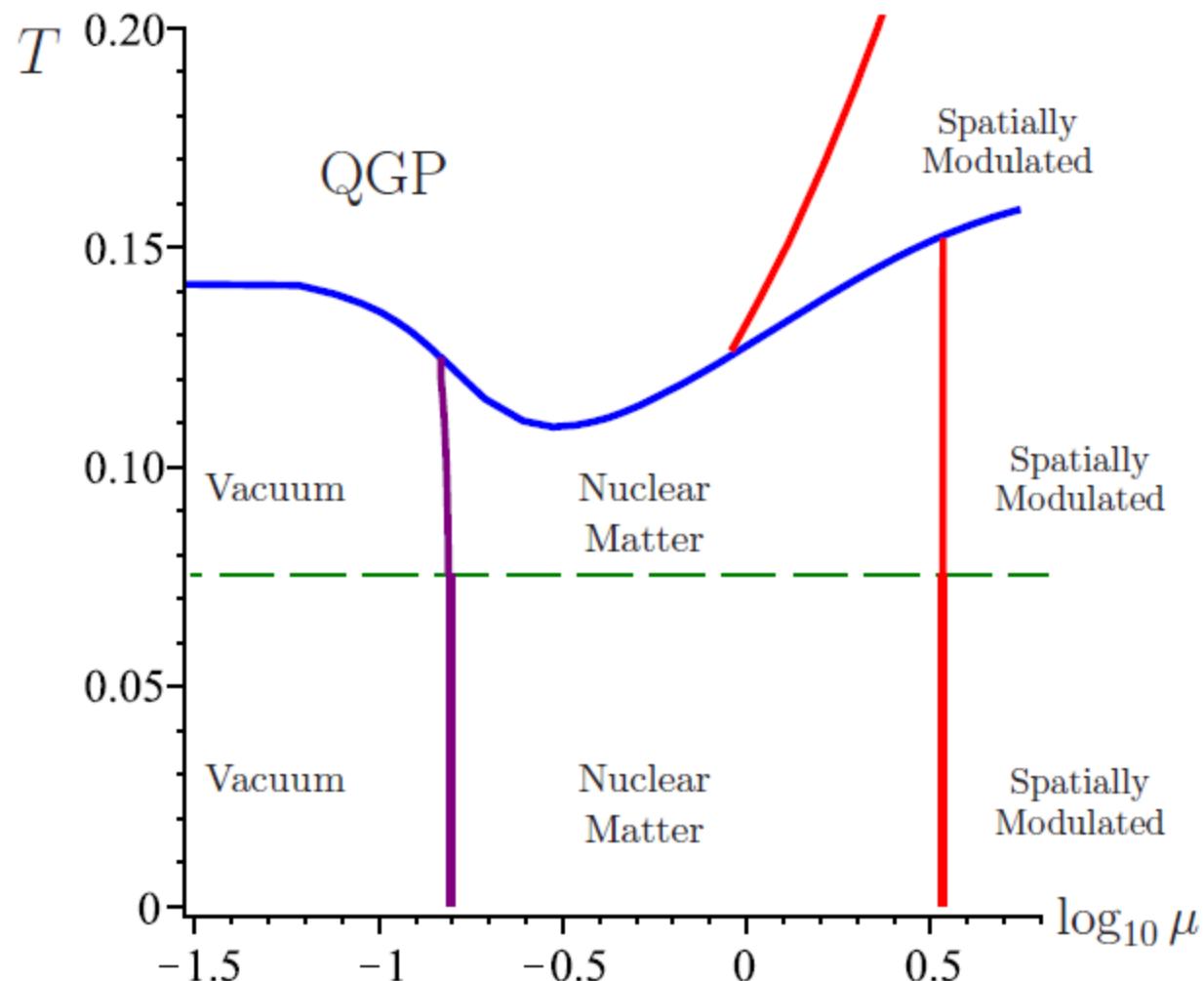
Bergman-Lifschytz-Lippert (2007)



Inhomogeneous States in SSM



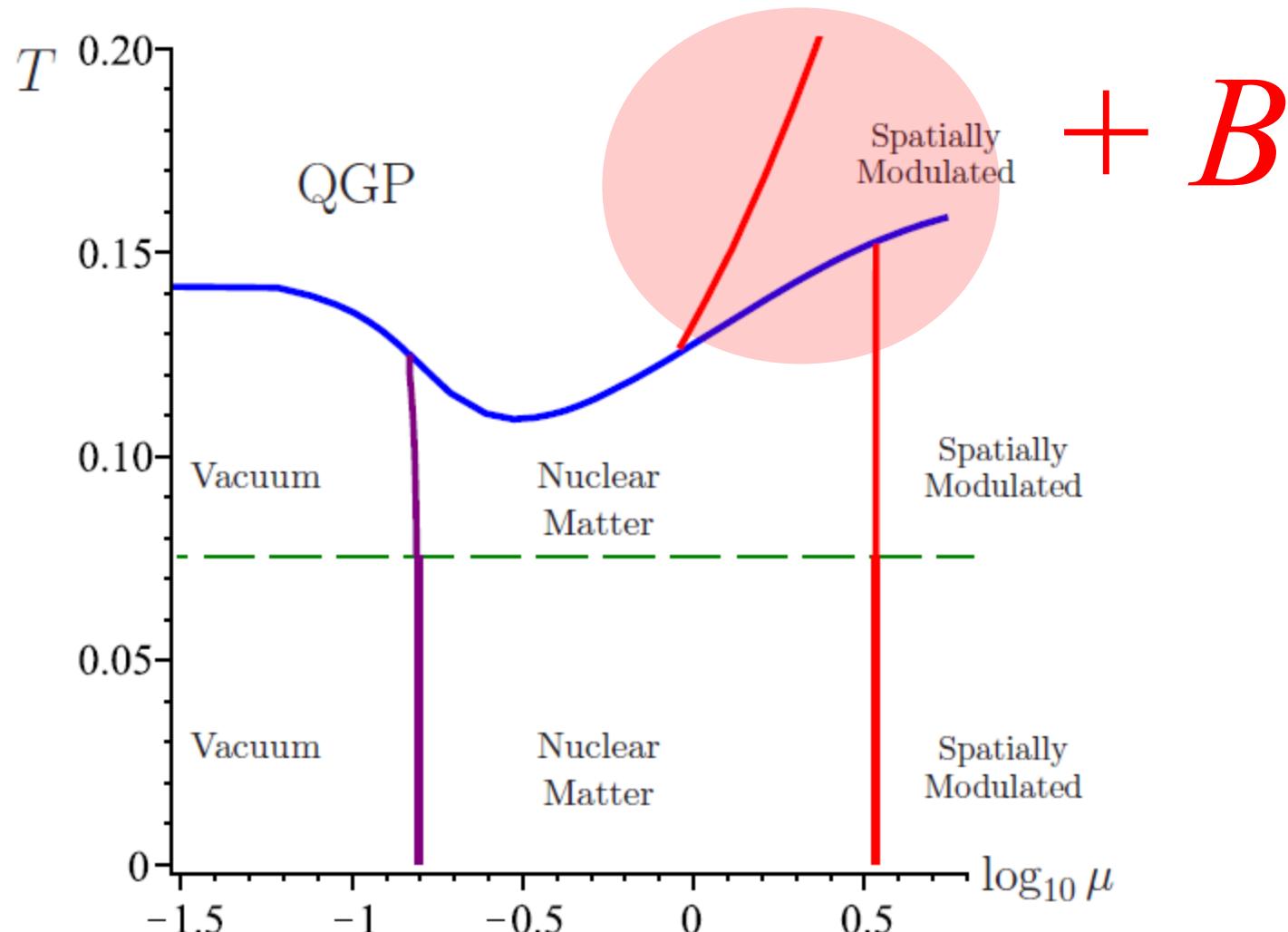
Chuang-Dai-Kawamoto-Lin-Yeh (2010)



Inhomogeneous States in SSM



Chuang-Dai-Kawamoto-Lin-Yeh (2010)



Sketch of the Calculations

Action for the flavor sector

$$S = N \int d^4x du u^{1/4} \sqrt{-\det(g_{\alpha\beta} + F_{\alpha\beta})} + \frac{\alpha}{3!} N \int d^4x du \epsilon^{\mu_1\mu_2\mu_3\mu_4\mu_5} A_{\mu_1} F_{\mu_2\mu_3} F_{\mu_4\mu_5}$$

DBI action

Chern-Simons action

Density	$\bar{a}_0(u)$	$\bar{a}_0(\infty) = \mu$
Current	$\bar{a}_z(u)$	
Magnetic field	$\bar{F}_{xy} = B(u) = B$	

$$S \sim \int du u^{5/2} \sqrt{(1 - {a_0}'^2 + f {a_z}'^2)(1 + B^2 u^{-3})} + 4\alpha B \int du \bar{a}_z \bar{a}_0'$$

Sketch of the Calculations

Equations of motion (w.r.t. a_0 and a_z)

$$\rho = u \bar{a}_0' \sqrt{\frac{u^3 + B^2}{1 - \bar{a}_0'^2 + f \bar{a}_z'^2}} - 4\alpha B \bar{a}_z$$

$$0 = u f \bar{a}_z' \sqrt{\frac{u^3 + B^2}{1 - \bar{a}_0'^2 + f \bar{a}_z'^2}} - 4\alpha B \bar{a}_0$$

Asymptotic solutions:

$$\bar{a}_z(u \rightarrow \infty) \simeq -\frac{8\alpha}{3} B \mu u^{-3/2}$$

$\sim j_5$

$$\bar{a}_0(u \rightarrow \infty) \simeq \mu - \frac{3\alpha}{2} \rho u^{-3/2}$$

Introducing spatial modulations:

$$F_{\alpha\beta} \rightarrow F_{\alpha\beta} + f_{\alpha\beta}(\omega, k)$$

\longrightarrow

$$D^{-1}(\omega=0, k=k_c) = 0$$
$$f_{\alpha\beta} = \partial_\alpha a_\beta - \partial_\beta a_\alpha$$

Sketch of the Calculations

Equations of motion (w.r.t. a_x , a_y , a_z)

From the condition that they are normalizable or

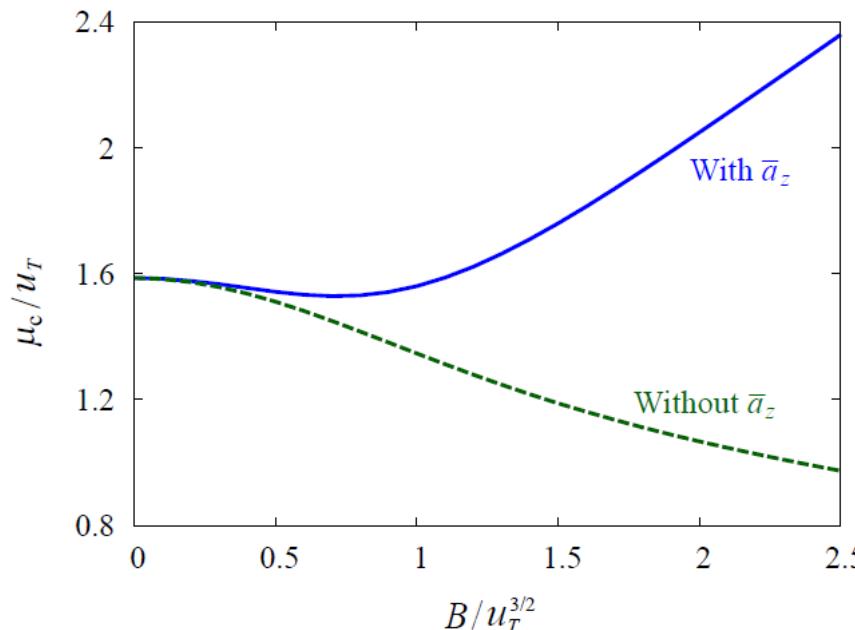
$$a_x(u=\infty) = a_y(u=\infty) = a_z(u=\infty) \rightarrow 0$$

→ k_x, k_y, k_z are found for $\rho > \rho_c$

Without B

$$\rho_c = 3.715 u_T^{5/2} \quad (\text{Ooguri-Park 2010})$$

With B



The presence of j_5 completely changes the results!

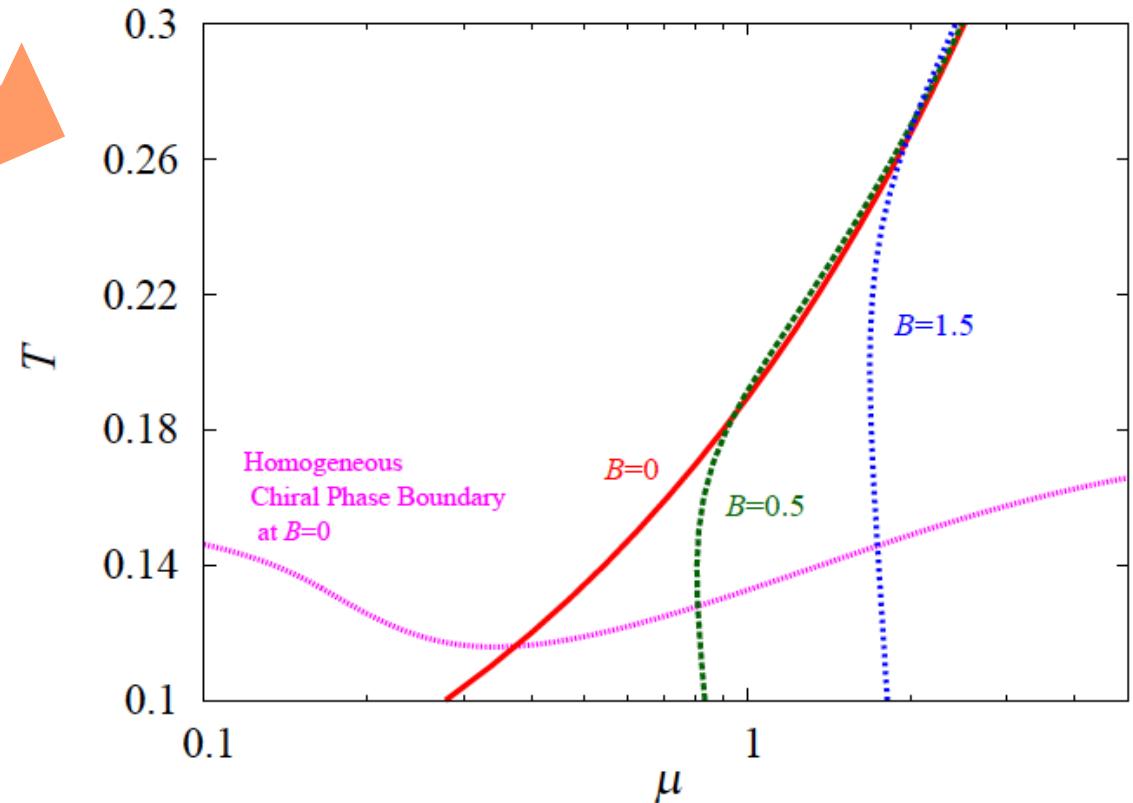
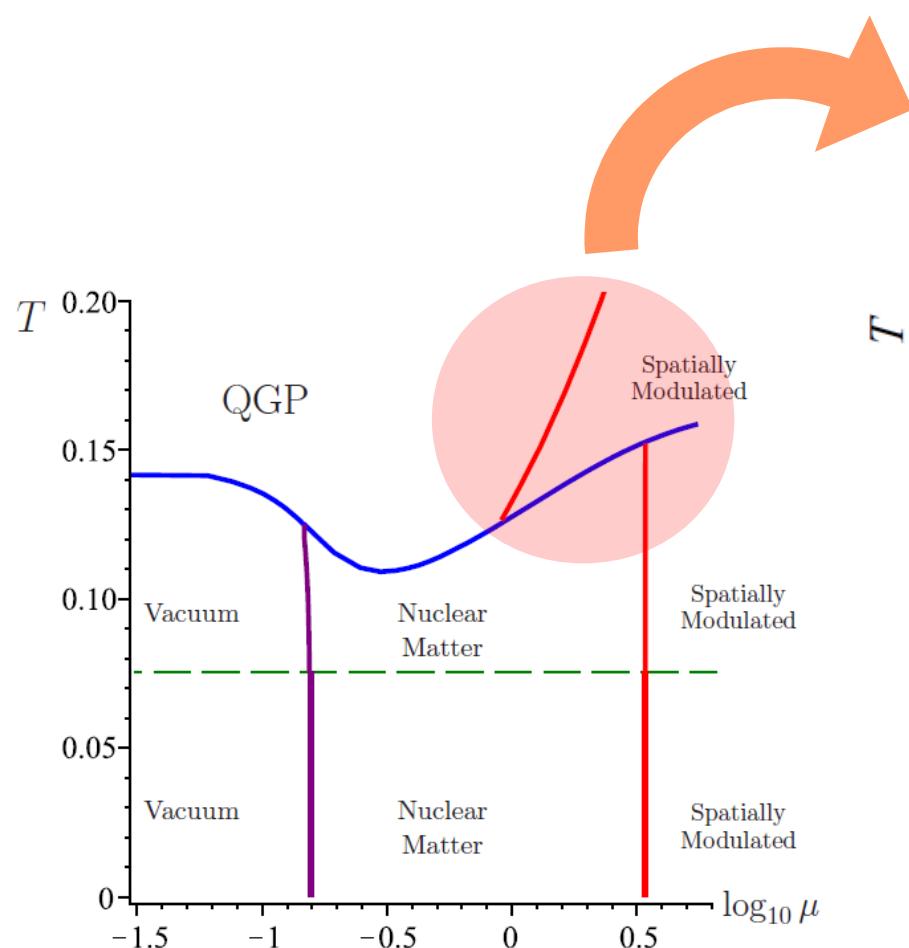
(Fukushima-Morales 2013)

Phase Diagram with B



Smaller with increasing B

Fukushima-Morales (2013)



Same tendency as what happened
to the p -wave pion condensation

Summary and...



Sakai-Sugimoto model is a powerful tool to investigate the *chiral sector of large- N_c QCD*.

→ Phase Diagram ??

Axial-vector interaction disfavors the spatial modulations in the same way as discussed by nuclear physicists long long time ago.

Full structure of the phase diagram with an extra axis of strong B ??

Earlier chiral phase transition? (Magnetic Inhibition)

Deconfinement?

(Fukushima-Hidaka 2012)